

Summary

Tilman Plehn

Simulation

Precision

Analysis

Interpretation

# Understanding Data from First Principles

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Bad Neuenahr, September 2018

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# Theory Summary

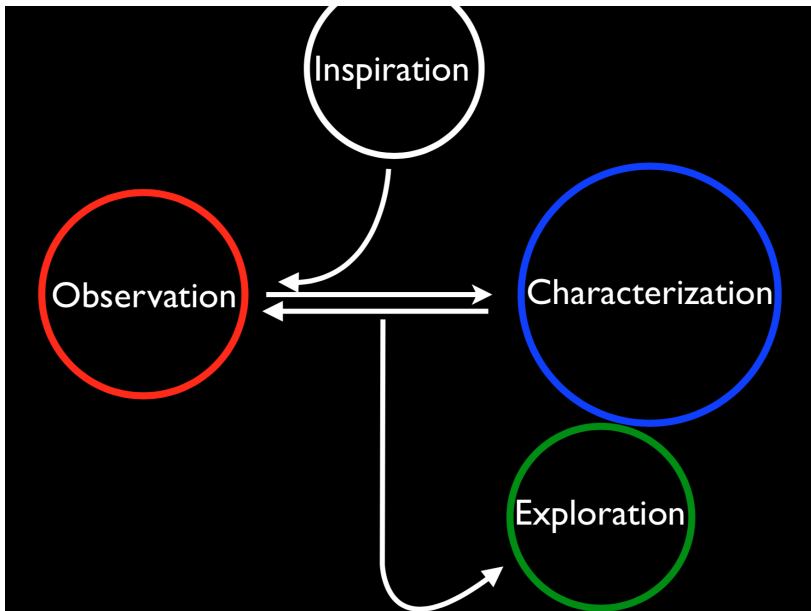
Tim M.P. Tait

University of California, Irvine



Top 2014  
October 3 2014

# How things used to look



# An era of data

## Inspiration milestones for the TeV scale

- 1964 Higgs boson [confirmed]
  - 1967 model of leptons [Weinberg, confirmed]
  - 1974 supersymmetry [Wess-Zumino]
  - 1984 composite Higgs [Kaplan, Georgi, Dimopoulos]
  - 1998 large extra dimensions [Arkani-Hamed, Dimouloulos, Dvali]
  - 1999 small extra dimensions [Randall, Sundrum]
  - 2000 little Higgs [Arkani-Hamed, Schmaltz]
- ⇒ brilliant people, but nothing beats data

## From data to renormalizable Lagrangians

- LHC data is described by QFT [fundamental physics, Standard Model is not a model]
- perturbative  $SU(3) \times SU(2) \times U(1)$  works [as does resummation]
- BSM physics exists [neutrino, matter-antimatter, coupling unification]
- bottom-up search for weakly coupled BSM physics [precision physics]
- bottom-up search for heavy BSM physics [EFT, precision physics]
- ⇒ once we see something, we build a new QFT description

# An era of data

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## Role of theory

- simulation tools [amazing progress]
- precision predictions [more amazing progress]
- analysis ideas [not so amazing progress]
- interpretation frameworks [hardly amazing progress]

# Pushing back modelling

## The way we compare data and theory [likelihood-free inference]

### 1- simulate events

hard process: perturbative field theory

QCD jets: resummed perturbative field theory

hadronization: QCD-inspired modelling

pile-up, detector,...: whatever works

### 2a compare simulated events with measured events

### 2b other interface, but the same: fiducial or unfolding,....

### 3- theory-inspired distributions etc just illustration? What's the weakest link?

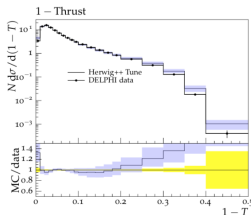
## First-principle simulations [Marek Schönherr]

– QCD whenever possible

– be clear about input

– include error bars

⇒ **keep track of modelling**



### Tune uncertainties

Tune performed by minimising

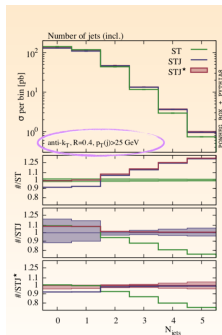
$$\chi^2(\vec{x}) = \frac{1}{N} \sum_{i \in \mathcal{O}} w_b \frac{(\text{MC}_i(\vec{x}) - \text{Data}_i)^2}{\sigma_{i,\text{Data}}^2}$$

Define **Eigentunes** as set of eigenvectors in  $\chi^2$  potential, normalised to some predefined  $\Delta\chi^2$  (same as PDF)

# Jets everywhere

## LHC physics is multi-jet physics [Rikkert Frederix]

- multi-jet simulation key to successful LHC simulations
  - $t$ -channel single top with forward jet structure relevant for  $t\bar{t}$  rejection
  - simulation of second jet important  
combination with parton shower crucial to describe data
  - cool technology: map phase spaces using a neural net
- ⇒ **deep learning the latest addition to our toolbox**



- ♦ For  $N_{\text{jets}} \geq 0,1$  bins **ST** is NLO accurate; for  $N_{\text{jets}} \geq 2$  bin the **STJ** is NLO accurate
- ♦ **STJ\*** is NLO accurate in the first three bins
- ♦ Excellent agreement among results where expected
- ♦ Due to POWHEG methodology the uncertainty bands for the higher-multiplicity bins artificially small

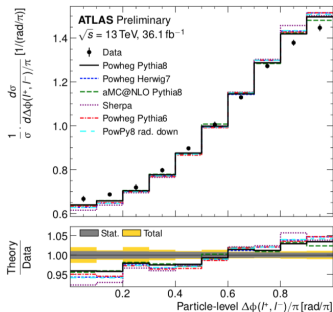
# Top pairs

## Precision predictions in top sector in great shape

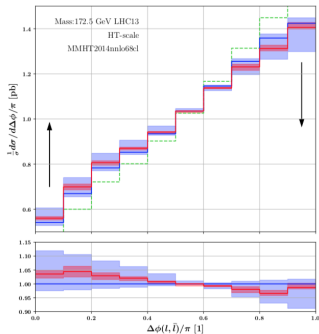
- NLO QCD also off-shell
- NNLO QCD distributions with decays [numerics limiting factor?]
- NLO e-w distributions feasible [more complex than NLO QCD]

## Kinematic distributions to NNLO [Rene Poncelet]

- closing gap between of traditional fixed-order rates and MC
- even better: **anomalous distributions**



## NWA @ NNLO predictions





# Top pairs

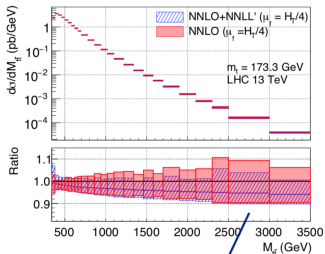
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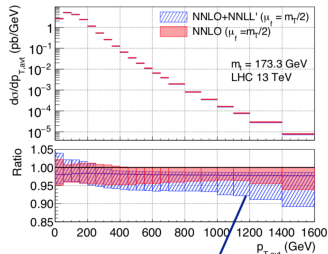
## Combining with soft resummation II II in Yannl

# NNLO+NNLL'

Czakon, Ferroglia, Heymes, Mitov, Pecjak,  
Scott, Wang, **LLY**: 1803.07623



Resummation reduces  
scale variation



Resummation  
softens the spectrum

# Top pairs with jets

## Pushing multi-leg fixed order [Stefano Pozzorini]

- $t\bar{t}b\bar{b}$  (most) interesting part of  $t\bar{t}jj$  [top-Higgs overlap]
  - phase space complex, hard process maybe  $t\bar{t}$
  - $\alpha_s^4$  leading to large  $K$ -factor and scale dependence
  - massive  $b$ -jet radiation? multi-scale process with cut-off?
- ⇒ **a tale of logarithms?**

### Natural scale choice for inclusive $\sigma_{t\bar{t}b\bar{b}}$

- for  $m_b = m_t$  the natural choice is  $\mu_R = m_t$
- natural generalisation  $\mu_R = \sqrt{m_b m_t} \Rightarrow$  **good convergence** for  $1 \leq m_t/m_b \leq 36$

$m_b$ [GeV]	4.75	15.7	52.1	172.5
$\sqrt{m_b m_t}$ [GeV]	28.7	52.1	94.8	172.5
$K(N_b \geq 0)$	1.14	1.24	1.32	1.35

⇒  $\langle \mu_{R,\text{def}} \rangle \simeq 66 \text{ GeV}$  should be reduced by factor 2–3 to match  $\sqrt{m_b m_t} = 28.7 \text{ GeV}$

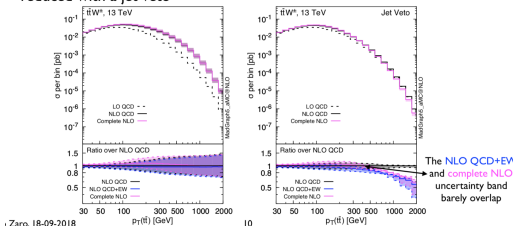
# Top pairs with stuff

## Combining strong and electroweak corrections [Marco Zaro]

- tree-level Feynman diagrams with  $\alpha\alpha_S^2$ ,  $\alpha^2\alpha_S$ ,  $\alpha^3$  [remember  $V+jets$ ]
  - QCD and ew corrections mixing classes
  - $t$ - $W$  scattering at NLO mediated by  $y_t$
  - ‘hard process’ phase-space-dependent,  $p_{Tt}$  formally safe, but especially bad
  - large  $K$ -factor controlled by jet veto
- ⇒ reproducing data with MC is dangerous, what’s our story?

## $p_T(\bar{t}t)$ and the effect of the jet veto

- QCD corrections to  $t\bar{t}W$  are dominated by real emissions recoiling against the  $t\bar{t}$  pair, with the  $W$  collinear to the emission or soft
- This leads to giant K-factors for the  $p_T(\bar{t}t)$  distribution, which are greatly reduced with a jet veto



# Top mass from first principles

## Top mass as link from Lagrangian to data [Paolo Nason, Andre Hoang]

- new physics inspiration:  
electroweak precision data  
vacuum stability  
hierarchy problem

impact from future LHC measurement?

- actually, field theory question  
related to large- $n_F$  loop diagrams [renormalons]

Tick the correct statements:

- Direct top mass measurements measure the Pole Mass.
- Direct top mass measurements measure the Monte Carlo Mass.
- Direct top mass measurements measure the Monte Carlo Mass. but you can pretend that it is the pole mass, just inflate the error a bit.
- The top is the only SM particle with more than one mass.
- You should use only leptons to avoid hadronization uncertainty.
- You should use at least NLO calculations to measure the pole mass.
- The top pole mass has renormalons, you should stay away from it.
- The MC mass differs from the pole mass by
  - terms of order  $m\alpha_s$ ;  terms of order  $\Lambda_{\text{QCD}}$ ;  terms of order  $\alpha_s\Gamma_t$ .
- The Pole Mass renormalon ambiguity is
  - $\approx 1\text{GeV}$ ;   $\approx 250\text{ MeV}$ ;   $\approx 200\text{ MeV}$ ;   $\approx 110\text{ MeV}$ .

# Top mass from first principles

## Top mass as link from Lagrangian to data [Paolo Nason, Andre Hoang]

- new physics inspiration:  
electroweak precision data  
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impact from future LHC measurement?

- actually, field theory question
- progress from cross talk: QCD vs Monte Carlo  
pole mass vs shower cut-off controllable
- ⇒ pushing back modelling, again!

$$m_t^{\text{MC}} = m_t^{\text{pole}} + \Delta_m^{\text{pert}} + \Delta_m^{\text{non-pert}} + \Delta_m^{\text{MC}}$$

QCD contribution:

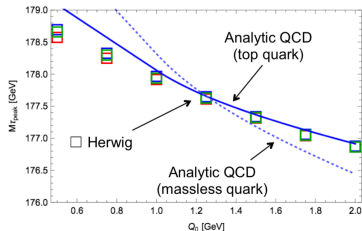
- Perturbative correction
- Depends on MC parton shower setup

Non-perturbative contribution:

- Effects of hadronization model
- May depend on parton shower setup

Monte Carlo shift:

- Contribution arising from systematic MC uncertainties
- E.g. color reconnection, b-jet modeling, finite width, ...
- Should be covered by MC uncertainty or better negligible

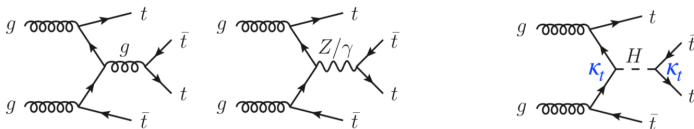


# Top Yukawa measurement

## Width-independent top Yukawa measurement [Qing-Hong Cao]

- off-shell  $t\bar{t}H, H \rightarrow t\bar{t}$
- avoid loop-induced production with its model dependence
- go for seriously off-shell

⇒ **cancellation feature of Standard Model**

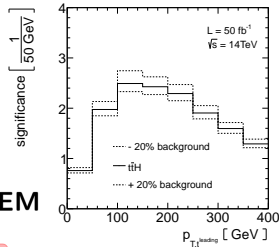


$$\sigma(t\bar{t}t\bar{t}) = \sigma^{\text{SM}}(t\bar{t}t\bar{t})_{g/Z/\gamma} + \kappa_t^2 \sigma_{\text{int}}^{\text{SM}} + \kappa_t^4 \sigma^{\text{SM}}(t\bar{t}t\bar{t})_H$$

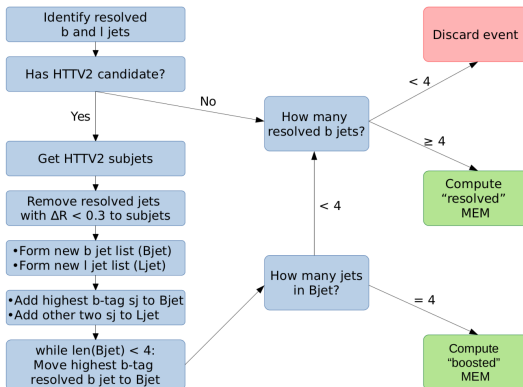
# Boosted MEM

## Combining MEM with top tagging [Maren Meinhard]

- search for  $t\bar{t}H$  combinatorics-limited [calling for MEM]
  - signal-background significance from boosted regime
- ⇒ **never stop trying new things!**



## Combining substructure methods with the MEM



# Machine learning

So what if we really do not know what we are searching for? [Tao Liu]

- yes, we have been doing ML forever, but we are not cutting edge anymore
- established ML answer: auto-encoder
- application: search for anomalous kinematics

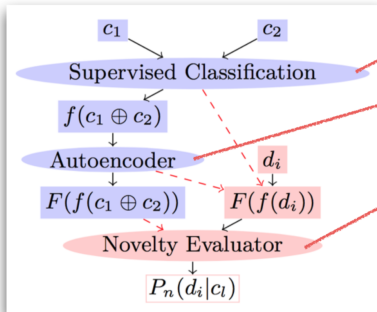
⇒ **just a diagnosis tool, don't be Chefarzt about it**



## Novelty (Anomaly) Detection - Algorithm

[J. Hajer, Y.-Y. Li, TL and H. Wang, arXiv: 1807.10261]

(For relevant studies, also see [Heimel et. al., 1808.08979; Farina et. al., 1808.08992])



- Step 1: (SM/background) feature learning
- Step 2: dimension reducing of feature space (auto-encoder)
- Step 3: novelty evaluating of testing data
- Analyze detection sensitivity based on novelty response of testing data

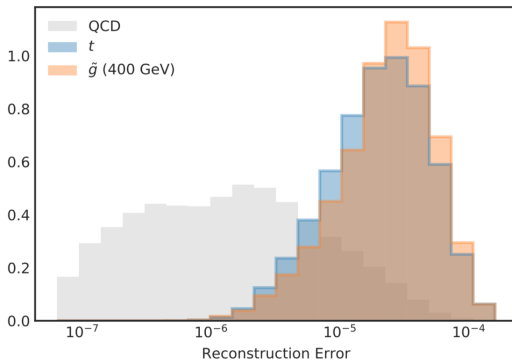


# More machine learning

If events are too hard, use jets [David Shih; coordinated with Heimele, Kasieczka, TP, Thompson]

- lots of training data in all phase space regions
  - check for jets which do not look like QCD
  - uncertainties the next big challenge
  - de-correlate/control the jet mass for control regions
- ⇒ training on data, searching in data means no modelling, yet again

Can use reconstruction error as an anomaly threshold.



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# Sub-jet physics

# Top EFT

## Bottom-up analysis in QFT framework: effective theory

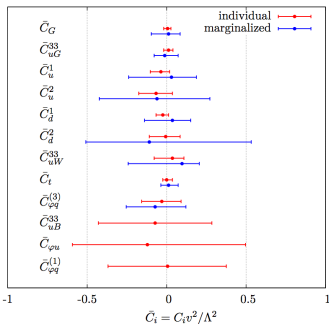
- 59 operators, assuming good-taste symmetries
- affecting rates (coupling modifiers) and kinematics (Lorentz structures)
- best framework for global analysis, leading to SMEFT
- renormalization no problem, at least in top sector
- for details, check endless Higgs-related discussions

⇒ where are the ATLAS/CMS results??

### Constraining top quark effective theory in the LHC Run II era

The TOPFITTER Collaboration

Andy Buckley, Christoph Englert, James Ferrando, David J. Miller,  
Liam Moore, Michael Russell, and Chris D. White



# Top-bottom link

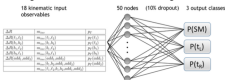
## Understanding EFT analyses [Seth Moortgat]

- upgrading  $t\bar{t}b\bar{b}$  from background to signal
  - categorize EFT effects in terms of physics
  - flavor-EFT fit established and very constraining
  - $SU(2)$  doublet means  $t_L$  talks to bottom sector
- ⇒ **combining machine learning and EFT!**

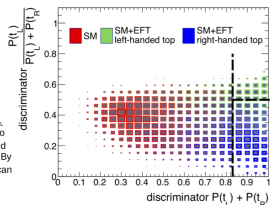
### 4. Learning the effective operators

**Case study:** can a NN learn to distinguish between operators with left-handed top-quark currents ( $t_L$ ) and right-handed top-quark currents ( $t_R$ )? If so, can we use this to improve limits on the Wilson coefficients?

A shallow neural network was constructed that combines 18 kinematical variables to predict one of **three output classes**. From the network outputs, two discriminators are built; one to distinguish between SM and EFT in general, and another one to distinguish  $t_L$  from  $t_R$ .



**Figure (right):** the x-axis represents the SM vs EFT discriminator, whereas the y-axis represents the  $t_L$  vs  $t_R$  discriminator. SM contributions are located to the left, whereas the EFT contributions are located to the upper right ( $t_L$ ) and lower right ( $t_R$ ) corners. By combining limits in dedicated signal regions, we can improve the limits/confidence intervals!



Scenario 1: observation of the SM

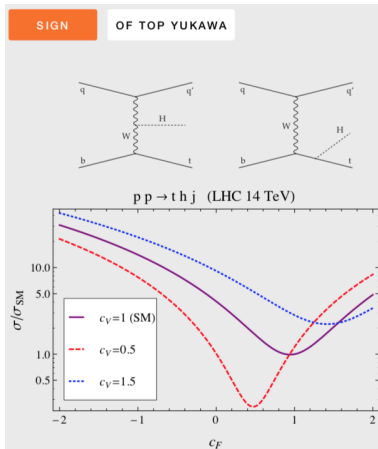
Scenario 2: observation of an EFT signal

# Top-Higgs link

## Single-top plus Higgs production [Roberto Franceschini]

- tough analysis, like single top, but with less rate
- correlation  $t\bar{t}H$  and  $WWH$  couplings [orthogonal to  $H \rightarrow \gamma\gamma$ ]

⇒ in case you find single top too easy



# Best of EFT worlds

## CLIC: precision and energy [Francesco Riva]

- EFT sensitive to  $v^2/\Lambda^2$  in rate [LEP1]
  - same EFT sensitive to  $p^2/\Lambda^2$  in tails [LHC]
  - learn from LHC for the future:  $e^+e^- \rightarrow t\bar{t}$  at high energies
- ⇒ in our era of data it's too far away for me...
- ⇒ ...but: shows that EFT analyses only need first and last bins

$$\sim c_{(\Psi)^4} \frac{E^2}{M^2}$$

- ▶ maximal E-growth
- ▶ interference

What will we learn?

CLIC:  $\sigma \sim 20 \text{ fb} @ 3 \text{ TeV}$      $2 \text{ ab}^{-1}$      $\longrightarrow \delta\sigma/\sigma|_{stat} \sim 1\%$

▶  $M \sim 30 \frac{\sqrt{c_{(\Psi)^4}}}{g} \text{ TeV}$

# Even more data in the future

## Data-driven analyses with full theory

- what limits our first-principle simulations?
- which precision calculations do we need next?
- which analysis ideas get us closer to data?
- how do we interpret searches without models?
- ATLAS/CMS: who is the audience for your papers?
- how can someone understand modern analyses?

⇒ we need to work on our story; sorry, Tim, but BSM inspiration is broken

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## All 'field in crisis' talk is bullshit, but HL-LHC has a problem [Jan Kieseler]

- 1- particle physics has plenty of data
- 2- our standard theory works, if anything, too well
- 3- at least I have a big goal: dark matter
- 4- dead new physics models are good news
  - **but we should** fight for the best students
    - let young people with new ideas run
    - become inventive again ourselves
    - avoid death by boredom