

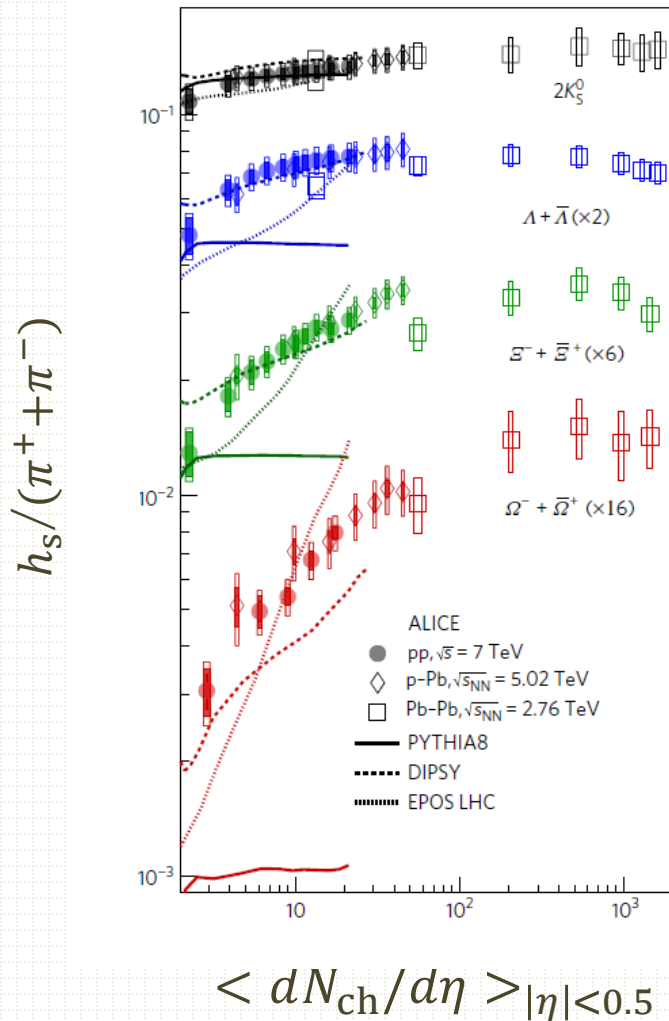
Strangeness enhancement from dynamical initialization with core-corona picture

Yuuka Kanakubo
Sophia Univ.

In Collaboration with: Michito Okai (Sophia Univ.)
Yasuki Tachibana (Wayne state Univ.)
Tetsufumi Hirano (Sophia Univ.)

Introduction & Motivation

Strangeness enhancement in small systems

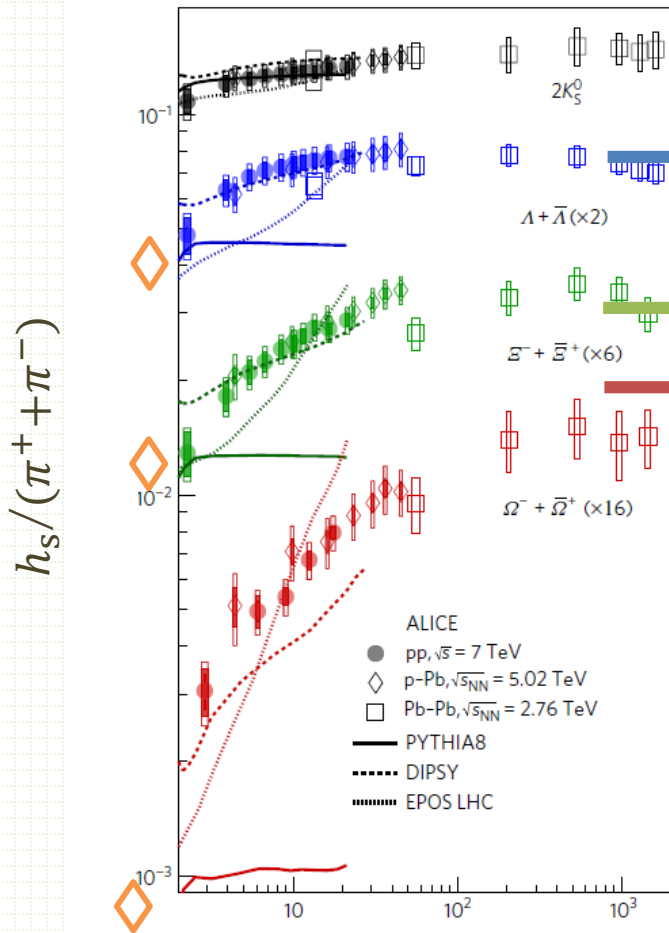


- Rapid enhancement of multi-strange hadrons relative to charged pions in small systems
- Continuous increase as a function of multiplicity towards the ratio in Pb+Pb.

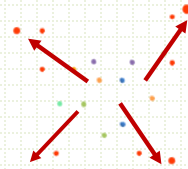


Indication of QGP creation
in high multiplicity small colliding
systems

Strangeness enhancement in small systems

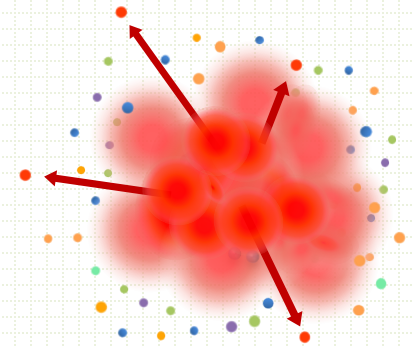


$\langle dN_{ch}/d\eta \rangle_{|\eta| < 0.5}$



low

← multiplicity →



high

◇ String fragmentation

★ Particlization from chemically equilibrated matter

A. Andronic *et al.*, (2017).

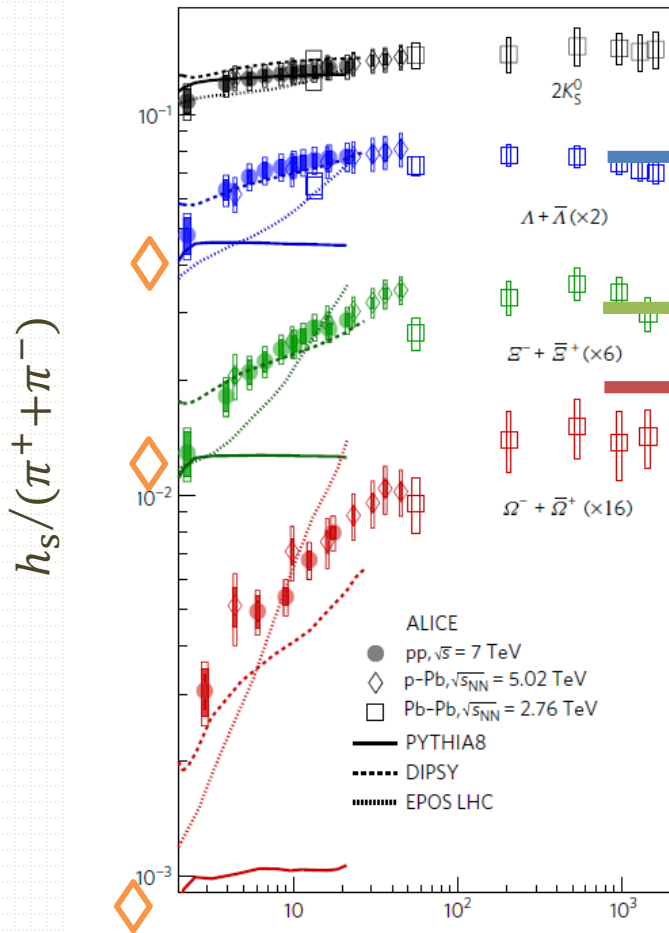
$$P \propto \exp\left(-\pi \frac{m_q^2}{\kappa}\right)$$

$$P \propto d_i \exp\left(-\frac{m_i}{T_{ch}}\right)$$

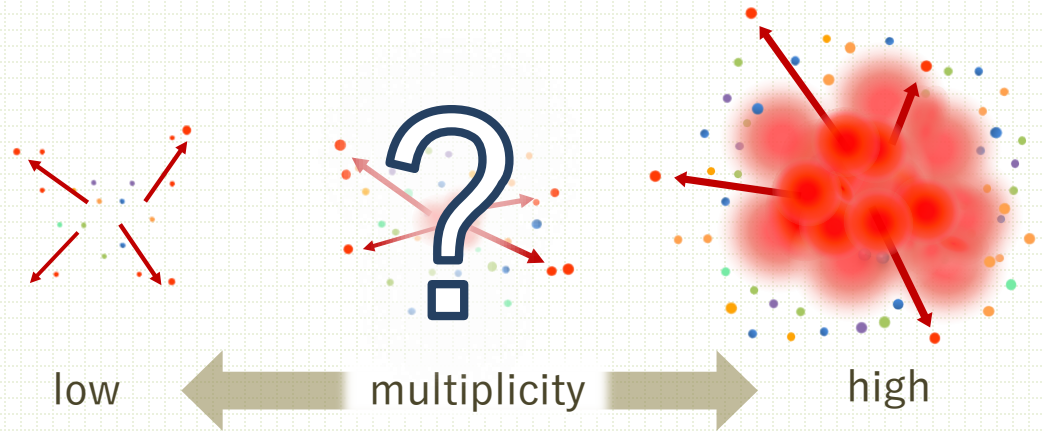
Different hadronization mechanism

➔ Two different ratios of strangeness products

Strangeness enhancement in small systems



$\langle dN_{ch}/d\eta \rangle_{|\eta| < 0.5}$



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Different hadronization mechanism

➔ Two different ratios of strangeness products

Purpose

Interpretation of the strangeness enhancement from dynamical initialization with core-corona picture



Hydrodynamic equation with source terms

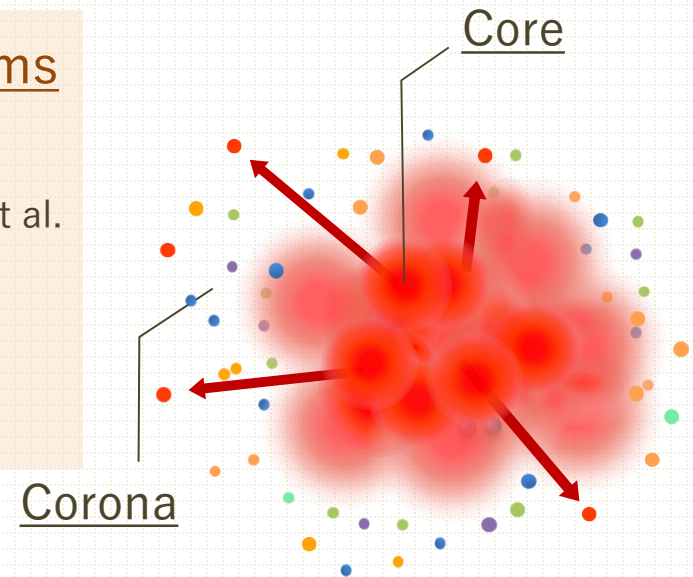
Dynamical initialization

M. Okai et al, C. Shen et al, Y. Akamatsu et al.

+

Core-corona picture

J. Aichelin et al, F. Becattini et al, T. Pierog et al



Final hadron products from
“core” and “corona” separately

→ Final strangeness ratio from superposition of
core and corona products

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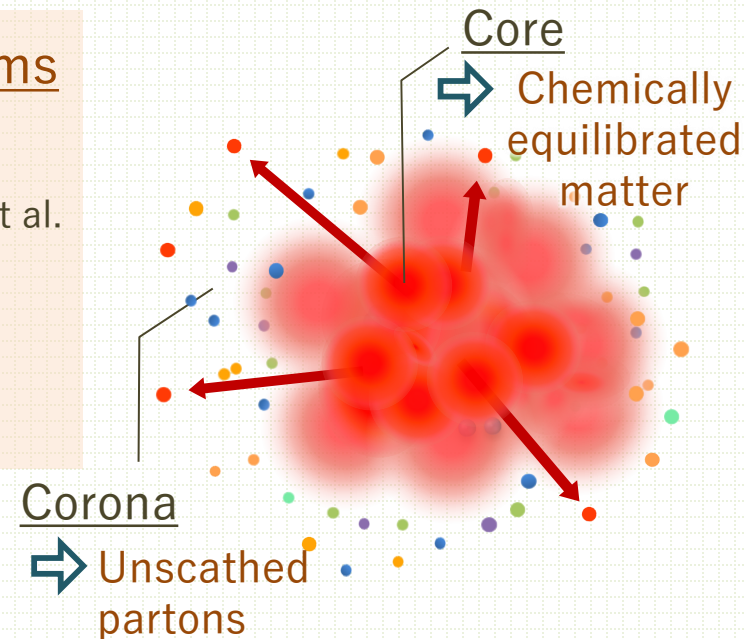
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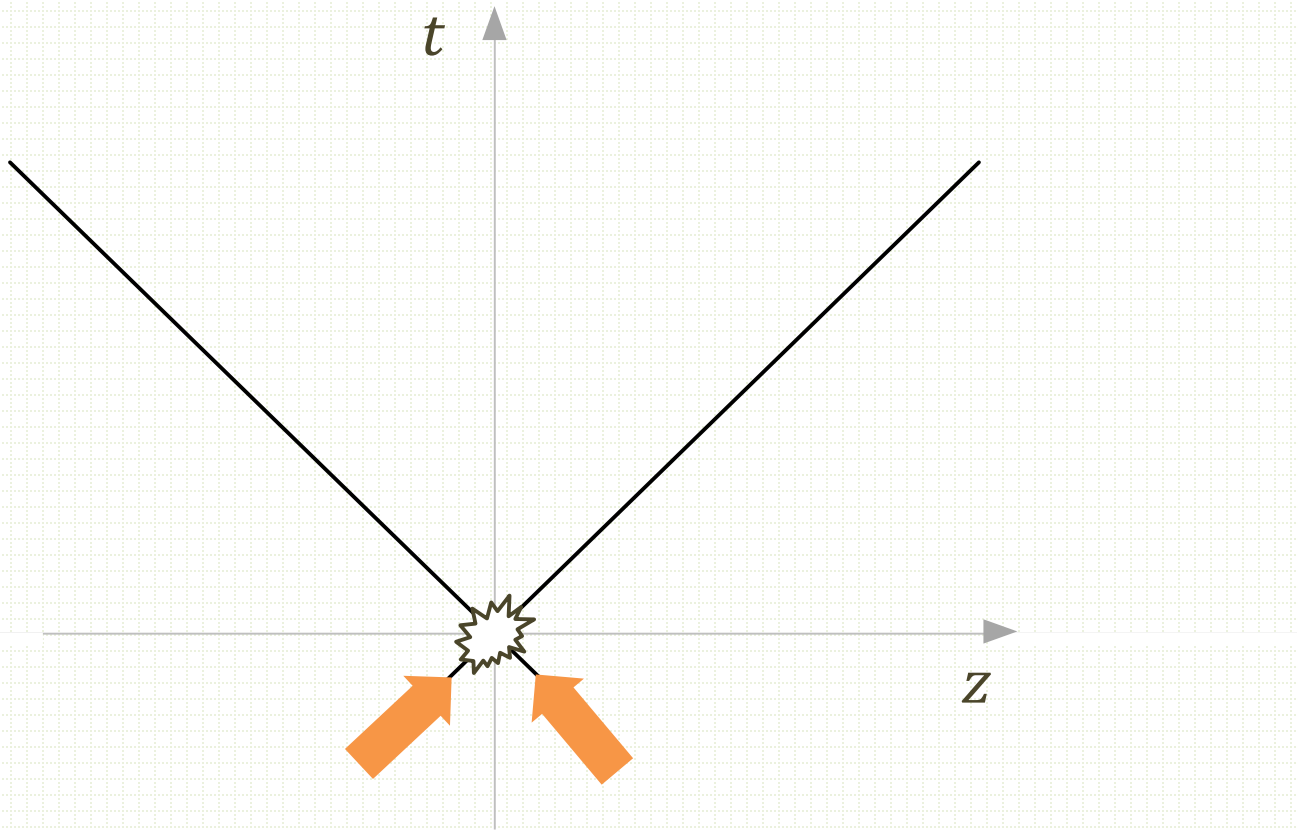
→ Final strangeness ratio from superposition of
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Model

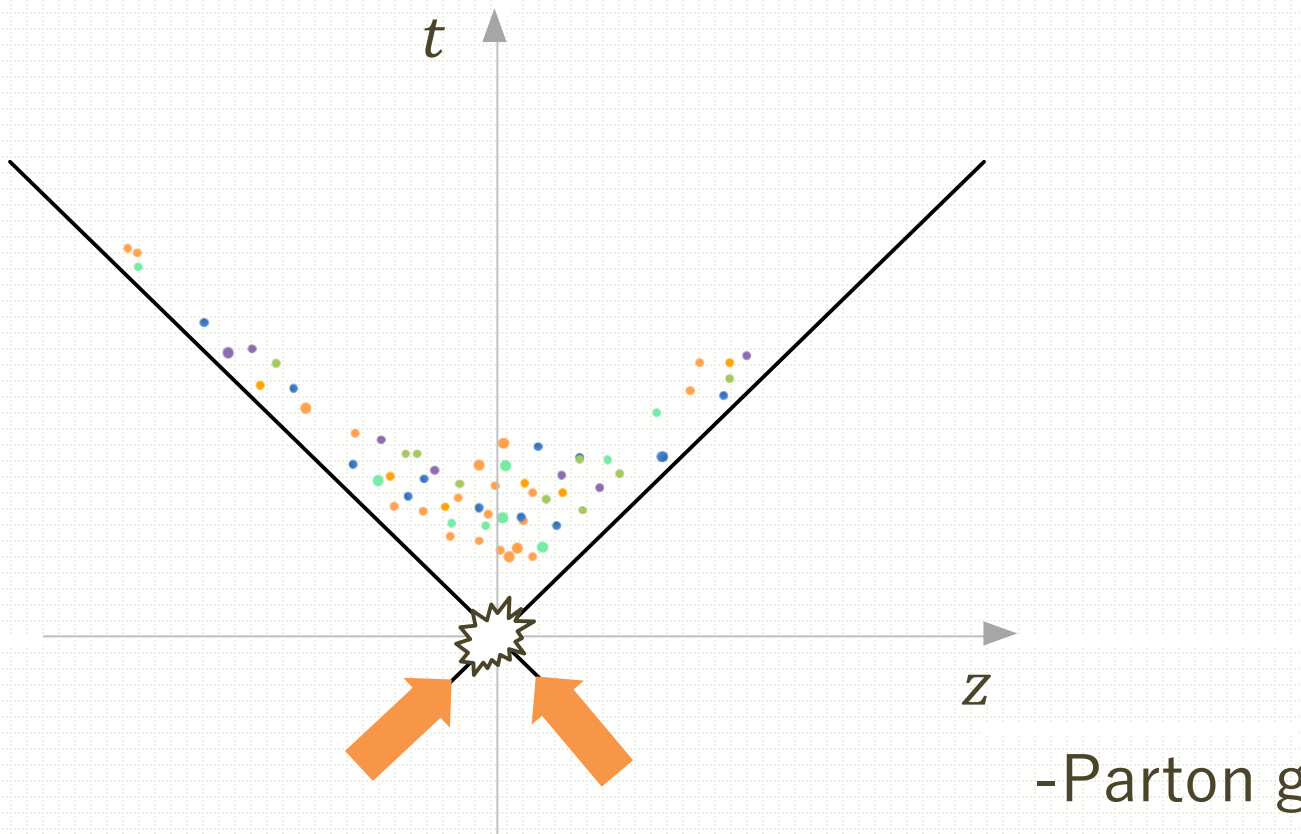
Model

Time evolution of high energy nuclear collisions



Model

Time evolution of high energy nuclear collisions



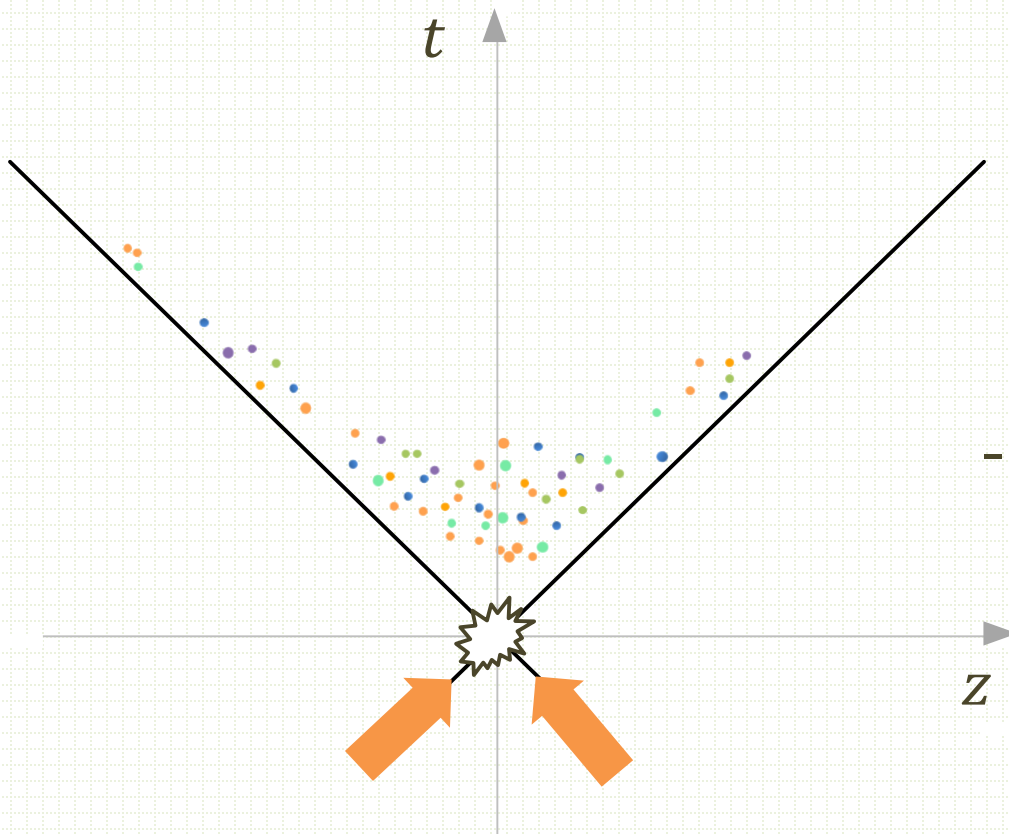
-Parton generation

PYTHIA ver 8.230

T. Sjöstrand *et al.*,
Comput. Phys. Commun. 191, 159 (2015).

Model

Time evolution of high energy nuclear collisions



-Evolution of QGP fluids

Dynamical initialization

Ideal hydro, Lattice EoS

M.Okai K.Kawaguchi, Y.Tachibana, T.Hirano,
Phys. Rev. C 95, 054914 (2017).

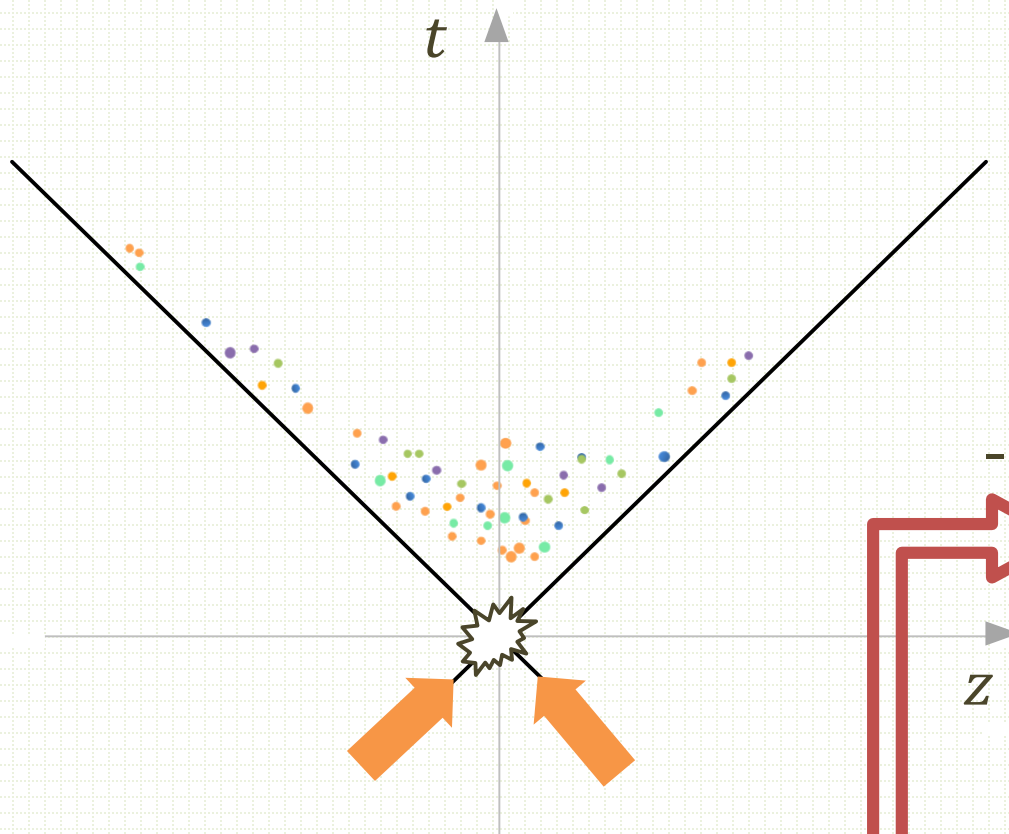
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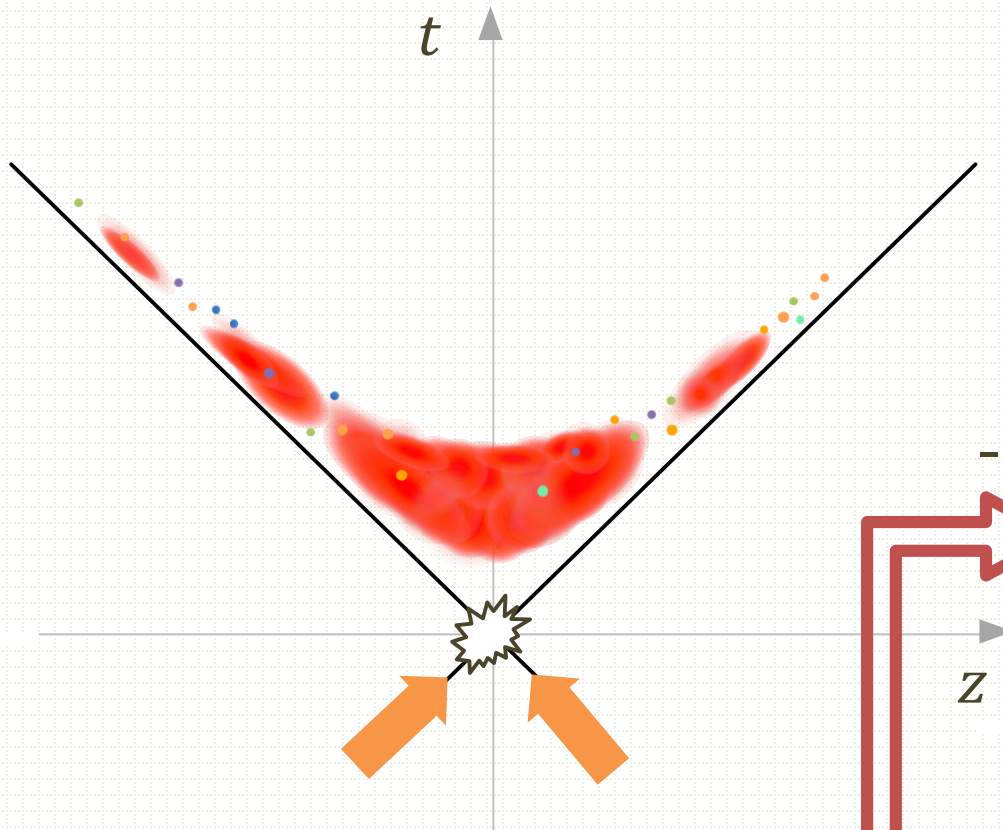
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+ Core-corona picture

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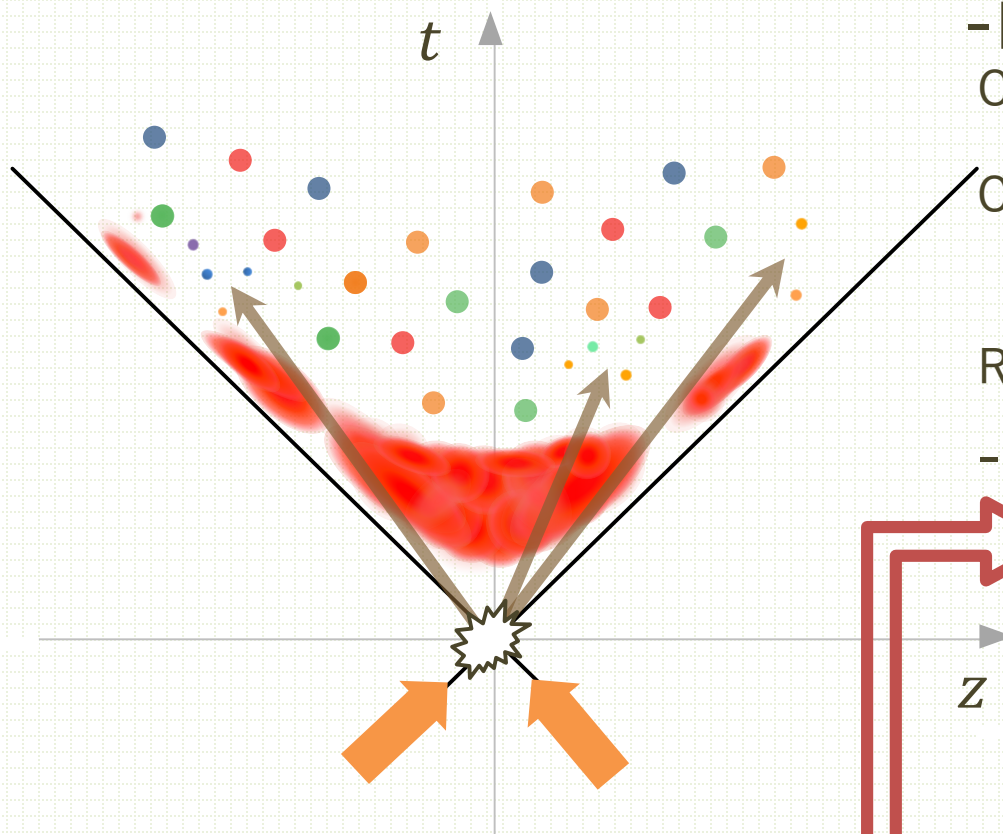
PYTHIA ver 8.230

T. Sjöstrand *et al.*,
Comput. Phys. Commun. 191, 159 (2015).

+ Core-corona picture

Model

Time evolution of high energy nuclear collisions



-Hadronization

Corona \rightarrow String fragmentation
(PYTHIA)

Core \rightarrow Particlization at $T_{ch} = 160$ MeV
with Cooper-Frye formula

F. Cooper and G. Frye, Phys.
Rev. D10, 186 (1974).

Resonance factor: A. Andronic *et al.*, (2017).

-Evolution of QGP fluids

Dynamical initialization

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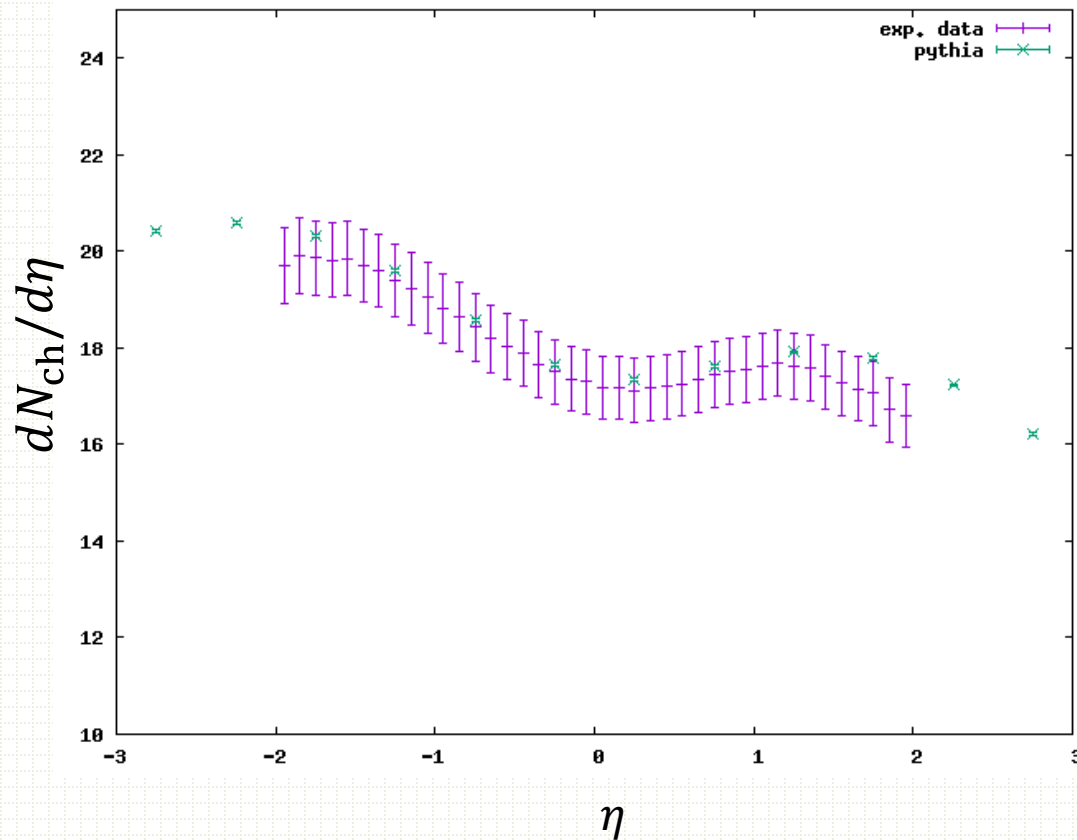
PYTHIA ver 8.230

T. Sjöstrand *et al.*,
Comput. Phys. Commun. 191, 159 (2015).

+ Core-corona picture

Pseudo-rapidity distribution from PYTHIA

p+Pb $\sqrt{s_{NN}} = 2.76$ TeV, NSD



Simulation for heavy ion collision available.

C. Bierlich, G. Gustafson, L. Lönnblad, JHEP PoS DIS2016, 051 (2016)

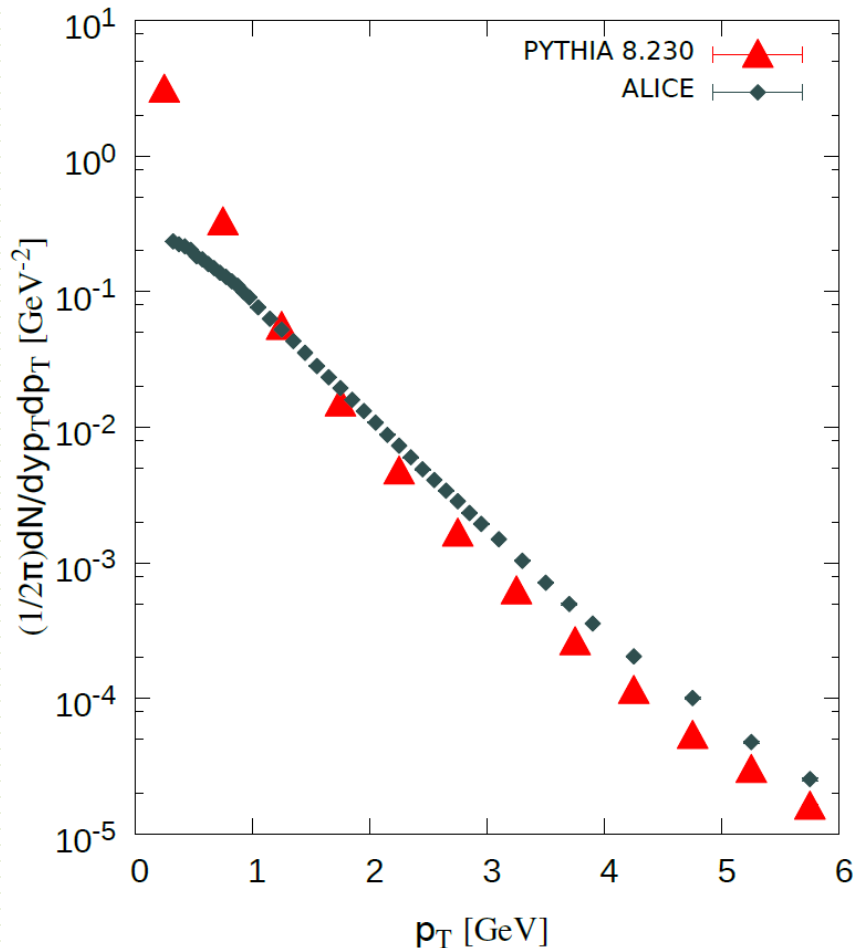


Reproduction of $dN_{ch}/d\eta$

ALICE Collaboration.
Phys. Rev. Lett. 110 (2013) 032301

Transverse momentum distribution from PYTHIA

$p+Pb$ $\sqrt{s_{NN}} = 2.76$ TeV, NSD, mid rapidity



p_T spectra of $p + \bar{p}$

Exp. data \rightarrow hard
PYTHIA \rightarrow soft



No blue-shifts contrary
to experimental data



Need of some
transverse dynamics

Dynamical initialization

M.Okai, K.Kawaguchi,
Y.Tachibana, T.Hirano,
Phys. Rev. C 95, no. 5,
054914 (2017).

Energy-momentum
tensor of QGP fluids

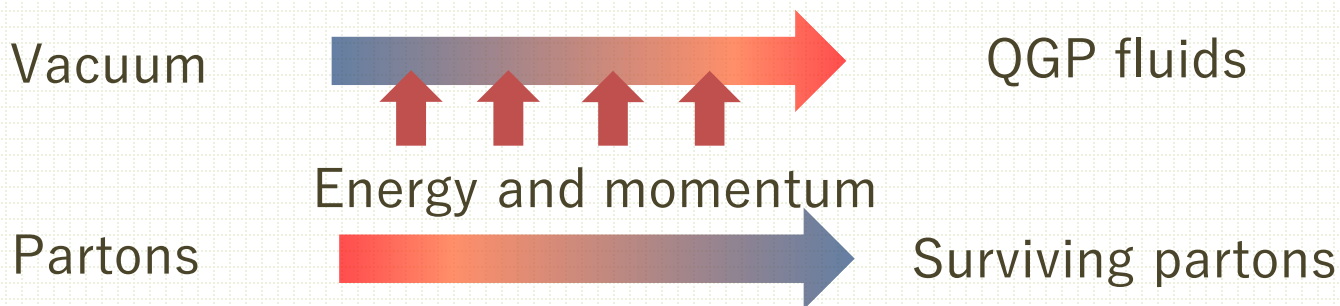
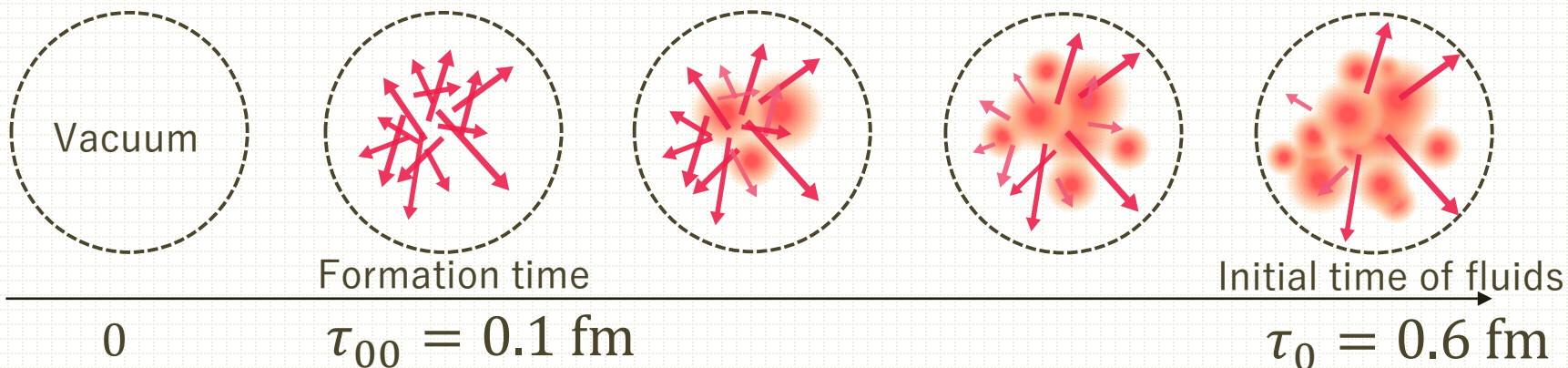
$$\partial_\mu T^{\mu\nu} = J^\nu$$

Energy and momentum
deposited from initial partons

$$J^\mu = - \sum_i \frac{dp_i^\mu}{dt} G(\mathbf{x} - \mathbf{x}_i(t))$$

G : Gaussian function
for smearing

“Fluidization rate”



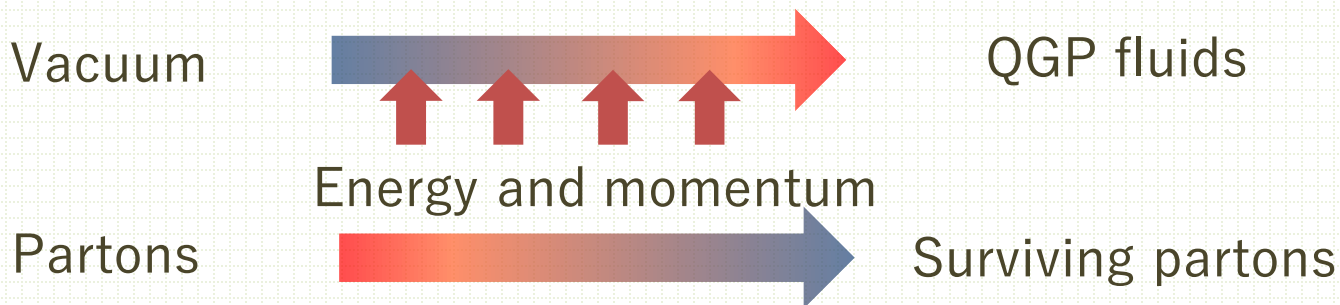
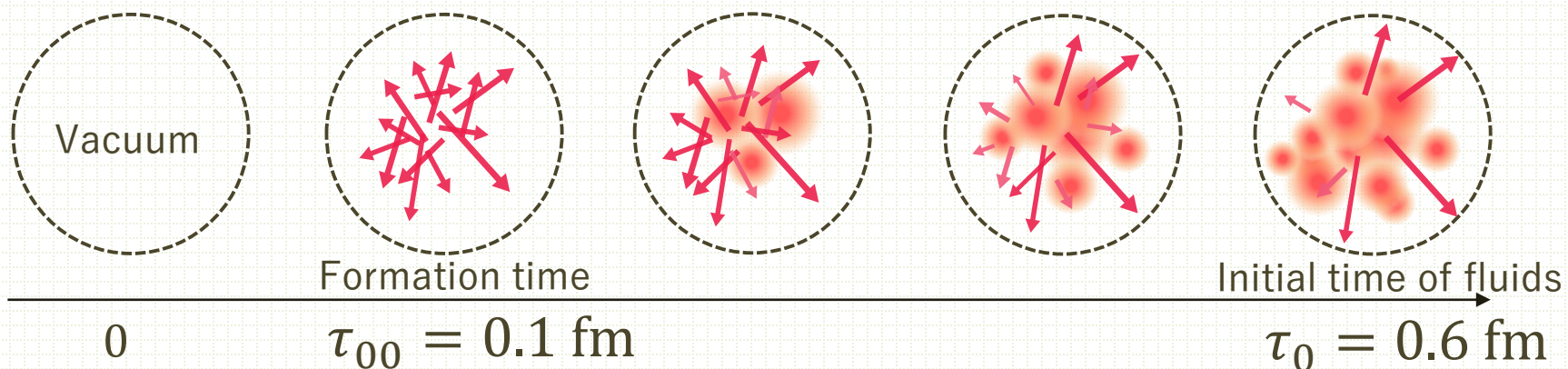
Dynamical initialization

Energy-momentum tensor of QGP fluids $\partial_\mu T^{\mu\nu} = J^\nu$ Energy and momentum deposited from initial partons

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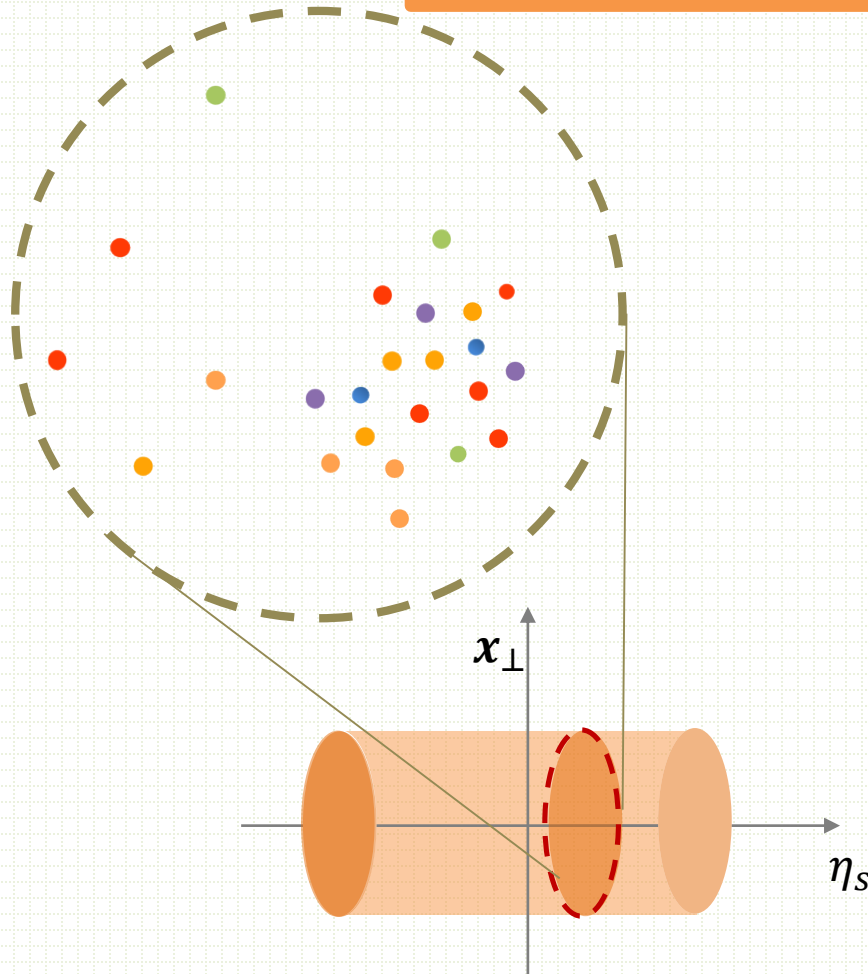
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“Fluidization rate”



Fluidization rate with core-corona picture

$$\frac{dp_i^\mu(t)}{dt} = -a_0 \frac{\rho_i(x_i(t))}{p_{T,i}^2} p_i^\mu(t)$$



– $\rho_i(x_i(t))$: Density distribution

$$\rho_i(x_i(t))d^3x = \sum_{j \neq i} G(x_i - x_j(t))d^3x$$

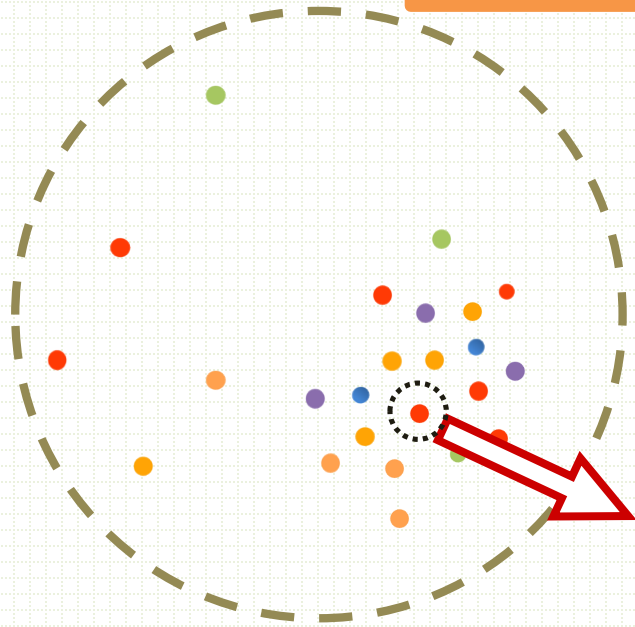
– a_0 : Free parameter to control core-corona picture

“Mean free path”

$$\frac{1}{\lambda_i} = \rho_i \sigma_i \approx \frac{\rho_i}{p_{T,i}^2}$$

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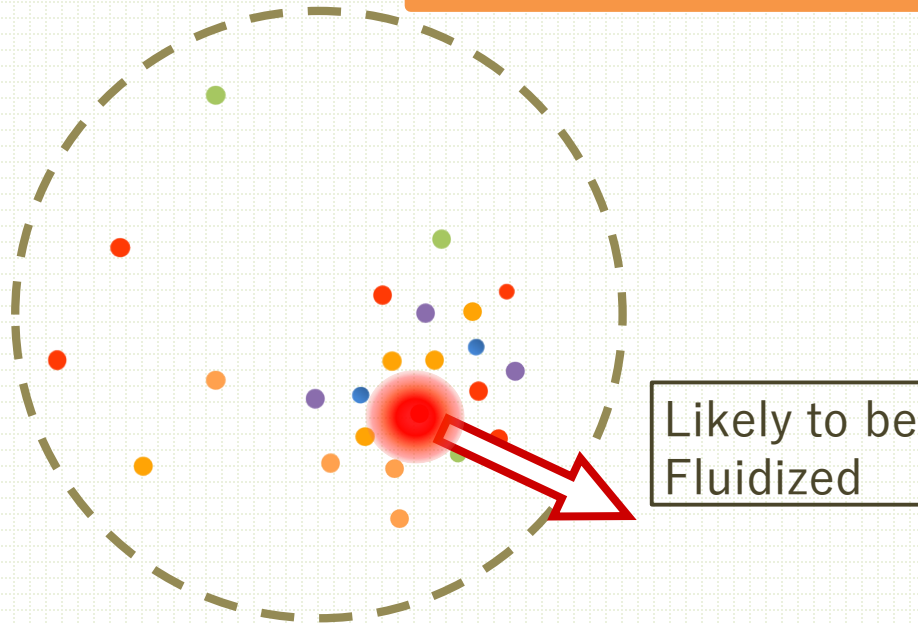
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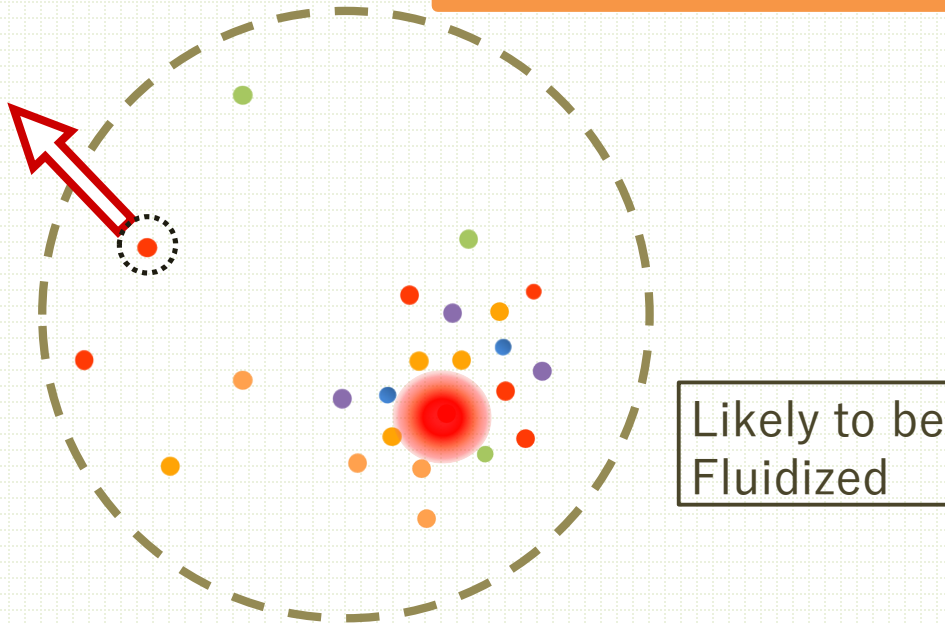
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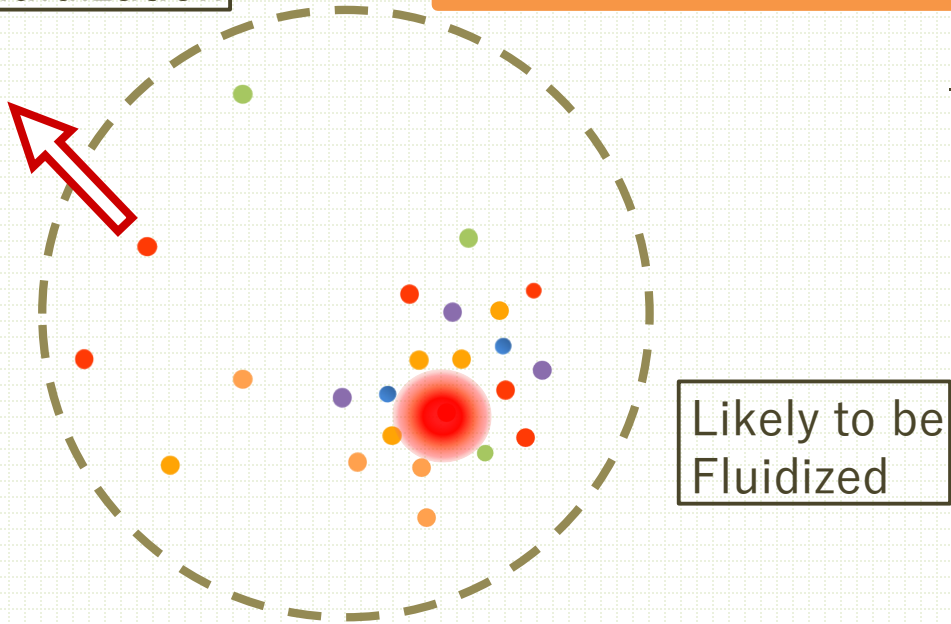
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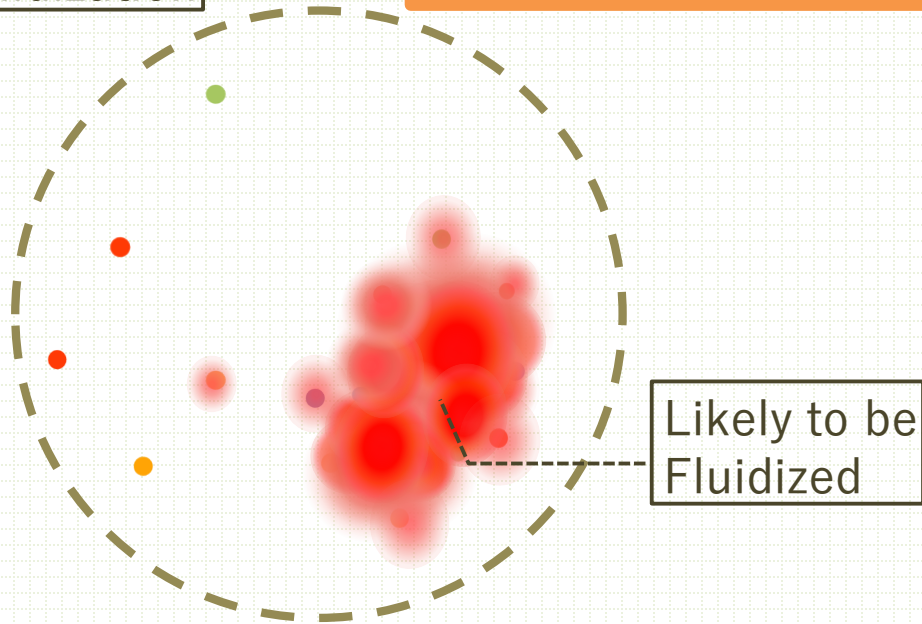
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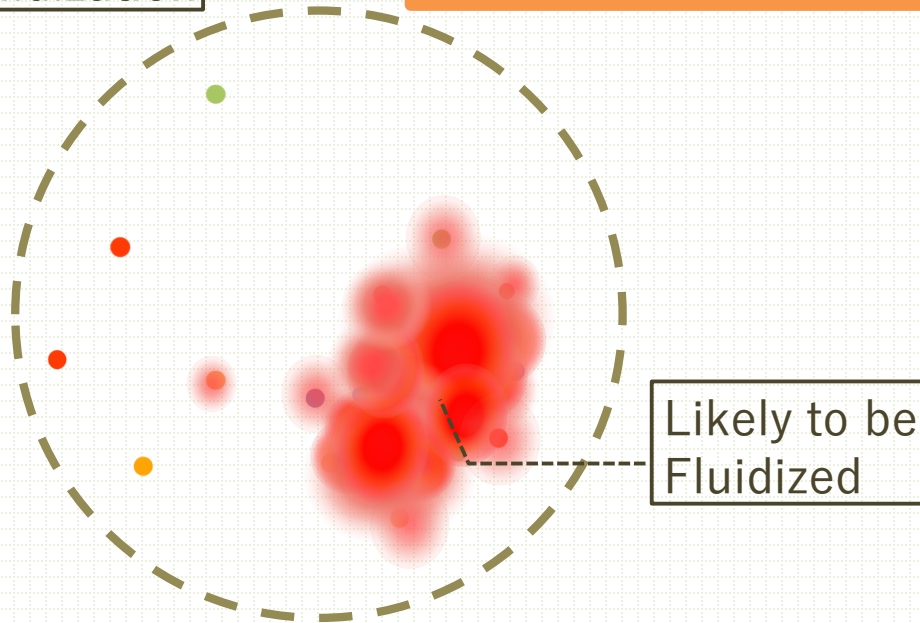


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Likely to be Fluidized



“Mean free path”

$$\frac{1}{\lambda_i} = \rho_i \sigma_i \approx \frac{\rho_i}{p_{T,i}^2}$$

● Fluidization depends on the initial parton distribution.

● dense region

dilute region



QGP fluids



Surviving partons

Results

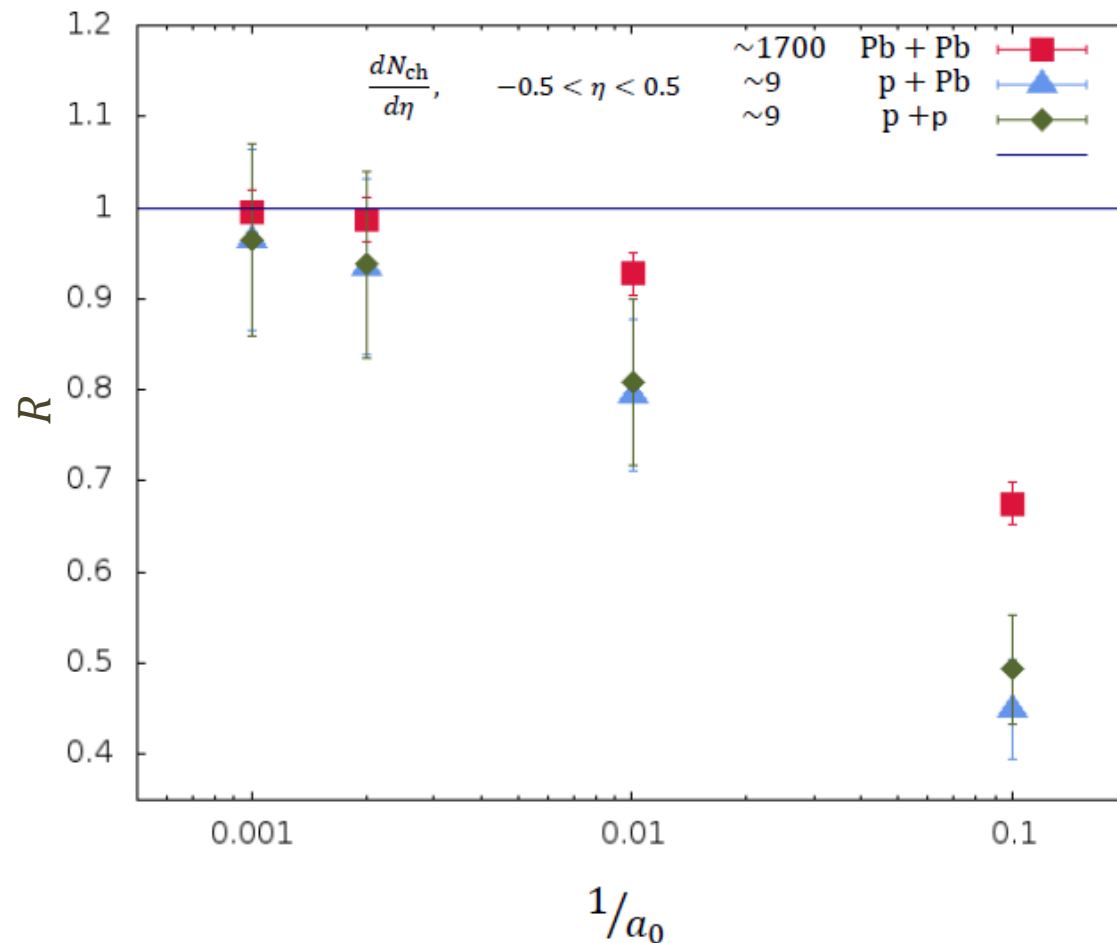
Estimation of a_0

a_0 dependence of fluidized energy rate R in each system (typical multiplicity events)

$$R = E_{\text{generated fluid}}/E_{\text{input}}$$

a_0 : Free parameter to control core-corona

$$\frac{dp_i^\mu(t)}{dt} = -a_0 \frac{\rho_i(x_i(t))}{p_{T,i}^2} p_i^\mu(t)$$



$$a_0 = 10^2$$



$R \approx 1$ Pb+Pb
 \rightarrow Almost fluidized

$R < 1$ p+p, p+Pb
 \rightarrow Supressed fluidization

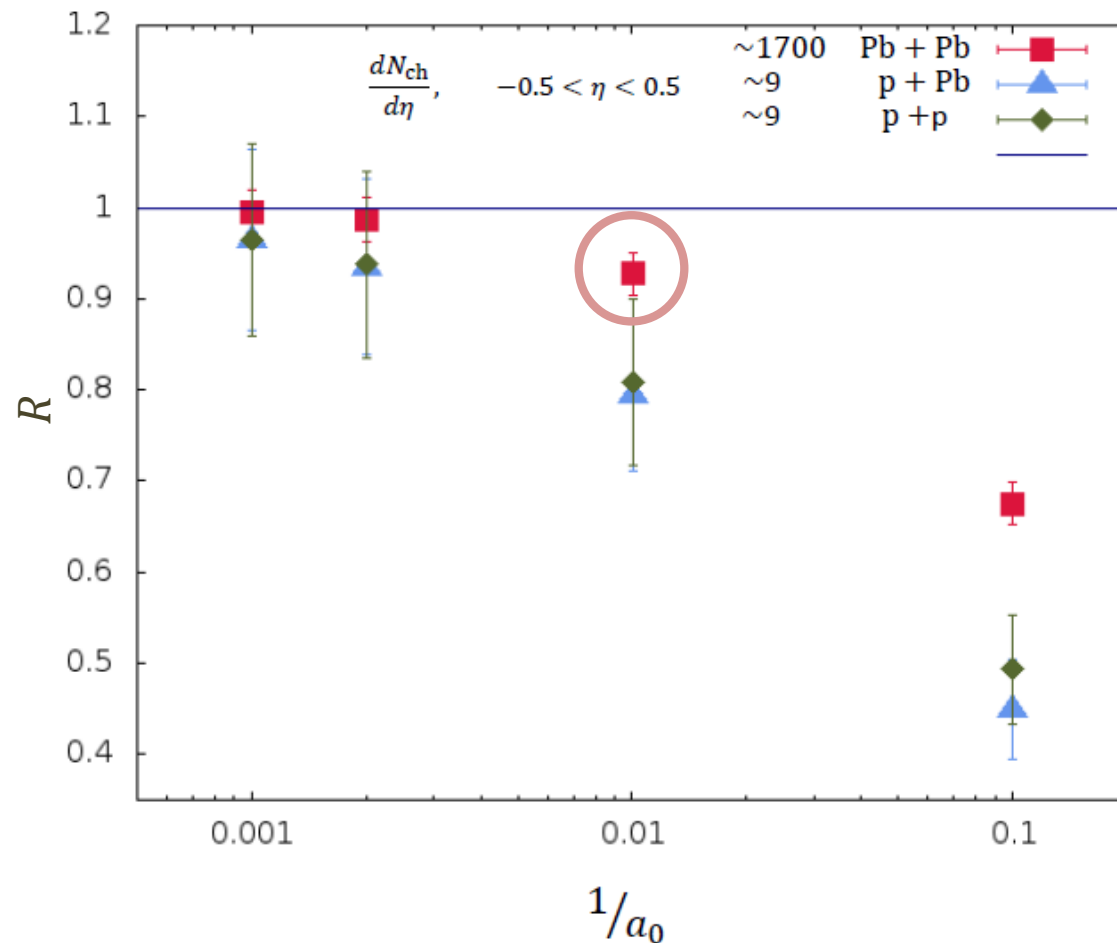
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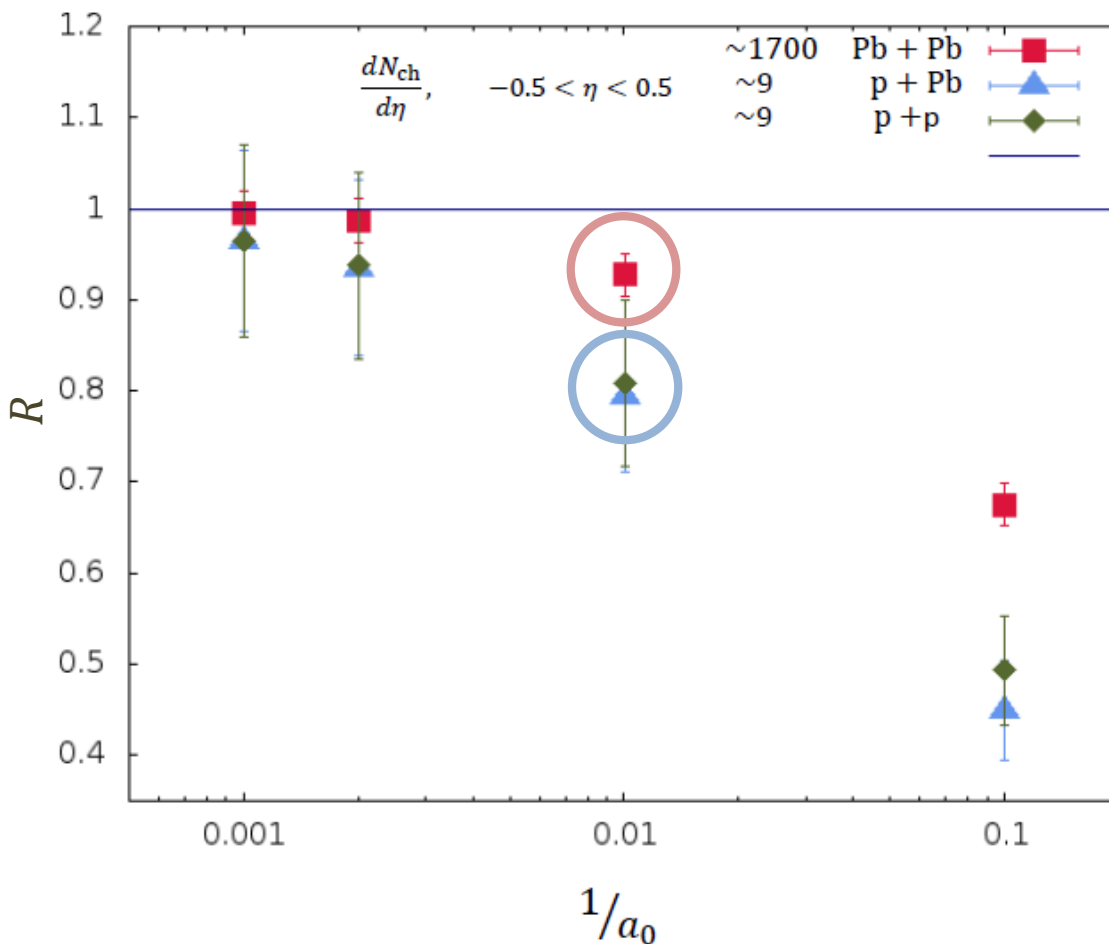
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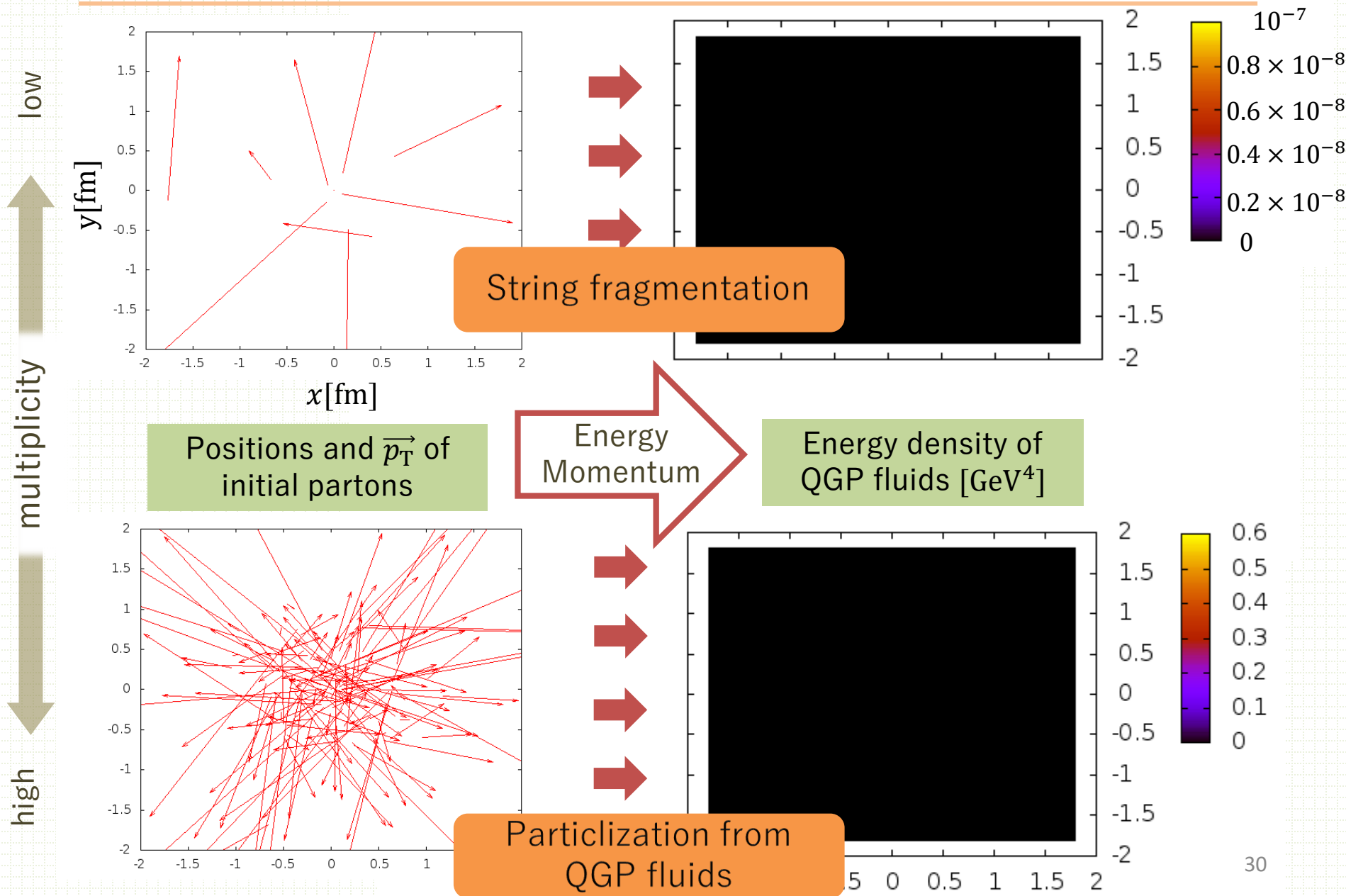
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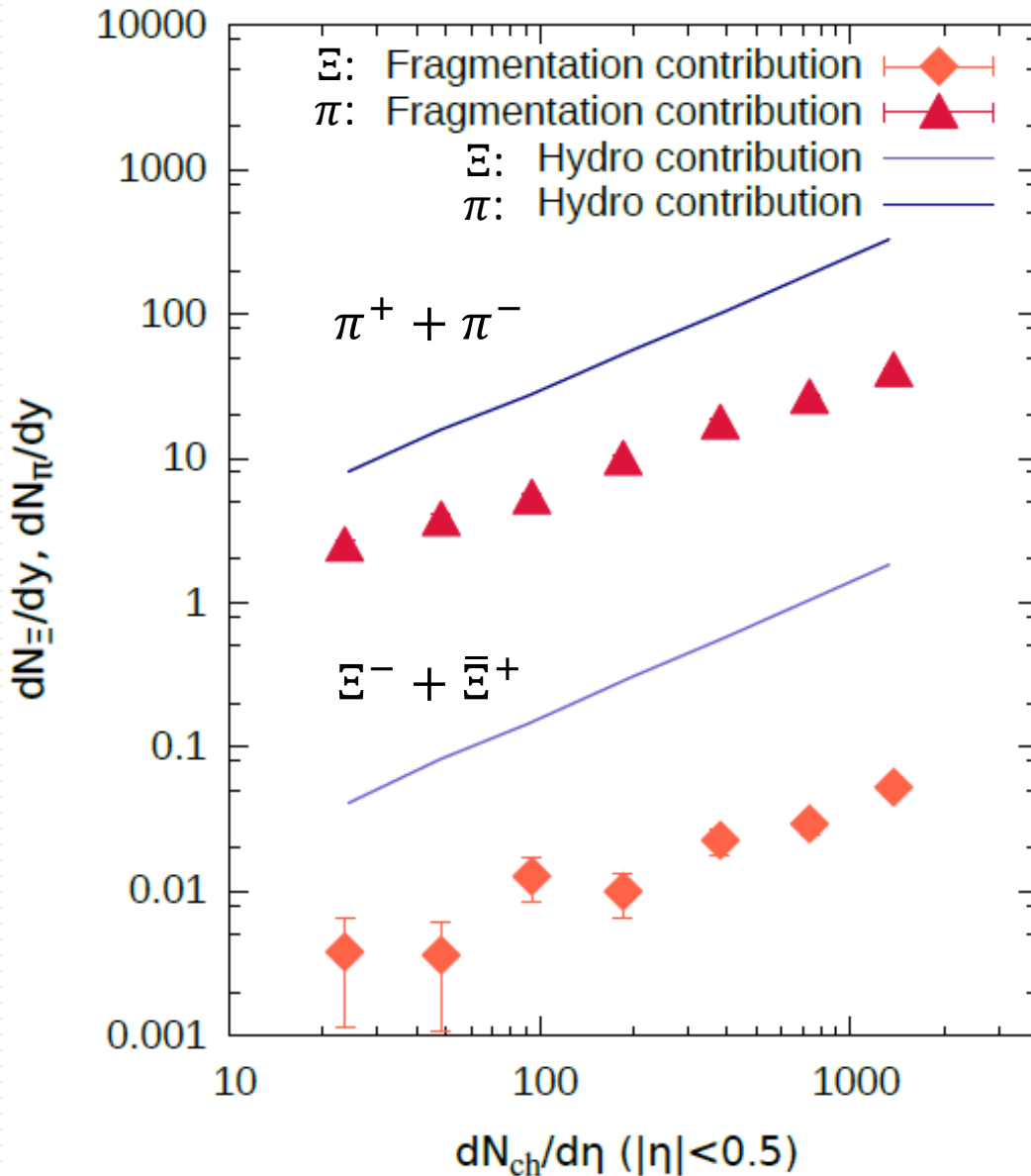
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→ Supressed fluidization

Dynamical initialization with core-corona in p+p



Yield vs multiplicity in Pb+Pb

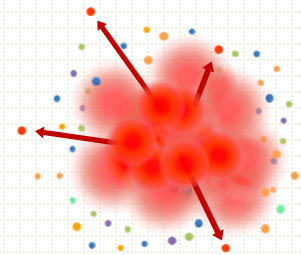


Pb+Pb, $\sqrt{s_{NN}} = 2.76$ TeV

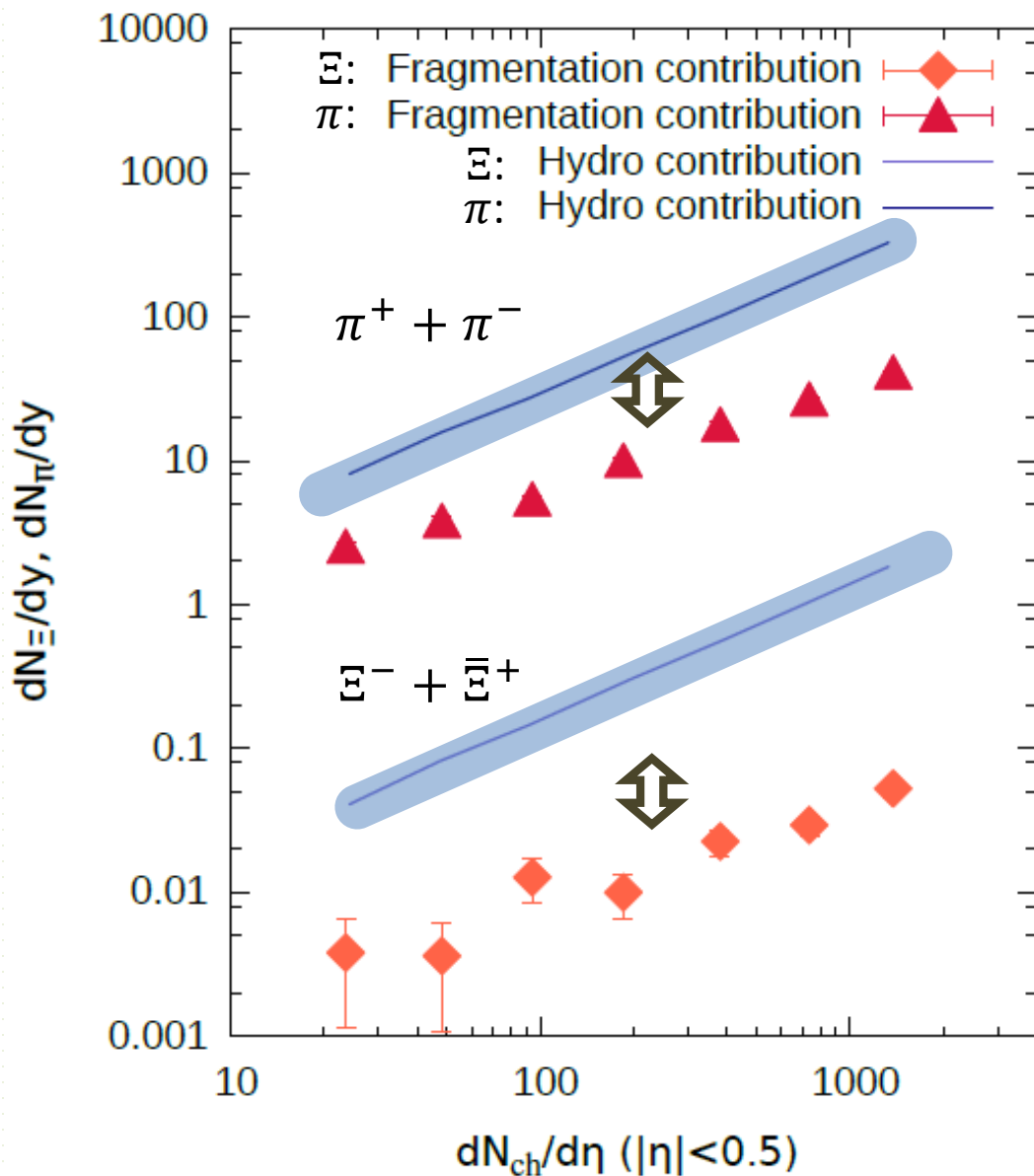


Initial parton distribution:

Dense



Yield vs multiplicity in Pb+Pb

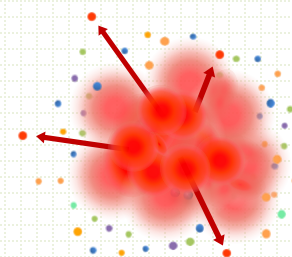


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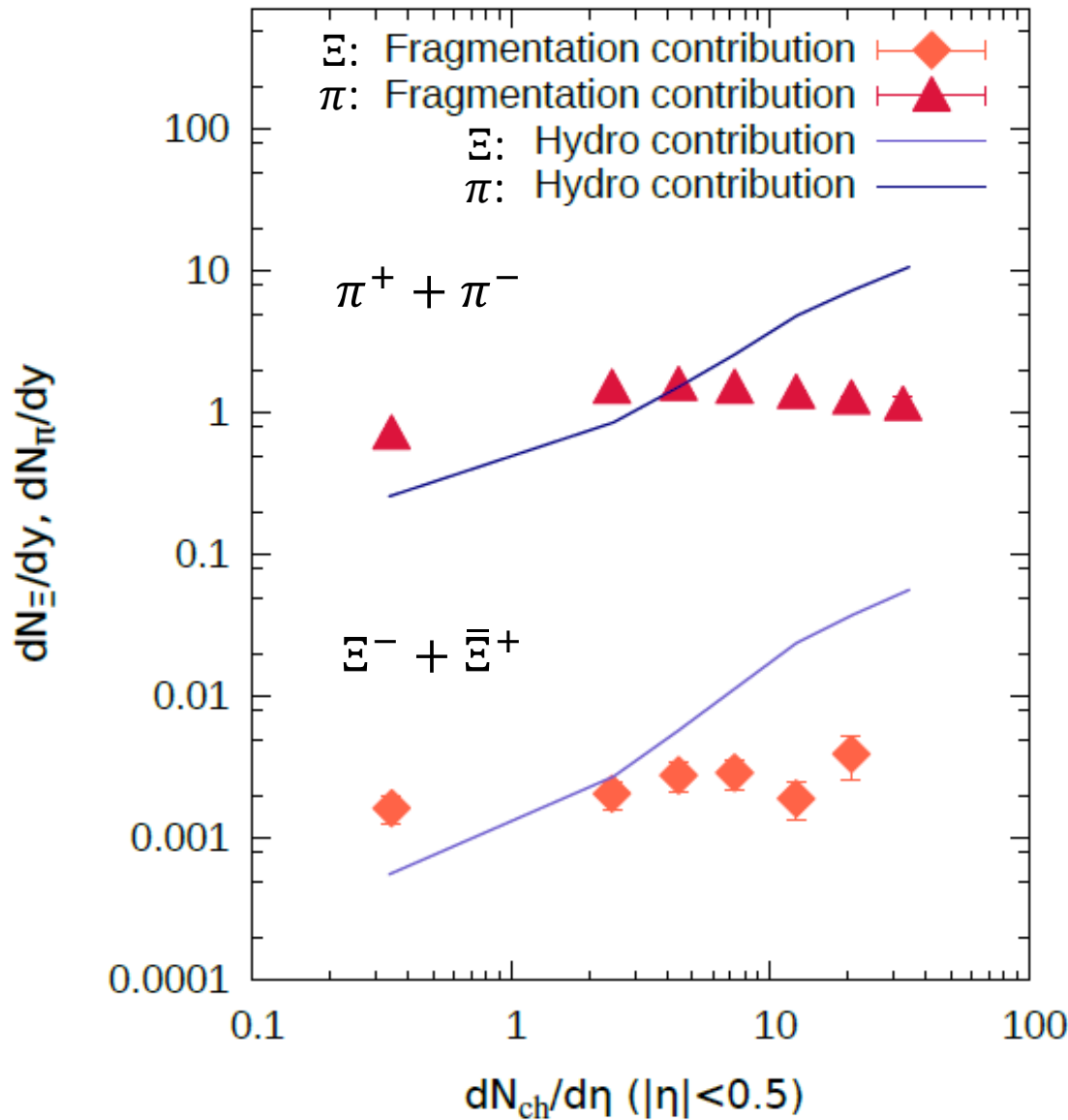
Initial parton distribution:

Dense



Fluids contribution
is dominant
in most events.

Yield vs multiplicity in p+p

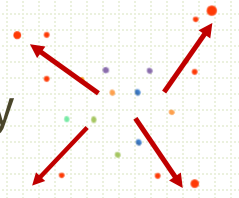


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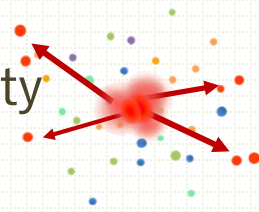


Initial parton distribution:

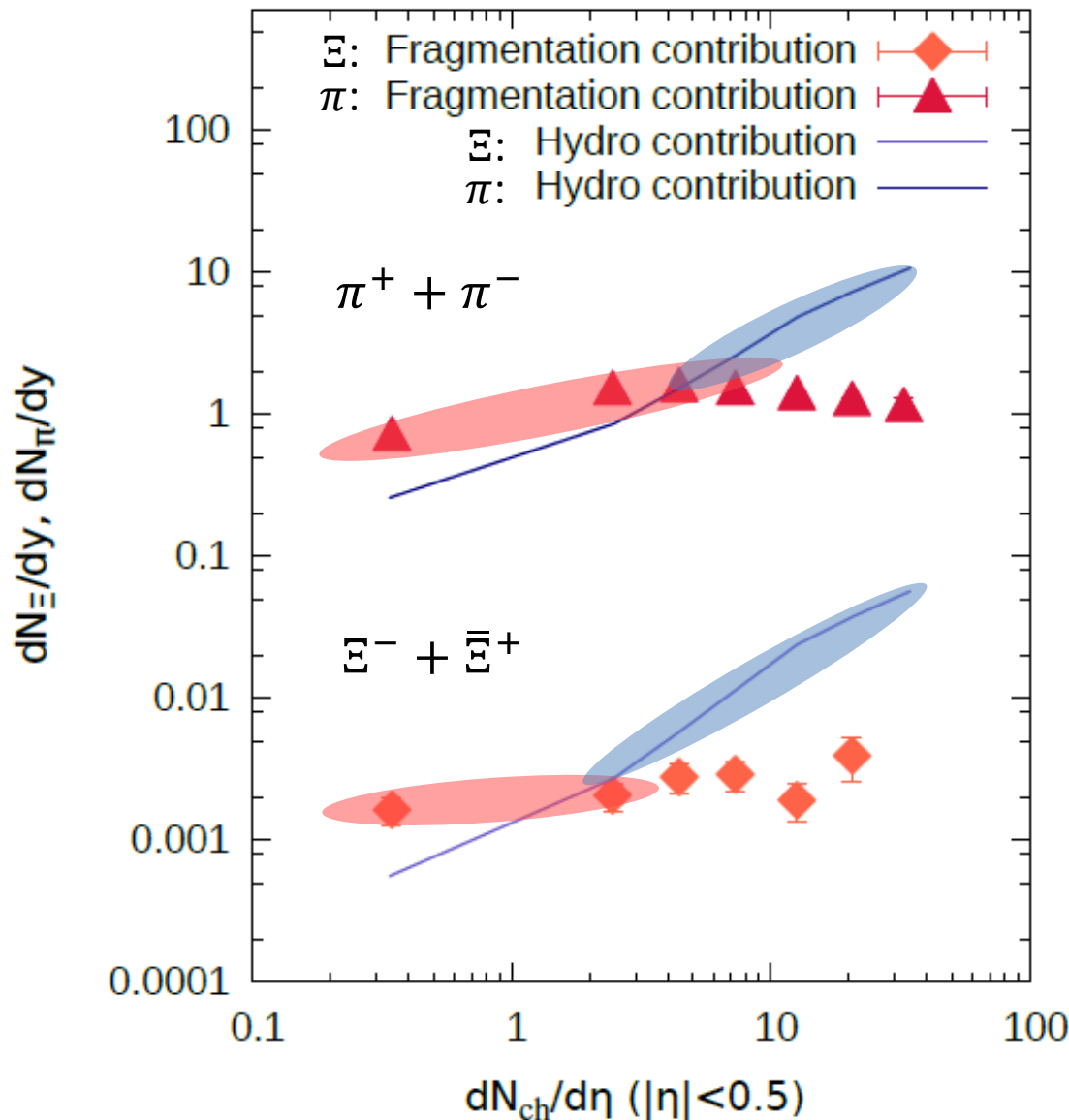
-Dilute
in low multiplicity
events.



-Dense
in high multiplicity
events.



Yield vs multiplicity in p+p

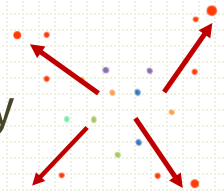


p+p $\sqrt{s_{NN}} = 2.76$ TeV

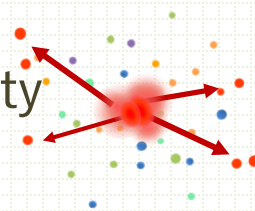


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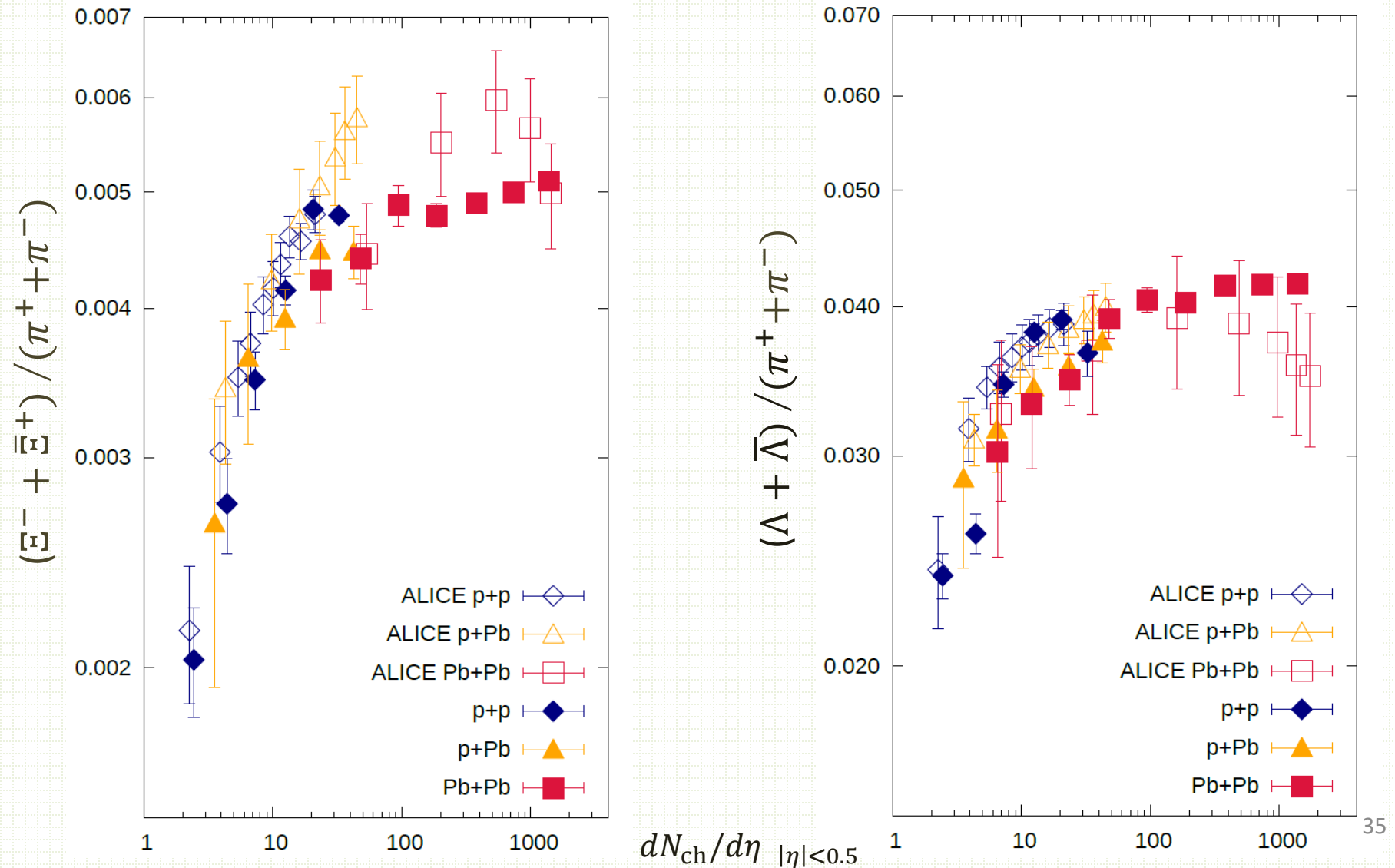
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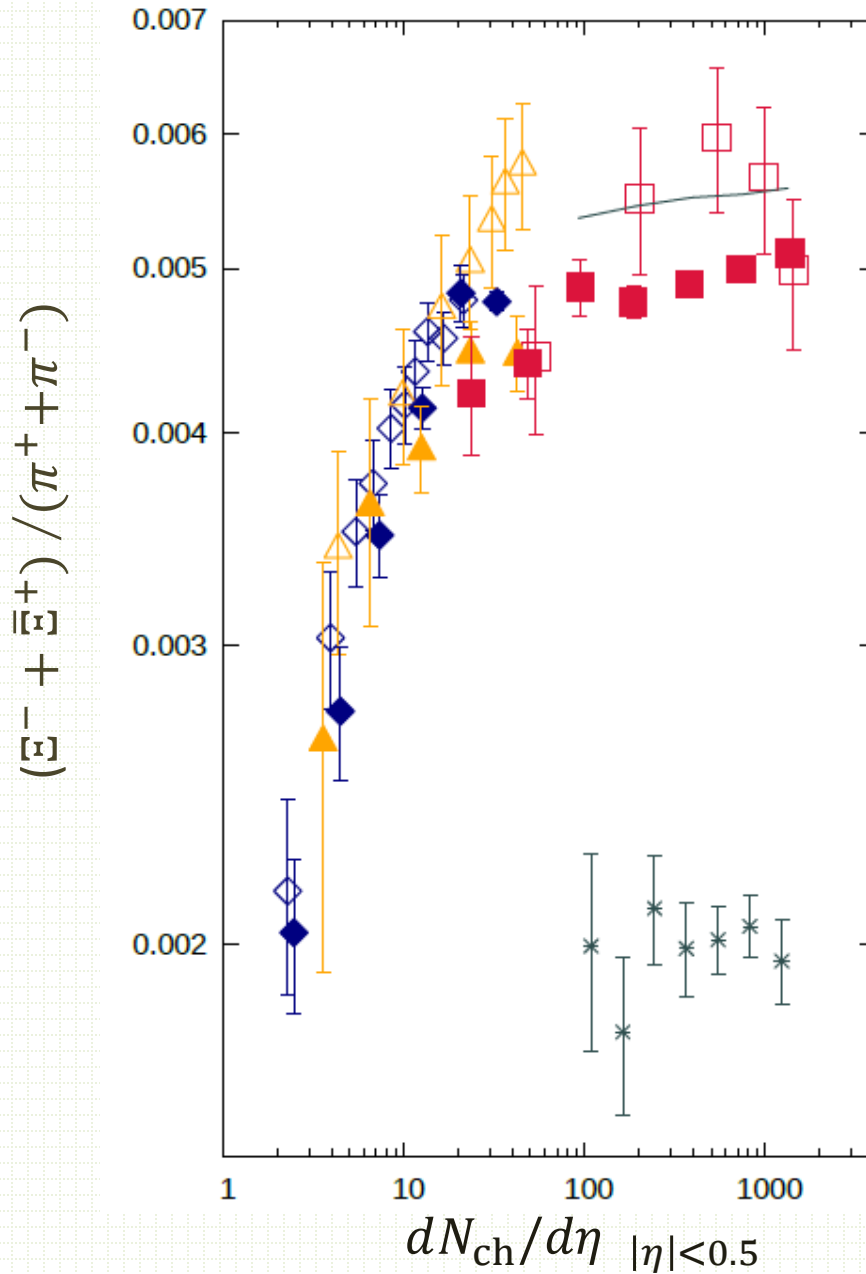
Competition between
fluids and fragmentation
→ depend on multiplicity

Ratio of yields to pions vs multiplicity

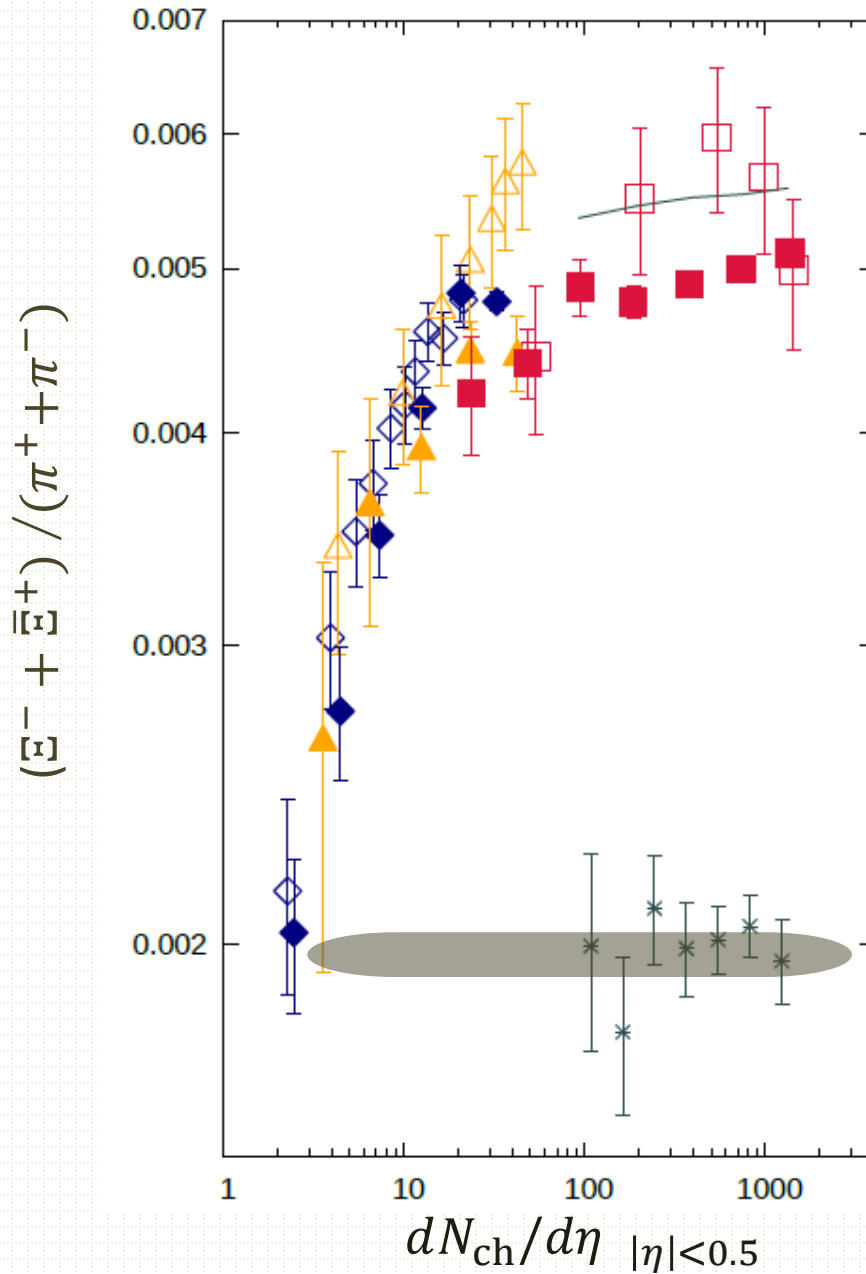
Results from dynamical initialization with core-corona picture



Continuous change of ratio

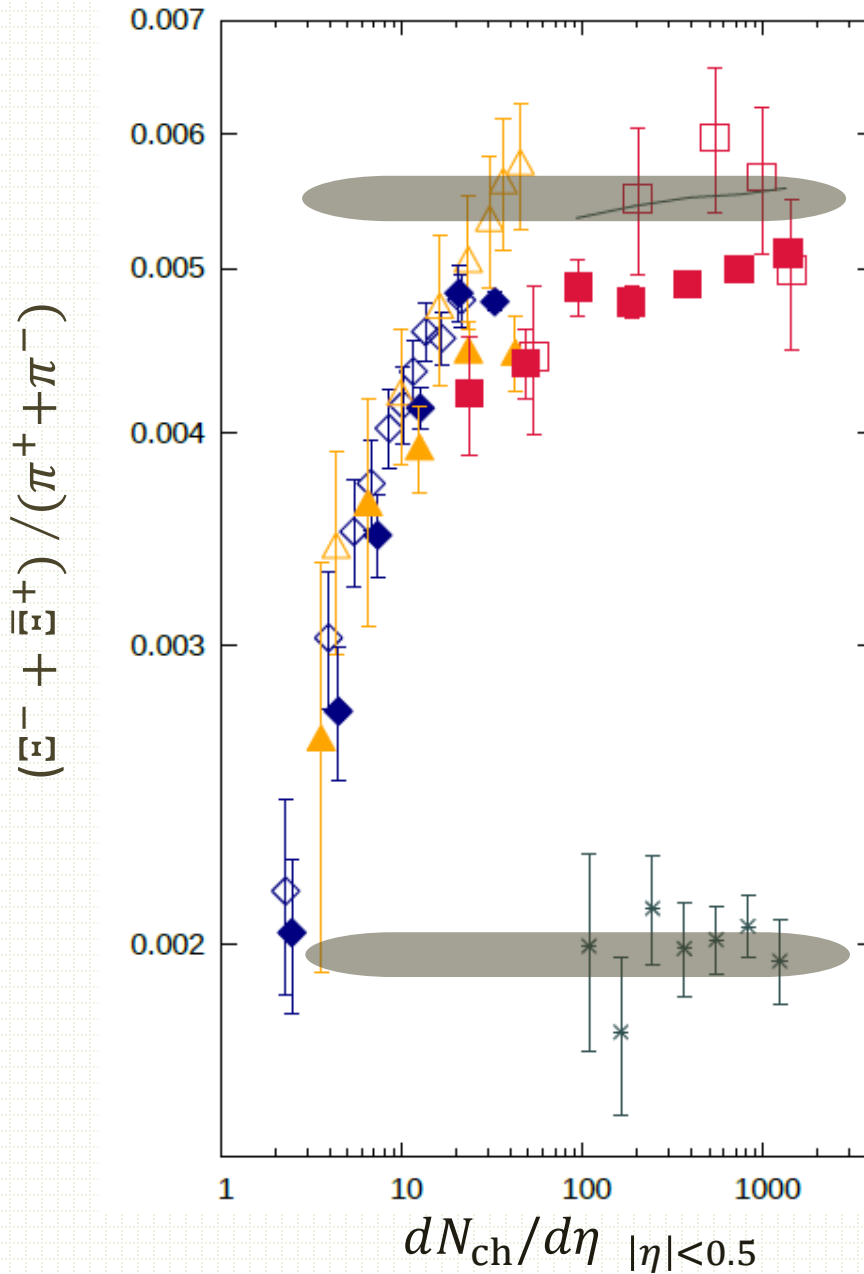


Continuous change of ratio



String fragmentation limit:
hadron products only from
string fragmentation

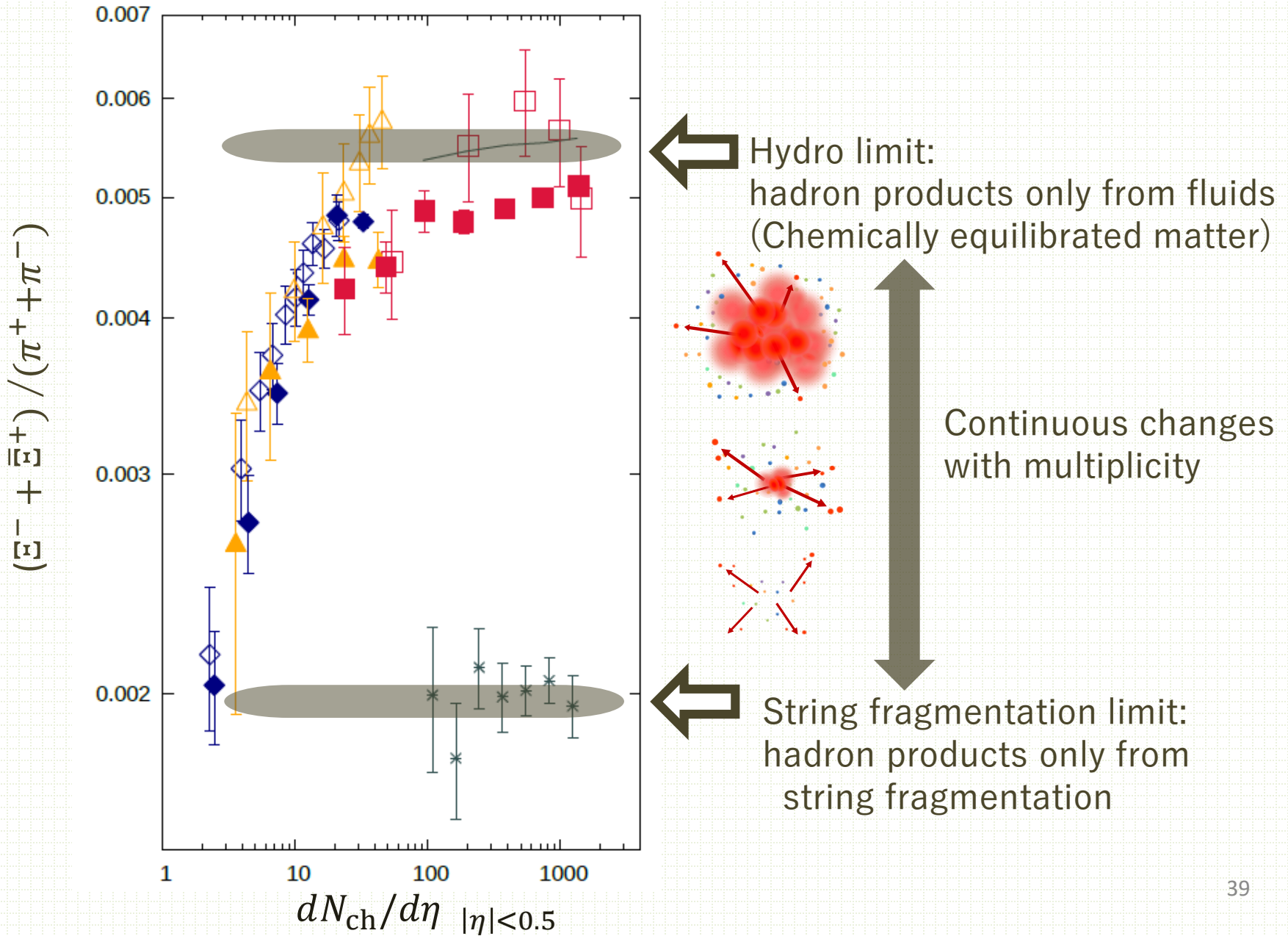
Continuous change of ratio



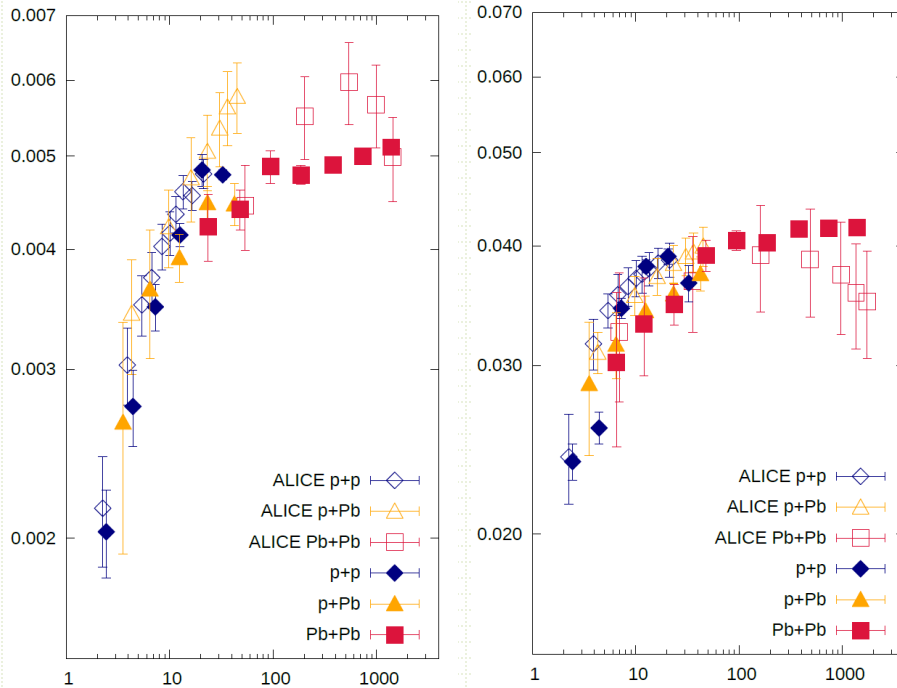
← Hydro limit:
hadron products only from fluids
(Chemically equilibrated matter)

← String fragmentation limit:
hadron products only from
string fragmentation

Continuous change of ratio



Interpretation of results



$$\begin{aligned} p+p & \sqrt{s_{NN}} = 7 \text{ TeV} \\ p+Pb & \sqrt{s_{NN}} = 5.02 \text{ TeV} \\ Pb+Pb & \sqrt{s_{NN}} = 2.76 \text{ TeV} \end{aligned}$$



$$dN_{ch}/d\eta < \sim 100: \\ \rightarrow \text{Increase}$$

$$dN_{ch}/d\eta \geq \sim 100: \\ \rightarrow \text{Saturation}$$

Almost independent of collision energy and system size.

Strong implication:

QGP fluids are partly created even in small systems if multiplicity is high.

Summary & Outlook

Summary

- Dynamical initialization + core-corona picture
 - Low p_T partons in dense region tend to be fluids.
 - Two different hadronization processes
 - Particlization of chemically equilibrated matter (fluids)
 - String fragmentation
- Analysis of strangeness enhancement in p+p, p+Pb and Pb+Pb
 - Reproduced the tendency of ALICE data.
 - Continuous change of strangeness ratio with multiplicity
 - Almost no energy or system size dependence

→ **Multiplicity** is the key for QGP creation.

→ **Core-corona** picture is important for interpretation of strangeness enhancement in small systems.

Outlook

More possibilities to explain collectivity
in small systems with dynamical
initialization with core-corona picture

- Collective flow in small systems
 - Radial flow
 - Flow harmonics
 - Ridge structure
 - ...
- Core-corona effects in large systems
 - Jet quenching
 - Collective flow in peripheral collisions
 - ...

Future work: Sophistication of our model for more detailed study