Application of Global Bayesian Analysis to Large and small systems

Weiyao Ke

Duke Univeristy

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Presentation is based on recent works by Jonah Bernhard, Scott Moreland, Steffen Bass and WK, arXiv:1804.06469, arXiv:1806.04802, PRC 96, 044912 (2017).

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Questions

How to determine the properties of a physics system with observations of the final states?
 How accurate are our statements about its properties?



A schematic example: extract one parameter from one observation

A system with a state x_{true} leads to an observation y. $y = f_{phy}(x_{\text{true}})$.

- Measurement of y comes with uncertainty $\rightarrow \text{Exp}$ uncertainty $y_{exp} \pm \sigma_{exp}$.
- Our model f_M of the true function f_{phy} is imperfect \rightarrow Theory uncertainty σ_{th} .

A Bayesian point of view

• What is the probability of x to be the truth,

$$P(x = x_{true} | f_M, y_{exp}) \propto \text{Likelihood}(f_M(x) = y_{exp} | f_M, x) \times \text{Prior}(x)$$
$$\log \text{Likelihood} = -\frac{(f_M(x) - y_{exp})^2}{2\sigma^2} + C, \sigma^2 = \sigma_{exp}^2 + \sigma_{th}^2 \quad (+\cdots)$$

• Bayes' theorem P(A|B) = P(B|A)P(A)/P(B).



A much more complex situation for heavy ion physics

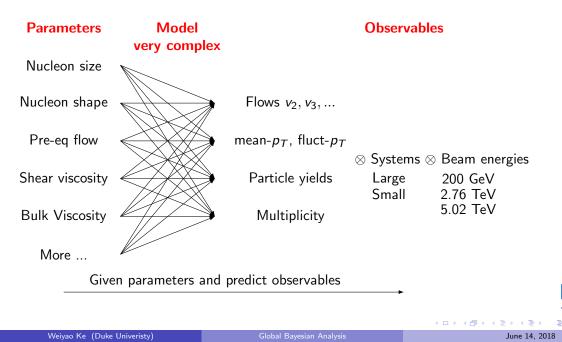
Parameters		Observable	es
Nucleon size			
Nucleon shape	Flows <i>v</i> ₂ , <i>v</i> ₃ ,		
Pre-eq flow	mean- p_T , fluct- p_T	⊗ Systems ⊗	Beam energies
Shear viscosity	Particle yields	Large Small	200 GeV 2.76 TeV
Bulk Viscosity	Multiplicity	5.02 TeV	

More ...



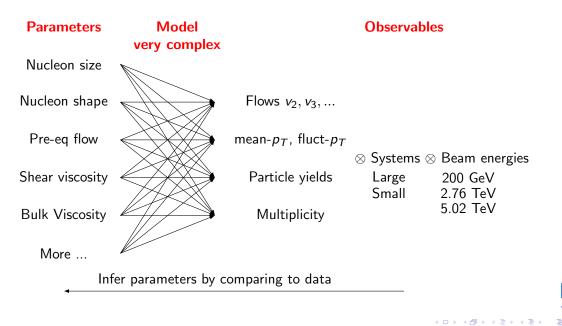
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A much more complex situation for heavy ion physics

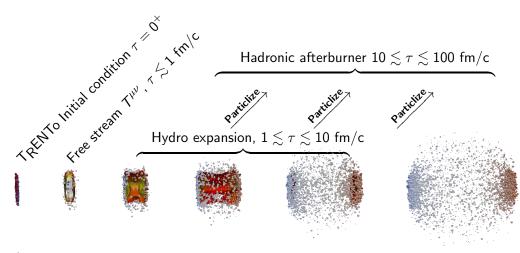


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A much more complex situation for heavy ion physics



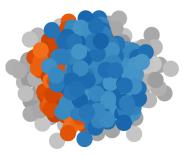
Work through a real example: extract η/s , ζ/s from AA collisions @LHC **Physics models: a multistage hybrid simulation.**



*Thesis work by JE Bernhard arXiv:1804.06469.

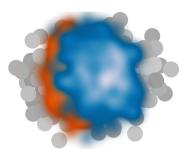


Prepare two nuclei nucleon positions sampled from a Woods-Saxon function \downarrow Determine binary collisions $P_{coll}(b) = 1 - \exp(-\sigma_{eff}T_{pp}(b))$ σ_{eff} fixed by fitting $\sigma_{DD, inel}$





Prepare two nuclei nucleon positions sampled from a Woods-Saxon function **Determine binary collisons** $P_{\rm coll}(b) = 1 - \exp(-\sigma_{\rm eff} T_{pp}(b))$ $\sigma_{\rm eff}$ fixed by fitting $\sigma_{\rm pp, inel}$ Calculate participant density $T_{A,B}(x) = \sum \gamma_i \rho(x-x_i).$ $i \in N_{\text{part}}$

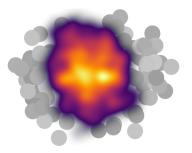




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Entropy deposition Generalized mean ansatz:

$$rac{dS}{ au dx_{\perp}^2 d\eta_s} \propto \left(rac{T_A^p + T_B^p}{2}
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Prepare two nuclei nucleon positions sampled from a Woods-Saxon function \downarrow Determine binary collisons $P_{coll}(b) = 1 - \exp(-\sigma_{eff} T_{pp}(b))$ σ_{eff} fixed by fitting $\sigma_{pp, inel}$ \downarrow Calculate participant density

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Entropy deposition Generalized mean ansatz:

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ight)^{1/p}.$$

Features

- Interpolates a family of entropy deposition mappings at mid-rapidity.
- Gaussian proton shape.
- Parameters: nucleon width (w), fluctuation (σ), entropy deposition (p).



Physics model: pre-equilibrium free streaming and hydro

Pre-equilibrium stage

- Before $\tau_{\rm fs}$ free stream initial energy density and match it to hydrodynamic $T^{\mu\nu}$ at $\tau = \tau_{\rm fs}$.
- Shall use more realistic model in the future such as an effective kinetic theory.
- A tunable parameter τ_{fs} .

Hydrodynamics

- 2+1D relativistic viscous hydrodynamics (OSU hydro).
- EoS: interpolate between HRG and Lattice results (HotQCD).
- Parametrization of η/s and ζ/s

$$\frac{\eta}{s} = (\eta/s)_{\min} + (\eta/s)_{\text{slope}} (T - T_c) \left(\frac{T}{T_c}\right)^{(\eta/s)_{\text{curv}}}, \frac{\zeta}{s} = \frac{(\zeta/s)_{\max}}{\left(\frac{T - T_0}{(\zeta/s)_{\text{width}}}\right)^2 + 1}$$

Particularization and hadronic rescattering

Particularization: frzout by JE Bernhard

- Hadrons are sampled from fluid cell with Cooper-Frye prescription at $T = T_{sw}$.
- Include the effect of hadron finite widths.
- Non- δf type shear and bulk corrections.
- Parameter: particularization temperature ($T_{\rm sw}$).

Hadronic rescattering

• An Ultra-relativistic Quantum Molecular Dynamics model (UrQMD).



A summary of parameters

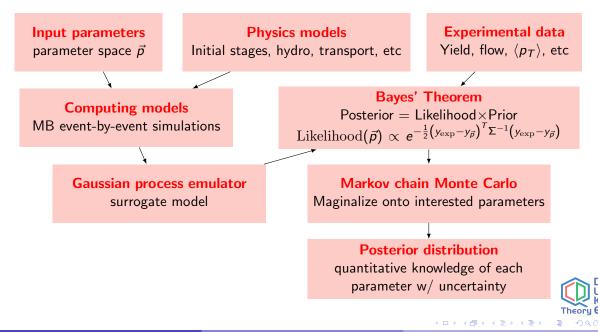
Group	Parameters	Explanation
IC	р	Entropy deposition parameter
	Norm(s)	Overall entropy normalizations
	σ fluct	Gamma fluctuation std of N-N collision
	<i>w</i> [fm]	Gaussian nucleon width
	d min [fm]	Nucleon-nucleon minimum distance in nucleus
Pre-eq	$ au_{ m fs}~[{ m fm/c}]$	Freestreaming time
QGP	η/s min	value of $\eta/s(T_c)$
	η/s slope [GeV $^{-1}$]	slope parameter of $\eta/s(T)$
	η/s curv	curvature $\eta/s(T)$
	ζ/s max	max value of $\zeta/s(T)$
	ζ/s width [GeV]	width parameter of $\zeta/s(T)$
	$\zeta/s T_0$ [GeV]	peak temperature of $\zeta/s(T)$
Particularization	T switch [GeV]	switching temperature from hydro to UrQMD.
		Theory

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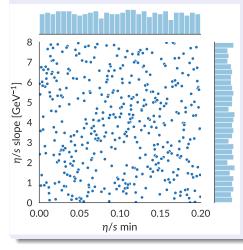
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A sophisticated Bayesian parameter estimation



Calibration to LHC 2.76 TeV and 5.02 TeV Pb+Pb data (new!)

Sample design input parameters



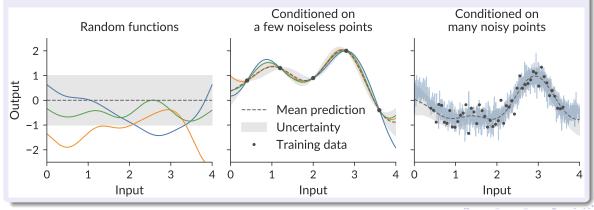
- For a high dimensional (n = 14) space, grid interpolate is impossible N ~ O(10ⁿ).
- A smarter sampling technique: Latin Hyper cube sampling.
- 1. Optimized random samples.
- 2. Maximize the mini. dist. between pairs of points.
- 3. The marginalization on any parameter is uniform.
- 4. $N \sim \mathcal{O}(10 \times n)$ for smooth function over the space.
- Perform full model calculation on design points.



Calibration to LHC 2.76 TeV and 5.02 TeV Pb+Pb data (new!)

Surrogate model: Gaussian process emulator

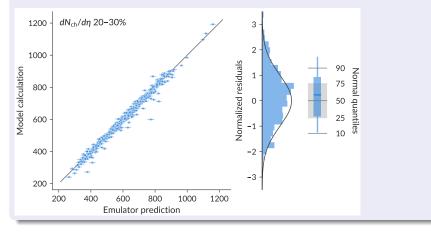
- After full model evaluation (\vec{Y}_i) on design points (\vec{x}_i) : $\vec{x}_i \longrightarrow \vec{Y}_i$.
- How to get arbitrary $\vec{Y}(\vec{x})$ for "arbitrary" parameters \vec{x} ?
- Non-parametric interpolation through Gaussian process emulator.
- $\bullet~$ A mean prediction + interpolation uncertainty.



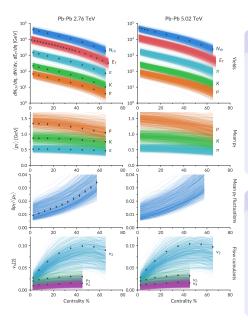
Calibration to LHC 2.76 TeV and 5.02 TeV Pb+Pb data (new!)

Validate the performance of Gaussian process emulator

- Predict observables with newly generated inputs and compare with actually calculation.
- Validate that the trained emulator predictions reproduce the dependence of \vec{Y} on \vec{x} .
- Emulator do have uncertainty!



Calibration to LHC 2.76 TeV and 5.02 TeV Pb-Pb data (New)



Calculation of design points (prior)

- 500 design points.
- Observables:

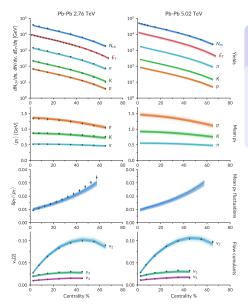
multiplicity, transverse energy, p, π, K yield and mean- p_T , mean- p_T EbE fluctuation, v_2, v_3, v_4 .

• Parameter design ranges are wide enough to spread over all data points.

The likelihood function

$$\begin{split} \Delta y &= (y_{\exp} - y_{\mathrm{p}}) \\ \log \mathrm{Likelihood} &= C - \frac{1}{2} \det \Sigma - \frac{1}{2} \left(\Sigma^{-1} \right)_{ij} \Delta y_i \Delta y_j. \\ \Sigma &= \sigma_{\mathrm{stat}}^2 + \sigma_{\mathrm{sys}}^2 + \sigma_{\mathrm{GP}}^2 + \sigma_{\mathrm{model}}^2 \end{split}$$

Calibration to LHC 2.76 TeV and 5.02 TeV Pb-Pb data (New)



After model-to-data comparison (posterior, from emulator)

- After comparing to experimental data, the emulator can predict observables from the calibrated distribution of parameters.
- Achieve a global agreement with data.



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Parameters calibration

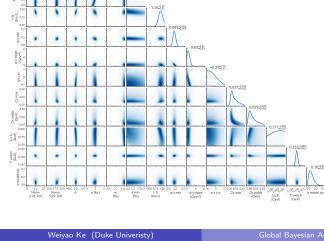
Marginalized posterior distribution

- Diagonal: single parameter distributions (integrate over N - 1).
- Off diagonal: pairwise correlations (integrate over N 2).
- Marginalization of the other N 1 parameters concentrates on the single parameter with uncertainties from other parts of the model folded in.



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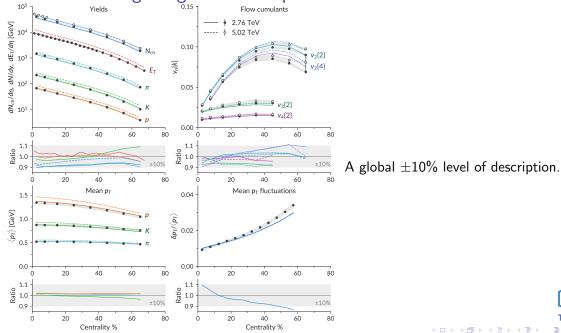
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0.006

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Calculation using a high-likelihood parameter set



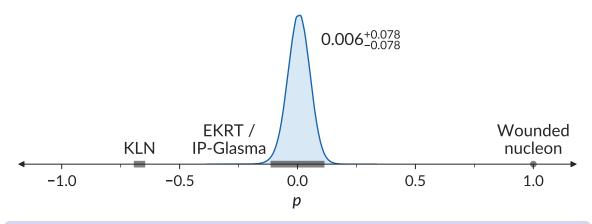
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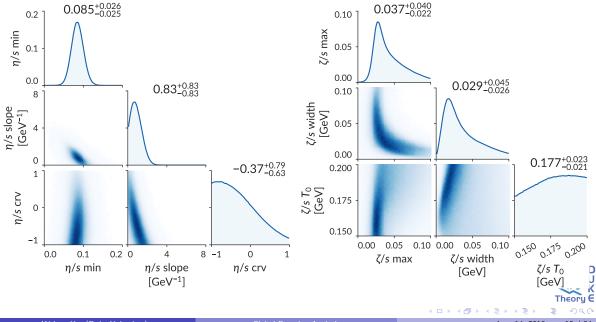
Focus on the initial condition



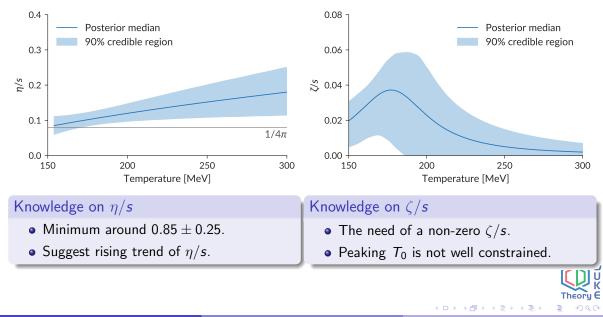
Knowledge on initial entropy deposition

• Support the eccentricity scaling ($\epsilon_{2,3}$ as functions of b) given by saturation physics based/motivated initial condition models.

Focus on transport properties



Focus on transport properties



Extend to small systems (from a final state point of view)

Collective features observed in small systems

- Long range pseudorapidity two-particle correlation.
- Finite v_2 , v_3 from multi-particle correlations.
- Competing of two pictures,
 - Initial state effect.
 - Final state effect
 - 1. Pressure driven and the formation of QGP in high-multiplicity events.
 - 2. Escape mechanism.

Study small system from a final state perspective

- Developments of T_RENTo towards a more realistic description of small system.
- Perform analysis on both pA and AA \rightarrow Can we describe both in a single framework?
- Can we study other observables that may help to distinguish IS and FS effects?



Developments in $\mathsf{T}_{\mathsf{R}}\mathsf{ENTo}$ for studying possible final state effects

The $\mathsf{T}_R\mathsf{ENTo}$ family

Original: round proton, boost-invariant



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Developments in $\mathsf{T}_{\mathsf{R}}\mathsf{ENTo}$ for studying possible final state effects

The T_RENTo family

 $\mbox{Original: round proton, boost-invariant} \quad \longrightarrow \quad \mbox{Add proton shape fluctuation}$

• A recent simultaneous calibration on both p-Pb and Pb-Pb collisions (2+1D).



Developments in $\mathsf{T}_{\mathsf{R}}\mathsf{ENTo}$ for studying possible final state effects

The TRENTo family

Original: round proton, boost-invariant \longrightarrow Add proton shape fluctuation $\downarrow \qquad \qquad \downarrow \qquad \qquad \downarrow$ Add local rapidity dependence \longrightarrow Combine both features (in progress)

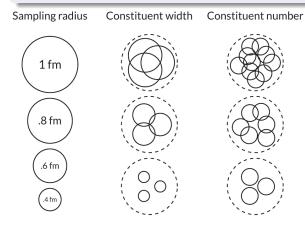
- A recent simultaneous calibration on both p-Pb and Pb-Pb collisions (2+1D).
- Opportunity of calibration using 3+1D simulations and predicts for possible final state effect in small systems.



Adding proton shape fluctuation to $\mathsf{T}_{\mathsf{R}}\mathsf{ENTo}$

Subnucleonic fluctuation

- In p-A collisions, initial geometry is sensitive to proton subnucleonic fluctuations.
- $\bullet\,$ Extend the default Gaussian round proton in ${\sf T}_{\sf R}{\sf ENTo}$ by adding subnucleon constituents.



Proton degrees of freedom

- A proton width parameter.
- Width of Gaussian constituents.
- Number of constituents.



Adding proton shape fluctuation to $\mathsf{T}_{\mathsf{R}}\mathsf{ENTo}$

Collision of Gaussian round proton

• Proton-proton collision probability at fixed b,

$$P_{
m coll}(b) = 1 - \exp(-\sigma_{
m eff}T_{
m pp}(b)), T_{
m pp}(b) = \int dx_{\perp}^2
ho(x+b/2)
ho(x-b/2)$$

•
$$\sigma_{\rm eff}$$
 fixed by requiring $\sigma_{\rm pp, inel} = \int d\vec{b}^2 P_{\rm coll}(b)$.

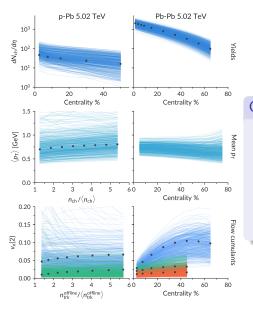
Collision of fluctuating proton

 \bullet Proton-proton ${\it P}_{\rm coll}$ determined from constituent collisions,

$$egin{aligned} & P_{ ext{coll}}(b) = \left\langle 1 - \prod_{i,j=1}^{\#} [1 - P_{ ext{coll}}(b_{ij})]
ight
angle_{x_i,x_j ext{ given } b} \end{aligned}$$



Calibration on p-Pb and Pb-Pb systems simultaneously

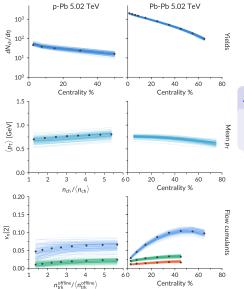


Calculation of design points (Prior)

- Observables:
 - Pb+Pb: multiplicity, v_2 , v_3 , v_4 p+Pb: multiplicity, v_2 , v_3 , mean- p_T
- The fluctuating proton can generate huge eccentricity fluctuations.
- Prior range well-cover the data from both systems.



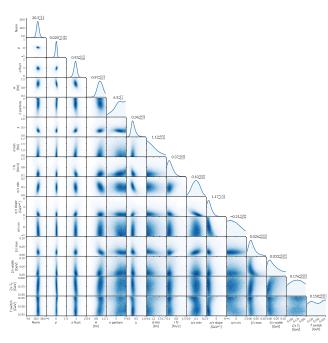
Calibration on p-Pb and Pb-Pb systems simultaneously



After compare to data (Posterior, from emulator)

• The calibrated model emulator very well predicts the observable for both p-A and A-A.





Parameters calibrated on pA, AA

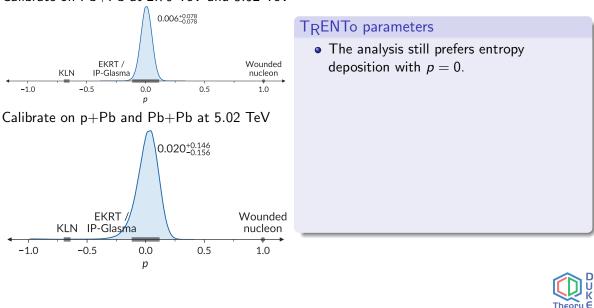
Marginalized posterior distribution

- What does it say about initial condition with proton fluctuations.
- Are the extracted transport coefficients different from using only AA data?

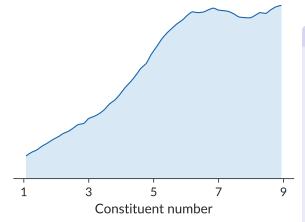


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The constrained TRENTo with proton fluctuations Calibrate on Pb+Pb at 2.76 TeV and 5.02 TeV



The constrained T_RENTo with proton fluctuations

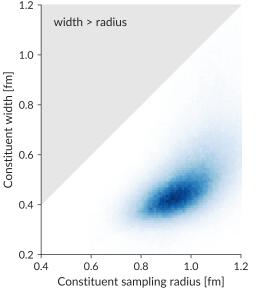


T_RENTo parameters

- The analysis still prefers entropy deposition with p = 0.
- No particular preference of # of constituents. But a round proton (n = 1 case) is disfavored.



The constrained T_RENTo with proton fluctuations

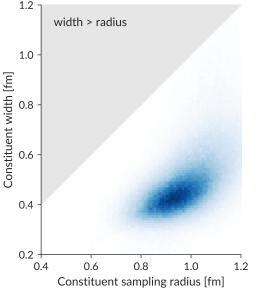


T_RENTo parameters

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- No particular preference of # of constituents. But a round proton (n = 1 case) is disfavored.
- The joint distribution of constituent width and its sampling radius within a proton is well constrained.



The constrained T_RENTo with proton fluctuations

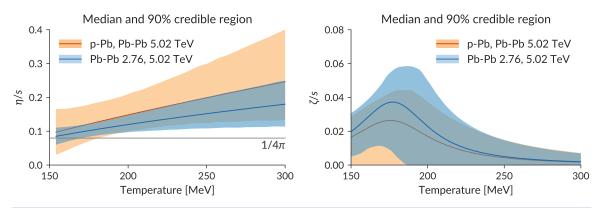


T_RENTo parameters

- The analysis still prefers entropy deposition with p = 0.
- No particular preference of # of constituents. But a round proton (n = 1 case) is disfavored.
- The joint distribution of constituent width and its sampling radius within a proton is well constrained.
- Need to check the resulting proton eccentricities.



Do we need a different $\eta/s, \zeta/s$ to describe p-A?



Comparing transport properties from two calibration.

- The η/s, ζ/s that describe pA and AA @ 5.02TeV is consistent with those describing a lot more AA observables @ 2.76 and 5.02 TeV.
- No extra handle on $\eta/s, \zeta/s$ by including pA.

How can a 3+1D simulation help us understand the nature of pA

Looking at event-plane decorrelations

- In Pb-Pb collisions, the event-planes decorrelate over pseudorapidity due to initial participant plane twisting and fluctuations.
- If medium in pA also undergoes pressure driven expansion, event-plane decorrelation should be described in the same framework as AA.
- In an initial state approach, the decorrelation has a different origin.



Extending T_RENTo (round proton version) to finite rapidity

$$\frac{dS}{dx_{\perp}^{2}dy} \propto s_{0}(\vec{x}_{\perp}, y = 0) \times f(y, x_{\perp}).$$

$$T_{R}ENTo \times rapidity profile$$
• $f(y, x_{\perp})$ has three degrees of freedom (first 3 y-moments):
mean $\mu(x_{\perp})$, std $\sigma(x_{\perp})$, skewness $\gamma(x_{\perp})$

$$f(y, x_{\perp}) \propto \mathcal{F}^{-1} \exp\left\{i\mu k - \frac{1}{2}(\sigma k)^{2} - \frac{i}{6}\gamma(\sigma k)^{3} + ...\right\}$$
• μ, σ, γ parametrized in nuclear thickness functions $T_{A}(x_{\perp}), T_{B}(x_{\perp}).$

$$\overline{\frac{\text{mean } (\mu)}{\frac{\mu_{0}}{2} \log \frac{T_{A}}{T_{B}}} \sigma_{0} \frac{\gamma_{0}(T_{A} - T_{B}) \text{ or } \gamma_{0} \frac{T_{A} - T_{B}}{T_{A} + T_{B}}}$$



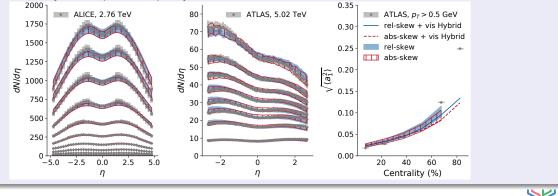
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Image: A matrix

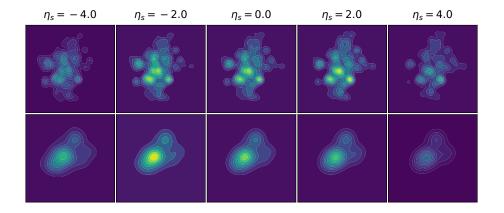
Extending T_RENTo (round proton version) to finite rapidity

Calibrated to p-Pb and Pb-Pb

- This calibration did not include proton substructure.
- Parameters constrained by charged particle pseudorapidity density of Pb-Pb and p-Pb and event-by-event psuedorapidity fluctuation in Pb-Pb.



Rapidity evolution in the presence of proton shape-fluctuation.



How much of initial geometry info gets transformed into final state particles. TRENTo \rightarrow 3+1D Free streaming \rightarrow 3+1D viscous hydrodyanmics \rightarrow UrQMD.

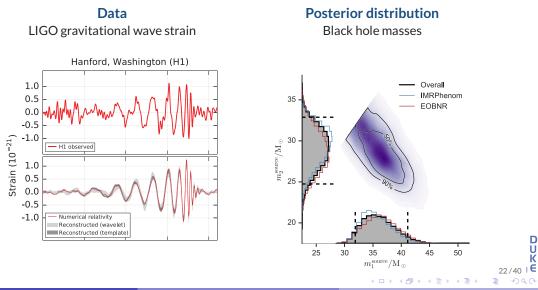


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Image: Image:

Summary

Example application: Gravitational waves





- Global Bayesian model-to-data comparison is a powerful quantitative tool to learn QGP properties from experimental data.
- It has been successfully applied to large systems to infer initial condition model and QGP transport coefficients.
- Pb-Pb and p-Pb collisions can be simultaneously described at mid-rapidity by including proton shape fluctuations.
- Looking at event-plane decorrelations in future 3+1D analysis.



Outlooks

Bayes factor

- A quantitative measurement of model performance.
- Can be used to select a better-performing model.
- Introduce natural penalty for over complex model (too many parameters).
- How to calculate:

$$\frac{P(M_1)}{P(M_2)} = \frac{\int P(\vec{p}|\mathrm{Exp}, M_1) d\vec{p}}{\int P(\vec{p}|\mathrm{Exp}, M_2) d\vec{p}}$$

• It has been applied to comparing hydro+UrQMD v.s. hydro + partial chemical equilibrium EoS.



$^{16}O + ^{16}O$

Estimate fluid cell temperature distribution at $\tau = \tau_{hvdro}$ using T_RENTo initial condition.

