



Wir schaffen Wissen – heute für morgen

NLO QCD corrections to Higgs boson pair production via gluon fusion

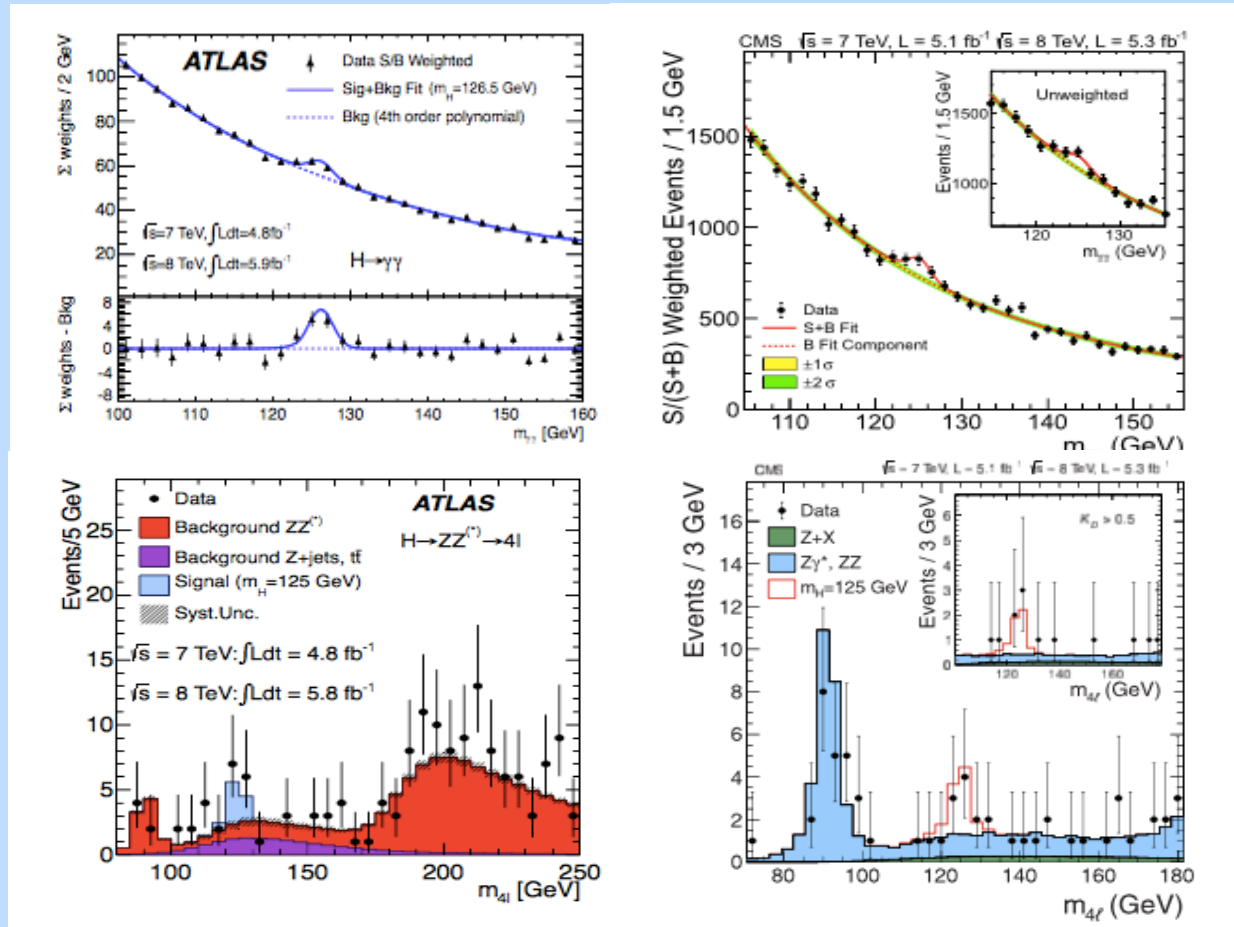
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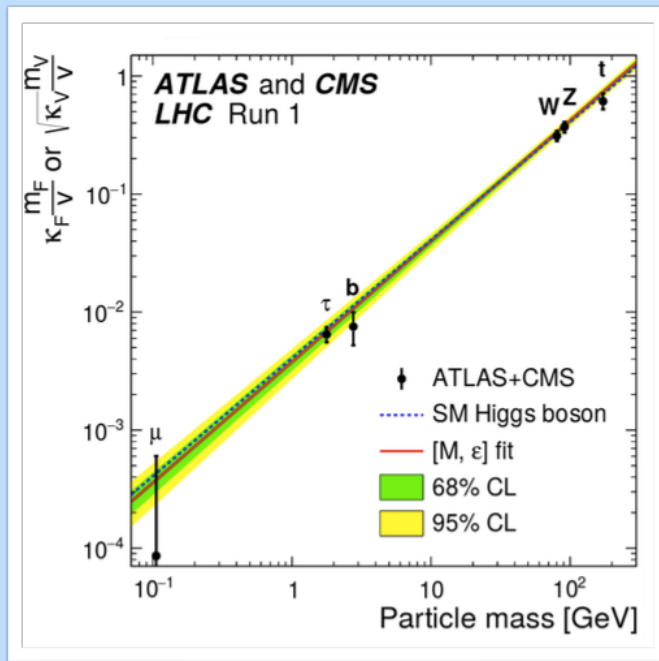
In collaboration with M. Spira, M. Mühlleitner, J. Baglio,
F. Campanario, J. Streicher

- Motivation
- Objective
- Previous work
- NLO Cross section
 - Virtual Corrections
 - Real Corrections
- Numerical analysis
- Next steps
- Conclusions

Detection of a Higgs boson with a mass ~ 125 GeV



- Higgs mass, coupling strengths, spin and CP already determined
- self-coupling strength still unknown



$$V(\phi) = \frac{\lambda}{2} \left\{ |\phi|^2 - \frac{v^2}{2} \right\}^2$$

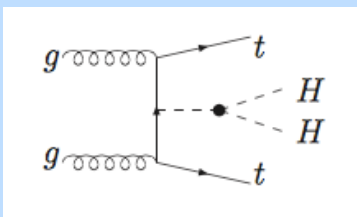
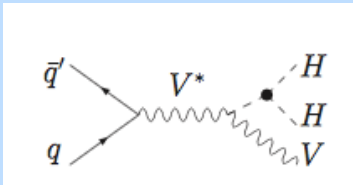
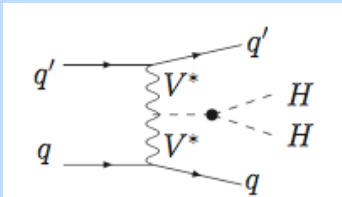
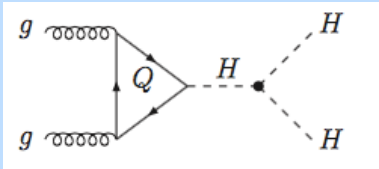
$$\lambda = \frac{M_H^2}{2v^2}$$

$$\lambda_{H^3} = 3 \frac{M_H^2}{v}$$

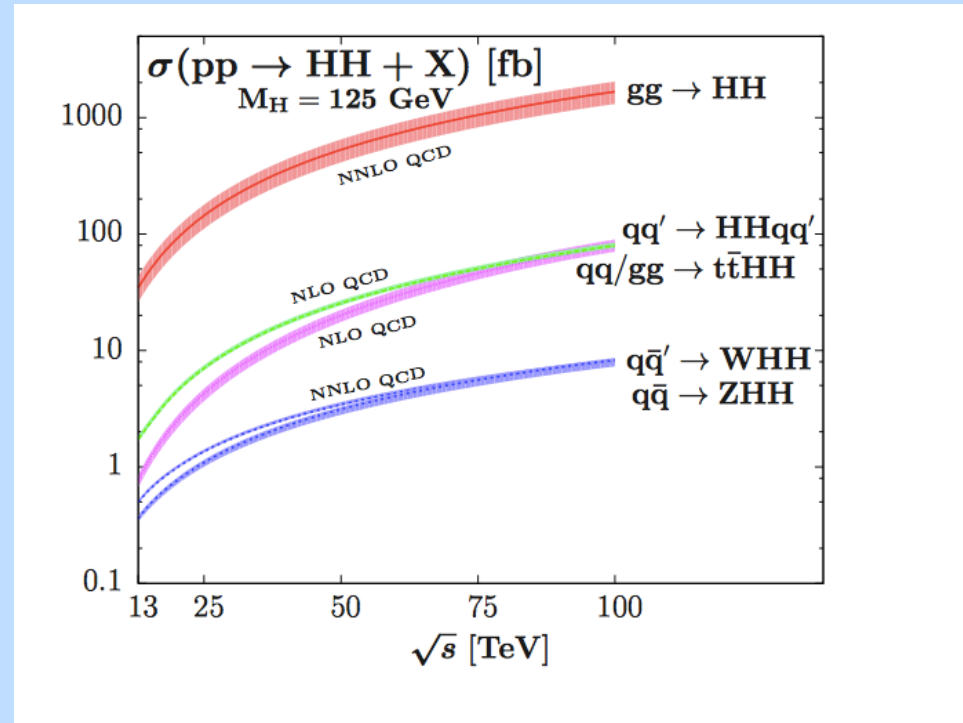
$$\lambda_{H^4} = 3 \frac{M_H^2}{v^2}$$

Higgs boson pair production

Production channels

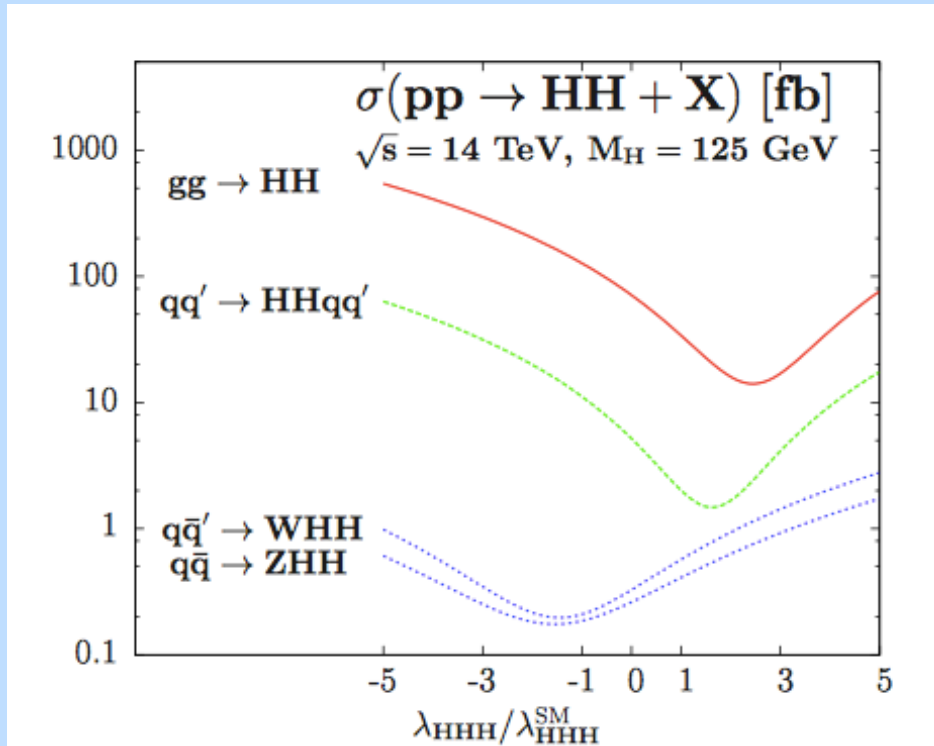


Cross sections



Baglio, Djouadi, Quevillon

Uncertainties:

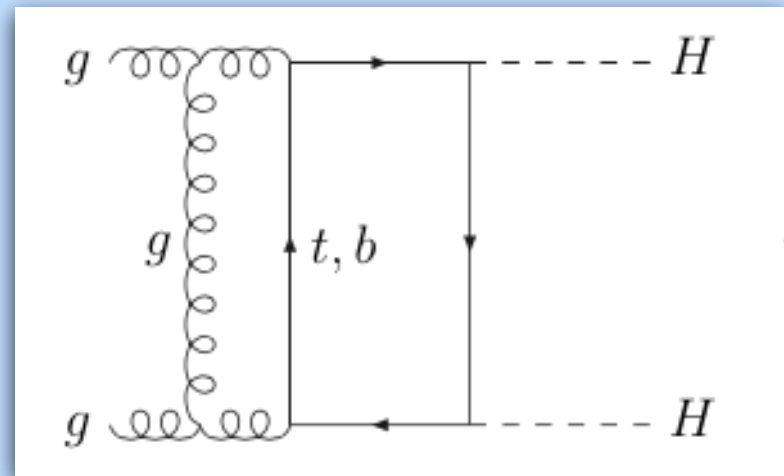


$$gg \rightarrow HH : \frac{\Delta\sigma}{\sigma} \sim -\frac{\Delta\lambda}{\lambda}$$

Baglio, Djouadi, Gröber,
Mühlleitner, Quevillon, Spira

Gluon fusion $gg \rightarrow HH$: loop induced

Complete calculation of the NLO QCD corrections (2-loop) considering the top- and bottom mass dependences in the context of the Standard Model \rightarrow 2-loop integrals with 3 kinematical parameter ratios



- Virtual & real (N)NLO QCD corrections in large top mass limit: $\sim 100\%$

Dawson, Dittmaier, Spira
de Florian, Mazzitelli
Grigo, Melnikov, Steinhauser

- Large top mass expansion: $\sim \pm 10\%$

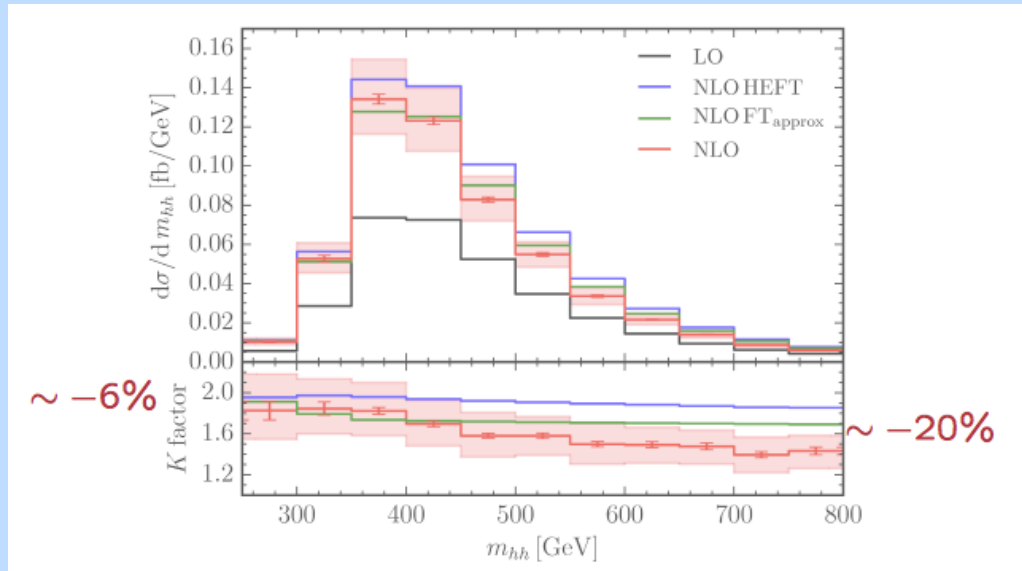
$$\sigma = \sigma_0 + \frac{\sigma_1}{m_t^2} + \dots + \frac{\sigma_4}{m_t^8}$$

Grigo, Hoff, Melnikov,
Steinhauser

- NLO mass effects of the real NLO correction alone
 $\sim -10\%$

Frederix, Frixione, Hirschi, Maltoni,
Mattelaer, Torrielli, Vryonidou, Zaro

- NLO QCD corrections including the full top mass dependence

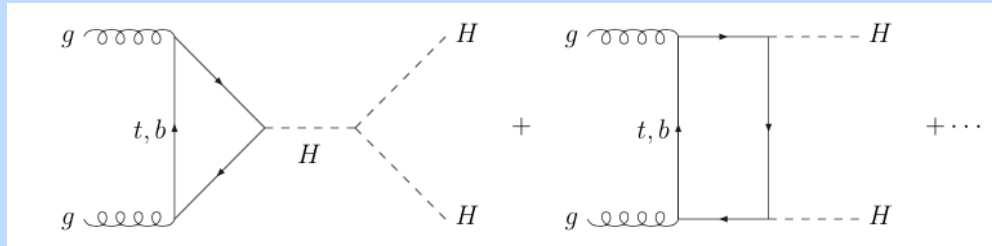


Borowka, Greiner, Heinrich, Jones, Kerner, Schlenk, Schubert, Zirke

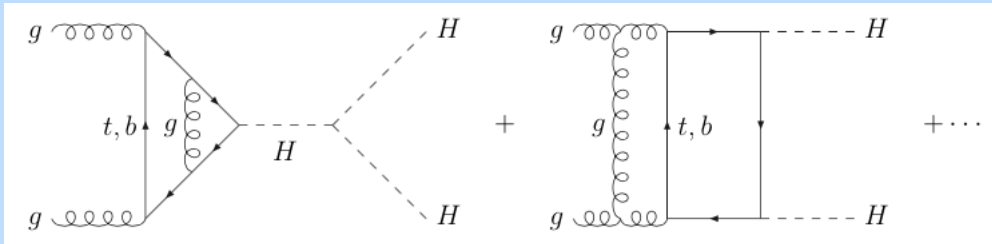
→ - 13.7 % mass effects

$$\sigma_{\text{NLO}}(pp \rightarrow HH + X) = \sigma_{\text{LO}} + \Delta\sigma_{\text{virt}} + \Delta\sigma_{gg} + \Delta\sigma_{gq} + \Delta\sigma_{q\bar{q}},$$

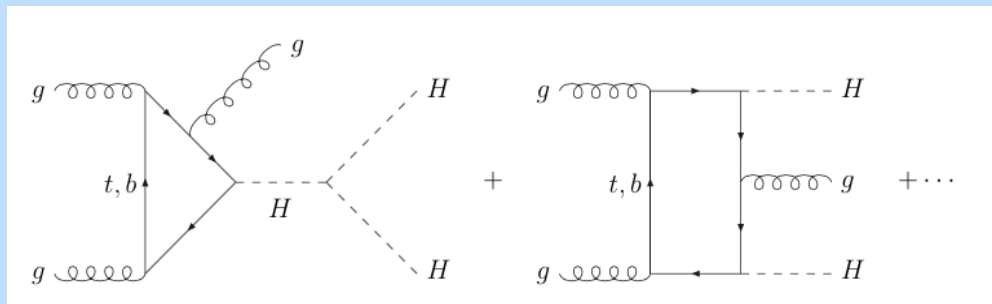
$\sigma_{\text{LO}}:$



$\Delta\sigma_{\text{virt}}:$



$\Delta\sigma_{ij}:$



$$\sigma_{\text{NLO}}(pp \rightarrow HH + X) = \sigma_{\text{LO}} + \Delta\sigma_{\text{virt}} + \Delta\sigma_{gg} + \Delta\sigma_{gq} + \Delta\sigma_{q\bar{q}},$$

$$\begin{aligned} \sigma_{\text{LO}} &= \int_{\tau_0}^1 d\tau \frac{d\mathcal{L}^{gg}}{d\tau} \hat{\sigma}_{\text{LO}}(Q^2 = \tau s) \\ \Delta\sigma_{\text{virt}} &= \frac{\alpha_s(\mu)}{\pi} \int_{\tau_0}^1 d\tau \frac{d\mathcal{L}^{gg}}{d\tau} \hat{\sigma}_{\text{LO}}(Q^2 = \tau s) C \\ \Delta\sigma_{gg} &= \frac{\alpha_s(\mu)}{\pi} \int_{\tau_0}^1 d\tau \frac{d\mathcal{L}^{gg}}{d\tau} \int_{\tau_0/\tau}^1 \frac{dz}{z} \hat{\sigma}_{\text{LO}}(Q^2 = z\tau s) \left\{ -z P_{gg}(z) \log \frac{M^2}{\tau s} \right. \\ &\quad \left. + d_{gg}(z) + 6[1 + z^4 + (1-z)^4] \left(\frac{\log(1-z)}{1-z} \right)_+ \right\} \\ \Delta\sigma_{gq} &= \frac{\alpha_s(\mu)}{\pi} \int_{\tau_0}^1 d\tau \sum_{q,\bar{q}} \frac{d\mathcal{L}^{gq}}{d\tau} \int_{\tau_0/\tau}^1 \frac{dz}{z} \hat{\sigma}_{\text{LO}}(Q^2 = z\tau s) \left\{ -\frac{z}{2} P_{gq}(z) \log \frac{M^2}{\tau s(1-z)^2} \right. \\ &\quad \left. + d_{gq}(z) \right\} \\ \Delta\sigma_{q\bar{q}} &= \frac{\alpha_s(\mu)}{\pi} \int_{\tau_0}^1 d\tau \sum_q \frac{d\mathcal{L}^{q\bar{q}}}{d\tau} \int_{\tau_0/\tau}^1 \frac{dz}{z} \hat{\sigma}_{\text{LO}}(Q^2 = z\tau s) d_{q\bar{q}}(z) \end{aligned}$$

$$C \rightarrow \pi^2 + \frac{11}{2} + C_{\Delta\Delta}, \quad d_{gg} \rightarrow -\frac{11}{2}(1-z)^3, \quad d_{gq} \rightarrow \frac{2}{3}z^2 - (1-z)^2, \quad d_{q\bar{q}} \rightarrow \frac{32}{27}(1-z)^3$$

Set-up

- Associate the 47 two-loop **box** diagrams to similar topologies
- Two independent calculations in the collaboration
- Use dimensional regularisation: $D = 4 - 2\epsilon$
- Perform Feynman parametrisation \rightarrow 6-dimensional integrals

Divergences

- Extraction of the ultraviolet divergences of the matrix elements using endpoint subtractions of the 6-dimensional Feynman integrals

$$\int_0^1 dx \frac{f(x)}{(1-x)^{1-\epsilon}} = \int_0^1 dx \frac{f(1)}{(1-x)^{1-\epsilon}} + \int_0^1 dx \frac{f(x) - f(1)}{(1-x)^{1-\epsilon}}$$
$$= \frac{f(1)}{\epsilon} + \int_0^1 dx \frac{f(x) - f(1)}{1-x} + \mathcal{O}(\epsilon)$$

- Extraction of the infrared and collinear divergences using a proper subtraction of the integrand
- Integration by parts due to numerical instabilities at the thresholds

$$M_{HH}^2 > 4m_Q^2 \Rightarrow m_Q^2 \rightarrow m_Q^2(1 - i\bar{\epsilon}) \quad \text{wi } \bar{\epsilon} \ll 1$$

$$\int_0^1 dx \frac{f(x)}{(a+bx)^3} = \frac{f(0)}{2a^2} - \frac{f(1)}{2(a+b)^2} + \int_0^1 dx \frac{f'(x)}{2(a+bx)^2}$$

different $\bar{\epsilon} \rightarrow$ Richardson extrapolation $\rightarrow \bar{\epsilon} = 0$ (NWA)

Renormalisation

α_s and m_q need to be renormalised

→ α_s in \overline{MS} with $N_F = 5$

→ m_q on shell

$$\delta\sigma = \delta\alpha_s \frac{\delta\sigma_{LO}}{\delta\alpha_s} + \delta m_t \frac{\delta\sigma_{LO}}{\delta m_t}$$

Subtraction of the heavy-top limit → virtual mass effects only (infrared finite)

$$\Delta C_{mass} = C^0 - C_{HTL}^0$$

Adding back the results of HPAIR (heavy-top limit)

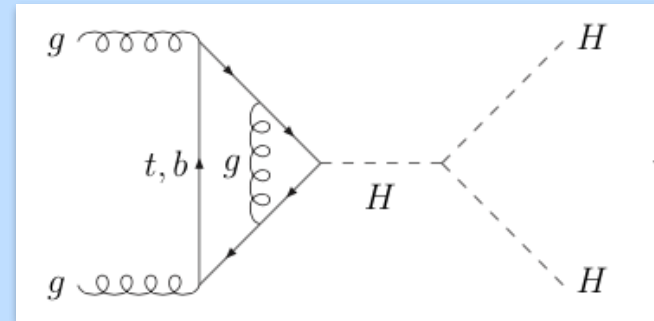
$$C = C_{HTL} + \Delta C_{mass}$$

↑
HPAIR

Remaining Diagrams

Triangular diagrams

← single Higgs case

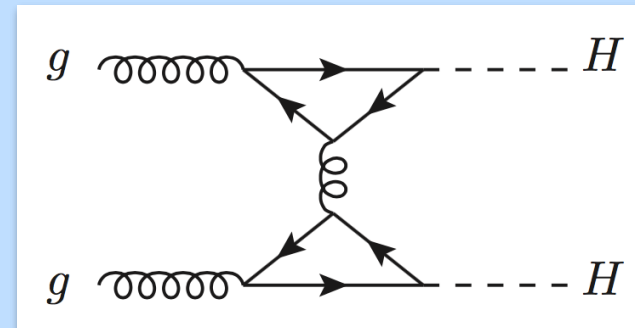


One-particle reducible diagrams

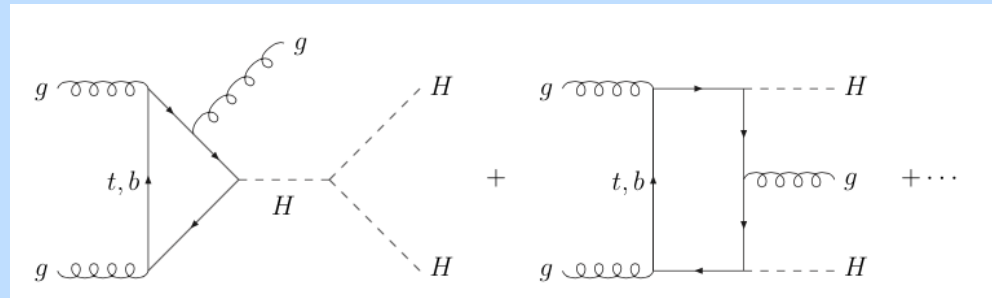
→ analytical results for $C_{\Delta\Delta}$

$(H \rightarrow Z\gamma)$

see e.g. [Degrassi, Giardino, Gröber](#)



Processes: $gg, q\bar{q} \rightarrow HHg$; $gq \rightarrow HHq$



- Full matrix elements generated with LoopTools
- Matrix elements in the heavy-top limit rescaled locally by massive LO matrix elements \rightarrow subtracted \rightarrow free of divergences
- Adding back the results of HPAIR (heavy-top limit)
- Integration of real corrections straight forward

Grid for differential cross section $\frac{d\sigma}{dQ^2}$

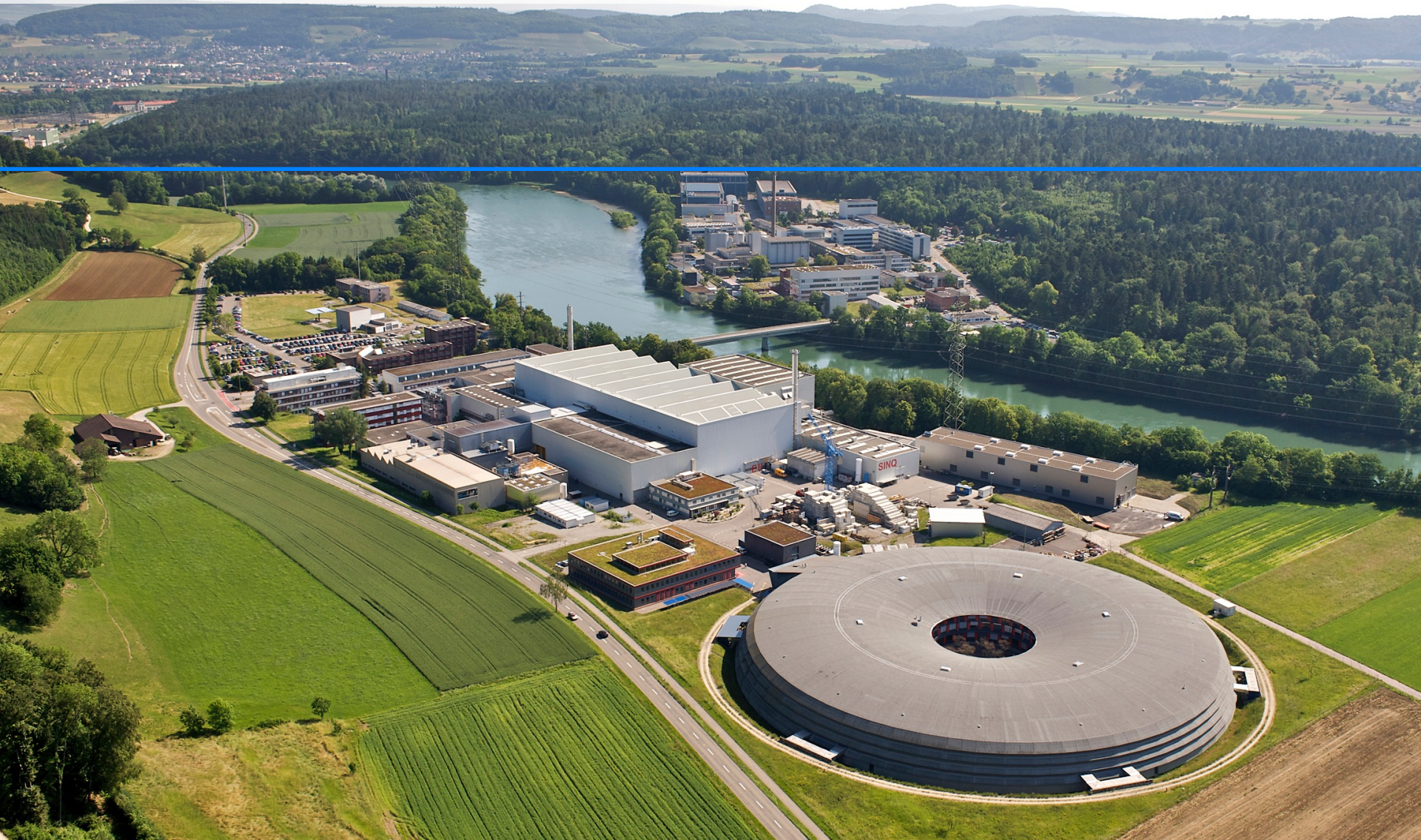
- first with 50 GeV steps in Q (up to about 1 – 1.5 TeV)
- gradual refinement of the grid

Cross checks

- two independent codes using different Feynman parametrisations
- cancellation of divergencies
- comparison with heavy top limit
- comparison with full top mass dependent results of Munich group
- comparison with previously available real corrections

(Frederix, Frixione, Hirschi, Maltoni, Mattelaer, Torrielli, Vryonidou, Zaro)

- Managed to calculate two-loop integral with three free parameter ratios without reduction to master integrals
- Full NLO calculation close to finalisation
- Only matter of CPU time until we get differential cross sections
- NLO mass effects expected in the 10-20% range
- Compared to Munich group:
 - completely different methods
 - no fixed masses \longrightarrow easy to implement arbitrary top and Higgs masses
 - bottom-loops (presently checking stability)



Divergences

- Extraction of the ultraviolet divergences of the matrix elements using endpoint subtractions of the 6-dimensional Feynman integrals

$$\int_0^1 dx \frac{f(x)}{(1-x)^{1-\epsilon}} = \int_0^1 dx \frac{f(1)}{(1-x)^{1-\epsilon}} + \int_0^1 dx \frac{f(x) - f(1)}{(1-x)^{1-\epsilon}}$$

$$= \frac{f(1)}{\epsilon} + \int_0^1 dx \frac{f(x) - f(1)}{1-x} + \mathcal{O}(\epsilon)$$

- Extraction of the infrared and collinear divergences using a proper subtraction of the integrand

denominator: $N = ar^2 + br + c$ $a = \mathcal{O}(\rho)$ $\rho_s = \frac{\hat{s}}{m_Q^2}$

$N_0 = br + c$ $b = 1 + \mathcal{O}(\rho)$

$c = -\rho_s x(1-x)(1-s)t$

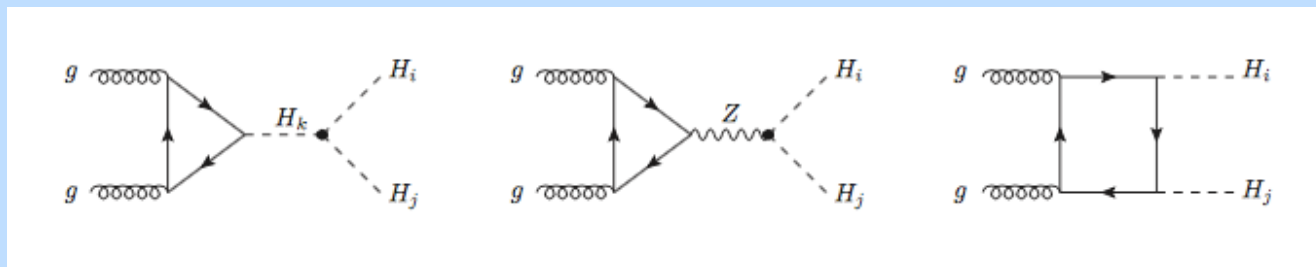
$$\int_0^1 d\vec{x} dr \frac{rH(\vec{x}, r)}{N^{3+2\epsilon}} = \int_0^1 d\vec{x} dr \left\{ \left(\frac{rH(\vec{x}, r)}{N^{3+2\epsilon}} - \frac{rH(\vec{x}, 0)}{N_0^{3+2\epsilon}} \right) + \frac{rH(\vec{x}, 0)}{N_0^{3+2\epsilon}} \right\}$$

↑
Taylor expansion in ϵ

↑
analytical r-integration

Extension to BSM-Higgs scenarios

- Example 2: 2HDM \rightarrow h, H, A
pairs of different Higgs bosons
mixing \rightarrow modified Yukawa couplings
pseudoscalar: treatment of γ_5 (e.g. 't Hooft-Veltman, Larin, Kreimer)
resonant contributions: $H \rightarrow hh, AA$



→ Use Vegas for numerical integration P. Lepage

→ Calculation of differential cross section $\frac{d\sigma}{dQ^2}$, ($Q^2 = m_{HH}^2$)

$$Q^2 \frac{d\Delta\sigma_{virt}}{dQ^2} = \tau \frac{d\mathcal{L}^{gg}}{d\tau} \hat{\sigma}_{virt}(Q^2) \Big|_{\tau = \frac{Q^2}{s}}$$

↑
partonic cross section
(seven-dimensional integrals)

→ Thresholds: $m_Q^2 \rightarrow m_Q^2(1 - i\bar{\epsilon})$
different $\bar{\epsilon} \rightarrow$ Richardson extrapolation $\rightarrow \bar{\epsilon} = 0$ (NWA)

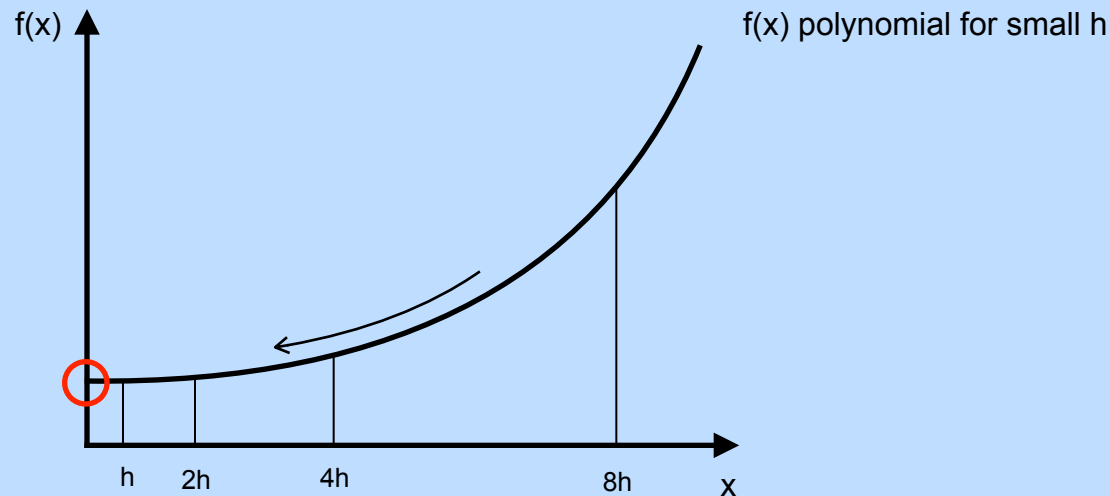
→ Integration of real corrections straight forward

Richardson Extrapolation

$$M_2[f(h), f(2h)] = 2f(h) - f(2h) = f(0) + \mathcal{O}(h^2)$$

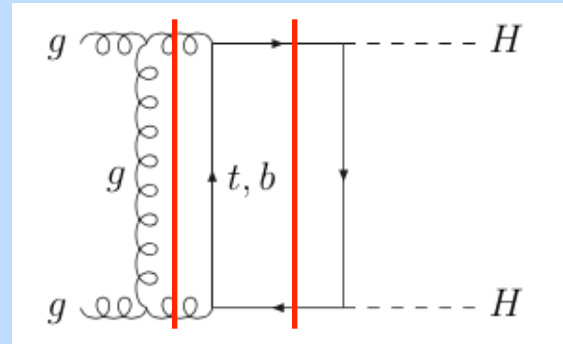
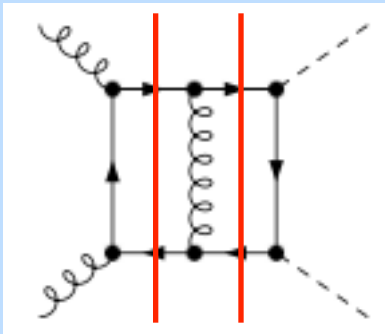
$$M_4[f(h), f(2h), f(4h)] = (8f(h) - 6f(2h) + f(4h))/3 = f(0) + \mathcal{O}(h^3)$$

$$M_8[f(h), f(2h), f(4h), f(8h)] = (64f(h) - 56f(2h) + 14f(4h) - f(8h))/21 = f(0) + \mathcal{O}(h^4)$$



Divergences

- gluon rescattering: threshold for $M_{HH}^2 > 0$

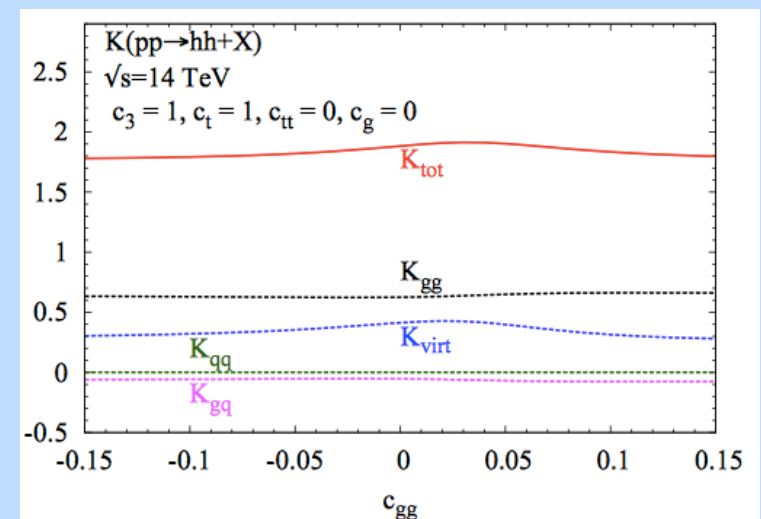
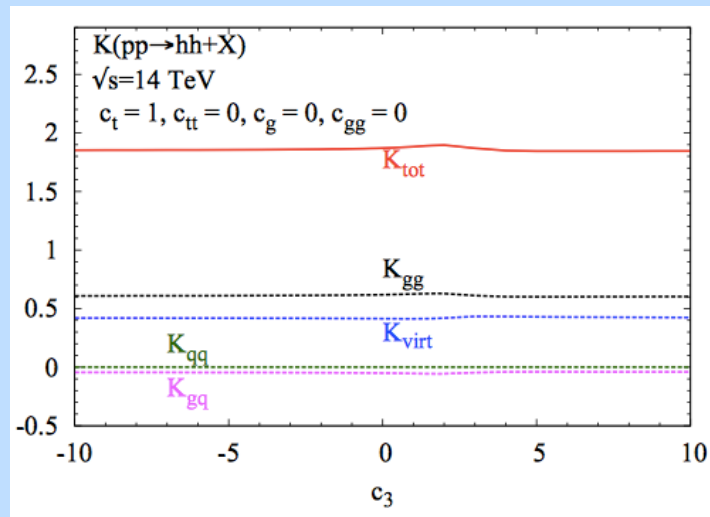
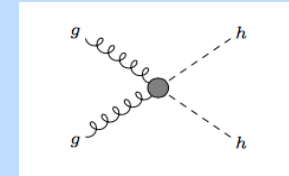
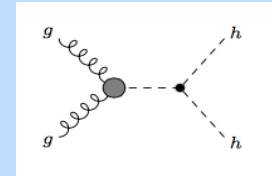
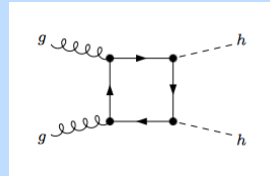
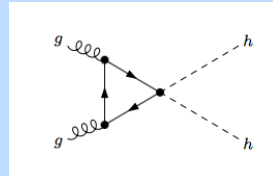
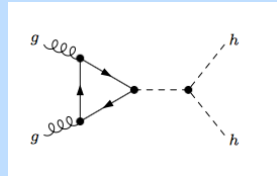


- tried contour deformation
- several 'proper' subtractions of the integrand

Extension to BSM-Higgs scenarios

- Example 1: dim6

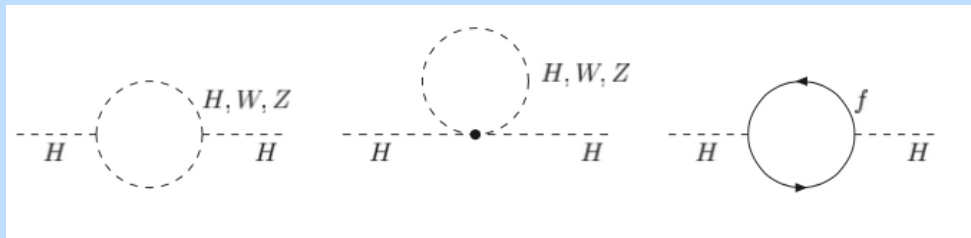
$$\mathcal{L}_{eff} = -m_t \bar{t}t \left(c_t \frac{h}{v} + c_{tt} \frac{h^2}{2v^2} \right) - c_3 \frac{1}{6} \left(\frac{3M_h^2}{v} \right) h^3 + \frac{\alpha_s}{\pi} G^{a\mu\nu} G_{\mu\nu}^a \left(c_g \frac{h}{v} + c_{gg} \frac{h^2}{2v^2} \right)$$



Gröber, Mühlleitner, Spira, Streicher

Hierarchy problem:

Quantum corrections to Higgs mass:



$$\Delta M_H^2 \sim \Lambda^2 \sim \mathcal{O}(M_{GUT}^2) \gg M_H^2$$

mass difference can be absorbed in counter term:

$$M_H^2 \rightarrow M_H^2 + \Delta M_H^2 - \delta M_H^2$$

fine tuning of ~ 28 digits