

# High-energy ep/eA collisions with the eSTARlight Monte Carlo

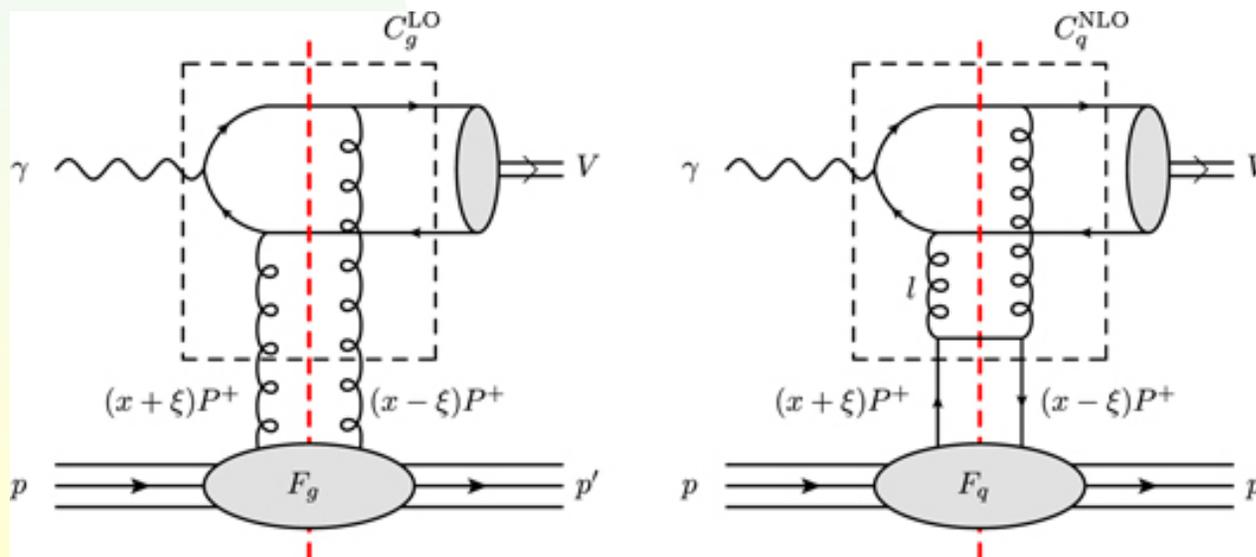
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Presented at the workshop on “Electrons for the LHC - LHeC/FCCeh and Perle Workshop”

- Vector meson photo & electro-production
- The eSTARlight Monte Carlo
- ep and eA collisions
  - ◆ Current Status
  - ◆ Future plans
- Conclusions

# Exclusive vector meson photo & electro production

- An important probe of nuclear structure
- LHeC/FCC-eh probe gluons at very low Bjorken-x
  - ◆ Proton or nuclear targets
- Proceeds via Pomeron or Reggeon (meson) exchange
  - ◆ Pomeron  $\sim\sim$  2 gluons  $\sim\sim$  BFKL gluon ladder
    - ✦ Directly sensitive to gluon distribution
    - ✦ 2<sup>nd</sup> gluon complicates gluon density extraction
      - Much discussion....



Jones, Martin, Ryskin, Teubner

# From now to LHeC & FCC-eh

- By the time LHeC or FCC-eh run, we will have good data on vector meson photoproduction cross-sections at a wide range of energies from UPCs
- ep/A colliders can focus on more data-hungry analyses, and those requiring precise control of  $Q^2$ , like measuring the spatial dependence of shadowing.
  - ◆ Scan in  $Q^2$  with  $M_\nu$  fixed
  - ◆ Larger reach in  $t$  -> better measurement of the transverse interaction profile (like GPD, but for nuclei)
  - ◆ Better (but not perfect) separation of coherent and incoherent photoproduction.
    - ◆ Nuclear excitation in an independent reaction
      - Growing importance at higher luminosities...
    - ◆ Missing photons from nuclear de-excitation

# eSTARlight

- Monte Carlo for photoproduction and electroproduction of vector mesons at an EIC
  - ◆ Here, photoproduction is  $Q^2 < 1 \text{ GeV}^2$ , while electroproduction is  $Q^2 > 1 \text{ GeV}^2$
- Physics model follows STARlight UPC event generator, but covers photons with arbitrary  $Q^2$
- A fast, complete, reasonably accurate model of vector meson production, not a sophisticated theoretical calculation
  - ◆ For detector simulations....
  - ◆ Electron (or positron)  $\rightarrow \gamma^* \rightarrow$  vector meson  $\rightarrow$  final state
  - ◆ Vector meson polarization and decay angular distribution
  - ◆ Based on data where possible, phenomenology elsewhere
    - ✦ \* Some extrapolations required
- Designed to be easily extensible

# Initial and final states

- Electron (or positron) on protons/ light ions ( $Z < 7$ , Gaussian distribution)/ heavy ions (Woods-Saxon distribution)
- Final states:  $\rho$ ,  $\omega$ ,  $\phi$ ,  $\rho'$  (i. e.  $\pi\pi\pi\pi$ ),  $\rho$  + direct  $\pi\pi$ , with interference,  $J/\psi$ ,  $\psi'$ ,  $Y(1S)$ ,  $Y(2S)$ ,  $Y(3S)$  states
  - ◆ Simple states decayed in STARlight
  - ◆ Complex final states via PYTHIA interface
  - ◆ Easily extensible
- Incoherent photonuclear interactions w/ DPMJET
  - ◆ Real photon approximation
- eSTARlight tracks outgoing electron & proton/nucleon
- eSTARlight outputs photon 4-vector

# Electronuclear interactions

$$\sigma(e + X \rightarrow e + X + V.M.) = \int dQ^2 \int dE_\gamma \frac{dN_\gamma(E_\gamma, Q^2)}{dE_\gamma dQ^2} \sigma_{\gamma X}(W, Q^2)$$

- Convolution of photon flux from electron with cross-section; both depend on  $Q^2$
- Photon flux depends on virtuality

$$\frac{d^2 N}{d(Q^2) dE_\gamma} = \frac{\alpha}{\pi} \frac{1}{E_\gamma |Q^2|} \left[ 1 - \frac{E_\gamma}{E_e} + \frac{1}{2} \left( \frac{E_\gamma}{E_e} \right)^2 - \left( 1 - \frac{E_\gamma}{E_e} \right) \left| \frac{Q_{min}^2}{Q^2} \right| \right]$$

# Cross-sections

- Parameterized from HERA data

$$\sigma_{\gamma p}(W) = \sigma_P \cdot W^\epsilon + \sigma_M \cdot W^\eta$$

$$\sigma_{\gamma p} = \left( \frac{1}{1 + Q^2/M_v^2} \right)^n \sigma_{\gamma p}(W)$$

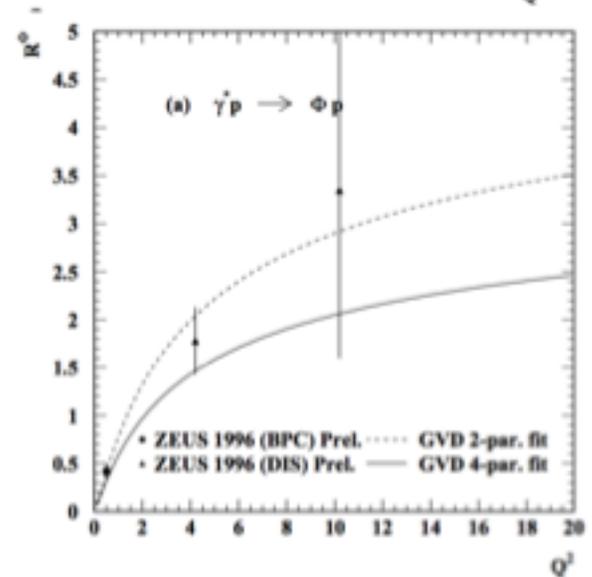
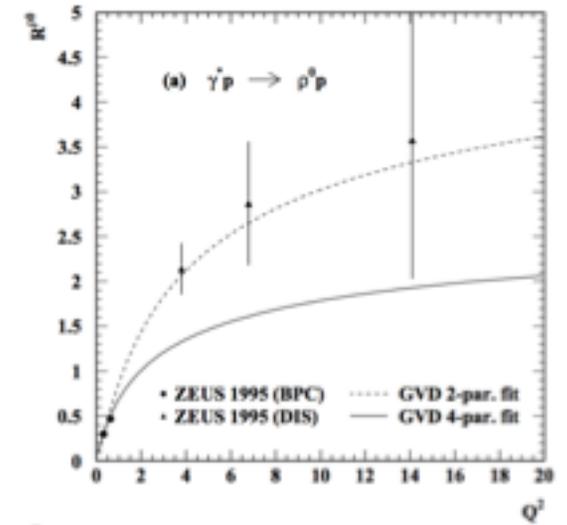
- $n = c_1 + c_2(Q^2 + M_v^2)$
- $\sigma_{\gamma p}$  parameterized from HERA data

Meson	$c_1$	$c_2$ ( $10^{-2} \text{GeV}^{-2}$ )
$\rho$	$2.09 \pm 0.10$	$0.73 \pm 0.18$
$\phi$	$2.15 \pm 0.17$	$0.74 \pm 0.46$
$J/\psi$	$2.36 \pm 0.20$	$0.29 \pm 0.43$

- ◆ Pomeron exchange + Reggeon exchange
  - ✦ Reggeon exchange matters for  $\rho/\omega$ , at lower photon energies
- $Q^2$  dependence included via a power-law
  - ◆ Data on power  $n$  is not available for all mesons; we use the 'closest' meson
- More accurate parameterization used for heavy mesons, to better model near-threshold production

# Vector meson decays

- Vector mesons retain the spin of the incident photon
- For  $Q^2 \rightarrow 0$ , s-channel helicity conservation means that the vector mesons are transversely polarized to the beam direction
  - ◆ As  $Q^2$  rises, longitudinal polarization rises
- The  $Q^2$  dependence of the transverse:longitudinal polarization ratio is not well known
- Parameterize HERA data in terms of spin-matrix elements:
- Only known for some mesons; use most 'similar' meson where needed



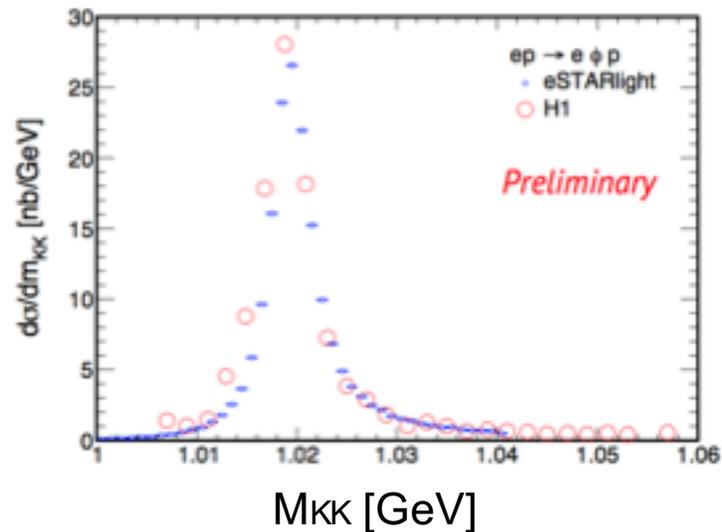
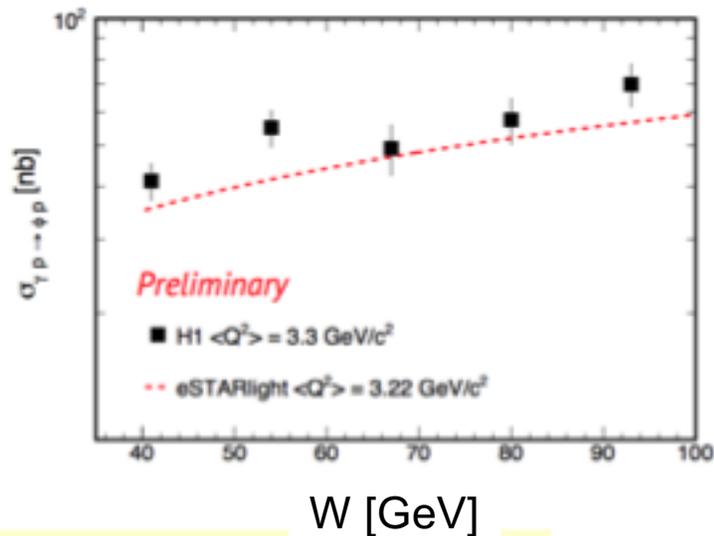
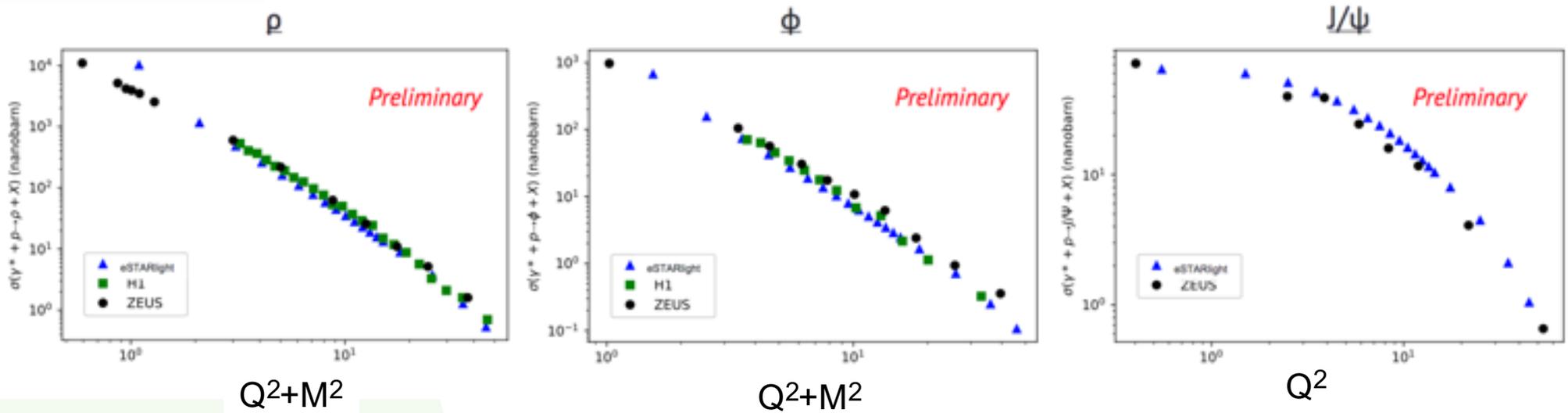
$$R_v = \frac{1}{\epsilon} \frac{r_{00}^{04}}{1 - r_{00}^{04}}$$

# Comparison with HERA data

- HERA shows  $\gamma^*p$  cross-sections

- Remove the photon flux from the eSTARlight calculations

$$\sigma_{\gamma p} = \frac{\int dE_\gamma \int dQ^2 \frac{d^2 N}{dE_\gamma d(Q^2)} \sigma_{\gamma p}(E_\gamma, Q^2)}{\int dE_\gamma \int dQ^2 \frac{d^2 N}{dE_\gamma d(Q^2)}}$$



# From $\gamma p$ to $\gamma A$

- With a quantum Glauber calculation, generalized vector meson dominance and the optical theorem:

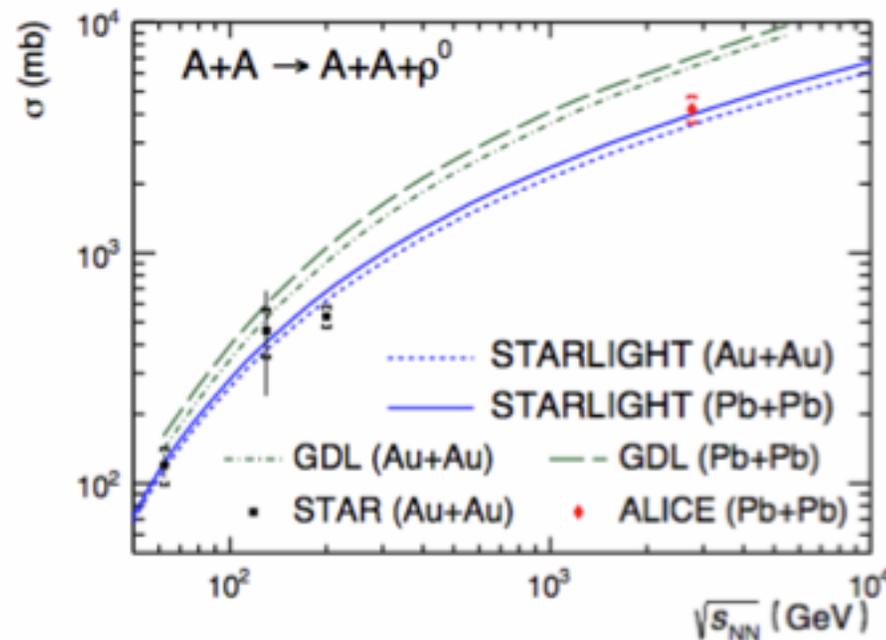
$$\sigma_{tot}(VA) = \int d^2b \left[ 2 \cdot \left( 1 - e^{-\sigma_{tot}(Vp)T_{AA}(b)/2} \right) \right]$$

$$\sigma(\gamma A \rightarrow VA) = \left. \frac{d\sigma(\gamma A \rightarrow VA)}{dt} \right|_{t=0} \int_{t_{min}}^{\infty} dt |F(t)|^2$$

- For heavy mesons (small dipoles),  $d\sigma/dt|_{t=0} \sim A^2$
- For the  $\rho^0$  (smallish dipoles),  $d\sigma/dt|_{t=0} \sim A^{4/3}$

# Glauber calculations

- Quantum Glauber calculation does not match STAR and ALICE UPC data; a classical Glauber does well.
  - ◆ Can add a correction for nuclear inelastic shadowing
  - ◆ eSTARlight uses quantum Glauber; classical Glauber is an option
  - ◆ Only important for  $\rho/\omega$



ALICE, JHEP 1509, 095 (2015).

L. Frankfurt et al. Phys. Lett. **B752**, 51 (2018)

# Accelerator parameters

- The calculations that follow use:

Accelerator	Collision System	Electron Energy	Ion Energy
eRHIC	ep eA	18 GeV	275 GeV 100 GeV/A
JLEIC	ep eA	10 GeV	100 GeV 40 GeV/A
LHeC	ep eA	60 GeV	7 TeV 2.8 TeV/A
Fcc-eh	ep eA	60 GeV	50 TeV 20 TeV/A
HERA	ep	27.5 GeV	920 GeV

I will not cover PERLE here, but 1 GeV electrons colliding with 7 TeV protons is at 167 GeV center of mass energy. This can be interesting, but the extreme energy imbalance would require a very backward detector.

# Rates at EICs

- Assumed integrated luminosity  $10 \text{ fb}^{-1}/\text{A}$

		Photo-production ( $Q^2 < 1 \text{ GeV}^2$ )					Electro-production ( $Q^2 > 1 \text{ GeV}^2$ )				
LHeC	ep	100 G	5.6 G	470 M	78 M	1.2 M	260 M	37 M	29 M	6.3 M	180 K
	ePb	110 G	8.2 G	720 M	140 M	2.0 M	100 M	16 M	27 M	7.2 M	250 K

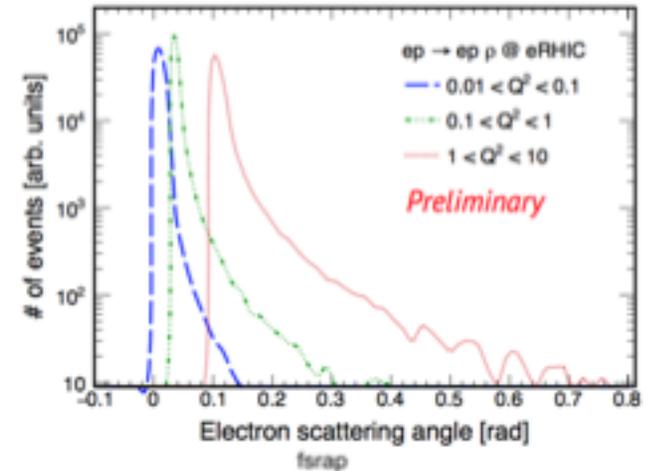
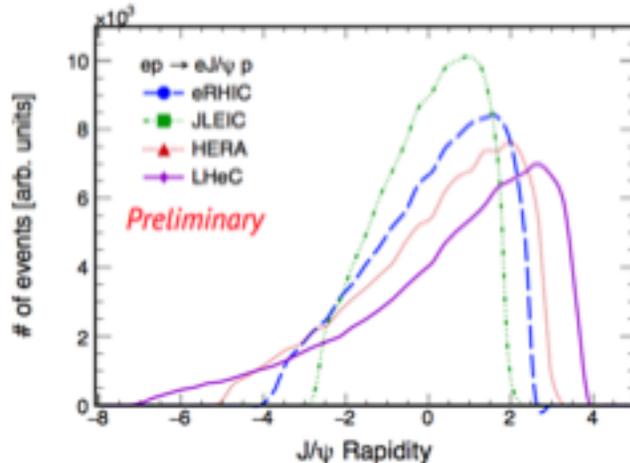
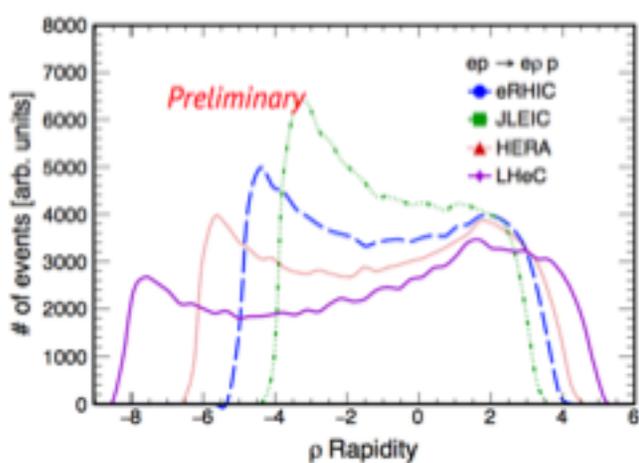
- Photoproduction at LHeC
  - High rates ( $>10^9/\text{year}$ ) for light mesons
  - Good rates ( $>10^6/\text{year}$  for c-cbar,  $>10^5/\text{year}$  for b-bbar)
- Electroproduction
  - Rates from  $\sim <1\%$  of photoproduction (light mesons), rising to 15% of photoproduction rates for the Upsilon
- FCC-eh luminosity  $15^*$  LHeC,  $\sigma(\gamma p \rightarrow J/\psi p)$   $12^*$  higher
  - Rate 180 times higher

# Implication for physics program

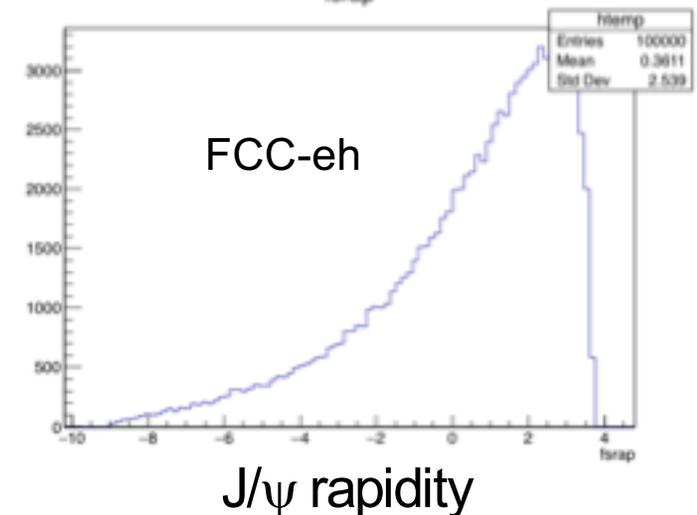
- FCC-eh & LHeC rates allow multi-dimensional binning and tomographic studies for light,  $c$ - $c$ bar &  $b$ - $b$ bar mesons.
  - ◆ At LHeC, statistics may be somewhat limited for  $b$ -bar mesons.
  - ◆  $\psi(3770)$ ,  $\psi(4040)$  (even with the small branching ratios)
  - ◆ Exotica, including, at FCC-eh  $b$ - $b$ bar exotica
- A host of  $\rho'$ ,  $\omega'$ , and  $\phi'$ , etc. states should be accessible
  - ◆ Meson spectroscopy & rare decays, and probe nucleons with different dipoles
- At FCC-eh & LHeC, exclusive  $t$ - $t$ bar production should be visible
  - ◆ Top quarks decay before toponium can form
  - ◆ At FCC-eh,  $t$ - $t$ bar dipole lives long enough to scatter from target
  - ◆ Usual VM width based cross—section calculation doesn't work
  - ◆ Worth further investigation

# Rapidity and Angular distributions

- Vector meson production over a wide rapidity range
  - N. b. unscaled distributions here



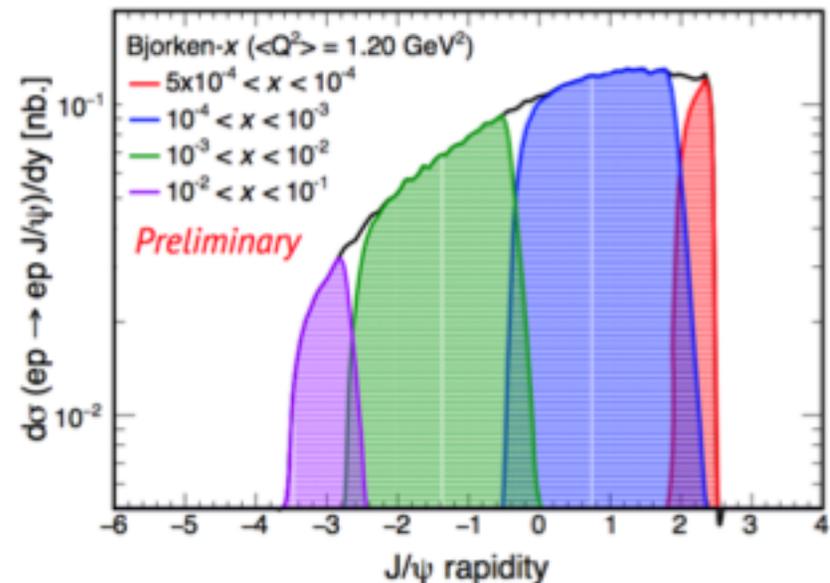
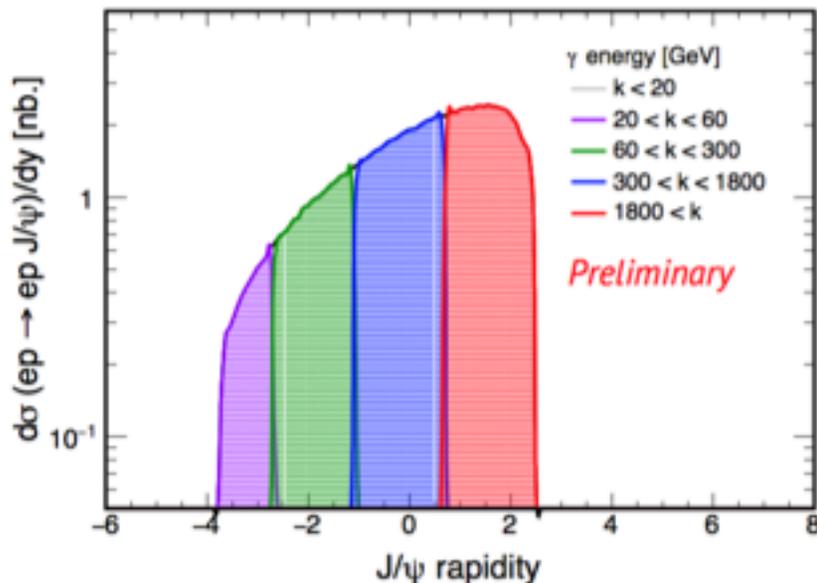
- $\rho^0$  'double peak' is from Reggeon exchange (near threshold) and Pomeron exchange at large  $k$ /rapidity
  - $\phi$  or  $J/\psi$  are not produced via Reggeon exchange
- Electrons scattering angle is small (no surprise)



Rapidity peak similar to LHeC;  
Tail at larger  $-y$

# ep production vs. photon energy, Bjorken-x

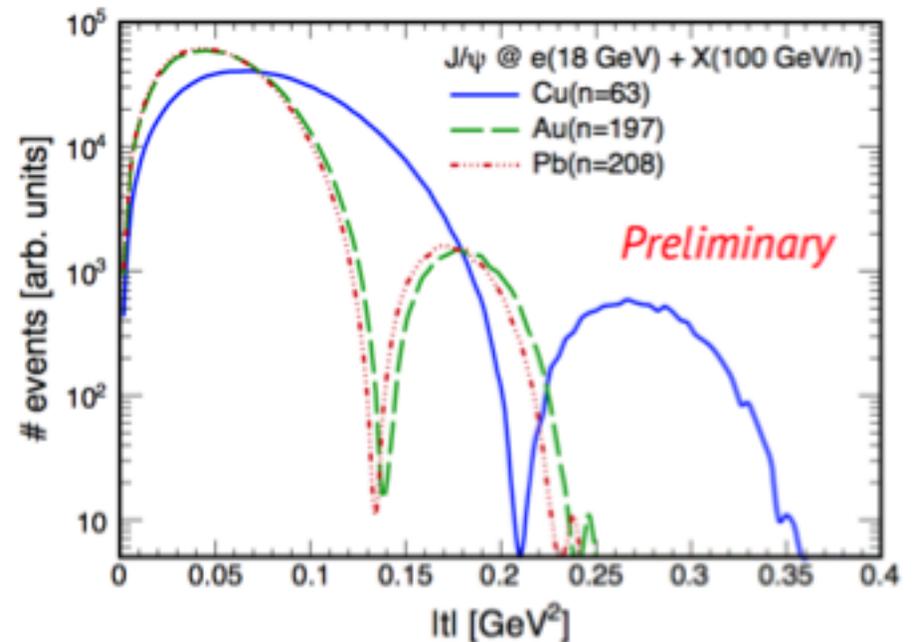
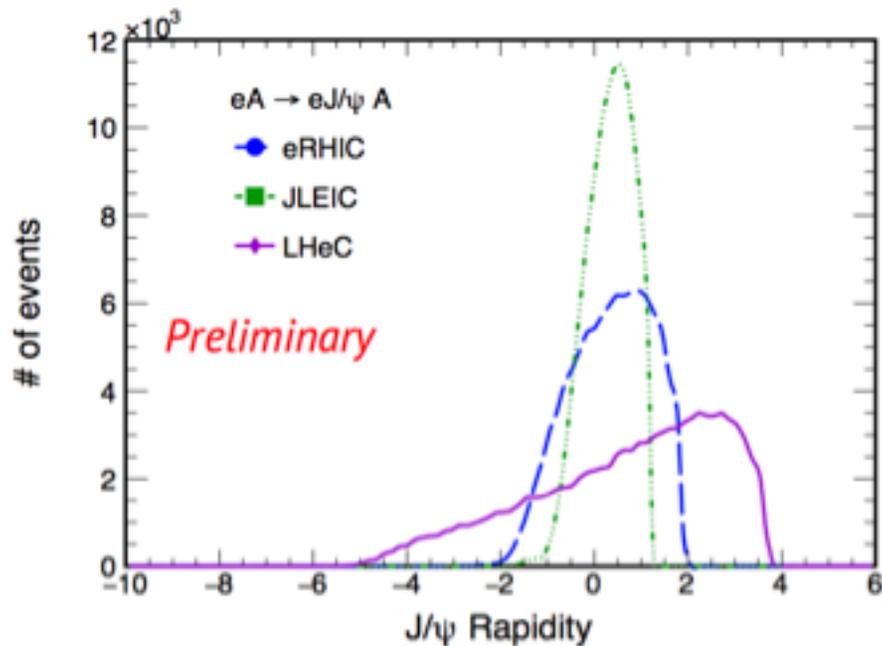
- Photon energy maps into rapidity
  - ◆ For photoproduction,  $k = Mv/2 \ln(y)$
  - ◆ Electroproduction shifts this slightly to the right



- Photon energy also maps onto Bjorken-x
- For maximum energy/Bjorken-x reach, need to detect vector mesons forward, with  $y \sim 2.5$
- Near threshold, production is at large negative rapidity
  - ◆ Could shift to mid-rapidity by lowering beam energy

# Production in eA

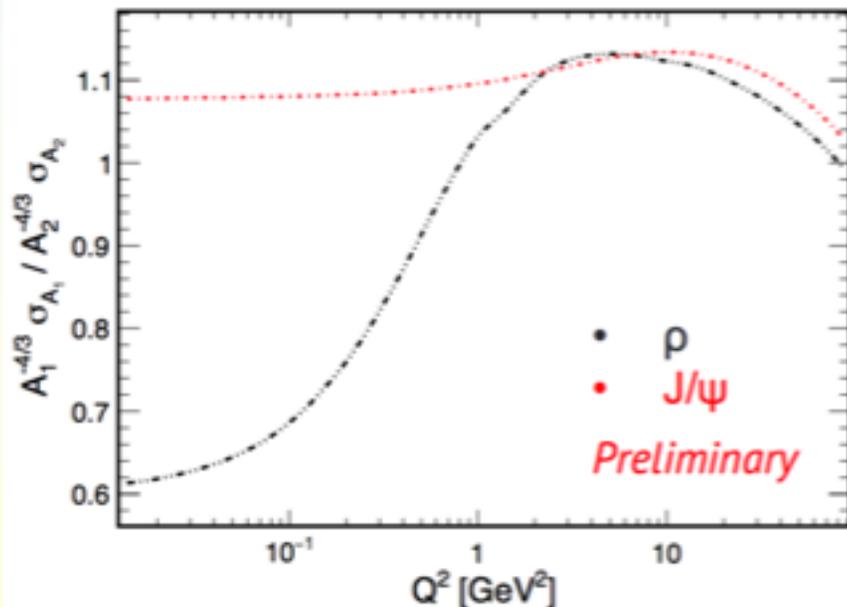
- Smaller  $\gamma$ -nucleon center of mass energy
  - ◆ Narrower rapidity range
- Lower Pomeron  $p_z$   $\rightarrow$  production is more central
- Expect clean diffractive minima
  - ◆ Unlike in UPCs, photon momentum can be removed



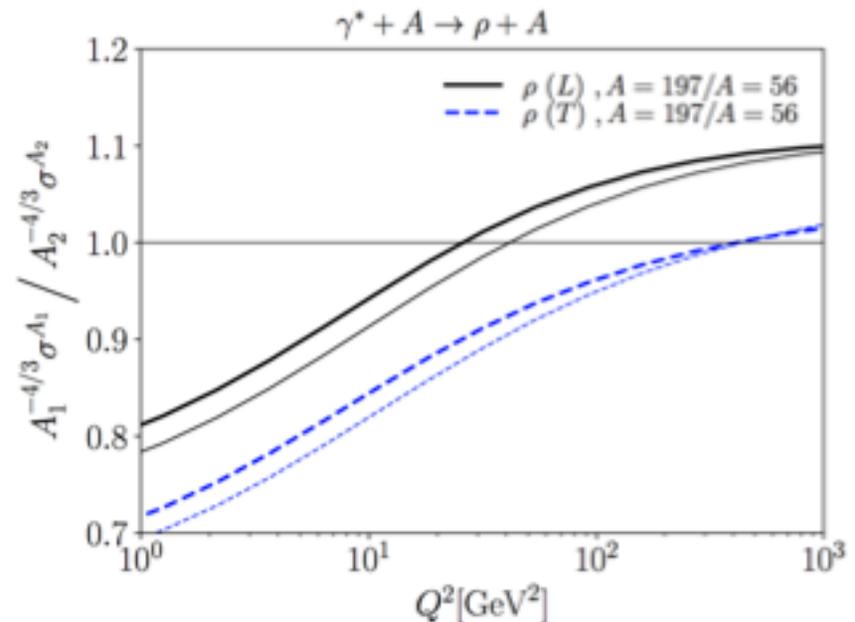
# Cross-section vs. A & Q<sup>2</sup>: shadowing

- Without shadowing (i. e. for small dipoles),  $\sigma \sim A^{4/3}$ 
  - ◆  $A^2$  for forward scattering cross-section,  $A^{-2/3}$  for phase space
  - ◆ With shadowing, the growth in  $\sigma$  with A is smaller
- eSTARlight reproduces this well
- Shadowing should increase slowly with increasing energy
  - ◆ Eventually – reach ‘black disk’ limit

eSTARlight

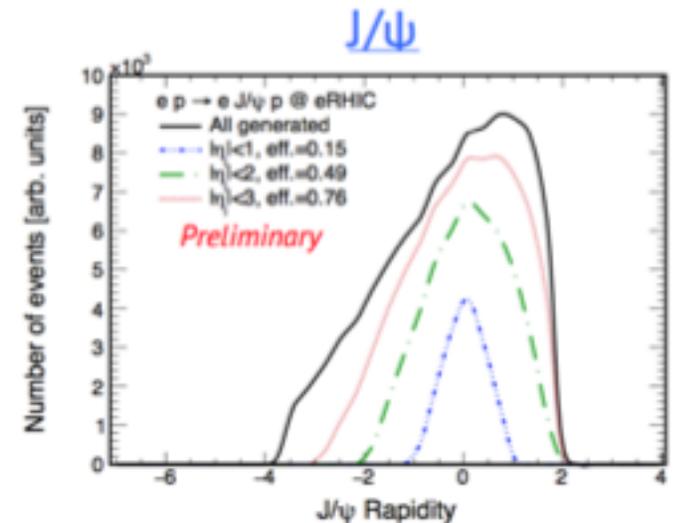
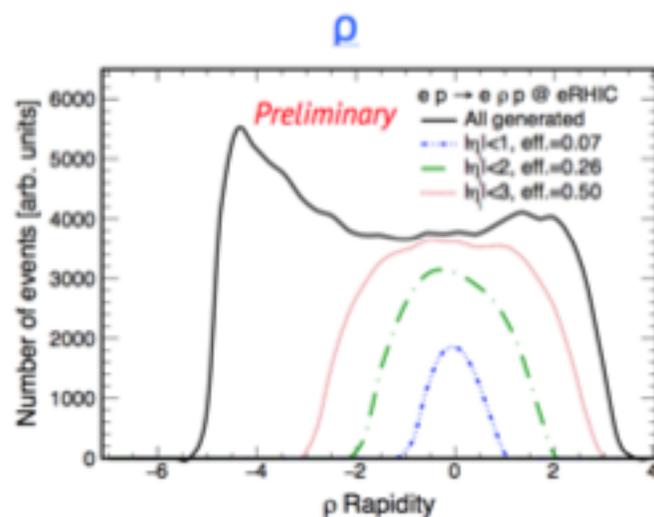
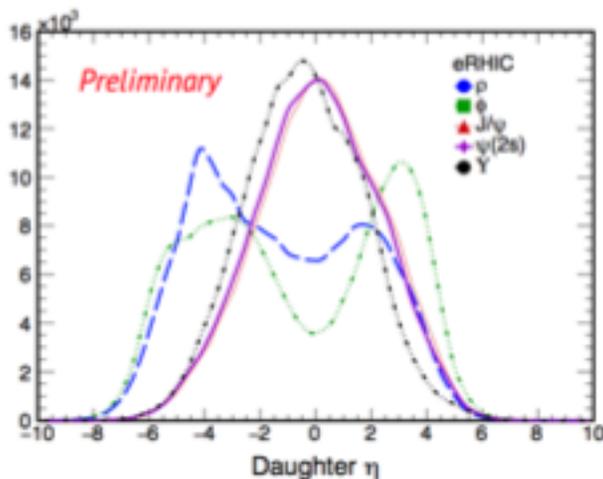


Mantysaari & Venugopalan



# Final state particle distributions

- The vector meson daughter particles generally follow the rapidity distribution of their vector meson parents
- The final state matters: VM  $\rightarrow$  spin 0 spin – (e. g.  $\pi\pi$ ) has a very different angular distribution from VM  $\rightarrow$  spin  $\frac{1}{2}$  spin  $\frac{1}{2}$ 
  - ◆ Clebsch Gordon coefficients



- Large detector acceptance is key to high acceptance.
  - ◆ Otherwise, we waste beam

# eSTARlight code base & future plans

- Straightforward C++ code hosted on hepforge
  - ◆ <http://estarlighthepforge.net>
  - ◆ Easy to download and install
  - ◆ If you need a hepforge account, please request one
  - ◆ Please try it, and provide feedback
- Optional inclusion of PYTHIA8 (for complex decays) and DPMJET3 for arbitrary eA interactions (w/ $Q^2=0$  for DPMJET)
- Future plans: additional mesons, charge exchange reactions ( $\gamma p \rightarrow X^+ n$ ), exotica, top quarks, spin effects, GPDs?
- We welcome interested parties as co-developers

# Conclusions

- UPCs at hadron colliders and an EIC are complementary. UPCs have a larger photon energy/Bjorken- $x$ , but lack good control of  $Q^2$
- The EIC will also offer the luminosity to collect enormous data samples  $d\sigma_{\text{coherent}}/dpt^2$ , to study the effective shape of the nucleus, as a function of  $Q^2$
- STAR has made a preliminary study of shape changes with varying  $Q^2$ , using dipion  $M_{\pi\pi}$  to select events with different dipole size
- We have developed the eSTARlight Monte Carlo event generator which simulates production of vector mesons at an EIC
  - ◆ It covers arbitrary ranges of  $Q^2$
  - ◆ Initial runs show the importance of a wide detector acceptance. Forward acceptance is needed to probe the highest energy photons
- The eSTARlight code is available on hepforge. Please try it.
- We welcome both feedback and co-development efforts to add features to the code.