



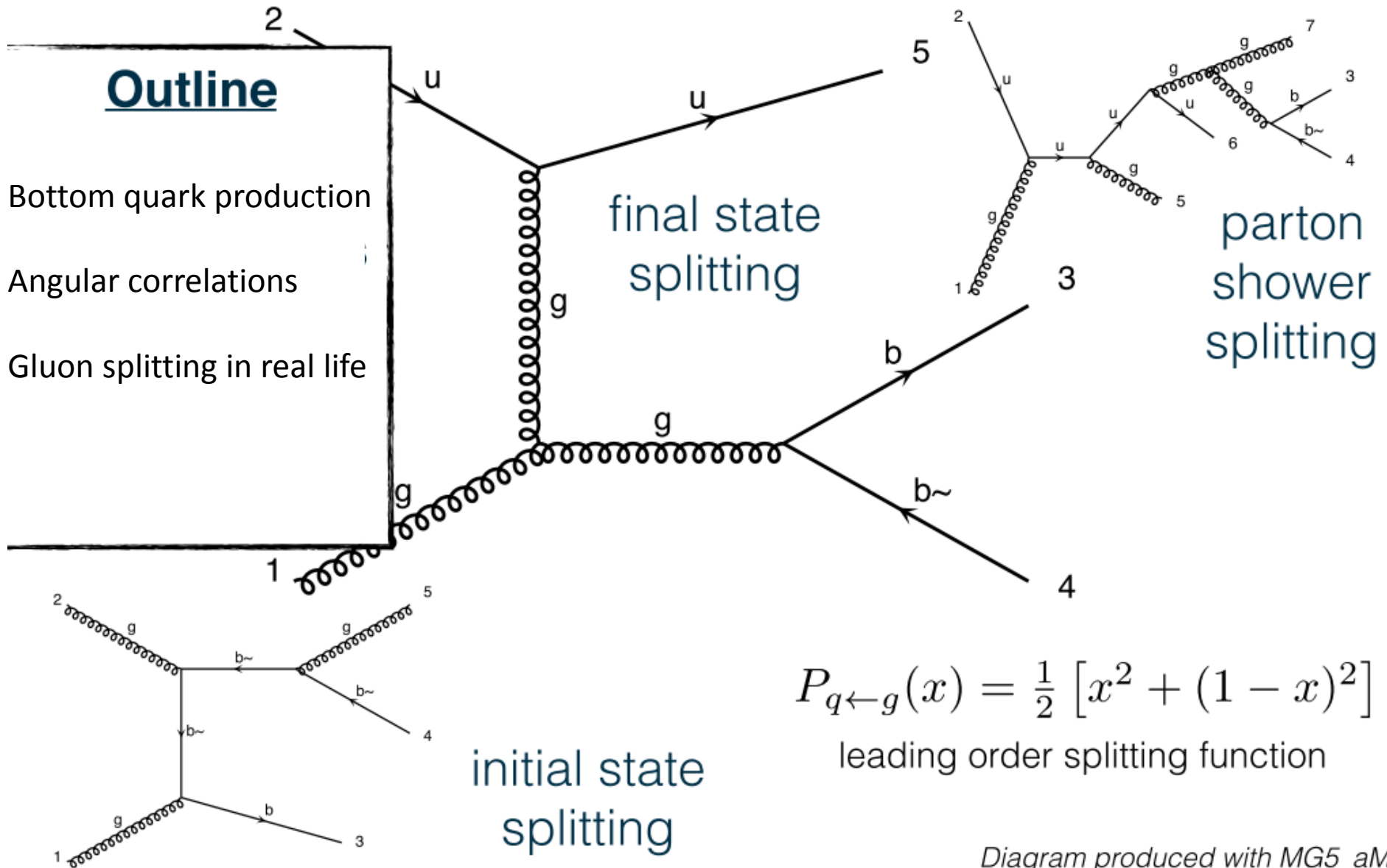
There is something about... gluon splitting

Luca Perrozzi (ETH Zurich)
CMS Heavy flavour tagging workshop
Bruxelles, April 10th 2018

Gluon Splitting to bottom quarks at the LHC

Outline

- Bottom quark production
- Angular correlations
- Gluon splitting in real life



$$P_{q \leftarrow g}(x) = \frac{1}{2} [x^2 + (1-x)^2]$$

leading order splitting function

Diagram produced with MG5_aMC

Bottom quark production at the LHC

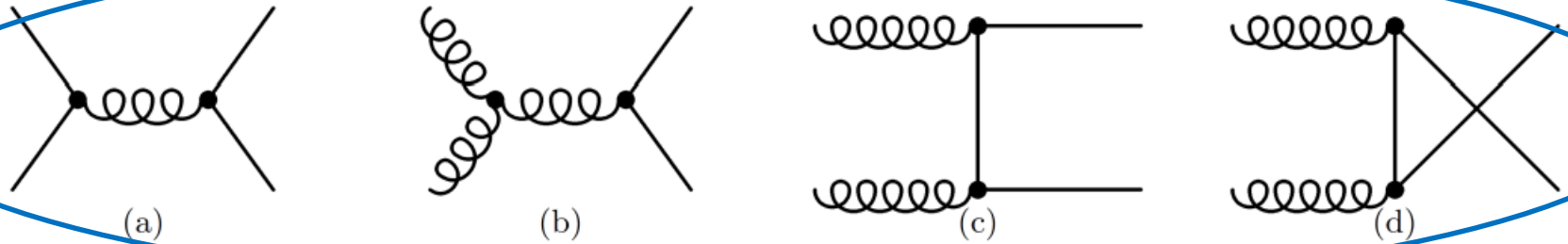


Figure 3.5: Leading order diagrams for heavy-quark pair production: (a) quark-antiquark annihilation $q\bar{q} \rightarrow Q\bar{Q}$, (b)-(d) gluon-gluon fusion $gg \rightarrow Q\bar{Q}$.

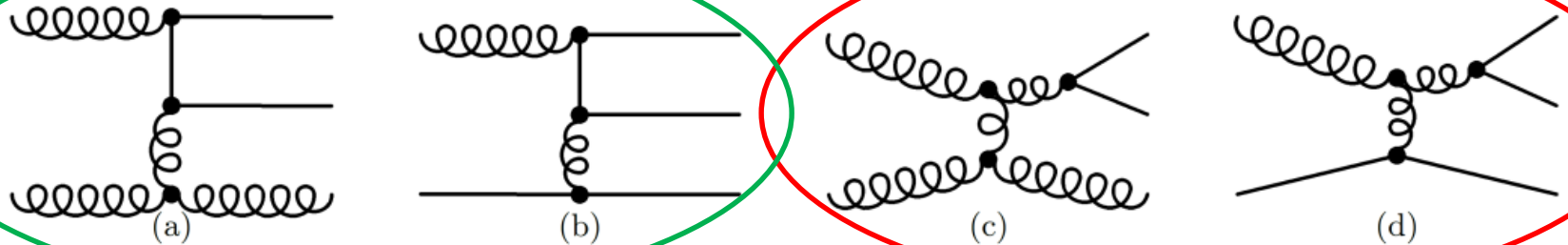
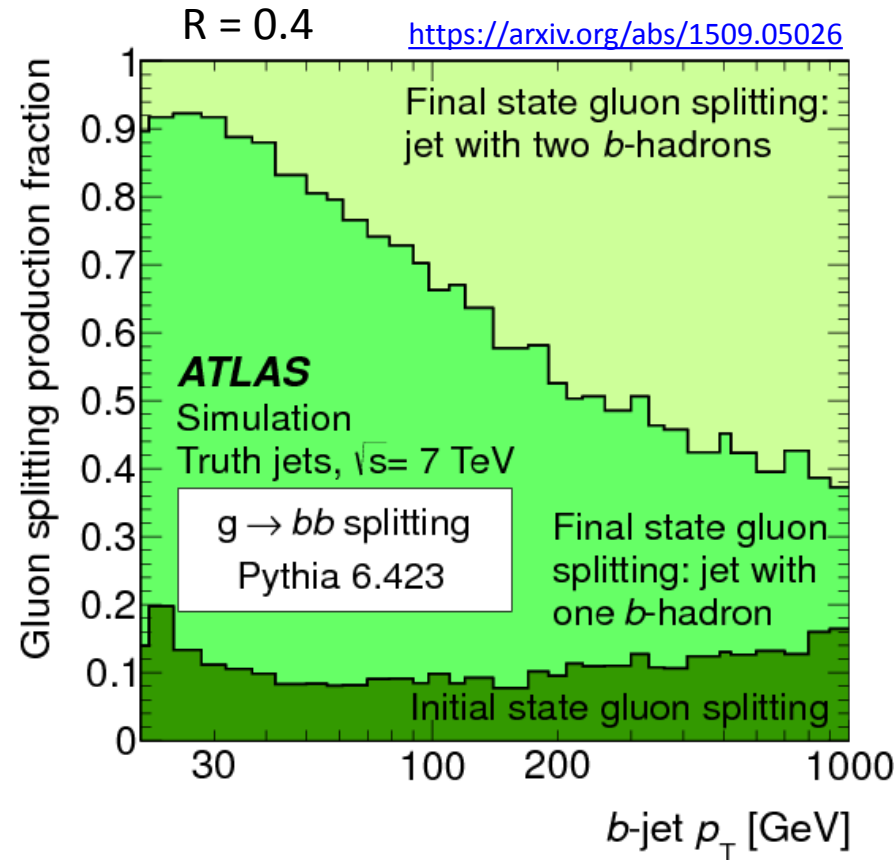
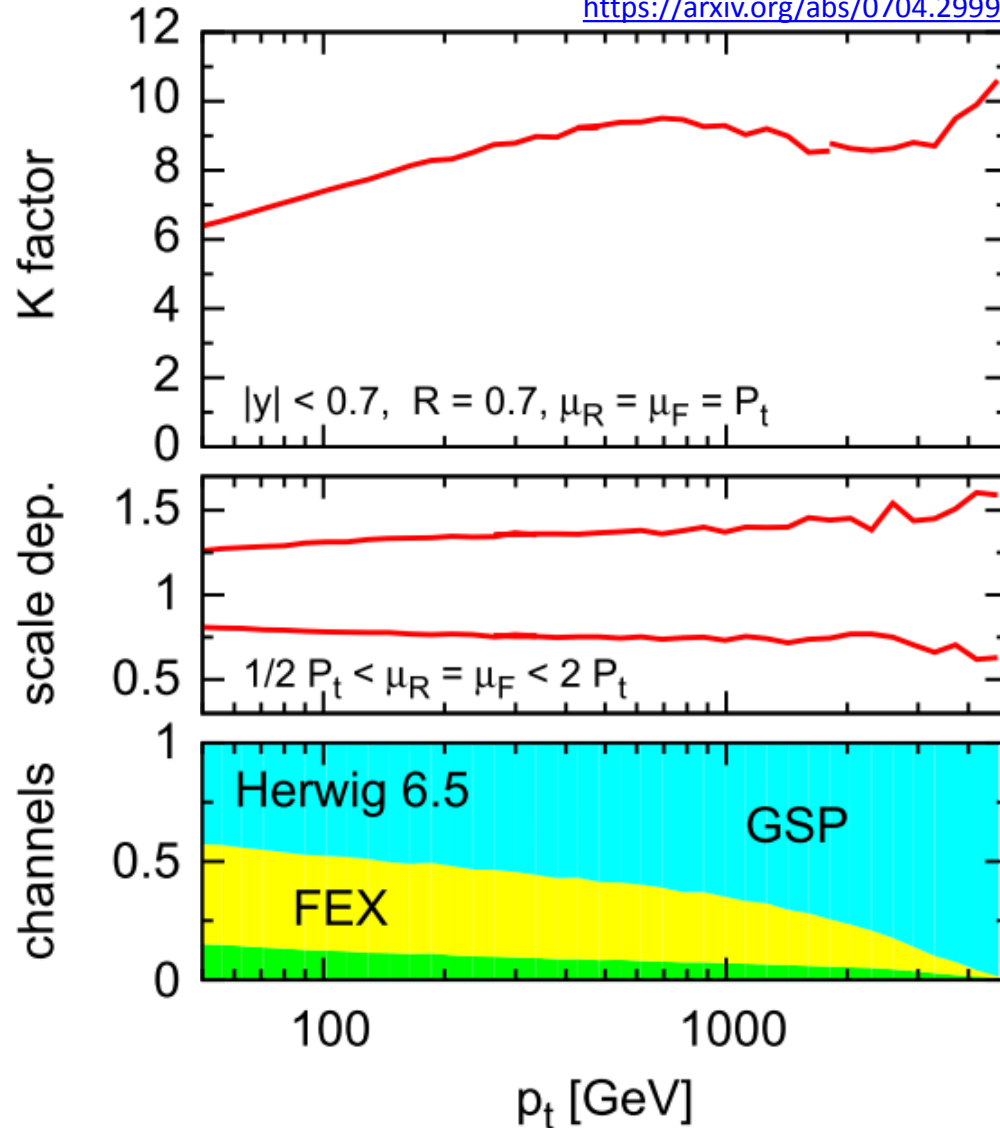


Figure 3.6: Next-to-leading order diagrams for heavy-quark pair production: (a),(b) flavor excitation; (c),(d) gluon splitting.

Bottom quark production at the LHC

LHC (14 TeV)

<https://arxiv.org/abs/0704.2999>



Bottom quark production at the LHC

- General properties of the (initial state) gluon splitting:
 - Low $\Delta\phi$
 - Low $\Delta\eta$
 - Mostly (but not only) similar p_T between the two quarks

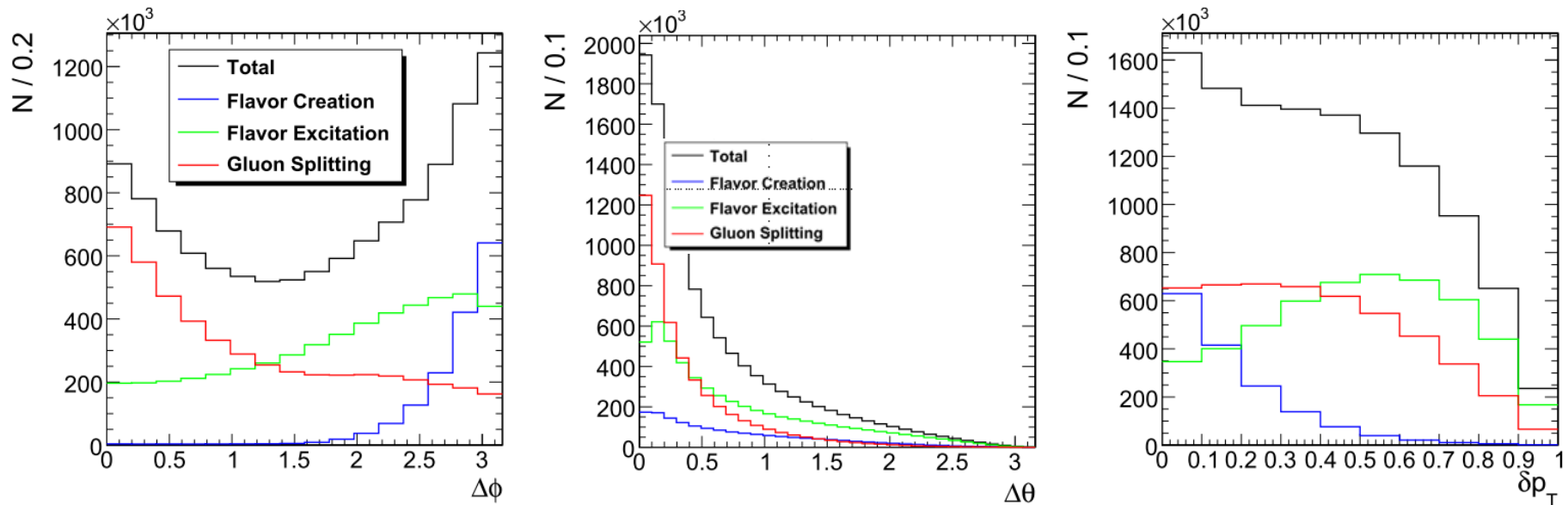


Figure 3.7: Correlated variables between simulated b -quark pairs produced from proton collisions at 10 TeV are also shown for the different mechanisms [108]: $\Delta\phi$ (top left), $\Delta\eta$ (top right), and the quark momentum asymmetry in the transverse plane δp_T (bottom).

Rich and diverse probe of the SM

Perturbative QCD

-essentially the only (nearly) direct measurement of a parton splitting function

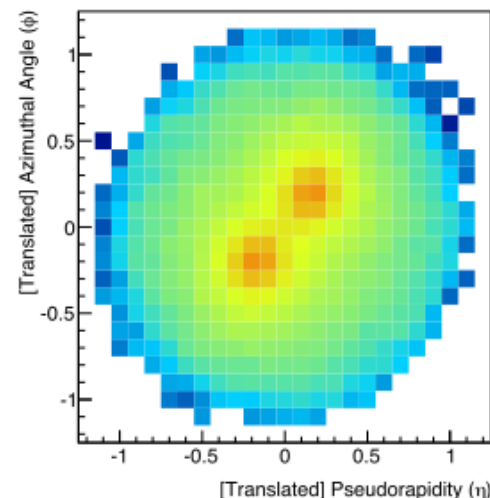
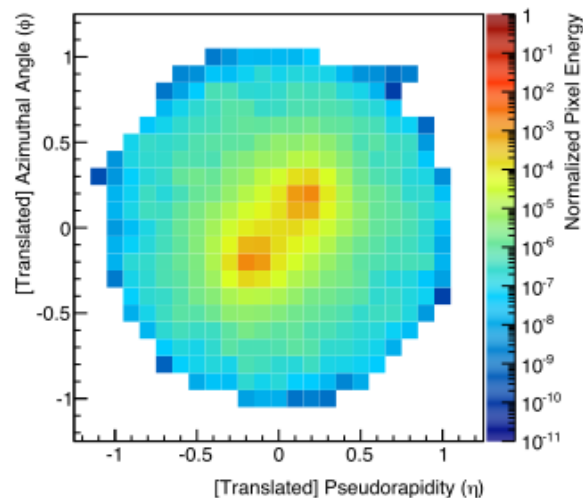
-pure source of gluon jets (though complicated by B-hadrons)

Non-perturbative QCD

-pure source of color octets →

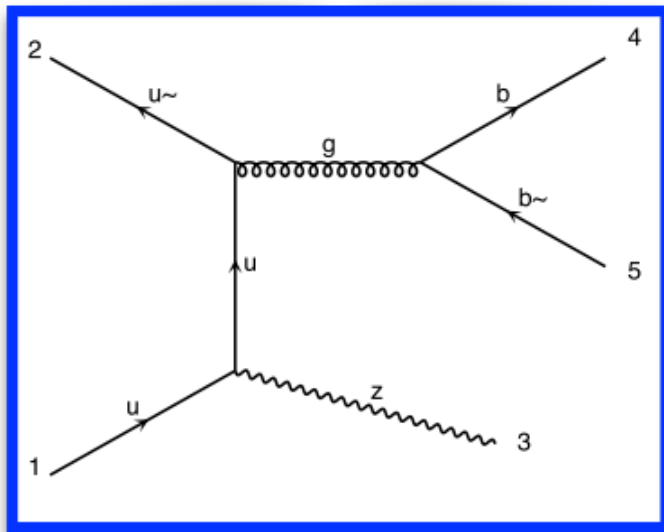
Higgs Boson (self-coupling)

-important background to many Higgs processes (VH, HH, BSM)



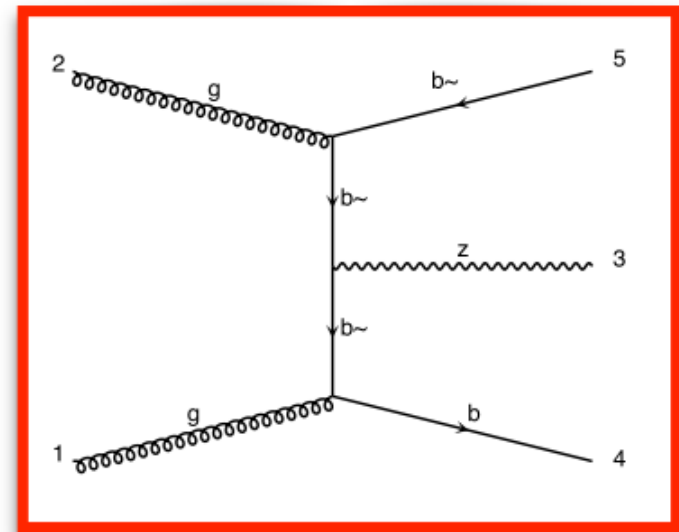
What have we learnt from LHC Run 1?

- So far, several 7 TeV measurements:
 - $b\bar{b}$
 - $Z+b(b)$ (gluon splitting diluted by gg -induced process)
 - $W+b(b)$ (only gluon splitting contributes)



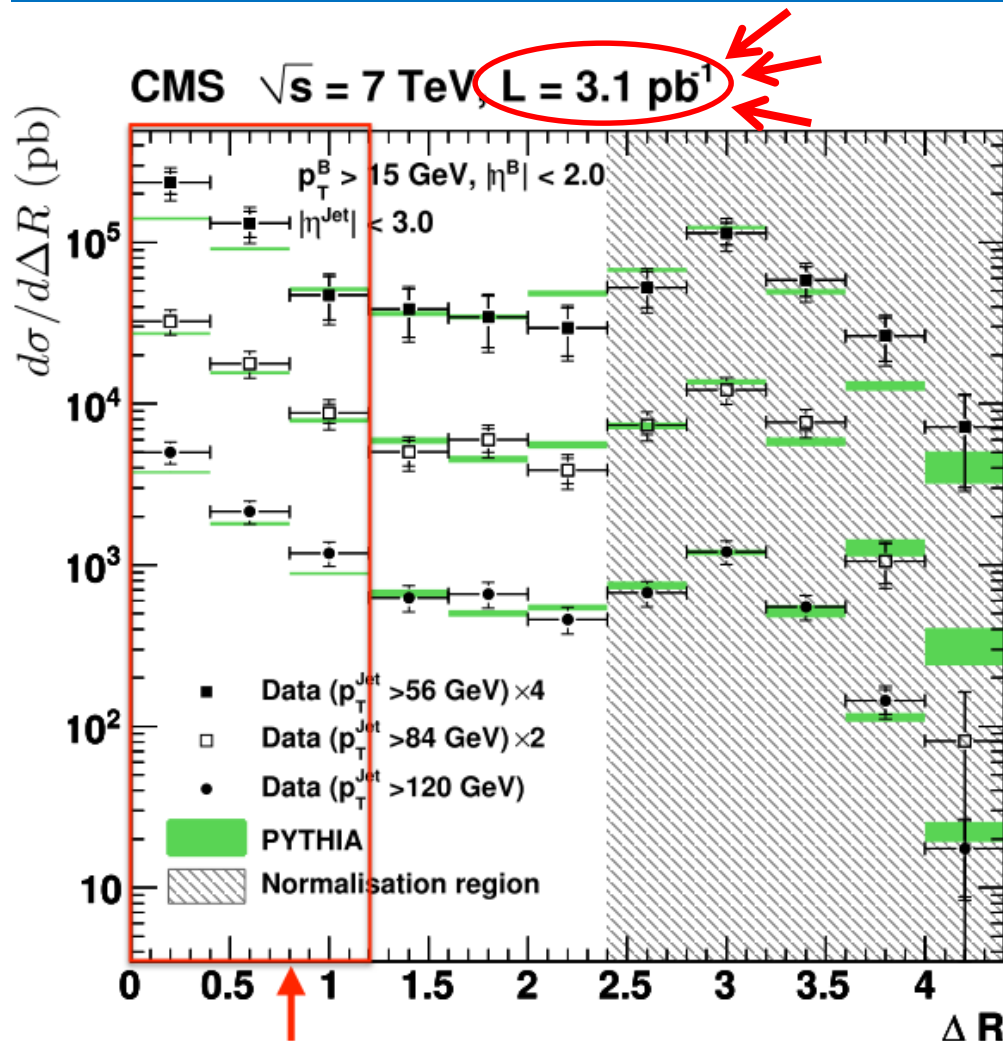
Not FS gluon
splitting →

← enhance
by looking at
low DR

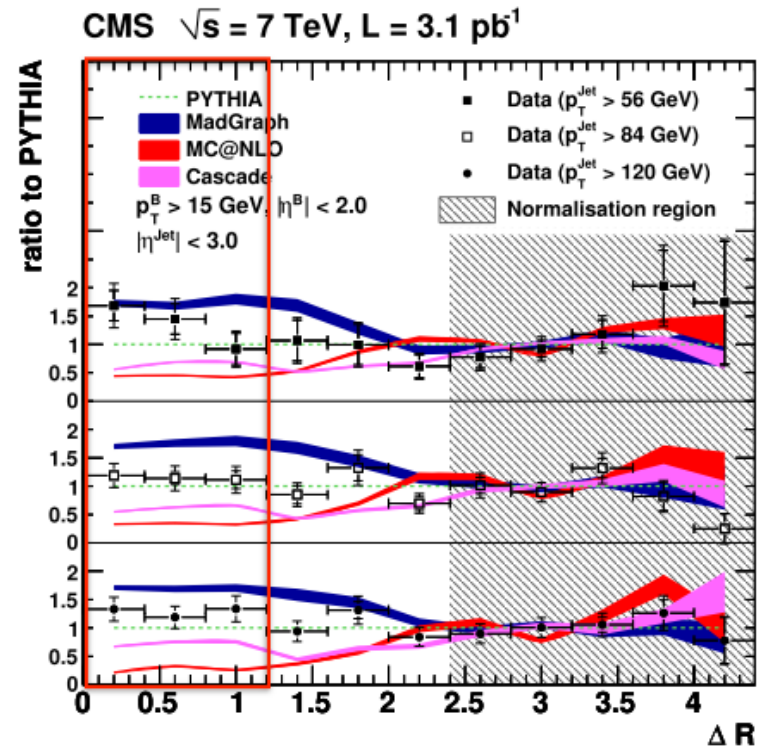


Why no $\gamma+b\bar{b}$? For gluon splitting,
this is more interesting than $Z+b\bar{b}$

CMS di-hadrons at 7 TeV



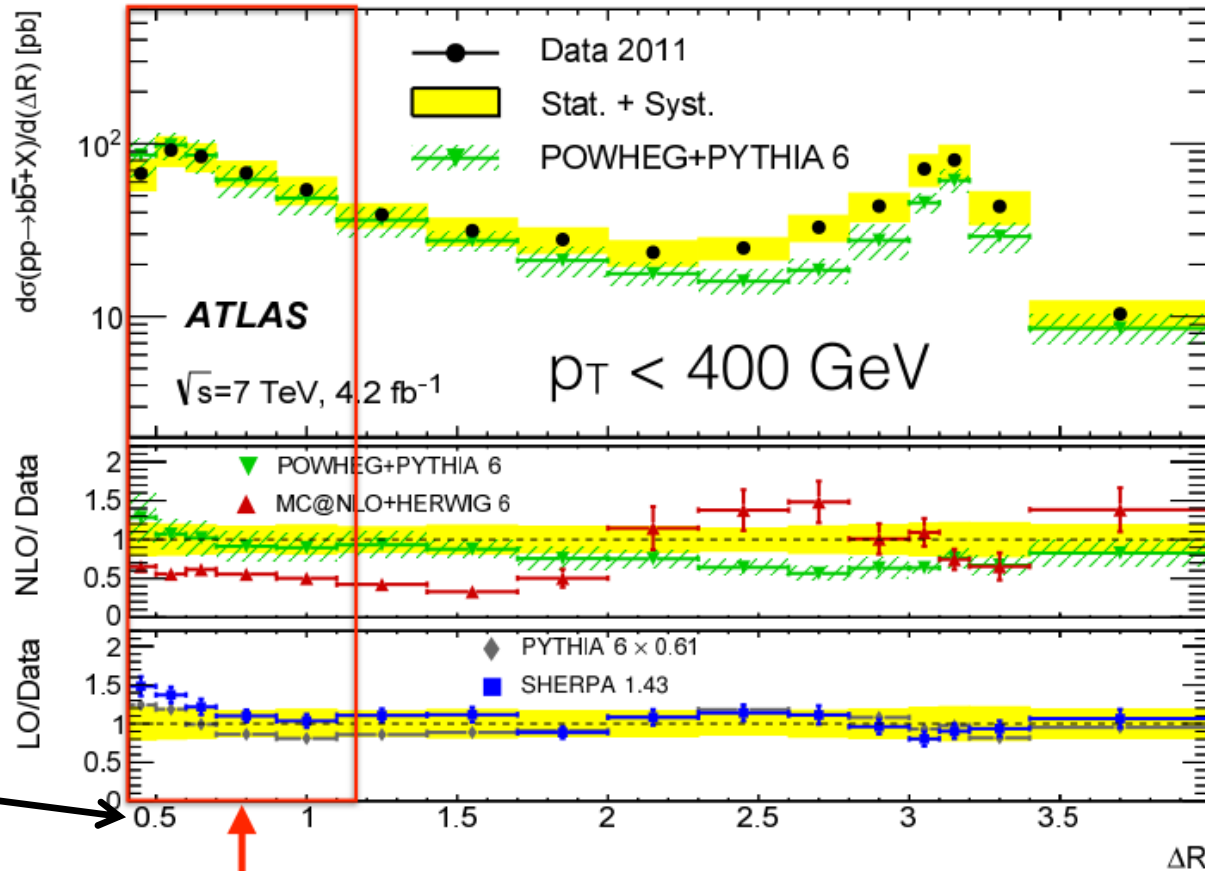
Starting to see the shape in the gluon-splitting dominated regime



As with the ATLAS result, significant differences with the MC (though Pythia is not so bad), though the comparisons there are by now outdated.

ATLAS di-bjets at 7 TeV

Higher p_T (though still relatively low) - trigger limited at low p_T

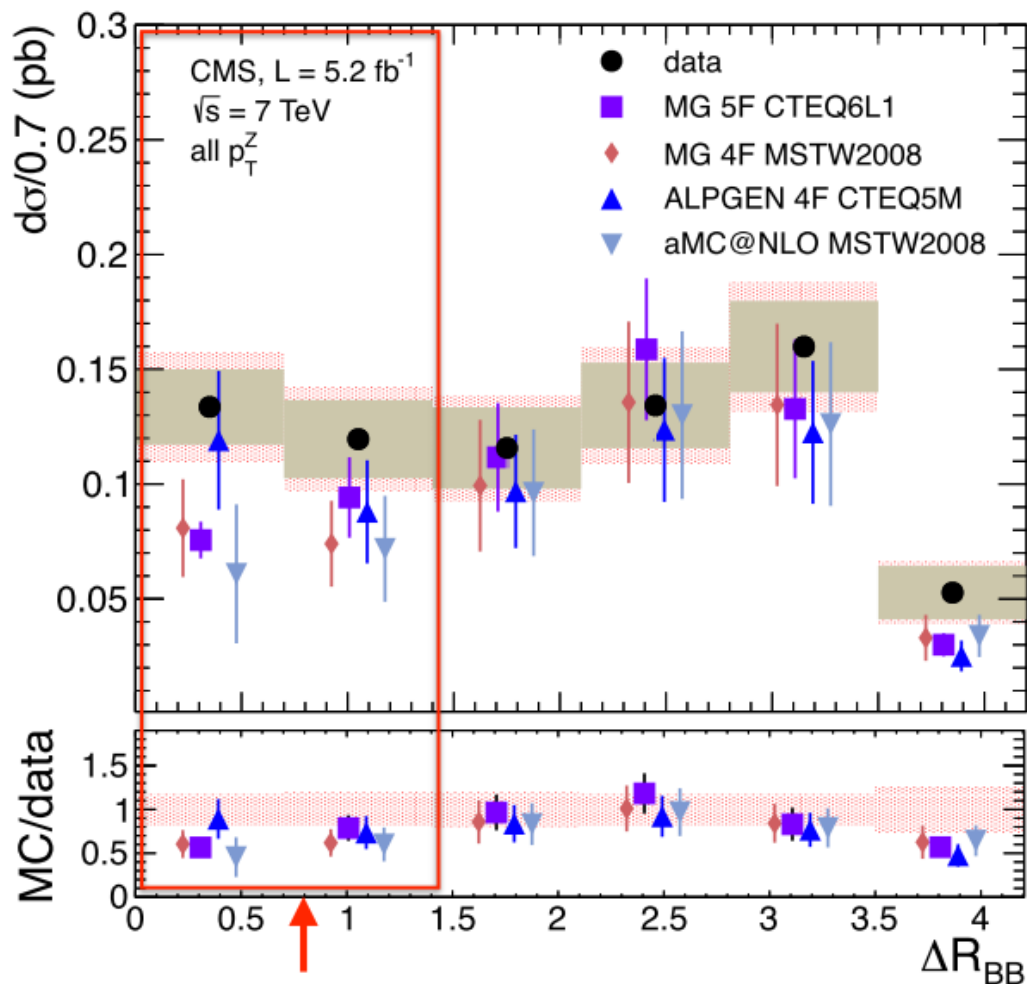


Penalty to pay when using jets

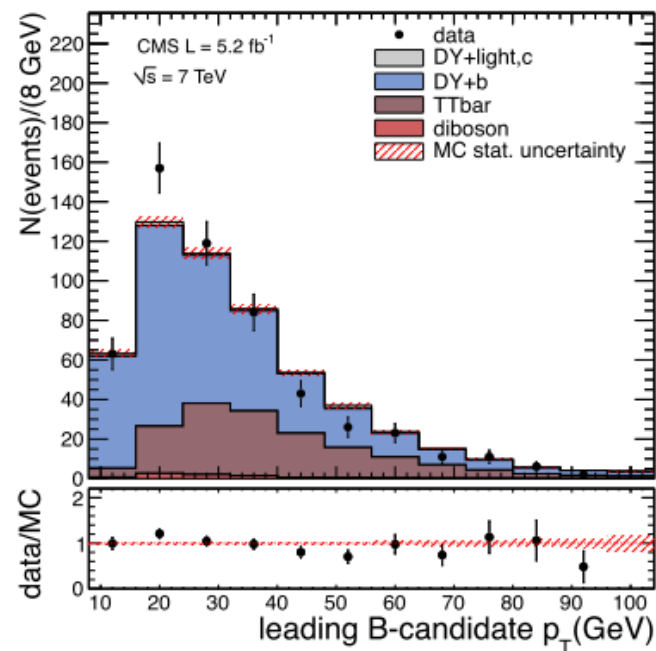
Starting to see the shape in the gluon-splitting dominated regime

(note that the theory comparisons here are rather dated)

CMS Z+b(b) at 7 TeV



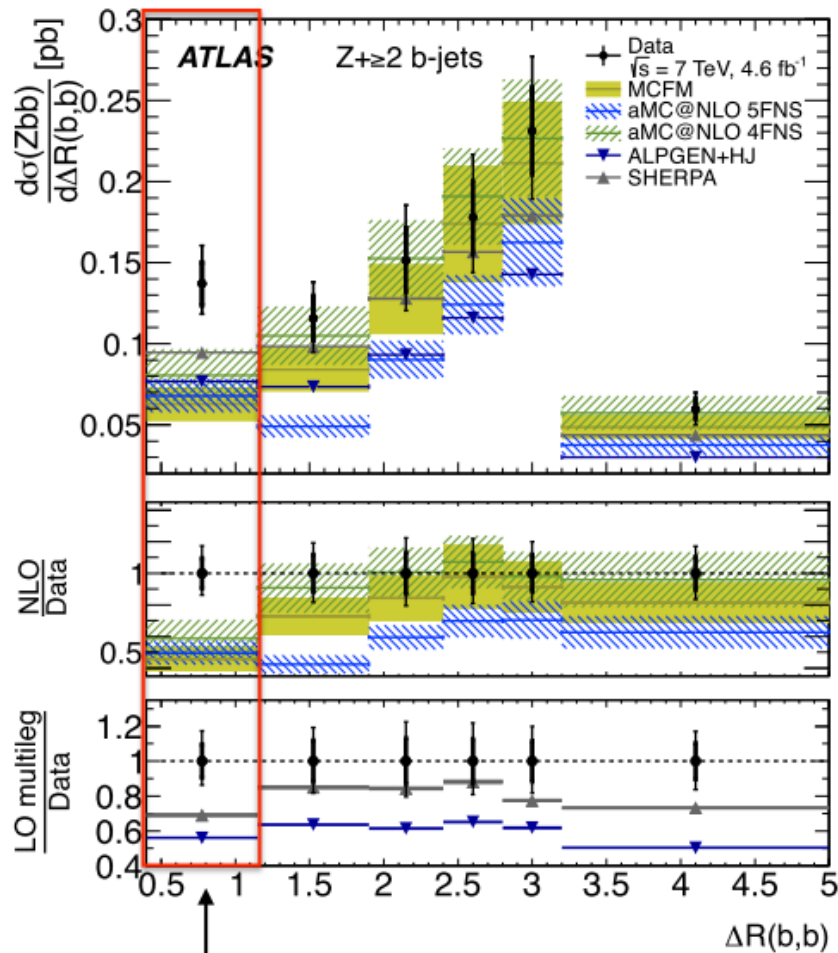
Dominated by
gluon splitting



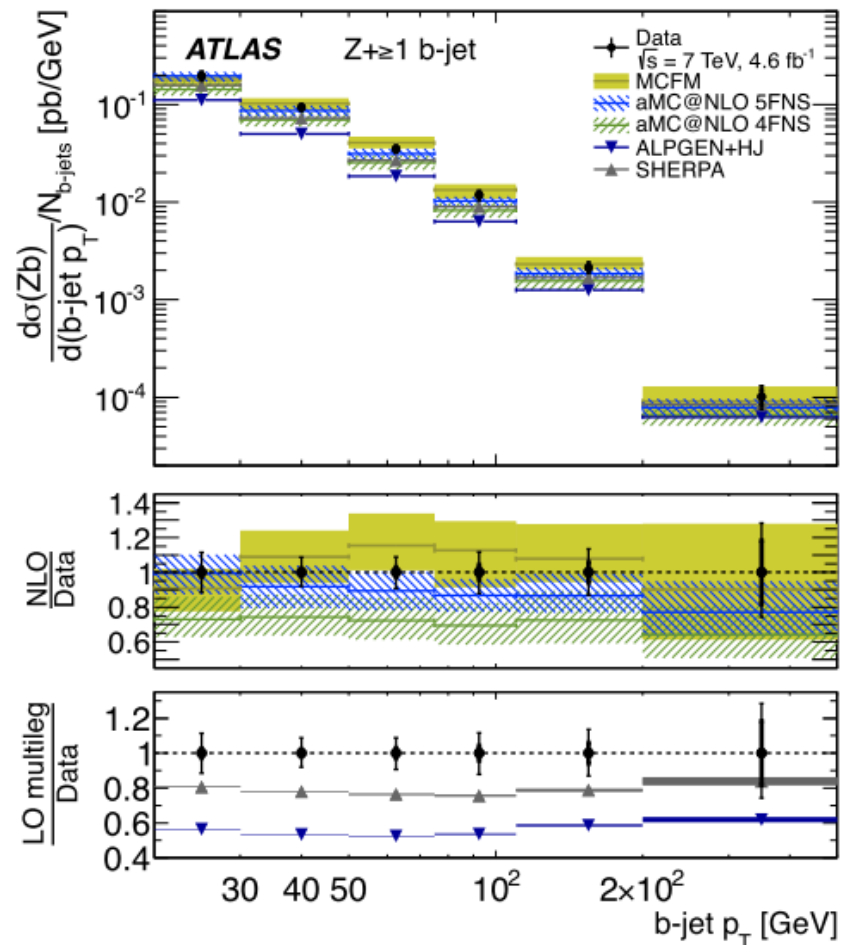
cfr Inclusive cross-section with two b-jets

<https://arxiv.org/abs/1402.1521>

ATLAS Z+b(b) at 7 TeV

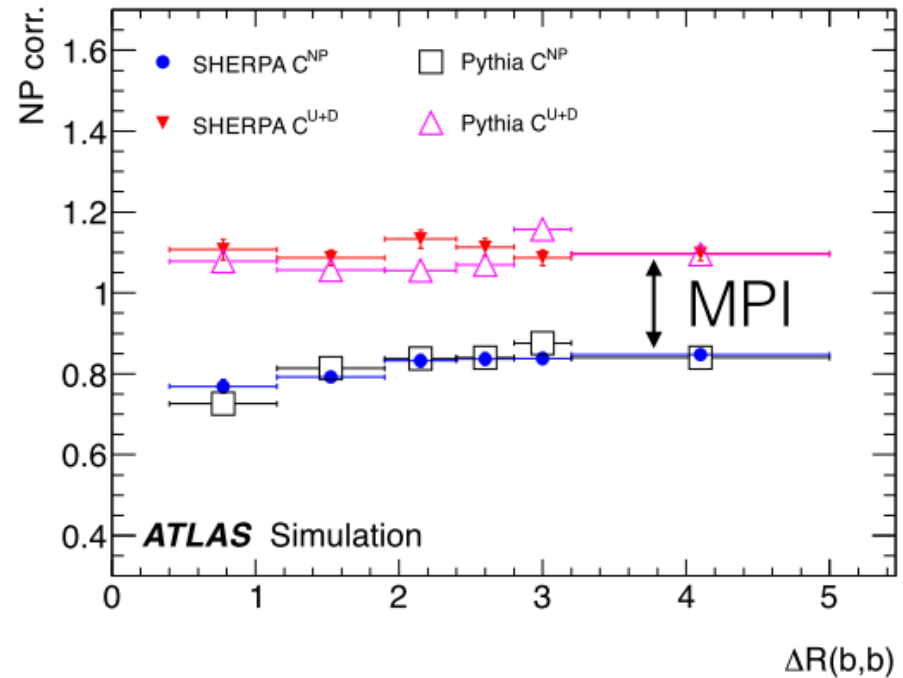
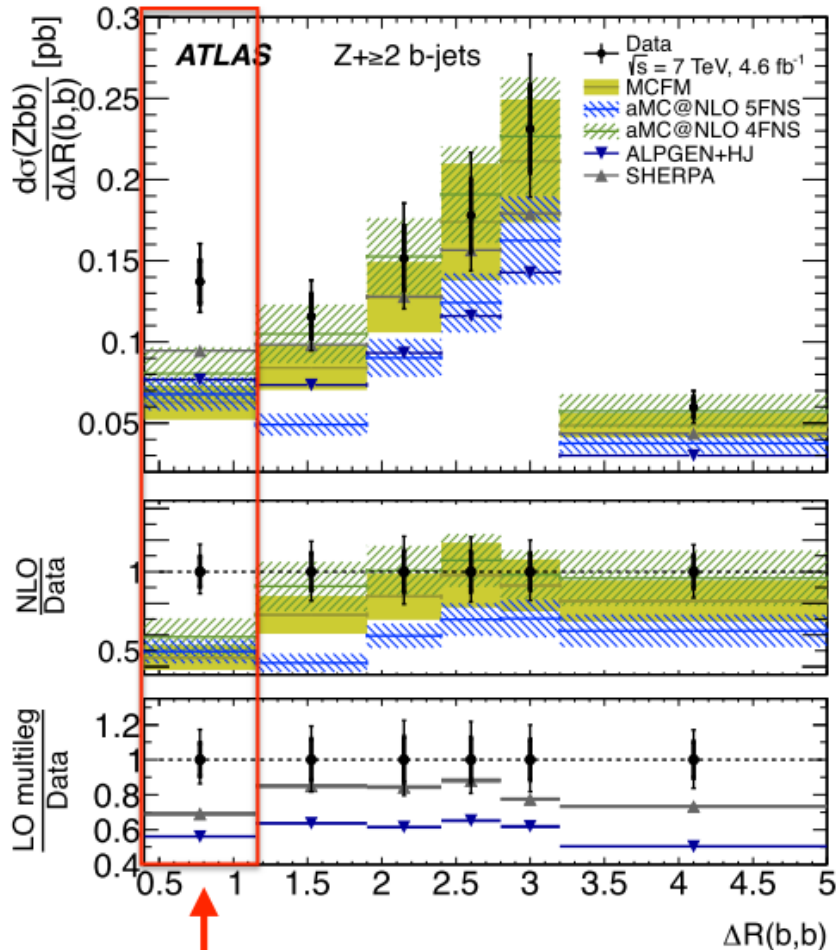


Approaching the PS phase space



Relatively low p_T regime

ATLAS Z+b(b) at 7 TeV



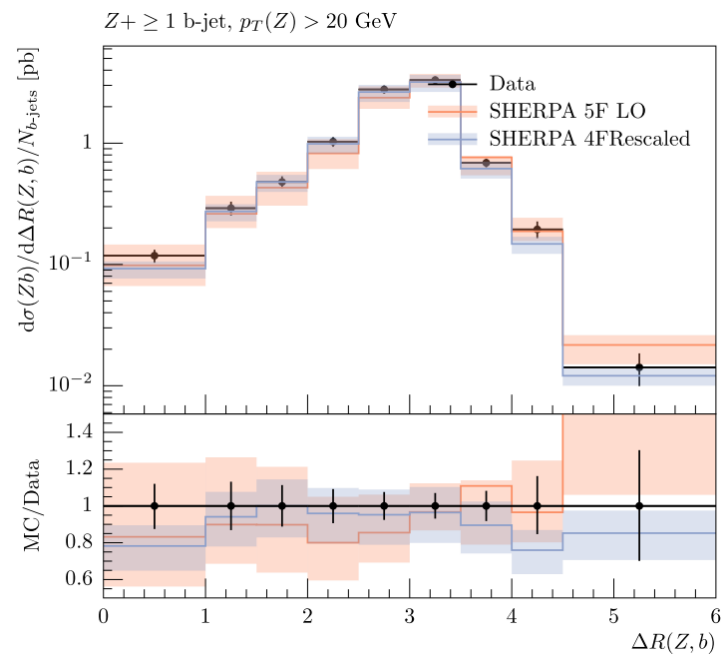
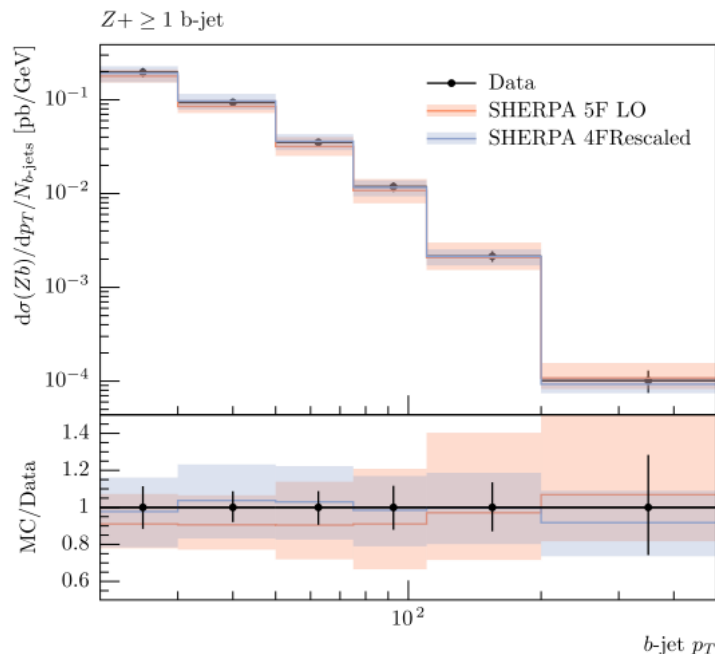
Interestingly, MPI is a relatively big effect here.

LH15 exercise on V+bb

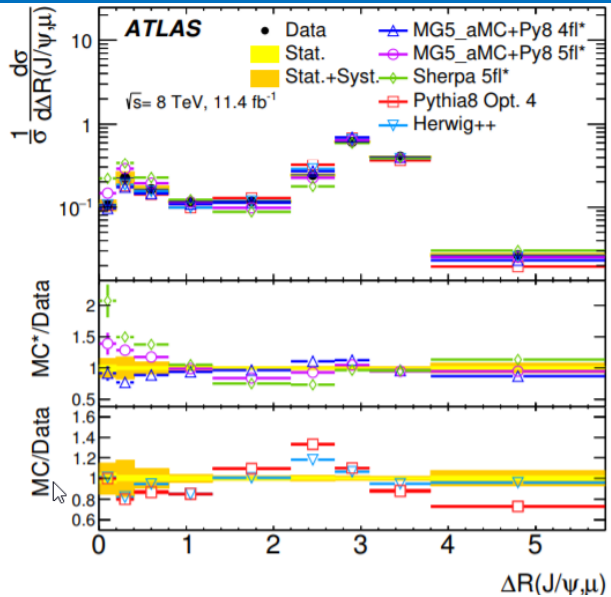
6 Study of associated production of vector bosons and b-jets at the LHC

6.6 Conclusions

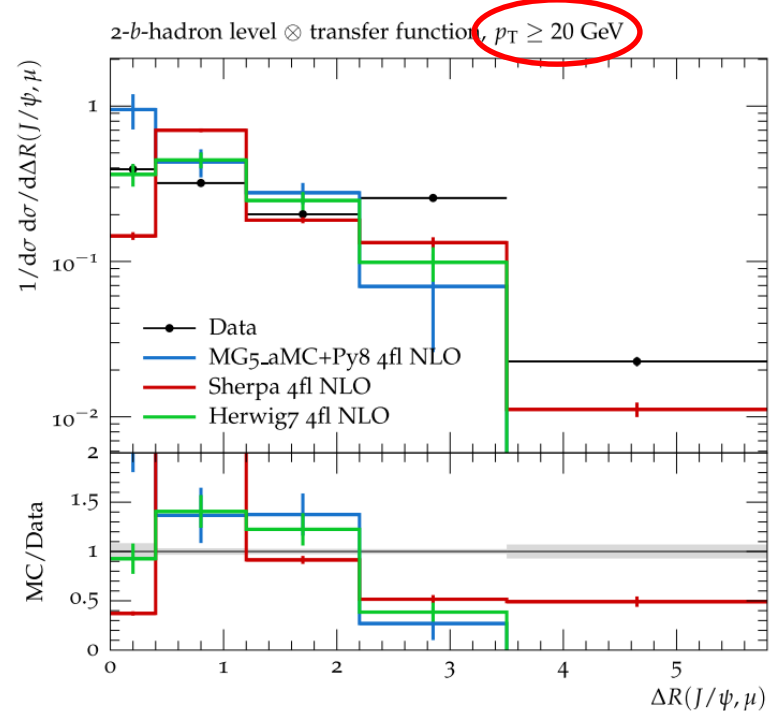
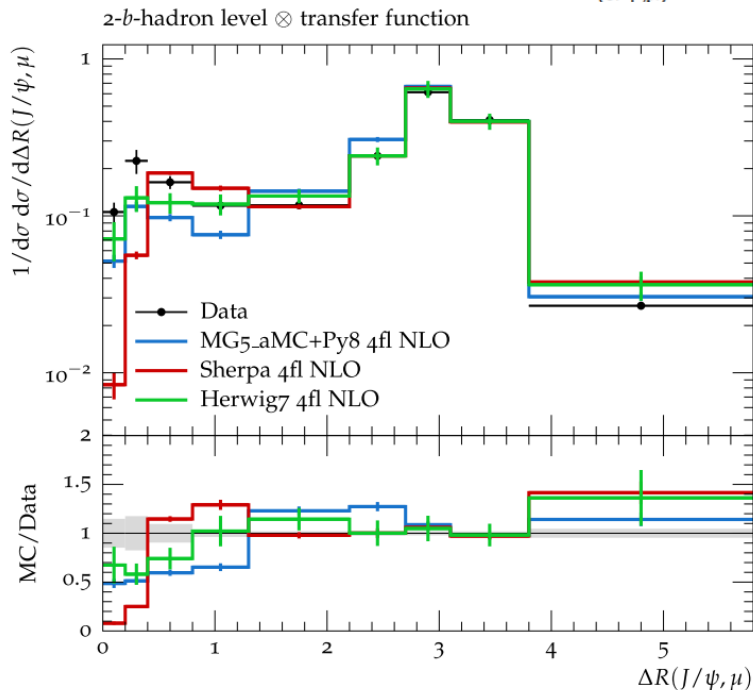
We presented a comparison of generators predictions using 4F and 5F scheme to most recent measurements of vector boson production in association with b-jets at the LHC. In the 4F scheme a good agreement is found among the different generators at NLO accuracy, and among different matrix-element to parton-shower matching algorithms. The agreement with data however is good only when two b-jets are tagged in the final state or, when one b-jet only is required, if a rescaling to the 5F integrated cross-section is applied. For Wbb, even taking into account the contribution from MPI, predictions seem to significantly undershoot the data. The Zb(b) production has been compared with predictions obtained in the 5F scheme with different setups,



LH17 exercise on hadron correlations



- Using the ATLAS di-hadron correlation measurement (<https://arxiv.org/abs/1705.03374>) in the channel $B(\rightarrow J/\psi[\rightarrow \mu\mu] + X)B(\rightarrow \mu + X)$
 - Go down to $\Delta R \sim 0$ but need to unfold fragmentation



Where is this going?

- V+bb background modeling strategies for VH(bb)

CMS

- V+(light-flavor) modeling**
CRs defined by inverting b-tagging requirements (anti-2-btag)
- Data driven pre-fit systematics
- V+(heavy-flavor) modeling**
CRs defined by inverting M(jj)-window

(b-tag CMVA_{min} shape fitted from CRs)

Process	0-lepton	1-lepton	2-lepton low- $p_T(V)$	2-lepton high- $p_T(V)$
W0b	1.14 ± 0.07	1.14 ± 0.07	—	—
W1b	1.66 ± 0.12	1.66 ± 0.12	—	—
W2b	1.49 ± 0.12	1.49 ± 0.12	—	—
Z0b	1.03 ± 0.07	—	1.01 ± 0.06	1.02 ± 0.06
Z1b	1.28 ± 0.17	—	0.98 ± 0.06	1.02 ± 0.11
Z2b	1.61 ± 0.10	—	1.09 ± 0.07	1.28 ± 0.09
tt	0.78 ± 0.05	0.91 ± 0.03	1.00 ± 0.03	1.04 ± 0.05

ATLAS

- V+(heavy-flavor) modeling**
W: dedicated CR (large m-top, low m-bb) - yield only, no shape
Z: no *dedicated* CR - full m-bb spectrum included in the SRs
Theory driven pre-fit systematics
 $V+hf = V+(bb, bc, bl, cc)$

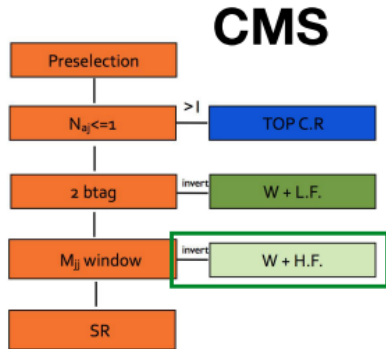
Process	Normalisation factor
$\bar{t}t$ 0- and 1-lepton	0.90 ± 0.08
$\bar{t}t$ 2-lepton 2-jet	0.97 ± 0.09
$\bar{t}t$ 2-lepton 3-jet	1.04 ± 0.06
W + HF 2-jet	1.22 ± 0.14
W + HF 3-jet	1.27 ± 0.14
Z + HF 2-jet	1.30 ± 0.10
Z + HF 3-jet	1.22 ± 0.09

Background reweighting corrections for V+jets:

- $f(p_T^V)$ differential correction (up to 10% at 400GeV) accounting for EW corrections
- $f(p_T^V)$ dedicated 1-lepton correction on W+light, W+b(b), ttbar, single-t
- $\Delta\eta(jj)$ correction from LO/NLO comparison (depending on #b-labeled jets)

Where is this going?

- More detailed example: W+b(b) background modeling for VH(bb)

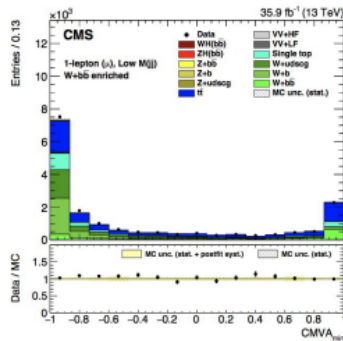


ATLAS

- standard 1-lepton selection + $m(bb) < 75\text{GeV}$
 $m(\text{top}) > 225\text{GeV}$
- Scale factor fitted directly in the SR
- extrapolation uncertainties from CR to SR obtained from theory
 - Sherpa 2.2.1 muR, muF, ckw, qsf scale variations
 - Sherpa 2.2.1 comparison with Madgraph_aMC@NLO 2.2.2

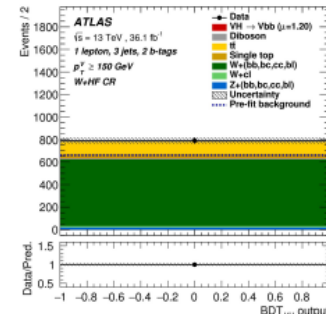
- Define dedicated control region (CR)
- Scale factors applied from CR to Signal Regions (SR)
- Systematic uncertainties fully correlated between CR and SR

Variable	W+HF
$p_T(j_1)$	>25
$p_T(j_2)$	>25
$p_T(j_j)$	>100
$p_T(V)$	>100
$CMVA_{\text{max}}$	$>CMVA_T$
N_{aj}	$=0$
$N_{a\ell}$	$=0$
$\sigma(p_T^{\text{miss}})$	>2.0
$\Delta\phi(\vec{p}_T^{\text{miss}}, \ell)$	<2
$M(jj)$	$<90, [150, 250]$



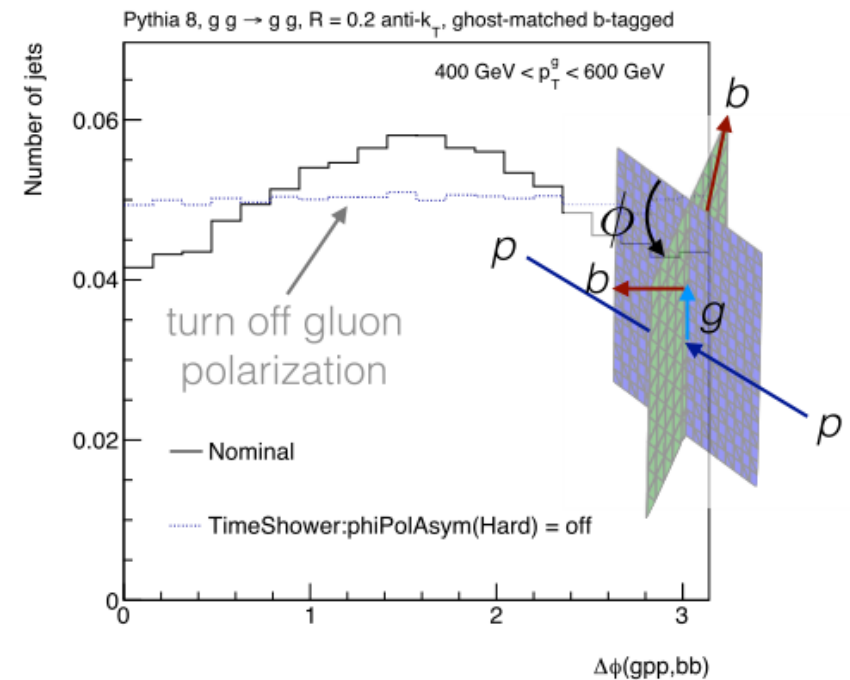
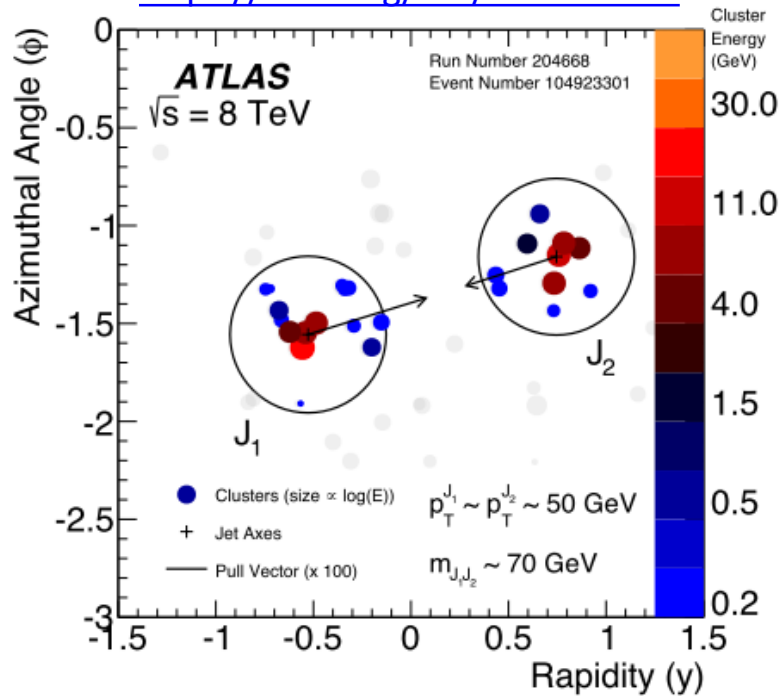
- Pre-fit theory modeling uncertainties

W + jets	
W + ll normalisation	32%
W + cl normalisation	37%
W + bb normalisation	Floating (2-jet, 3-jet)
W + bl-to-W + bb ratio	26% (0-lepton) and 23% (1-lepton)
W + bc-to-W + bb ratio	15% (0-lepton) and 30% (1-lepton)
W + cc-to-W + bb ratio	10% (0-lepton) and 30% (1-lepton)
0-to-1 lepton ratio	5%
W + HF CR to SR ratio	10% (1-lepton)
m_{bb}, p_T^V	S



What else can and should we measure?

<https://arxiv.org/abs/1506.05629>

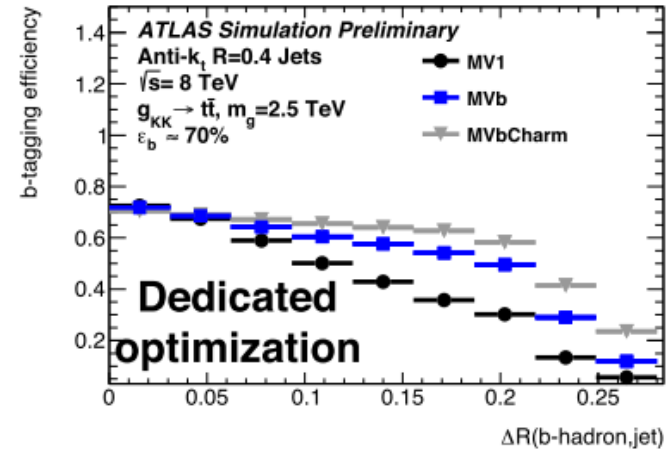
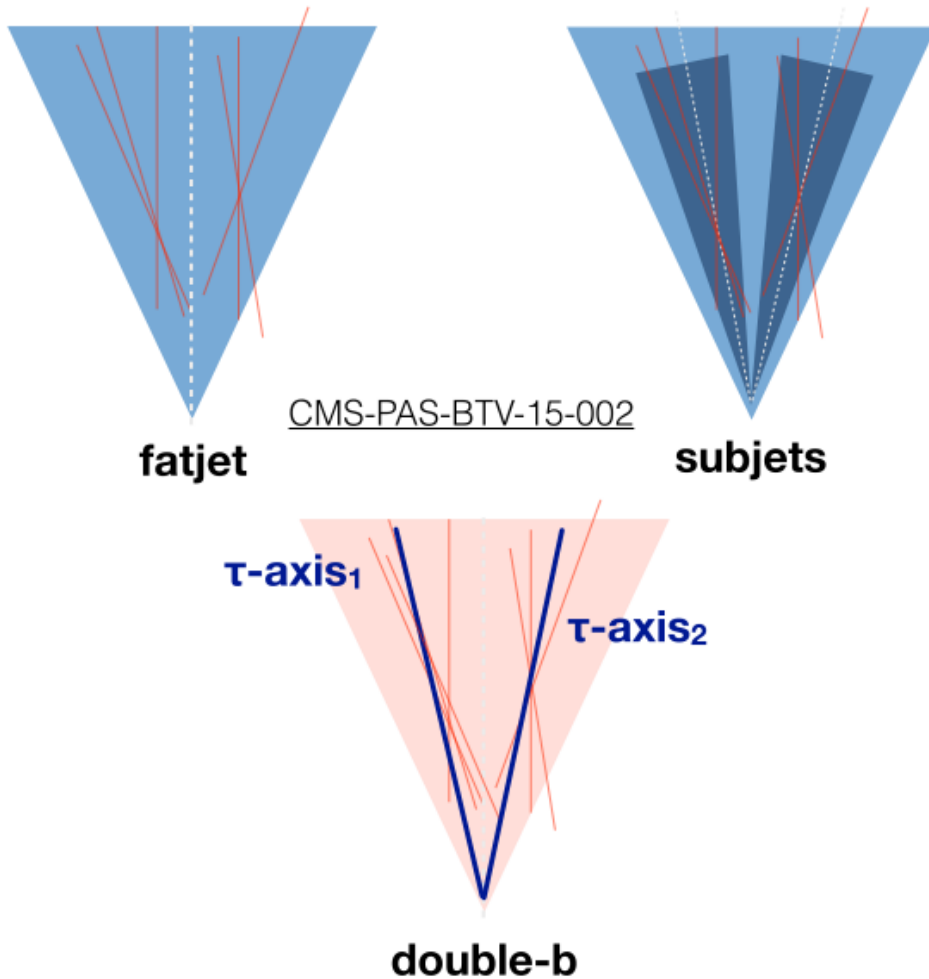


The singlet radiation pattern has been measured in W decays - should measure octet in $g \rightarrow bb$!

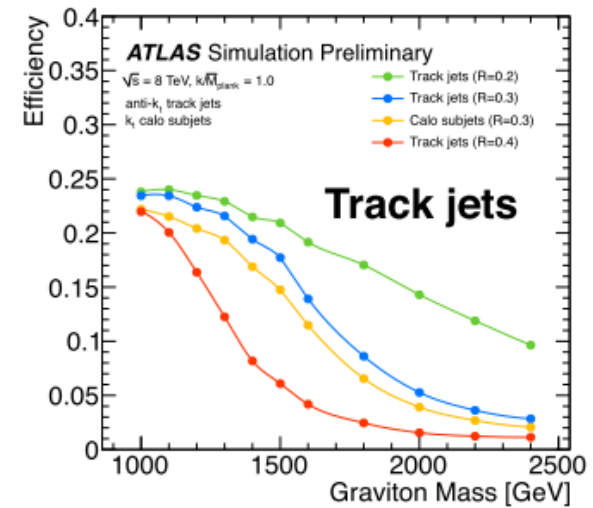
We should measure all aspects of the $g \rightarrow bb$ production (angles + energies)

What have we been up to since early Run 1?

Since 7 TeV, there has been a lot of work to improve b-tagging inside jets and to measure the efficiency in data.



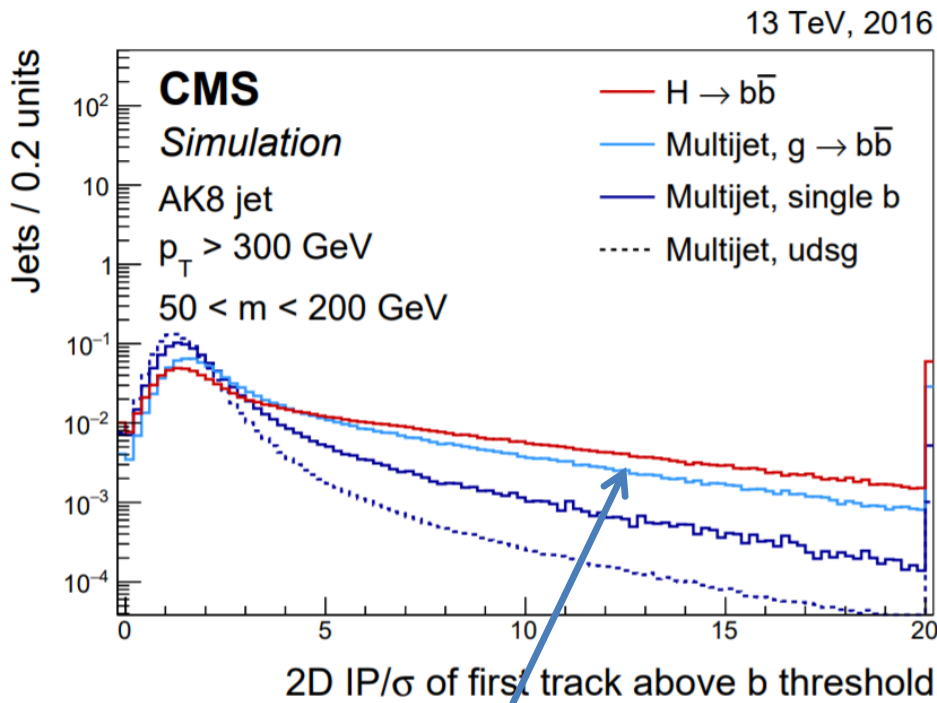
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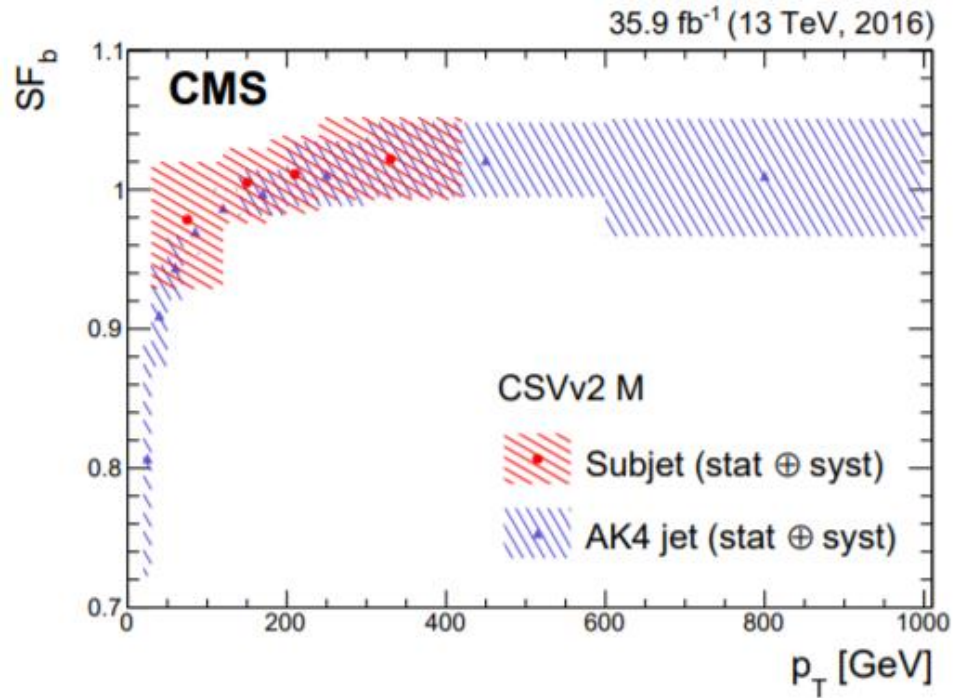
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CMS $g \rightarrow bb$ performance

<https://arxiv.org/abs/1712.07158>



Higgs and $g \rightarrow bb$
are very similar!



$\sim 10\%$ uncertainty;
stats limited

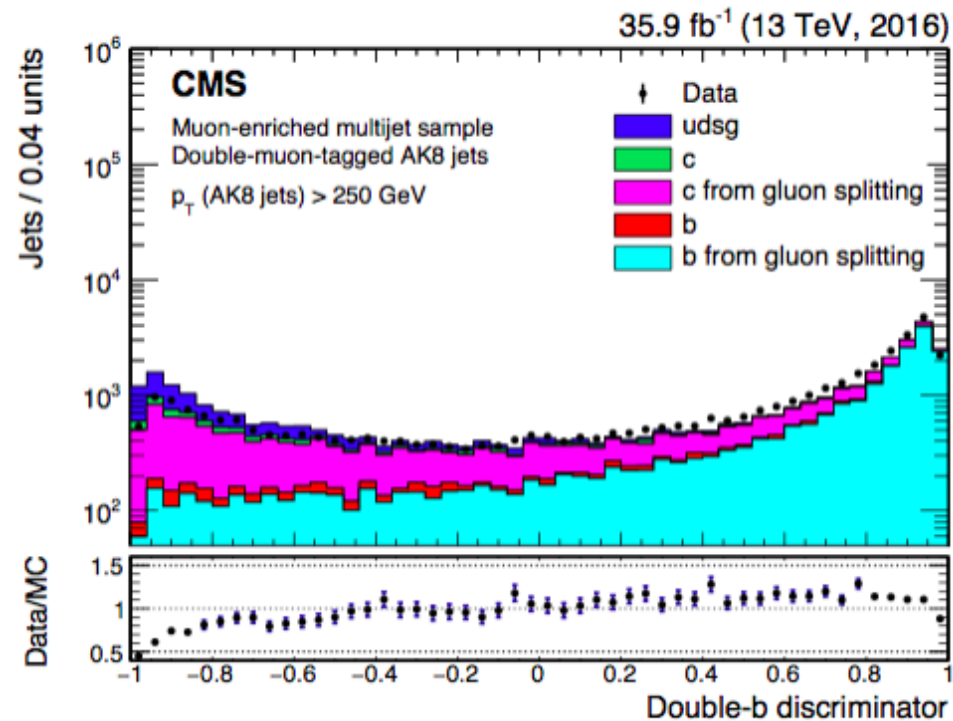
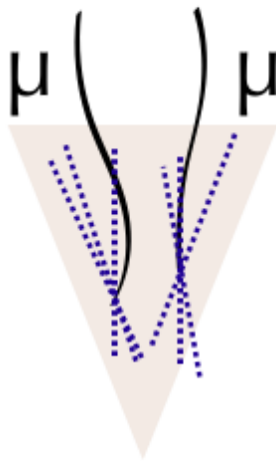
[8 TeV analysis](#)

Refining the calibration with more data

<https://arxiv.org/abs/1712.07158>

Jet selection has been designed to ensure jets are signal-like

- High AK8 p_T jet ($p_T > 250$ GeV)
- **double-muon** tagged jets (muon $p_T > 7$ GeV)
- **mass cut** (>50 GeV)



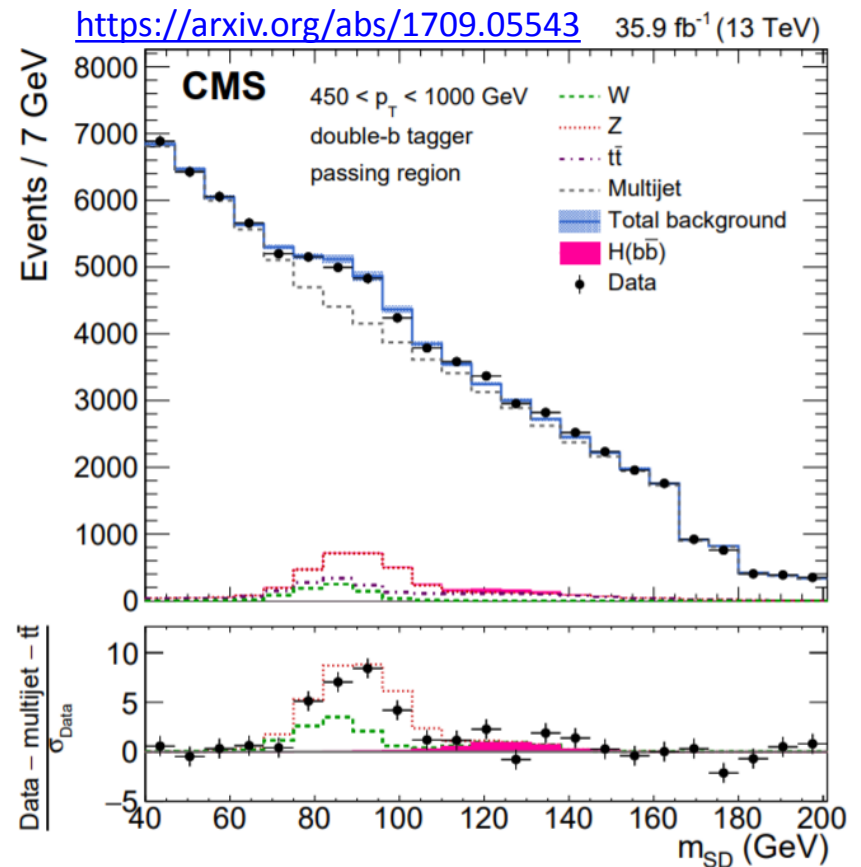
Where is this going?

The Run 2 efforts are motivated by (boosted) Higgs tagging.

For many of these searches, $g \rightarrow b\bar{b}$ is the main background.

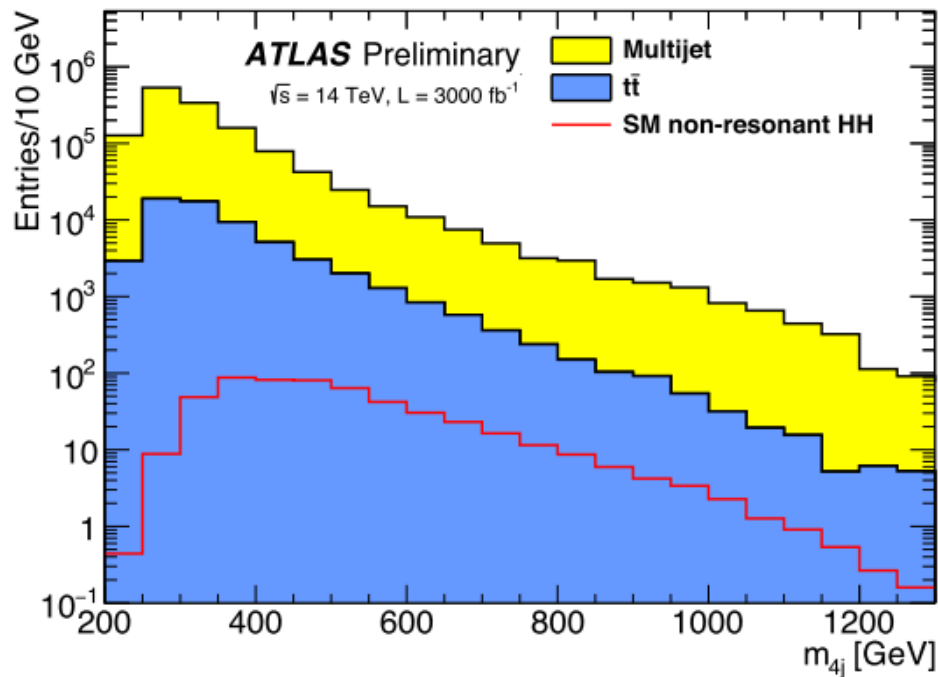
Data-driven techniques are used because the MC is not reliable (still needed to check closure)

Is this something that could change in the next 10 years?



<https://twiki.cern.ch/twiki/bin/view/CMSPublic/HighPtTrackingDP>

HH @ HL-LHC



<https://cds.cern.ch/record/2221658/>

For example, one of the most challenging and important measurements is the Higgs self-coupling.

The $g \rightarrow bb$ background is complicated, but maybe a better understanding could be a game-changer here!

>> CMS HH projections

<http://cms-results.web.cern.ch/cms-results/public-results/preliminary-results/FTR-15-002/>

Source	$\Delta\mu$
Luminosity	0.05
Jet Energy	0.09
b -tagging	0.34
Theoretical	0.10
Multijet	1.85
$t\bar{t}$	2.83

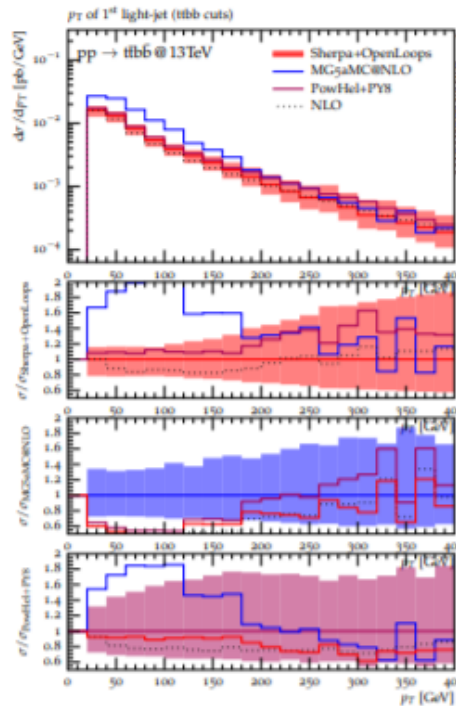
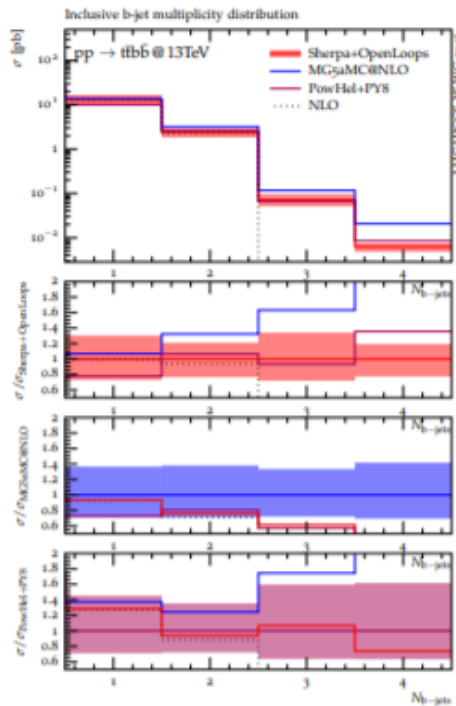
What about $t\bar{t}H(bb)$?

- after YR4 we decided to focus $t\bar{t}H/tH$ subgroup activities on **highest priority TH issues in EXP analyses**
- in the recent months we focussed on **TH uncertainties of $t\bar{t} + b$ -jet background, which dominate $t\bar{t}H(bb)$ systematics**
- $t\bar{t} + b$ -jet data help, but precise “extrapolation” to signal region calls for **$t\bar{t} + b$ -jet shape uncertainties at 10% level**
- $pp \rightarrow t\bar{t}b\bar{b}$ remains a nontrivial **multi-particle multi-scale QCD process**
- **better understanding of its QCD dynamics and NLOPS technicalities crucial for assessment of TH uncertainties**

<https://indico.cern.ch/event/407347/contributions/975965/attachments/1211342/1766869/hxswg16.pdf>

What about ttH(bb)?

YR4 highlights: Validation of NLO+PS tools, $t\bar{t} + b$ jets



↪ switch off top decays, hadronization, UE

↪ To better compare the effect of

- different matchings
- different parton showers
- different flavor scheme

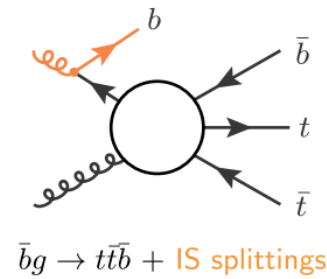
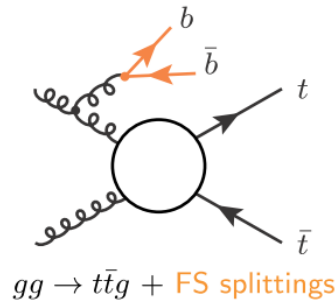
Discrepancies emerge that will have to be understood if we want to resolve the very large systematic uncertainties that affects experimental analyses

↪ **This is becoming a limiting factor.**

ttbb 5F at NLO

NLOPS ttbb 5F
(e.g. POWHEG hvq)

$t\bar{t}b\bar{b}$ described through $t\bar{t}j$ tree MEs plus $g \rightarrow b\bar{b}$ shower splittings



Precision vs accuracy

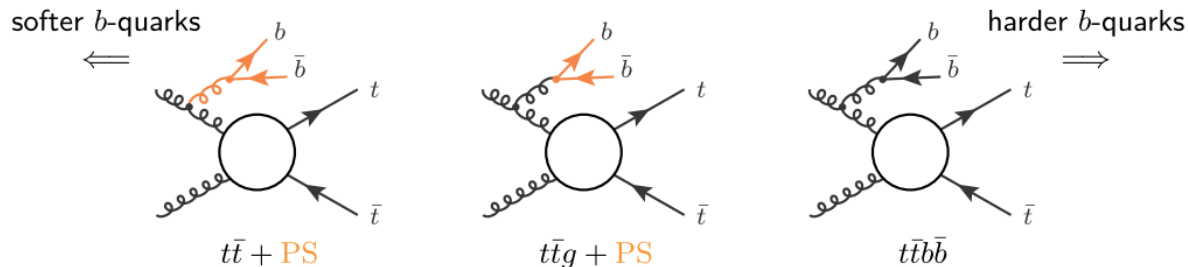
- precision lower than LO (parton shower allows for **accurate tuning to data**)

Calls for improved description based on $t\bar{t}b\bar{b}$ MEs

- crucial for **more realistic TH uncertainties**

(N)LO ttbar from merging
tt+0,1,2j 5F

$t\bar{t}b\bar{b}$ described through $t\bar{t} + 0, 1, 2 \text{ jet MEs}$ and $g \rightarrow b\bar{b}$ shower splittings

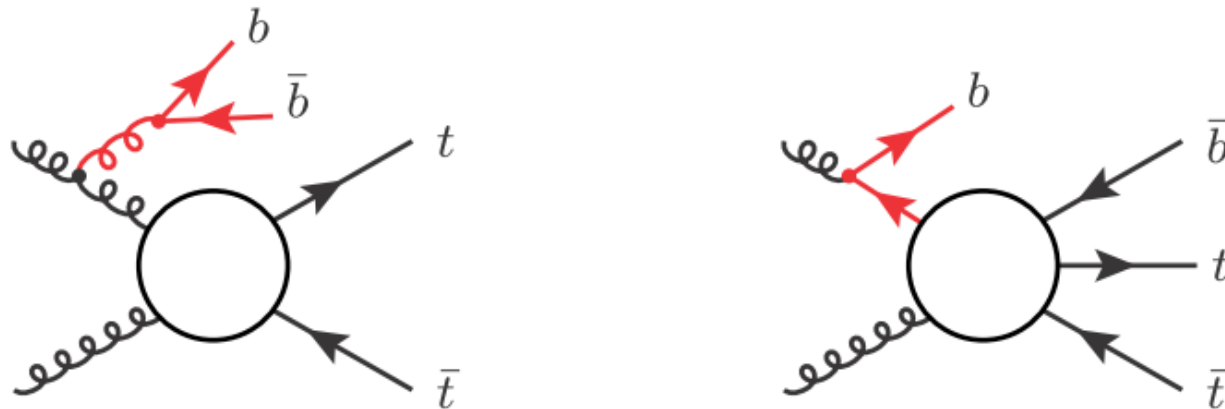


Precision and CPU cost strongly depend on choice of merging cut Q_{cut}

- separates ME regions ($k_T > Q_{\text{cut}}$) from shower regions ($k_T < Q_{\text{cut}}$)

Does this describe $t\bar{t}+b$ -jet production mostly through $t\bar{t}b\bar{b}$ MEs?

ttbb 4F at NLO



4F $t\bar{t}b\bar{b}$ MEs with $m_b > 0$ cover full b -quark phase space

- NLO precision for $t\bar{t} + 2 b$ -jet and 1 b -jet! [Cascioli et al '13]
- 80% LO uncertainty reduced to 20–30% at NLO
- collinear $g \rightarrow b\bar{b}$ splittings and m_b effects very important

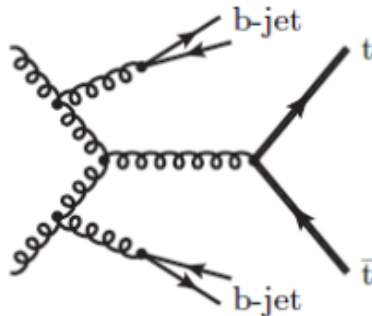
what about drawbacks of 4F scheme (e.g. no b -quark PDF)?

NLOPS $ttbb$ 4F SHERPA+OPENLOOPS

Convergence of 4F scheme but unexpected MC@NLO enhancement

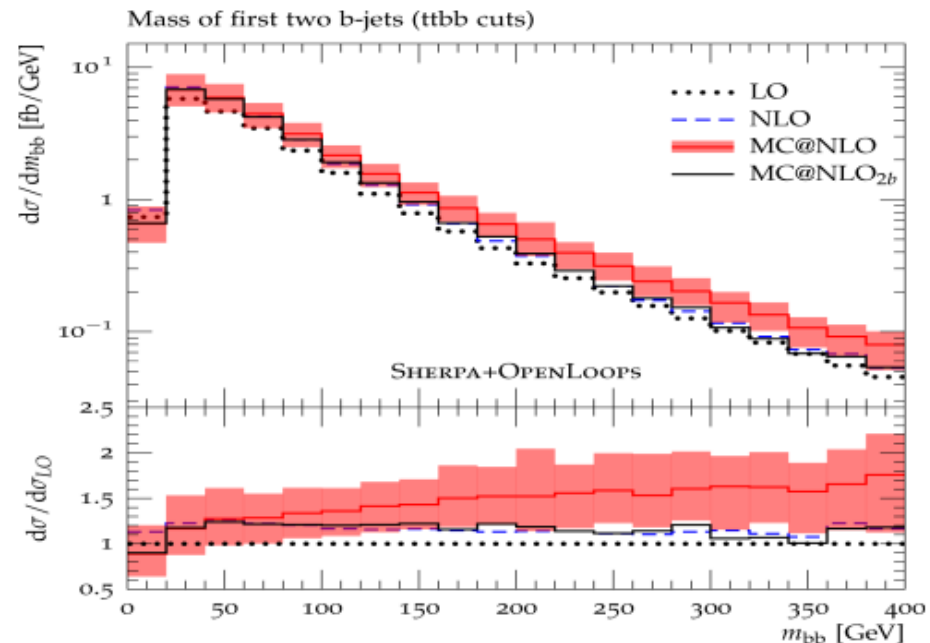
	ttb	$ttbb$	$ttbb (m_{bb} > 100)$
$\sigma_{\text{LO}} [\text{fb}]$	$2644^{+71\%+14\%}_{-38\%-11\%}$	$463.3^{+66\%+15\%}_{-36\%-12\%}$	$123.4^{+63\%+17\%}_{-35\%-13\%}$
$\sigma_{\text{NLO}} [\text{fb}]$	$3296^{+34\%+5.6\%}_{-25\%-4.2\%}$	$560^{+29\%+5.4\%}_{-24\%-4.8\%}$	$141.8^{+26\%+6.5\%}_{-22\%-4.6\%}$
$\sigma_{\text{NLO}}/\sigma_{\text{LO}}$	1.25	1.21	1.15
$\sigma_{\text{MC@NLO}} [\text{fb}]$	$3313^{+32\%+3.9\%}_{-25\%-2.9\%}$	$600^{+24\%+2.0\%}_{-22\%-2.1\%}$	$181^{+20\%+8.1\%}_{-20\%-6.0\%}$
$\sigma_{\text{MC@NLO}}/\sigma_{\text{NLO}}$	1.01	1.07	1.28

Large enhancement ($\sim 30\%$) in Higgs region from double $g \rightarrow b\bar{b}$ splittings



One $g \rightarrow b\bar{b}$ splitting from PS

\Rightarrow TH uncertainties related to matching, shower and 4F/5F schemes crucial!



Conclusions

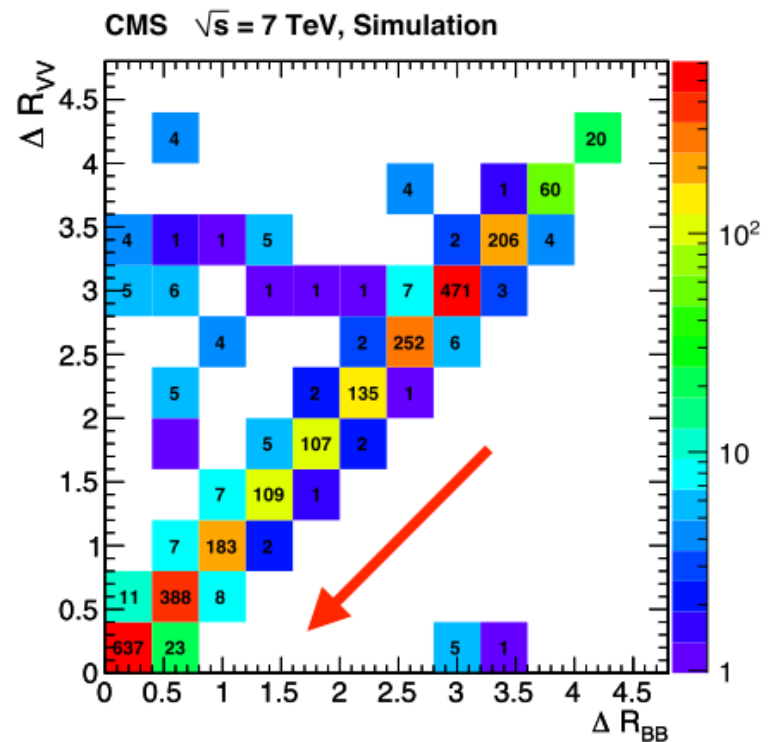
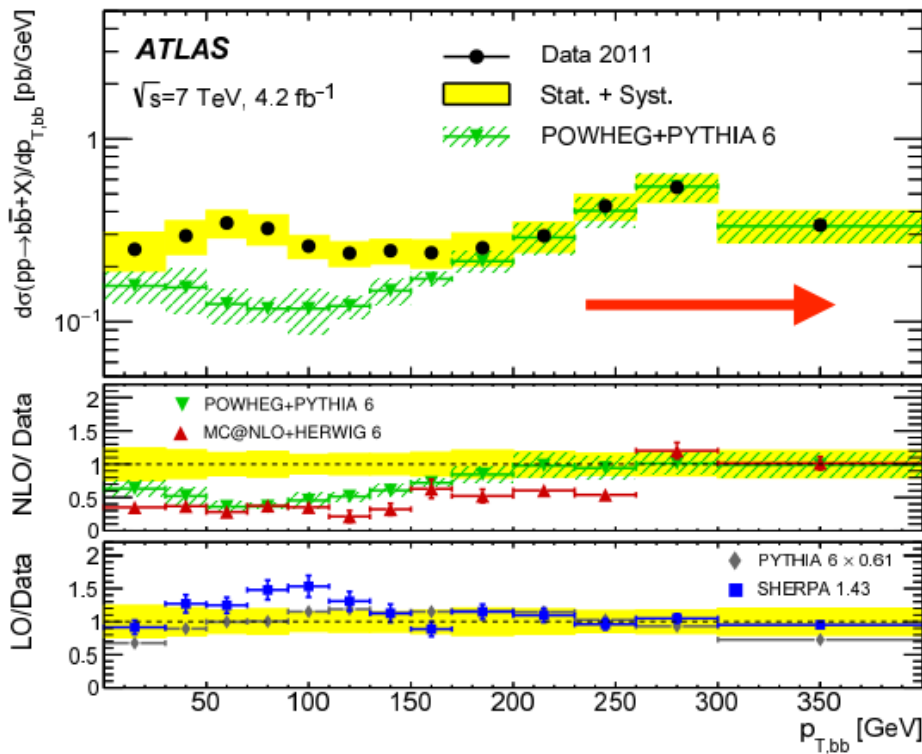
- **Usually, from the experimental perspective, gluon splitting is something you want to get rid of**
 - $VH(bb)$
 - $gg \rightarrow H(bb)$
 - $H(bb)H(bb)$
 - $ttH(bb)$
 - ...
- **But you must know very well**
 - Very abundantly produced at LHC
 - Enters in most of the analyses with final state b quarks
- **Very difficult to model theoretically**
 - Large uncertainties
- **Useful to calibrate b-tagging**
- **We should probably perform more SM measurements in different phase spaces, multi-differential etc**

Backup

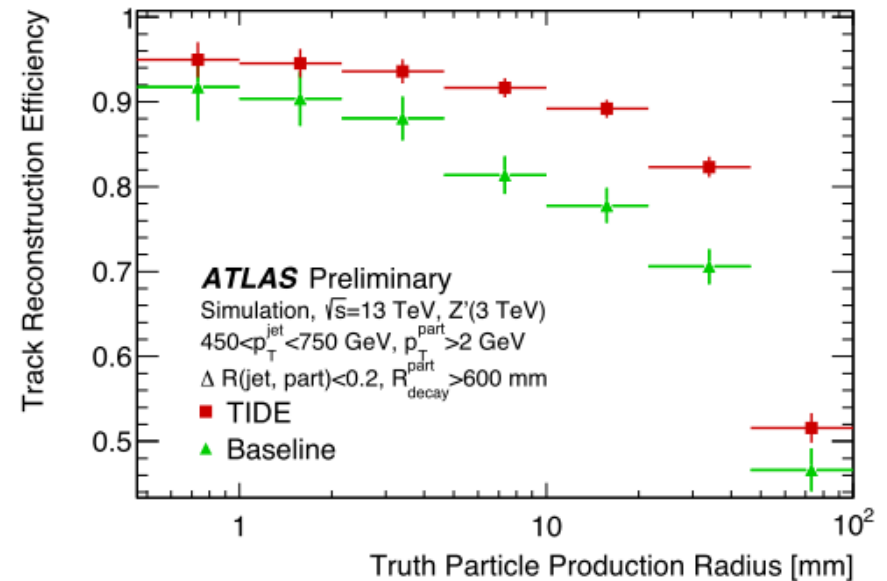
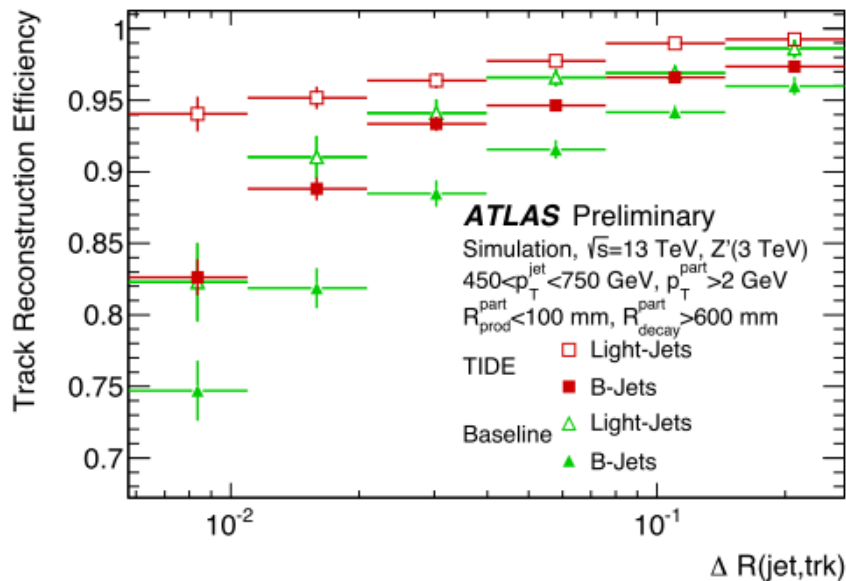
Partial list of references

- B. Nachman - Gluon splitting to bottom quarks at the LHC
 - Parton Radiation and Fragmentation from LHC to FCC-ee
<https://indico.cern.ch/event/557400/timetable/>
- J. A. McFayden, D. Napoletano, E. Re - Is there room for improvement in the description of the gluon splitting to heavy quarks?
 - Les Houches 2017: Physics at TeV Colliders Standard Model Working Group Report <https://arxiv.org/abs/1803.07977>
- R. M. Ralich - Study of b Quark Pair Production Mechanisms in pp Collisions with the CMS Experiment at LHC
 - PhD thesis <https://cdsweb.cern.ch/record/1311216>

What does it look like with more data, higher p_T , and state-of-the-art simulation?



Note: in order to maintain b-tagging performance, it is critical to have dedicated methods for tracking inside jets

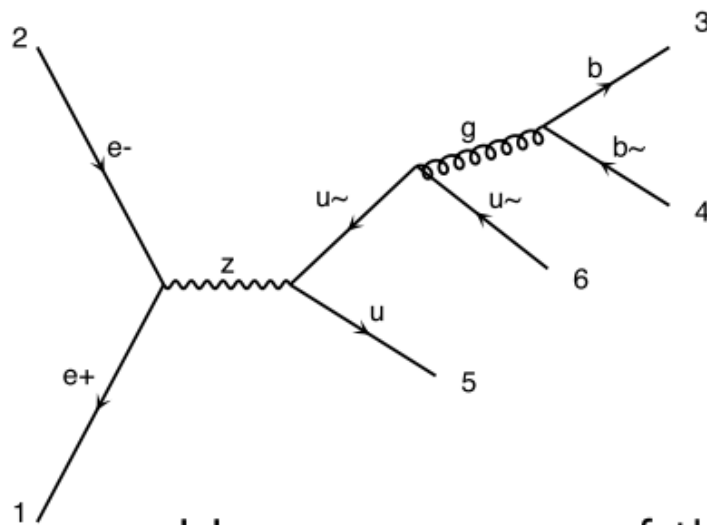


Algorithmic improvements are (much) cheaper than hardware ones - it is important to optimize performance when designing a new detector!

>> Similar studies in CMS

At a lepton collider, we would have an experimentally and theoretically clean environment for studying $g \rightarrow bb$

>> no pileup, UE, MPI, etc. <<

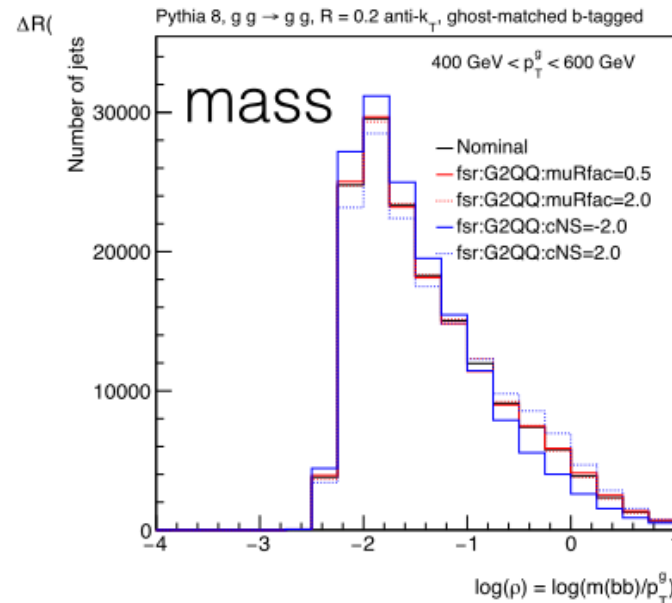
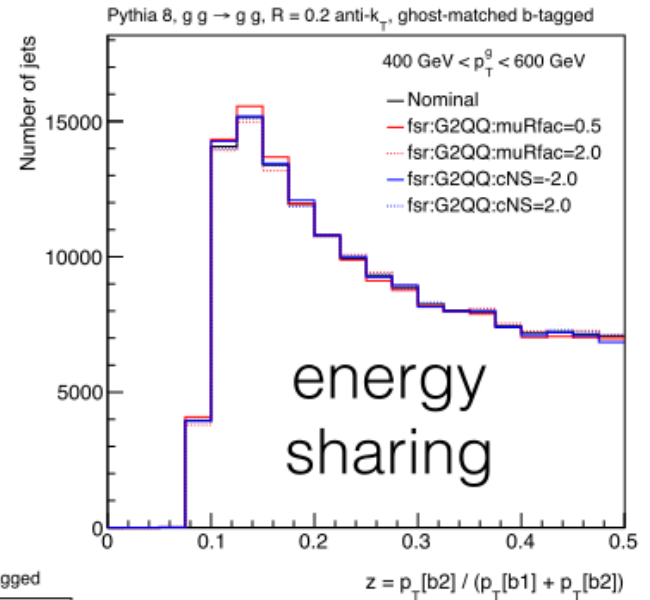
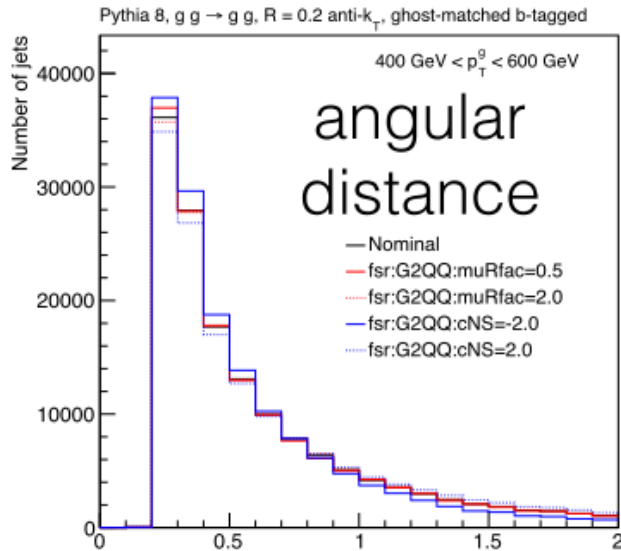


LEP measurements limited to inclusive rates of $g \rightarrow bb$

can probe pQCD to high precision by measuring properties of the splitting

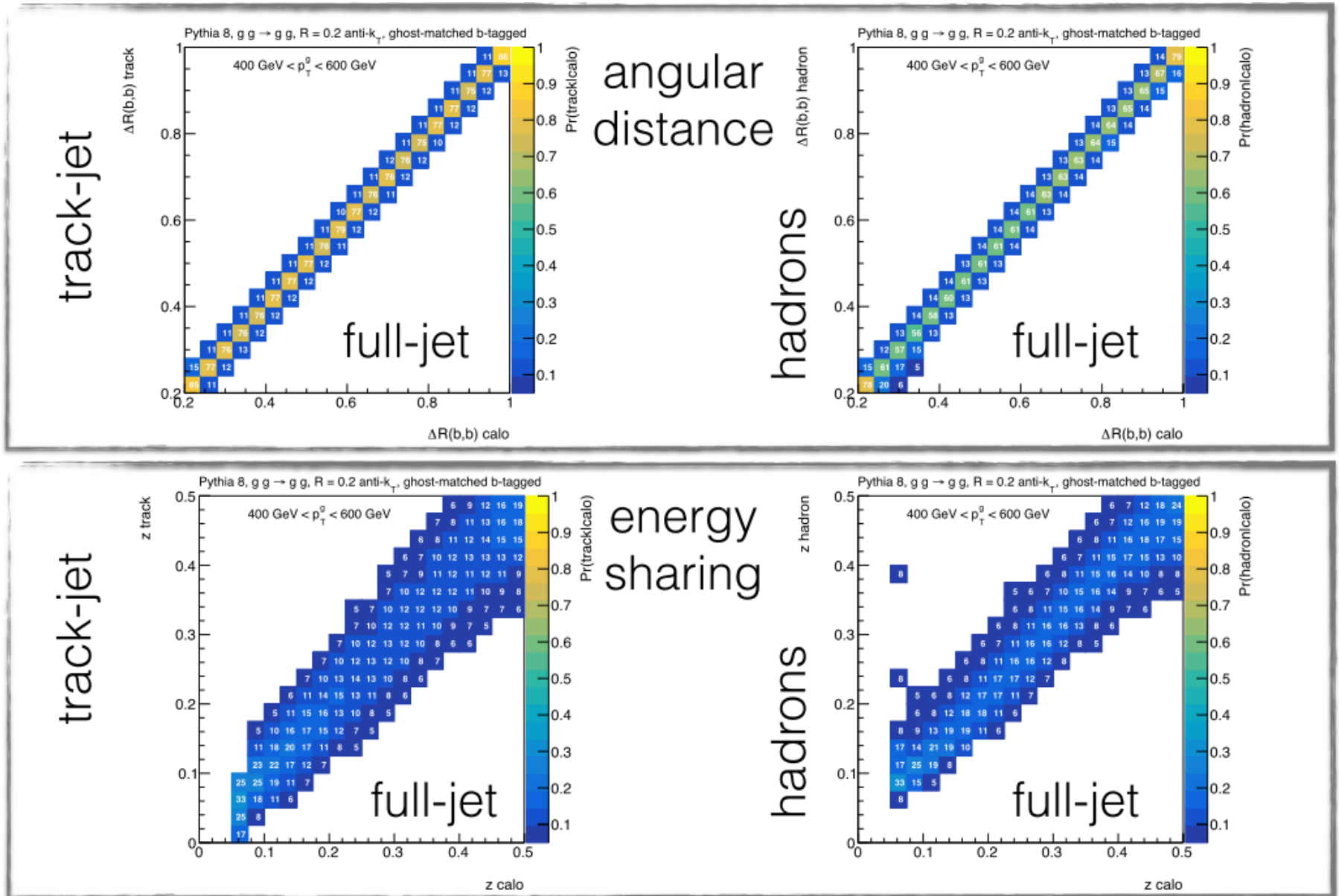
However, one of the most exciting prospects of FCC-ee is to perform a series of comparative measurements of $g \rightarrow bb$, $Z \rightarrow bb$, and $H \rightarrow bb$

What else can and should we measure?

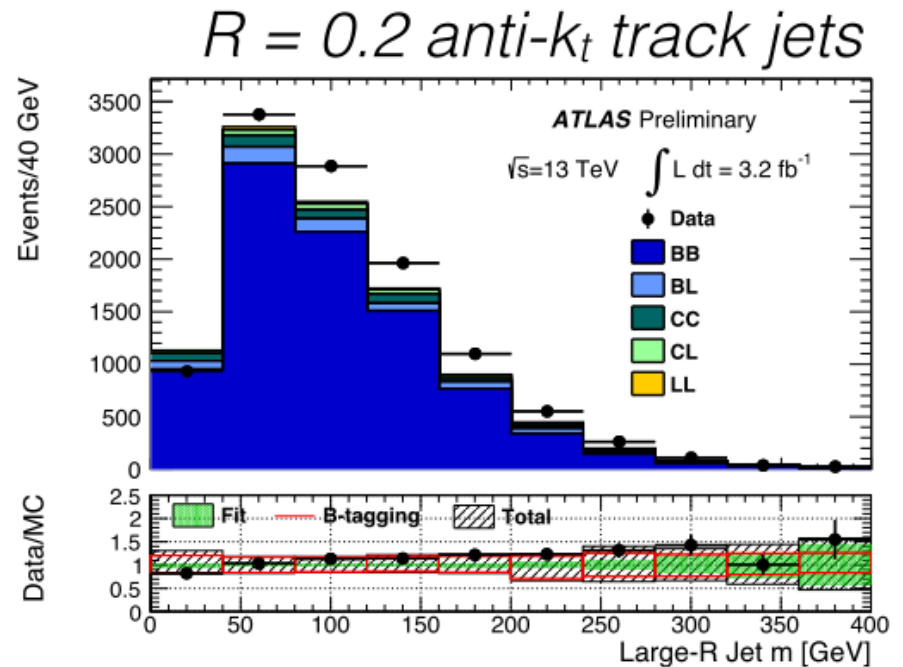
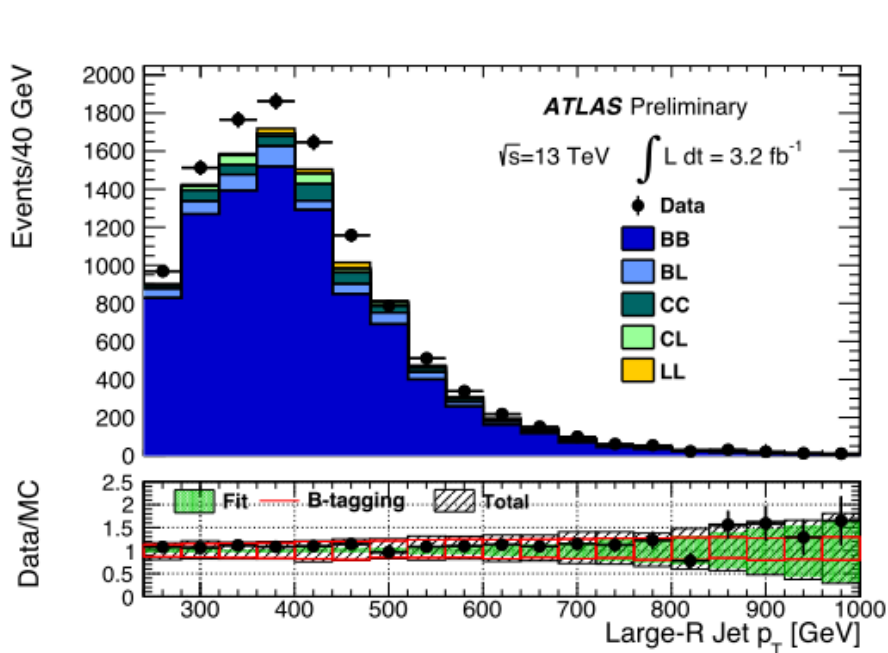


all the lines: tuning nobs in Pythia to see the sensitivity to the modeling of $g \rightarrow bb$

Caution: do we measure partons, hadrons, or (track) jets?



ATLAS $g \rightarrow bb$ performance



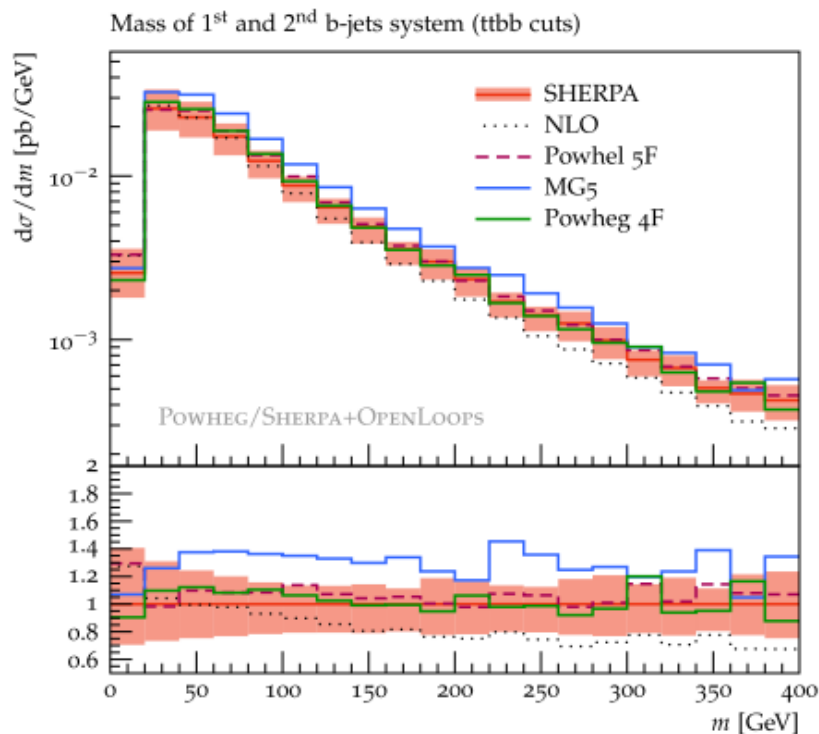
With the full Run 1 and Run 2 datasets, there are plenty of gluon jets for studying the modeling of double b-tagging
(in this case, use muons to increase purity)

[7 TeV double b-tagger](#)

[8 TeV analysis](#)

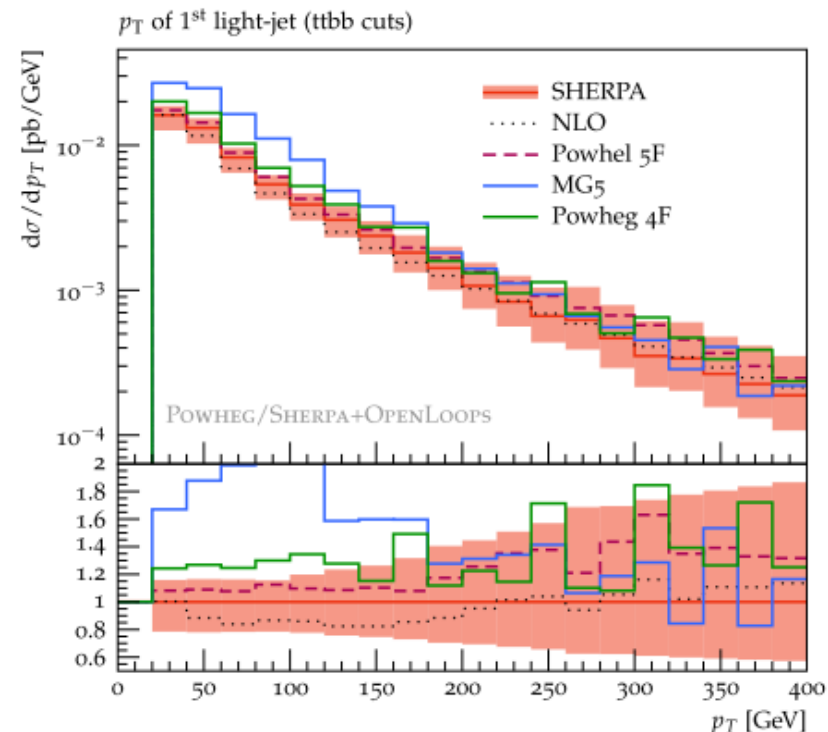
Further Powheg+Pythia8 results

m_{bb} with ttbb cuts



- “double $g \rightarrow b\bar{b}$ splittings” confirmed also by Powheg+PY8

p_{T,j_1} with ttbb cuts

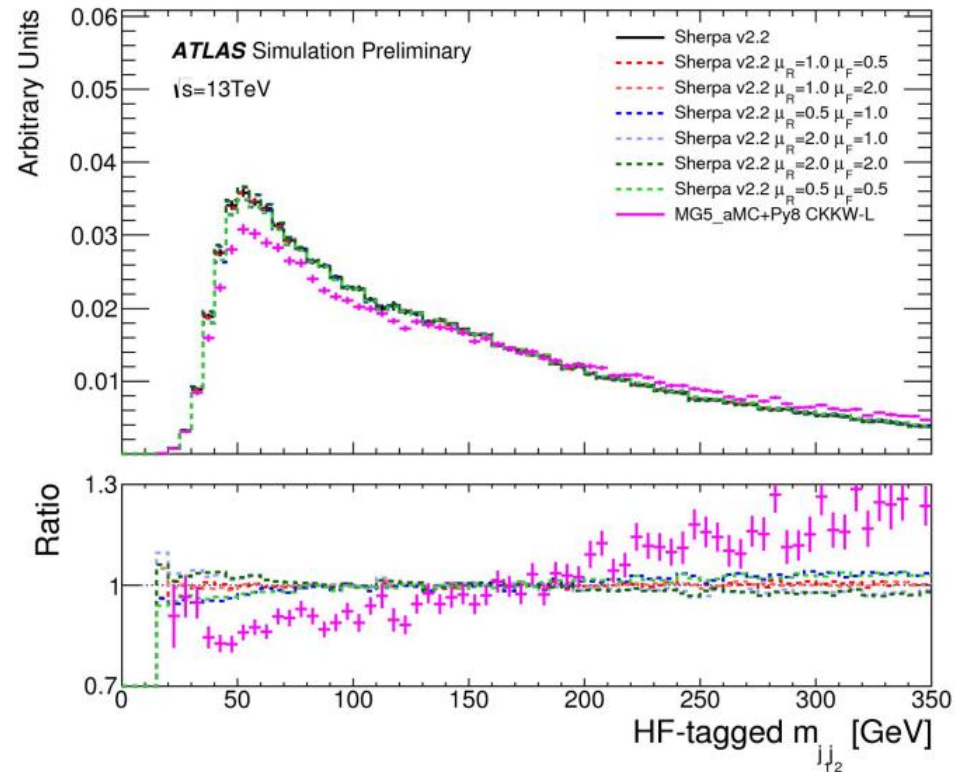


- Powheg+PY8 features enhancement in same direction as MG5+PY8
- but no strong distortion of spectrum

V+jets background modeling

[ATLAS PUB note](#) on V+jets modeling and MC simulation

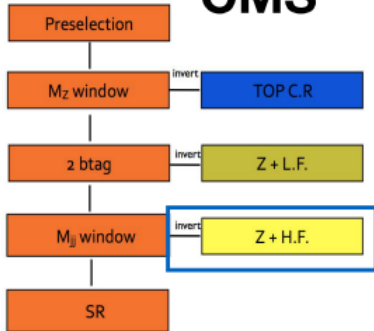
<https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/PUBNOTES/ATL-PHYS-PUB-2017-006/>



- **selection close to nominal VH(bb) analysis regions**
- no W+hf CR/SR separation

Z+heavy flavors (0-/2-lepton channel)

CMS

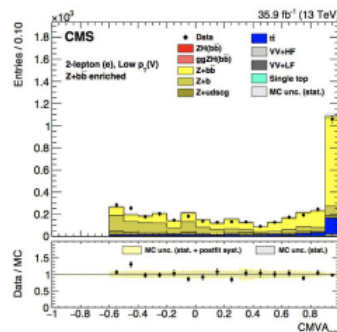


ATLAS

- no dedicated control region for Z+hf
- no $m(bb)$ window selection applied in the nominal analysis selection
- $m(bb)$ and p_{TV} shape systematic derived from data/MC in Z+hf enriched-region (2-lepton) \times (1-btag)
- (2-lepton) \times (2-btag) \times (remove events with $m(jj)$ around

- Define dedicated control region (CR)
- Scale factors applied from CR to Signal Regions (SR)
- Systematic uncertainties fully correlated between CR and SR

Variable	Z+HF
$p_T(jj)$	—
$p_T(V)$	[50, 150], >150
$CMVA_{max}$	> $CMVA_T$
$CMVA_{min}$	> $CMVA_L$
N_{Aj}	—
$N_{A\ell}$	—
p_T^{miss}	<60
$\Delta\phi(V, jj)$	>2.5
$M(\ell\ell)$	[85, 97]
$M(jj)$	\notin [90, 150]



- Pre-fit theory modeling uncertainties

Z + jets	
Z + ll normalisation	18%
Z + $c\ell$ normalisation	23%
Z + bb normalisation	Floating (2-jet, 3-jet)
Z + bc -to-Z + bb ratio	30 – 40%
Z + cc -to-Z + bb ratio	13 – 15%
Z + bl -to-Z + bb ratio	20 – 25%
0-to-2 lepton ratio	7%
m_{bb}, p_T^V	S

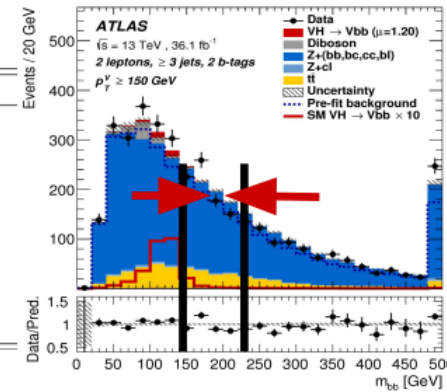


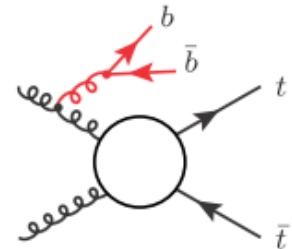
Table 8: The total numbers of events in each channel, for the rightmost 20% region of the event BDT output distribution, are shown for all background processes, for the SM Higgs boson VH signal, and for data. The yields from simulated samples are computed with adjustments to the shapes and normalizations of the BDT distributions given by the signal extraction fit. The signal-to-background ratio (S/B) is also shown.

Process	0-lepton	1-lepton	2-lepton low- p_T (V)	2-lepton high- p_T (V)
Vbb	216.8	102.5	617.5	113.9
Vb	31.8	20.0	141.1	17.2
V+udscg	10.2	9.8	58.4	4.1
$t\bar{t}$	34.7	98.0	157.7	3.2
Single top quark	11.8	44.6	2.3	0.0
VV(udscg)	0.5	1.5	6.6	0.5
VZ(bb)	9.9	6.9	22.9	3.8
Total background	315.7	283.3	1006.5	142.7
VH	38.3	33.5	33.7	22.1
Data	334	320	1030	179
S/B	0.12	0.12	0.033	0.15

Dominant topologies in 4F ttbb (FS vs IS)

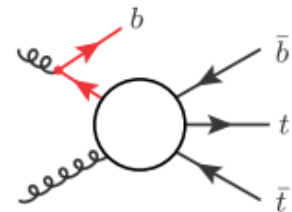
$t\bar{t}b\bar{b}$ topologies with FS $g \rightarrow b\bar{b}$ splittings

- dominant in full ttbb and ttb phase space
- notion of $g \rightarrow b\bar{b}$ splittings and IS/FS separation seems ill defined at large ΔR_{bb} , m_{bb} , $p_{T,b}$ due to sizable interferences

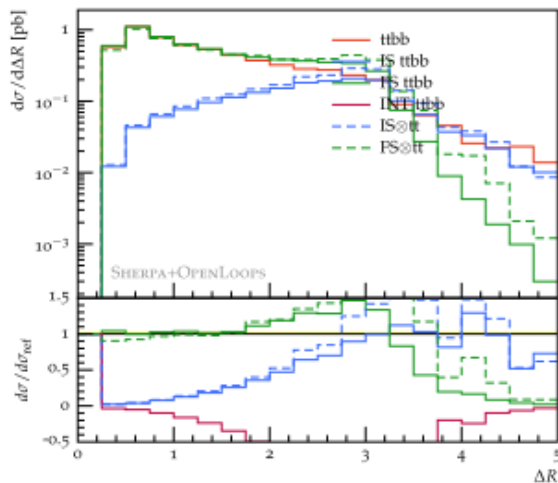


$t\bar{t}b\bar{b}$ topologies with IS $g \rightarrow b\bar{b}$ splittings

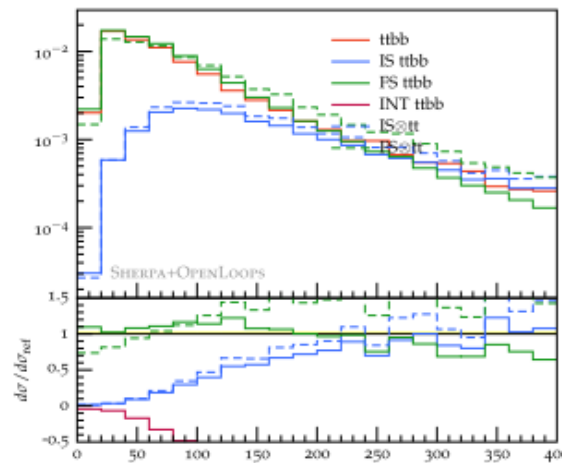
- mostly clearly subdominant (no need for 5F scheme resummation)



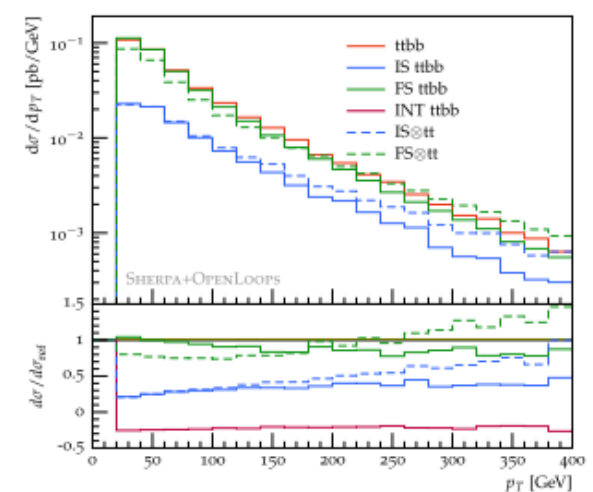
ΔR_{bb} with ttbb cut



m_{bb} with ttb cuts



$p_{T,b1}$ with ttb cuts



supports choice of 4F scheme with $m_b > 0$ and no b -quark PDF