

Wormholes and Entanglement

Juan Maldacena

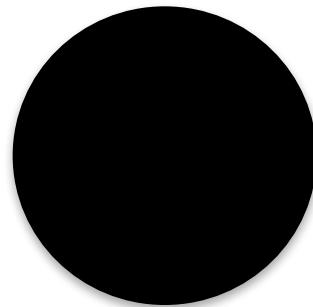
Institute for Advanced Study

Outline

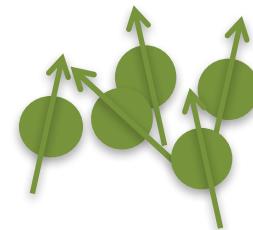
- Black holes as quantum systems
- Schwarzschild wormhole and entangled states.
- Traversable wormholes.
- Simple dynamics in Nearly- AdS_2
- Traversable wormhole solution in 4d.

Black holes as quantum systems

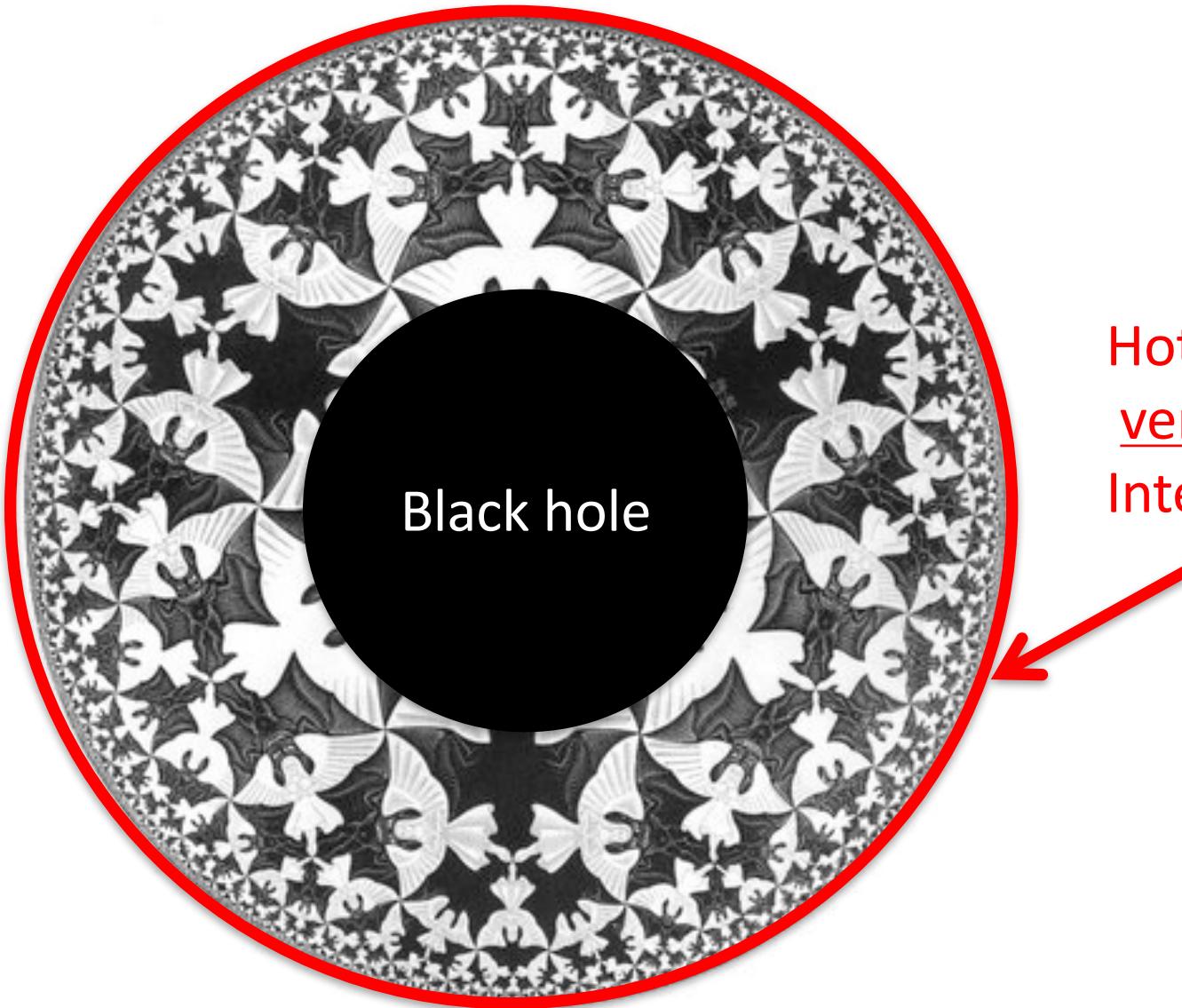
- A black hole seen from the outside can be described as a quantum system with S degrees of freedom (qubits).



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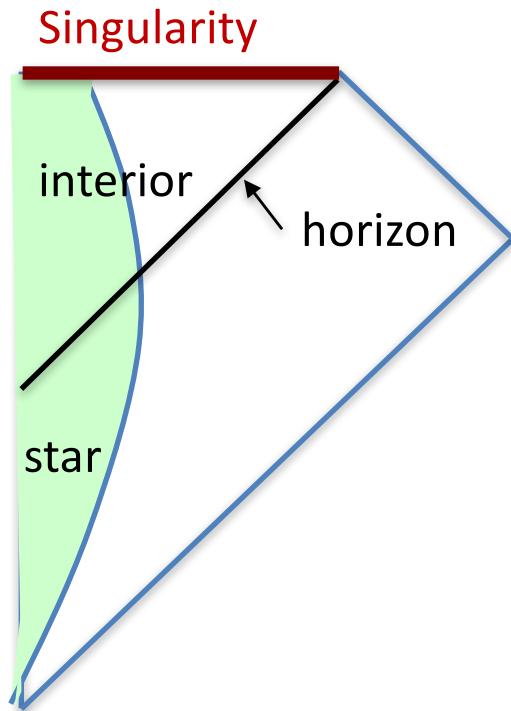


one reason is AdS/CFT...



Hot fluid made out of
very strongly
Interacting particles.

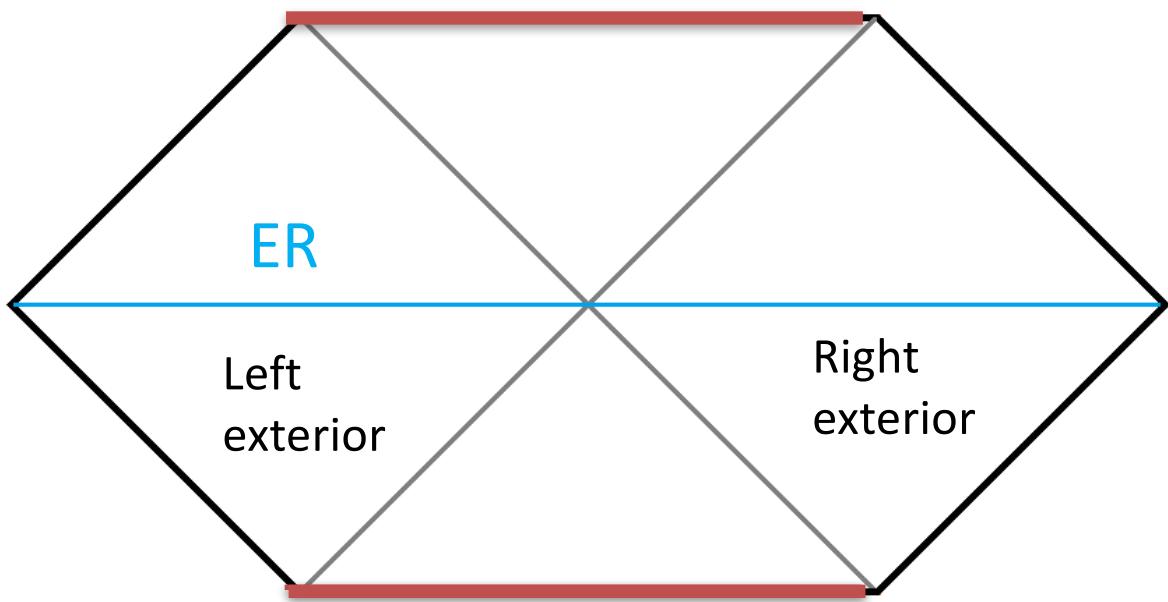
Geometry of a Black Hole made from collapse



Oppenheimer Snyder 1939

One exterior, one interior.

Full Schwarzschild solution

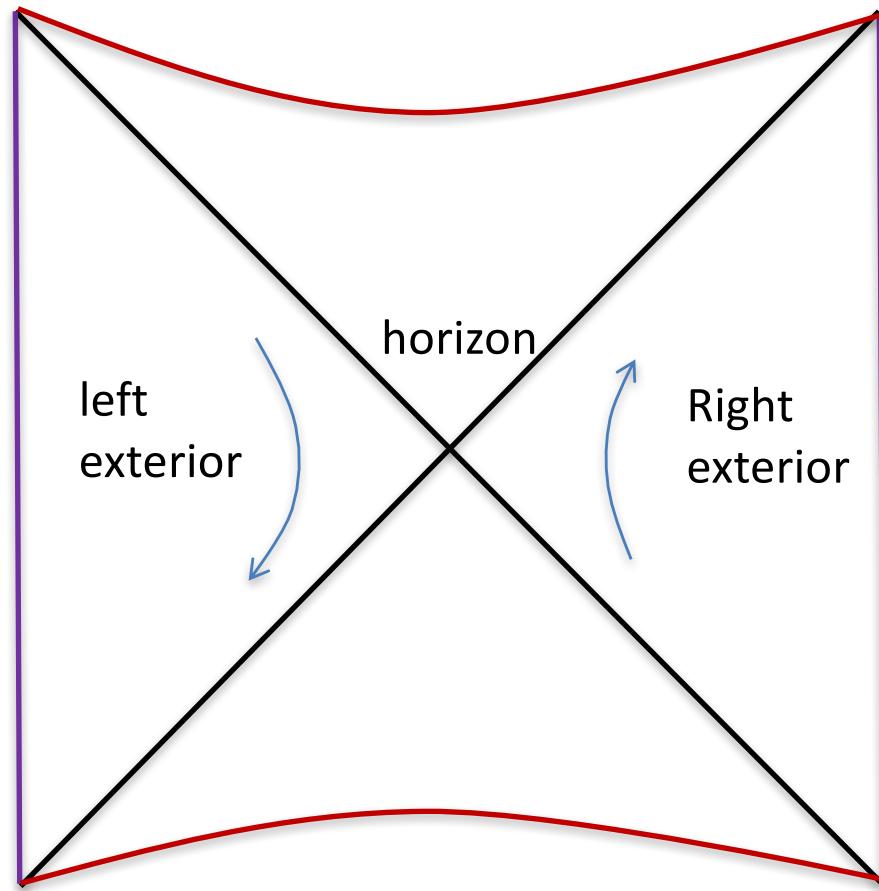


Eddington, Lemaitre, Einstein,
Rosen, Finkelstein,
Kruskal

Vacuum solution.
Two exteriors, sharing the interior.

The full Schwarzschild wormhole

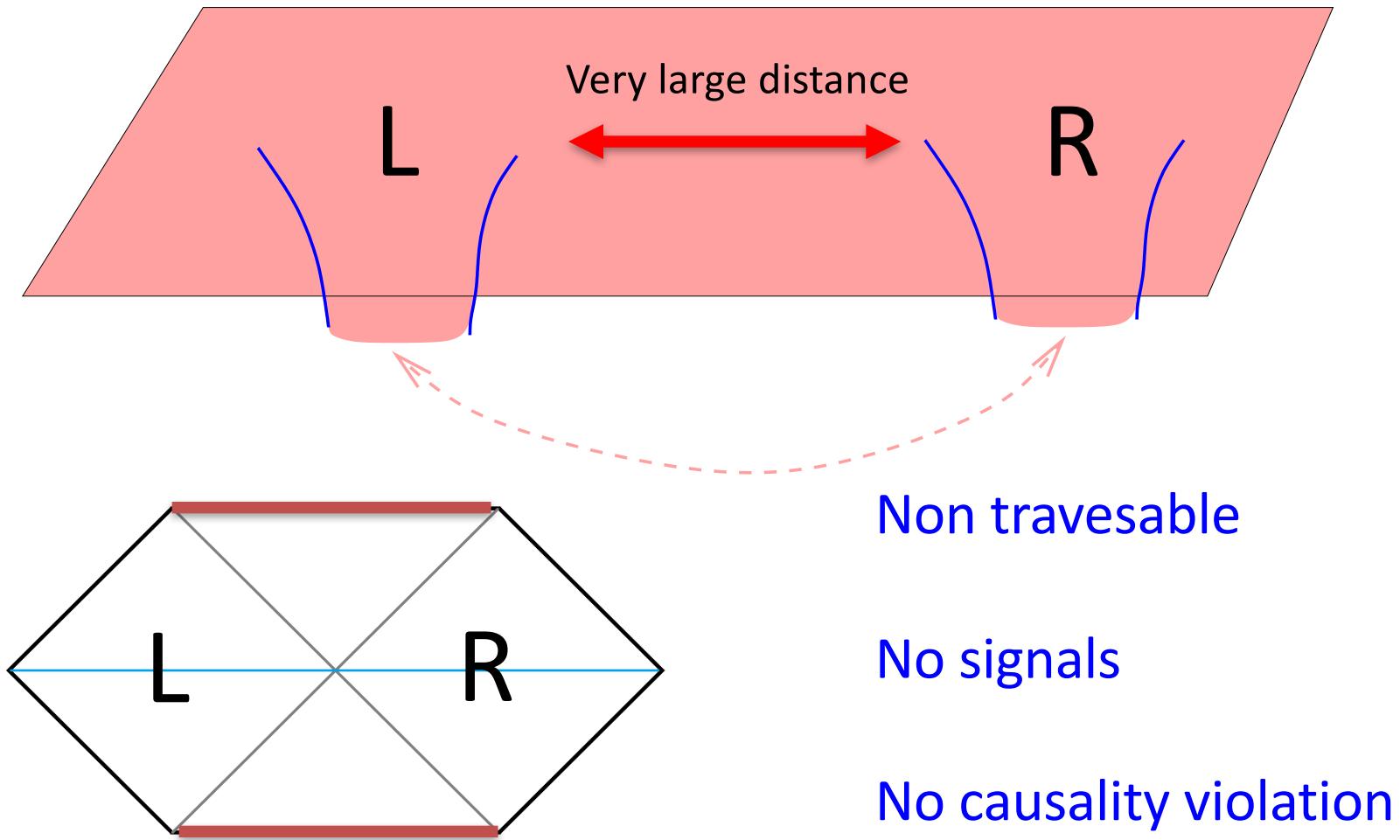
AdS asymptotics.



No need to postulate any exotic matter

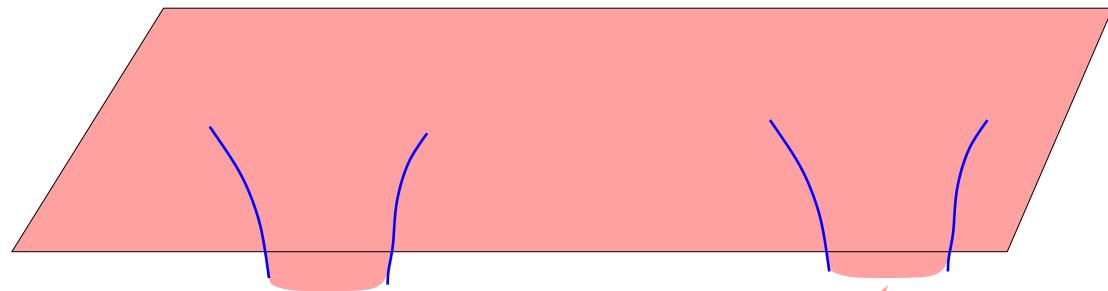
No matter at all !

Wormhole interpretation.



Fuller, Wheeler, Friedman, Schleich, Witt, Galloway, Wooglar

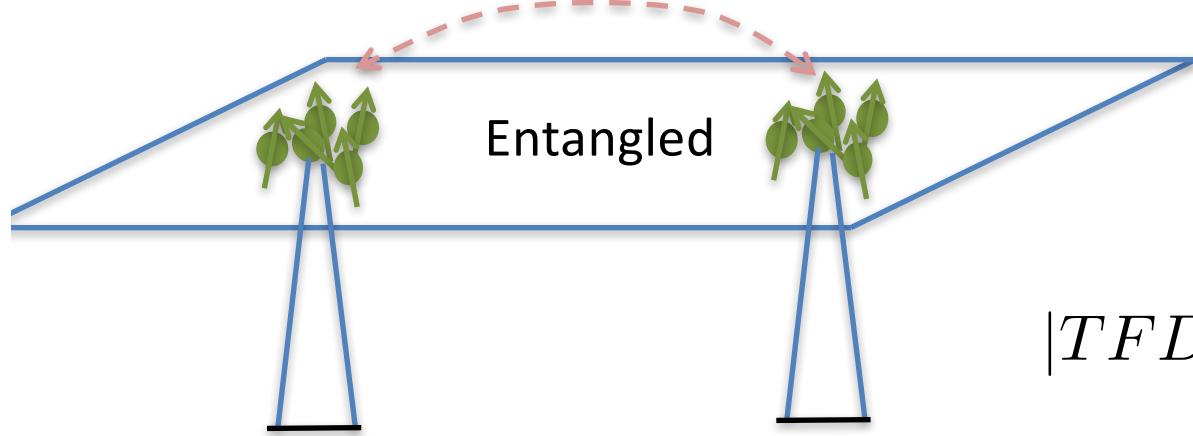
Wormhole and entangled states



Connected through the interior

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W. Israel
J.M.



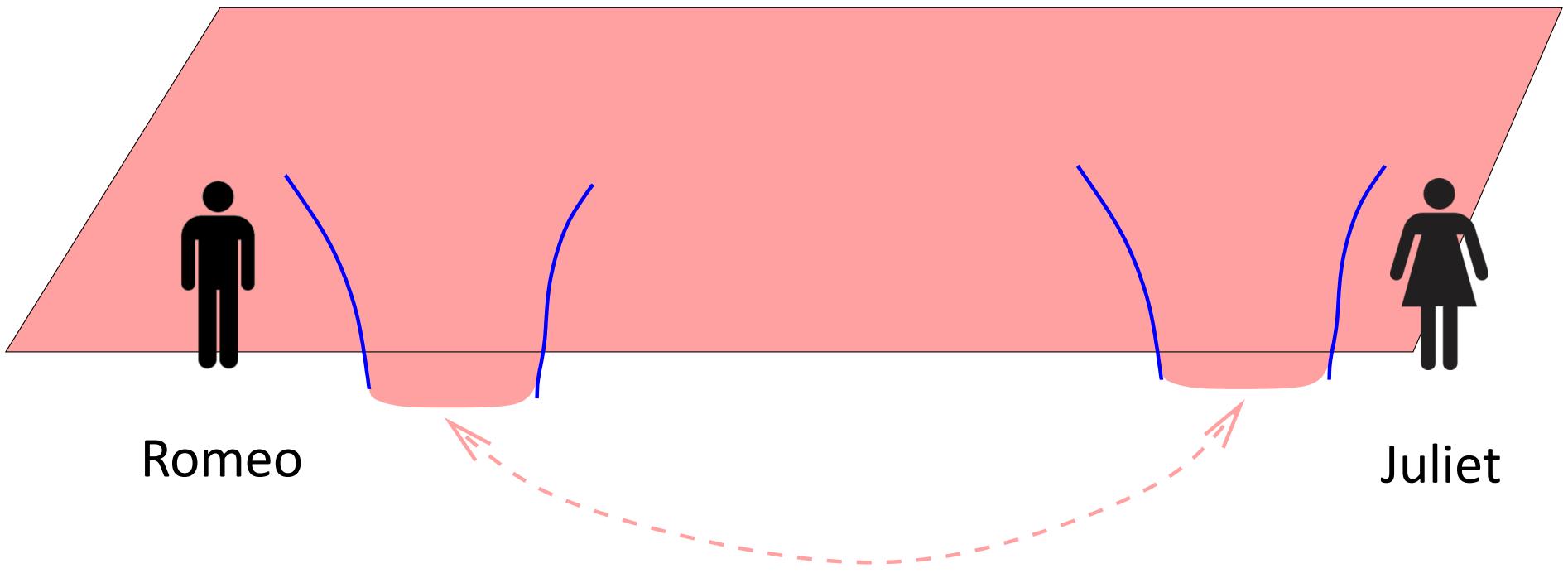
$$|TFD\rangle = \sum_n e^{-\beta E_n/2} |\bar{E}_n\rangle_L |E_n\rangle_R$$

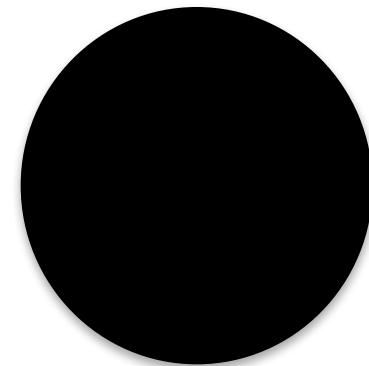
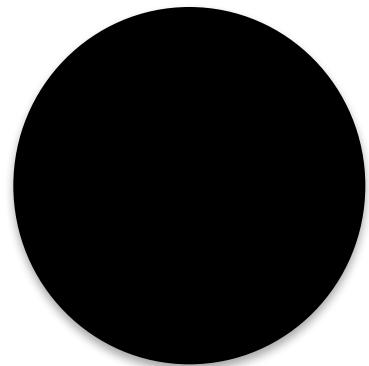
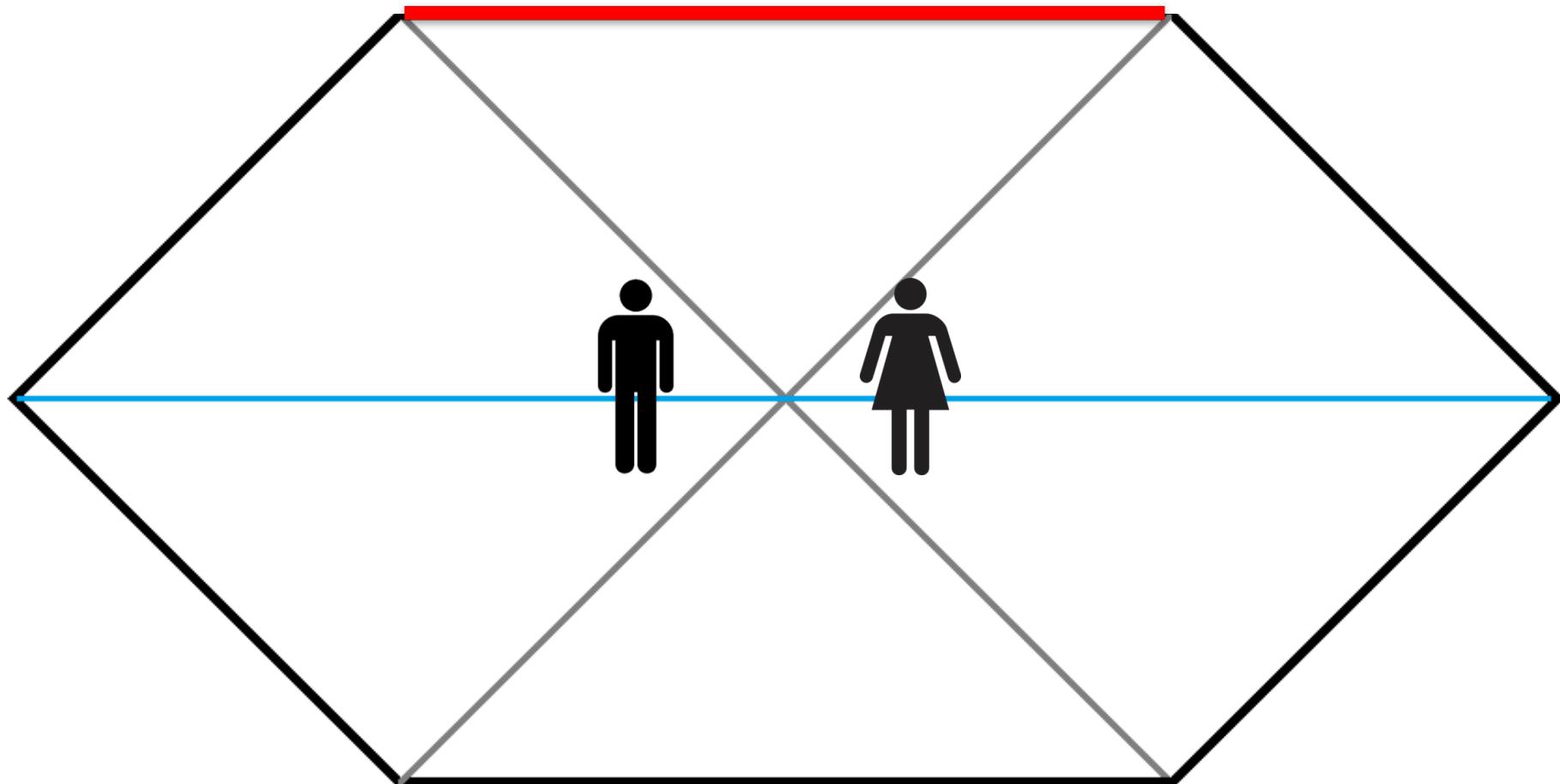
(in a particular entangled state)

ER = EPR

J.M. , Susskind

A forbidden meeting





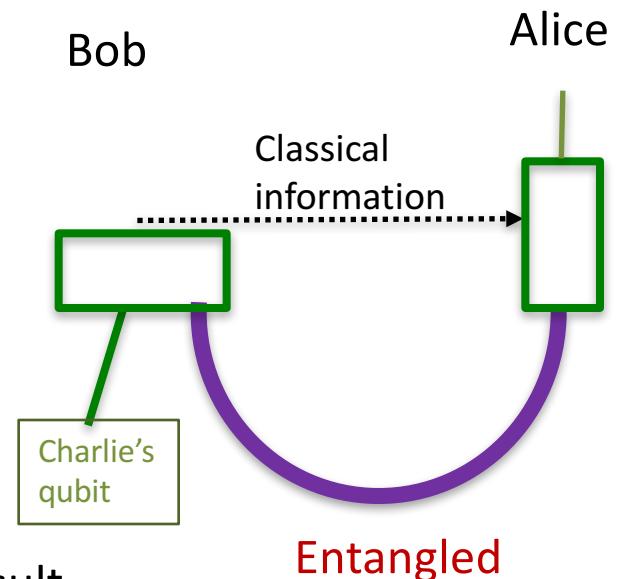
Is there a better way ?

- Traversable wormholes
- Quantum teleportation

Gao, Jafferis, Wall,
...
...

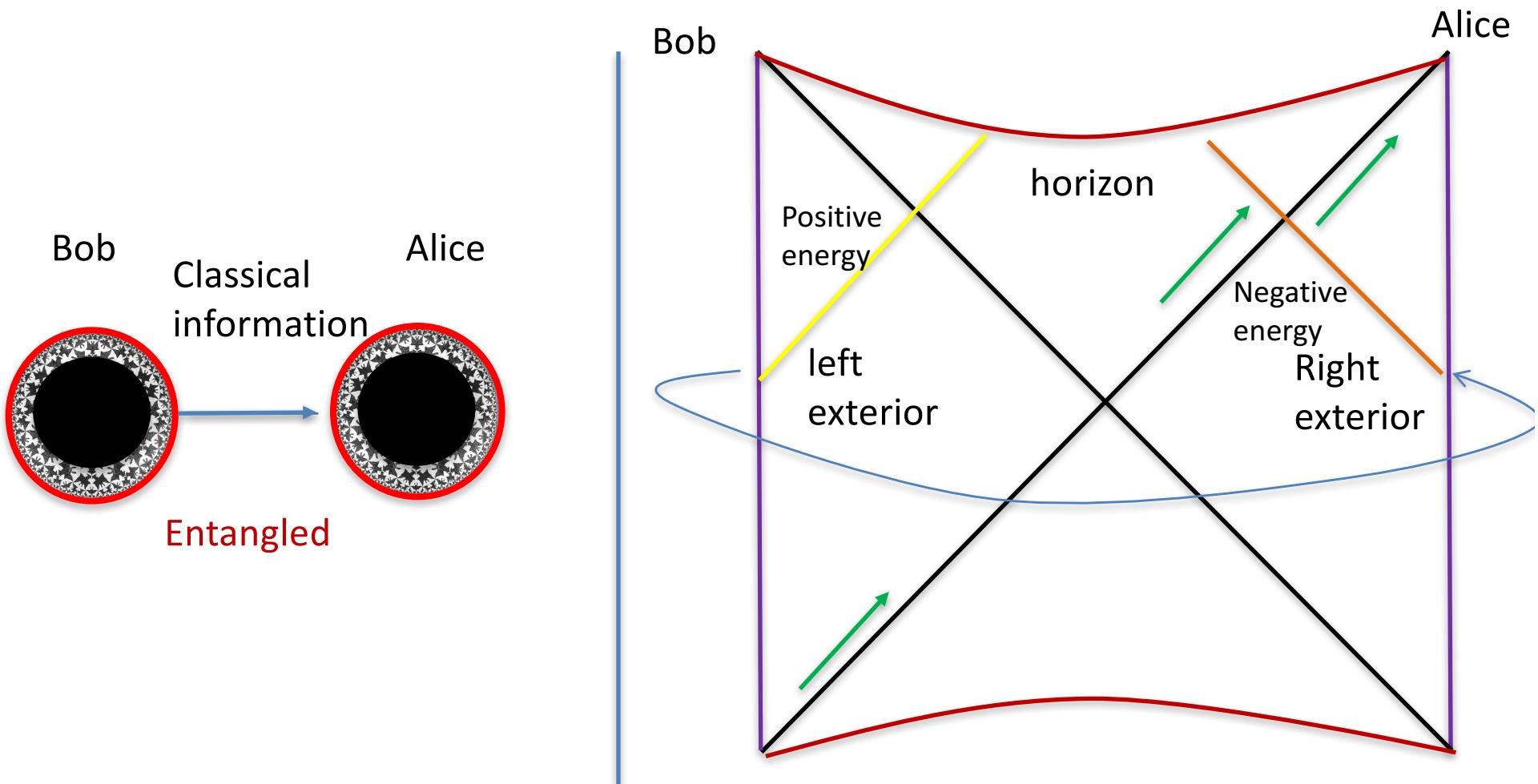
Quantum teleportation

- Bob and Alice share an entangled pair of qubits.
- Charlie gives Bob a qubit and he wants to send it to Alice.
- Bob does a joint measurement of Charlie's qubit and his qubit.
- Sends the the result to Alice as classical information
- Alice does an operation on his qubit that depends on Bob's result.
- Alice gets the qubit.
- Resources needed to send a qubit: One entangled qubit and 2 bits of classical information.



How does the qubit travel ? Would you like to be teleported ?

Teleportation through the wormhole

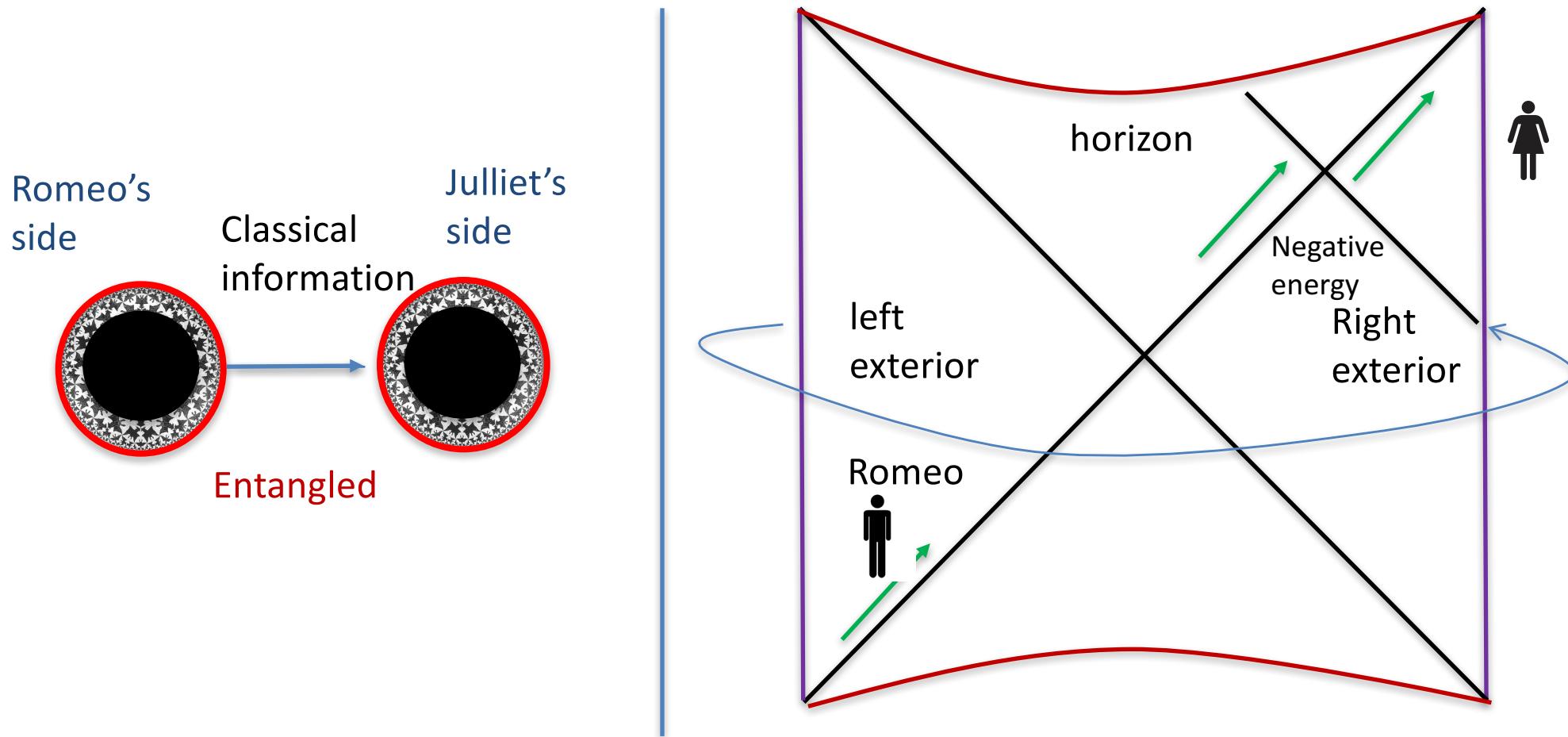


Gao, Jafferis, Wall

Quantum teleportation through the wormhole

- What does the message feel ?
- A shock wave.
- Can be arranged to be painless if it is spherically symmetric.

Teleportation through the wormhole



Gao, Jafferis, Wall

- Romeo can report he saw a smooth interior!

Conclusions

- Entangled black holes can display a geometric connection.
- Entanglement is related to wormholes.
ER=EPR.
- Teleportation is related to travel through wormholes.

Challenges

- We do not really know how to “derive” the spacetime corresponding to a given state.
- What ‘‘happens’’ at a black hole singularity or Big Bang singularity ?

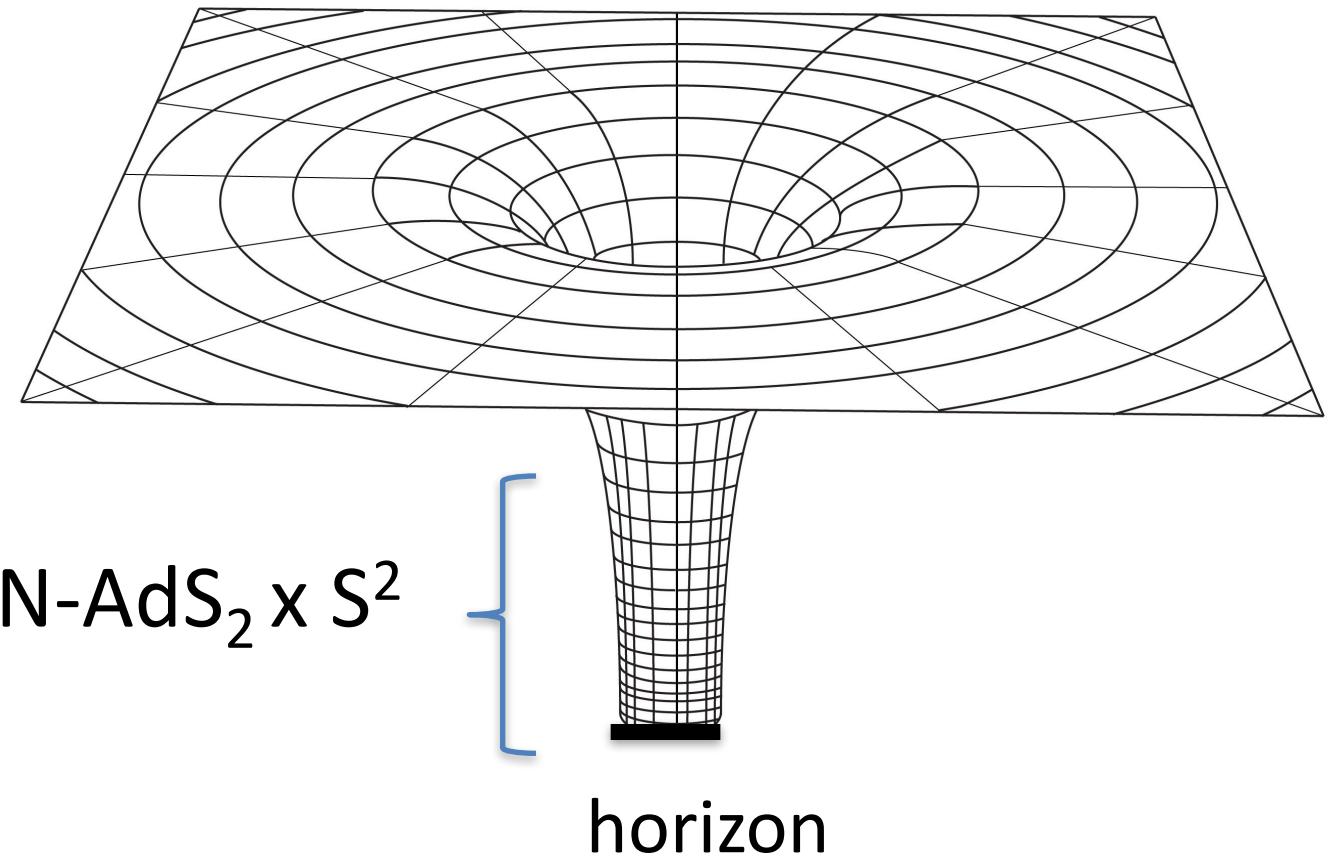
A simple case

Near extremal black holes
and
Nearly- AdS_2

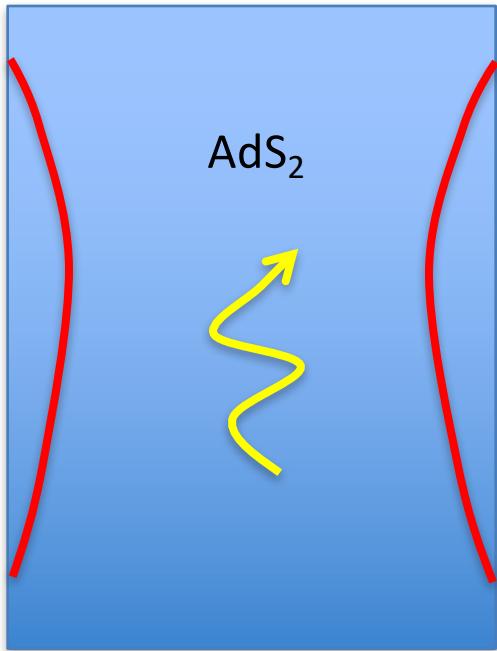
Near extremal black holes

$M \geq Q$

$M \sim Q$



The surprisingly simple gravitational dynamics of N-AdS₂



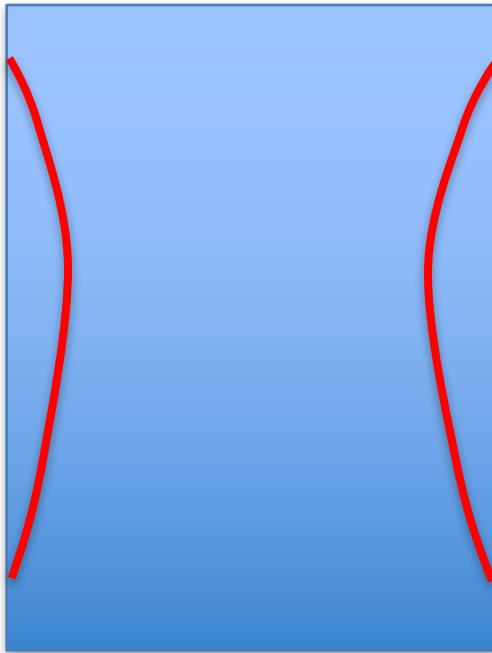
NAdS₂ = AdS₂ + location of **boundary**

Dynamics of the boundary is SL(2) invariant.

Proper time along the boundary = time of the asymptotically flat region = time of the quantum system

$$(H_{L\,Bdy} \times H_{\text{bulk}} \times H_{R\,Bdy})/SL(2, R)$$

Temperature



Distance between boundaries → temperature

Redshift between boundary and center is of order $\beta = 1/T$

Smaller temperature → larger redshift factor → longer wormhole

Traversability

Add a constant interaction between the dual quantum systems

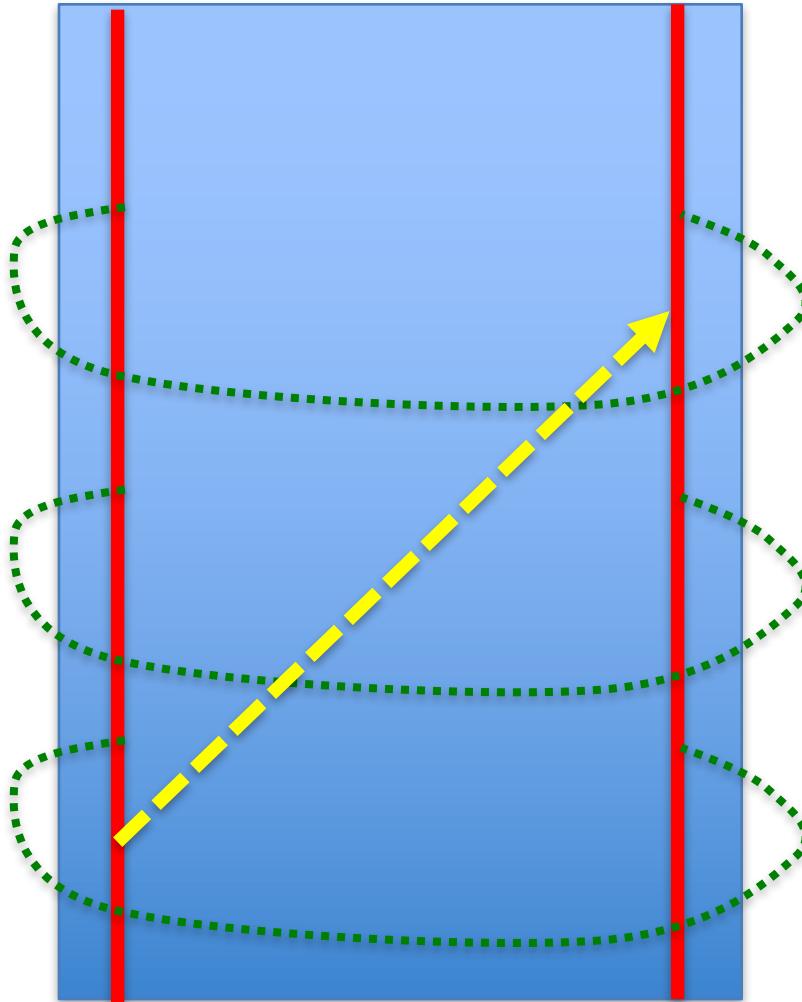
- Nearly- AdS_2 gravity
- Plus matter
- Plus boundary conditions connecting the two sides

$$S_{int} = \mu \int du \chi_L(u) \chi_R(u)$$

u is proper length along the boundary, or boundary time.

- This generates negative null energy and allows for an eternally traversable wormhole

NAdS₂ gravity + Interaction



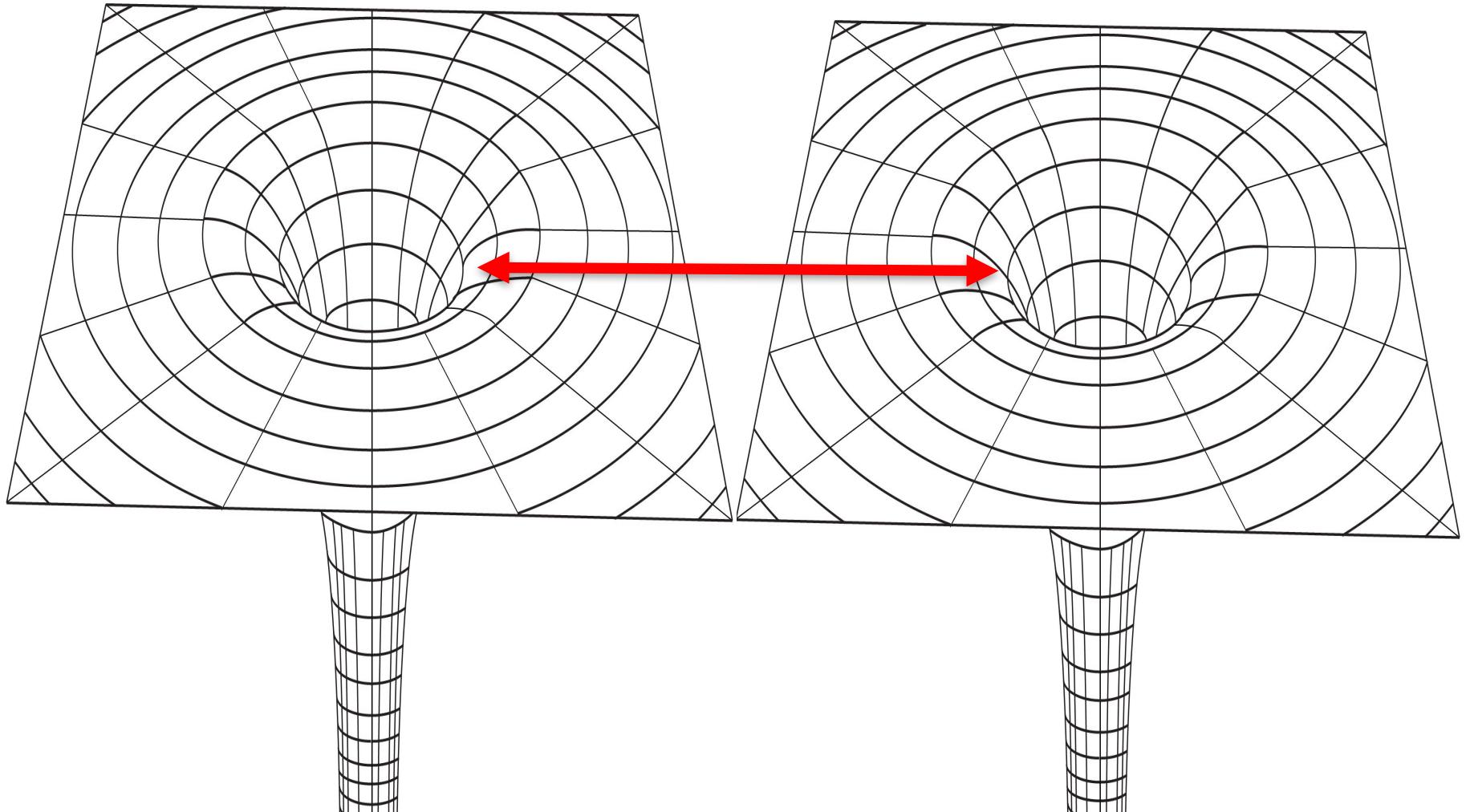
Interactions

Boundaries now move “straight up”

Signals can now propagate from one boundary to the other.

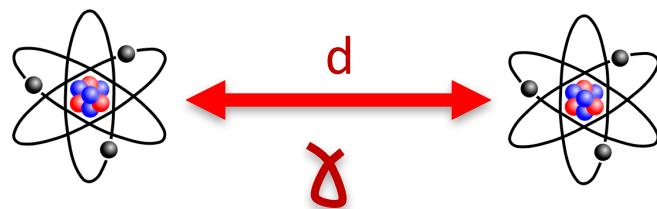
- So far we have introduced the interaction by hand.
- Can we understand how it arises?
- By bringing the black holes near each other!

Exchange of bulk fields can lead to the interaction between the quantum systems the describe the black hole



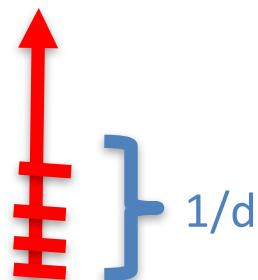
Analogy: Van der Waals interaction

Two neutral atoms exchanging photons.



$$H_{int} \propto \frac{\vec{d}_L \cdot \vec{d}_R}{d^3}$$

d small enough so that $1/d$ is larger than the gap between the ground state and the next states.



Entangle the two atoms.



This reasoning inspired the following
solution

Wormholes in 4 dimensions, in the Standard Model + gravity

Based on work with:

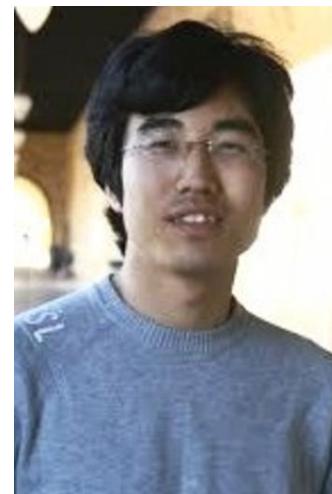


Alexey Milekhin

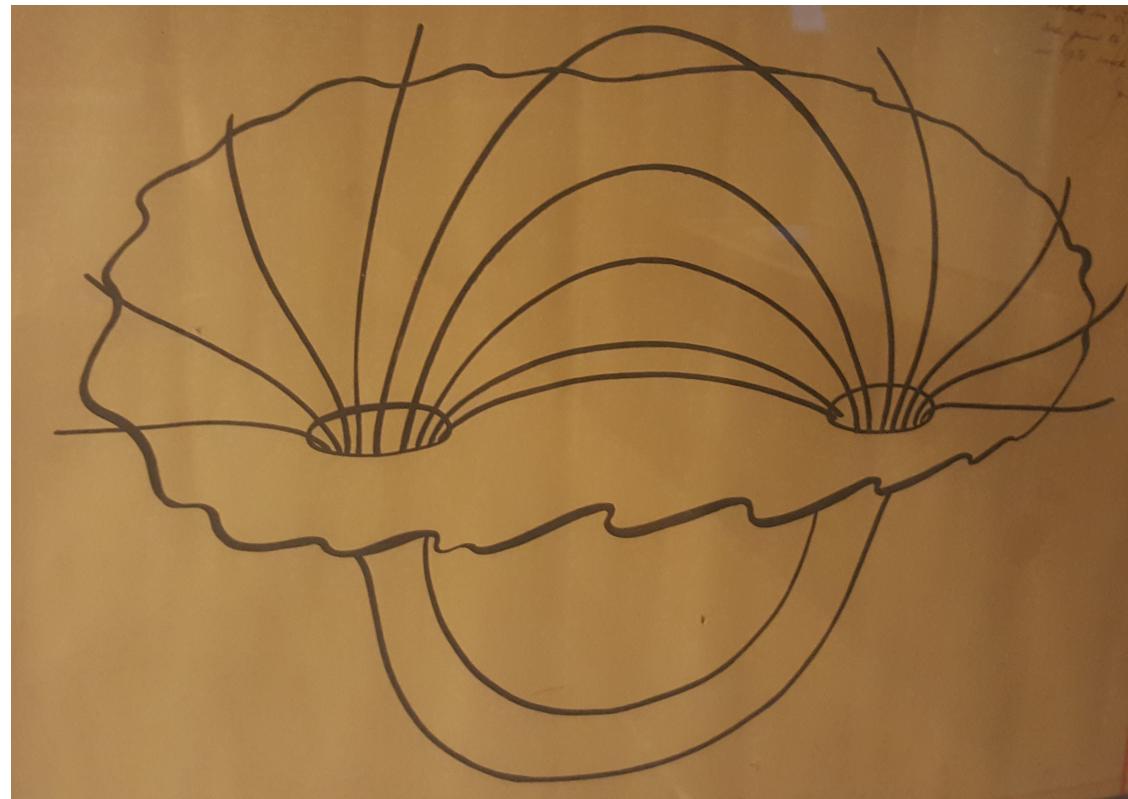


Fedor Popov

Related to previous work with Xiaoliang Qi



Drawing by John Wheeler, 1966



Charge without charge.
Mass without mass

Spatial geometry. Traversable wormhole

Recall classic results

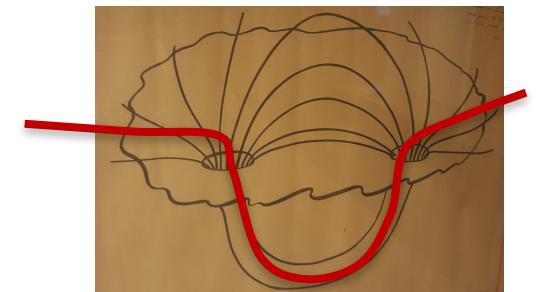
There are no science fiction wormholes!

- No wormhole allows you to travel faster than the speed of light in the ambient space.
- Forbidden by:
 - I) Einstein equations.
 - II) The Achronal Average Null Energy Condition

Friedman Schleich, Witt, Galloway, Woolgar
Gao Wald

Not yet proven in a general spacetime, but believed to hold in QFT

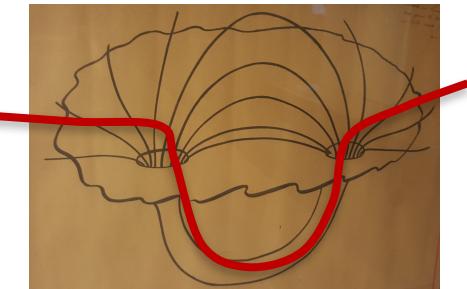
$$\int dx^- T_{--} \geq 0$$



Achronal = points along null line are not timelike separated = fastest line

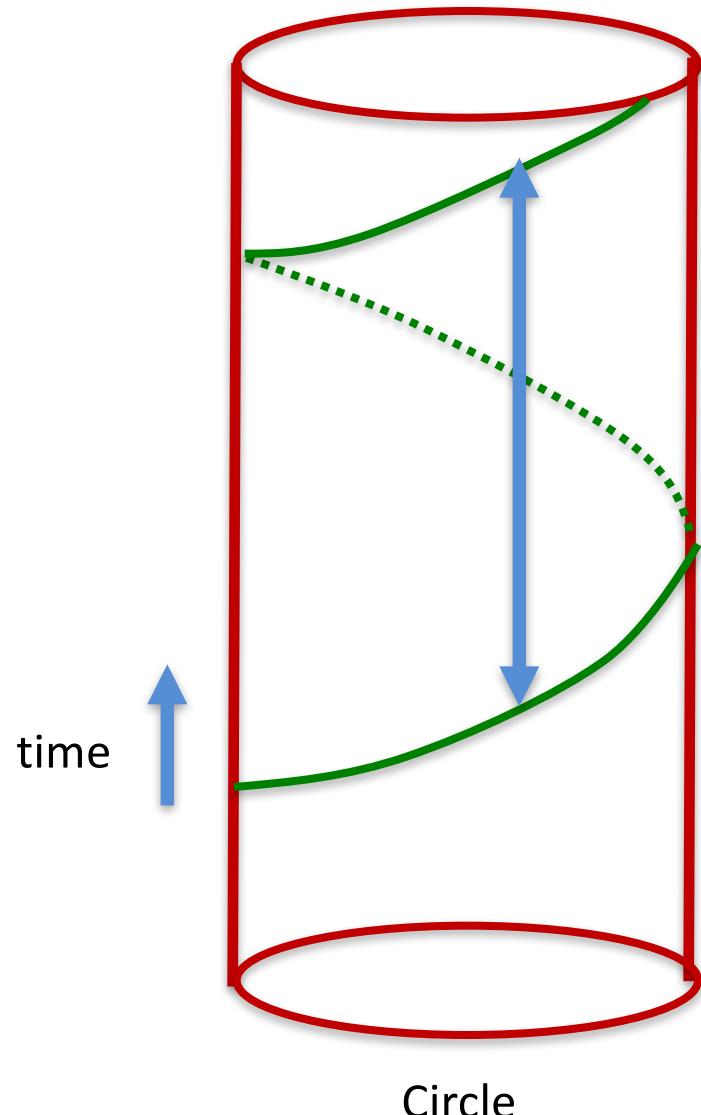
Longer wormholes?

- What if it takes longer to go through the wormhole ?
- Not possible in classical physics due to the Null Energy Condition.
Topological censorship: Friedman Schleich, Witt, Galloway, Woolgar
- In classical physics → we can go to the covering space, which is forbidden by the previous case.
- → We need quantum effects to find a solution.
Casimir-like energy.
- Can we do it in a controllable way ?



Negative null energy in QFT

Eg. Two spacetime dimensions



$$T_{--} < 0$$

$$E \propto -\frac{c}{L}$$

Negative Casimir energy

Quantum effect

The null energy condition does not hold for null lines that are not achronal!

Some necessary elements

- We need something looking like a circle to have negative Casimir energy.
- Large number of bulk fields to enhance the size of quantum effects.
- We will show how to assemble these elements in a few steps.

The theory

$$S = \int d^4x [R - F^2 + i\bar{\psi} \not{D} \psi]$$

Einstein + U(1) gauge field + massless charged fermion

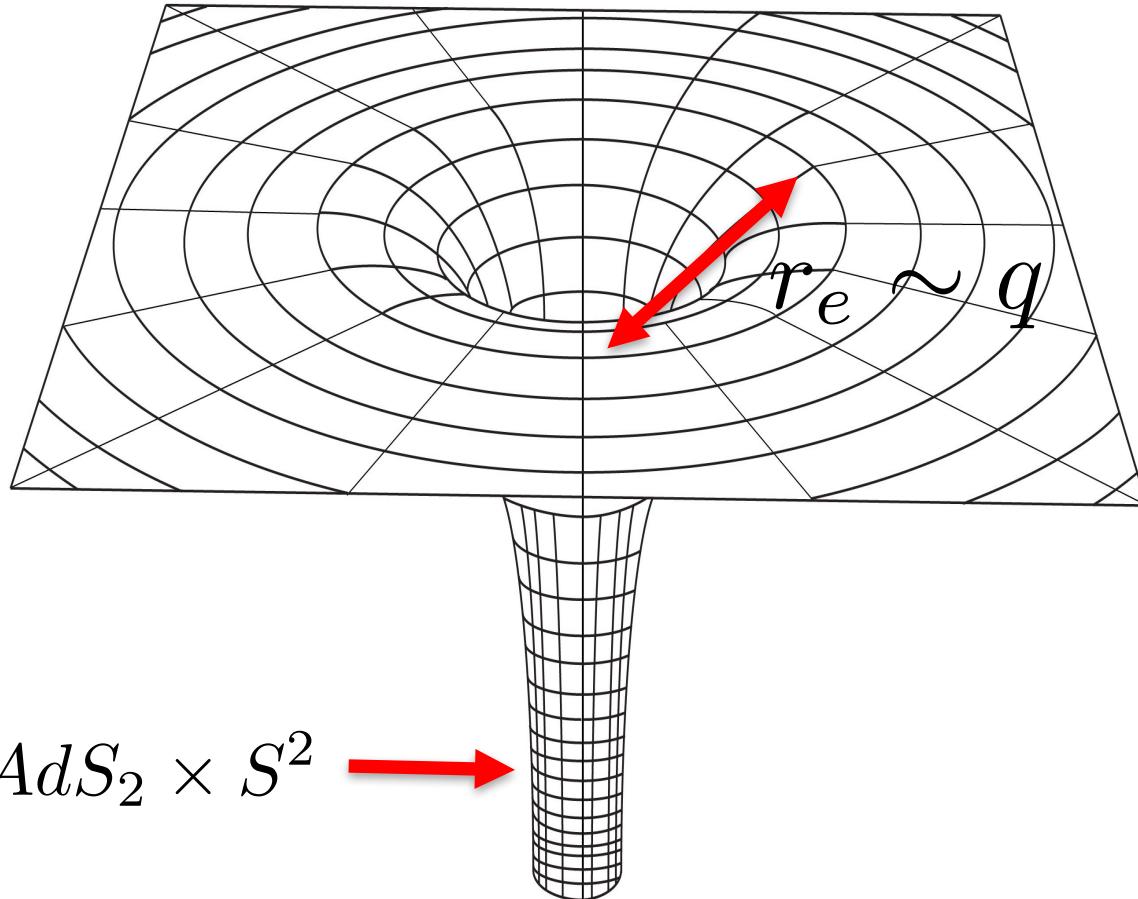
Could be the Standard Model at very small distances, with the fermions effectively massless. The U(1) is the hypercharge. SU(3) x SU(2) x U(1).

$$l_{\text{Planck}} = 1$$

The first solution: Extremal black hole

Magnetic charge q

$$\int_{S^2} F = q = \text{integer}$$

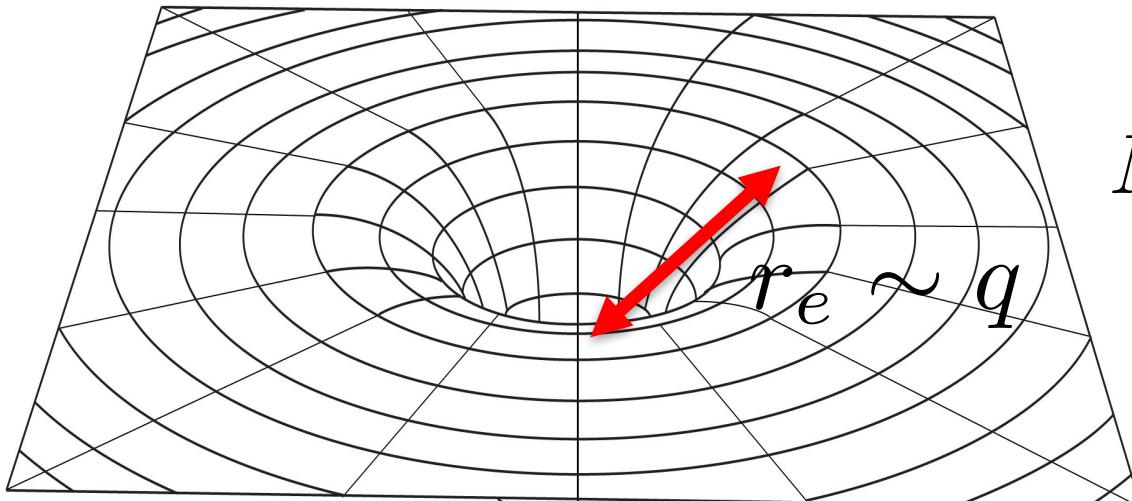


$$M = q$$

$$AdS_2 \times S^2 \rightarrow$$

$$l_{\text{Planck}} = 1$$

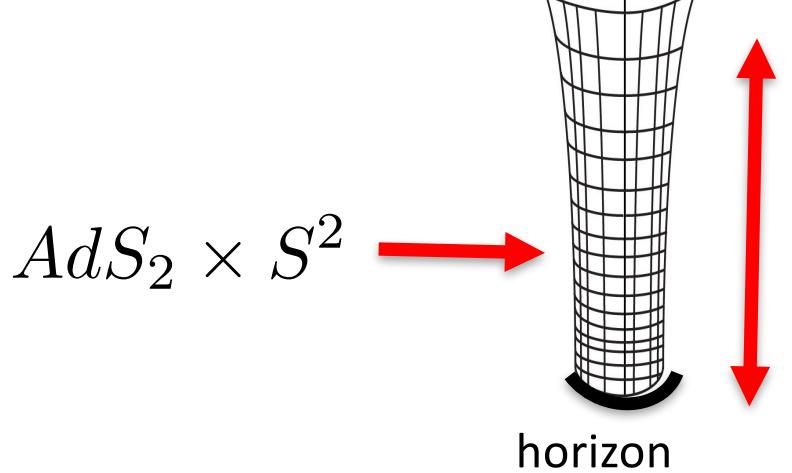
The next solution: Near Extremal black hole



$$M - q = q^3 T^2 = \frac{q^3}{\beta^2}$$

↳

Very small



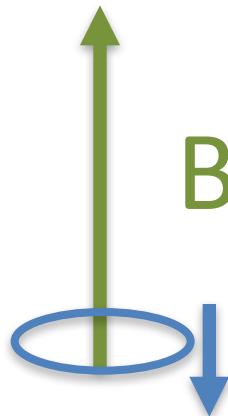
β is the “length” of the throat. Redshift factor between the top and the bottom

$$l_{\text{Planck}} = 1$$

Motion of charged fermions

- Magnetic field on the sphere.
- There is a Landau level with precisely zero energy.
- Orbital and magnetic dipole energies precisely cancel.
- Explained by an anomaly argument

Ambjorn, Olesen



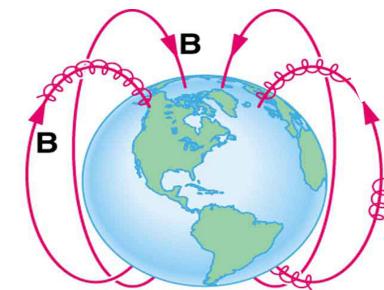
Massless fermions \rightarrow U(1) chiral symmetry

4d anomaly \rightarrow 2d anomaly \rightarrow there should be massless fermions in 2d.

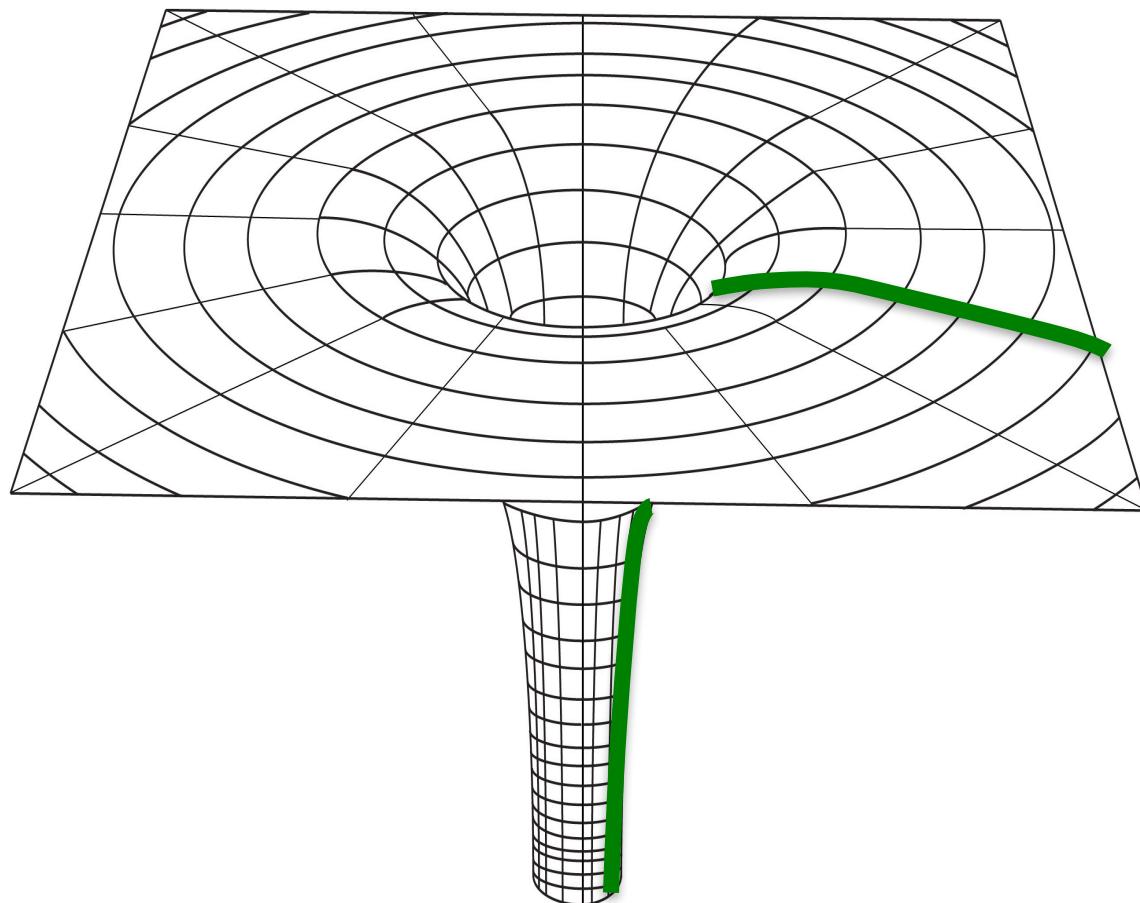
(Here we view F as non-dynamical).

Motion of charged fermions

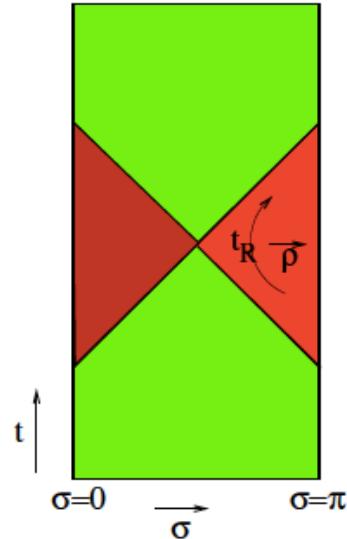
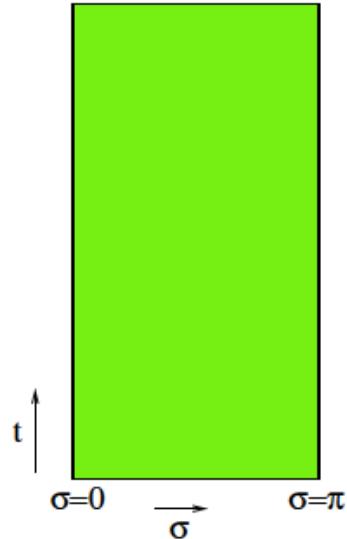
- Degeneracy = $q = \text{flux of the magnetic field on the sphere. Form a spin } j, \text{ representation of } \text{SU}(2), 2j + 1 = q.$
- We effectively get q massless two dimensional fermions along the time and radial direction.
- We can think of each of them as following a magnetic field line.



q massless two dimensional fields, along field lines.



AdS₂



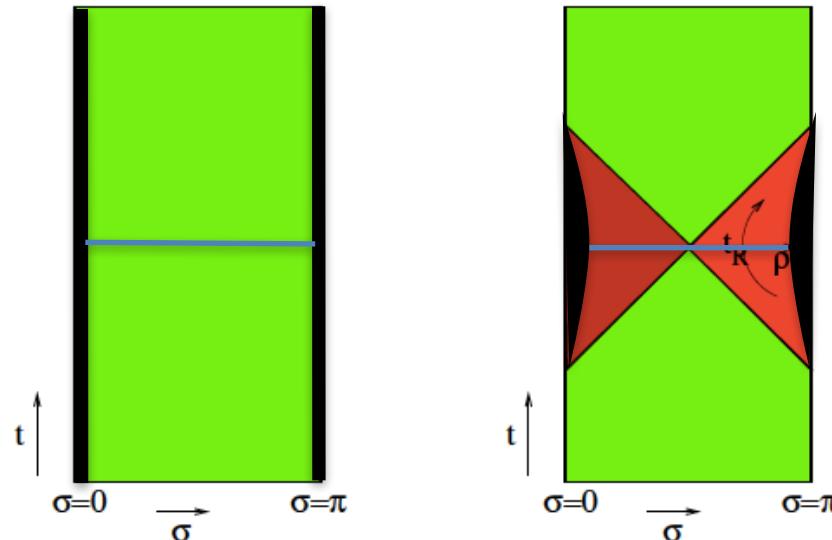
$$ds^2 = \frac{-dt^2 + d\sigma^2}{\sin^2 \sigma}$$

Global

$$ds^2 = -(r^2 - 1)d\tau^2 + \frac{dr^2}{(r^2 - 1)}$$

Thermal/Rindler

Nearly AdS₂



$$ds^2 = \frac{-dt^2 + d\sigma^2}{\sin^2 \sigma}$$

Global

$$ds^2 = -(r^2 - 1)d\tau^2 + \frac{dr^2}{(r^2 - 1)}$$

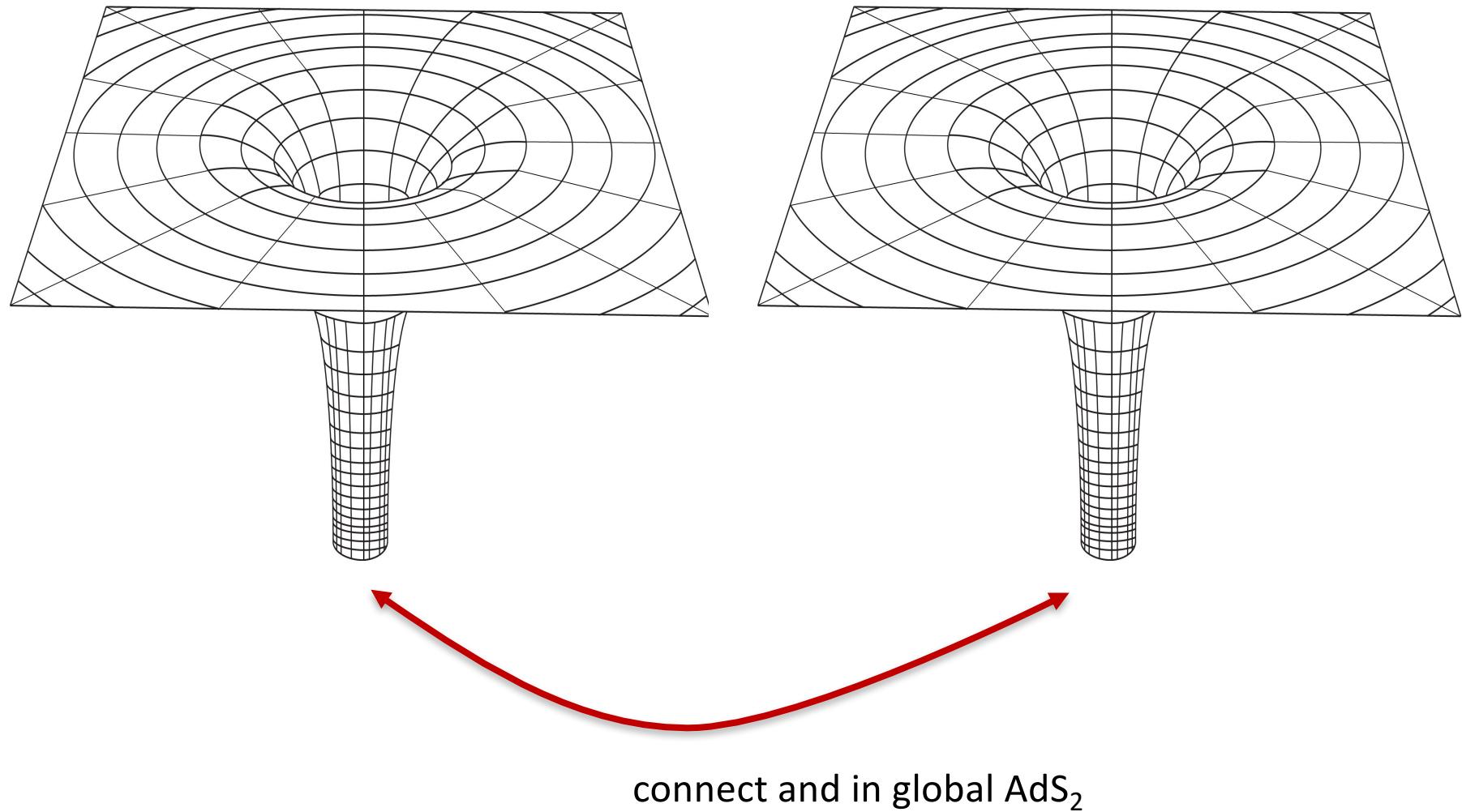
Thermal/Rindler

Connect them to flat space, so that t is an isometry.

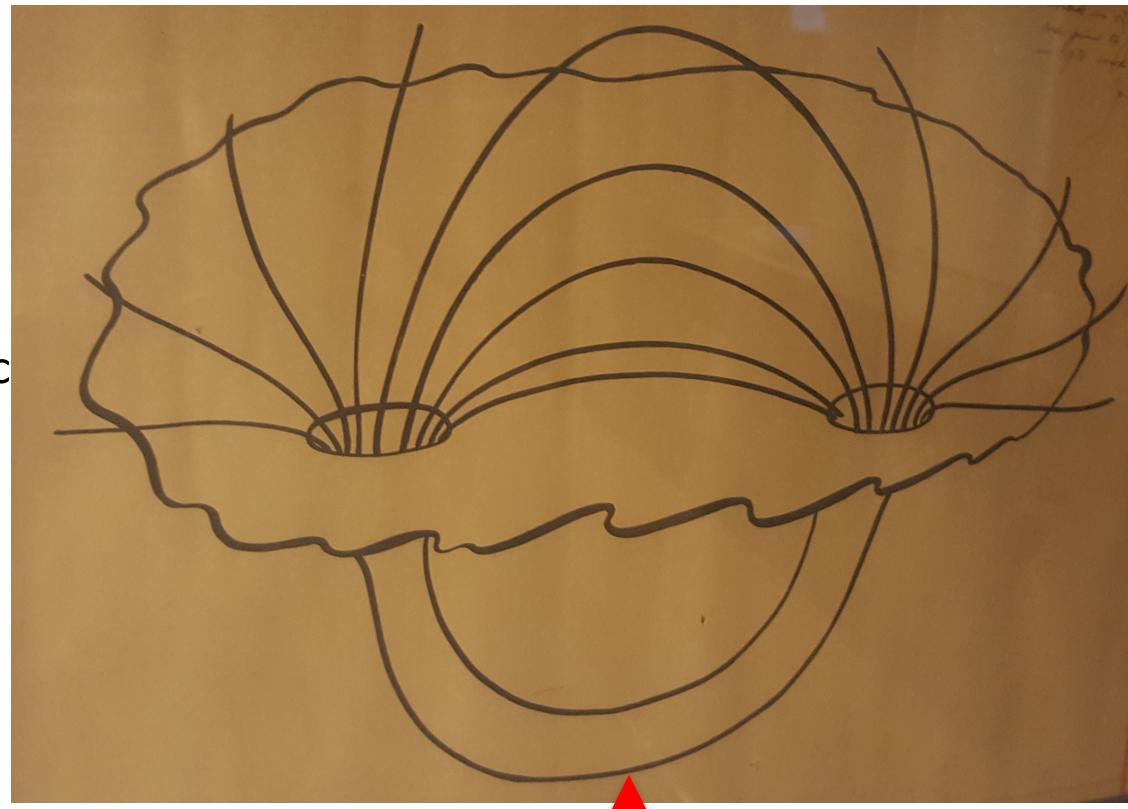
The acquire non-zero energy when the throat has finite length

$$M - q = q^3 T^2 = \frac{q^3}{\beta^2}$$

Connect a pair black holes



Connect a pair black holes



Positive magnetic
charge

Negative magnetic
charge

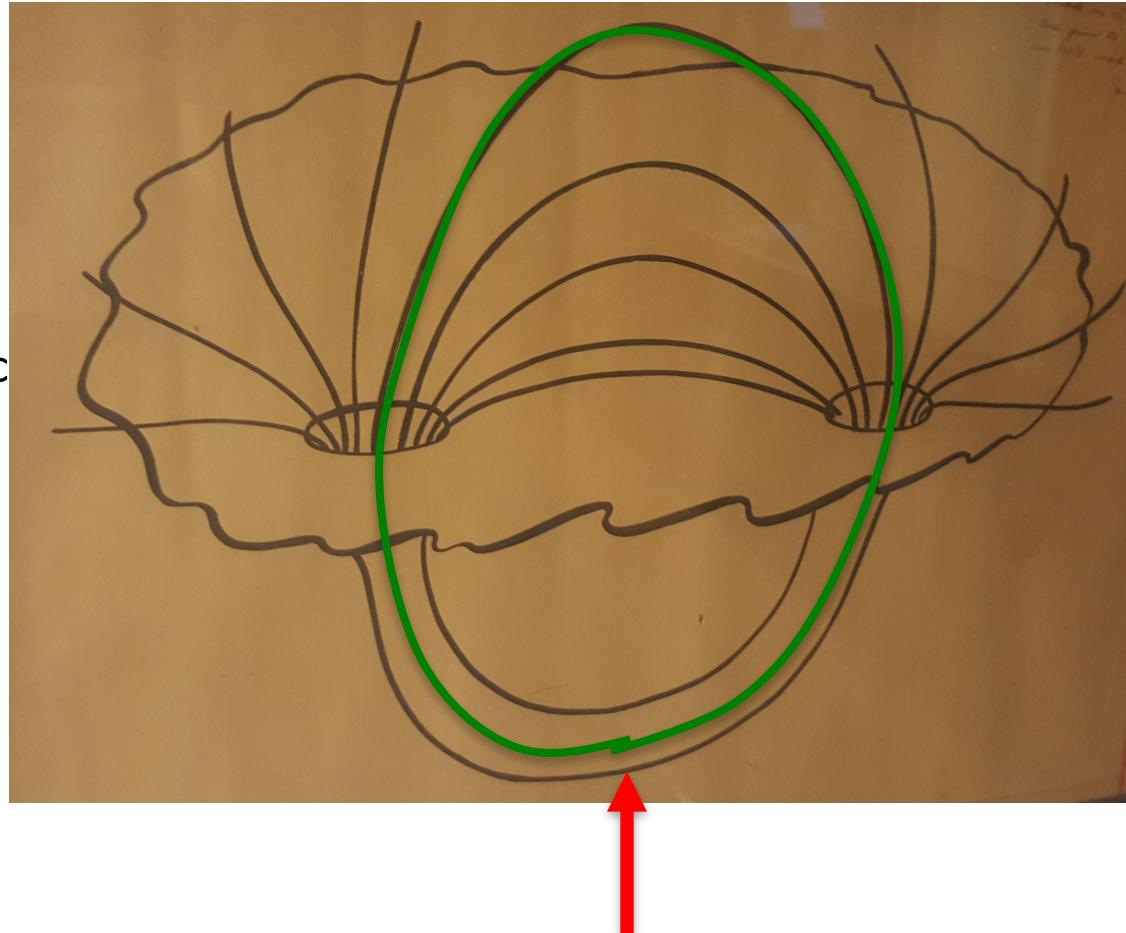
Not a solution yet.
Not a black hole.

Nearly $\text{AdS}_2 \times S^2$ wormhole of finite length

Fermion trajectories

Positive magnetic charge

Negative magnetic charge



Charged fermion moves along this closed circle.

Casimir energy

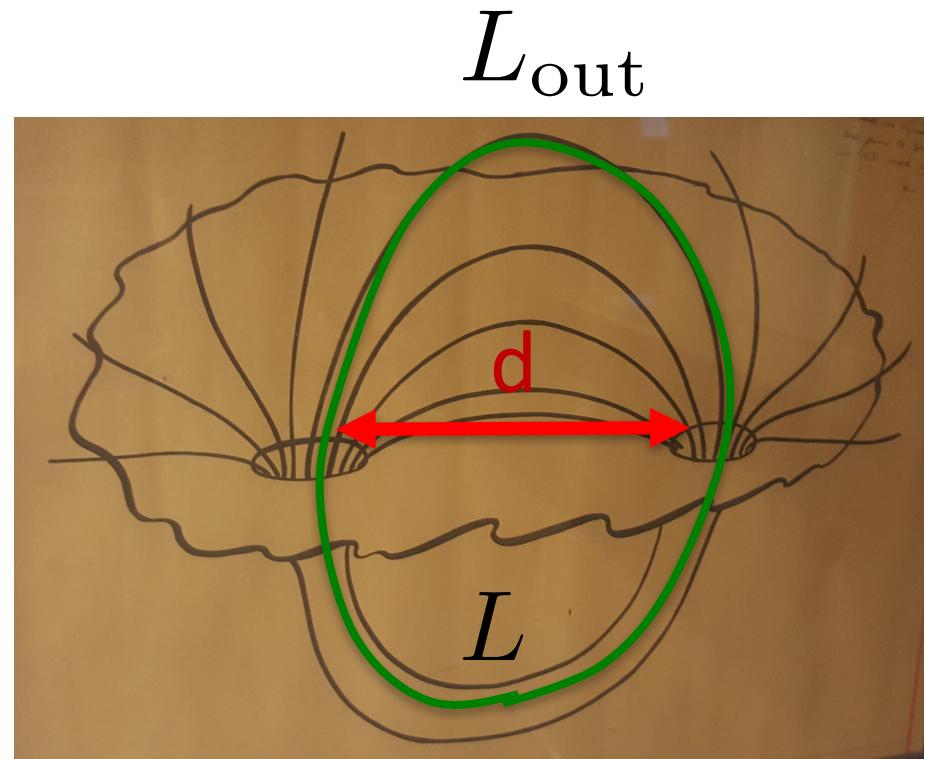
Assume: “Length of the throat” is larger than the distance.

$$L \gg L_{out} > d$$

Casimir energy is of the order of

$$E \propto -\frac{q}{L}$$

Full energy also need to take into account the conformal anomaly because AdS₂ has a warp factor.
That just changes the numerical factor.



Finding the solution

Solve Einstein equations in the throat region with the negative quantum stress tensor

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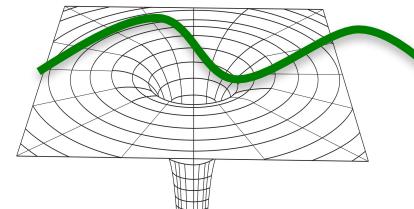
Balance the classical curvature + gauge field energy vs the Casimir energy.

$$M - q = \frac{q^3}{L^2} - \frac{q}{L}, \quad \frac{\partial M}{\partial L} = 0 \longrightarrow L \sim q^2$$

Now the throat is stabilized. Negative binding energy.

$$E_{\text{binding}} = M - q = -\frac{1}{q} = -\frac{1}{r_s}$$

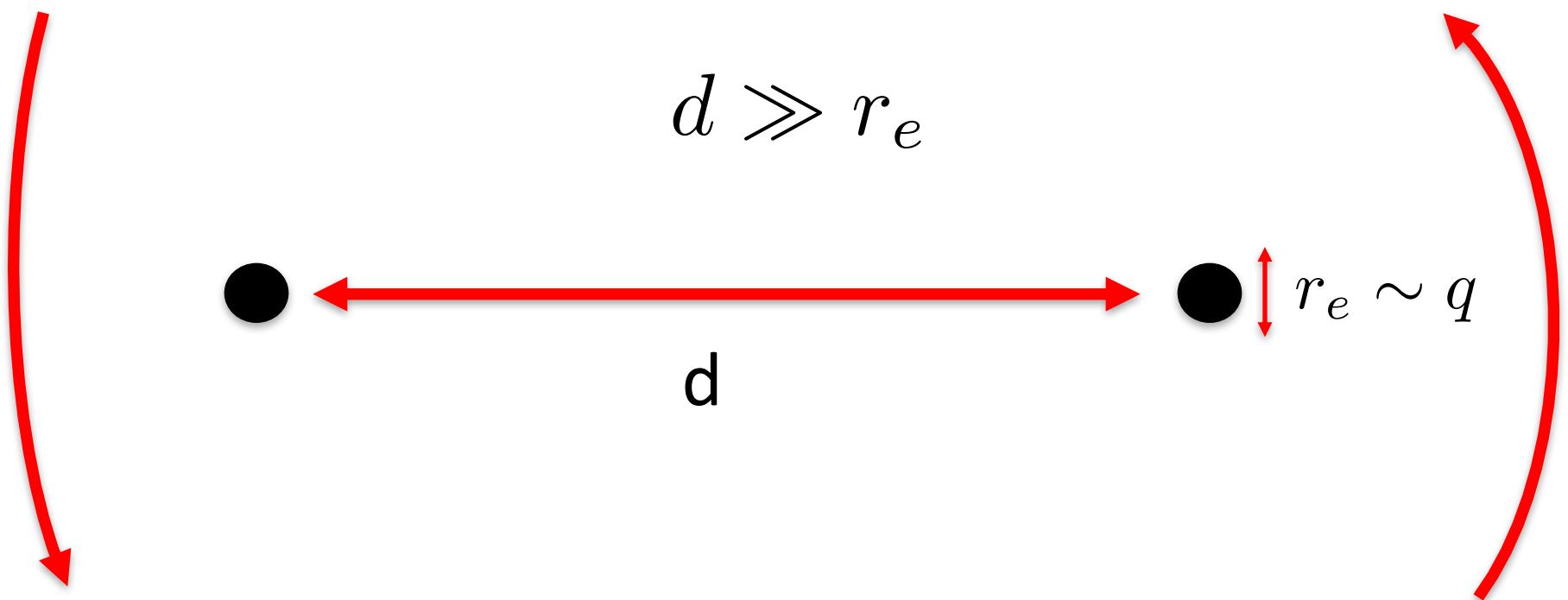
Very small. Only low energy waves can explore it



This is not yet a solution in the outside region:

The two objects attract and would fall on to each other

Adding rotation



$$\Omega = \sqrt{\frac{r_e}{d^3}}$$

Kepler rotation frequency

Throat is fragile

- We must make sure not to start sending matter into the throat that can accumulate there and produce a black hole.
- Rotation \rightarrow radiation \rightarrow effective temperature: $T \sim \Omega$
- We need that Ω is smaller than the energy gap of the throat
$$\Omega \ll \frac{1}{L}$$
- The configuration will only live for some time, until the black holes get closer..
- These issues could be avoided by going to AdS_4 ...

Some necessary inequalities

$$L \sim q^2$$

From stabilized throat solution

$$d \ll L \rightarrow d \ll q^2$$

Black holes close enough so that Casimir energy computation was correct.

$$\sqrt{\frac{q}{d^3}} = \Omega \ll \frac{1}{L} \rightarrow q^{\frac{5}{3}} \ll d$$

Black holes far enough so that they rotate slowly compared to the energy gap.

Kepler rotation frequency

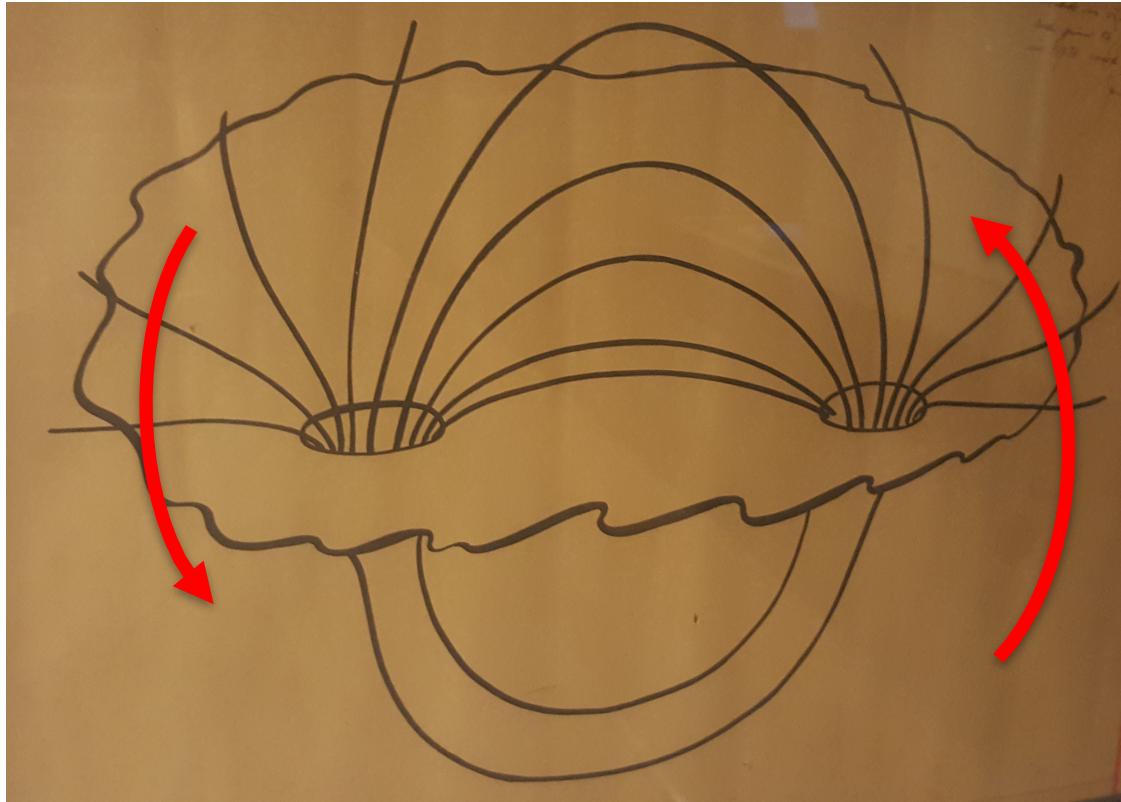
Unruh-like temperature less than energy gap

They are compatible

$$q^{\frac{5}{3}} \ll d \ll q^2$$

Other effects we could think off are also small :
can allow small eccentricity, add electromagnetic and gravitational radiation, etc. Has a finite lifetime.

Final solution



Looks like the exterior of two near extremal black holes. But they connected.
But there is no horizon!. Zero entropy solution.
It has a small binding energy.

Rotation \rightarrow temperature

$$T \sim \Omega \ll E_{\text{gap}} \sim \frac{1}{L}$$



Temperature does not create particles
in the throat

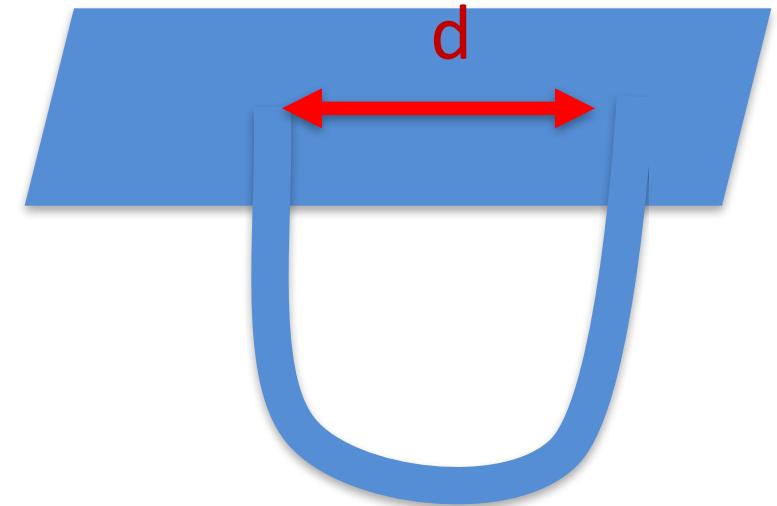
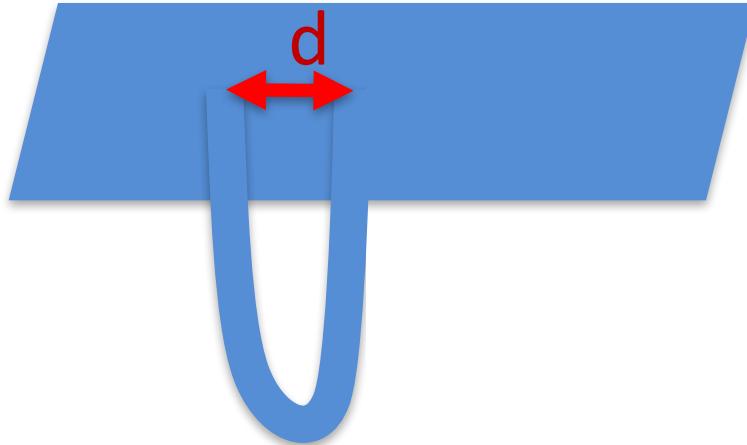
Two Black Holes : $F = -Tq^2$

Wormhole : $F = -E_{\text{binding}} = -\frac{1}{q}$

Wormhole is the stable thermodynamic phase for $T < 1/q^3$

For the solution we described so far: Wormhole is metastable.

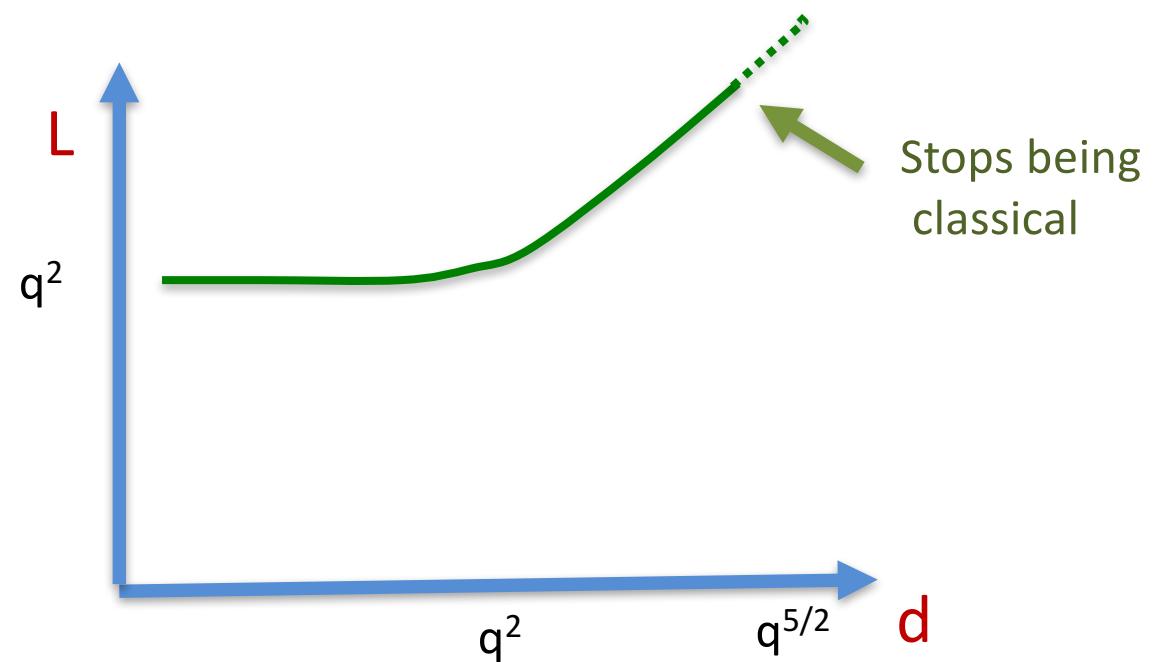
Length L as d increases



$$E_c \propto +\frac{1}{4} \frac{1}{L} - \frac{1}{L+d}$$

Conformal
anomaly

Casimir cylinder



Wormholes in the Standard Model

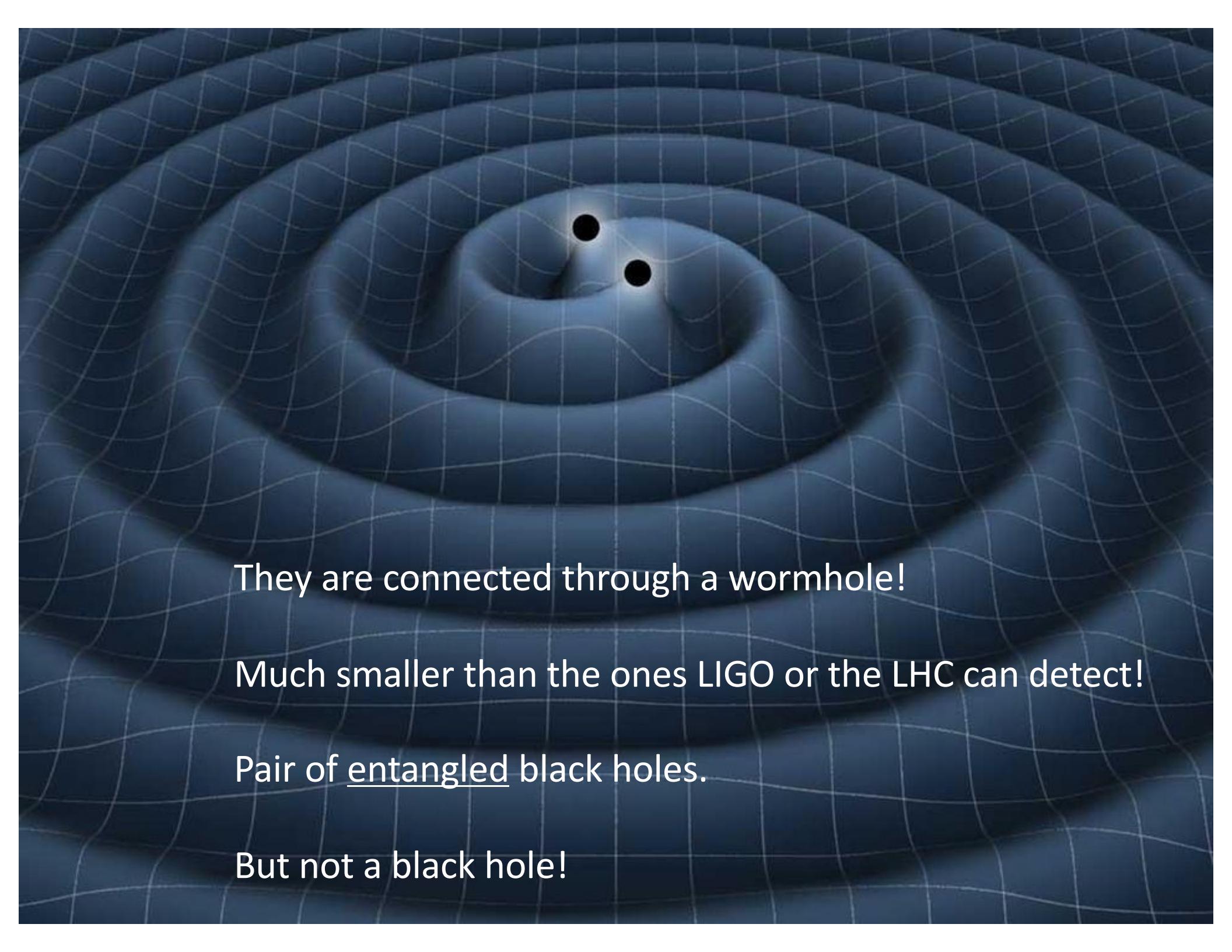
If nature is described by the Standard Model at short distances and d is smaller than the electroweak scale,

$$1 \ll q \ll 10^8$$


Distance d smaller than electroweak scale.

If the standard model is not valid \rightarrow similar ingredients might be present in the true theory.

That it can exist, does not mean that it is easily produced by some natural or artificial process.



They are connected through a wormhole!

Much smaller than the ones LIGO or the LHC can detect!

Pair of entangled black holes.

But not a black hole!

Entropy and entanglement

- Total spacetime has no entropy and no horizon.
- If we only look at one object → entanglement entropy = extremal black hole entropy
- Wormhole = two entangled black holes
- Total Hamiltonian
$$H = H_L + H_R + H_{\text{int}}$$

↑
Generated by fermions in exterior

Conclusions

- We displayed a solution of an Einstein Maxwell theory with charged fermions.
- It is a traversable wormhole in four dimensions and with no exotic matter.
- It balances classical and quantum effects.
- It has a non-trivial spacetime topology, which is forbidden in the classical theory.
- It has no horizon and no entropy.
- Can be viewed as a pair of entangled black holes.

Questions

- If we start from disconnected near extremal black holes: Can they be connected quickly enough ? → topology change.
- Could we turn it into a prediction from quantum gravity ?