

Long-range two particle correlations in small system

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Outline

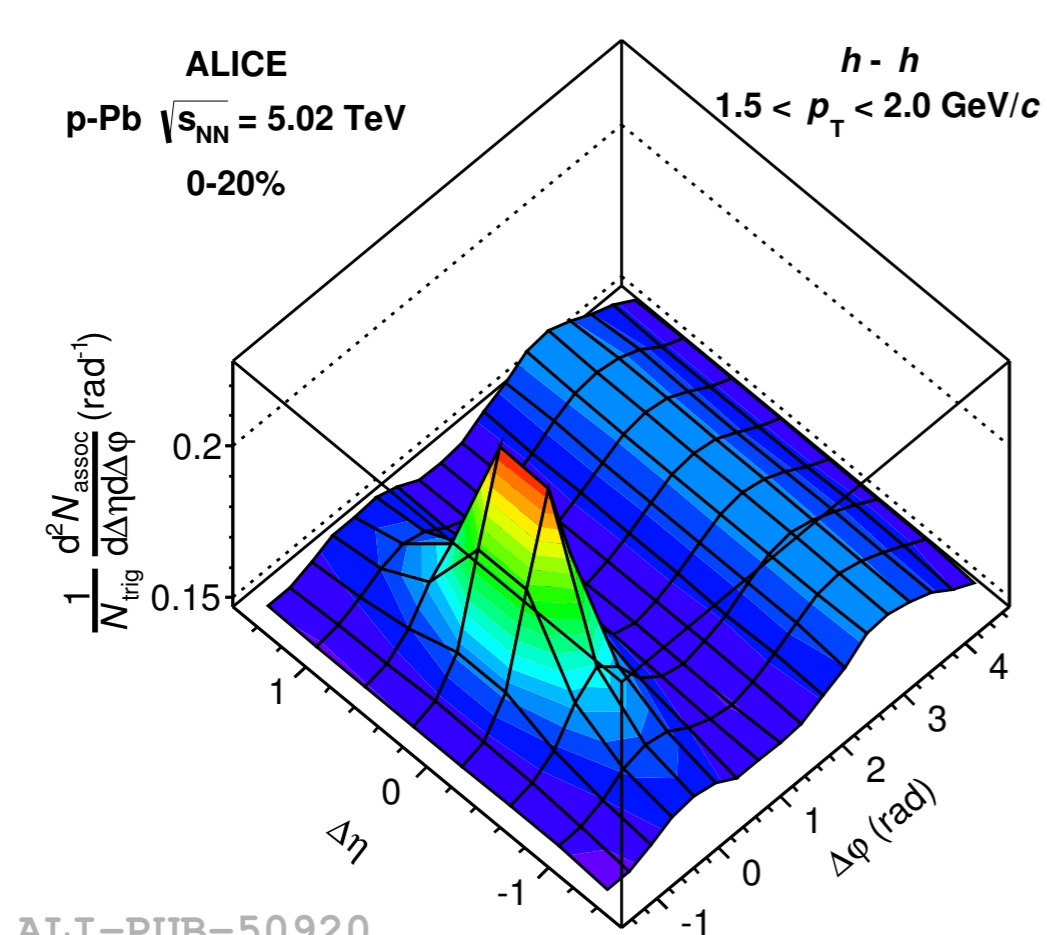
- Two particle correlation analysis in ALICE
- PID v_2
- Forward/backward v_2 in pPb
- Heavy flavor and J/Psi v_2 in pPb
- Model comparison for PID v_2
- Event plane decorrelation in PbPb and pPb.
- Outlook
- Summary

2 particle correlation analysis in ALICE

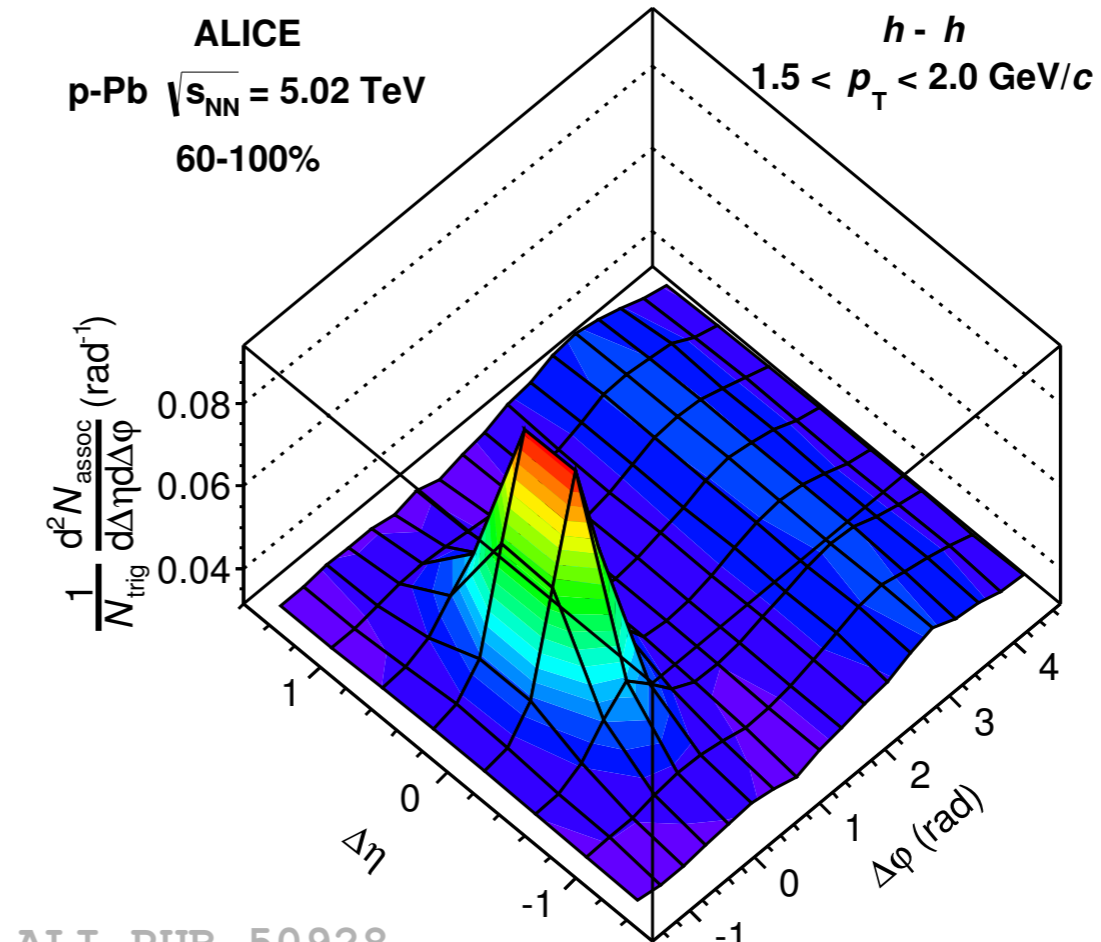
$$\frac{1}{N_{\text{trig}}} \frac{d^2 N_{\text{assoc}}}{d\Delta\eta d\Delta\phi} = \frac{S(\Delta\eta, \Delta\phi)}{B(\Delta\eta, \Delta\phi)} \quad S(\Delta\eta, \Delta\phi) = 1/N_{\text{trig}} d^2 N_{\text{same}}/d\Delta\eta d\Delta\phi$$

$$B(\Delta\eta, \Delta\phi) = \alpha d^2 N_{\text{mixed}}/d\Delta\eta d\Delta\phi$$

Phys. Lett. B 726 (2013) 164-177



ALI-PUB-50920



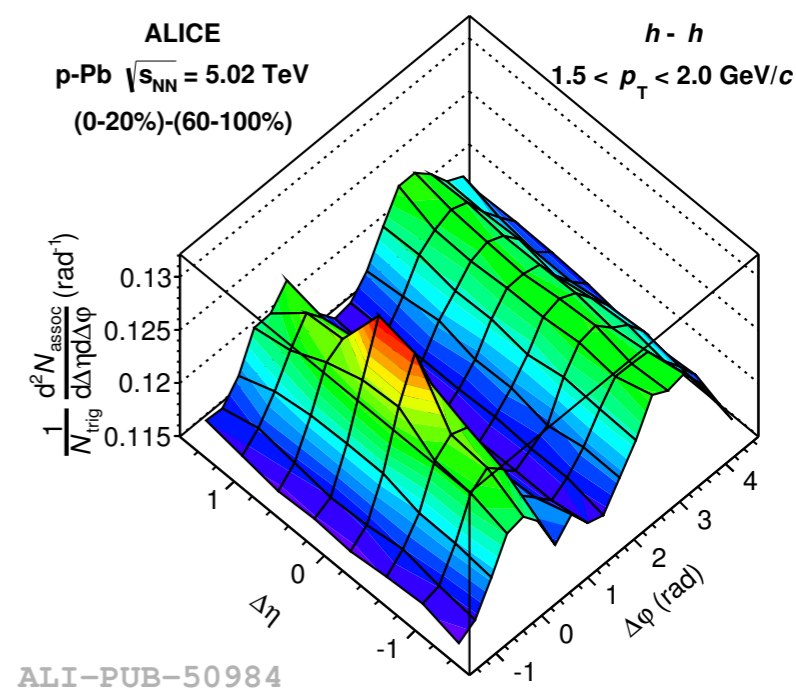
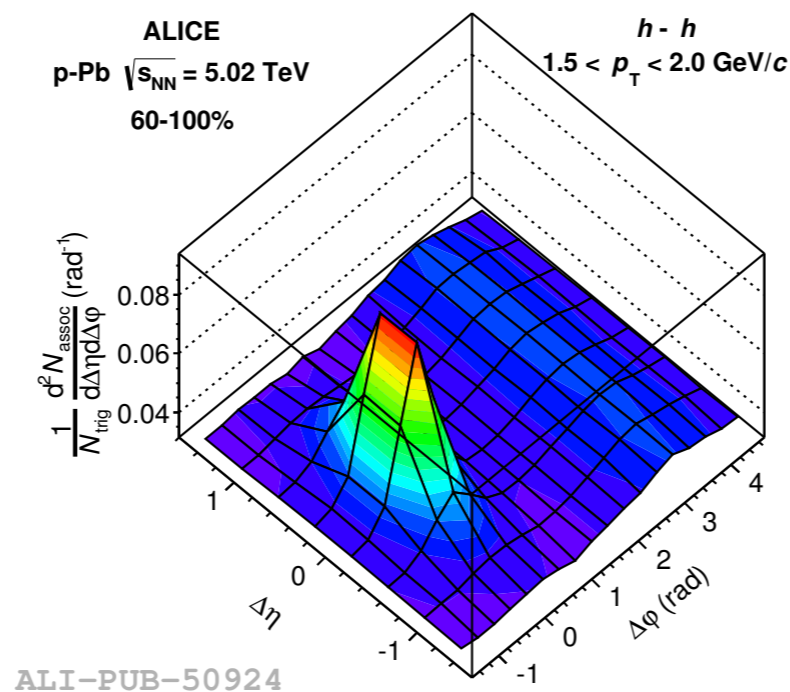
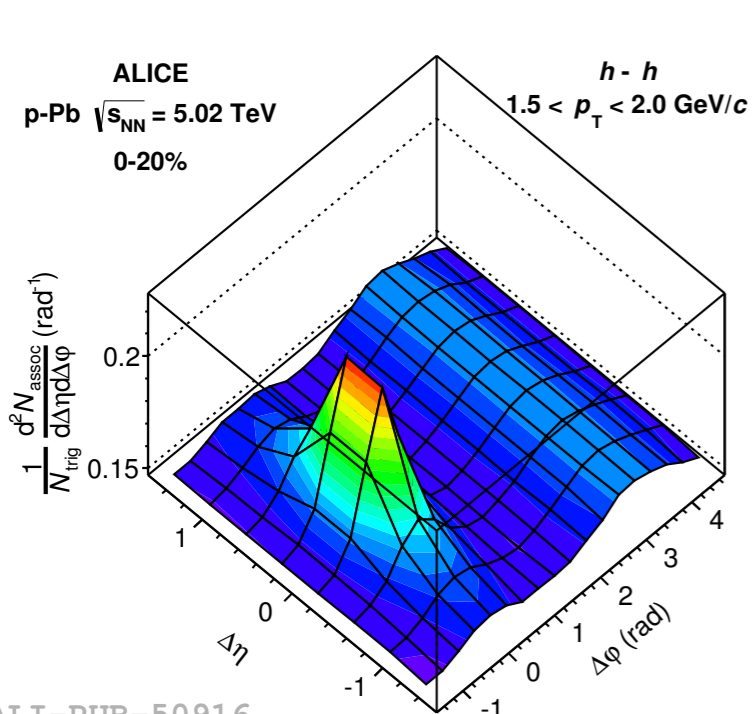
ALI-PUB-50928

- Associate yields per trigger in same events are divided by the mixed events to cancel acceptance effects.
- Jet peak at $(\Delta\eta, \Delta\phi) \sim (0, 0)$ and recoil jet at $\Delta\phi \sim \pi$.
- Near-side ridge is visible in 0-20%.

Non flow subtraction in ALICE

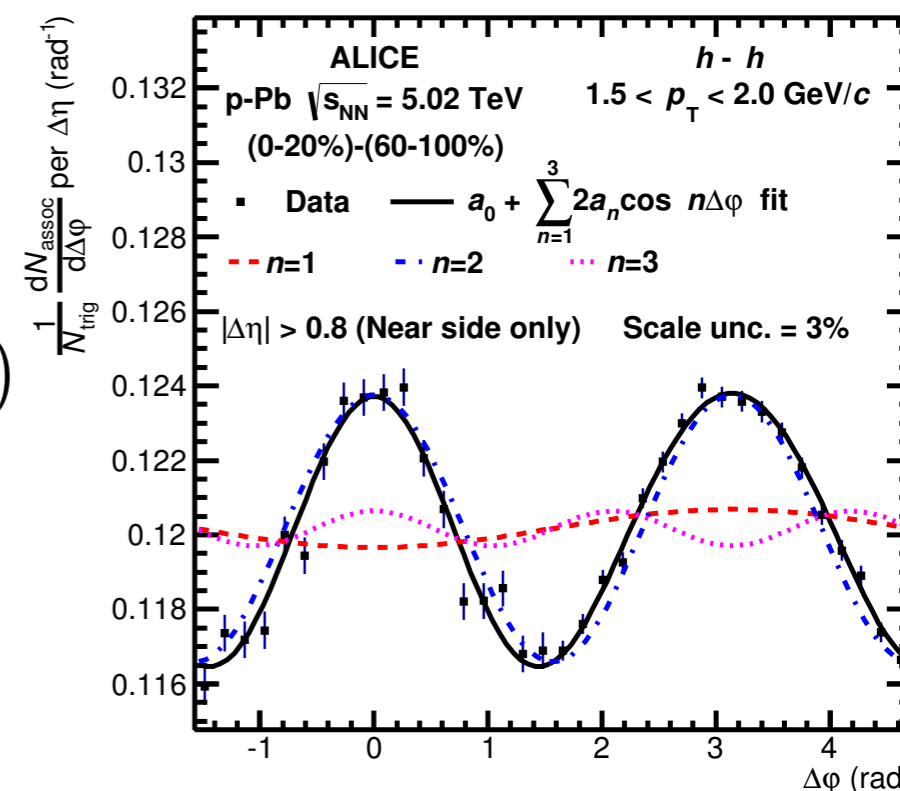
- Subtract 60-100% from 0-20%

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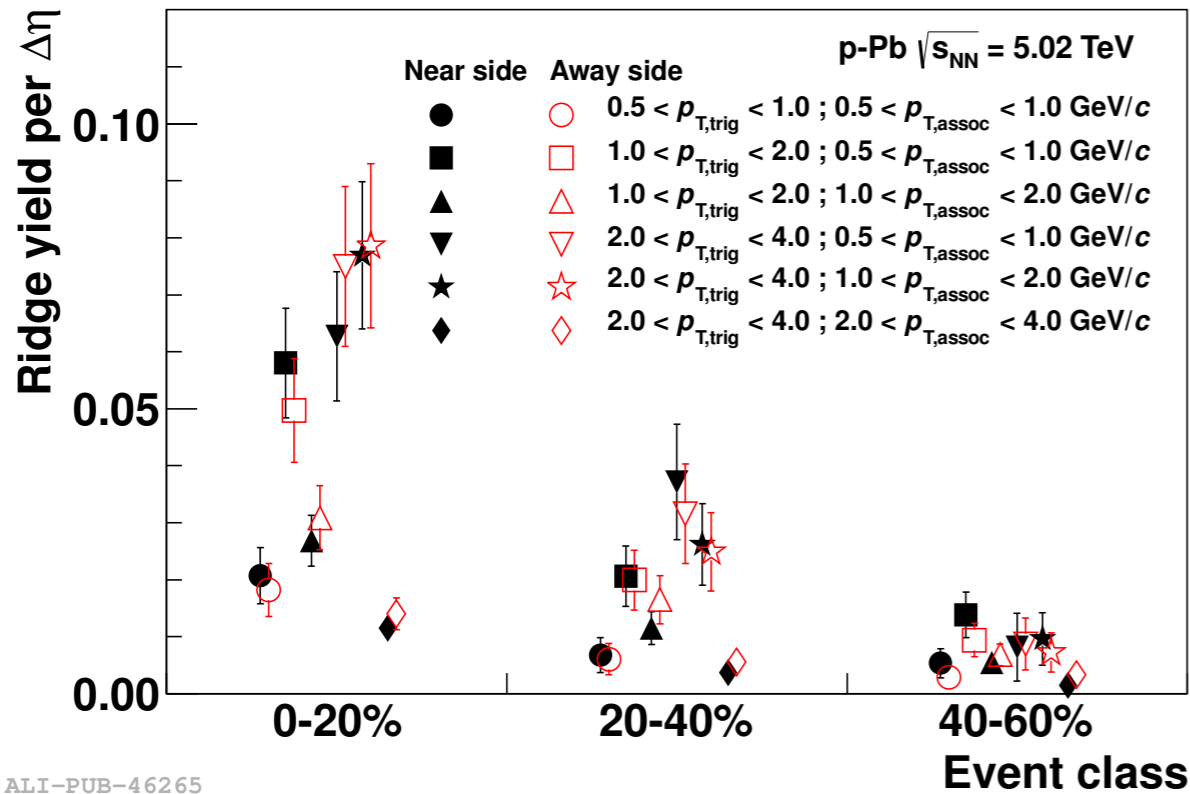
- Exclude $|\Delta\eta| < 0.8$ in near-side

$$\frac{1}{N_{trig}} \frac{dN_{asso}}{d\Delta\phi} = a_0 + \sum_{n=1}^3 2a_n \cos(n\Delta\phi)$$

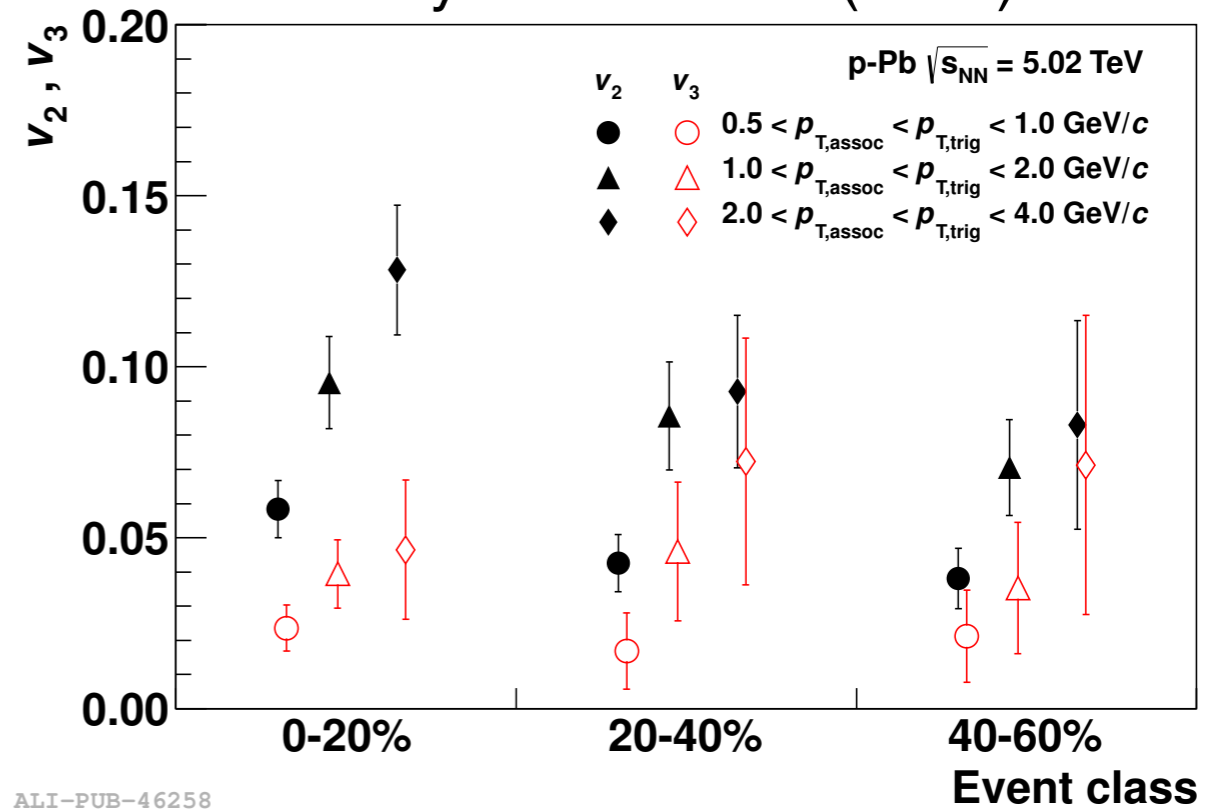


Ridge yield and flow coefficient

Phys. Lett. B 719 (2013) 29-41



ALI-PUB-46265



ALI-PUB-46258

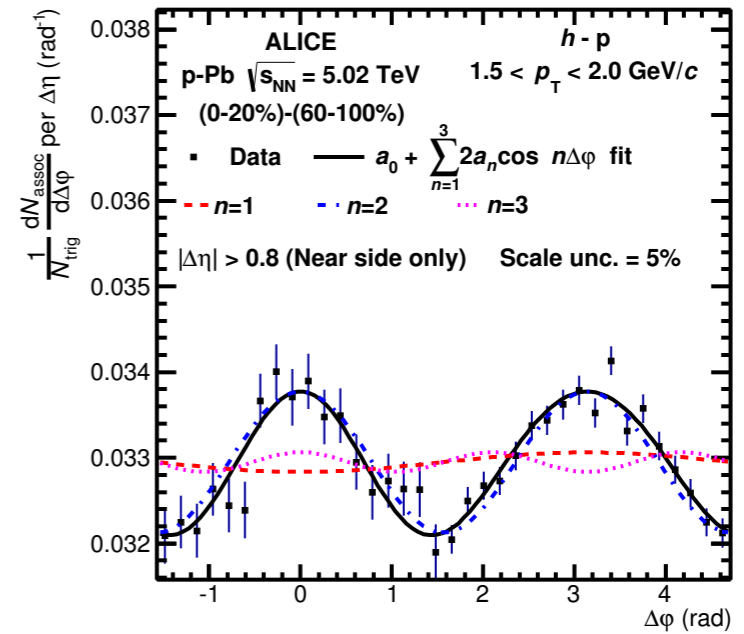
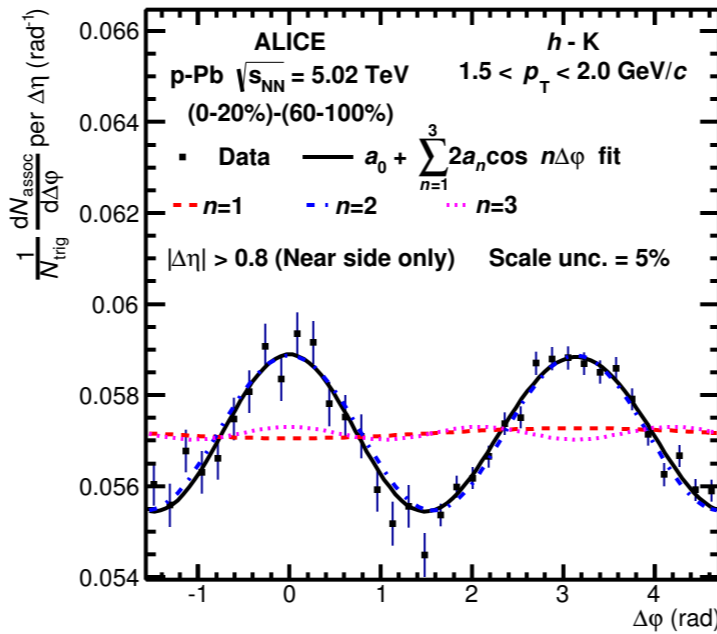
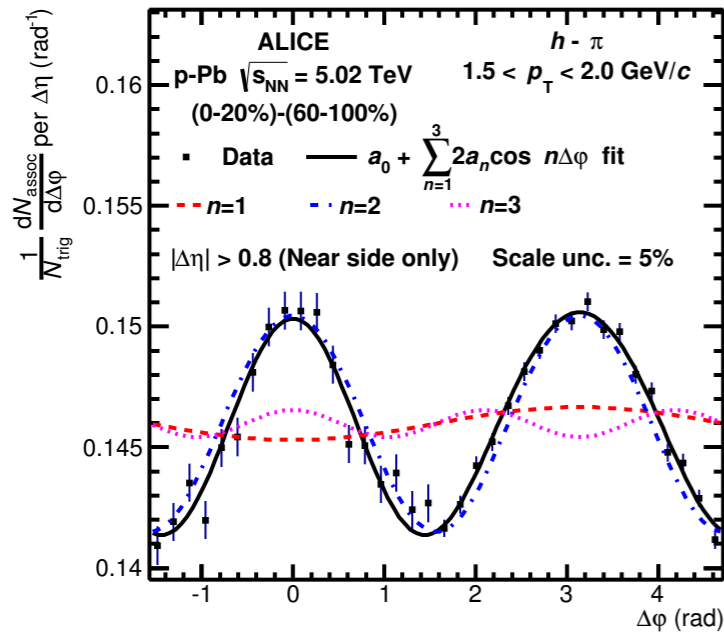
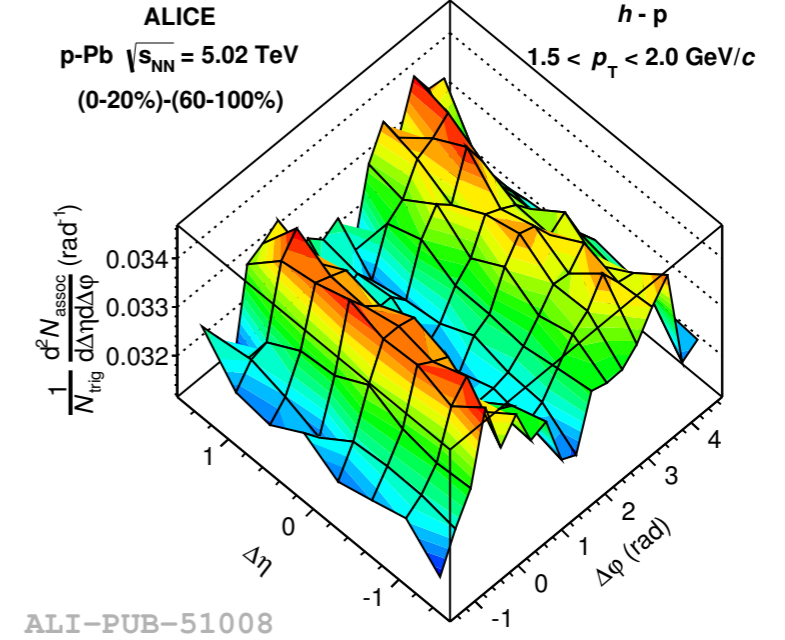
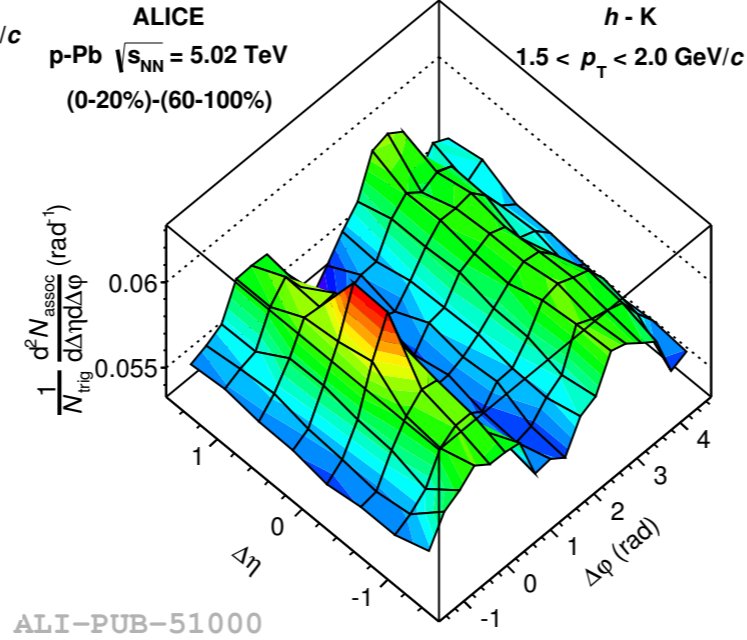
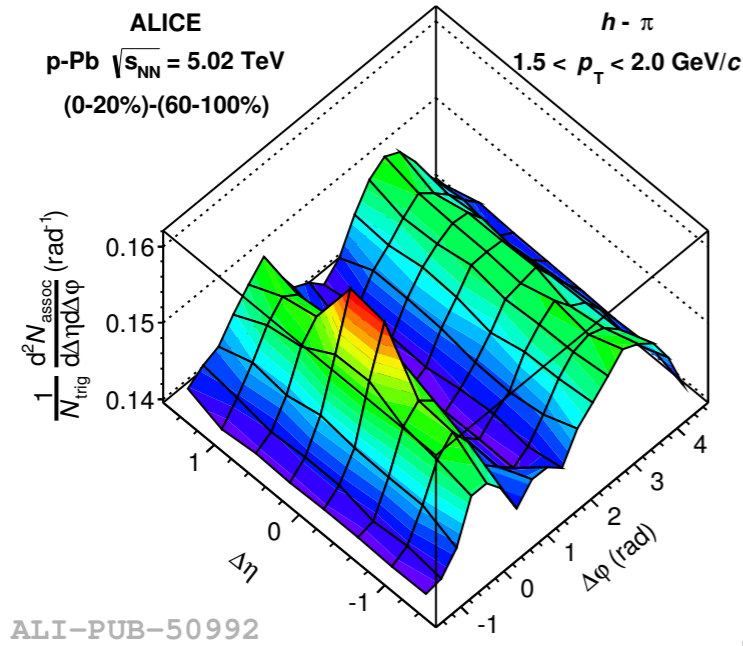
- Near-side yield is comparable to away-side yield.
 - Double ridge is reproduced by CGC.
- Extracted v_2, v_3 depend on p_T and centrality.

K. Dusling and R. Venugopalan, Phys.Rev. D87 (2013), 054014

$$V_{n\Delta}^{h-i}\{2PC\} = a_n^{h-i} / a_0^{h-i} \quad v_n^h\{2PC\} = \sqrt{V_{n\Delta}^{h-h}}$$

- v_2 and v_3 are qualitative agreement with hydro calculation.
- P. Bozek and W. Broniowski, Phys.Lett. B718 (2013) 1557-1561

Double ridge with PID



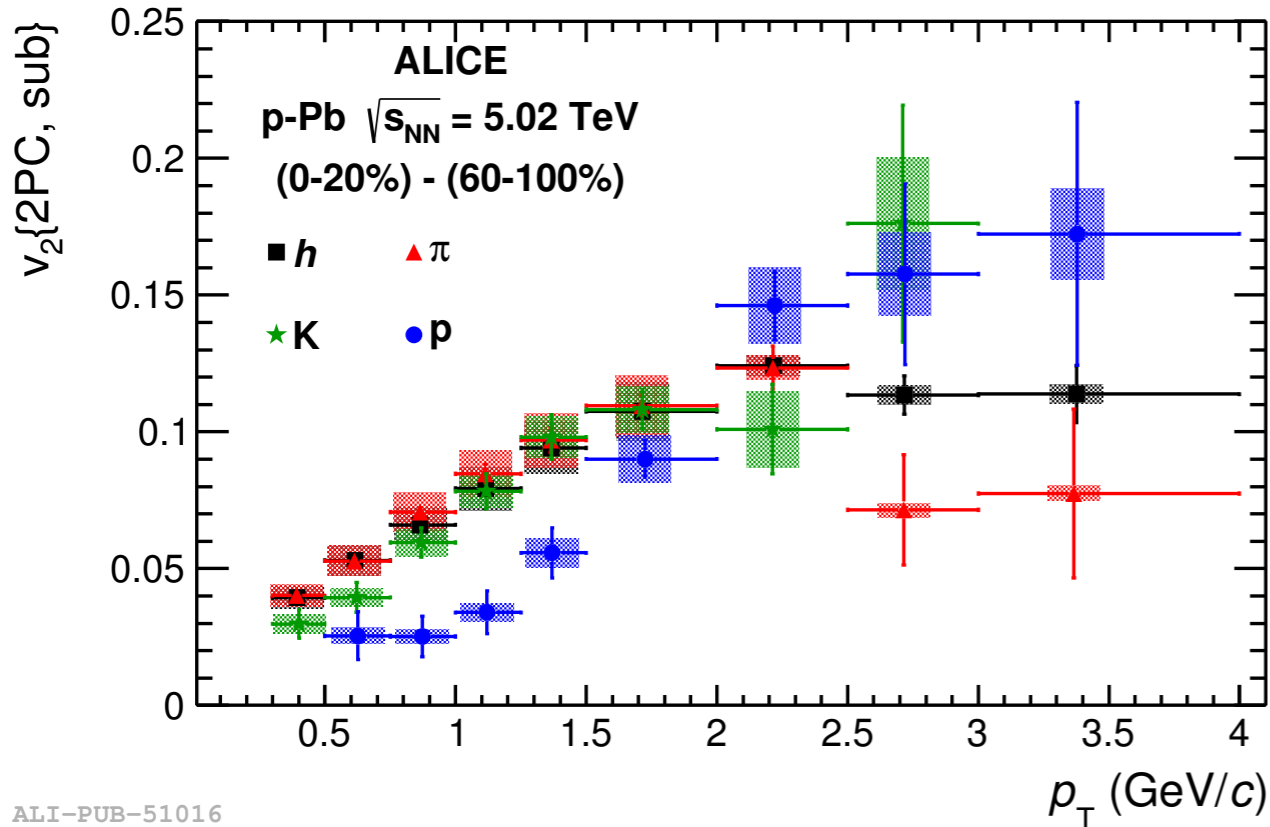
- Double ridge structure for π , K, and p.

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$$V_{n\Delta}^{h-i} \{2PC\} = a_n^{h-i} / a_0^{h-i} \quad v_n^i \{2PC\} = \frac{V_{n\Delta}^{h-i}}{\sqrt{V_{n\Delta}^{h-h}}}$$

v_2 of identified hadrons in pPb

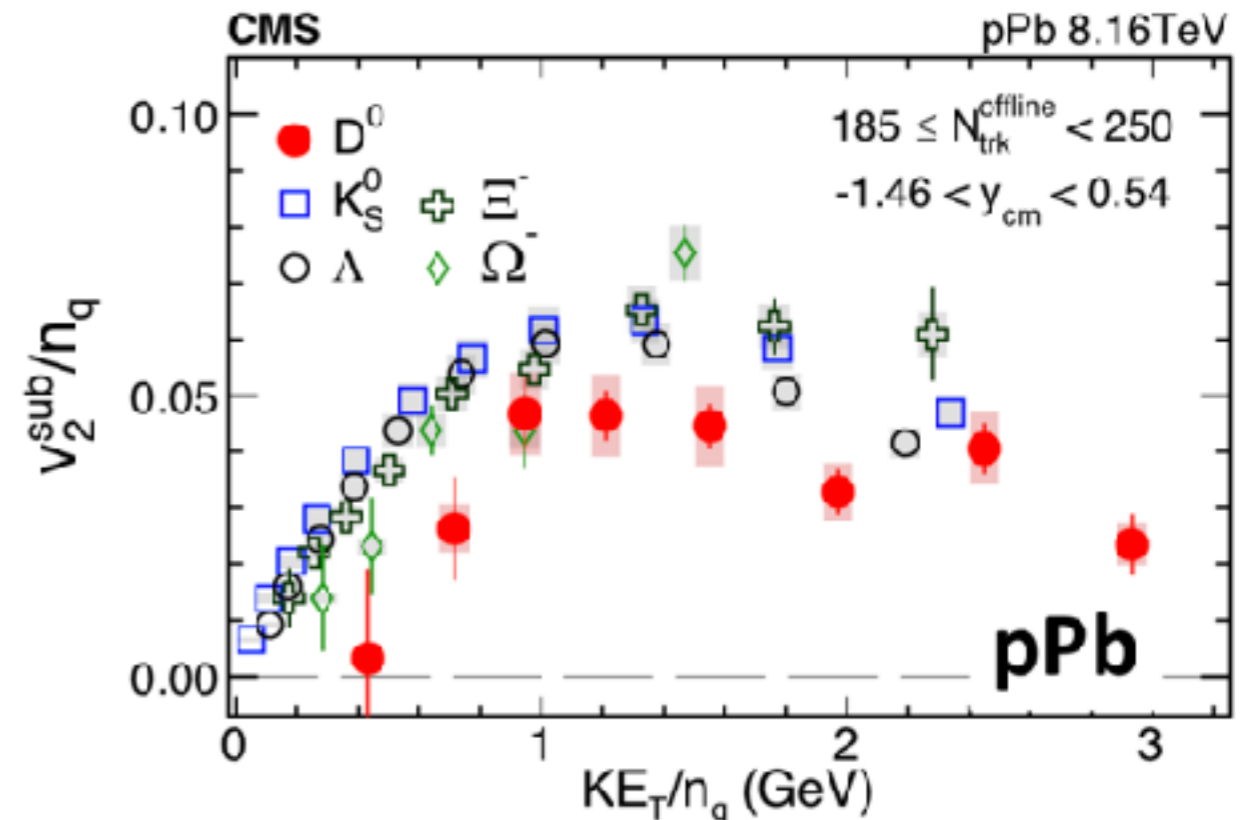
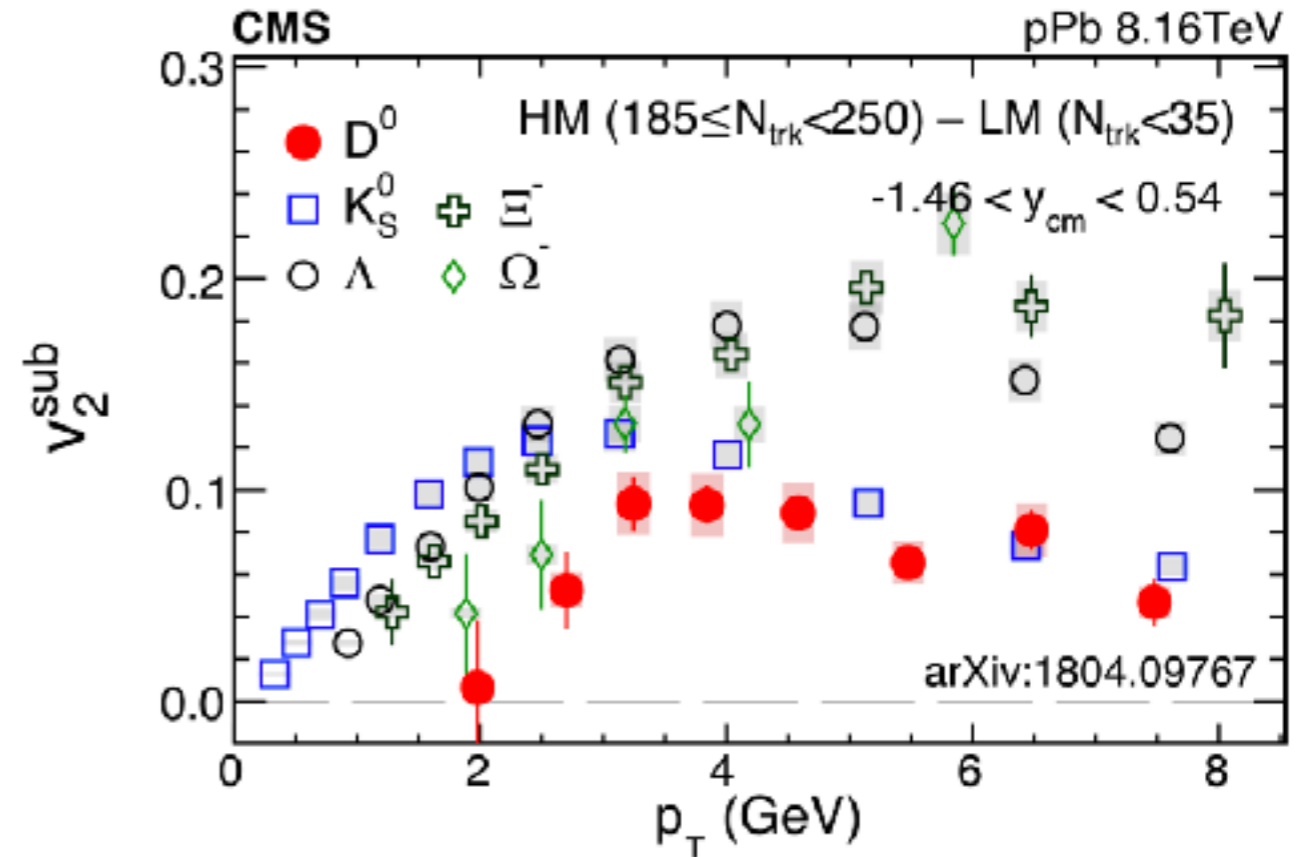
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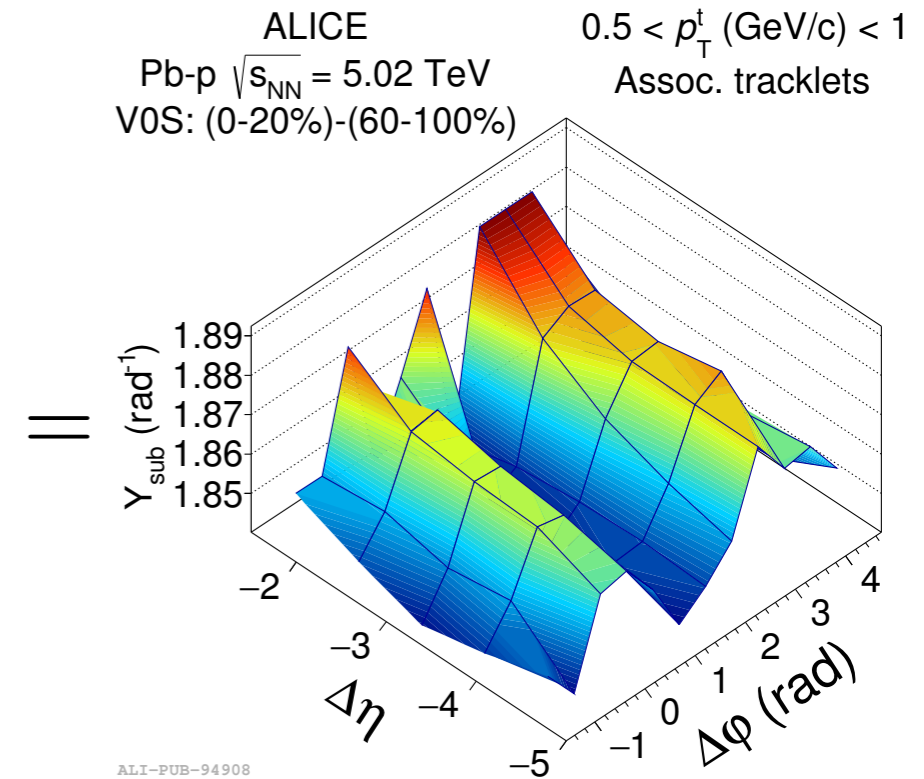
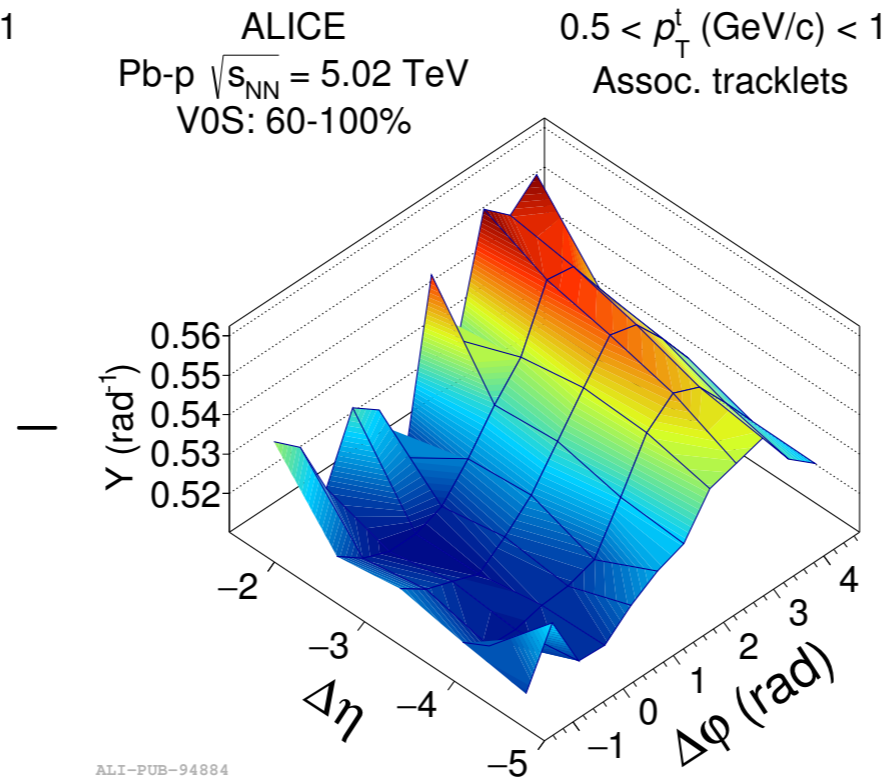
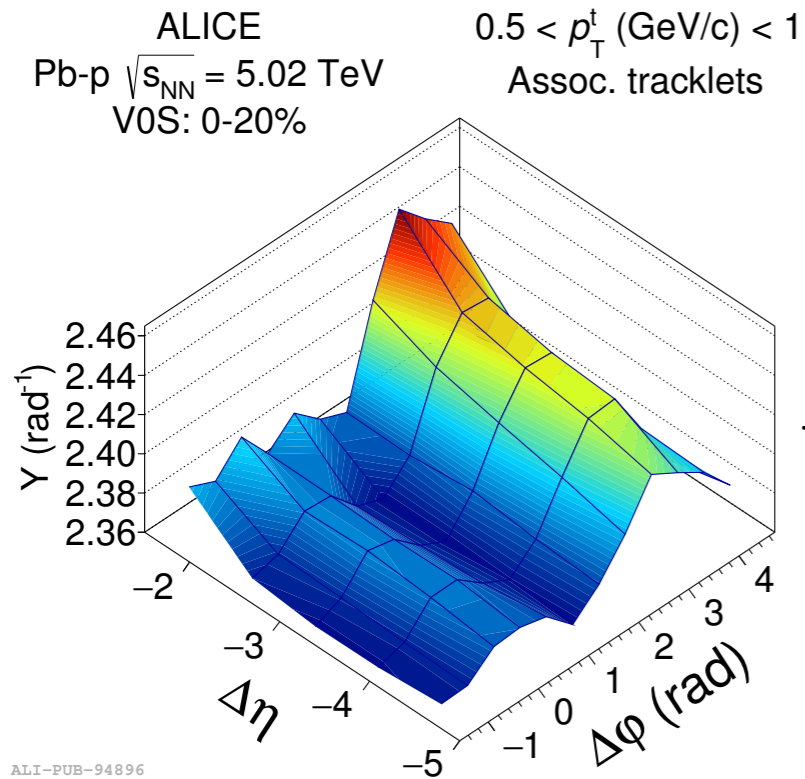
- Mass ordering at $p_T < 2$ GeV/c.
- Quark number scaling succeeded except D^0 .
 - Different trend from PbPb.

$$V_2^{sub} = V_2^{HM} - V_2^{LM} \frac{N_{assoc}^{LM}}{N_{assoc}^{HM}} \frac{Y_{jet}^{HM}}{Y_{jet}^{LM}}$$

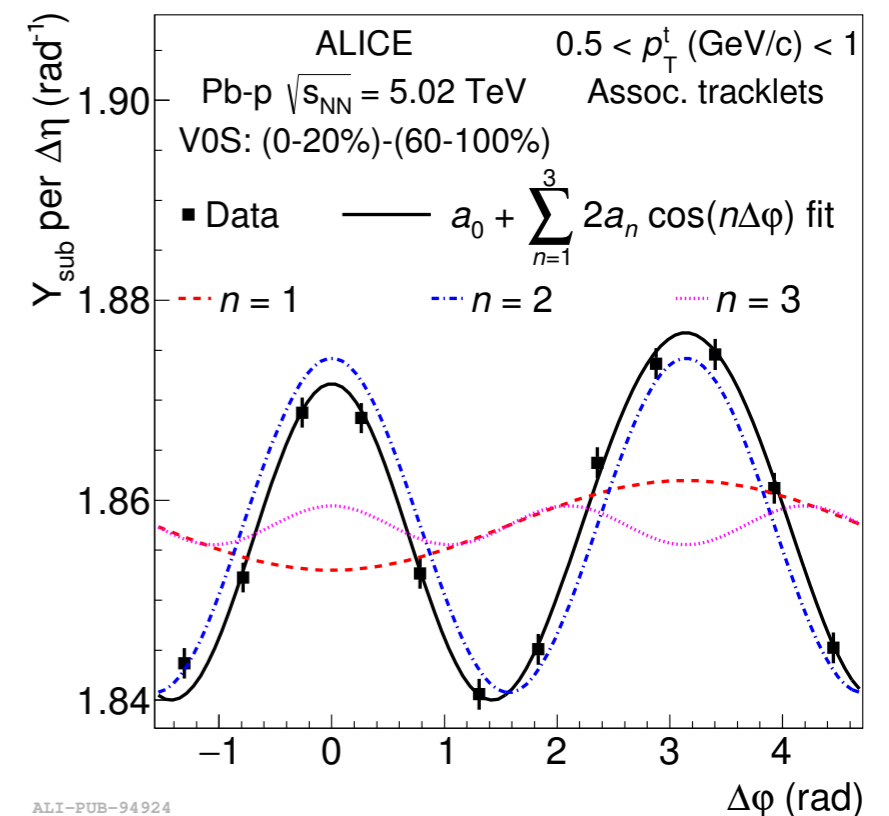


Forward/Backward-Central correlations

Phys. Lett. B 753 (2016) 126-139

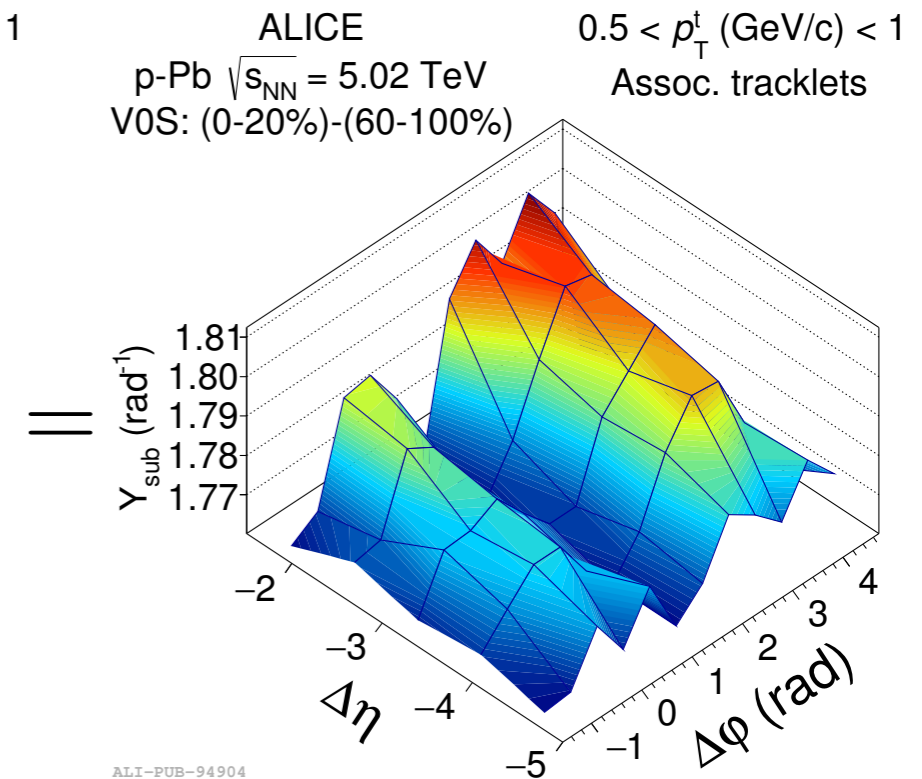
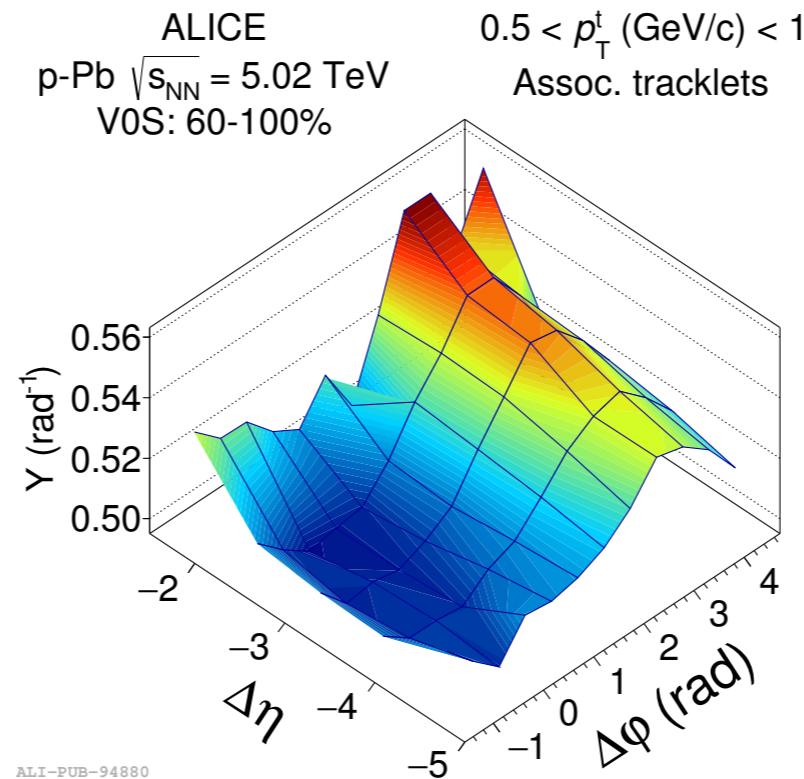
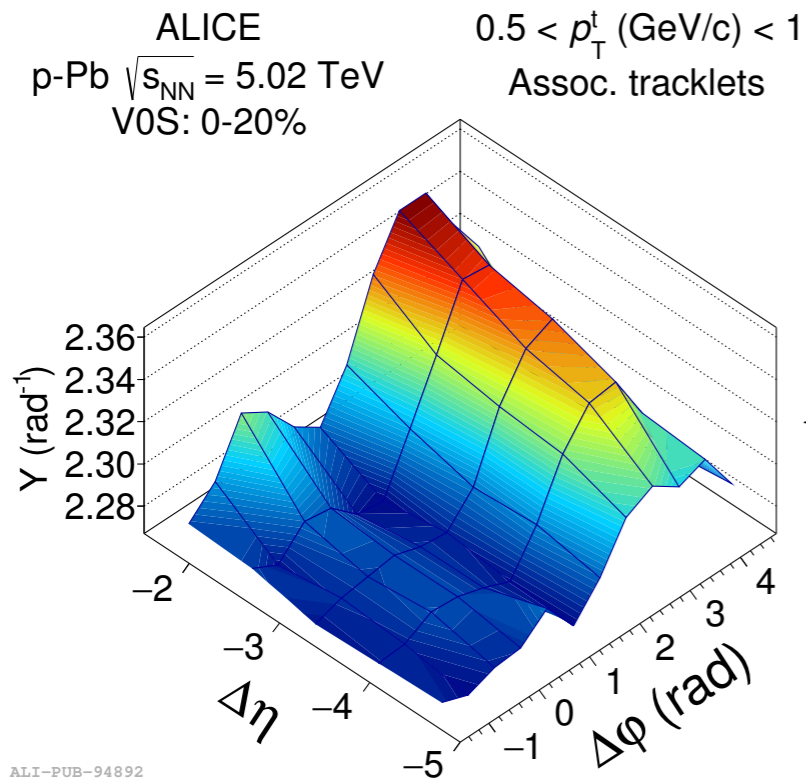


- Correlations between muons at forward rapidity at $-4 < \eta < -2.5$ in **Pb-going** and charged particle at mid rapidity.

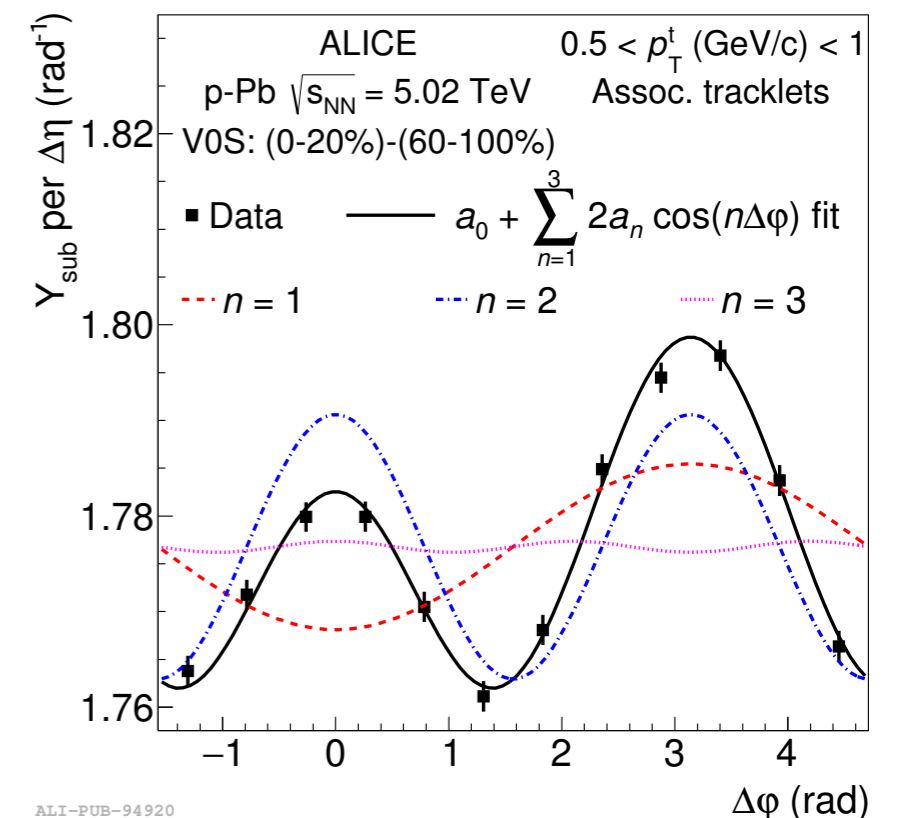


Forward/Backward-Central correlations

Phys. Lett. B 753 (2016) 126-139

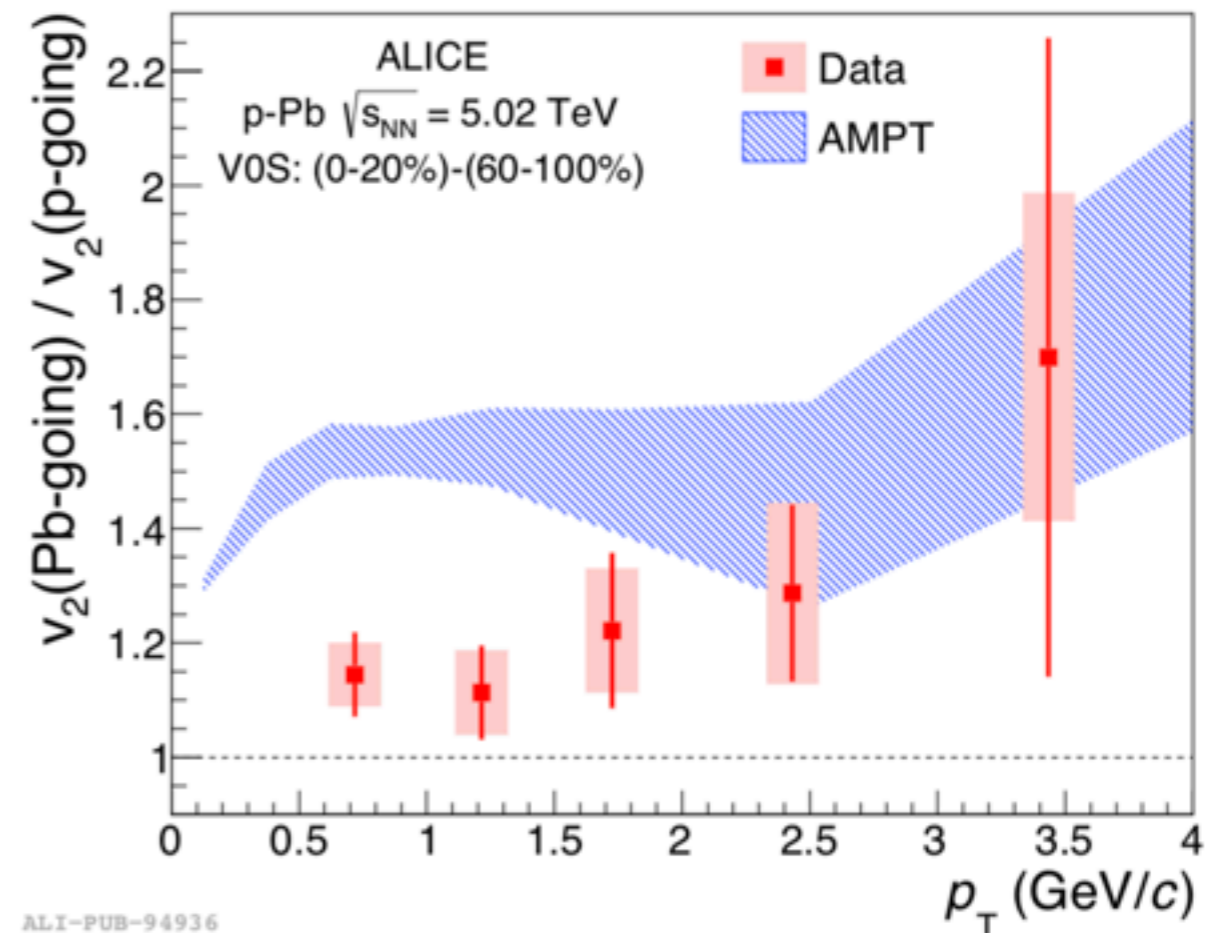
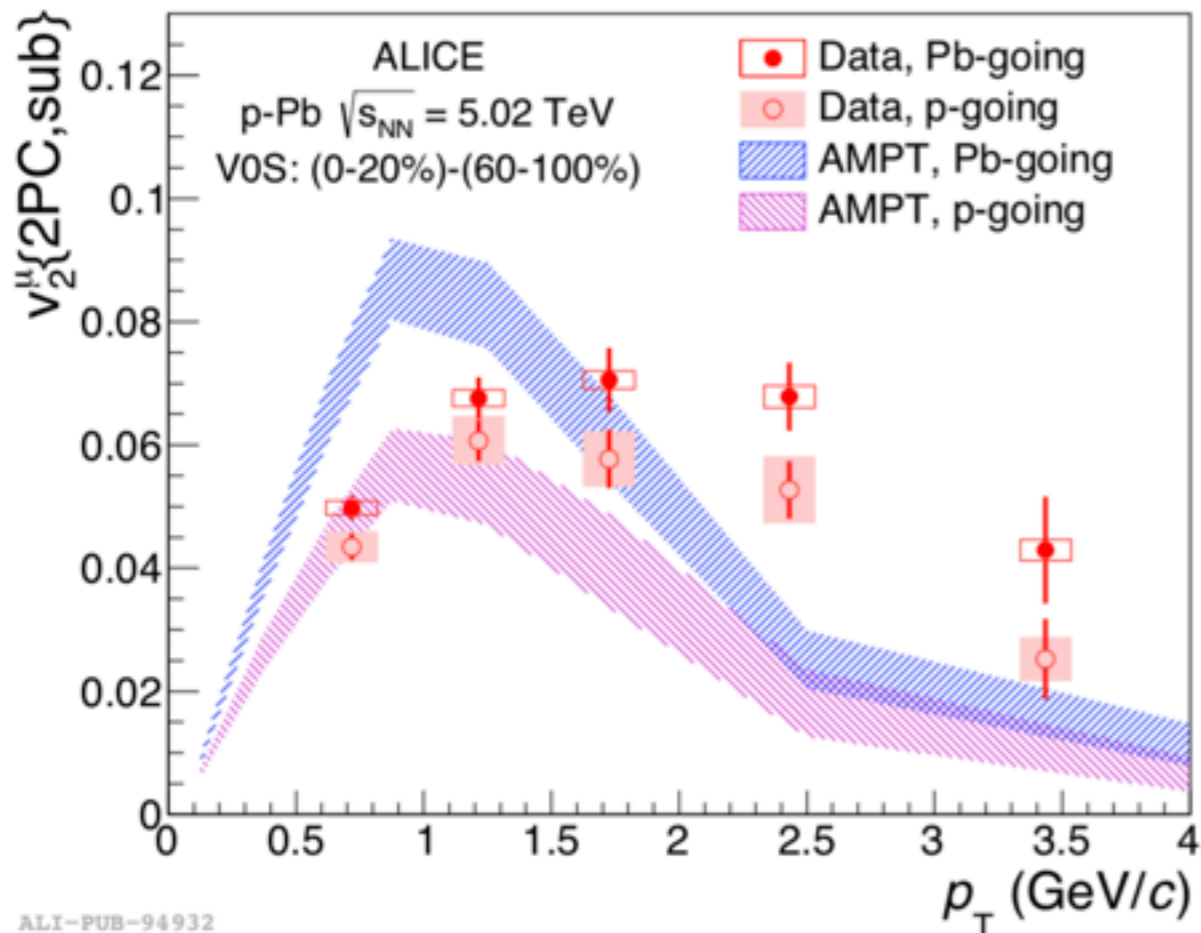


- Correlations between muons at forward rapidity at $-4 < \eta < -2.5$ in p-going and charged particle at mid rapidity.



v_2 at forward and backward rapidity in pPb collisions

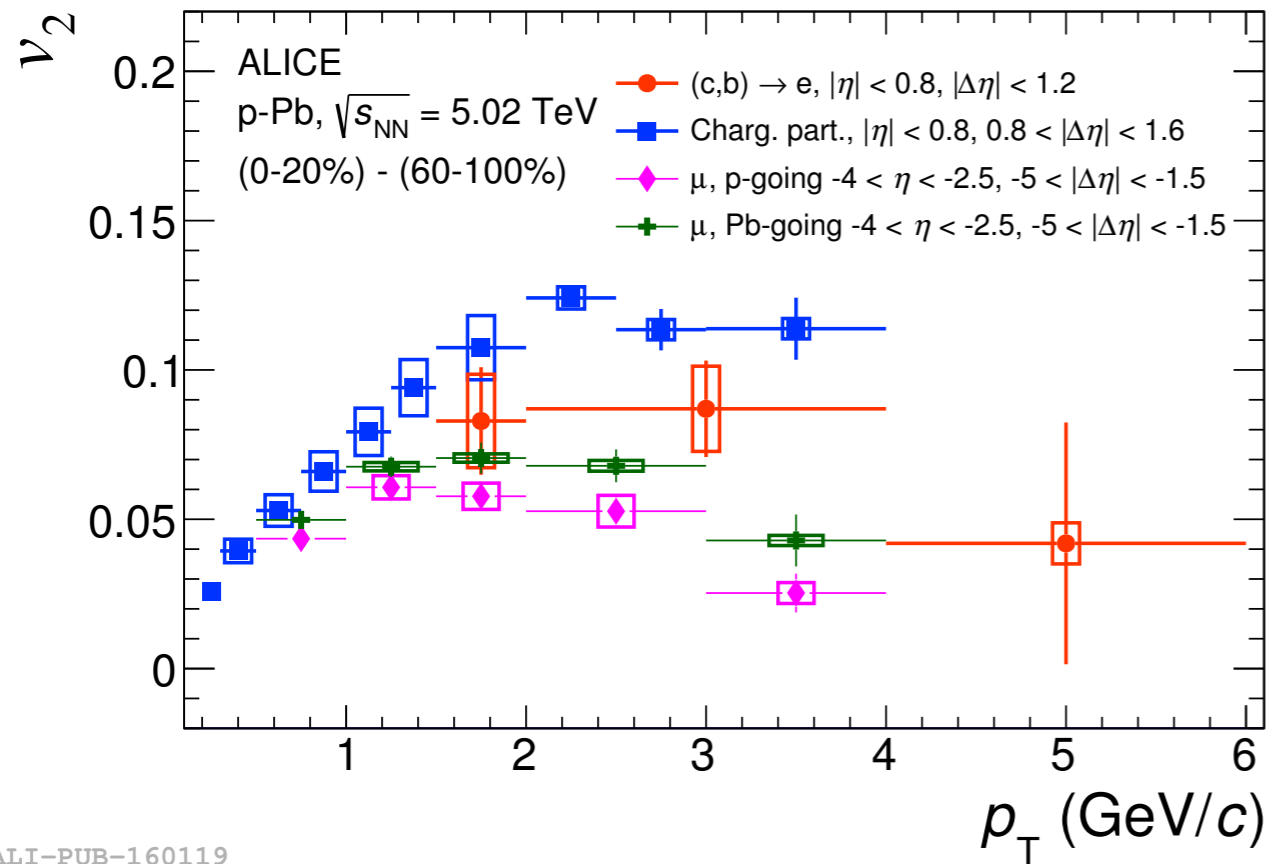
Phys. Lett. B 753 (2016) 126-139



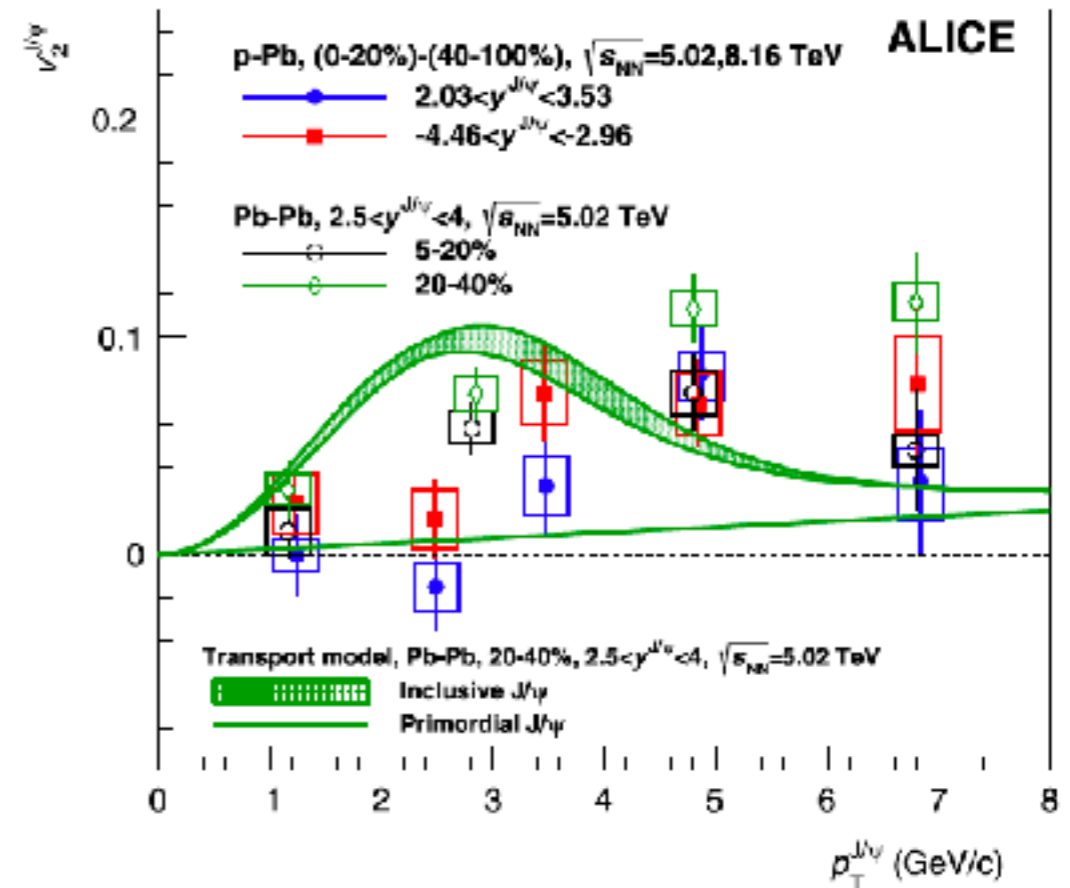
- v_2 in Pb going is larger by $16\% \pm 6\%$ than in p going.
- p going side probe small-x in Pb while large x in Pb going.
- Multiplicity in Pb going is larger than in p going.
- AMPT trend is similar to the data.

v_2 of Heavy flavor decay electron and J/ ψ

arXiv:1805.0436



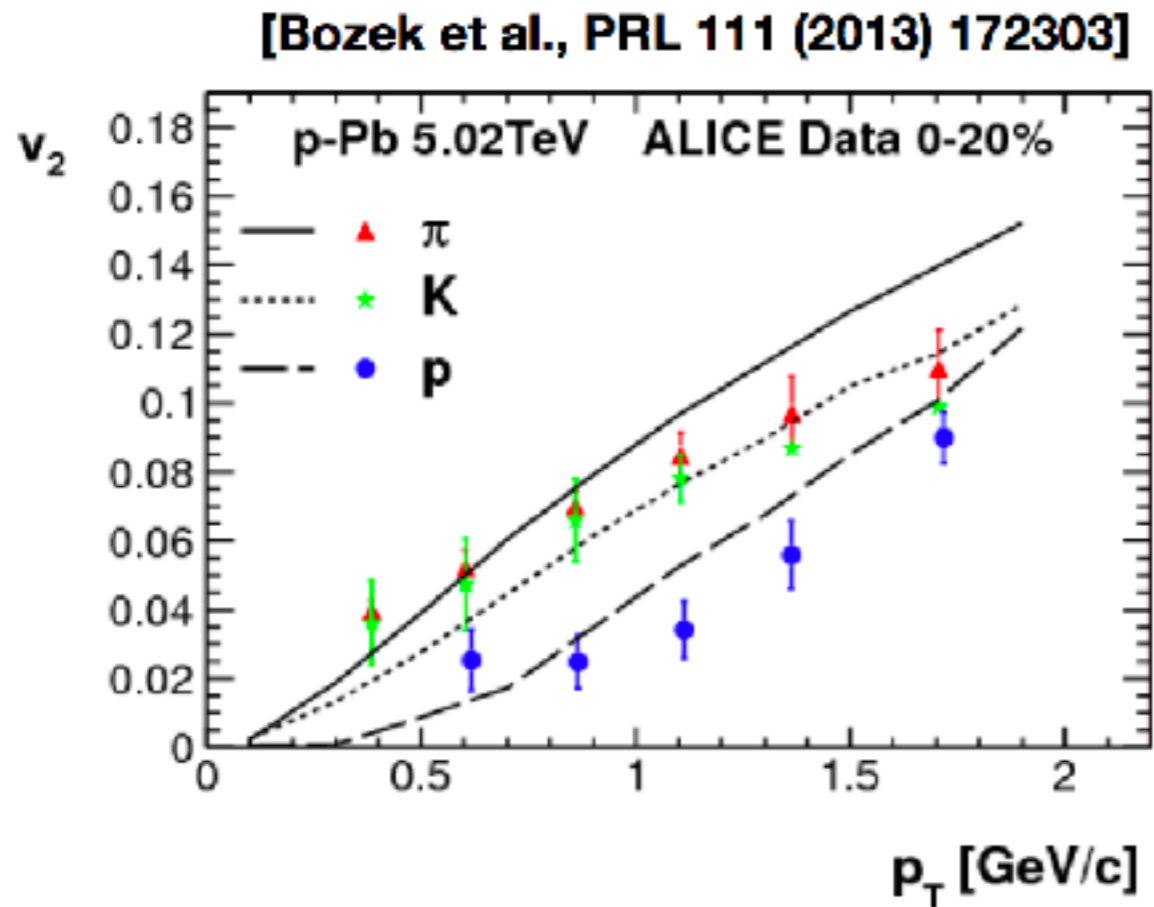
Physics Letters B 780 (2018)



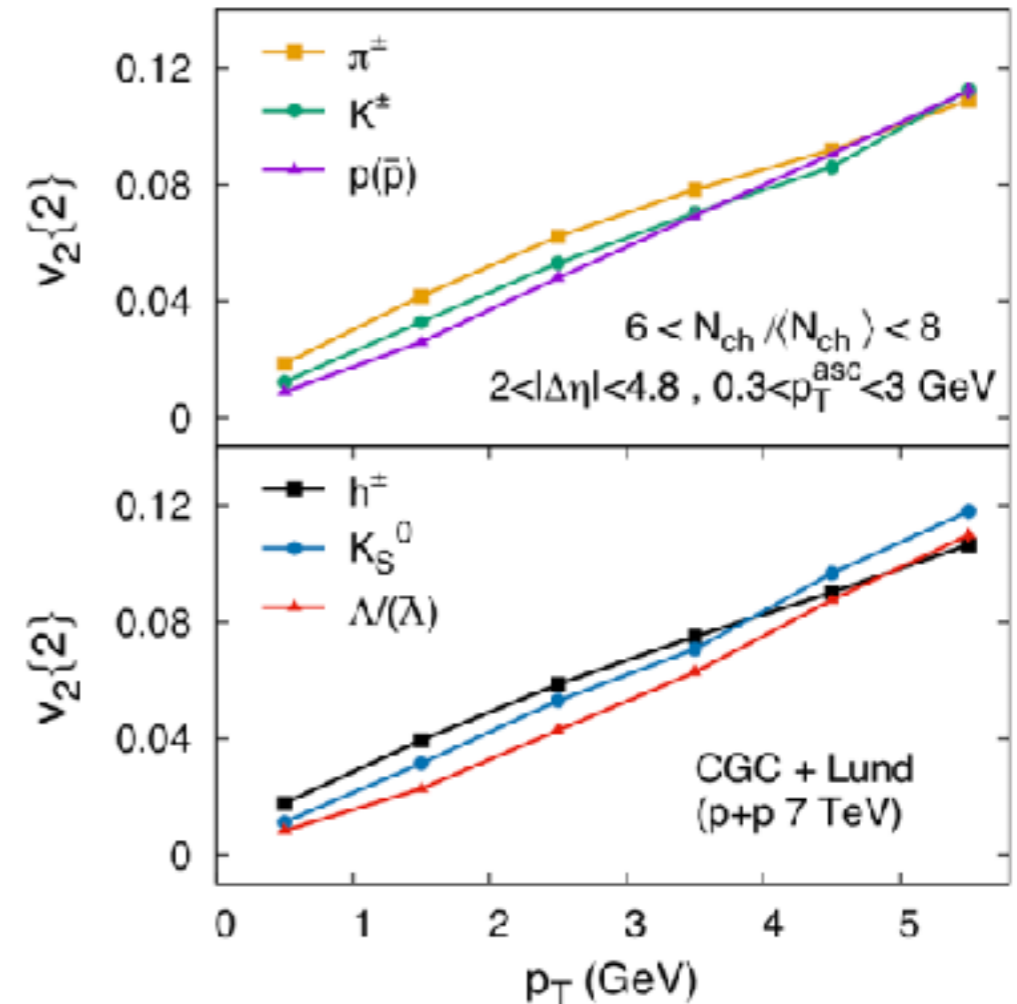
ALI-PUB-160119

- HFe v_2 is significant with more than 5σ at $1.5 < p_T < 4$.
- Significant J/ ψ v_2 . Evidence of charm flow.
- Large v_2 even at high p_T .

Model comparison



arXiv:1607.02496



- Hydro describes data in pPb.
- Initial effects only reproduce mass ordering in pp.
 - qualitatively similar to CMS results in pp for h , k_s^0 , and Λ .
- Need to disentangle initial and final state dynamics.
 - $v_2(\eta)$, correlation of different harmonics, event plane decorrelation.
 - $v_2(\eta)$ is sensitive to longitudinal dynamics.
 - Correlation of different EP are sensitive to transverse + longitudinal fluctuation.

Event plane decorrelations in PbPb

Phys. Rev. C 92 (2015)

- $r_2(\eta^a, \eta^b)$ is defined as

$$r_n(\eta^a, \eta^b) \equiv \frac{V_{n\Delta}(-\eta^a, \eta^b)}{V_{n\Delta}(\eta^a, \eta^b)} = \frac{\langle v_n(-\eta^a) v_n(\eta^b) \cos\{n[\Psi_n(-\eta^a) - \Psi_n(\eta^b)]\} \rangle}{\langle v_n(\eta^a) v_n(\eta^b) \cos\{n[\Psi_n(\eta^a) - \Psi_n(\eta^b)]\} \rangle}$$

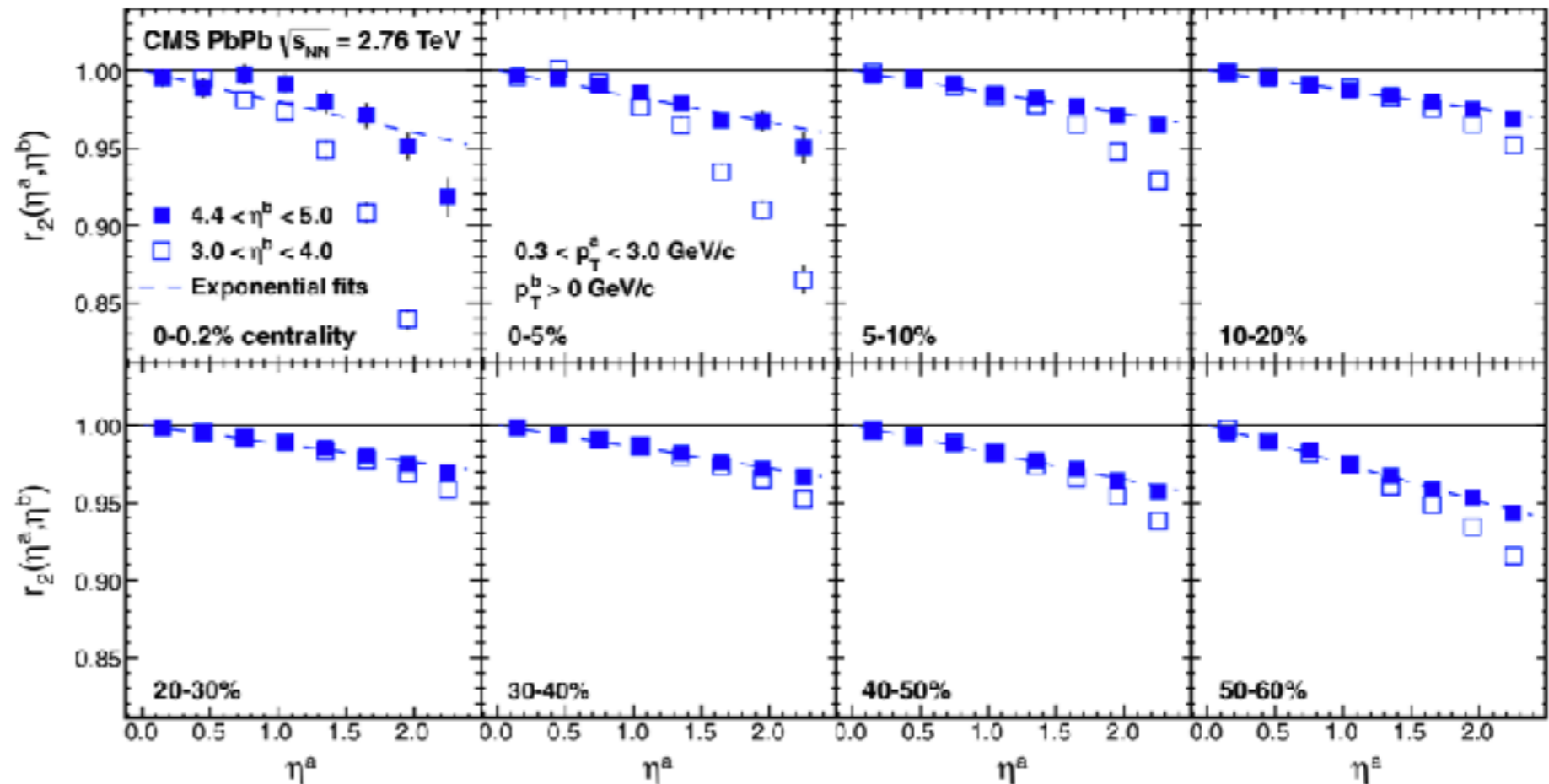
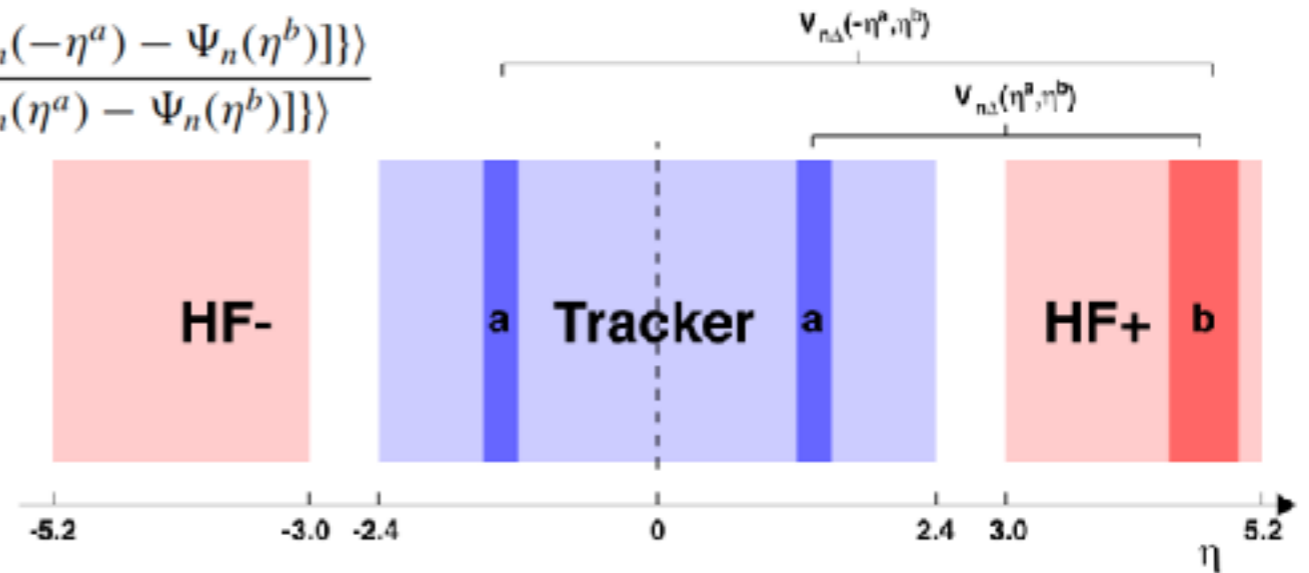
- Simple empirical parametrization is introduced

$$\cos\{n[\Psi_n(\eta^a) - \Psi_n(\eta^b)]\} = e^{-F_n^\eta |\eta^a - \eta^b|}$$

- $r_2(\eta^a, \eta^b)$ can be expressed as bellow.

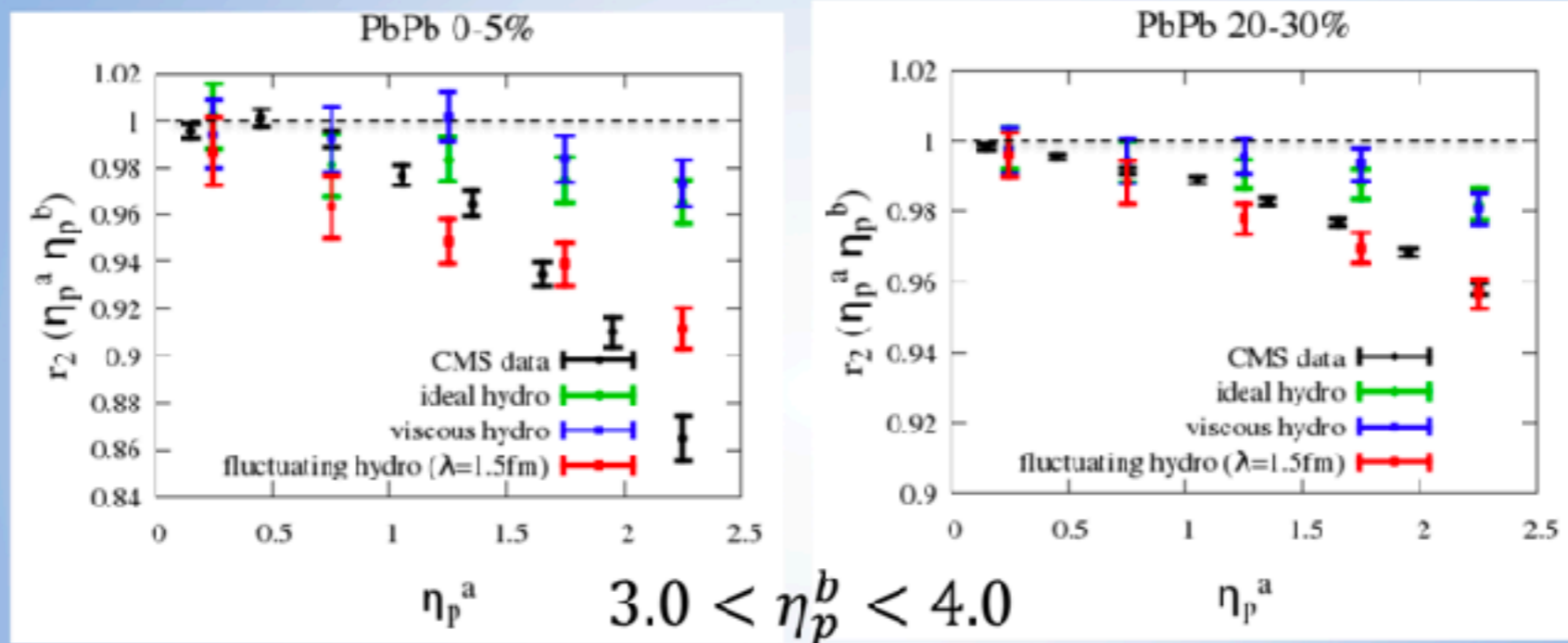
$$r_n(\eta^a, \eta^b) \approx e^{-2F_n^\eta \eta^a}$$

- $r_2(\eta^a, \eta^b)$ decrease significant decrease linearly as η^a increase.



Decorrelation from final effect

Factorization ratio $r_2(\eta_p^a, \eta_p^b)$



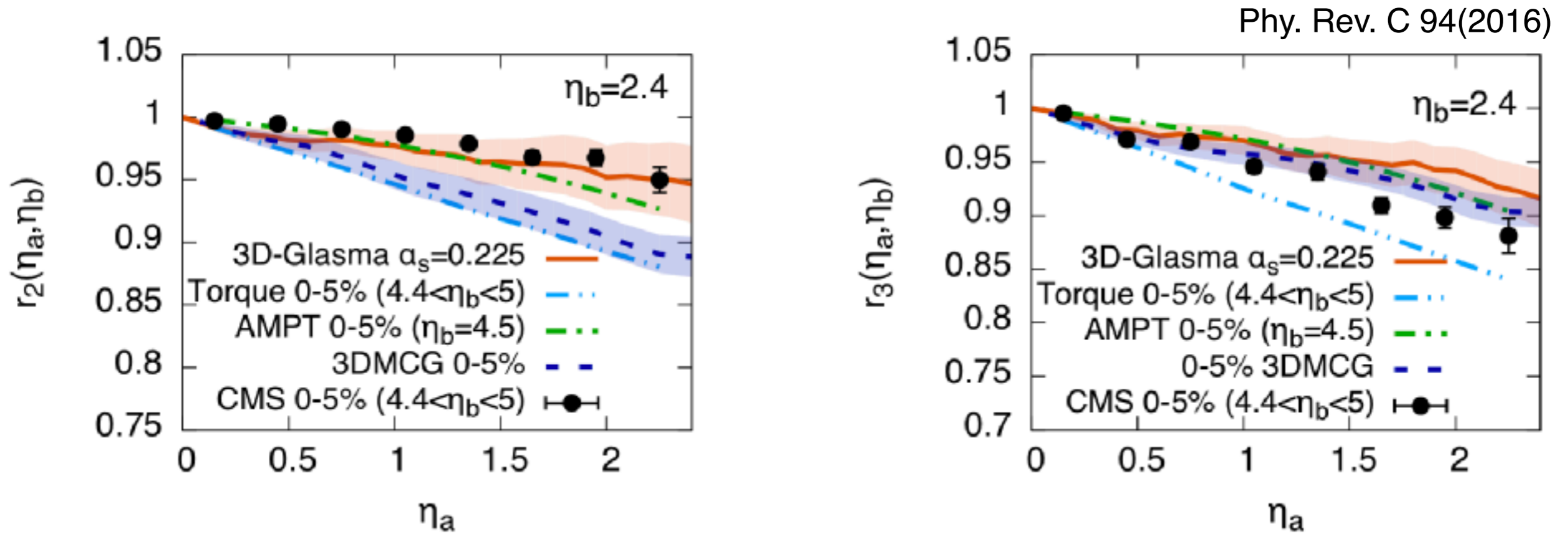
Ideal \approx Viscous $>$ CMS data \approx Fluctuating hydro
 Hydrodynamic fluctuations
 \rightarrow Factorization more broken

CMS, Phys. Rev. C 92, 034911 (2015).

QM2018
 A. Sakai

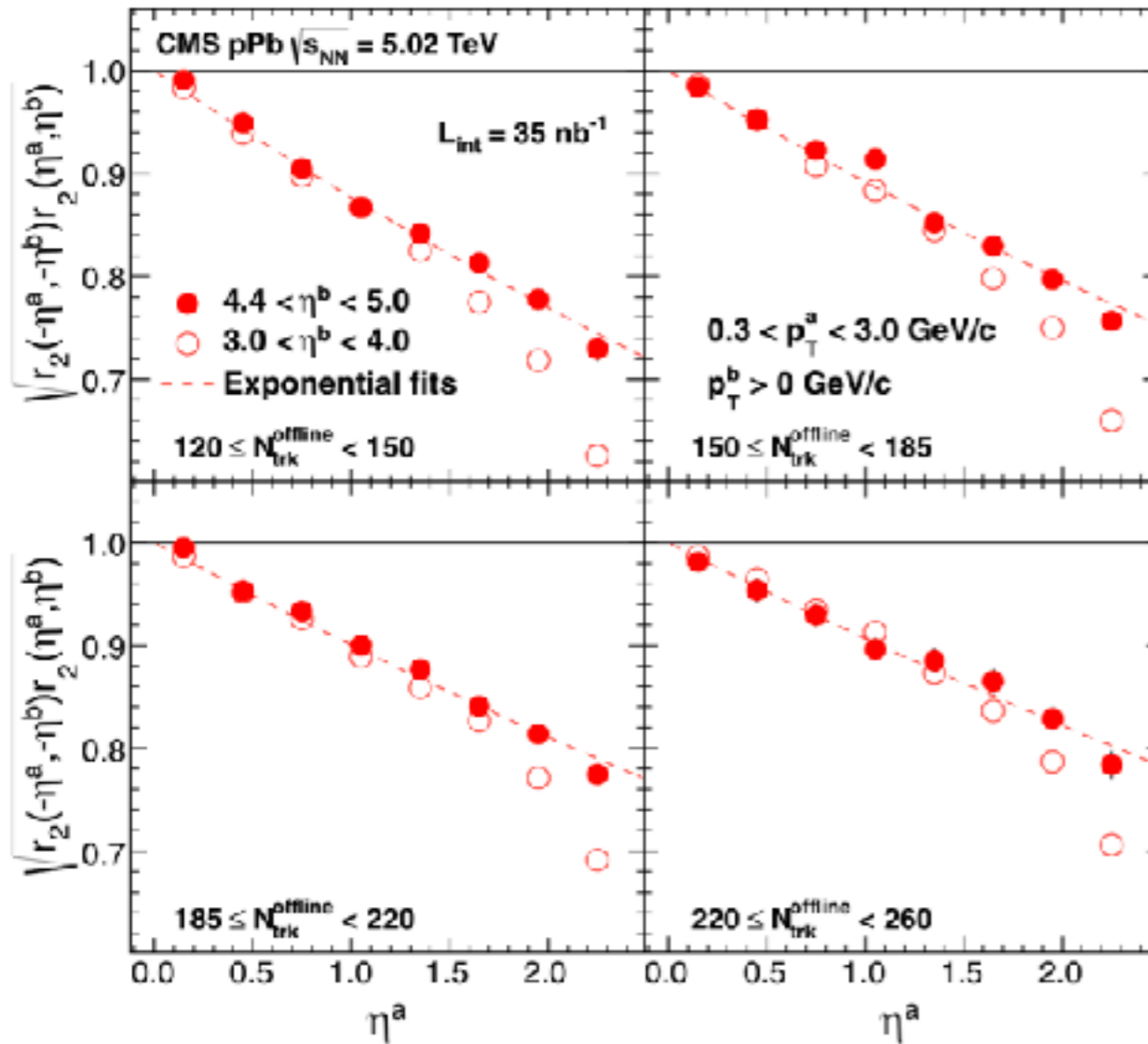
- Hydrodynamic fluctuation reproduce the CMS data.
- Importance to understand the final state fluctuation (much less in pPb collisions?)

Decorrelation from initial fluctuation

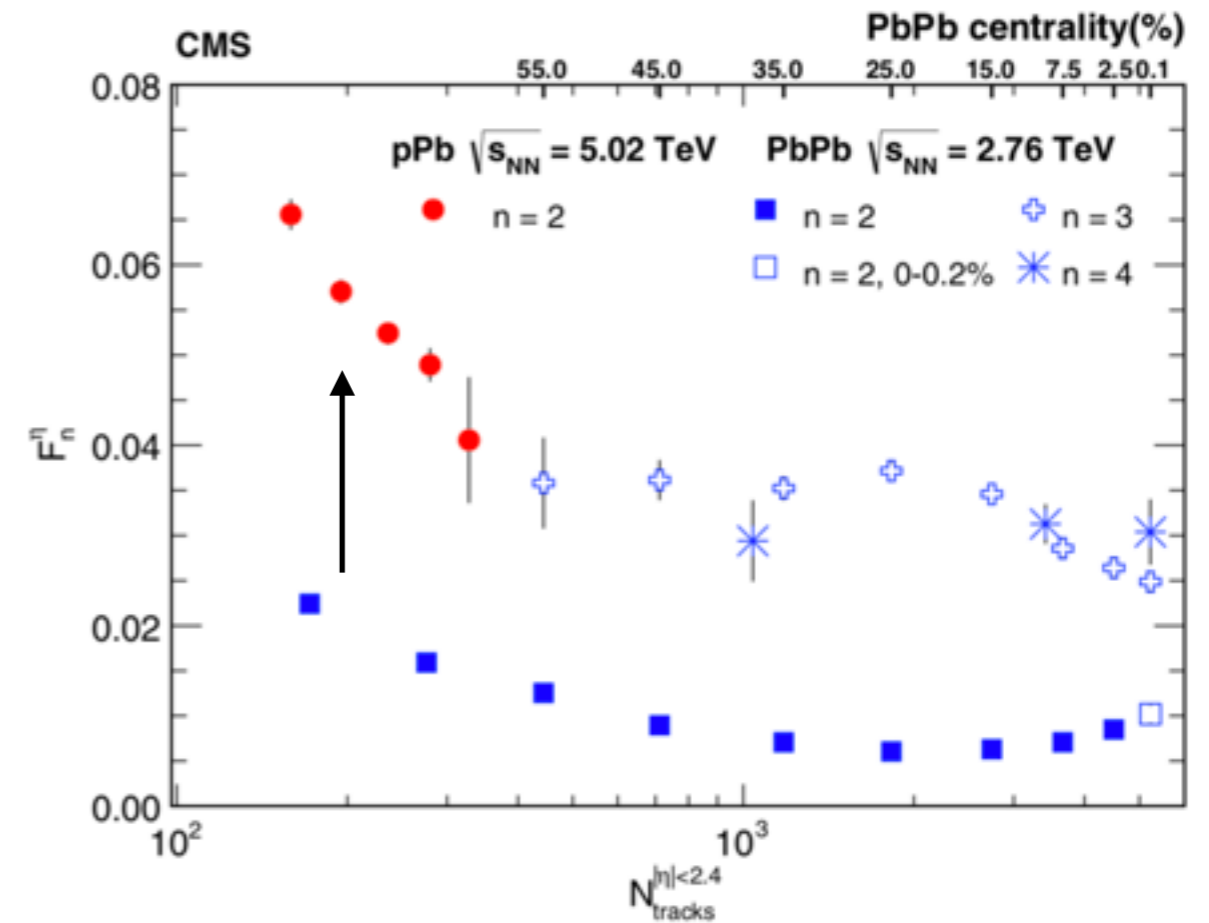


- $r_2(\eta^a, \eta^b)$ is reproduced with the model only with initial fluctuation.
- Both final fluctuation and initial fluctuation can describe data for PbPb.
- Initial fluctuation is much important for pPb.

Event plane decorrelation in pPb



Phys. Rev. C 92 (2015)



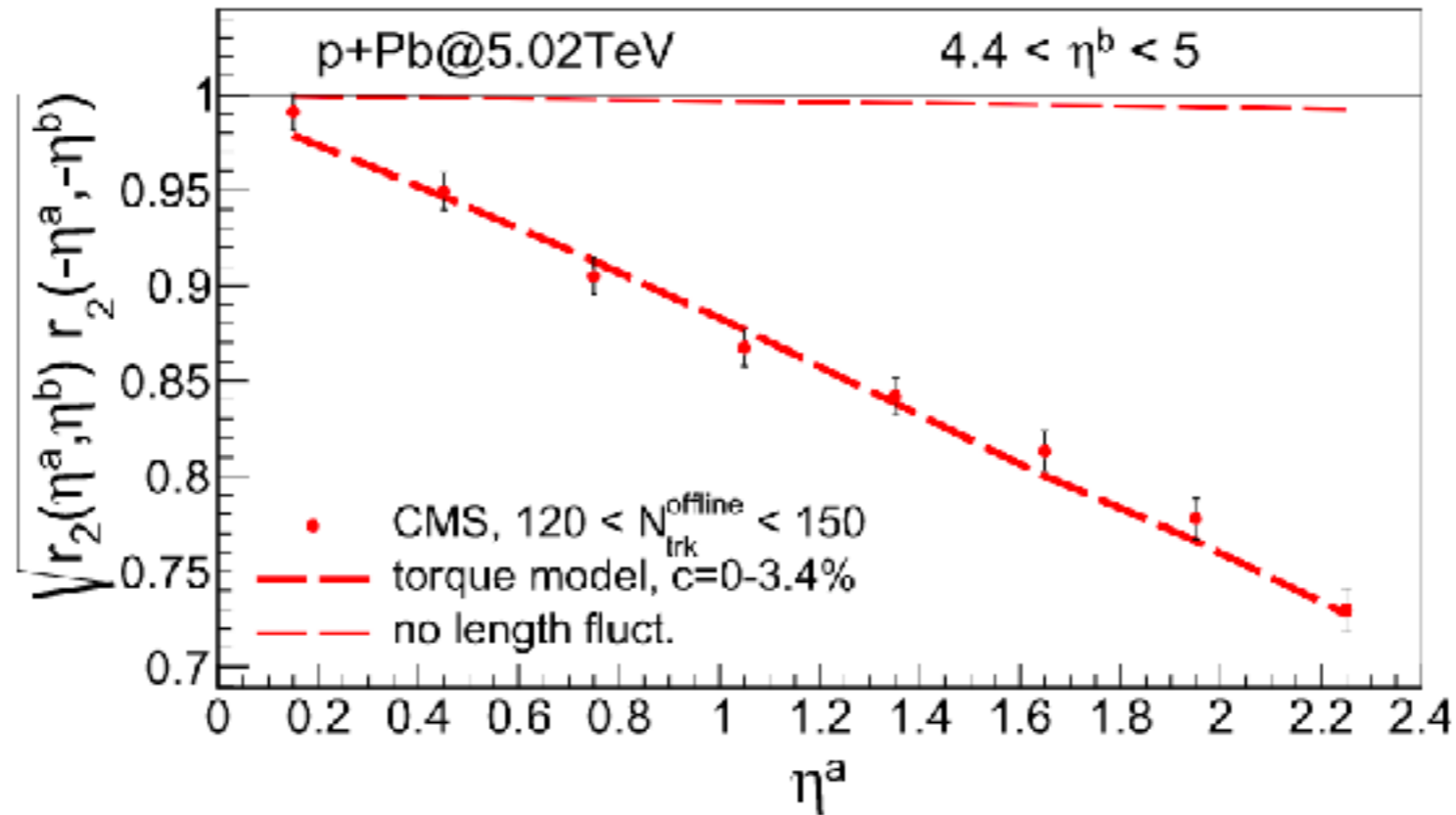
- Use following equation to remove magnitude v_2 .

$$\sqrt{r_n(\eta^a, \eta^b)r_n(-\eta^a, -\eta^b)} \approx \sqrt{\frac{\langle \cos\{n[\Psi_n(-\eta^a) - \Psi_n(\eta^b)]\} \rangle}{\langle \cos\{n[\Psi_n(\eta^a) - \Psi_n(\eta^b)]\} \rangle} \frac{\langle \cos\{n[\Psi_n(\eta^a) - \Psi_n(-\eta^b)]\} \rangle}{\langle \cos\{n[\Psi_n(-\eta^a) - \Psi_n(-\eta^b)]\} \rangle}} \approx e^{-2F_n^\eta \eta^a}$$

- F_n^η in pPb is larger than in PbPb.

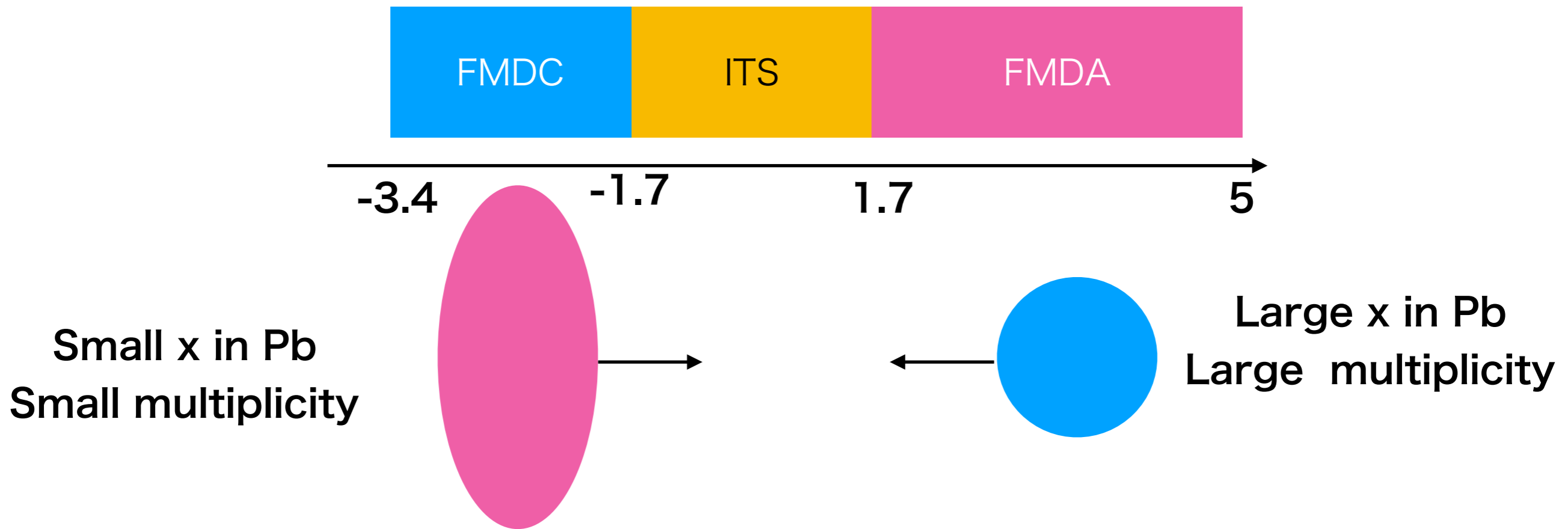
Event plane decorrelation in pPb

Physics Letters B 752 (2016)

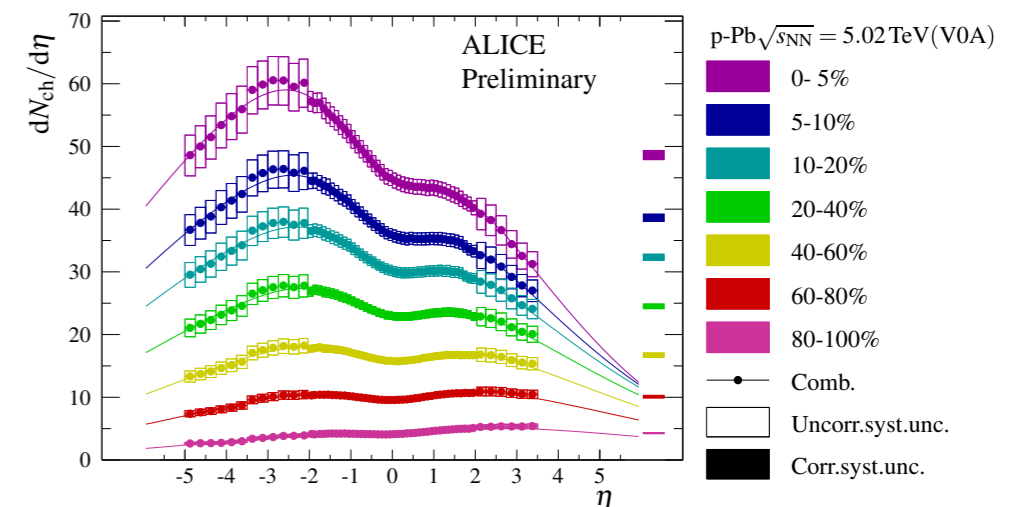


- Initial fluctuation, Toque model, describe data well.
- The model describe also r_2 and r_3 of PbPb for semi-central and peripheral collisions.
- Further accurate simulation with 3+1D hydro and nonflow.

Outlook: Central-Forward-Backward correlations in pPb collisions with ALICE



- $v_2(\eta)$ with Central-forward, Central-backward, and Forward-Backward.
 - Suppress non-flow by large eta gap.
- Correlation of Event plane
- Correlation of v_2 and v_3 at forward rapidity.



ALI-PREL-99853

Summary

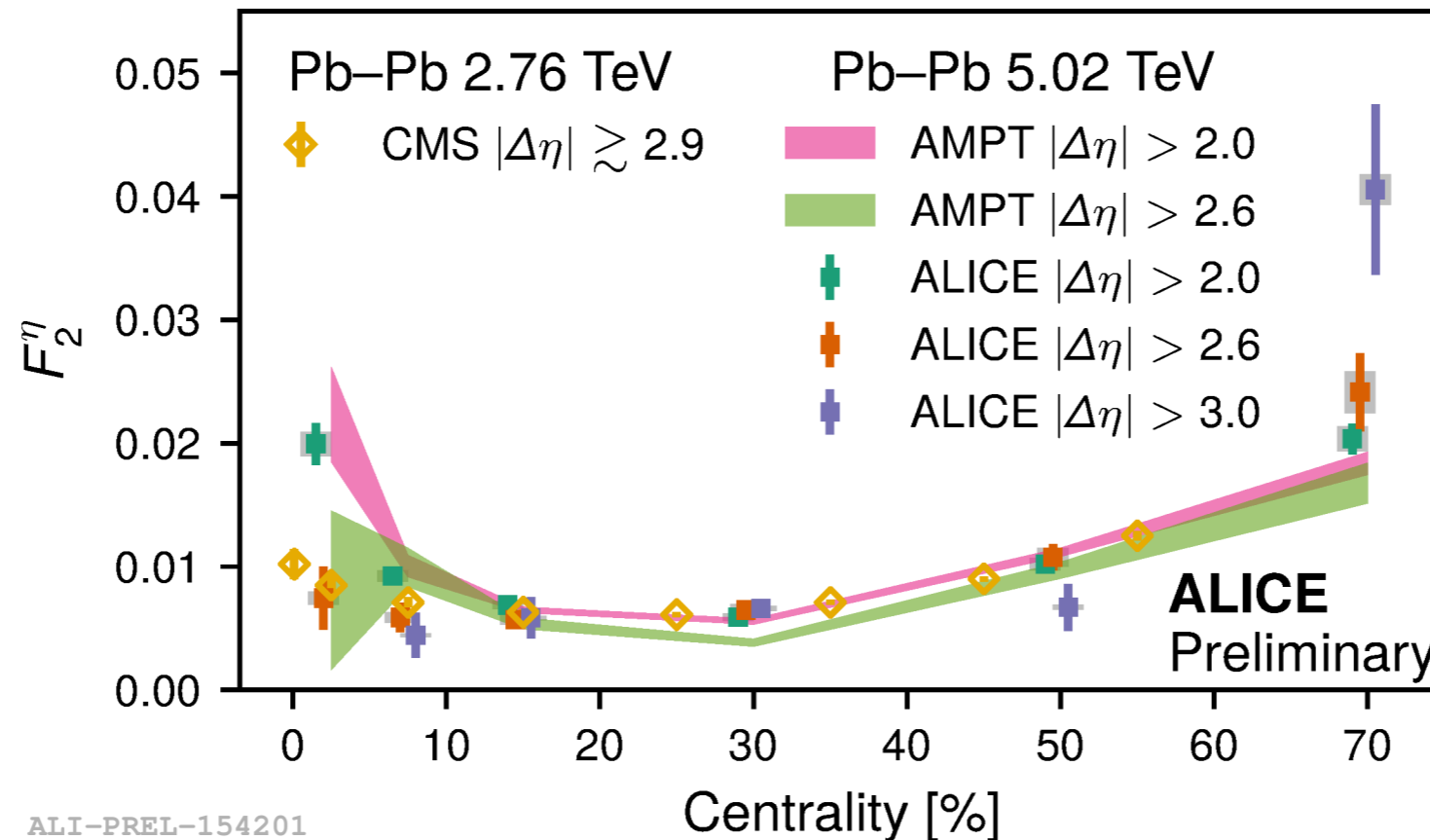
- Long range correlation is observed in p-Pb collisions.
 - Mass ordering in mid rapidity.
 - Similar to Pb-Pb collisions.
 - Hydro and CGC described data.
 - Significant v_2 of heavy flavor decay electrons.
 - Significant v_2 of J/ψ .
 - Correlation between forward muon and charged hadrons
 - v_2 in Pb going is larger by $16\% \pm 6\%$ than in p going.
 - AMPT trend similar to data at low p_T .
 - Event plane decorrelation
 - Empirical parameter F_n^η quantifies a $\Delta\eta$ -dependent event-plane decorrelation
 - Hydrodynamic fluctuation and initial fluctuation describe the data.
- Further measurements are needed to understand origin of ridge.

backup

Event plane decorrelation

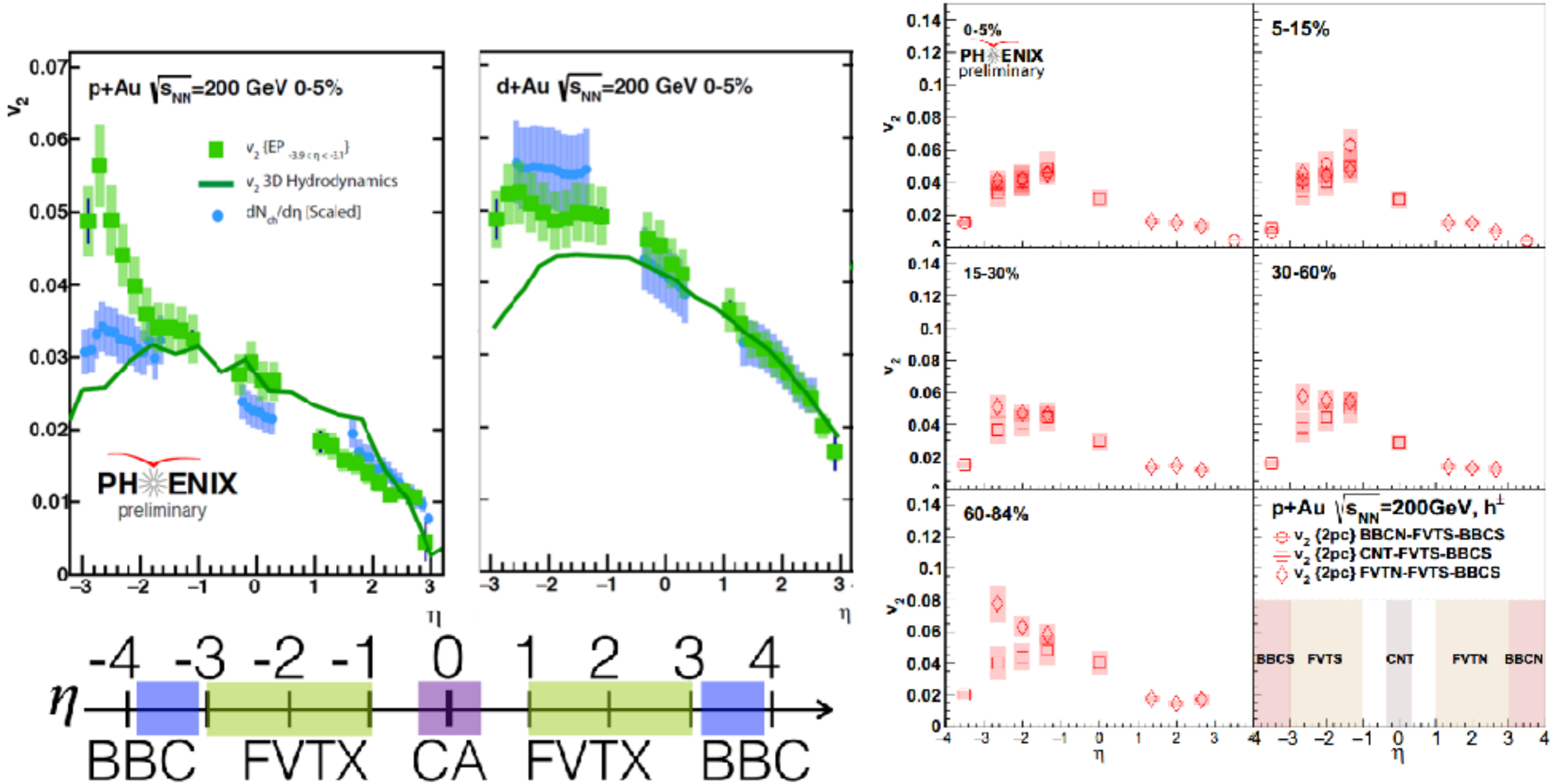
- $v_n(\eta^a, \eta^b)$ is assumed to be factorized into flow coefficients and a simple event plane decorrelation factor.

$$\hat{v}_n(\eta^a, \eta^b) = v_n(\eta^a)v_n(\eta^b)e^{-F_n^\eta|\eta^a - \eta^b|}$$



- Empirical parameter F_n^η quantifies a $\Delta\eta$ -dependent event-plane decorrelation.
- F_n^η depends on centrality.
- Consistent with CMS results.
- AMPT also describe the data.
 - Initial fluctuation is propagated to decorrelation.

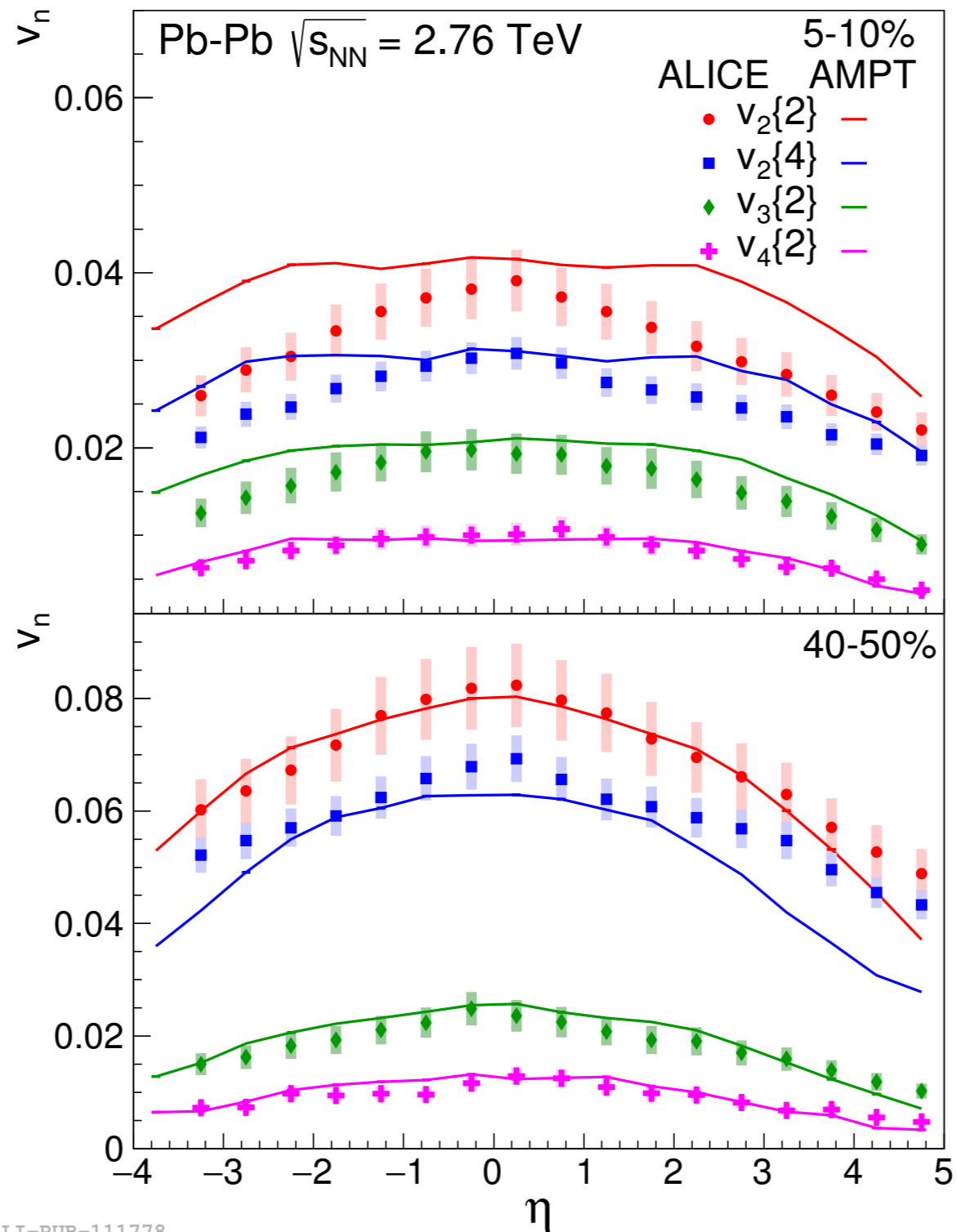
v_2 as a function of pseudo-rapidity in small system in RHIC



- Au going is larger than p/d going.
- 3D hydro describes d-Au well and p-Au except $\eta < -2$ in Au-going.
- 2 particle correlations with large η gap can suppress non-flow effects.

Investigate initial condition in PbPb

- The understanding of the initial condition is crucial for the subsequent hydrodynamic expansion.



Perfect factorization model

- The two-particle Fourier coefficients $v_{n,n}(\eta^a, \eta^b)$ are computed from the reduced two-particle distribution r_2 .

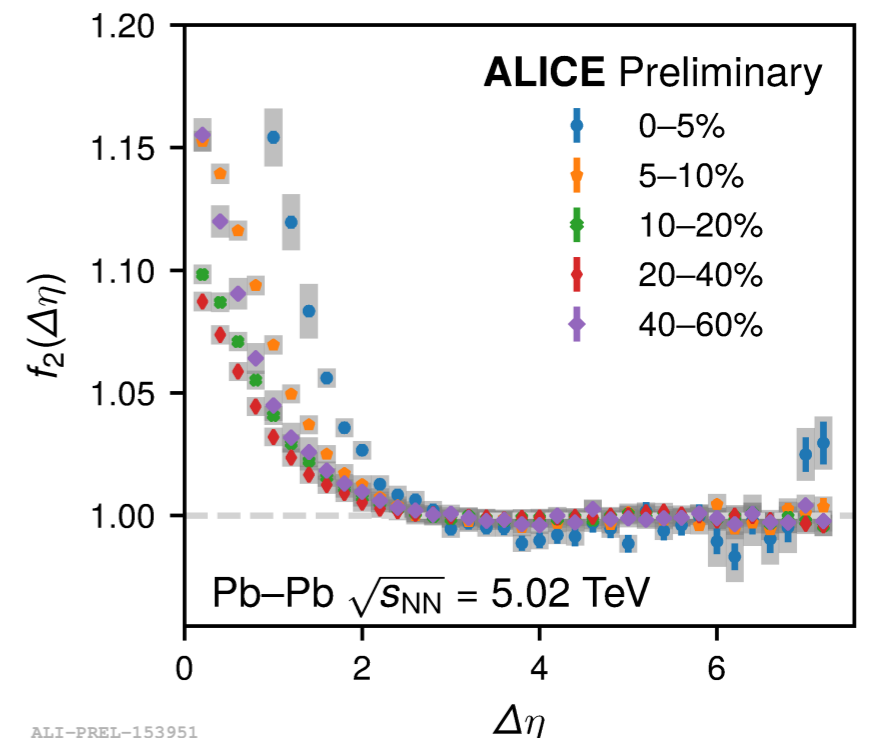
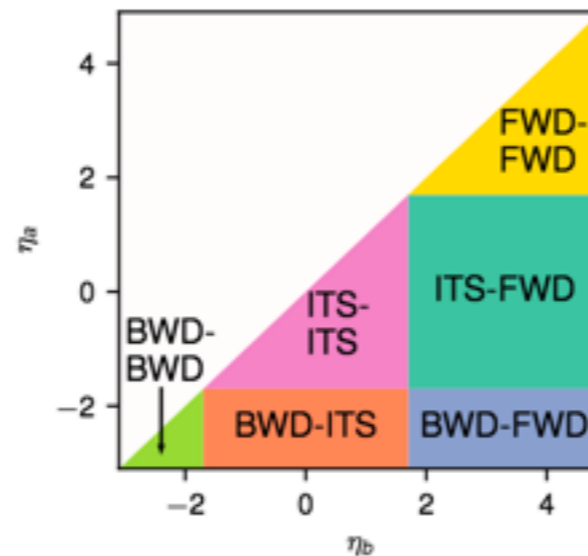
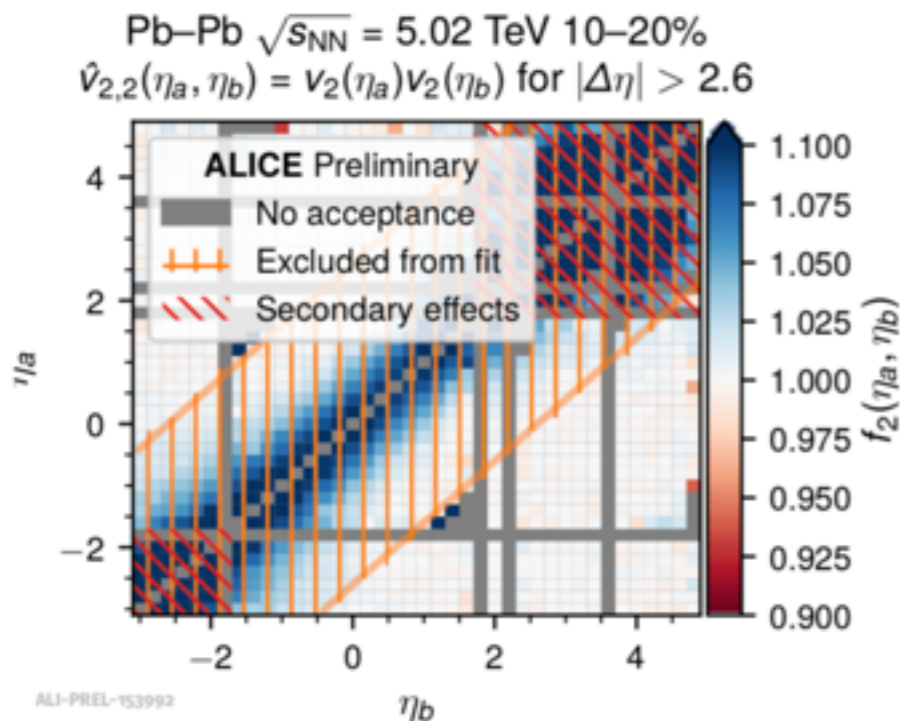
$$r_2(\eta^a, \eta^b, \phi^a, \phi^b) = \frac{\langle n_a n_b \rangle}{\langle n_a \rangle \langle n_b \rangle} \xrightarrow[\text{transform}]{\text{Fourier}} \hat{v}_{n,n}(\eta^a, \eta^b)$$

$$\hat{v}_{n,n}(\eta^a, \eta^b) = v_n(\eta^a) v_n(\eta^b) f_n(\eta^a, \eta^b)$$

$$f_n(\eta^a, \eta^b) = \frac{\hat{v}_{n,n}(\eta^a, \eta^b)}{v_2(\eta^a) v_2(\eta^b)}$$

Factorization hold $\longrightarrow f_n(\eta^a, \eta^b) = 1$

Factorization break $\longrightarrow f_n(\eta^a, \eta^b) \neq 1$



- Factorization assumption holds well for $|\Delta\eta| > 2.6$