NLO QCD corrections to off-shell $t\overline{t}\gamma$ production at the LHC

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arXiv:1803.09916 + arXiv:1809.08562 [hep-ph]

Outline

Motivations for $t\bar{t}\gamma$ at LHC

- the need for high precision
- state of the art

Predictions for $t\bar{t}\gamma$ at LHC 13 TeV with HELAC-NLO

- total cross sections, differential distributions
- dominant theoretical uncertainties

The $t\bar{t}\gamma/t\bar{t}$ cross section ratio at LHC 13 TeV

- analysis of correlations between $t\bar{t}\gamma$ and $t\bar{t}$
- estimate of theoretical uncertainties

Summary and outlook

Introduction

In the absence of convincing evidence for new resonances effects, precise measurements of the properties of SM particles are key to look for effects of New Physics at the LHC.

This is especially true for the top quark, where NP effects are expected to be more prominent due to its large mass scale

The LHC provides a unique opportunity for testing the properties of the top quark (and providing insights into NP) via abundant production of $t\bar{t}$ pairs.



- top-quark mass, charge
- spin correlations
- charge asymmetries
- top-quark EW couplings

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Motivations for $t\bar{t}\gamma$

Probe of the top quark charge

$$\ \hookrightarrow \ \sigma_{t\bar{t}\gamma}\sim Q_t^2 \ \ \text{@ LHC}$$

Indirect from $t\bar{t}$: $Q_t = Q_W - Q_b$







Probe of the effective $tt\gamma$ interaction

- \hookrightarrow dimension-six SMEFT
- \hookrightarrow top quark anomalous couplings

Baur *et al.* '05, Aguilar Saavedra '09 Schulze *et al.*'16, Maltoni *et al.* '16 ...

Also: $t\bar{t}\gamma \rightarrow$ irreducible background in direct BSM searches

Status of $t\bar{t}\gamma$

Experiment

- First evidence: CDF @ Tevatron CDF Collaboration '11
- Observation: ATLAS @ LHC 7 TeV ATLAS Collaboration '15
- Measurements: ATLAS/CMS @ LHC 8 TeV ATLAS and CMS Collaborations '17



Theory

• NLO QCD \rightarrow stable top quarks

Duan, Guo, Han, Ma, Wang and Zhang '09,'11; Maltoni, Pagani and Tsinikos '15

- NLO EW → stable top quarks
 Duan, Zhang, Wang, Song and Li '16
- NLO QCD → NWA: radiative decays + spin correlations

Melnikov, Scharf and Schulze '11

 NLO QCD + PS (PowHel) → LO top decays in PS Kardos and Trocsanyi '14

NLO predictions of $t\bar{t}\gamma$ have been so far restricted to on-shell tops

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$t\bar{t}\gamma$: the Narrow Width Approximation



- only the dominant double-resonant contributions are retained
- the photon can originate at the *production* (tt
 τ γ → ...) or at the *decay* stage (t → bWγ)
- sufficiently accurate for *inclusive* observables: $\Gamma_t/m_t \sim 0.8\%$

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$t\bar{t}\gamma$ in NWA

 $pp
ightarrow b ar{b} \ell^+
u_\ell j j \gamma ~$ @ 14 TeV

Melnikov, Scharf and Schulze, arXiv:1102.1967 [hep-ph]



Contributions from photon radiation in the decay stage are important

$$\sigma_{\text{NLO}} = 138.1 \text{ fb} = \underbrace{60.9}_{\gamma-Prod} + \underbrace{77.2}_{\gamma-Decay} \text{ fb}$$

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Beyond NWA

To further improve the accuracy of fixed-order predictions, one needs to release the approximation of intermediate top quarks produced on-shell.

Some technical details for $pp \rightarrow e^+ \nu_e \mu^- \bar{\nu}_\mu b \bar{b} \gamma$:

- LO: 628 diagrams for the gg channel @ $\mathcal{O}(\alpha^5 \alpha_s^2)$
- Real: 4348 diagrams for the gg channel @ $\mathcal{O}(\alpha^5 \alpha_s^3)$
- Virtual: 36032 one-loop diagrams for the gg channel @ $\mathcal{O}(\alpha^5 \alpha_s^3)$
- \hookrightarrow Compare with related benchmark processes: off-shell $t\bar{t}$ and $t\bar{t}j$



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Off-shell effects can reach tens of percents in tails

 \hookrightarrow Examples:



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The HELAC-NLO framework



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- Functionality extended to produce Ntuples of events
- Recomputing for different scales + PDFs is not practical \rightarrow use re-weighting

Predictions for $t\bar{t}\gamma$ production at LHC 13 TeV

Setup for LHC 13 TeV

Final state and parameters

- Fully leptonic decay channel: $pp \rightarrow e^+ \nu_e \mu^- \bar{\nu}_\mu b \bar{b} \gamma + X$
- All leptons and light quarks (including bottom) massless \rightarrow 5F scheme
- Top quark (pole mass): $m_t = 173.2 \text{ GeV}$
- Complex Mass Scheme: $m_t^2
 ightarrow m_t^2 i \, m_t \, \Gamma_t$ [Denner *et al.* '99, '05]

Kinematics

• exactly two b-jets , one photon , two charged leptons , missing p_T

• cuts: $p_{T \ell} > 30 \text{ GeV}$ $p_{T b} > 40 \text{ GeV}$ $p_T > 20 \text{ GeV}$ $p_{T,\gamma} > 25 \text{ GeV}$ $\Delta R_{bb} > 0.4$ $\Delta R_{\ell\ell} > 0.4$ $\Delta R_{\ell b} > 0.4$ $|y_{\ell}| < 2.5$ $|y_b| < 2.5$ $|y_{\gamma}| < 2.5$

- photon isolation condition: $\sum_{i} E_{T,i} \Theta(R R_{\gamma,i}) \le E_{T,\gamma} \left(\frac{1 \cos(R)}{1 \cos(R_{\gamma k})} \right), \quad R_{\gamma k} = 0.4$
- for the hard photon $\alpha = \alpha(0) = 1/137 \rightarrow$ predictions decreased by 3%

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Total cross sections

G.B., Hartanto, Kraus, Weber and Worek, arXiv:1803.09916 [hep-ph]



- at LO: $gg \sim 79\%$, $q\bar{q} \sim 21\%$,
- negative NLO corrections: -10%

- scale uncertainties: 35% @ LO $\hookrightarrow 14\%$ @ NLO

- estimate of non-factorizable contrib. via $\Gamma_t \rightarrow 0$ limit: 2.5% @ NLO



Differential cross sections

Focus on observables relevant for BSM searches

1. Transverse momentum of the photon: $p_{T,\gamma}$



G.B., Hartanto, Kraus, Weber and Worek, arXiv:1803.09916 [hep-ph]

- Scale uncertainties via envelope: $\left(\frac{\mu_R}{\mu_0}, \frac{\mu_F}{\mu_0}\right) = \{(2, 1), (0.5, 1), (1, 2), (1, 1), (1, 0.5), (2, 2), (0.5, 0.5)\}$

- Differential K-factor varies from -8% to -18% in plotted range
- Dynamical scale $\mu = H_T/4$ helps to improve perturbative stability

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Differential cross sections

Focus on observables relevant for BSM searches

2. Separation between photon and 2^{nd} hardest *b*-jet: $\Delta R_{b_2,\gamma}$



G.B., Hartanto, Kraus, Weber and Worek, arXiv:1803.09916 [hep-ph]

- Severe shape distortions for $\Delta R_{b_2,\gamma} > 4$ are genuine NLO effects
 - \hookrightarrow initial-state γ radiation from qg channel
- Similar effect observed for other ΔR observables

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Differential cross sections

Impact of different PDF sets on $p_{T,\gamma}$ and $\Delta R_{b_2,\gamma}$

G.B., Hartanto, Kraus, Weber and Worek, arXiv:1803.09916 [hep-ph]



 \hookrightarrow Global estimate of theoretical uncertainties [$p_{T,\gamma} \ge 25 \text{ GeV}$]:

$$\sigma_{
m NLO} (\mu = m_t/2) = (7.4 \pm 1.0^{\text{[scale]}} \pm 0.3^{\text{[PDF]}}) \text{ fb}$$

 $\sigma_{
m NLO} (\mu = H_T/4) = (7.5 \pm 0.5^{\text{[scale]}} \pm 0.3^{\text{[PDF]}}) \text{ fb}$

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The $t\bar{t}\gamma$ / $t\bar{t}$ cross section ratio at LHC 13 TeV

The ratio

Instead of considering the *absolute* $t\bar{t}\gamma$ cross section, normalize to $t\bar{t}$

$$\mathcal{R} = \frac{\sigma(pp \to e^+ \nu_e \mu^- \bar{\nu}_\mu b \bar{b} \gamma)}{\sigma(pp \to e^+ \nu_e \mu^- \bar{\nu}_\mu b \bar{b})}$$

Advantages

Experiment: more accurate measurement

→ common systematics cancelled in the ratio (e.g. b-jet reconstruction efficiency, luminosity ...)

- Theory: more accurate prediction
 - ↔ theoretical uncertainties (dominated by scale variation) can be dramatically reduced provided the two processes are *correlated*

Melnikov, Scharf, Schulze '11; Mangano, Rojo '12; G.B, Worek '14; Schulze, Soreq '16 ...

How strongly correlated are $t\bar{t}\gamma$ and $t\bar{t}$?

Looking for correlations

 $t\bar{t}\gamma$ vs $t\bar{t}$ @ LHC : distributions normalized to unit



G.B., Hartanto, Kraus, Weber and Worek, arXiv:1809.08562 [hep-ph]

Kinematics of *b*-jets and leptons in $t\bar{t}\gamma$ and $t\bar{t}$ show correlations

Compare to an example of *uncorrelated* processes: $t\bar{t}b\bar{b}$ vs $t\bar{t}jj$ @ LHC \rightarrow see HP2 Workshop '14



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Differential cross section ratios

G.B., Hartanto, Kraus, Weber and Worek, arXiv:1809.08562 [hep-ph]



Using *correlated* scales in $t\bar{t}\gamma$ and $t\bar{t}$ helps to constrain uncertainty bands



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Differential cross section ratios

G.B., Hartanto, Kraus, Weber and Worek, arXiv:1809.08562 [hep-ph]



 \hookrightarrow Global estimate of theoretical uncertainties $[p_{T,\gamma} \ge 25 \text{ GeV}]$:

$$\begin{split} \mathcal{R}(\mu = m_t/2) &= (\mathbf{4.56} \pm \mathbf{0.25}^{\text{[scale]}} \pm \mathbf{0.02}^{\text{[PDF]}}) \cdot 10^{-3} \\ \mathcal{R}(\mu = H_T/4) &= (\mathbf{4.62} \pm \mathbf{0.06}^{\text{[scale]}} \pm \mathbf{0.02}^{\text{[PDF]}}) \cdot 10^{-3} \end{split}$$

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Summary and outlook

- First NLO predictions of $pp \rightarrow t\bar{t}\gamma$ with fully leptonic decays, including off-shell and non-resonant effects at $\mathcal{O}(\alpha^5 \alpha_s^3)$
- Judicious dynamical scales quite effectively account for the multi-scale nature of the process
- NLO accuracy is important for a proper description of some observables relevant for BSM searches, e.g. ΔR_{bγ}
- Correlations between $t\bar{t}\gamma$ and $t\bar{t}$ production can be exploited to constrain theoretical uncertainties
- The $t\bar{t}\gamma/t\bar{t}$ cross section ratio has interesting potential in searches for BSM effects
- Next steps:
 - comparisons with NWA
 - pheno applications: SM parameter extraction,

constraining anomalous couplings

Backup slides

Differential cross sections for $pp \rightarrow e^+ \nu_e \mu^- \bar{\nu}_\mu b \bar{b} \gamma + X$ at LHC 13 TeV, based on the scale choice $\mu_R = \mu_F = m_t/2$



G.B., Hartanto, Kraus, Weber and Worek, arXiv:1803.09916 [hep-ph]

Differential cross sections for $pp \rightarrow e^+ \nu_e \mu^- \bar{\nu}_\mu b \bar{b} \gamma + X$ at LHC 13 TeV, based on the scale choice $\mu_R = \mu_F = \frac{H_T}{4}$



G.B., Hartanto, Kraus, Weber and Worek, arXiv:1803.09916 [hep-ph]

NLO cross sections for $pp \rightarrow e^+ \nu_e \mu^- \bar{\nu}_\mu b \bar{b} + X$ and $pp \rightarrow e^+ \nu_e \mu^- \bar{\nu}_\mu b \bar{b} \gamma + X$ at LHC 13 TeV, for various scale and PDF choices

(The errors refer to scale uncertainties)

PDF set, $\mu_R = \mu_F = \mu_0$	$\sigma^{\rm NLO}_{e^+\nu_e\mu^-\bar\nu_\mu b\bar b} {\rm [fb]}$	$\sigma_{e^+\nu_e\mu^-\bar{\nu}_\mu b\bar{b}\gamma}^{\rm NLO} \text{ [fb]}$ $p_{T,\gamma} > 25 \text{ GeV}$	$ \sigma^{\rm NLO}_{e^+\nu_e\mu^-\bar{\nu}_\mu b\bar{b}\gamma} \text{[fb]} $ $ p_{T,\gamma} > 50 \text{GeV} $
CT14, $\mu_0 = m_t/2$ CT14, $\mu_0 = H_T/4$	$1629.4_{-144.7(9\%)}^{+18.4(1\%)} \\ 1620.5_{-118.8(7\%)}^{+21.6(1\%)} $	$\begin{array}{c} 7.436^{+0.074}_{-1.034} \begin{array}{c} (1\%) \\ (14\%) \\ 7.496^{+0.099}_{-0.457} \begin{array}{c} (1\%) \\ (6\%) \end{array}$	$\begin{array}{c} 3.081^{+0.050~(2\%)}_{-0.514~(17\%)}\\ 3.125^{+0.040~(1\%)}_{-0.142~(4\%)}\end{array}$
MMHT14, $\mu_0 = m_t/2$ NNPDF3.0, $\mu_0 = m_t/2$	$\begin{array}{r} 1650.5_{-152.7\ (9\%)}^{+17.0\ (1\%)} \\ 1695.0_{-153.3\ (9\%)}^{+18.4\ (1\%)} \end{array}$	$7.490^{+0.080}_{-1.081} {}^{(1\%)}_{(14\%)}$ $7.718^{+0.078}_{-1.102} {}^{(1\%)}_{(14\%)}$	$\begin{array}{c} 3.093^{+0.053}_{-0.535} \begin{array}{(} 2\%) \\ 3.195^{+0.054}_{-0.550} \begin{array}{(} 2\%) \\ -0.550 \end{array} \end{array}$

Bevilacqua, Hartanto, Kraus, Weber and Worek, arXiv:1809.08562 [hep-ph]

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