SUSY 2018
Barcelona, Spain
July 23-27, 2018

Searches for Dark Matter mediators with the ATLAS detector

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on behalf the ATLAS collaboration



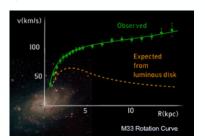


Dark matter: observations

Evidence of dark matter

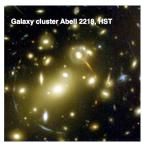
From astrophysics and cosmology observations at different scales, all consistent

Galaxy rotation



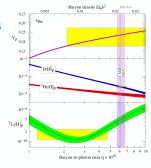
Galaxy clusters via Xrays and gravitational lensing, collisions



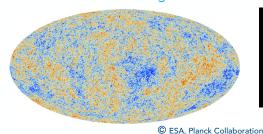




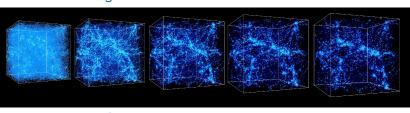
Nucleosynthesis



Cosmic microwave background



Large scale structure formation



© simulations were performed at the <u>National Center for Supercomputer</u>

Applications by Andrey Kravtsov (The University of Chicago) and <u>Anatoly Klypin</u>
(New Mexico State University).

Visualizations by Andrey Kravtsov.



(Schramm & Turner 1998).

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Dark matter: how to detect it?

Indirect detection

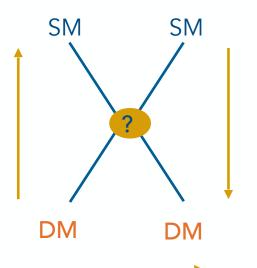
 Search for charged cosmic rays, gamma rays or neutrinos



© NASA / Sonoma State University, Aurore Simonnet

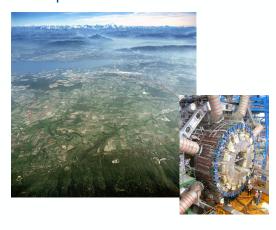
Direct detection

 Use scattering of DM on a nucleus



Collider search

 Produce DM particles from SM particles collisions



→ hyp: the dark matter is made up Weakly Interacting Massive Particles (WIMPS)



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Dark matter search at LHC

Search for particles from (UV) complete theories

- → simulate particles decays, dark matter reconstructed as missing transverse energy
 - Supersymmetry, Extra dimensions , Little Higgs, ...

Use of effective field theory (EFT)

→ many theories show common low energy behaviour, EFT allows more general searches



- describes new interactions with few operators
- → focus of LHC run 1 analyses

- → Advantages:
 - model independent
 - Allows to translate LHC results into direct search frames (with some care on the hypothesis)
- → Limitations:
- EFT valid only if M >> q → Run 1 limits M ~1 TeV => should not use E > 1 TeV
- Loose correlations that can be used in a complete theory



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Dark matter search at LHC: simplified models

Simplified models

- In between EFT and complete theory: add a single DM candidate (Dirac fermion) and a mediator
- → Allow to relax the q² limit but more model dependent
- → Allow to use other signatures to probe mediator and thus constrain the model



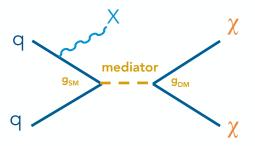
Common model and scenarios

- ATLAS/CMS + theory Dark Matter forum defined the DMSimp model (arXiv:1507.00966)
- Recommendations for benchmark scenarios (arXiv:1703.05703)
- Madgraph implementation (LO/NLO)

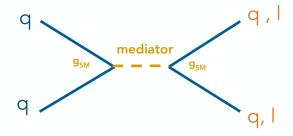
Dark matter search at LHC: simplified models

Simplified models

- In between EFT and complete theory: add a single DM candidate (Dirac fermion) and a mediator
- → Allow to relax the q² limit but more model dependent
- → Allow to use other signatures to probe mediator and thus constrains the model



→ MET + X signatures were the focus of Michaela Queitsch-Maitland's talk



- → di-fermion signatures are the focus of this talk
- Mediator could have spin 0 or 1, charged or neutral
 - recent results about spin 1, neutral mediator will be shown here



Simplified model: neutral (axial-)vector mediator

Model:



$$\begin{split} \mathcal{L}_{\text{vector}} &= g_{\text{q}} \sum_{q=u,d,s,c,b,t} Z'_{\mu} \bar{q} \gamma^{\mu} q + g_{\chi} Z'_{\mu} \bar{\chi} \gamma^{\mu} \chi \\ \mathcal{L}_{\text{axial-vector}} &= g_{\text{q}} \sum_{q=u,d,s,c,b,t} Z'_{\mu} \bar{q} \gamma^{\mu} \gamma^{5} q + g_{\chi} Z'_{\mu} \bar{\chi} \gamma^{\mu} \gamma^{5} \chi \end{split}$$

Scenarios:

 Chosen to show the complementarity of the DM production analyses (mono X) and the mediator-to-visible analyses (di X)

- Free parameters: $m(\chi)$, m(med), g_{DM} and g_{q} , g_{l}
- Minimal width computed according to couplings and considered particles mass
 - mediator decays considered = ones strictly necessary to maintain model self-consistency

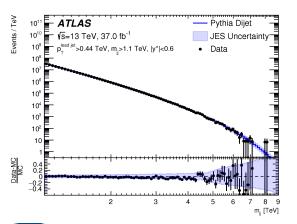
Scenarios	9 _q	9 _{DM}	9 _l
V1: vector model with only couplings to quarks	0.25	1.0	0.
V2: vector model with small couplings to leptons	0.1	1.0	0.01
A1: axial-vector model with only couplings to quarks	0.25	1.0	0.
A2: axial-vector model with equal coupling to quarks & leptons	0.1	1.0	0.1

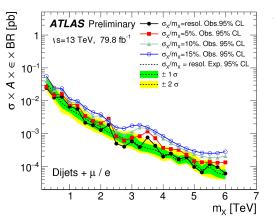


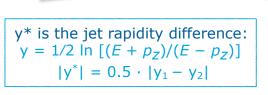
Di-jet: high mass

Analysis

- 2 small R-jets j selected with p_T > 440 GeV and 60 GeV
 - $|y^*|$ < 0.6 to reduce background from QCD process → m(jj) > 1.1 TeV
- Background obtained fitting the m(jj) distribution
- BumpHunter [Arxiv: 804.03496] algorithm used to quantify statistical significance of any local excess/deficit
- → No deviation from SM observed, limit set on cross-section and coupling to quarks



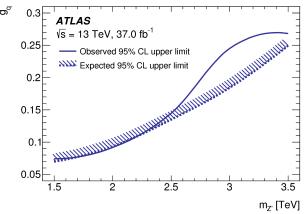




 $\sqrt{s} = 13 \text{ TeV}, L = 36.1 \text{ fb}^{-1}$

+ Phys. Rev. D 96 (2017) 052004

+ EXOT-2016-21





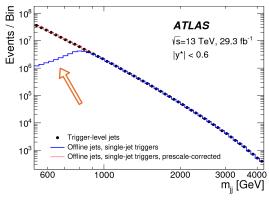
Di-jet: how to get sensitive to lower masses

Di-jet analysis uses single jet trigger: lowest unprescaled ATLAS trigger has p_T threshold at 380 GeV

→ restrain sensitivity of analysis on m_{ii} above 1.1 TeV

To recover sensitivity at lower mass, 2 strategies:

Trigger level analysis (TLA)







- Trigger on an other object: use di-jet + ISR topology
 - ISR = jet or photon, boost the di-jet resonance

→ gives access to lower mass range at the price of lower cross-section





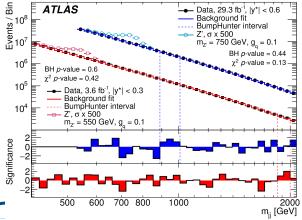
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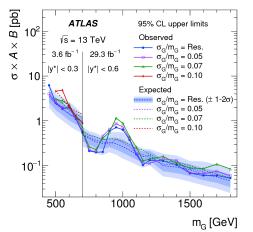
Di-jet: trigger level analysis

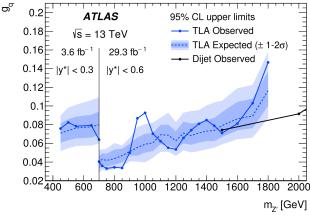
 \sqrt{s} = 13 TeV, L = 29.3 fb⁻¹

- ► EXOT-2016-20
- + arXiv:1804.03496

- → Analysis done at HLT trigger level
 - 2 (calibrated) jets selected with $p_T > 185$ GeV and 85 GeV (minimal info stored: 0.5% of an event)
 - 2 rapidity regions defined:
 - $_{-}$ ly*l < 0.6 for 700 GeV < m_{ii} < 1800 GeV for L1 trigger with E_T > 100 GeV threshold
 - $|y^*|$ < 0.3 for m_{ii} > 450 GeV for L1 trigger with E_T > 75 GeV threshold
 - sliding window fit, no deviation wrt SM observed, limits set









Di-jet (large R-jet) + ISR

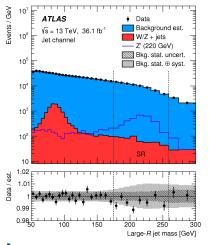
- $\sqrt{s} = 13 \text{ TeV}, L = 29.3 \text{ fb}^{-1}$
- + EXOT-2017-01
- + arXiv:1801.08769

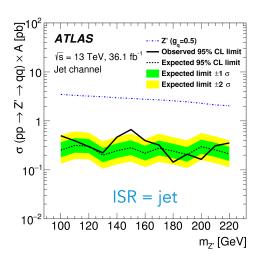
Di-jet + ISR

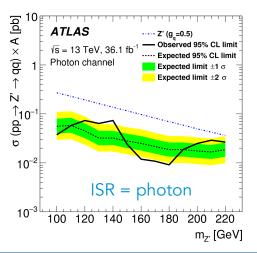
• Previous analysis done at $\sqrt{s} = 13$ TeV with L = 15.5 fb⁻¹ ATLAS-CONF-2016-070

Boosted di-jet +ISR

- → use high p_T ISR to lower trigger threshold for the event, ISR = jet or photon
 - at the price however of smaller cross-section
- → use boosted topology: select a single large-radius jet (consistent with di-jet substructure) recoiling against a hard jet or photon from ISR
- ISR jet (photon): large R-jet $p_T > 450$ (200) GeV, ISR $p_T > 450$ (155) GeV







ISR

ISR



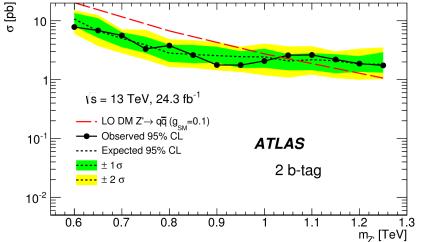
Di-bjet

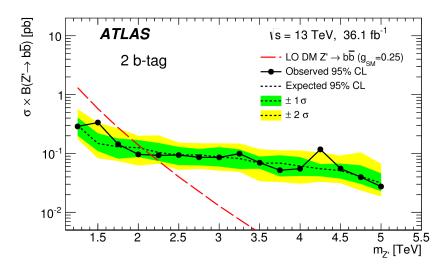
\sqrt{s} = 13 TeV, L = 36.1-24.3 fb⁻¹

- + EXOT-2016-33
- + arXiv:1805.09299

Analysis

- Low and high mass regions probed
 - \rightarrow high mass: m(jj) > 1.2 TeV, p_T(j) > 430 GeV, 80 GeV, using usual jet trigger
 - → low mass: 570 GeV < m(jj) < 1.5 TeV, using di-jet trigger (p_T threshold 150 and 50 GeV) with online b-tagging; $p_T(j) > 200$ GeV
- To reduce background and enhance s-channel process, cut on rapidity difference $y^* = (y_1 y_2)/2$ between jets: $|y^*| < 0.8$ (0.6) for high (low) mass



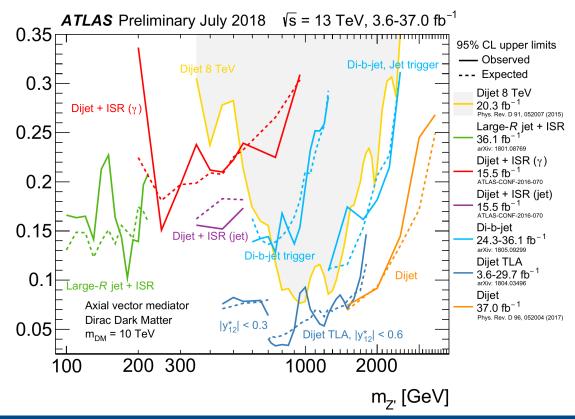


Di-jet analysis complementarity

 \sqrt{s} = 8-13 TeV, L = 3.6-37 fb⁻¹
+ ATLAS EXOTIC SUMMARY PLOTS

13

→ Limits on the DM of mediator coupling to quarks





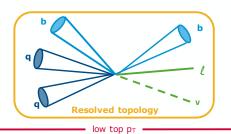
Top quark pair

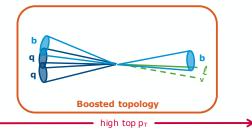
\sqrt{s} = 13 TeV, L = 36 fb⁻¹

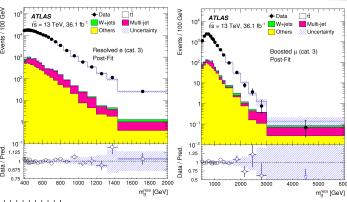
- + EXOT-2015-04
- + Eur. Phys. J. C 78 (2018) 565

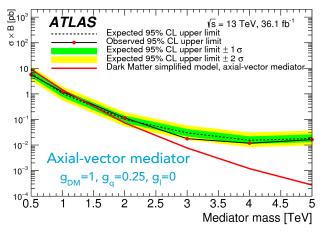
Analysis

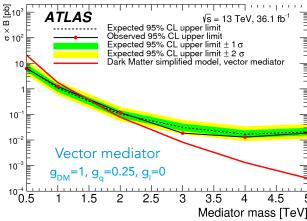
lepton+jets final state, resolved and boosted regimes











- → will be interesting to look at (pseudo-)scalar mediators
 - Interference!
 - done at 8 TeV:
 - + EXOT-2016-04/
 - Phys. Rev. Lett. 119 (2017) 191803



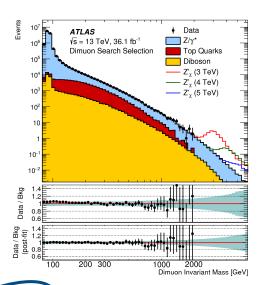
Di-lepton analysis

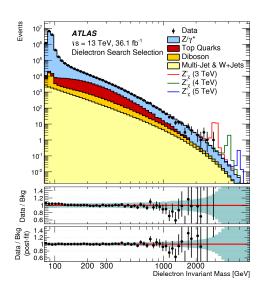
 \sqrt{s} = 13 TeV, L = 36 fb⁻¹

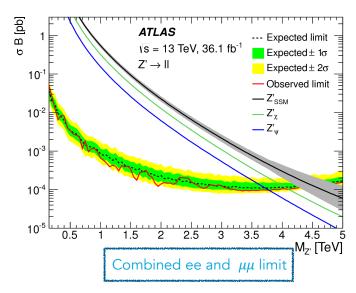
- + EXOT-2016-05
- + JHEP 10 (2017) 182

Analysis

- Select 2 same flavour leptons (opposite sign for muon channel)
 - $E_T(e) > 30 \text{ GeV}, p_T(\mu) > 30 \text{ GeV}$
- → Clean final state and much lower trigger threshold than for di-jet analyses



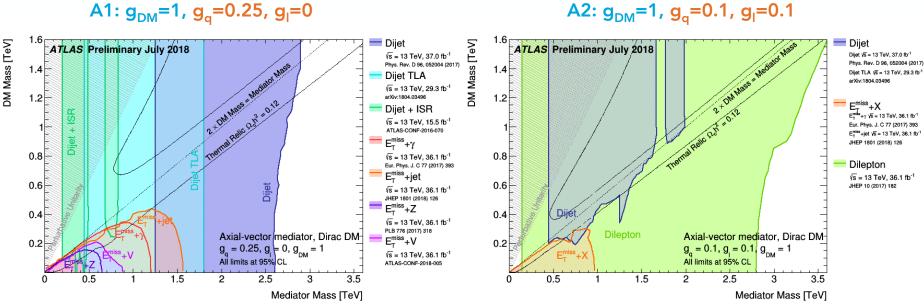




Complementarity with mono X + MET analyses

 \sqrt{s} = 8-13 TeV, L = 3.6-37 fb⁻¹
+ ATLAS EXOTIC SUMMARY PLOTS

- → Exclusions from mono-jet, mono-photon and di-jets for axial-vector coupling
 - Note: coupling dependence is important!
 - Resonance search independent of mediator mass for $m_{med} < 2 \cdot m_{DM}$

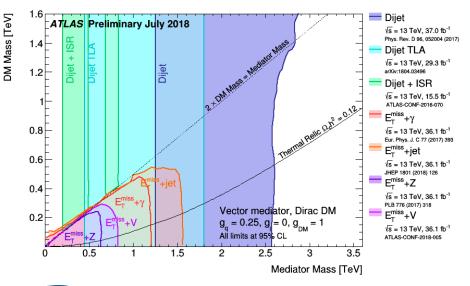


Complementarity with mono X +MET analyses

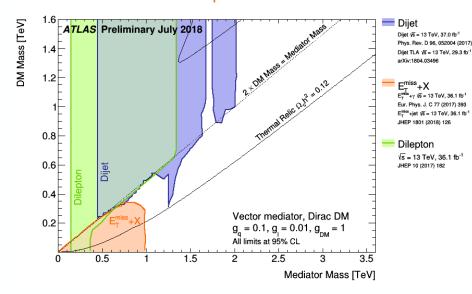
 \sqrt{s} = 8-13 TeV, L = 3.6-37 fb⁻¹
+ ATLAS EXOTIC SUMMARY PLOTS

- → Exclusions from mono-jet, mono-photon and di-jets for vector coupling
 - Note: couplings dependence is important!

V1:
$$g_{DM} = 1$$
, $g_{a} = 0.25$, $g_{l} = 0$



V2:
$$g_{DM} = 1$$
, $g_{q} = 0.1$, $g_{l} = 0.01$





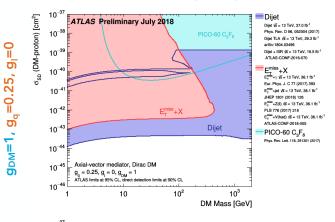
Complementarity with direct detection

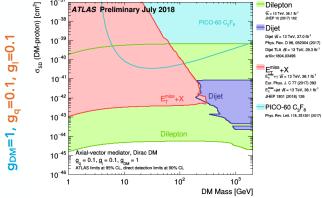
 \sqrt{s} = 8-13 TeV, L = 3.6-37 fb⁻¹ * ATLAS EXOTIC SUMMARY PLOTS

Beware! validity only for this choice of model...

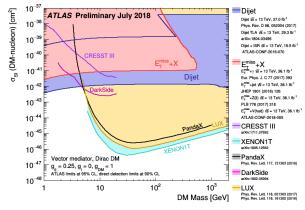
Note: limits at 95% CL for ATLAS, 90% CL for direct detection...

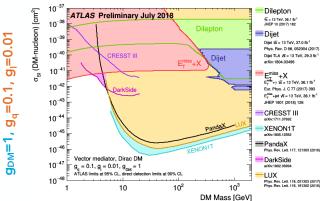
Axial-Vector mediator / Spin dependent





Vector mediator / Spin independent





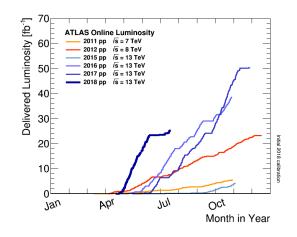


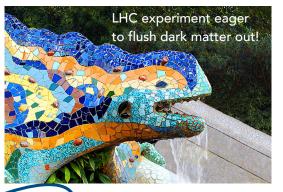
g₁=0.01

Summary and perspectives

Extensive program of dark matter search with the ATLAS experiment

- In the context of simplified model, analyses looking for DM mediator decaying into SM particles are powerful and show nice complementarity with DM MET signatures
- No sign of dark matter though, in the 36 fb-1 run 2 data scrutinised ... but much more to come (x4 data on tape)!
- These results are also complementary to dark matter direct searches





Much more to explore...

- Spin 1 charged mediator, FCNC, spin 0 mediator in simplified model with 1 mediator and 1 DM particle
- Extension of 2 HDM with a pseudo-scalar: very rich phenomenology







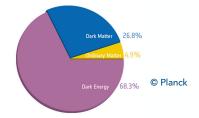
TO GO FURTHER...



Dark matter: what do we know about?

Properties

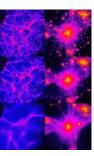
- It makes up 85% of the matter in the Universe
 - It is massive



- It interacts weakly with ordinary matter (at least through gravitation)
 - It is neutral
- It interacts weakly with itself



- It is stable (a minima very long-lived, order of the age of the universe)
 - □ Ruled out SM Z and Higgs
 - Need a symmetry to prevent it to decay ex T-parity
- It is "cold" ie non relativistic
 - ruled out SM neutrinos (also not enough massive)



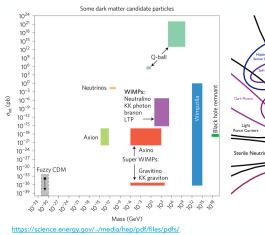
© Cold, Warm, and Hot dark matter simulations, credit ITP, University of Zurich.

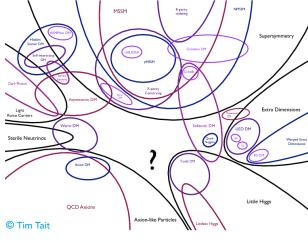
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Dark matter: which candidates? Associated theories?

Candidates

- WIMPs = Weakly Interacting Massive Particles
 - → WIMP "miracle" : weak cross-section + particle mass ~1 TeV ~ relic density
 - Susy neutralinos
 - Kaluza-Klein photon
- Very Weak Interacting Massive Particles
 - gravitinos
 - Axions: to solve the strong CP problem, unstable but long lived
 - Sterile neutrinos: to explain neutrino masses
 - Kaluza Klein gravitons
 - •
- Could be also a more complex sector with
 several particles and interactions





Theories

- Supersymmetry
 - Symmetry: R-parity
- Extra dimensions
 - Symmetry : KK parity

- Little Higgs
 - Symmetry: T-parity
- QCD axions
- •

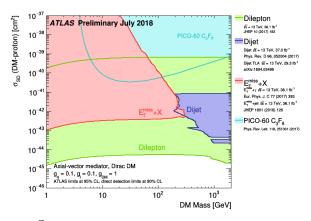
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Comparison with direct detection

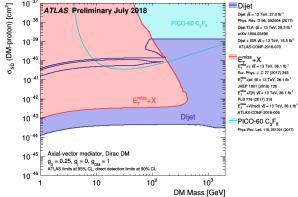
DM-proton

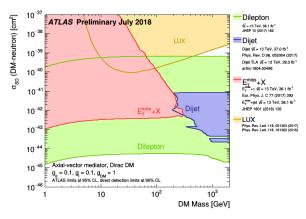
DM-neutron

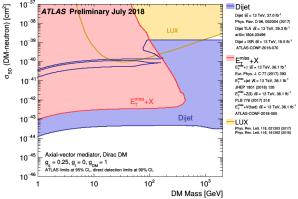
Axial-Vector mediator $g_q = 0.1$, $g_l = 0.1$, $g_{DM} = 1$



Axial-Vector mediator $g_a=0.25$, $g_i=0$, $g_{DM}=1$



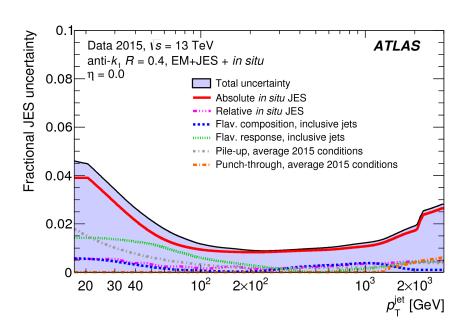


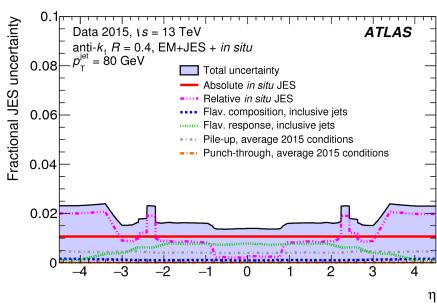


Jet Energy Scale



- + PERF-2016-04
- + Phys. Rev. D 96 (2017) 072002





- ▶ Jet energy scale uncertainty is 1% for central jets with pT of 500 GeV and grows to 3% for jets with pT of 2 TeV
- Di-jet mass resolution is 2.4% and 2.0% for dijet masses of 2 TeV and 5 TeV, respectively



Dijet analyses

Di-jet invariant mass fit

Used for previous analyses:

$$f(z) = p_1(1-z)^{p_2} z^{p_3} z^{p_4 \log z}$$

- where $z = m_{jj}/\sqrt{s}$ and the p_j are parameters
- New: sliding-window fitting technique fitting only restricted regions of the spectrum
 - windows ~ 1/2 total number of bins ⇒ whole signal contained
 - nore flexibility, 3 parameters sufficient
 - new method validated

TLA di-jet analysis

- Sliding window fit
- $f(x) = p_1(1 x) p^2 x^{p3+p4 \ln x+p5 \ln x2}$ or $f(x) = p_1 x^{-p2} e^{-p3x-p4x2}$

