

Higgs Cascade Decays and the LHC

SB, Freese, Shakya, Shah; 1703.07800 (Phys. Rev. D)

SB, Shah; 1807.XXXXX

SB, Shah, Freese; 180X.XXXXX

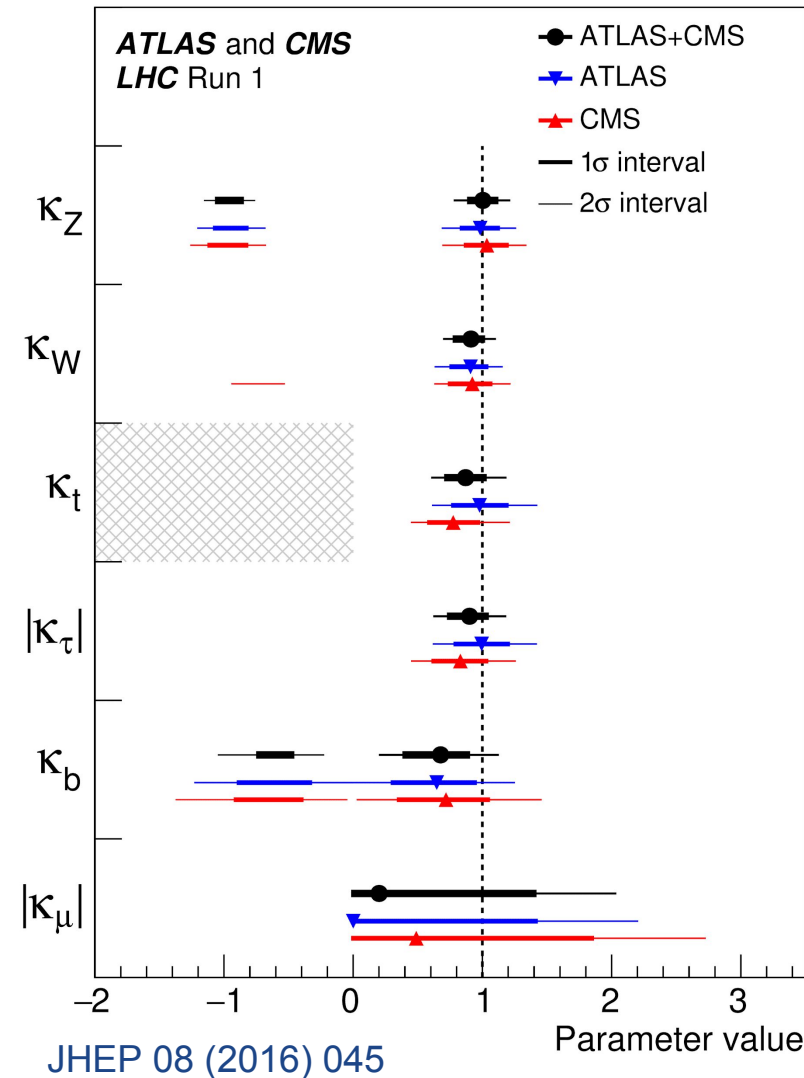
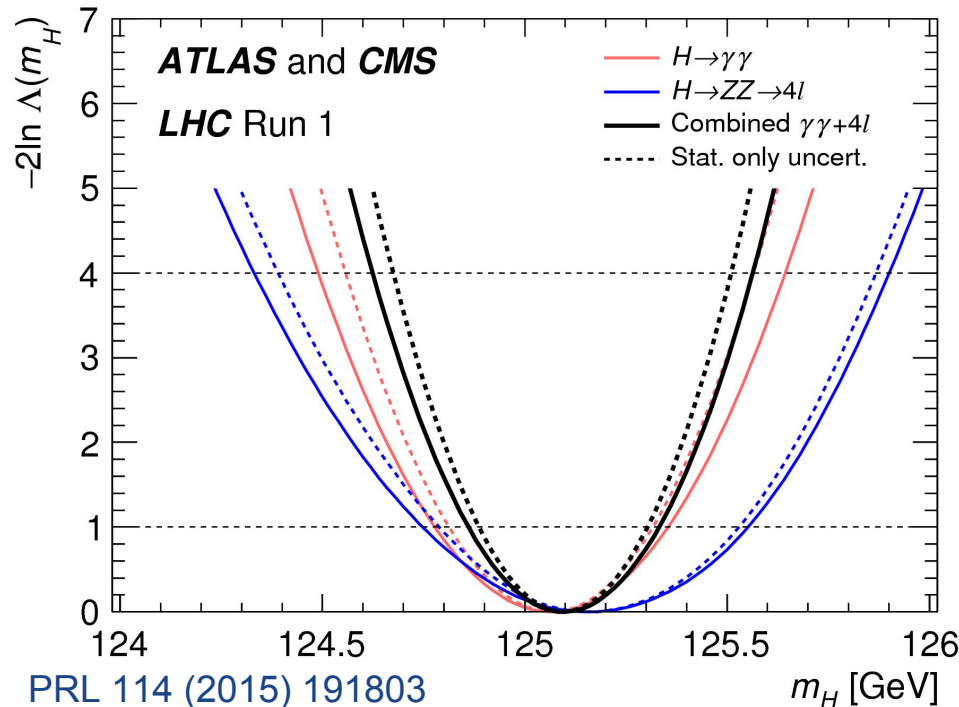
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SUSY 2018
Barcelona

All the Higgses that we know

- Mass ≈ 125 GeV
- Largest couplings within 10% of the SM expectation



Why do we want more Higgses?

Many shortcomings of the SM are intimately linked to the Higgs sector

- Hierarchy Problem
- Baryogenesis
- (Meta-)Stability Problem
- Flavor Structure

The Higgs sector is least well measured in the SM

Why 2HDM+S?

- 2HDM sector ubiquitous in BSM physics, e.g. SUSY
- SM gauge-singlets useful for
 - Baryogenesis [[Profumo+ 0705.2425](#), [Barger+ 0811.0393](#), [Cline+ 1210.4196](#)]
 - Constructing DM Models [[SB+ 1712.09873](#)]
- Same scalar sector as in well-motivated High-Scale models, e.g. the Next-to-Minimal Supersymmetric Standard Model

2HDM+S

- 2 SU(2)-doublets Φ_1 , Φ_2 , and 1 (complex) SM singlet S
- Most general CP-conserving scalar potential has 27 physical d.o.f., which can be re-parameterized by 6 physical masses, 4 mixing angles, 3 vevs, 10 trilinear and 4 quartic couplings

[SB, Shah 1807.XXXXX]

- Rotate to Higgs basis

$$\begin{aligned}\begin{bmatrix} G^+ \\ \frac{1}{\sqrt{2}} (H^{\text{SM}} + iG^0) \end{bmatrix} &= \sin \beta \Phi_1 + \cos \beta \Phi_2 \\ \begin{bmatrix} H^+ \\ \frac{1}{\sqrt{2}} (H^{\text{NSM}} + iA^{\text{NSM}}) \end{bmatrix} &= \cos \beta \Phi_1 - \sin \beta \Phi_2 \\ \frac{1}{\sqrt{2}} (H^{\text{S}} + iA^{\text{S}}) &= S ,\end{aligned}$$

- 7 degrees of freedom after EW symmetry breaking:

$$\{H^{\text{SM}}, H^{\text{NSM}}, H^{\text{S}}\} ; \quad \{A^{\text{NSM}}, A^{\text{S}}\} ; \quad H^\pm$$

Alignment and Mass Eigenstates

$$h_i = S_{h_i}^{\text{SM}} H^{\text{SM}} + S_{h_i}^{\text{NSM}} H^{\text{NSM}} + S_{h_i}^{\text{S}} H^{\text{S}}, \quad h_i = \{h_{125}, H, h\}$$

$$a_i = P_{a_i}^{\text{NSM}} A^{\text{NSM}} + P_{a_i}^{\text{S}} A^{\text{S}}, \quad a_i = \{A, a\}$$

- SM-like nature of the observed 125 GeV state requires a mass eigenstate approximately aligned with H^{SM}

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Alignment by decoupling

- Suppress mixing by choosing $\{H^{\text{NSM}}, H^{\text{S}}\}$ heavy wrt. H^{SM}

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Alignment without decoupling

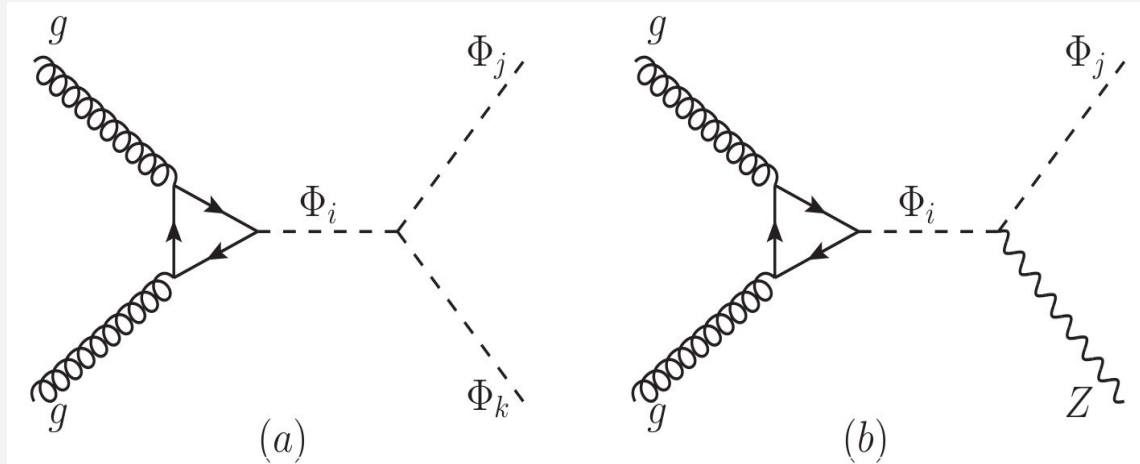
- Suppress mixing by conspiracy of parameters

⇒ Additional Higgses can remain light and thus accessible at LHC!

Cascade Decays

Complexity of Higgs sector allows for new decay modes

- $(h_i \rightarrow h_j h_k), (h_i \rightarrow a_j a_k), (a_i \rightarrow h_j a_k)$
- $(h_i \rightarrow Z a_j), (a_i \rightarrow Z h_j)$



Alignment suppressed

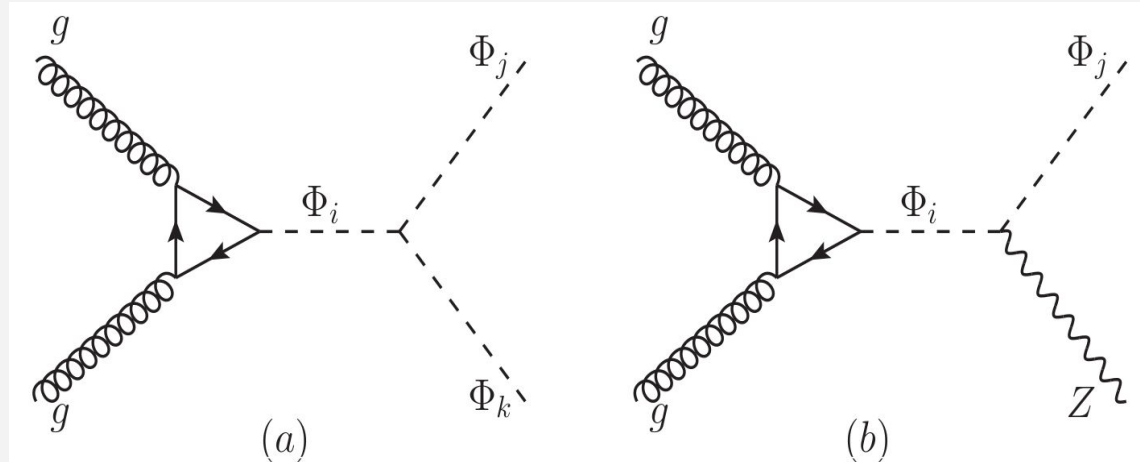
$$(H \rightarrow h_{125} h_{125})$$

$$(A \rightarrow Z h_{125})$$

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Dominant modes:

$$(H \rightarrow h_{125} h)$$

$$(H \rightarrow Z a)$$

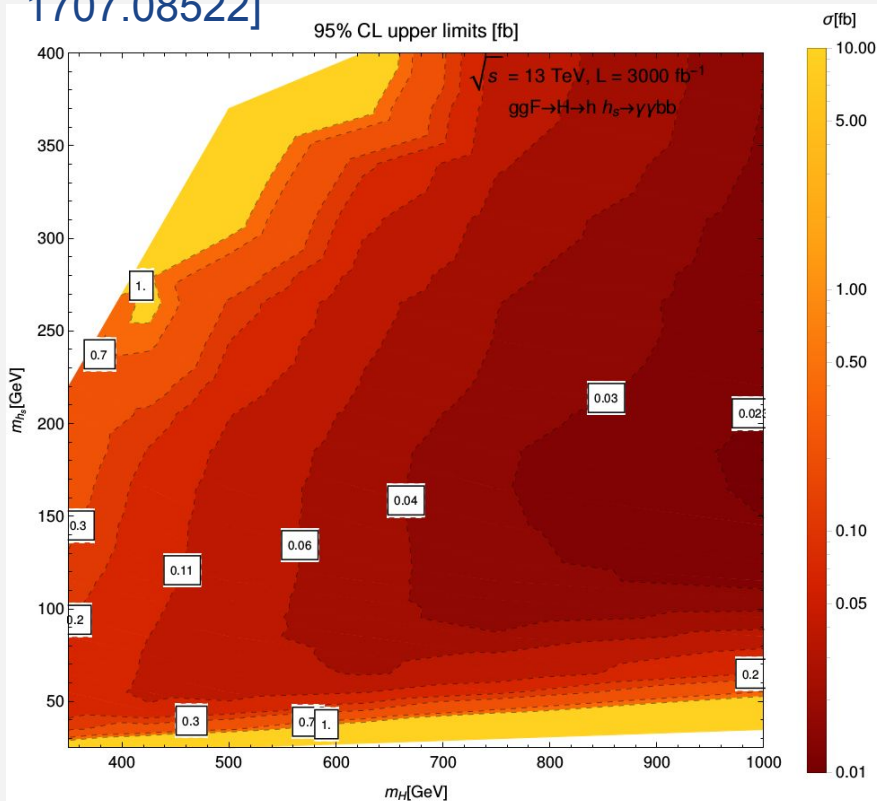
$$(A \rightarrow h_{125} a)$$

$$(A \rightarrow Z h)$$

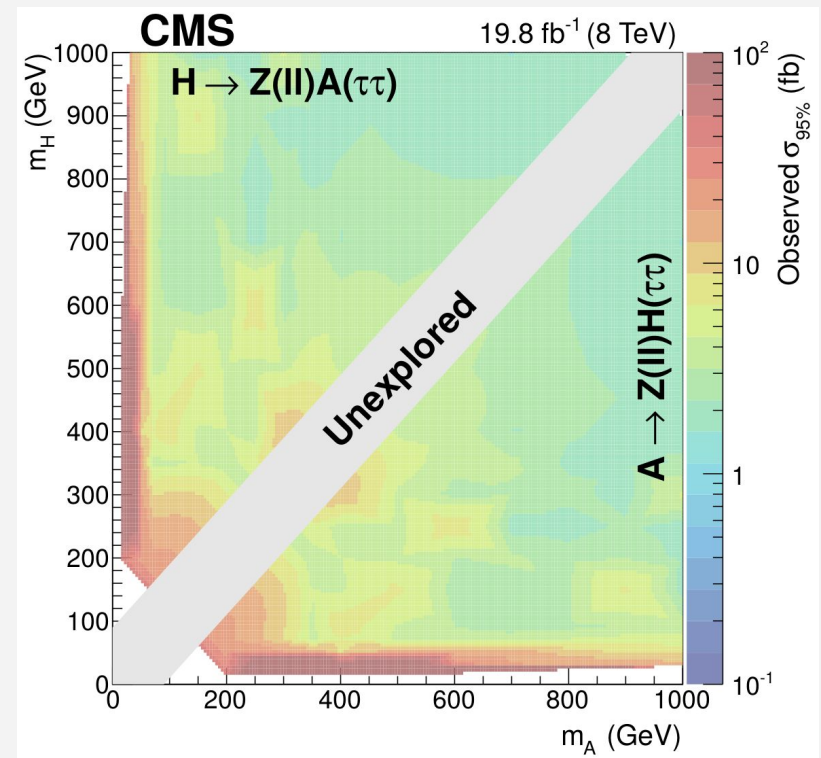
LHC Reach: SM decays of additional light states

Higgs + visible reach in
 $bbbb$ and $bb\gamma\gamma$ final states:

[Ellwanger, Rodriguez-Vazquez
1707.08522]



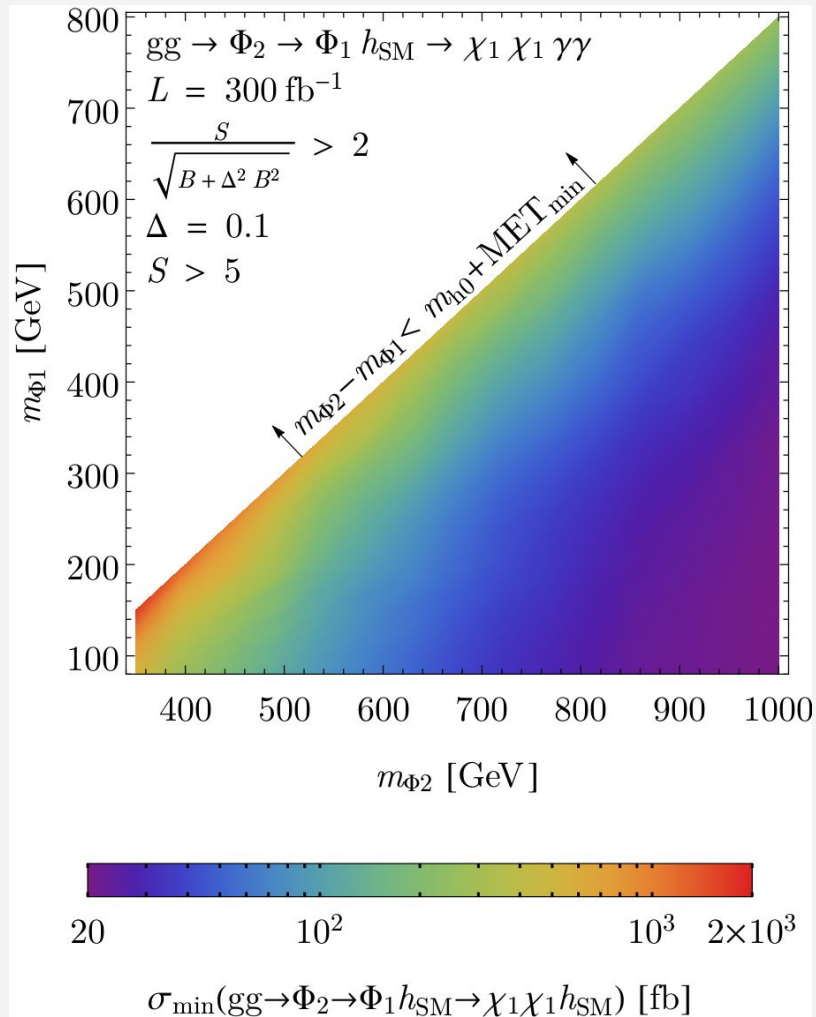
Z + visible reach in
 $\ell\ell bb$ and $\ell\ell\tau\tau$ final states
extrapolated from CMS limit
[CMS 1603.02991]



LHC Reach: MET (DM) decays of additional light states

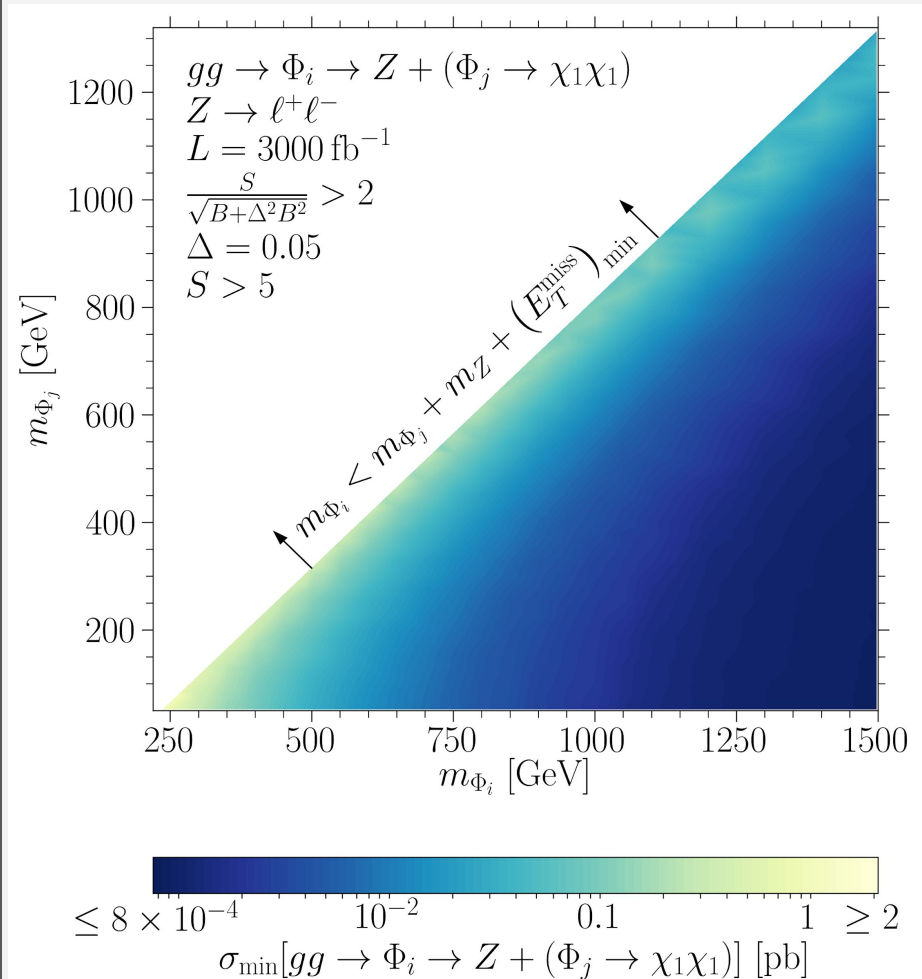
mono-Higgs reach:

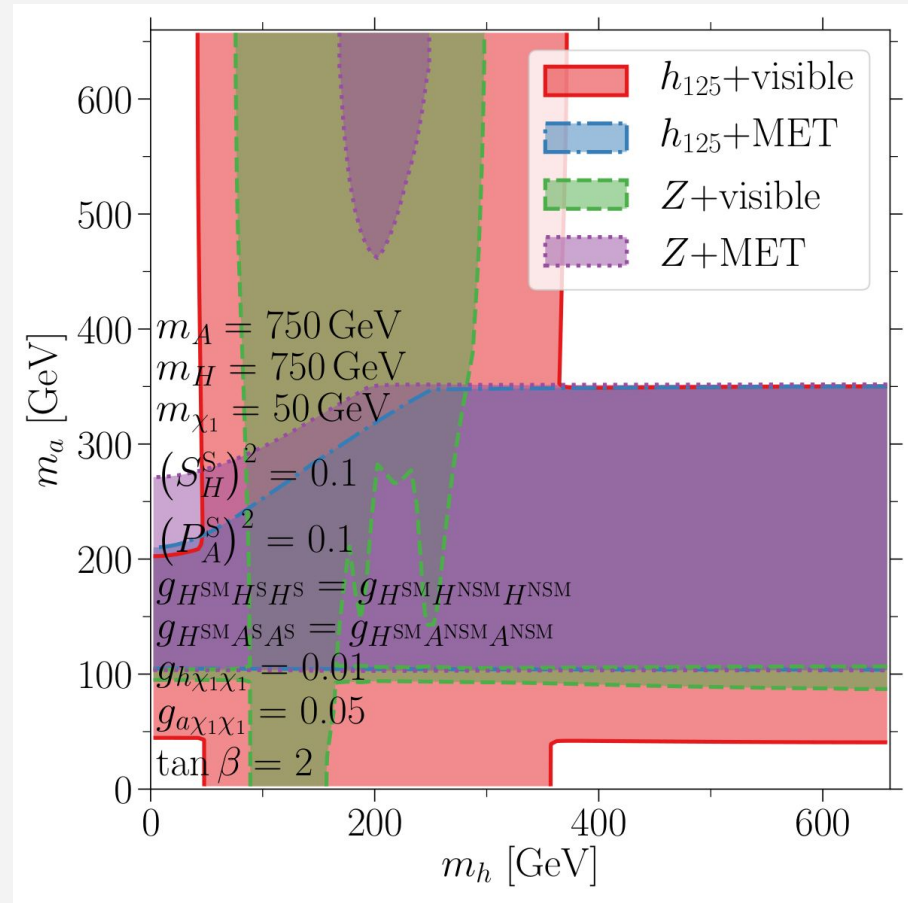
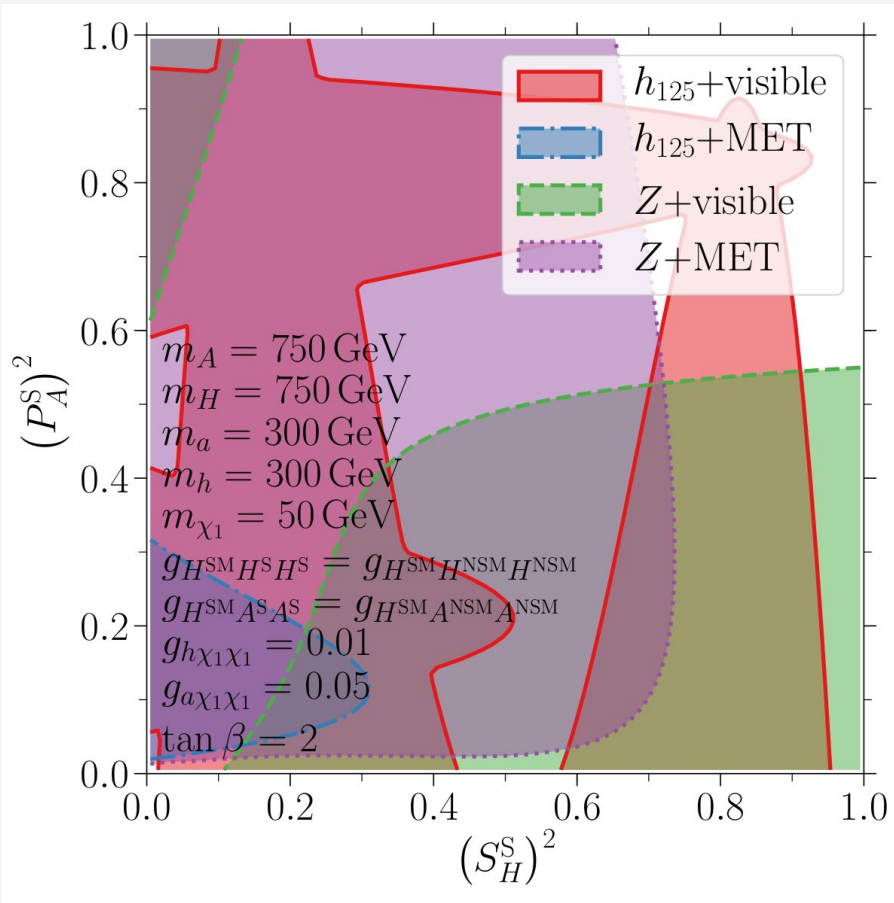
[SB, Freese, Shah, Shakya 1703.07800]



mono-Z reach

[SB, Shah 1807.XXXXXX]





Next-to-Minimal Supersymmetric Standard Model

- MSSM particle content + chiral superfield (singlet under SM)
- 2HDM+S Higgs sector, but only 6 free parameters (after fixing $v = 174$ GeV)

$$W \supset \lambda \hat{S} \hat{H}_u \cdot \hat{H}_d + \frac{\kappa}{3} \hat{S}^3$$

Acceptable pheno of 125 GeV Higgs
most easily obtained for $\lambda \sim 0.6$.
Alignment and ~ 125 GeV at tree level!

[Carena+ 1510.09137]

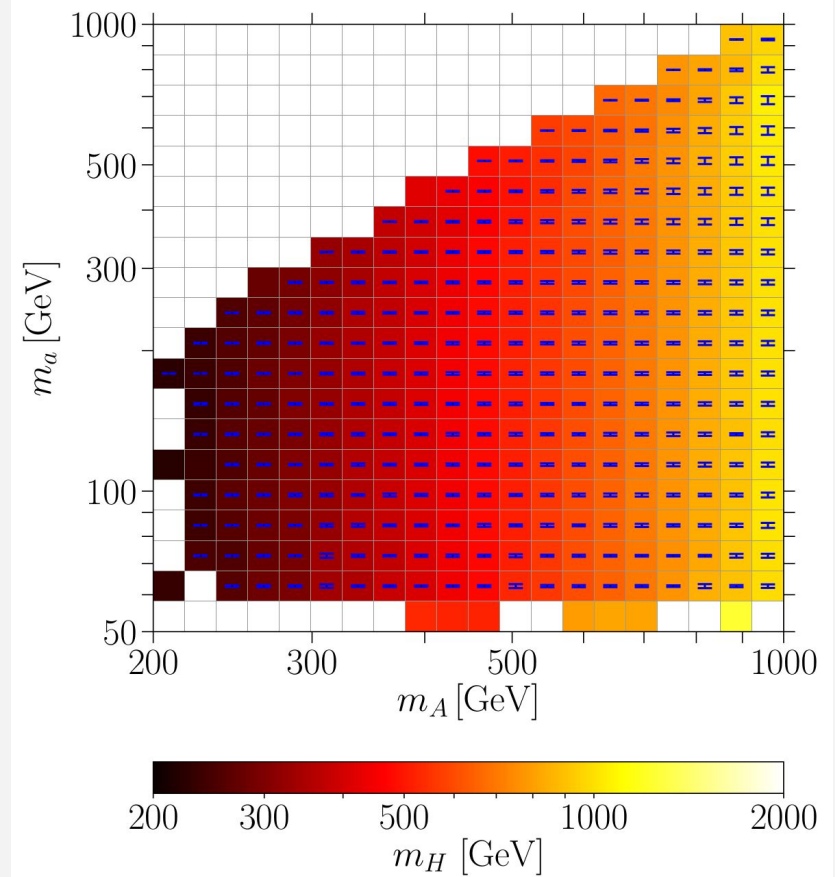
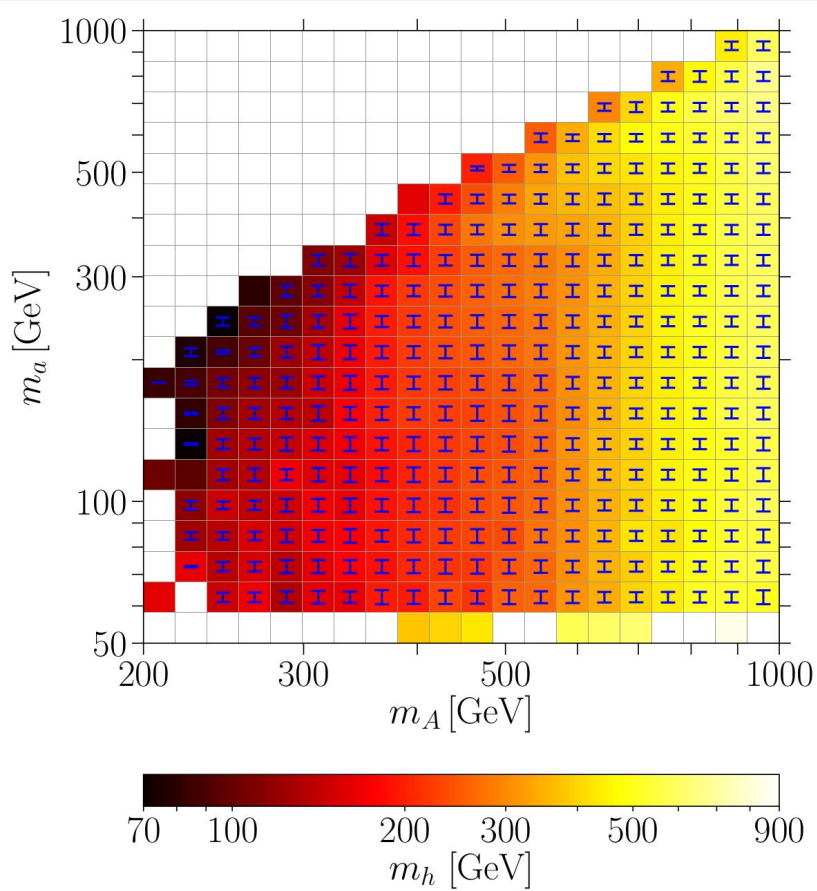
[SB, Freese, Shakya, Shah 1703.07800]

- Correlations of masses in Higgs & Neutralino sector allows for effective re-parametrization in terms of $\{\tan(\beta), m_A, m_a, P_A^S; \lambda, \kappa\}$

[SB, Shah, Freese 180X.XXXXX]

Correlation of Higgs Masses

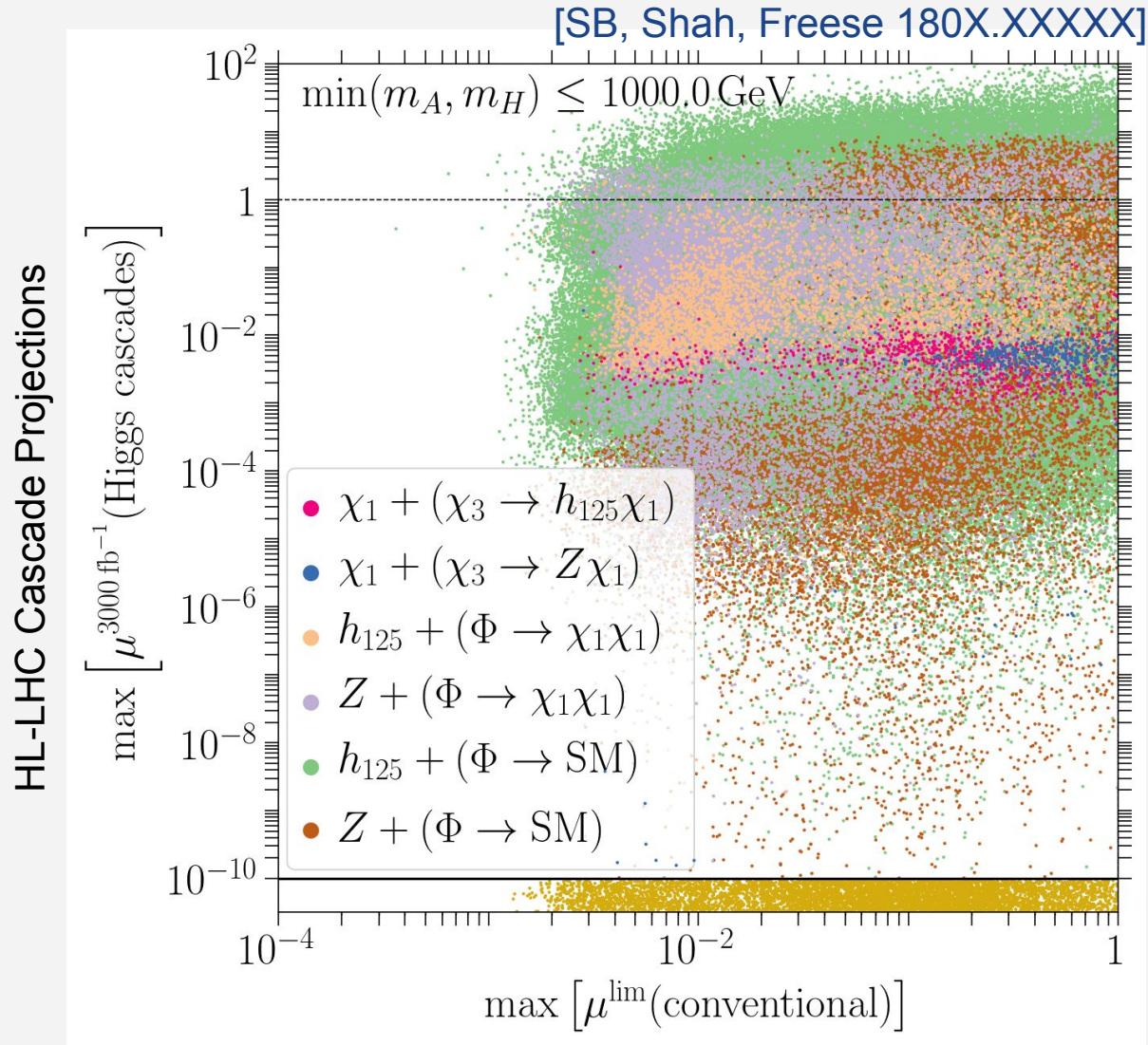
[SB, Shah, Freese 180X.XXXXX]



$$m_h^2 \simeq \frac{\kappa\mu}{\lambda} \left(A_\kappa + \frac{4\kappa\mu}{\lambda} \right) + \lambda^2 v^2 s_{2\beta}^4 \frac{M_A^2}{4\mu^2} - \frac{\lambda\kappa v^2}{2} s_{2\beta} (1 + 2c_{2\beta}^2) - \kappa^2 v^2 \frac{\mu^2}{M_A^2} c_{2\beta}^2 ,$$

$$m_a^2 \simeq 3\kappa v^2 \left[\frac{3}{2} \lambda s_{2\beta} - \left(\frac{1}{\lambda} \frac{\mu A_\kappa}{v^2} + 3\kappa \frac{\mu^2}{M_A^2} \right) \right] .$$

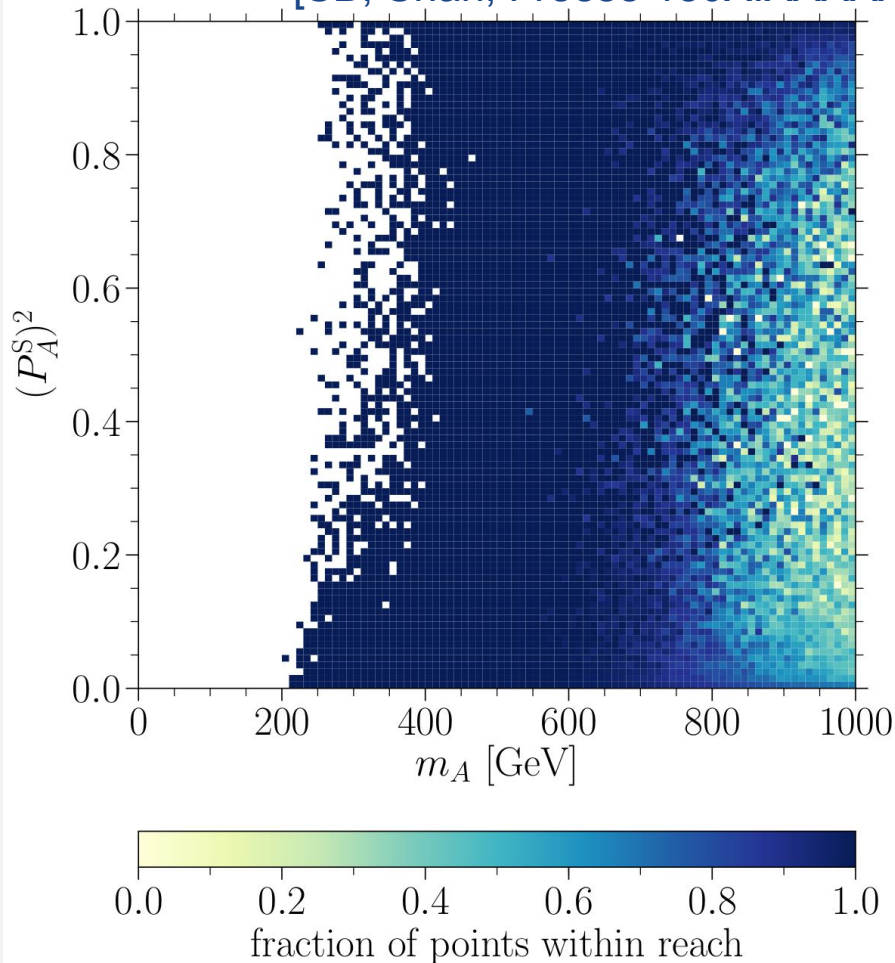
Power of Cascades + Conventional Searches



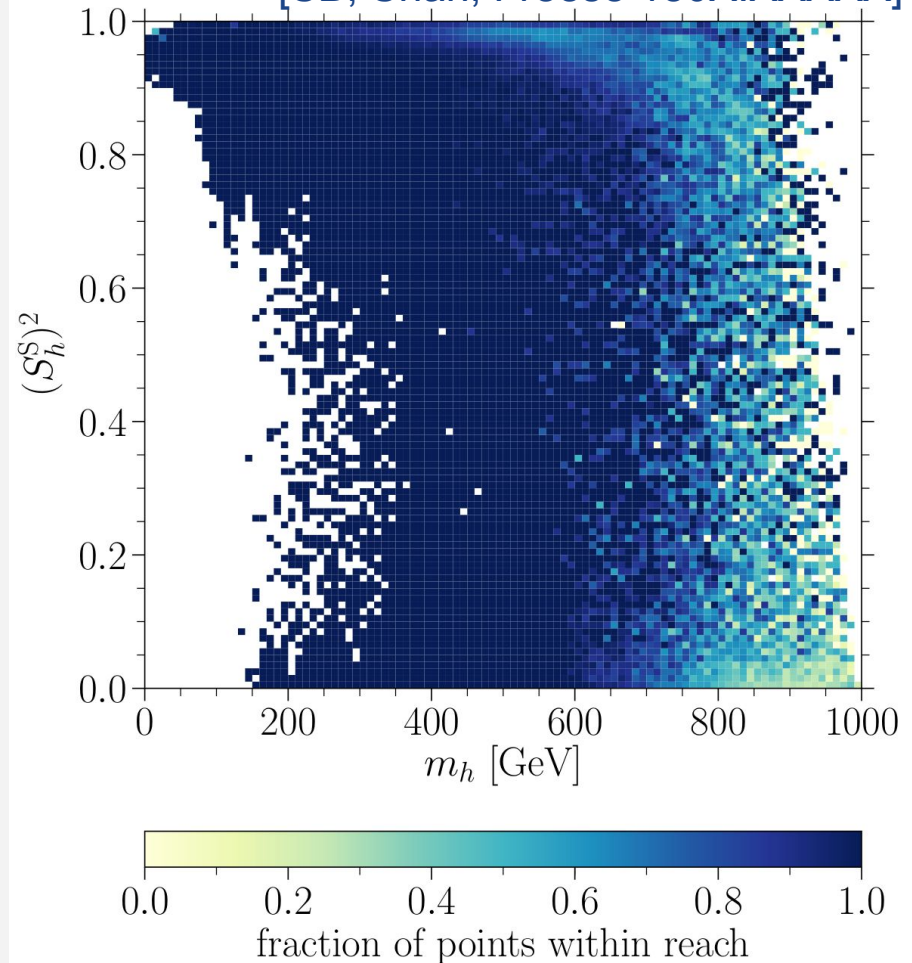
Current Bounds from conventional $pp \rightarrow H/A \rightarrow \text{SM SM}$ searches

Coverage of NMSSM parameter space

[SB, Shah, Freese 180X.XXXXXX]



[SB, Shah, Freese 180X.XXXXXX]



[Trusting LHC-experiments to improve conventional limits by factor 100
(30 -> 3000/fb + better bkg rejection) and Cascade projections by factor 10]

Conclusions

- Seemingly complicated Higgs sectors can be described effectively by approximate alignment + physically intuitive parameters such as masses and mixing angles
- Higgs Cascades provide a powerful tool to constrain models with Higgs sectors larger than 2HDM
 - Remain effective in large m_A - low $\tan(\beta)$ regime
 - Need to consider all Higgs Cascade modes to cover parameter space
 - In combination with conventional searches, much progress to complete coverage of the NMSSM parameter space with a Higgs sector below 1 TeV can be made!
- In case of discovery, measuring the effects of the singlet is crucial from differentiating the 2HDM+S (NMSSM) from a 2HDM (MSSM)

EXTRA SLIDES

	“standard”	“light subset”
$\tan \beta$	[1; 5]	[1; 5]
λ	[0.5; 2]	[0.5; 1]
κ	[−1; +1]	[−0.5; +0.5]
A_λ	[−1; +1] TeV	[−0.5; +0.5] TeV
A_κ	[−1; +1] TeV	[−0.5; +0.5] TeV
μ	[−1; +1] TeV	[−0.5; +0.5] TeV
M_{Q_3}	[1; 10] TeV	[1; 10] TeV

Table 1. NMSSM parameter ranges used in `NMSSMTools` scans.

- 10^8 points for each parameter set
- All sfermion mass parameters (except stop) set to 3 TeV
- $M_1 = M_2 = 1$ TeV , $M_3 = 2$ TeV

LHC searches used to constrain parameter scan

decay channel	NMSSM Higgs tested	Reference $\sqrt{s} = 8 \text{ TeV}$	Reference $\sqrt{s} = 13 \text{ TeV}$
$H \rightarrow \tau^+ \tau^-$	h_i, H_3, A_1, A_2	[52–54]	[55–58]
$H \rightarrow b\bar{b}$	h_1, H_3, A_1, A_2	–	[59]
$H \rightarrow \gamma\gamma$	h_i, H_3, A_1, A_2	[60–62]	[63–66]
$H \rightarrow ZZ$	h_1, H_3	[67]	[68–76]
$H \rightarrow WW$	h_i, H_3	[77–79]	[80–85]
$H \rightarrow h_{\text{SM}} h_{\text{SM}} \rightarrow b\bar{b} \tau^+ \tau^-$	h_i, H_3	[86–88]	[89–91]
$H \rightarrow h_{\text{SM}} h_{\text{SM}} \rightarrow b\bar{b} \ell \nu_\ell \ell \nu_\ell$	h_i, H_3	–	[92, 93]
$H \rightarrow h_{\text{SM}} h_{\text{SM}} \rightarrow b\bar{b} b\bar{b}$	h_i, H_3	[94, 95]	[96–99]
$H \rightarrow h_{\text{SM}} h_{\text{SM}} \rightarrow b\bar{b} \gamma\gamma$	h_i, H_3	[100, 101]	[102–104]
$A \rightarrow Z h_{\text{SM}} \rightarrow Z b\bar{b}$	A_1, A_2	[105, 106]	[107, 108]
$A \rightarrow Z h_{\text{SM}} \rightarrow Z \tau^+ \tau^-$	A_1, A_2	[86, 105]	–
$h_{\text{SM}} \rightarrow AA \rightarrow \tau^+ \tau^- \tau^+ \tau^-$	A_1, A_2	[109]	–
$h_{\text{SM}} \rightarrow AA \rightarrow \mu^+ \mu^- b\bar{b}$	A_1, A_2	[109]	–
$h_{\text{SM}} \rightarrow AA \rightarrow \mu^+ \mu^- \tau^+ \tau^-$	A_1, A_2	[109]	–
$h_{\text{SM}} \rightarrow AA \rightarrow \mu^+ \mu^- \mu^+ \mu^-$	A_1, A_2	–	[110]
$A/H \rightarrow Z h_i / A_1$	$A_2 / H_3, h_i / A_1$	[44]	–

Table 2. Direct Higgs searches at the LHC used for this work. $h_i = h_2$ (h_1) if the (second) lightest scalar is SM-like.