

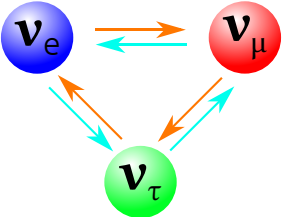
Review on Sterile Neutrinos

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Electrons for the LHC workshop, Orsay
Thursday, June 28, 2018

Neutrino oscillations & the Standard Model



Three Generations of Matter (Fermions) spin 1/2

	I	II	III	
mass -	2.4 MeV	1.27 GeV	173.2 GeV	0
charge -	2/3	2/3	2/3	0
name -	u up	c charm	t top	g gluon
	Left Right	Left Right	Left Right	0
Quarks	d down	s strange	b bottom	γ photon
	Left Right	Left Right	Left Right	0
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	Z weak force
	0	0	0	91.2 GeV
Leptons	e electron	μ muon	τ tau	H Higgs boson
	Left Right	Left Right	Left Right	126 GeV
	-1	-1	-1	0
	0.511 MeV	105.7 MeV	1.777 GeV	0
	Left Right	Left Right	Left Right	0
	0	0	0	spin 0
	0	0	0	80.4 GeV
	0	0	0	W weak force
	0	0	0	+1
	0	0	0	

courtesy M. Shaposhnikov

- ▶ No right-handed neutrinos in the Standard Model (SM).
 - ▶ No mass matrix, no mixing of the neutrino flavour states.
- ⇒ Neutrino oscillations are evidence of physics beyond the SM.

Lowscale seesaw

Benchmark model, defined in Antusch, OF; JHEP **1505** (2015) 053

Similar to e.g.: Mohapatra, Valle (1986); Malinsky, Romao Valle (2005); Shaposhnikov (2007);

- ▶ Lowscale seesaw Lagrangian, two sterile neutrinos N_i with protective symmetry:

$$\mathcal{L}_N = -\frac{1}{2} \overline{N_R^1} M (N_R^2)^c - y_{\nu\alpha} \overline{N_R^1} \tilde{\phi}^\dagger L^\alpha + \text{H.c.}$$

- ▶ The mass matrix after electroweak symmetry breaking:

$$M_\nu = \begin{pmatrix} 0 & m_D & m'_D \\ (m_D)^T & 0 & M \\ (m'_D)^T & M & \mu \end{pmatrix},$$

- ▶ **Perturbations** $\Rightarrow m_\nu$ and HNL mass splitting (ΔM)
- ▶ m'_D : Linear seesaw, $\Delta M^{\text{NO}} = 0.0416 \text{ eV}$, $\Delta M^{\text{IO}} = 0.000753 \text{ eV}$
- ▶ μ : inverse seesaw, $\Delta M \sim \frac{m_{\nu_j}}{|\theta|^2}$.

Heavy neutrino interactions

- ▶ Heavy neutrinos are mostly sterile, interactions via mixing.
- ▶ **Charged current (CC):**

$$j_{\mu}^{\pm} = \frac{g}{2} \theta_{\alpha} \bar{\ell}_{\alpha} \gamma_{\mu} N$$

- ▶ **Neutral current (NC):**

$$j_{\mu}^0 = \bar{\nu}_{\alpha} \gamma_{\mu} \theta_{\alpha} N$$

- ▶ Higgs boson **Yukawa** interaction:

$$\mathcal{L}_{\text{Yukawa}} = \sum_{\alpha=e,\mu,\tau} \theta_{\alpha} \frac{\sqrt{2} M}{v_{\text{EW}}} \nu_{\alpha} \phi^0 \bar{N}$$

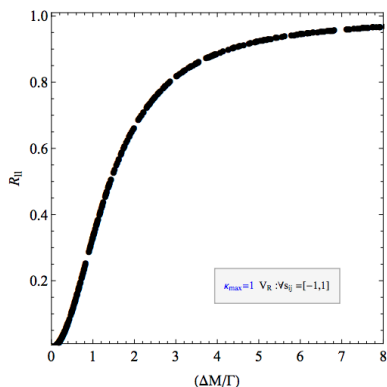
- ▶ Simplification: light neutrino mass eigenstates $\equiv \nu_e, \nu_{\mu}, \nu_{\tau}$

Traditionally searched for via
Lepton Number Violating signatures

e.g. $\mu^\pm\mu^\pm + J$ at pp (SS dimuons) or $e^+ + J$ at ep

However...

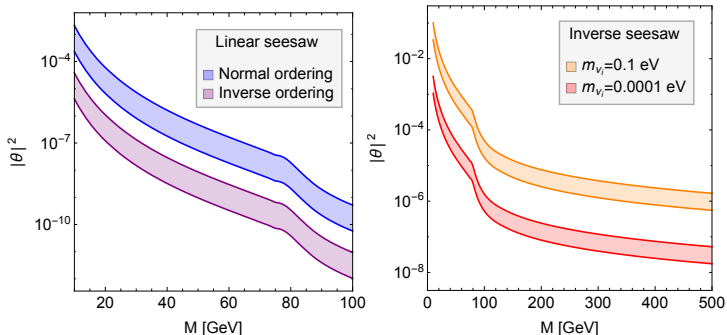
Lepton number violation and ΔM



Anamiati, Hirsch, Nardi; 1607.05641

- ▶ With $m'_D = \mu = 0$, no LNV in this class of models.
- ▶ $R_{\ell\ell}$ (=LNV/LNC) function of ΔM and Γ_N .
- ▶ For mass splitting \sim decay width, $R_{\ell\ell} \in [0, 1]$.
- ▶ Zero mass splitting \Rightarrow zero LNV.

Parameter space with LNV

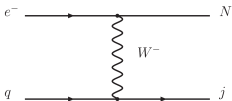


Antusch et al.; [1709.03797]

$$R_{\ell\ell} = \frac{\Delta M^2}{2\Gamma^2 + \Delta M^2}$$

- ▶ The colored bands separate parameterspace w/ and w/o LNV.
 - ▶ Upper contour: $R_{\ell\ell} = 0.1$, lower contour: $R_{\ell\ell} = 0.9$
- ⇒ No LNV for $M > 100$ GeV!

Interlude: Searching for N with $M > 100$ GeV

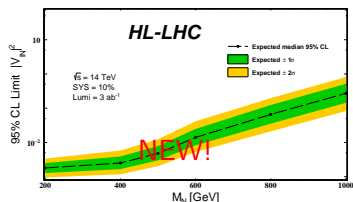


Name	Final State	Channel [production,decay]	$ \theta_\alpha $ dependency	LNV/LFV
lepton-trijet	$jjj\ell_\alpha$	$[\mathbf{W}_t^{(q)}, W]$	$\frac{ \theta_e\theta_\alpha ^2}{\theta^2}$	\checkmark/\checkmark
jet-dilepton	$j\ell_\alpha^+\ell_\beta^-\nu$	$[\mathbf{W}_t^{(q)}, \{W, Z(h)\}]$	$\left\{ \frac{ \theta_e\theta_\alpha ^2}{\theta^2}, \theta_{e1} ^{2(*)} \right\}$	\times/\checkmark
trijet	$jjj\nu$	$[\mathbf{W}_t^{(q)}, Z(h)]$	$ \theta_e ^2$	\times
monojet	$j\nu\nu\nu$	$[\mathbf{W}_t^{(q)}, Z]$	$ \theta_e ^2$	\times

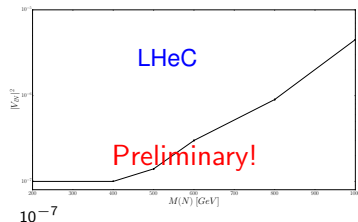
lepton-quadrjet	$jjjj\ell_\alpha$	$[\mathbf{W}_t^{(*)}, W]$	$\frac{ \theta_e\theta_\alpha ^2}{\theta^2}$	\checkmark/\checkmark
dilepton-dijet	$\ell_\alpha\ell_\beta\nu jj$	$[\mathbf{W}_t^{(*)}, \{W, Z(h)\}]$	$\left\{ \frac{ \theta_e\theta_\alpha ^2}{\theta^2}, \theta_{e1} ^{2(*)} \right\}$	\times/\checkmark
trilepton	$\ell_\alpha^-\ell_\beta^-\ell_\gamma^+\nu\nu$	$[\mathbf{W}_t^{(*)}, \{W, Z(h)\}]$	$\left\{ \frac{ \theta_e\theta_\alpha ^2}{\theta^2}, \theta_{e1} ^{2(*)} \right\}$	\times/\checkmark
quadrjet	$jjjj\nu$	$[\mathbf{W}_t^{(*)}, Z(h)]$	$ \theta_e ^2$	\times
lepton-dijet	$\ell_\alpha^- jj\nu\nu$	$[\mathbf{W}_t^{(*)}, Z(h)]$	$ \theta_e ^2$	\times
dijet	$jj\nu\nu\nu$	$[\mathbf{W}_t^{(*)}, Z]$	$ \theta_e ^2$	\times
monolepton	$\ell_\alpha^- \nu\nu\nu\nu$	$[\mathbf{W}_t^{(*)}, Z]$	$ \theta_e ^2$	\times

see Antusch et al.; [1612.02728]

Lepton flavor violation



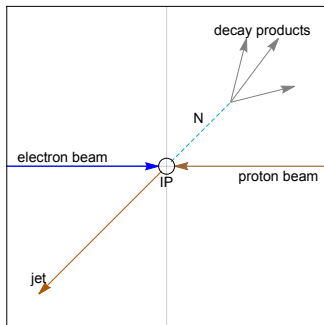
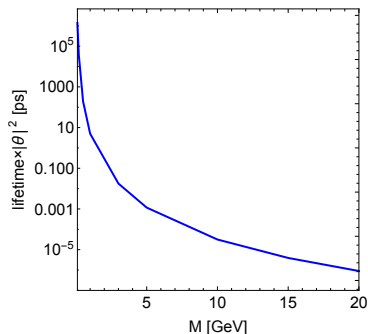
Antusch et al.; [1805.11400]



- ▶ LHC search: $e^\pm \mu^\mp + J$
- ▶ LHeC search: $\mu^- + J$ or $\tau^- + J$
- ▶ Include many bkg for pp, and $(\mu\nu\nu + J)$, +single τ s for ep.
- ▶ Preliminary results for LHeC improve sensitivity significantly.

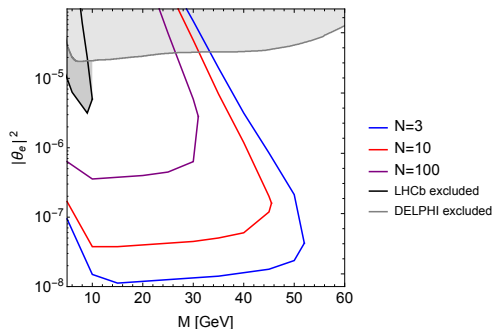
Now: new phenomena from N with $M < m_W$

Displaced vertex searches



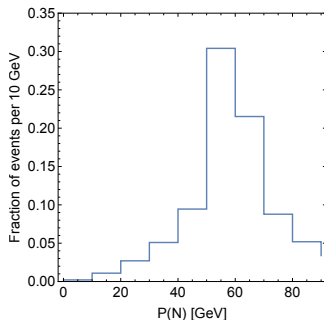
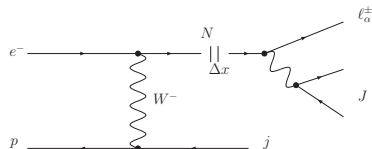
- ▶ $M < m_W$ and $|\theta|^2 < 10^{-5}$ leads to macroscopic lifetimes
- ▶ Interaction point from recoiling jet with high precision
- ▶ Secondary vertex with “large” displacement unique signal
- ▶ “Large”: a few times tracking resolution ($8 \mu\text{m}$).

Sensitivities at the LHeC



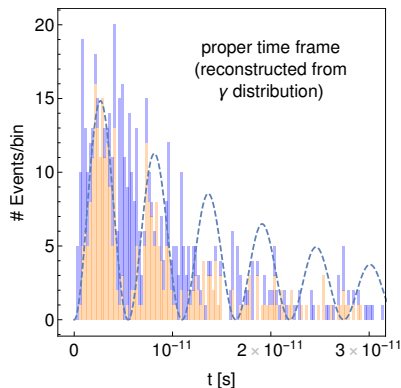
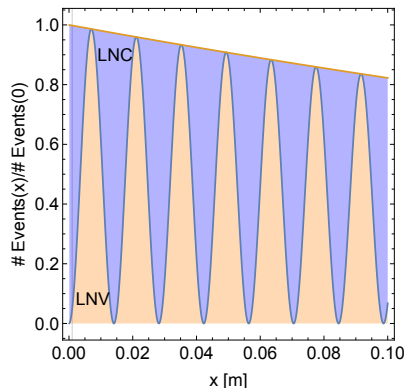
- ▶ Present bounds by LHCb and DELPHI.
- ▶ Large numbers of displaced N decays possible.
- ▶ In principle this parameter space can be probed by LHC.
- ▶ LHeC complementary: testing electron and tau flavor.

Heavy neutrino oscillations



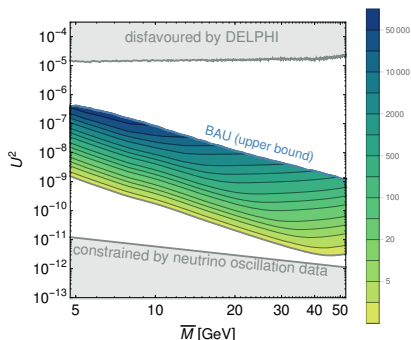
- ▶ Final lepton can be ℓ^+ due to Δx -dependent interference.
- ▶ Interference with oscillation length dependent on theory.
- ▶ Example: $\lambda_{\text{osc}}^{\text{lin,IO}} = 3.29 \cdot 10^{-3} \sqrt{\gamma^2 - 1} \text{ m}$.
- ▶ Lorentz boosts make this oscillation more visible.

Oscillation, a theorist's take



- ▶ Left: distribution of events as a function of displacement.
- ▶ Right: Monte Carlo simulation with 500 events, assuming a reconstruction error of 10%.
- ▶ Both assume a mass of 7 GeV and mixing close to existing bounds.

Why we should care



Antusch, Cazzato, Drewes, Fischer, Garbrecht, Gueter, Klaric; 1710.03744

- ▶ Active-sterile mixing angles θ_α fixed by the light neutrino data.
- ▶ Ratios of θ_α measurable at LHeC with high accuracy.
- ▶ Test minimal type I seesaw hypothesis.
- ▶ Together with ΔM also tests the compatibility with leptogenesis.

Conclusions

- ▶ Sterile neutrinos are well motivated extensions of the SM.
 - ▶ Present constraints allow for abundant production at LHeC.
 - ▶ Most sensitive searches for sterile neutrinos with masses
 - below m_W via displaced vertices
 - above $\mathcal{O}(100)$ GeV via lepton-flavour violating signatures.
 - ▶ Studying heavy neutrino Oscillations to gain insight into leptogenesis and the neutrino mass mechanism.
- ⇒ LHeC physics case + Baryon Asymmetry of the Universe.

Personal opinion:

Not upgrading LHC with an electron beam = wasted opportunity!

Antusch, OF; [1709.00880]