

Probing QCD Wigner/GTMD distribution in diffractive di-jet production

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Outline

- Nucleon tomography
- Phase space distributions in QCD
- Connection to experimental observables

Electron-Ion Collider (EIC)

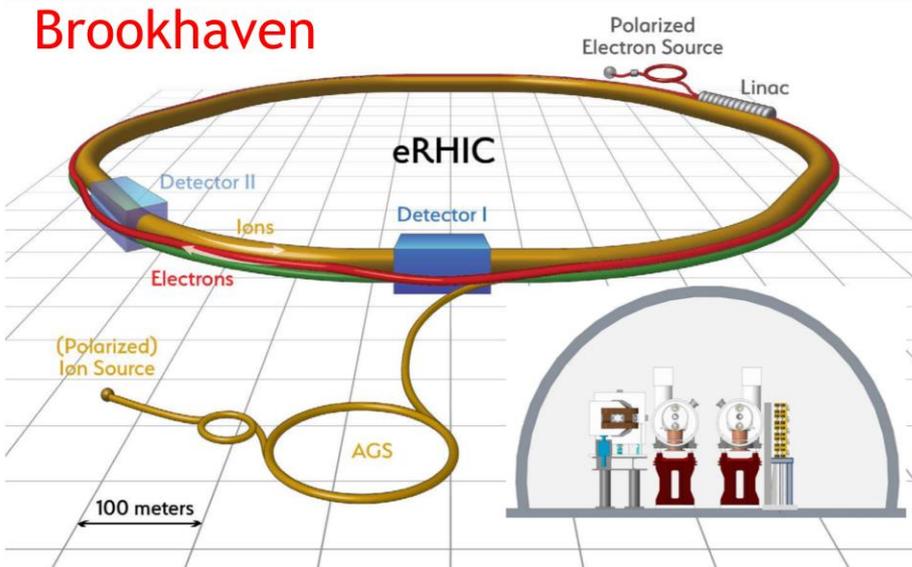
A future (2025~) high-luminosity polarized ep , eA collider dedicated to the study of the nucleon and nucleus structure.

Center-of-mass energy
Luminosity

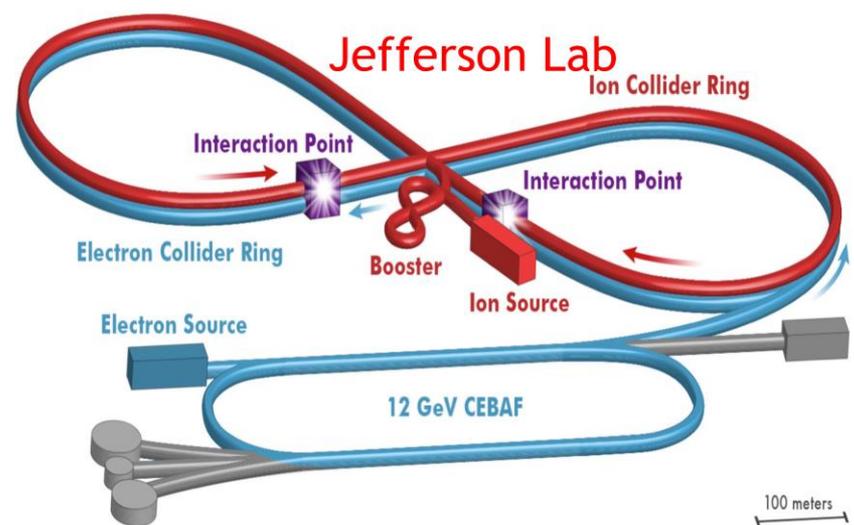
$$20 \lesssim \sqrt{s} \lesssim 140 \text{ GeV}$$
$$\sim 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$$

2 possible realizations

Brookhaven



Jefferson Lab

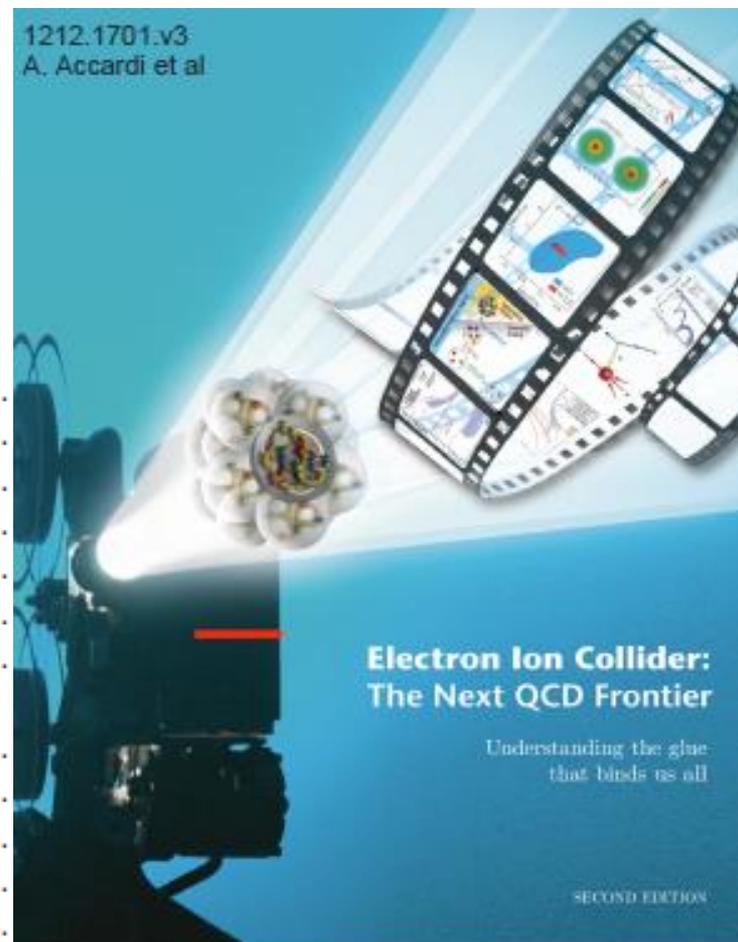


Nucleon tomography



2 Spin and Three-Dimensional Structure of the Nucleon

- 2.1 Introduction
- 2.2 The Longitudinal Spin of the Nucleon
 - 2.2.1 Introduction
 - 2.2.2 Status and Near Term Prospects
 - 2.2.3 Open Questions and the Role of an EIC
- 2.3 Confined Motion of Partons in Nucleons: TMDs
 - 2.3.1 Introduction
 - 2.3.2 Opportunities for Measurements of TMDs at the EIC
Semi-inclusive Deep Inelastic Scattering
Access to the Gluon TMDs
 - 2.3.3 Summary
- 2.4 Spatial Imaging of Quarks and Gluons
 - 2.4.1 Physics Motivations and Measurement Principle

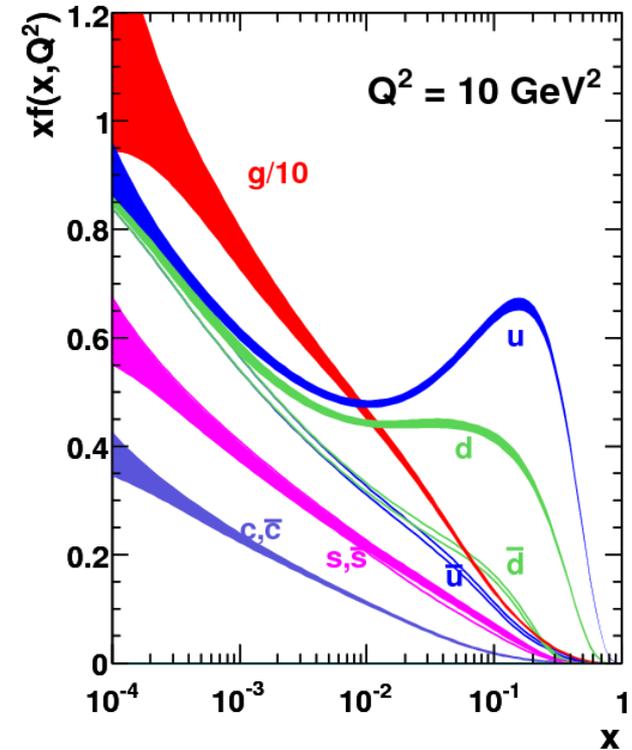
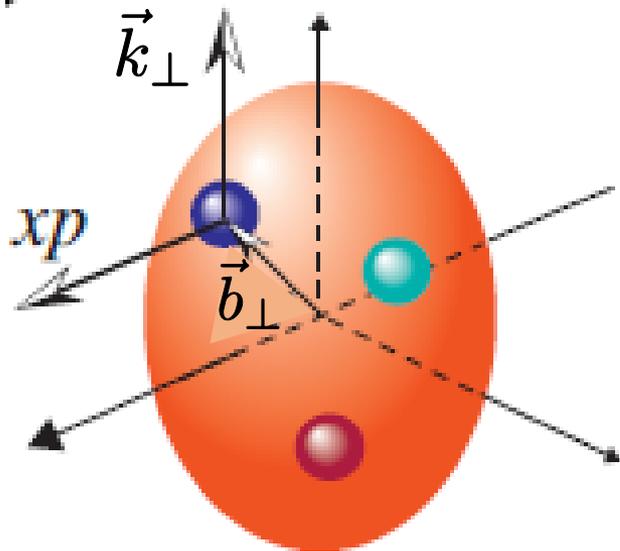


1D tomography: Parton distribution function (PDF)

$$f(x) = \int \frac{dz^-}{4\pi} e^{ixP^+z^-} \langle P | \bar{q}(-\frac{z^-}{2}) \gamma^+ q(\frac{z^-}{2}) | P \rangle$$

Probability distribution of quarks and gluons with **longitudinal** momentum fraction

$$x = \frac{p_{parton}^+}{P_{proton}^+}$$



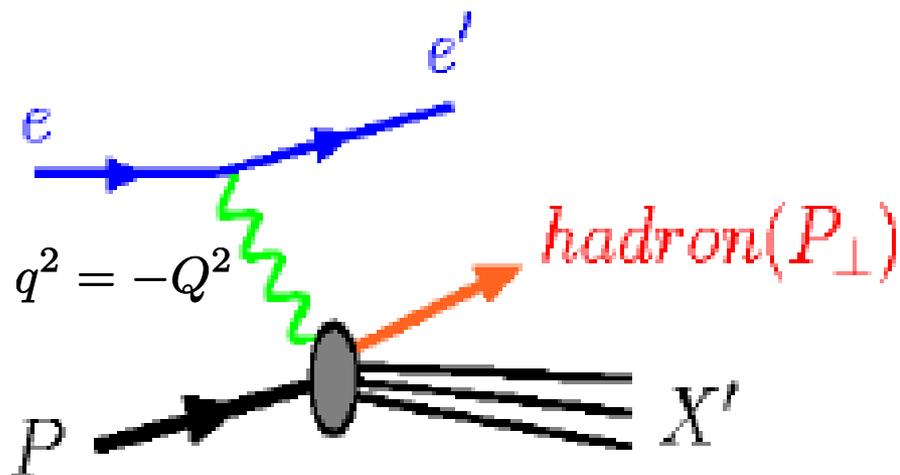
The nucleon is much more complicated!
Partons also have transverse momentum \vec{k}_\perp
and are spread in impact parameter space \vec{b}_\perp

3D tomography:

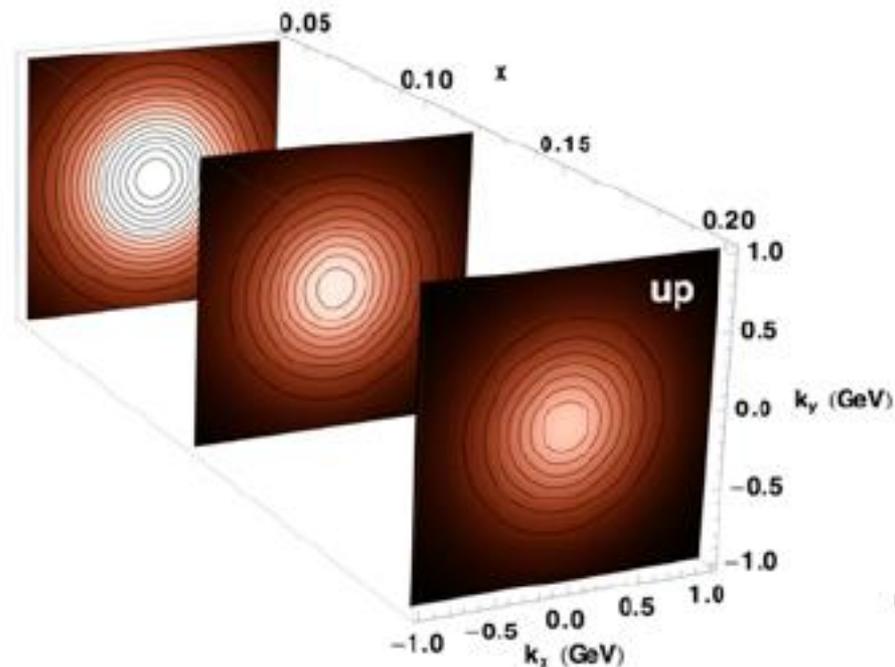
Transverse momentum dependent distributions (TMD)

$$f(x, \vec{k}_\perp) = \int \frac{dz^- d^2 z_\perp}{16\pi^3} e^{ixP^+ z^- - i\vec{k}_\perp \cdot \vec{z}_\perp} \langle P | \bar{q}(-z/2) \gamma^+ W q(z/2) | P \rangle$$

Relevant in semi-inclusive DIS (SIDIS), etc.

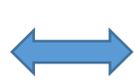


$$Q \gg P_\perp \gtrsim \Lambda_{QCD}$$



3D tomography: Generalized parton distributions (GPD)

$$f(x, \vec{\Delta}_{\perp}) \sim \int \frac{dz^{-}}{4\pi} e^{ixP^{+}z^{-}} \langle P - \frac{\Delta}{2} | \bar{q}(-z/2) \gamma^{+} q(z/2) | P + \frac{\Delta}{2} \rangle$$

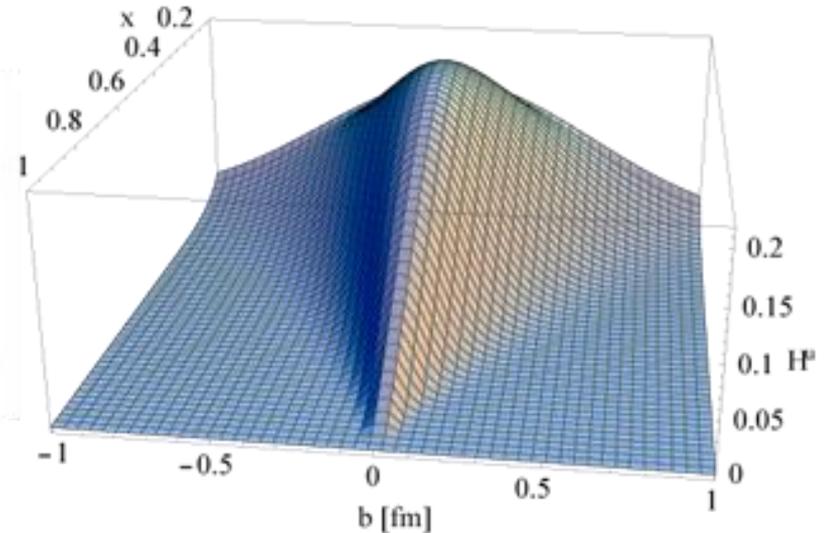
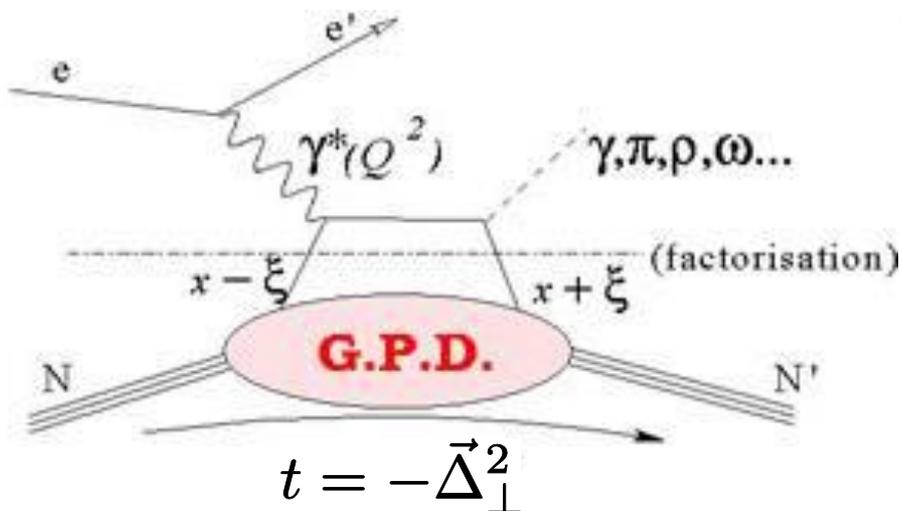


$$f(x, \vec{b}_{\perp})$$

distribution of partons in **impact parameter** space

Fourier transform

Deeply Virtual Compton Scattering (DVCS)

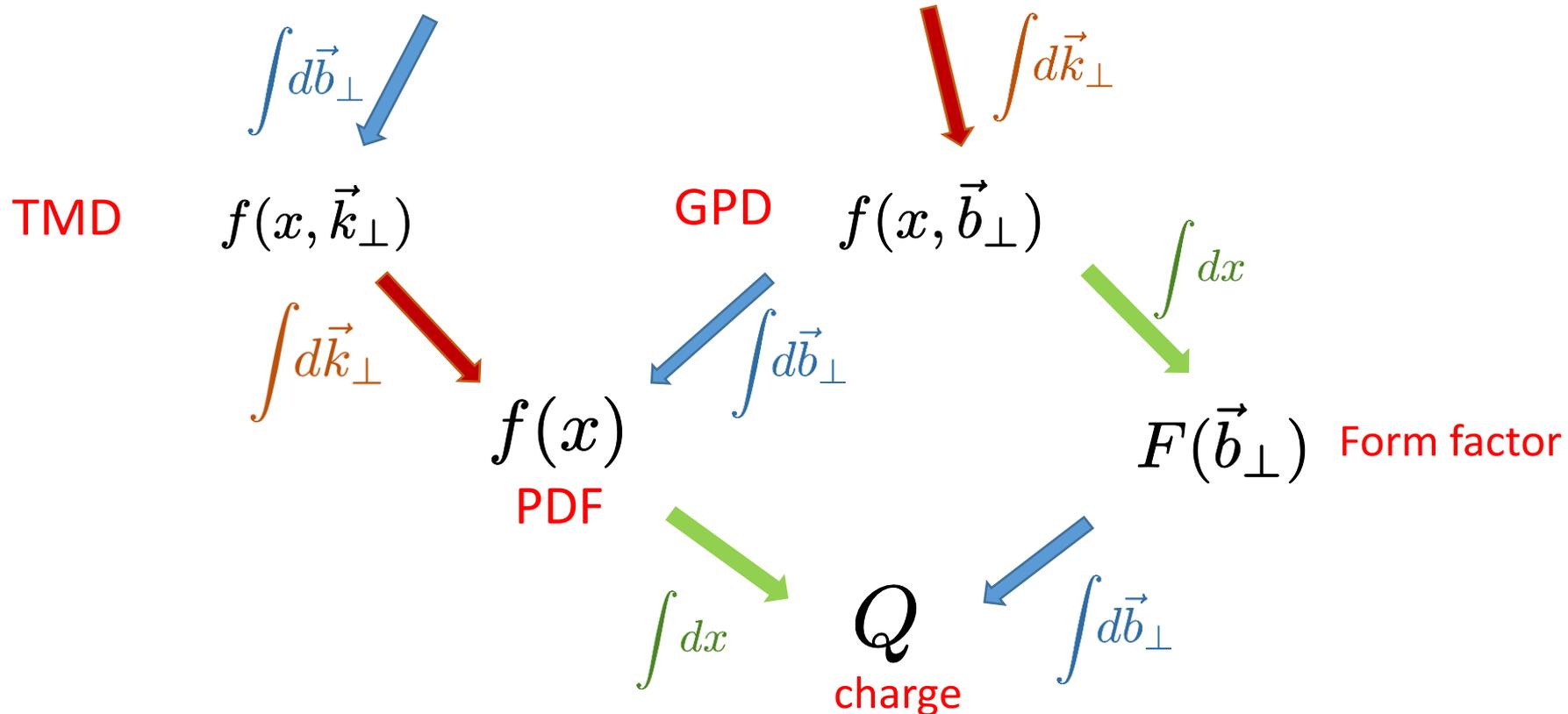


5D tomography:

Wigner distribution— the “mother distribution”

Belitsky, Ji, Yuan (2003);
Lorce, Pasquini (2011)

$$W(x, \vec{k}_\perp, \vec{b}_\perp) = \int \frac{d^2 \Delta_\perp}{(2\pi)^2} e^{i\vec{b}_\perp \cdot \vec{\Delta}_\perp} \int \frac{dz^- d^2 z_\perp}{16\pi^3} e^{ixP^+ z^- - i\vec{k}_\perp \cdot \vec{z}_\perp} \langle P - \frac{\Delta}{2} | \bar{q}(-z/2) \gamma^+ q(z/2) | P + \frac{\Delta}{2} \rangle$$



5D tomography: GTMD and Husimi

GTMD Meissner, Metz, Schlegel (2009)

Husimi Hagiwara, YH (2015)

$$G(x, \vec{k}_\perp, \vec{\Delta}_\perp)$$

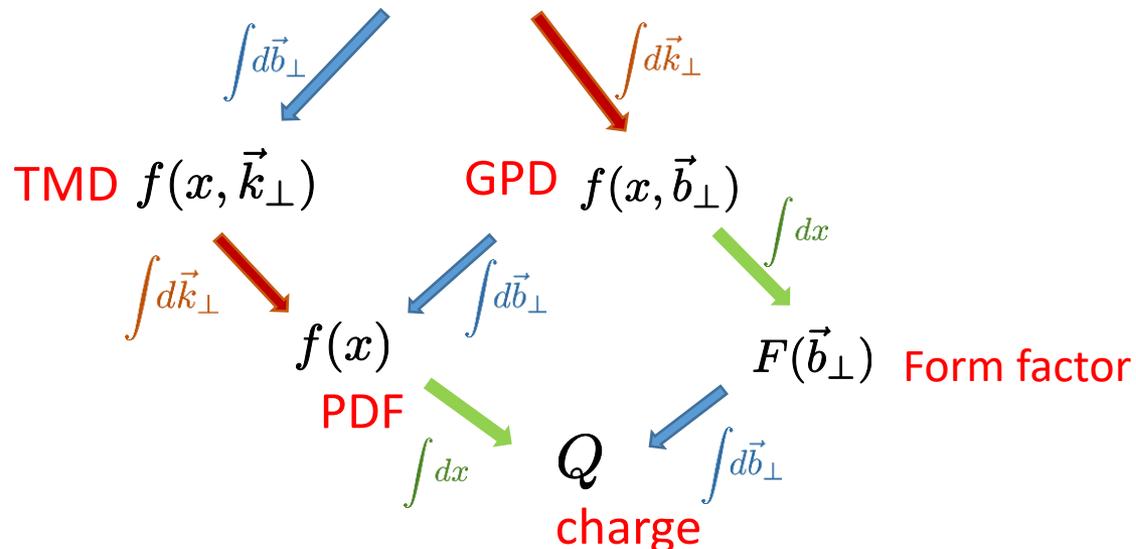
$$H(x, \vec{k}_\perp, \vec{b}_\perp)$$

$$\vec{b}_\perp \leftrightarrow \vec{\Delta}_\perp$$

Wigner

Gaussian smearing in k, b

$$W(x, \vec{k}_\perp, \vec{b}_\perp)$$



Gluon Wigner distribution—there are two of them

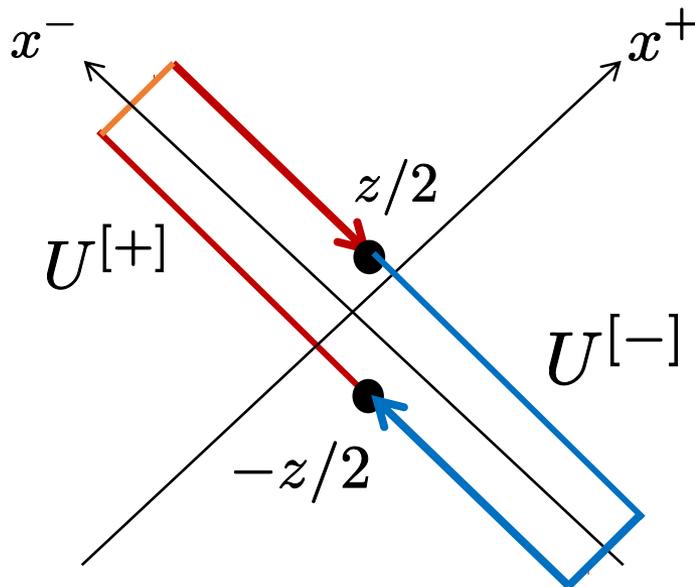
$$xW(x, \vec{k}_\perp, \vec{b}_\perp) = \int \frac{d^2\Delta_\perp}{(2\pi)^2} e^{i\vec{b}_\perp \cdot \vec{\Delta}_\perp} \int \frac{dz^- d^2z_\perp}{16\pi^3} e^{ixP^+z^- - i\vec{k}_\perp \cdot \vec{z}_\perp} \langle P - \Delta/2 | F^{+i}(-z/2) F_i^+(z/2) | P + \Delta/2 \rangle$$



There are **two** ways to make it gauge invariant

[Bomhof, Mulders \(2008\)](#)

[Dominguez, Marquet, Xiao, Yuan \(2011\)](#)



Weizsacker-Williams (WW) distribution

$$\text{Tr}[F(-z/2)U^{[+]}F(z/2)U^{[+]}]$$

Dipole distribution

$$\text{Tr}[F(-z/2)U^{[+]}F(z/2)U^{[-]}]$$

Wigner in 2012 EIC white paper?

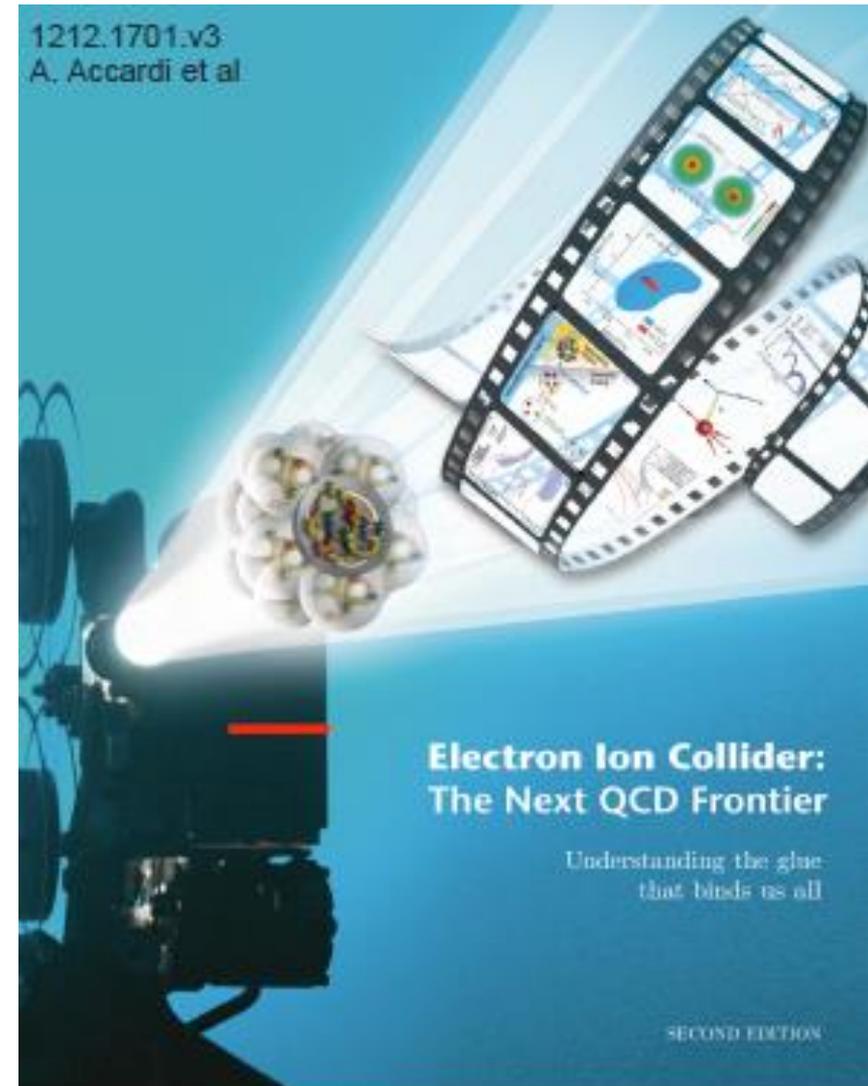
Almost no account.

Only briefly mentioned in two places.

*Although **there is no known way to measure Wigner distributions** for quarks and gluons, they provide a unifying theoretical framework for the different aspects of hadron structure.*

A lot of progress since then

Wigner \neq TMD+GPD



Wigner distribution and orbital angular momentum

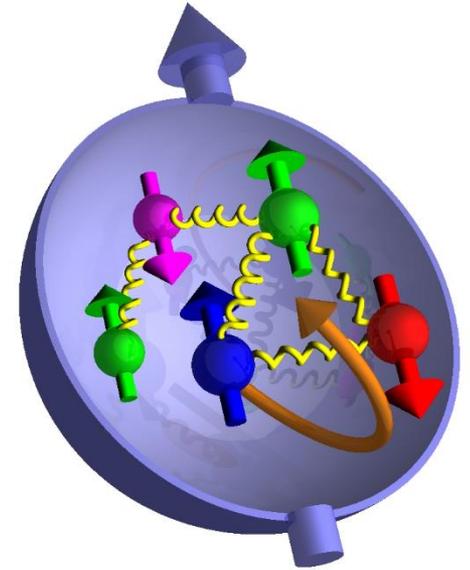
Jaffe-Manohar decomposition

$$\frac{1}{2} = \frac{1}{2} \Delta\Sigma + \Delta G + L^q + L^g$$

Quarks' helicity

Glucos' helicity

Canonical
Orbital
angular momentum
(OAM)



$$L^{q,g} = \int dx \int d^2b_{\perp} d^2k_{\perp} (\vec{b}_{\perp} \times \vec{k}_{\perp})_z W^{q,g}(x, \vec{b}_{\perp}, \vec{k}_{\perp})$$

Lorce, Pasquini, (2011);
YH (2011)

'PDF' for OAM

$$L^{q,g}(x) = \int d^2b_{\perp} d^2k_{\perp} (\vec{b}_{\perp} \times \vec{k}_{\perp})_z W^{q,g}(x, \vec{b}_{\perp}, \vec{k}_{\perp})$$

`Entropy' of partons

Hagiwara, YH, Xiao, Yuan (2018)

Phase space distribution naturally defines an entropy.

Use the QCD Husimi distribution

$$S(x) \equiv - \int d^2 b_{\perp} d^2 k_{\perp} x H(x, b_{\perp}, k_{\perp}) \ln x H(x, b_{\perp}, k_{\perp})$$

$$S(x) \underset{x \rightarrow 0}{\sim} \frac{N_c}{\alpha_s} Q_s^2(x) S_{\perp} \propto A \left(\frac{1}{x} \right)^{\# \alpha_s}$$

cf. Kutak (2011)
Kovner-Lublinsky (2015)

Measure of `complexity' of the multiparton system.

Saturation of entropy due to the Pomeron loop effect?

Connection to the `jet entropy' in the final state? Neill, Waalewijn (2018)

Wigner distribution: Is it measurable?

In quantum optics, yes!

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PHYSICAL REVIEW LETTERS

1 MARCH 1993

Measurement of the Wigner Distribution and the Density Matrix of a Light Mode Using Optical Homodyne Tomography: Application to Squeezed States and the Vacuum

D. T. Smithey, M. Beck, and M. G. Raymer

Department of Physics and Chemical Physics Institute, U

A. Faridani

Department of Mathematics, Oregon State Uni

(Received 16 Novembe

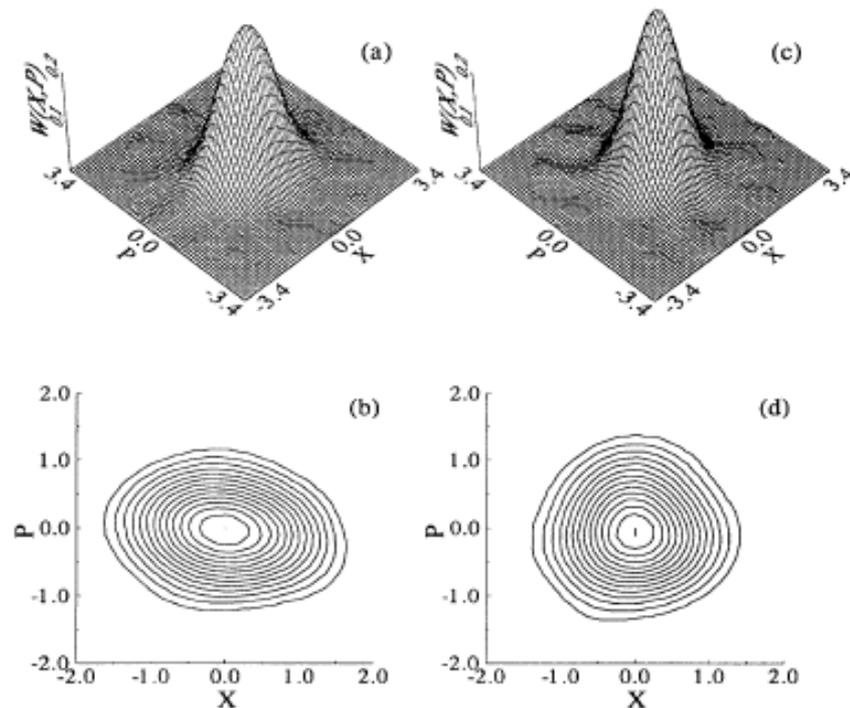


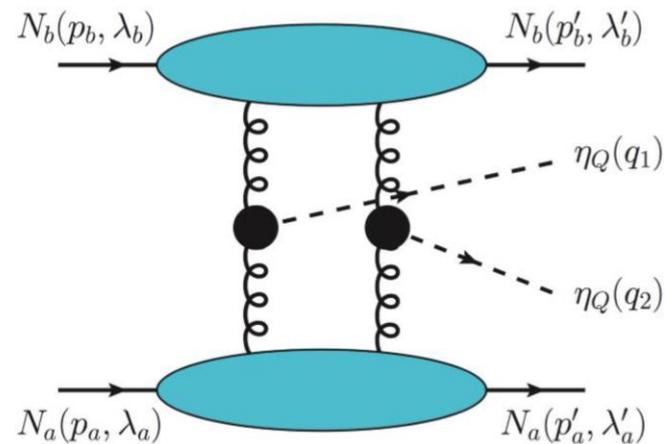
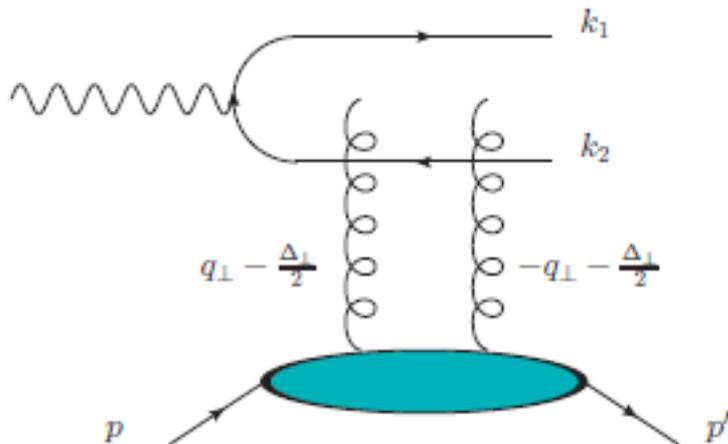
FIG. 1. Measured Wigner distributions for (a),(b) a squeezed state and (c),(d) a vacuum state, viewed in 3D and as contour plots, with equal numbers of constant-height contours. Squeezing of the noise distribution is clearly seen in (b).

Measuring Wigner/GTMD in experiments

- ① Tag two hadrons (jets) in the final state, together with the recoiling proton

YH, Xiao, Yuan (2016); Bhattacharya, Metz, Zhou (2017)

$G(x, k_{\perp}, \Delta_{\perp})$
 Relative momentum between the two hadrons
 Recoiling proton momentum



Measuring Wigner/GTMD in experiments

② Go to small-x (forward particle production)

Approximate $e^{ixP^+z^-} \approx 1$

$$xW(x, \vec{k}_\perp, \vec{b}_\perp) \approx \frac{2N_c}{\alpha_s} \int \frac{d^2\vec{r}_\perp}{(2\pi)^2} e^{i\vec{k}_\perp \cdot \vec{r}_\perp} \left(\frac{1}{4} \vec{\nabla}_b^2 - \vec{\nabla}_r^2 \right) S_x(\vec{b}_\perp, \vec{r}_\perp)$$

“Dipole S-matrix” $S_x(\vec{b}_\perp, \vec{r}_\perp) = \left\langle \frac{1}{N_c} \text{Tr} U \left(\vec{b}_\perp - \frac{\vec{r}_\perp}{2} \right) U^\dagger \left(\vec{b}_\perp + \frac{\vec{r}_\perp}{2} \right) \right\rangle_x$

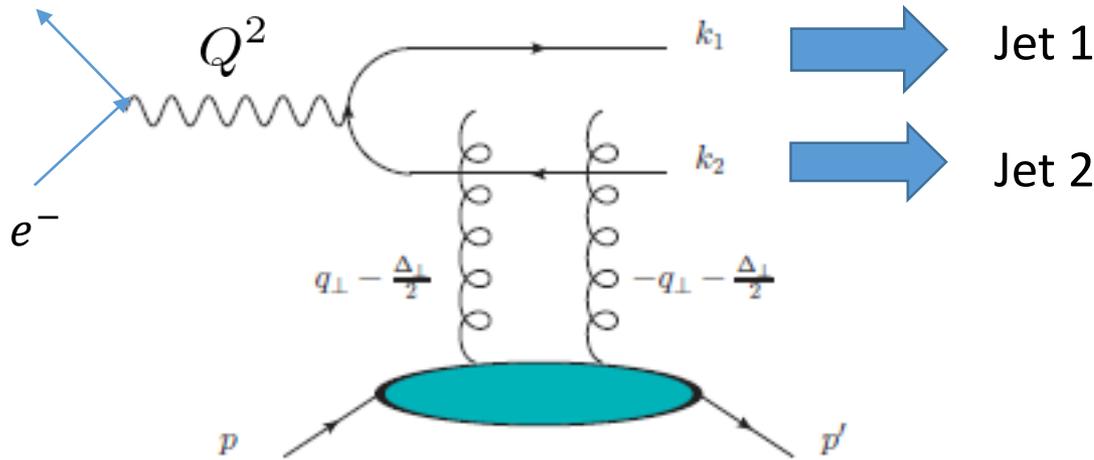
cos 2φ correlation expected

$$W(x, \vec{k}_\perp, \vec{b}_\perp) = W_0(x, k_\perp, b_\perp) + 2 \cos 2(\phi_k - \phi_b) W_1(x, k_\perp, b_\perp) + \dots$$

‘Elliptic Wigner’ distribution

Probing dipole Wigner (GTMD) in diffractive dijet production

YH, Xiao, Yuan (2016), see also, Altinoluk, Armesto, Beuf, Rezaeian (2015)



$$\vec{\Delta}_\perp = -(\vec{k}_{1\perp} + \vec{k}_{2\perp})$$

$$\vec{P}_\perp = \frac{1}{2}(\vec{k}_{2\perp} - \vec{k}_{1\perp})$$

GTMD

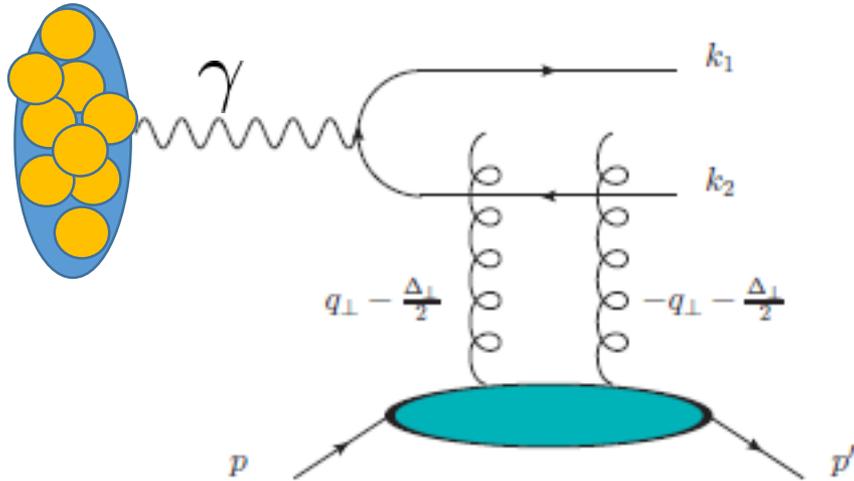
$$\frac{d\sigma}{dy_1 dy_2 d^2\vec{\Delta}_\perp d^2\vec{P}_\perp} \propto z(1-z)[z^2 + (1-z)^2] \int d^2q_\perp d^2q'_\perp S(q_\perp, \Delta_\perp) S(q'_\perp, \Delta_\perp)$$

$$\times \left[\frac{\vec{P}_\perp}{P_\perp^2 + \epsilon^2} - \frac{\vec{P}_\perp - \vec{q}_\perp}{(P_\perp - q_\perp)^2 + \epsilon^2} \right] \cdot \left[\frac{\vec{P}_\perp}{P_\perp^2 + \epsilon^2} - \frac{\vec{P}_\perp - \vec{q}'_\perp}{(P_\perp - q'_\perp)^2 + \epsilon^2} \right]$$

$$\sim d\sigma_0 + 2 \cos 2(\phi_P - \phi_\Delta) d\tilde{\sigma}$$

$$\epsilon^2 = z(1-z)Q^2$$

In ultra-peripheral collisions, too!



Hagiwara, YH, Pasechnik, Tasevsky, Teryaev (2017)

Q^2 preferably small



Use the Weizacker-Williams photons in UPC!

$$\frac{d\sigma^{pA}}{dy_1 dy_2 d^2\vec{k}_{1\perp} d^2\vec{k}_{2\perp}} = \omega \frac{dN}{d\omega} \frac{N_c \alpha_{em} (2\pi)^4}{P_{\perp}^2} \sum_f e_f^2 2z(1-z)(z^2 + (1-z)^2) (A^2 + 2 \cos 2(\phi_P - \phi_{\Delta}) AB)$$

photon flux $\propto Z^2$

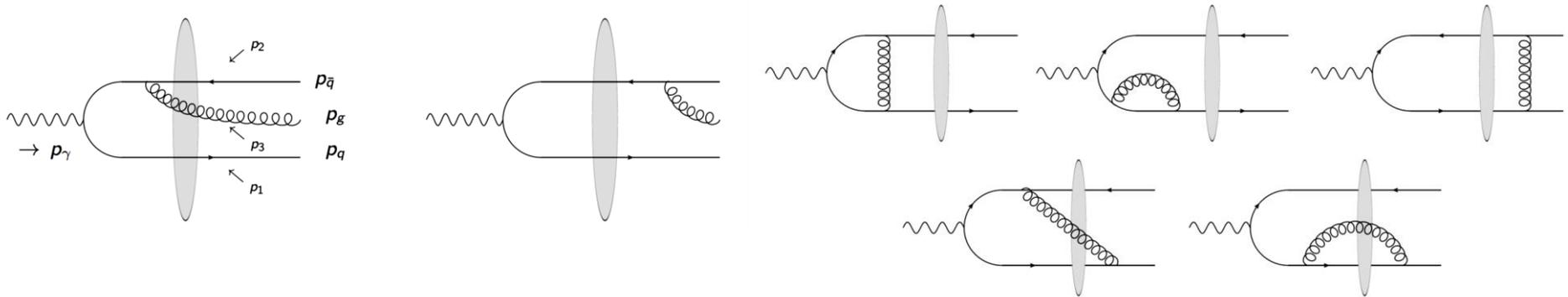
Inversion can be done analytically.

$$S_0(P_{\perp}, \Delta_{\perp}) = \frac{1}{P_{\perp}} \frac{\partial}{\partial P_{\perp}} A(P_{\perp}, \Delta_{\perp}).$$

$$S_1(P_{\perp}, \Delta_{\perp}) = \frac{\partial B(P_{\perp}, \Delta_{\perp})}{\partial P_{\perp}^2} - \frac{2}{P_{\perp}^2} \int^{P_{\perp}^2} \frac{dP'_{\perp}}{P'_{\perp}} B(P'_{\perp}, \Delta_{\perp})$$

Factorization at NLO

Boussarie, Grabovsky, Szymanowski, Wallon (2016)



$$\Phi_L^{(0)} = \frac{2x\bar{x}p_V^+Q}{(\bar{x}\vec{p}_1 - x\vec{p}_2)^2 + x\bar{x}Q^2},$$

$$\Phi_T^{(0)} = -\frac{(x - \bar{x})p_V^+(\bar{x}\vec{p}_{1\perp} - x\vec{p}_{2\perp}) \cdot \vec{\epsilon}_{\gamma T}}{(\bar{x}\vec{p}_1 - x\vec{p}_2)^2 + x\bar{x}Q^2}$$

No end point singularity, even for a transverse photon and even in the **photoproduction limit** and even at NLO.

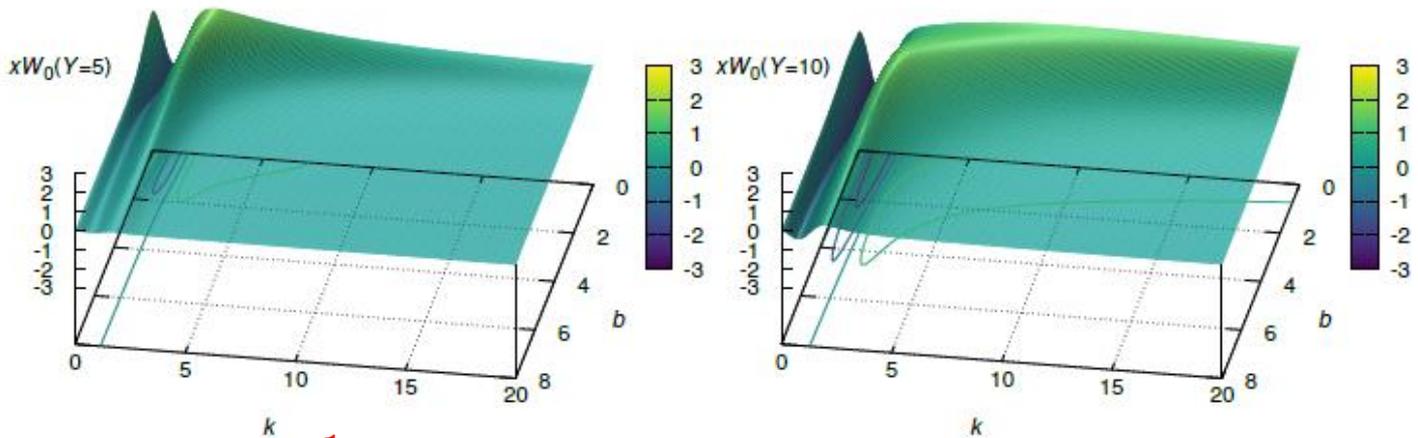
With null transverse momenta in the t channel, one could encounter $x \in \{0, 1\}$ end point singularities as $\frac{1}{x\bar{x}Q^2}$ thus **breaking collinear factorization**.

Dipole Wigner from Balitsky-Kovchegov equation

Hagiwara, YH, Ueda (2016)

Peak at the **saturation momentum**

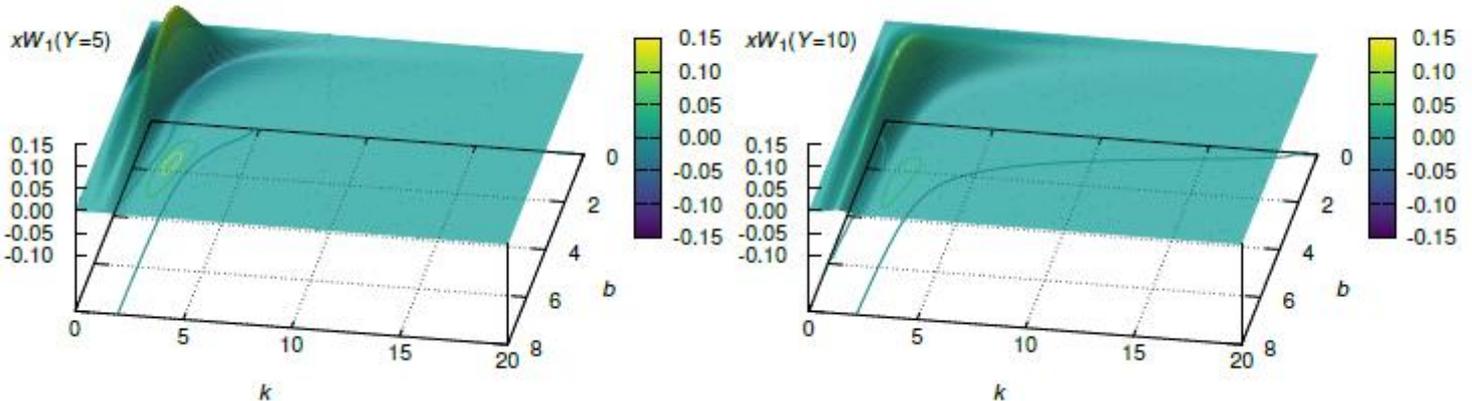
$W_0(k_{\perp}, b_{\perp})$



$$Y = \ln \frac{1}{x} = 5$$

$$Y = 10$$

$W_1(k_{\perp}, b_{\perp})$



Elliptic part small in magnitude (a few percent effect). No geometric scaling.

Elliptic Wigner in DVCS

YH, Xiao, Yuan (2017)

Gluson transversity GPD

$$\begin{aligned} & \frac{1}{P^+} \int \frac{d\zeta^-}{2\pi} e^{ixP^+\zeta^-} \langle p' | F^{+i}(-\zeta/2) F^{+j}(\zeta/2) | p \rangle \\ &= \frac{\delta^{ij}}{2} x H_g(x, \Delta_\perp) + \frac{x E_{Tg}(x, \Delta_\perp)}{2M^2} \left(\Delta_\perp^i \Delta_\perp^j - \frac{\delta^{ij} \Delta_\perp^2}{2} \right) + \dots, \end{aligned}$$

$$x E_{Tg}(x, \Delta_\perp) = \frac{4N_c M^2}{\alpha_s \Delta_\perp^2} \int d^2 q_\perp q_\perp^2 S_1 \quad \leftarrow \text{Elliptic GTMD}$$

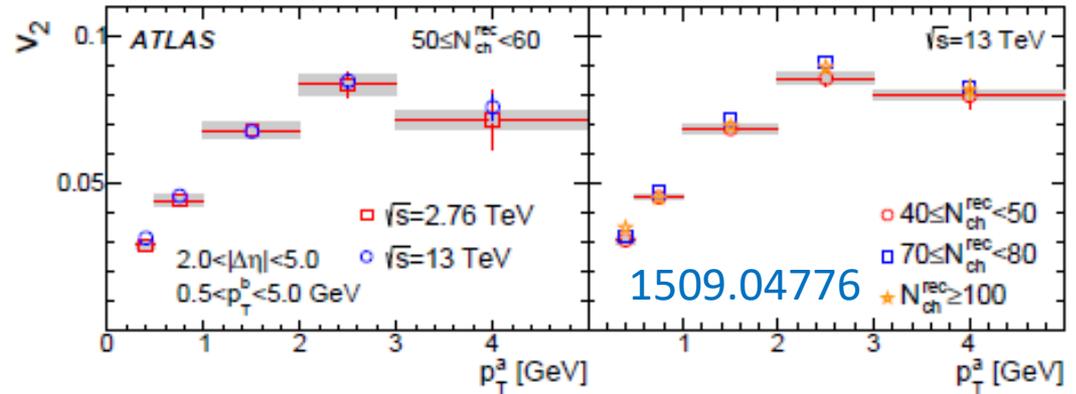
$$\begin{aligned} \frac{d\sigma(ep \rightarrow e'\gamma p')}{dx_B dQ^2 d^2 \Delta_\perp} &= \frac{\alpha_{em}^3}{\pi x_{Bj} Q^2} \left\{ \left(1 - y + \frac{y^2}{2} \right) (\mathcal{A}_0^2 + \mathcal{A}_2^2) + 2(1 - y) \mathcal{A}_0 \mathcal{A}_2 \cos(2\phi_{\Delta l}) \right. \\ &\quad \left. + (2 - y) \sqrt{1 - y} (\mathcal{A}_0 + \mathcal{A}_2) \mathcal{A}_L \cos \phi_{\Delta l} + (1 - y) \mathcal{A}_L^2 \right\} \end{aligned}$$

Elliptic Wigner in high-multiplicity pp and pA

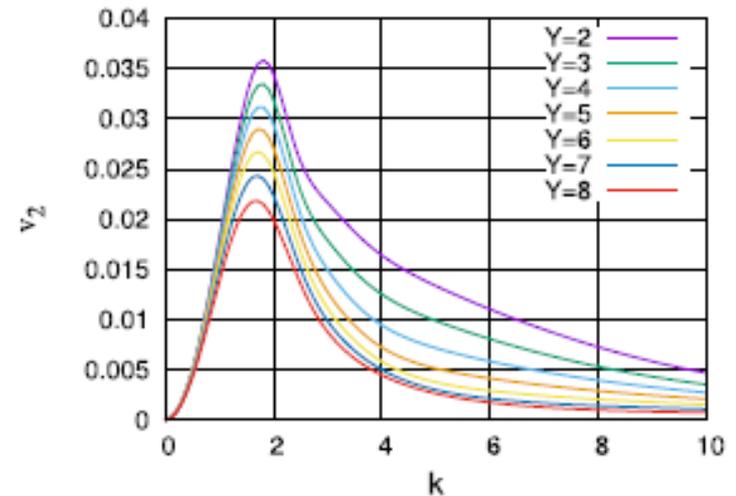
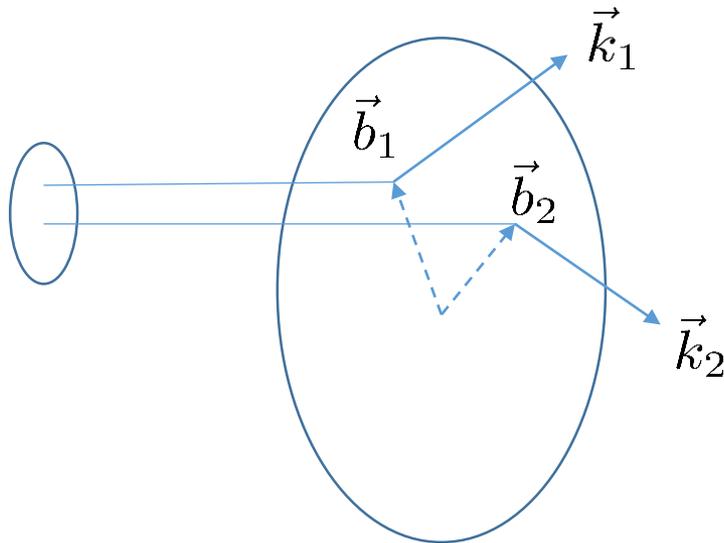
Kopeliovich et al. (2008),
 Levin Rezaeian (2011),
 Hagiwara, YH, Xiao, Yuan (2017)

Elliptic flow v_2 observed in
 high-multiplicity pp and pA.

Initial state or final state effect?



Double parton scattering + elliptic Wigner = elliptic flow

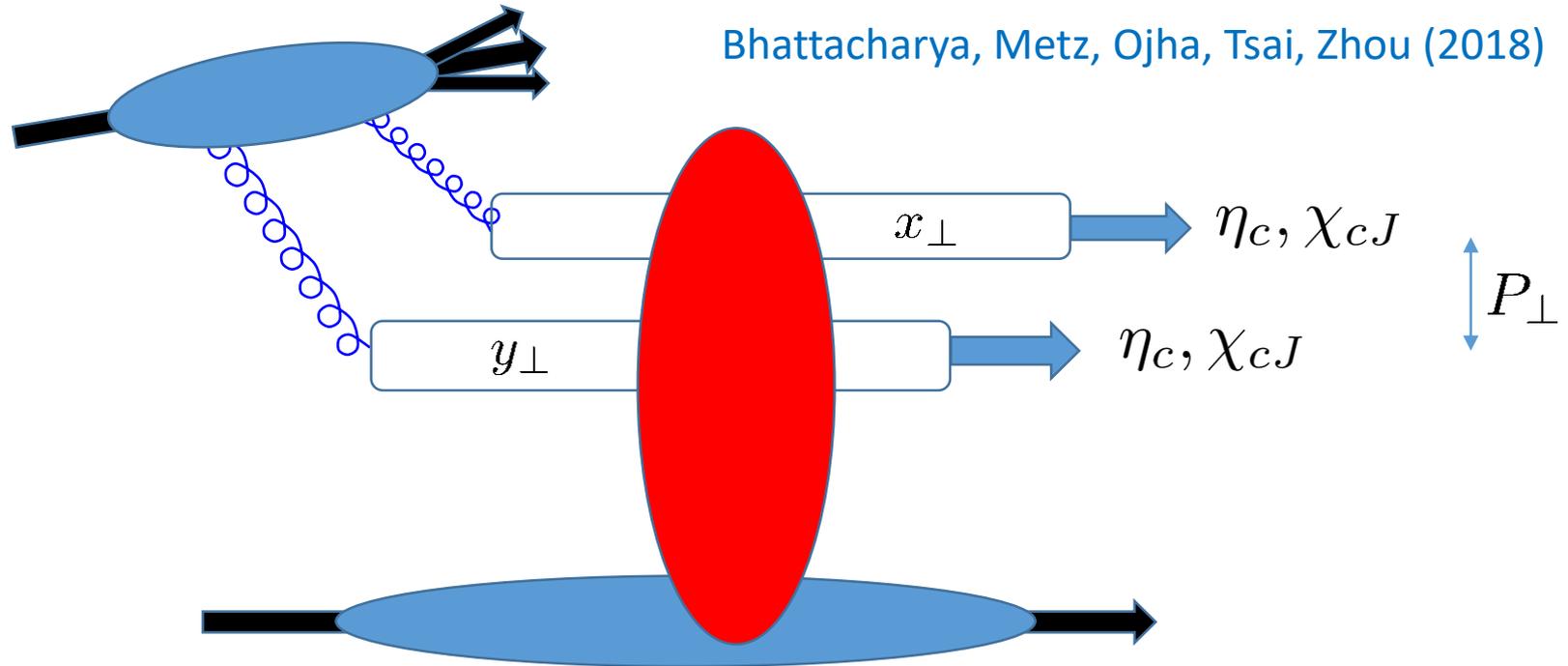


Probing the gluon WW GTMD in pp

Boussarie, YH, Xiao, Yuan (2018)

See also,

Bhattacharya, Metz, Ojha, Tsai, Zhou (2018)



Diffractive production of a $C = +1$ quarkonium pair

Amplitude proportional to

$$\int d^2(x_\perp - y_\perp) e^{iP_\perp \cdot (x_\perp - y_\perp)} \langle P' | U_x \vec{\partial} U_x^\dagger U_y \vec{\partial} U_y^\dagger | P \rangle$$

Very simple result in the case of χ_{c1}, χ_{c1} production, in the limit $P_{\perp} \gg \Delta_{\perp}$

$$\begin{aligned}
 & \frac{d\sigma}{dY_1 dY_2 d^2 \Delta_{\perp} d^2 P_{\perp}} \\
 &= \frac{x_1 x_2 F(x_1, x_2)}{64 m^{18} N_c^4 (N_c^2 - 1)^2} \alpha_s^4 \langle \mathcal{O}_{\chi_1} \rangle^2 P_{\perp}^4 \left(G(P_{\perp}, \Delta_{\perp}) + \frac{P_{\perp}^2}{2M^2} G_2 \right)^2
 \end{aligned}$$

Gluon dPDF
Long distance matrix element (LDME)

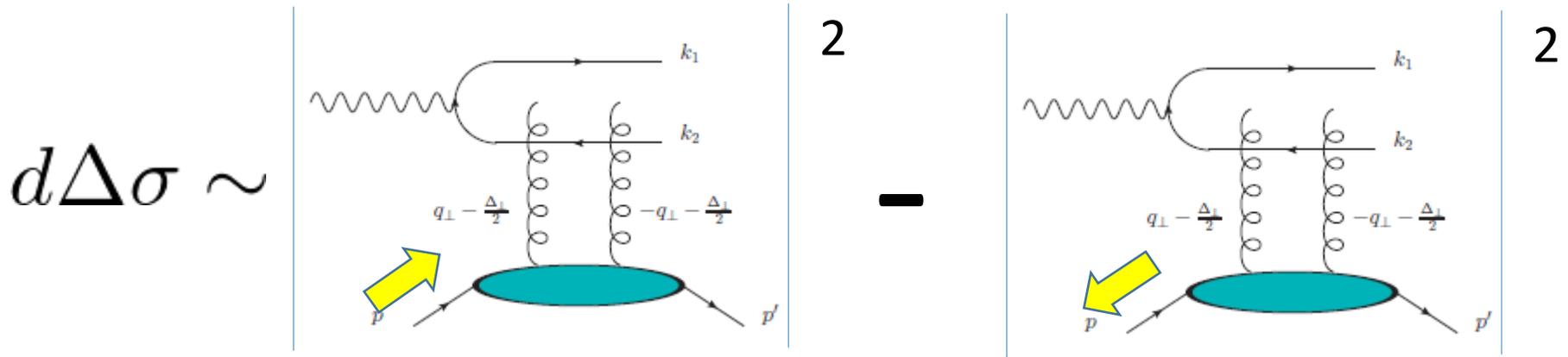
WW gluon GTMD
Linearly polarized WW gluon GTMD

No convolution in P_{\perp} .

Caveat: only color-singlet production included

Towards measuring the orbital angular momentum

Longitudinal single spin asymmetry in dijet production



Large-x gluon	Ji, Yuan, Zhao
Small-x gluon	YH, Nakagawa, Xiao, Yuan, Zhao
quark	Bhattacharya, Metz, Ojha, Tsai, Zhou

Sensitive to the OAM distribution

$$W(x, \vec{k}_\perp, \vec{b}_\perp) = W_0(x, k_\perp, b_\perp) + S^+ \sin(\phi_k - \phi_b) W_{OAM}(x, k_\perp, b_\perp) + \dots$$

Net angular momentum comes from the large-x region. At small-x, expect $L_g(x) \approx -\Delta G(x)$

Conclusions

- Let's get 5 dimensional. Even richer physics than TMD and GPD combined.
- Wigner/GTMD measurable in ep, pp, pA, including the elliptic part and spin-dependent part (connection to OAM).
- Need more foundational works. Proper definition, evolution,...

cf. [Ecchevarria, Idilbi, Kanazawa, Lorce, Metz, Pasquini, Schlegel \(2016\)](#)