

# **Are Lepton-Flavour Violating Decays Within Experimental Reach?**

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## **Disclaimer**

*I'll provide arguments for LFV in B decays.*

*By well-known pre-LHC arguments, LFV is expected in leptonic decays as well.*

*Will comment on this in last slide.*

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- *RK(\*) hint at Lepton-Universality Violation (LUV) in  $b \rightarrow s\ell\ell$ , the effect being in muons, rather than electrons*

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*That's what one expects of new interactions above the EW scale*

## **$b \rightarrow s$ anomalies: EFT understanding**

- Consider the following Hamiltonian

$$H_{\text{SM+NP}}(\bar{b} \rightarrow \bar{s} \mu \mu) = -\frac{4G_F}{\sqrt{2}} V_{tb}^* V_{ts} \frac{\alpha_{\text{em}}}{4\pi} \left[ \bar{b}_L \gamma^\lambda s_L \cdot \left( C_9^{(\mu)} \bar{\mu} \gamma_\lambda \mu + C_{10}^{(\mu)} \bar{\mu} \gamma_\lambda \gamma_5 \mu \right) \right]$$

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About equal size & opposite sign  
in the SM (at the  $m_b$  scale)

- In the SM, one has (by accident)  $C_9 \simeq -C_{10}$  at the  $m_b$  scale



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- A fully quantitative test by a global fit leads to the same conclusions.

See [Ghosh, Nardecchia, Renner, 2014] [global fits]

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- Note:** primed fields

- Fields are in the “gauge” basis (= primed)
- They need to be rotated to the mass eigenbasis 
- So, in general, this rotation induces LUV and LFV effects

$$b'_L \equiv (d'_L)_3 = (U_{L3i}^d) (d_L)_i$$

$$\tau'_L \equiv (\ell'_L)_3 = (U_{L3i}^\ell) (\ell_L)_i$$

*mass basis*

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The current  $BR(B^+ \rightarrow K^+ \mu e)$  limit yields the weak bound

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We expect this U-matrix factor to be  $O(1)$  – or larger – in at least one of these channels, just because  $U_L^\ell$  is unitary

In particular, note that  $BR(B^+ \rightarrow K^+ \mu \tau)$  scales as  $|(U_L^\ell)_{33}| / |(U_L^\ell)_{32}|^2$

## More quantitative LFV predictions? Need flavour models

- **Ex. 1 [DG, Lane, 2015]**

- *SM Yukawa couplings are brought to diagonal form through rotations like  $(U_L^\ell)^\dagger Y_\ell U_R^\ell = \hat{Y}_\ell$*

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*In any of these models one gets effects of  $O(10^{-8})$*

*This confirms the general argument given previously*

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hence in general less promising

See recent reappraisal in [Crivellin et al., 1601.00970]

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- *Obviously, LFV modes don't suffer from this shortcoming*

*Competitive playground for NA62, and even for LHCb...*

*[work in progress]*

## LFV in K decays

- The interaction advocated in [Glashow et al.]

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( BNL E871 Collab., 1998 )

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- Actually, under well-motivated TH assumptions, one gets even larger signals [work-in-prog.]

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*So, BSM LUV  $\Rightarrow$  BSM LFV (i.e. not suppressed by  $m_\nu$ )*

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- *Theory arguments in classical papers such as [Hall-Kostelecky, Raby, NPB 1986]*

*[Barbieri-Hall, PLB 1994]*

## Models with LUV and no LFV

**Ex. 1:** [Alonso, Grinstein, Martin-Camalich, 2015]

- Take Minimal Flavour Violation (MFV) in the lepton sector
  - By def, in MFV the only sources of flavour violation are the SM ones, i.e. the SM Yukawas
  - Tricky to define MFV in the lepton sector:  
we don't know whether LH  $\nu$  are Dirac or Majorana and whether RH  $\nu$  exist at all.  
Must-read ref: Cirigliano-Grinstein-Isidori-Wise, NPB 2005
- Bottom line: In such scenarios, LFV couplings are related to LH  $\nu$  masses.  
(Neglecting CPV in the LH  $\nu$  mass matrix, the above statement is generic within MLFV.)  
 Low-energy LFV processes are generally small, being suppressed by LH  $\nu$  masses

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*Plausible mechanism? Fine-tuning in model space?*