

Neutrino physics: present and future



Werner Rodejohann (MPIK)
FPCP '18
July 17



16th Conference on Flavor Physics & CP Violation

The aim of this conference is to review developments in flavor physics and CP violation, in both theory and experiment, exploiting the potential to study new physics at the LHC and future facilities. The topics include CP violation, rare decays, CKM elements with heavy quark decays, flavor phenomena in charged leptons and neutrinos, and also interplay between flavor and LHC high p_T physics

FPCP 2018

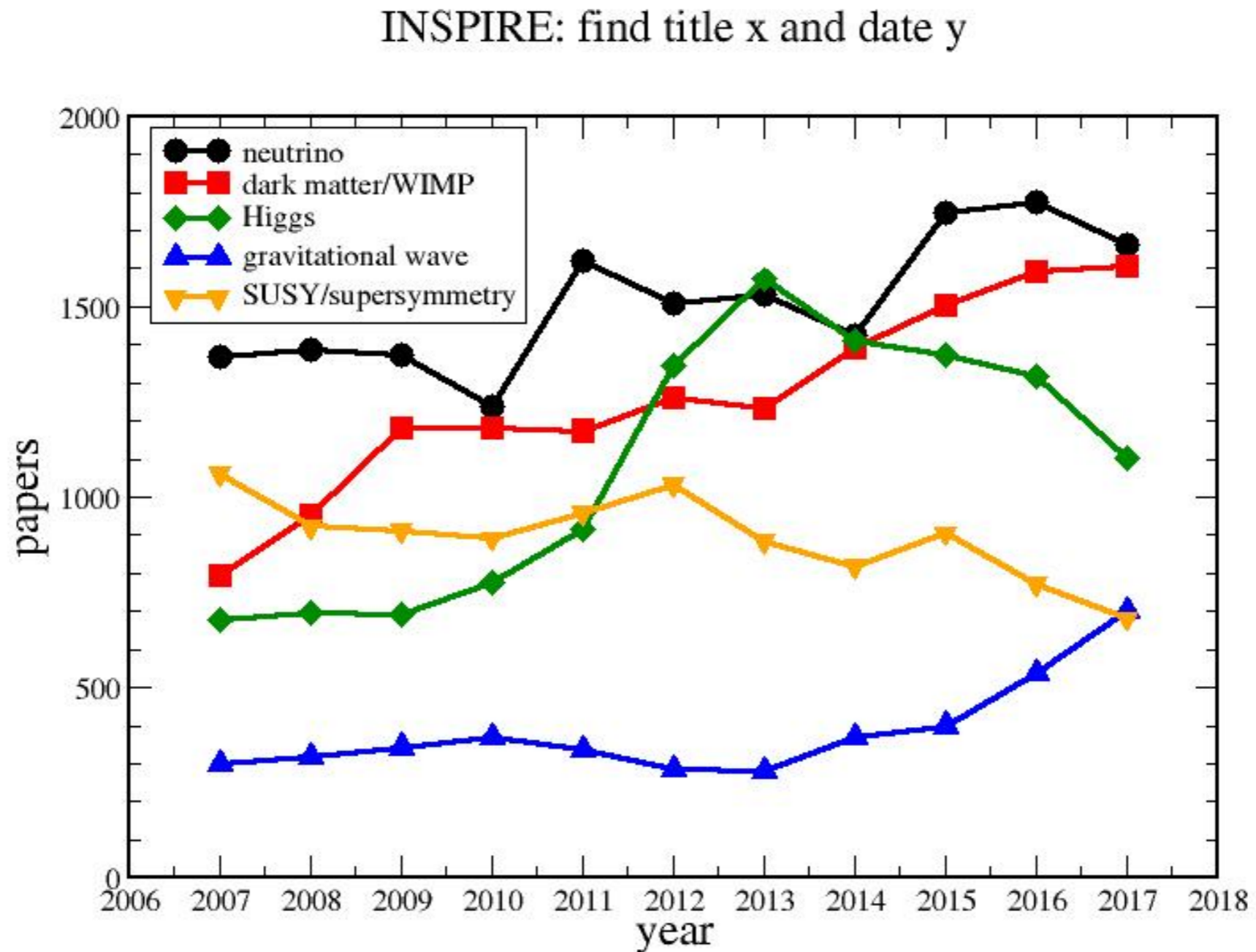
July 14-18, 2018, HYDERABAD, INDIA

University of Hyderabad IIT Hyderabad

Outline

- ❖ Neutrino mixing:
 - what have we learned?
 - what remains to be done?
- ❖ Neutrino mass:
 - what have we learned?
 - what remains to be done?
- ❖ New windows:
 - Coherent elastic neutrino-nucleus scattering

Neutrinos still a hot topic



Neutrinos oscillate and leptons mix

- ❖ we know that: $0 \neq \Delta m^2_{21} \neq \Delta m^2_{31}$
 - \Rightarrow all three masses different, at least two are non-zero
 - **hierarchy mild and neutrino mass much much smaller than all other masses**
- ❖ we know that: $U_{\text{PMNS}} = U_l^\dagger U_\nu \neq \mathbb{1}$
 - \Rightarrow charged lepton and neutrino mass matrices diagonalized with different matrices; Nature distinguishes ν_e, ν_μ, ν_τ
 - **mixing completely different from quark mixing**

Low Energy Paradigm

At low energies, neutrino mass matrix m_ν :

$$\mathcal{L} = \frac{1}{2} \nu^T m_\nu \nu \quad \text{with} \quad m_\nu = U \text{diag}(m_1, m_2, m_3) U^T$$

with PMNS matrix

$$U = \begin{pmatrix} c_{12} c_{13} & s_{12} c_{13} & s_{13} e^{-i\delta} \\ -s_{12} c_{23} - c_{12} s_{23} s_{13} e^{i\delta} & c_{12} c_{23} - s_{12} s_{23} s_{13} e^{i\delta} & s_{23} c_{13} \\ s_{12} s_{23} - c_{12} c_{23} s_{13} e^{i\delta} & -c_{12} s_{23} - s_{12} c_{23} s_{13} e^{i\delta} & c_{23} c_{13} \end{pmatrix} P$$

changes number of parameters in SM':

Species	#	Σ
Quarks	10	10
Leptons	3	13
Charge	3	16
Higgs	2	18
strong CP	1	19



Species	#	Σ
Quarks	10	10
Leptons	3 12	13 22
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3 Majorana neutrino paradigm \Rightarrow needs to be tested!

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with PMNS

$$U = \begin{pmatrix} c_{12} c_{13} & & s_{13} e^{-i\delta} \\ -s_{12} c_{23} - c_{12} s_{23} s_{13} & c_{13} & s_{23} c_{13} \\ s_{12} s_{23} - c_{12} c_{23} s_{13} e^{i\delta} & -s_{23} c_{13} & c_{23} c_{13} \end{pmatrix} P$$

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→

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Plus: mechanisms to generate m_ν have new particles, new energy scales, new concepts, new...

neutrino paradigm ⇒ needs to be tested!

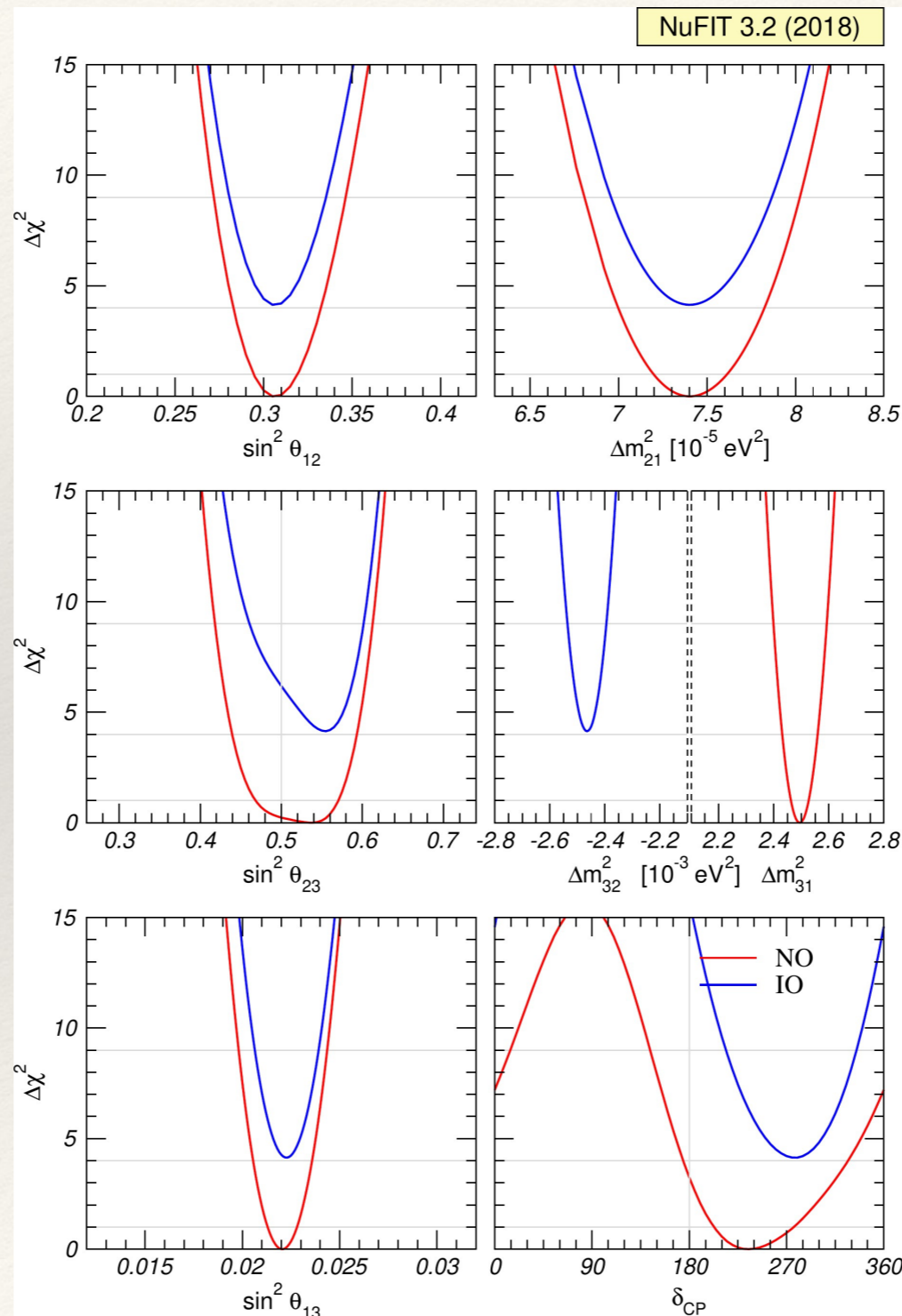
Low Energy Paradigm

❖ 3 Tasks:

- determine new parameters
- interpret / explain values of new parameters
- check for inconsistencies in standard picture

Determine Parameters

- ❖ We know:
 - θ_{12} and Δm_{21}^2
 - θ_{23} and $|\Delta m_{31}^2|$
 - θ_{13}
- ❖ We have limits:
 - m_1, m_2, m_3
- ❖ We don't know:
 - $\text{sgn}(\Delta m_{31}^2)$
 - δ, α, β

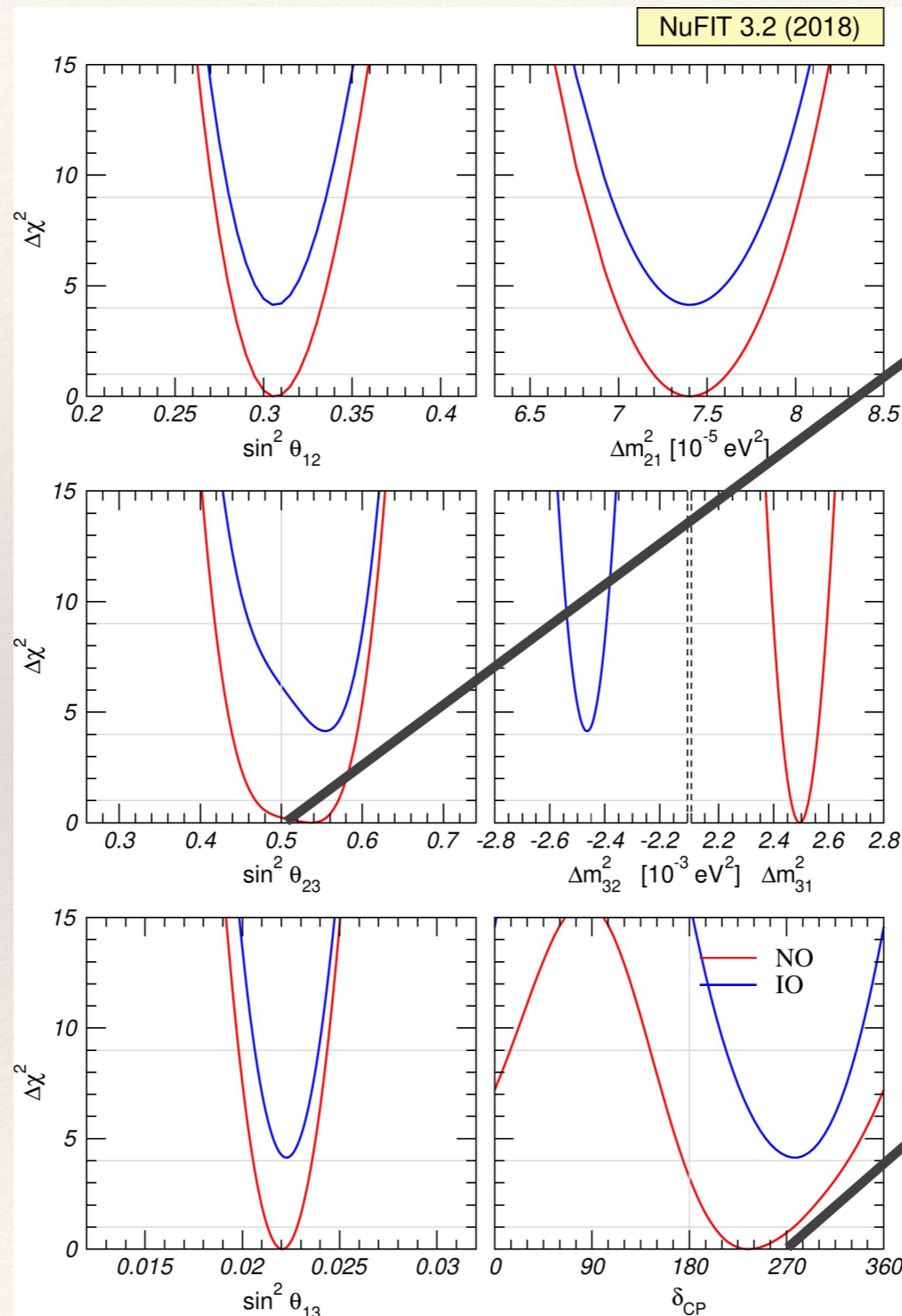


Robust fit results by
Valencia (1708.01186),
Bari (1804.09678),
NuFIT

See exptl. talks by
Sekiguchi (T2K),
Bhatnagar (NOvA),
Wu (reactors)

Determine Parameters

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maximal θ_{23} ?

$\delta = 3\pi/2$?

Normal Ordering preferred at $\approx 3\sigma$

Determine Parameters

❖ We know:

- θ_{12} and Δm_{21}^2
- θ_{23} and $|\Delta m_{31}^2|$
- θ_{13}

❖ We have limits:

- m_1, m_2, m_3

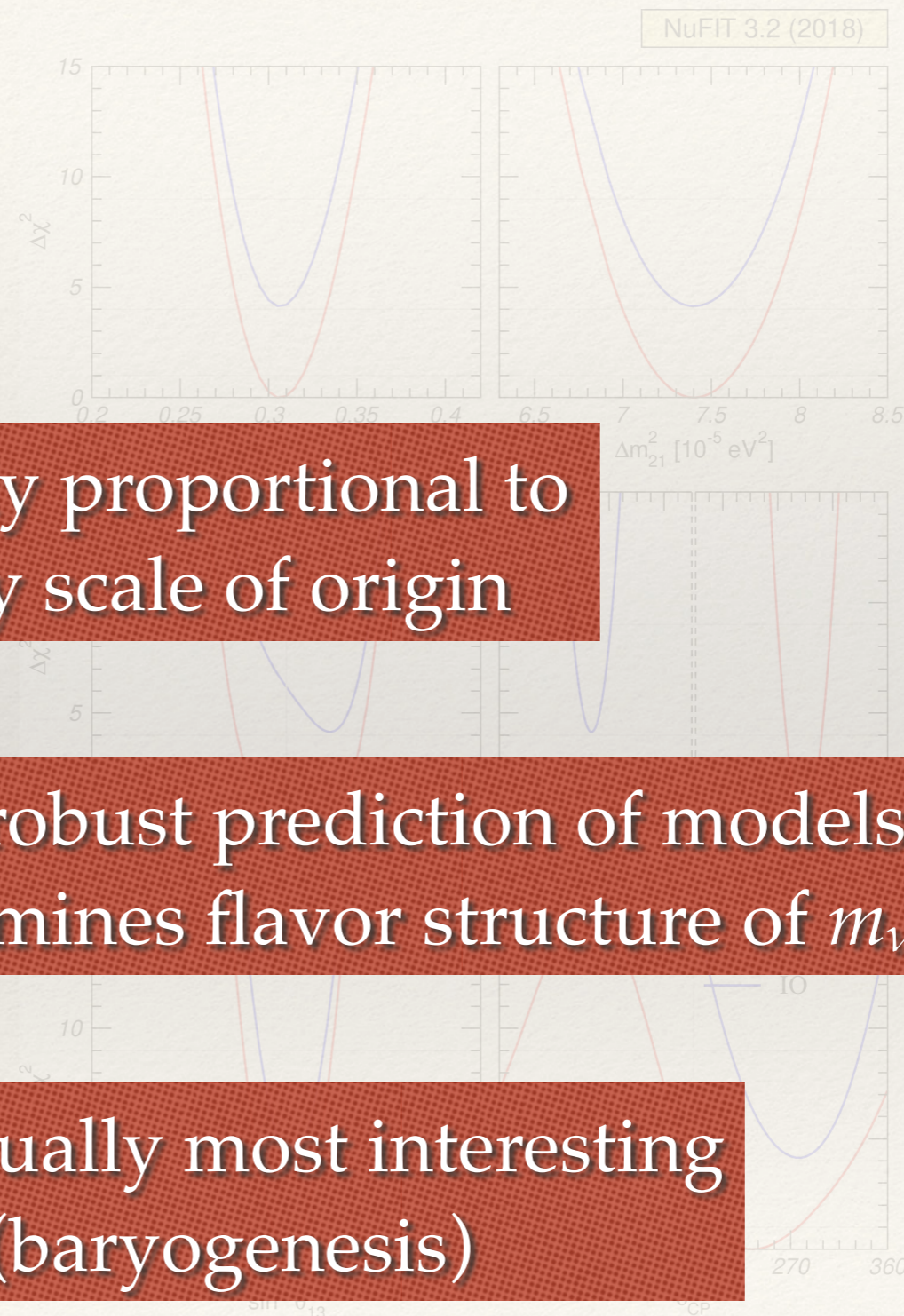
❖ We don't know:

- $\text{sgn}(\Delta m_{31}^2)$
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inversely proportional to
energy scale of origin

most robust prediction of models;
determines flavor structure of m_ν

conceptually most interesting
(baryogenesis)



Determine Parameters

❖ We know:

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❖ We have limits:

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❖ We don't know:

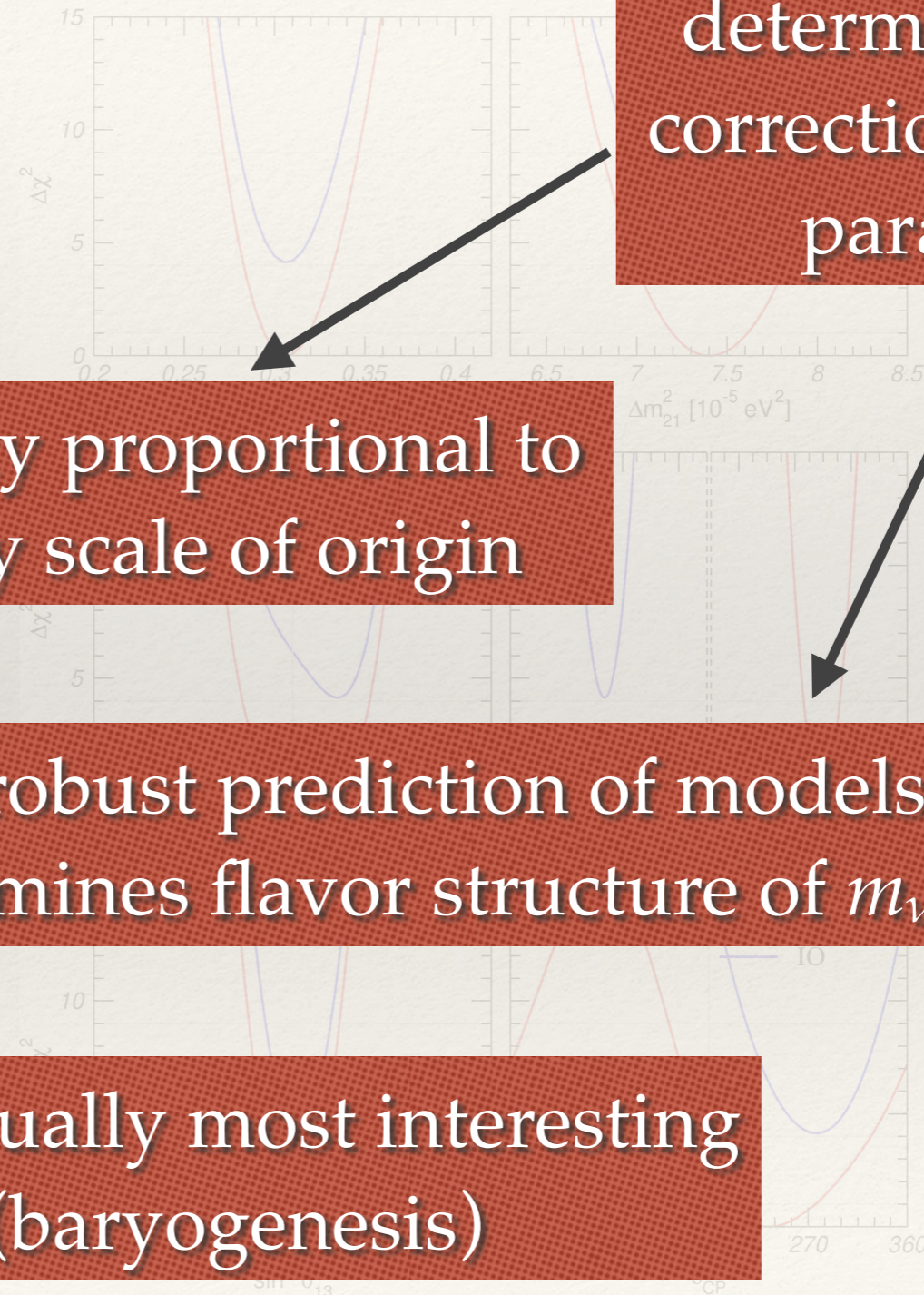
- $\text{sgn}(\Delta m^2_{31})$
- δ, α, β

determines size of correction to mixing parameters

inversely proportional to energy scale of origin

most robust prediction of models; determines flavor structure of m_ν

conceptually most interesting (baryogenesis)



Oscillation Parameters

parameter	best fit $\pm 1\sigma$	
Δm_{21}^2 [10^{-5}eV^2]	$7.55^{+0.20}_{-0.16}$	2.4%
$ \Delta m_{31}^2 $ [10^{-3}eV^2] (NO)	2.50 ± 0.03	1.3%
$ \Delta m_{31}^2 $ [10^{-3}eV^2] (IO)	$2.42^{+0.03}_{-0.04}$	1.3%
$\sin^2 \theta_{12}/10^{-1}$	$3.20^{+0.20}_{-0.16}$	5.5%
$\sin^2 \theta_{23}/10^{-1}$ (NO)	$5.47^{+0.20}_{-0.30}$	4.7%
$\sin^2 \theta_{23}/10^{-1}$ (IO)	$5.51^{+0.18}_{-0.30}$	4.4%
$\sin^2 \theta_{13}/10^{-2}$ (NO)	$2.160^{+0.083}_{-0.069}$	3.5%
$\sin^2 \theta_{13}/10^{-2}$ (IO)	$2.220^{+0.074}_{-0.076}$	3.5%
δ/π (NO)	$1.32^{+0.21}_{-0.15}$	10%
δ/π (IO)	$1.56^{+0.13}_{-0.15}$	9%

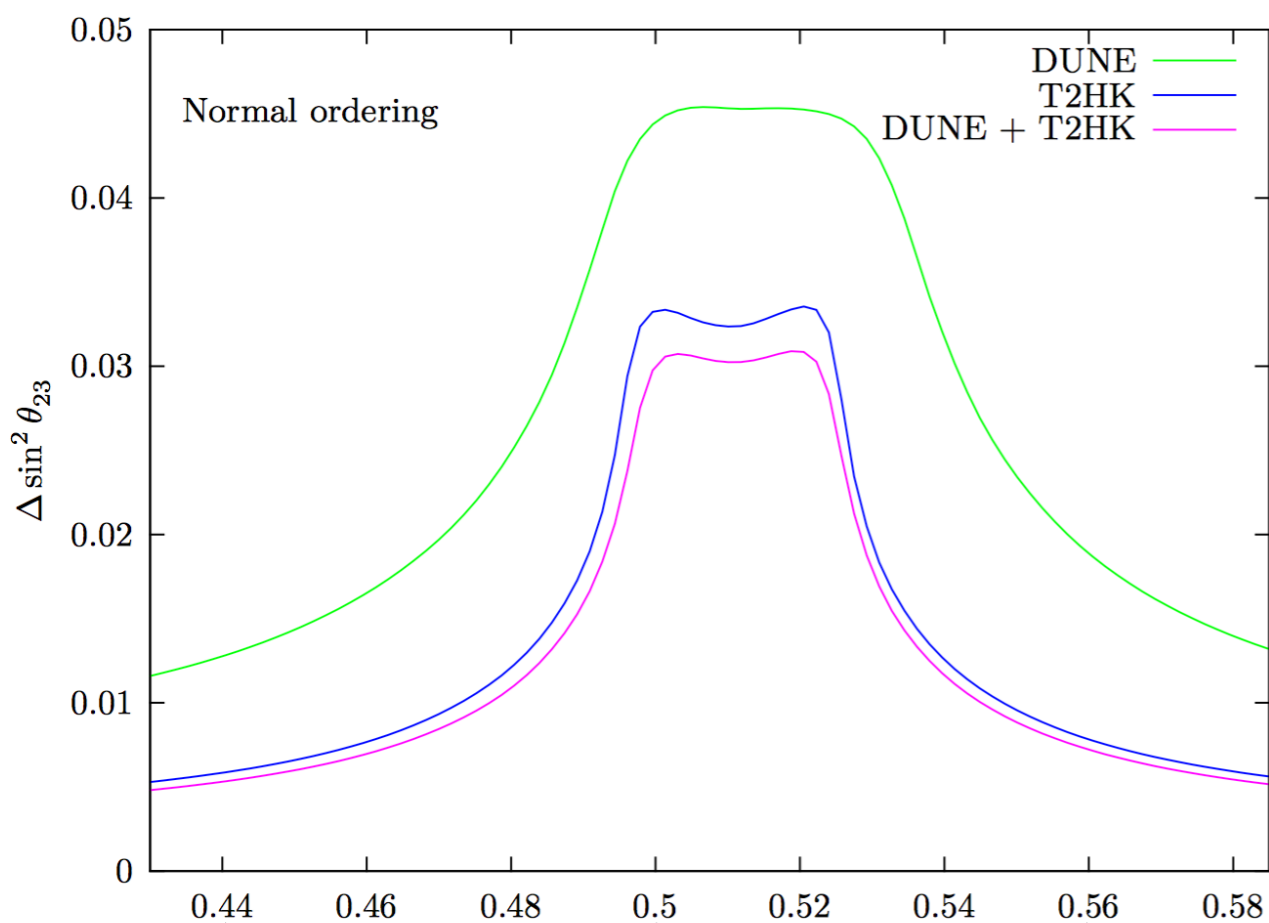
relative 1 σ uncertainty

	Current	JUNO
Δm_{12}^2	$\sim 3\%$	$\sim 0.6\%$
Δm_{23}^2	$\sim 5\%$	$\sim 0.6\%$
$\sin^2 \theta_{12}$	$\sim 6\%$	$\sim 0.7\%$
$\sin^2 \theta_{23}$	$\sim 20\%$	N/A
$\sin^2 \theta_{13}$	$\sim 14\% \rightarrow \sim 4\%$	$\sim 15\%$

Tortola, talk at Neutrino2018

Achievable Precision

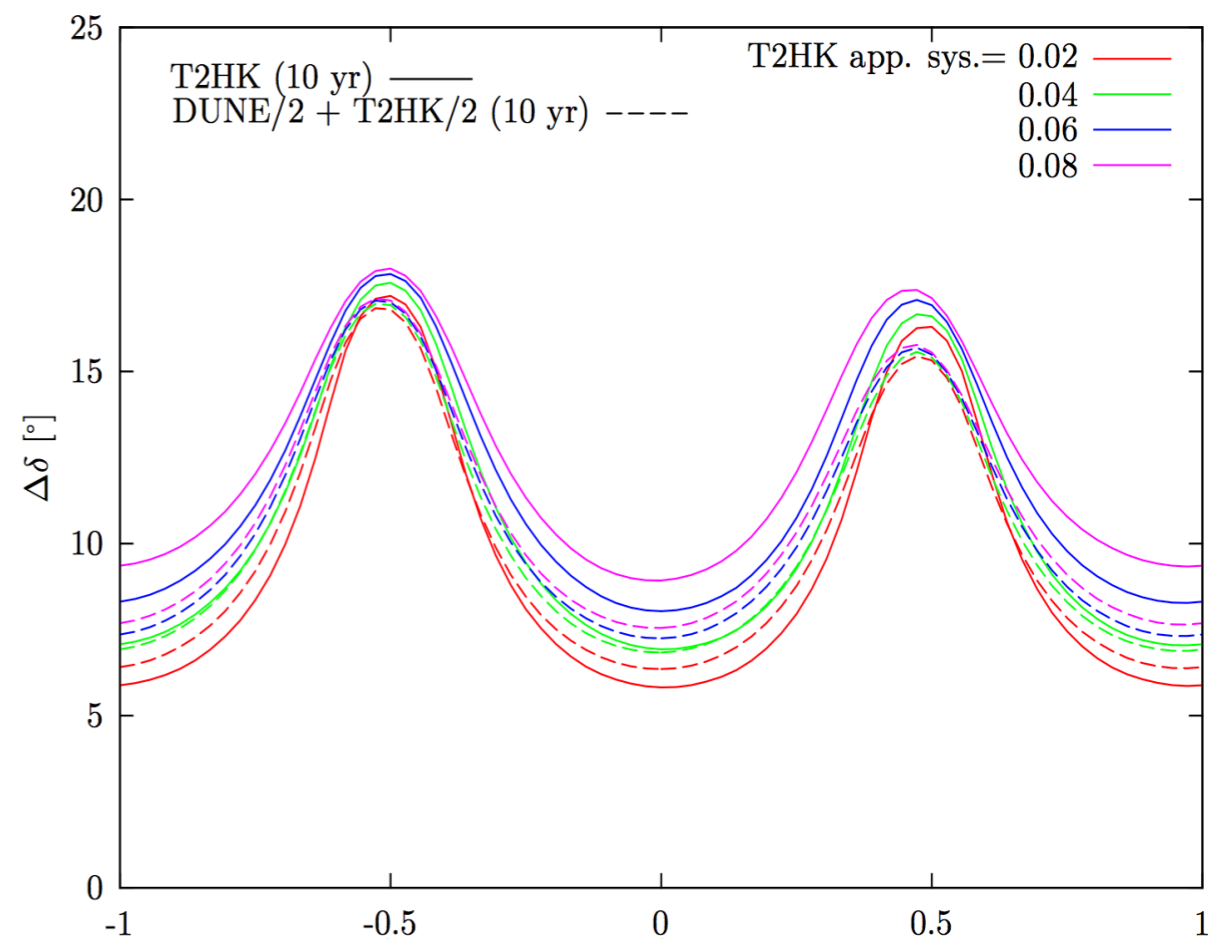
Ballet et al., 1612.07275



$(41.6 \pm 0.3)^{\circ}$

$(45 \pm 1.7)^{\circ}$

$(48.5 \pm 0.6)^{\circ}$



$\cos \delta:$

(0 ± 0.29)

(1 ± 0.006)

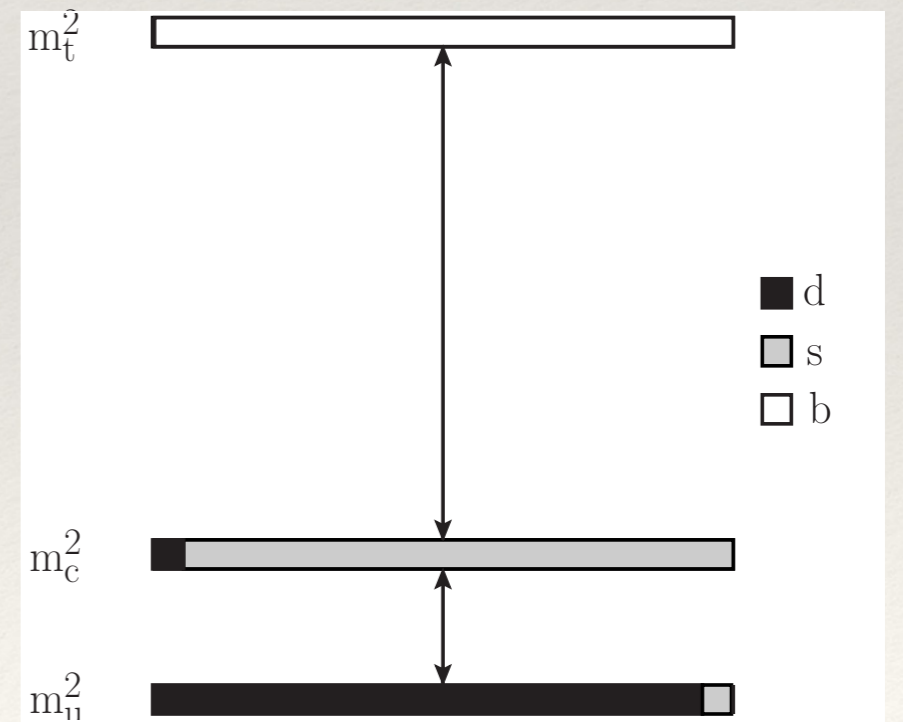
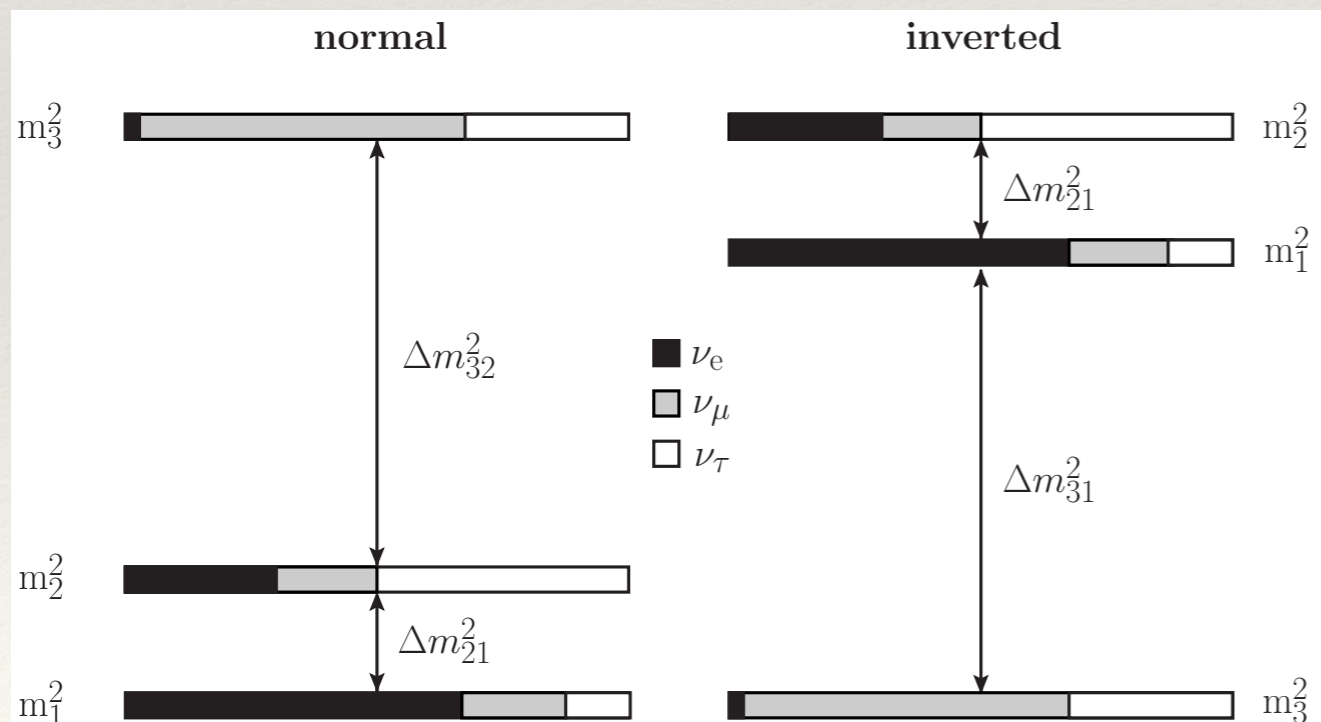
see talks by Tanaka, Indumathi

Implications of Lepton Mixing

NuFIT 3.2 (2018)

$$|U|_{3\sigma} = \begin{pmatrix} 0.799 \rightarrow 0.844 & 0.516 \rightarrow 0.582 & 0.141 \rightarrow 0.156 \\ 0.242 \rightarrow 0.494 & 0.467 \rightarrow 0.678 & 0.639 \rightarrow 0.774 \\ 0.284 \rightarrow 0.521 & 0.490 \rightarrow 0.695 & 0.615 \rightarrow 0.754 \end{pmatrix}$$

$$V_{\text{CKM}} = \begin{pmatrix} 0.97434^{+0.00011}_{-0.00012} & 0.22506 \pm 0.00050 & 0.00357 \pm 0.00015 \\ 0.22492 \pm 0.00050 & 0.97351 \pm 0.00013 & 0.0411 \pm 0.0013 \\ 0.00875^{+0.00032}_{-0.00033} & 0.0403 \pm 0.0013 & 0.99915 \pm 0.00005 \end{pmatrix}$$

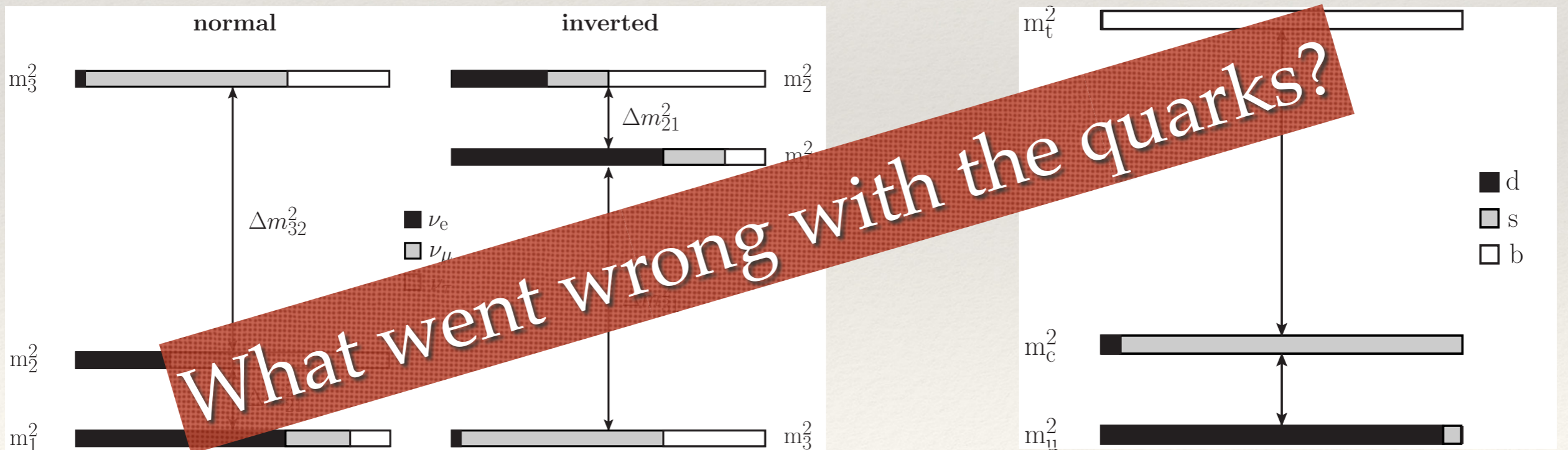


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Flavor Symmetries

- ❖ Nature seems to prefer large lepton mixing:

$$U_{\text{TBM}} = \begin{pmatrix} \sqrt{\frac{2}{3}} & \sqrt{\frac{1}{3}} & 0 \\ -\sqrt{\frac{1}{6}} & \sqrt{\frac{1}{3}} & -\sqrt{\frac{1}{2}} \\ -\sqrt{\frac{1}{6}} & \sqrt{\frac{1}{3}} & \sqrt{\frac{1}{2}} \end{pmatrix}$$

generated by rather special mass matrix

$$(m_\nu)_{\text{TBM}} = \begin{pmatrix} A & B & B \\ \cdot & \frac{1}{2}(A+B+D) & \frac{1}{2}(A+B-D) \\ \cdot & \cdot & \frac{1}{2}(A+B+D) \end{pmatrix}$$

mixing angles
independent from
masses!!

- ❖ completely different from quark sector (GST-relation):

$$M = \begin{pmatrix} 0 & a \\ a & b \end{pmatrix} \Rightarrow \tan \theta_C \simeq \sqrt{\frac{m_d}{m_s}}$$

Flavor Symmetries

❖ preferred solution: Discrete Non-Abelian Symmetries

Altarelli, Feruglio, 1002.0211

Group	d	Irr. Repr.'s	Presentation
$D_3 \sim S_3$	6	1, 1', 2	$A^3 = B^2 = (AB)^2 = 1$
D_4	8	1 ₁ , ..., 1 ₄ , 2	$A^4 = B^2 = (AB)^2 = 1$
D_7	14	1, 1', 2, 2', 2''	$A^7 = B^2 = (AB)^2 = 1$
A_4	12	1, 1', 1'', 3	$A^3 = B^2 = (AB)^3 = 1$
$A_5 \sim PSL_2(5)$	60	1, 3, 3', 4, 5	$A^3 = B^2 = (BA)^5 = 1$
T'	24	1, 1', 1'', 2, 2', 2'', 3	$A^3 = (AB)^3 = R^2 = 1, B^2 = R$
S_4	24	1, 1', 2, 3, 3'	$BM : A^4 = B^2 = (AB)^3 = 1$ $TB : A^3 = B^4 = (BA^2)^2 = 1$
$\Delta(27) \sim Z_3 \rtimes Z_3$	27	1 ₁ , ..., 1 ₉ , 3, $\bar{3}$	
$PSL_2(7)$	168	1, 3, $\bar{3}$, 6, 7, 8	$A^3 = B^2 = (BA)^7 = (B^{-1}A^{-1}BA)^4 = 1$
$T_7 \sim Z_7 \rtimes Z_3$	21	1, 1', $\bar{1}'$, 3, $\bar{3}$	$A^7 = B^3 = 1, AB = BA^4$

Type	L_i	ℓ_i^c	ν_i^c	Δ
A1	$\underline{3}$	1, 1', 1''
A2				1, 1', 1'', $\underline{3}$
B1	$\underline{3}$	1, 1', 1''	$\underline{3}$...
B2				1, $\underline{3}$
C1				...
C2	$\underline{3}$	$\underline{3}$...	1
C3				1, $\underline{3}$
C4				1, 1', 1'', $\underline{3}$
D1				...
D2	$\underline{3}$	$\underline{3}$	$\underline{3}$	1
D3				1'
D4				1', $\underline{3}$
E	$\underline{3}$	$\underline{3}$	1, 1', 1''	...
F	1, 1', 1''	$\underline{3}$	$\underline{3}$	1 or 1'
G	$\underline{3}$	1, 1', 1''	1, 1', 1''	...
H	$\underline{3}$	1, 1, 1
I	$\underline{3}$	1, 1, 1	1, 1, 1	...
J	$\underline{3}$	1, 1, 1	$\underline{3}$...

Many possible groups, within each group many models...

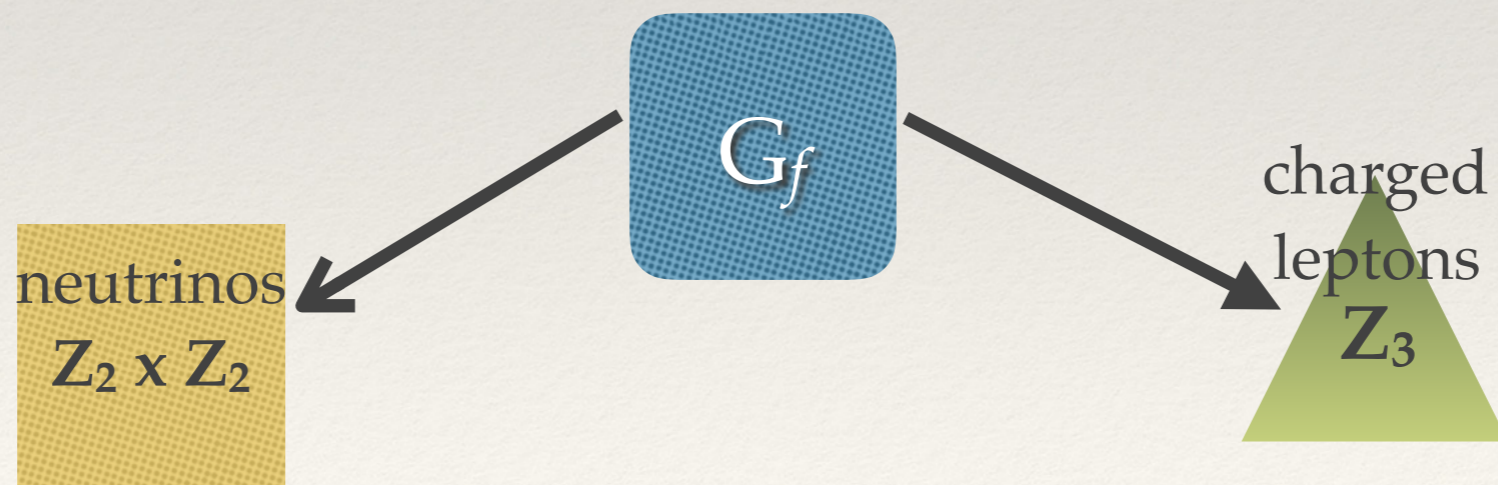
⇒ can distinguish only classes of models

Flavor Symmetries

Lesson 1: put different generations in same irrep of group:

$$\begin{pmatrix} L_e \\ L_\mu \\ L_\tau \end{pmatrix} = \begin{pmatrix} \begin{pmatrix} \nu_e \\ e^- \end{pmatrix}_L \\ \begin{pmatrix} \nu_\mu \\ \mu^- \end{pmatrix}_L \\ \begin{pmatrix} \nu_\tau \\ \tau^- \end{pmatrix}_L \end{pmatrix} \sim 3_f$$

Lesson 2: flavor group broken to different subgroups:



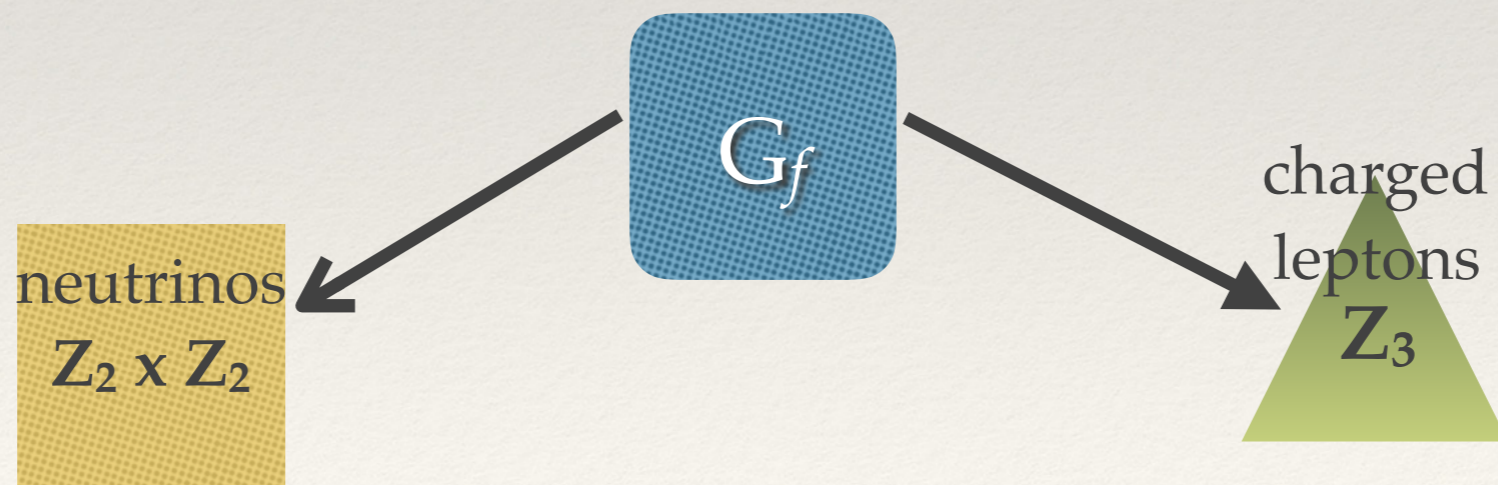
Flavor Symmetries

Lesson 1: put different generations in same irrep of group:

related to 3 generations?

$$\begin{pmatrix} L_e \\ L \end{pmatrix} \begin{pmatrix} \begin{pmatrix} \nu_e \\ e^- \end{pmatrix} \\ \begin{pmatrix} \nu_\tau \\ \tau^- \end{pmatrix} \end{pmatrix}_L \sim 3_f$$

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Flavor Symmetries

Lesson 1: put different generations in same irrep of group:

$$\begin{pmatrix} L_e \\ L \end{pmatrix} \begin{pmatrix} \begin{pmatrix} \nu_e \\ e^- \end{pmatrix} \\ \begin{pmatrix} \nu_\tau \\ \tau^- \end{pmatrix} \end{pmatrix} \sim 3_f$$

related to 3 generations?

Lesson 2: flavor group broken to different subgroups:



How to predict the CP phase

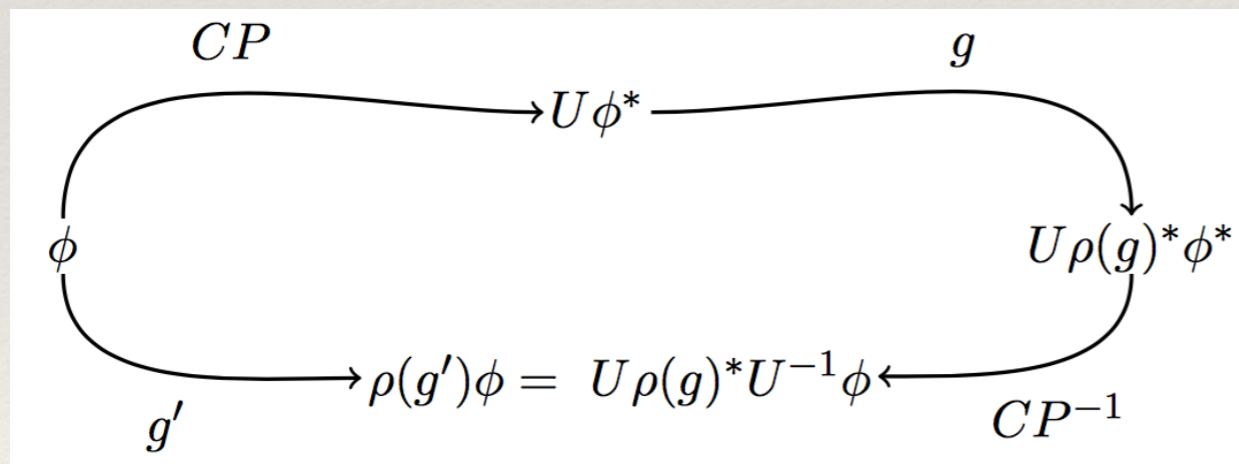
- ❖ μ - τ reflection symmetry: $\nu_e \leftrightarrow \nu_e^C$ and $\nu_\mu \leftrightarrow \nu_\tau^C$ (see talk by Deepthi)

$$m_\nu = \begin{pmatrix} x & z_1 & z_1^* \\ \cdot & z_2 & y \\ \cdot & \cdot & z_2^* \end{pmatrix}$$

gives $\delta = \pm\pi/2$ and $\theta_{23} = \pi/4$

Ma; Grimus, Lavoura; Joshipura, Patel; He, WR, Xu

- ❖ combine CP and flavor symmetry, typically gives $\delta = \pm\pi/2, \pm\pi, 0$



(implies consistency relation: generalized CP transformation can be interpreted as a representation of an automorphism of the discrete group)

Grimus; Chen; Feruglio, Hagedorn, Ziegler; Holthausen, Schmidt, Lindner; Ding, King, Stuart; Meroni, Petcov; Branco, King, Varzielas,...

Abelian Flavor Symmetries

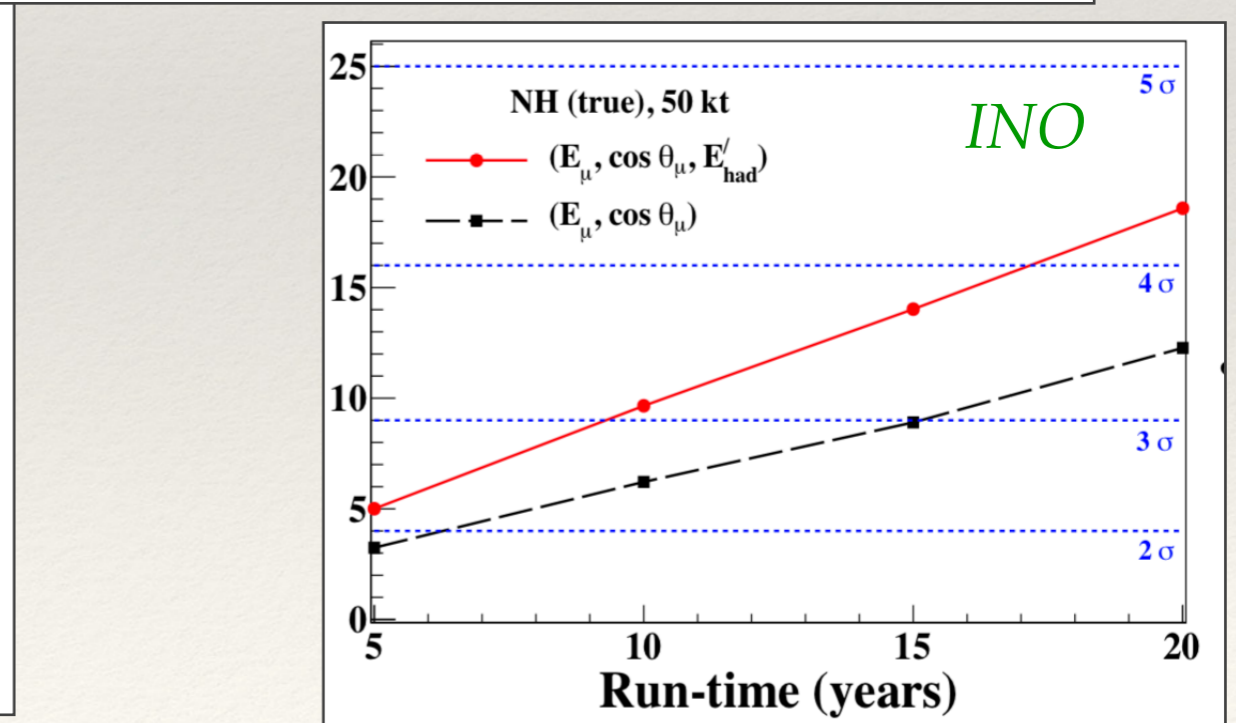
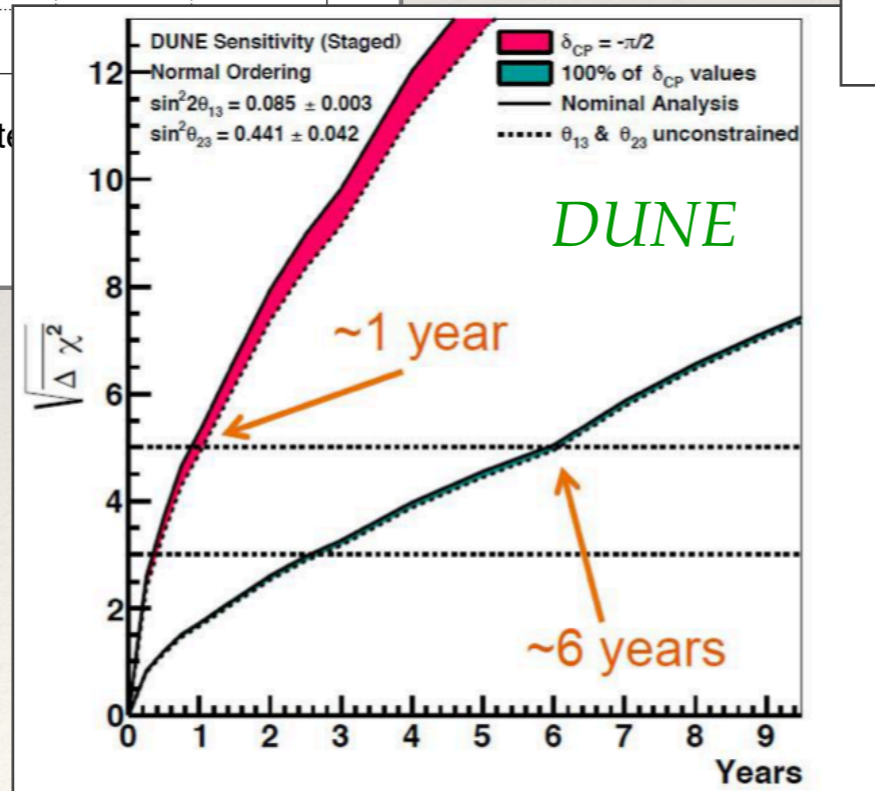
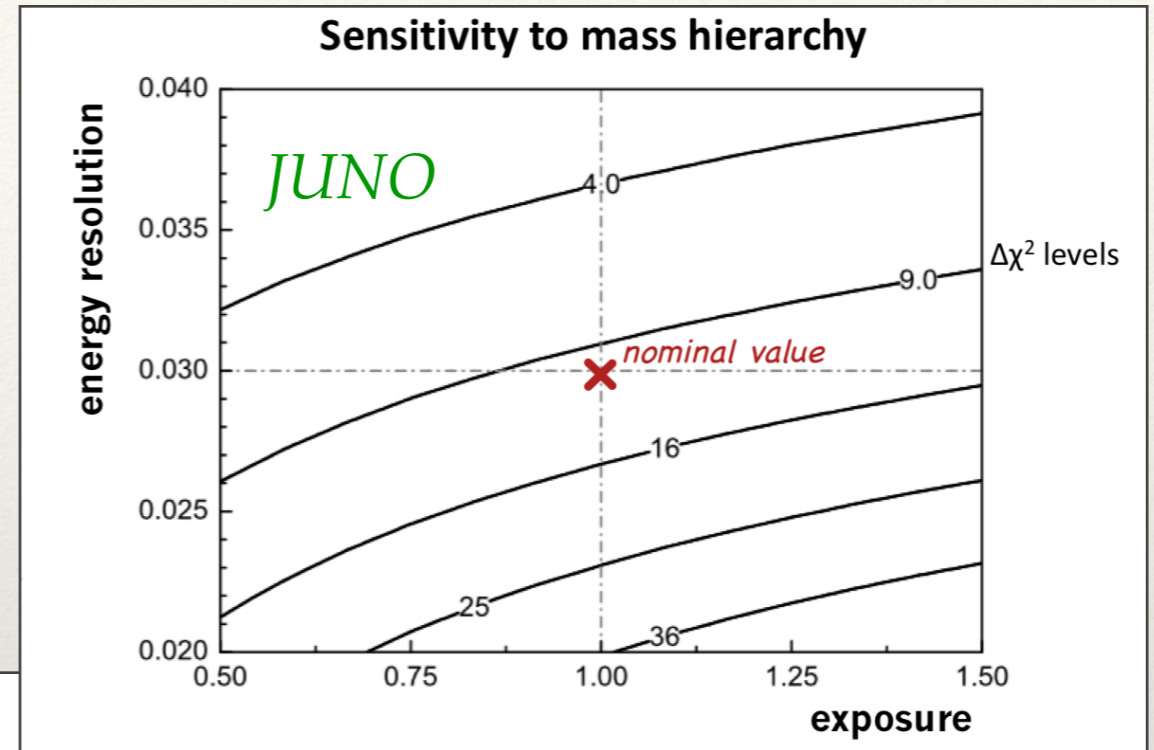
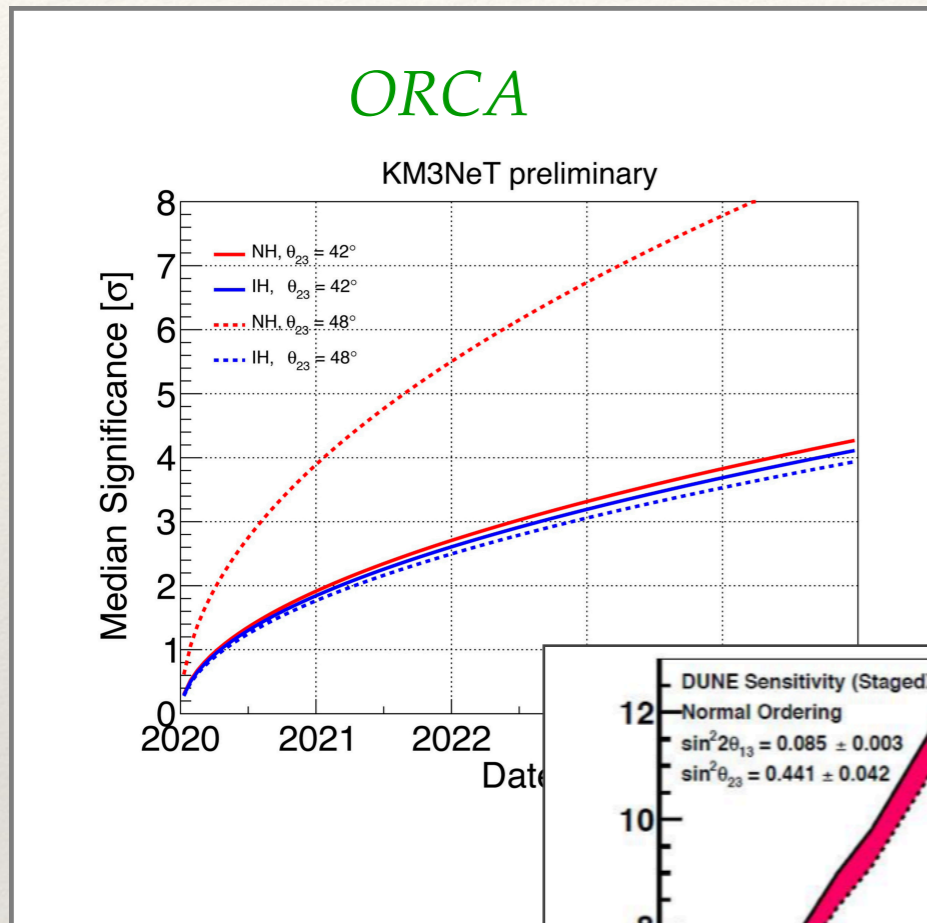
- ❖ Less predictive but less complicated: Abelian flavor symmetry, e.g. $L_\mu - L_\tau$
 - anomaly-free, has Z' and can explain $(g - 2)_\mu$
 - can be extended to quark sector to explain anomalies in $B \rightarrow K^* \mu\mu$ and $\text{BR}(B \rightarrow K\mu\mu) / \text{BR}(B \rightarrow Kee)$ [*Crivellin, Ambrosio, Heeck, 1501.00993*] (making predictions for $h \rightarrow \mu\tau$, LFV, etc.)
 - masses a and $\pm b$, $\theta_{23} = \pi/4$, $\theta_{13} = 0$

$$(m_\nu)^{L_\mu - L_\tau} = \begin{pmatrix} a & 0 & 0 \\ \cdot & 0 & b \\ \cdot & \cdot & 0 \end{pmatrix}$$

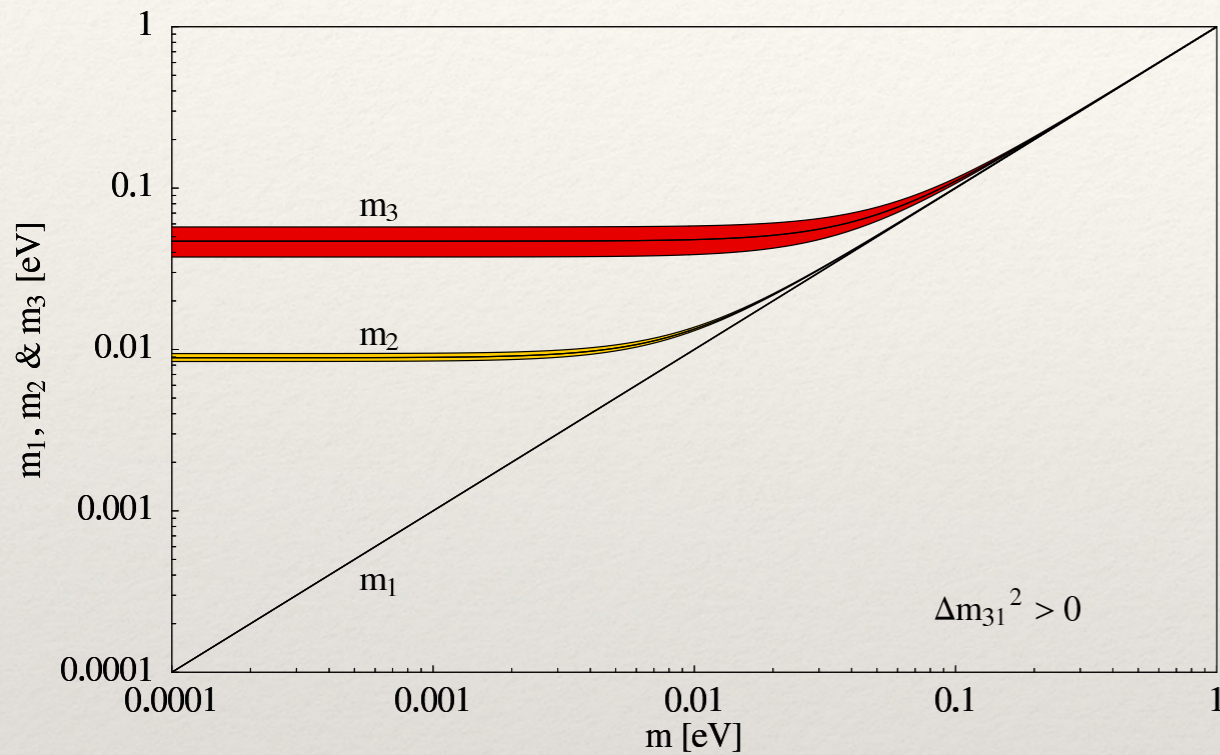
Heeck, WR, 1107.5238

Future of Mass Ordering

see talks by Tanaka, Wu, Indumathi



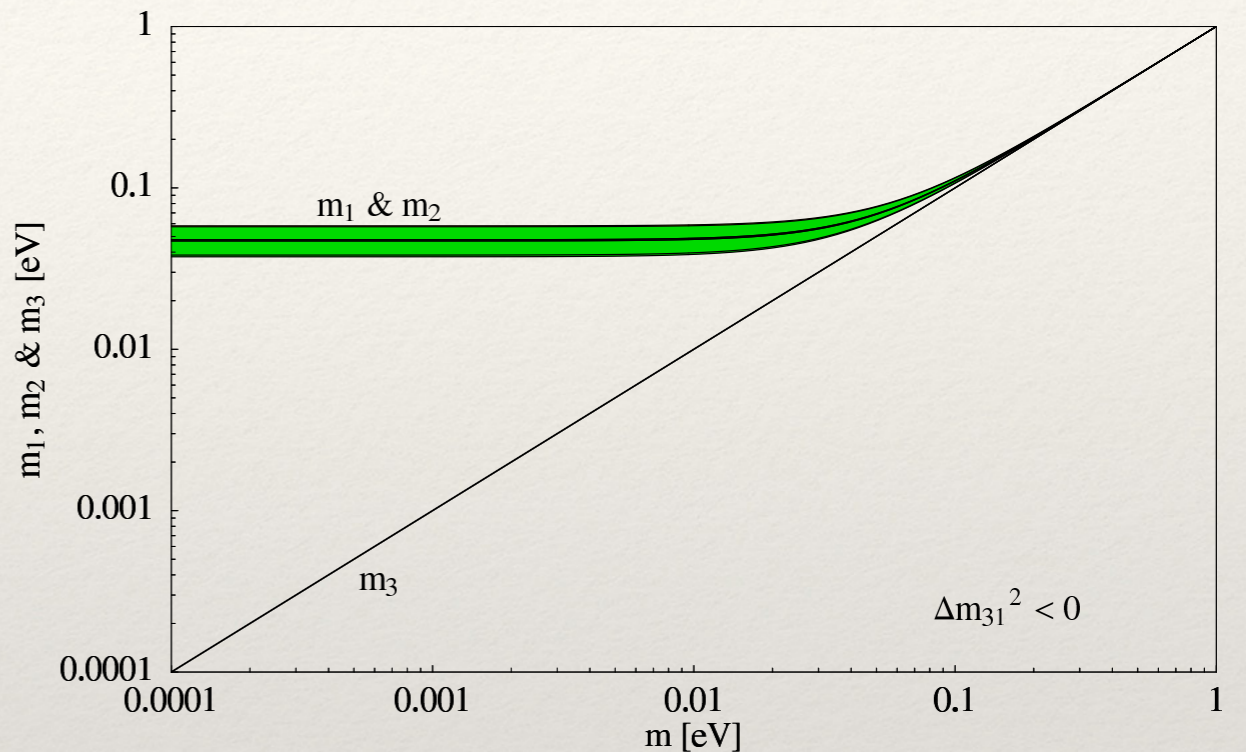
Masses and Ordering



mild hierarchy in normal ordering:

$$m_3 / m_2 \approx (\Delta m_{\text{atm}}^2 / \Delta m_{\text{sol}}^2)^{1/2} \approx 5$$

$$(m_\nu)_{\text{NH}} \sim \begin{pmatrix} \epsilon^2 & \epsilon & \epsilon \\ \epsilon & 1 & 1 \\ \epsilon & 1 & 1 \end{pmatrix}$$



strong tuning in inverted ordering:

$$m_2 / m_1 \approx 1 + \frac{1}{2} \Delta m_{\text{sol}}^2 / \Delta m_{\text{atm}}^2$$

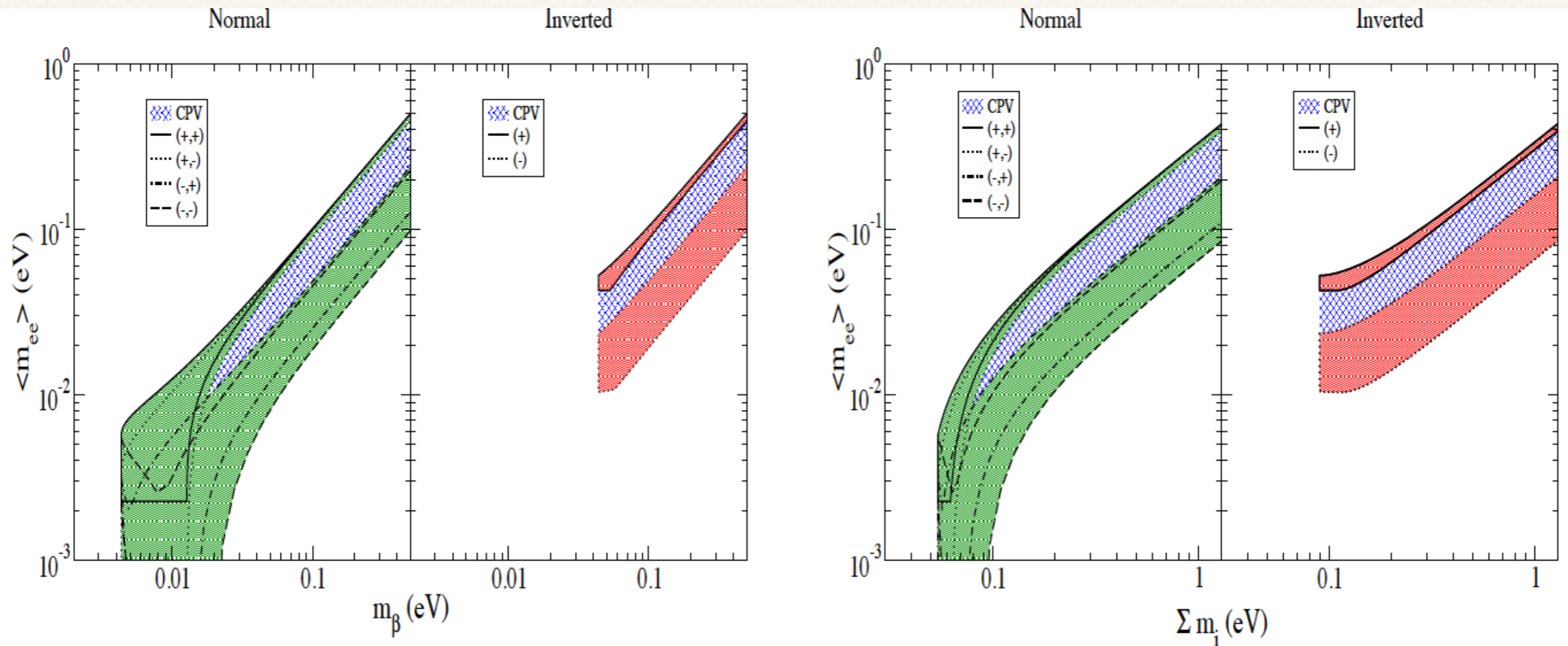
$$(m_\nu)_{\text{IH}} \sim \begin{pmatrix} 1 & 1 & 1 \\ 1 & 1 & 1 \\ 1 & 1 & 1 \end{pmatrix}$$

plus almost democratic structure of mass matrix

Neutrino Mass Observables

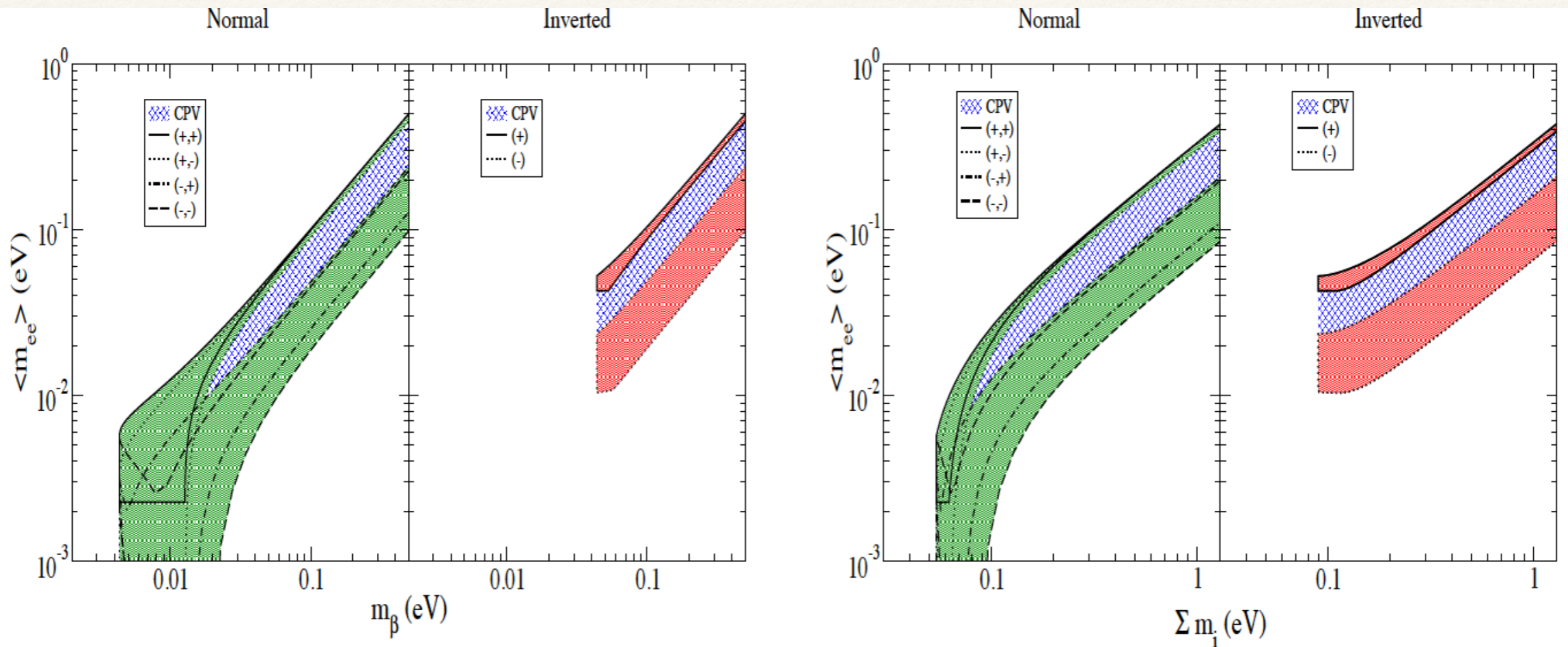
Method	Observable	current	near	far	pro	con
Kurie	$(\sum U_{ei} ^2 m_i^2)^{1/2}$	2.3 eV	0.3 eV	0.1 eV?	model-indep.; clean	final; weakest
cosmo	$\sum m_i$	0.25 eV	0.1 eV	0.05 eV?	best; NH/IH	model-dep.; systematics
$0\nu\beta\beta$	$\sum U_{ei}^2 m_i$	0.2 eV	0.05 eV	0.01 eV?	fundamental; NH/IH	model-dep.; NMEs

Neutrino Mass Observables



complete complementarity
of observables

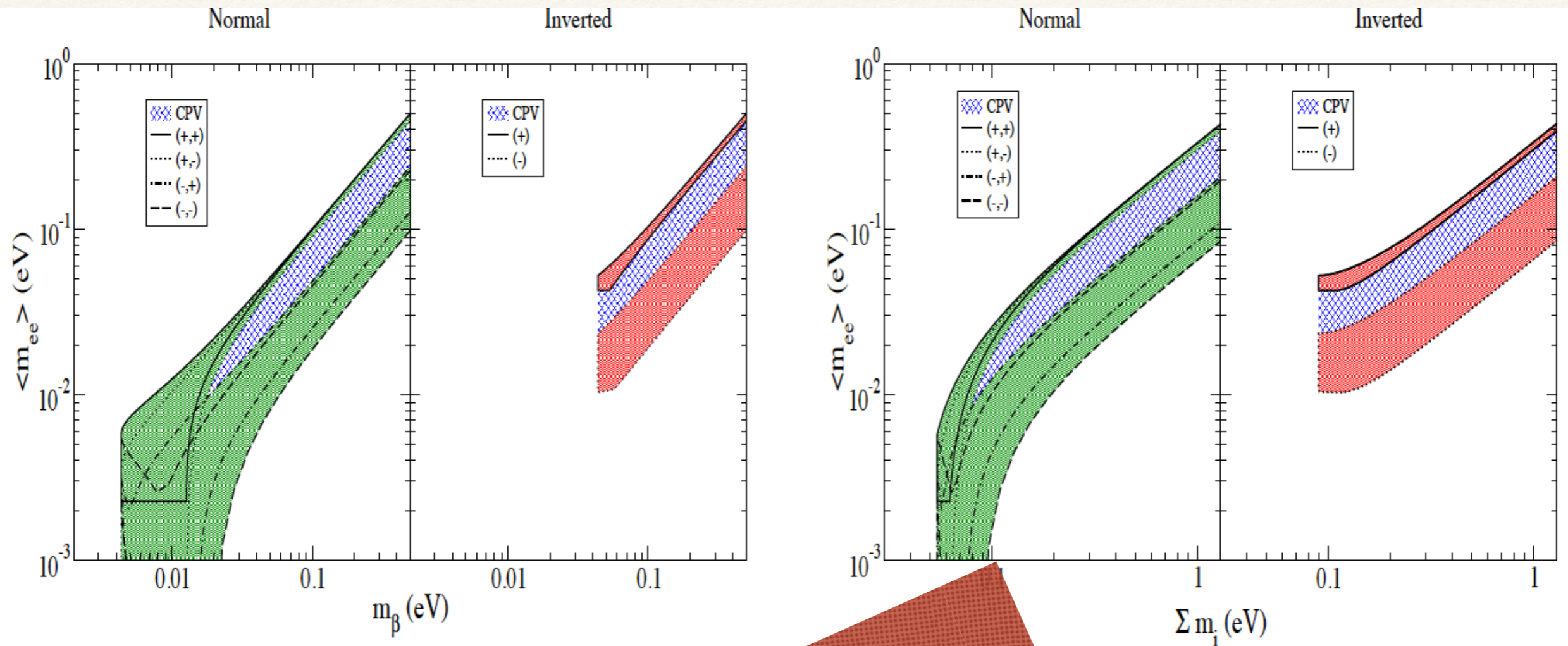
Neutrino Mass Observables



complete observability of CPV and parity variables

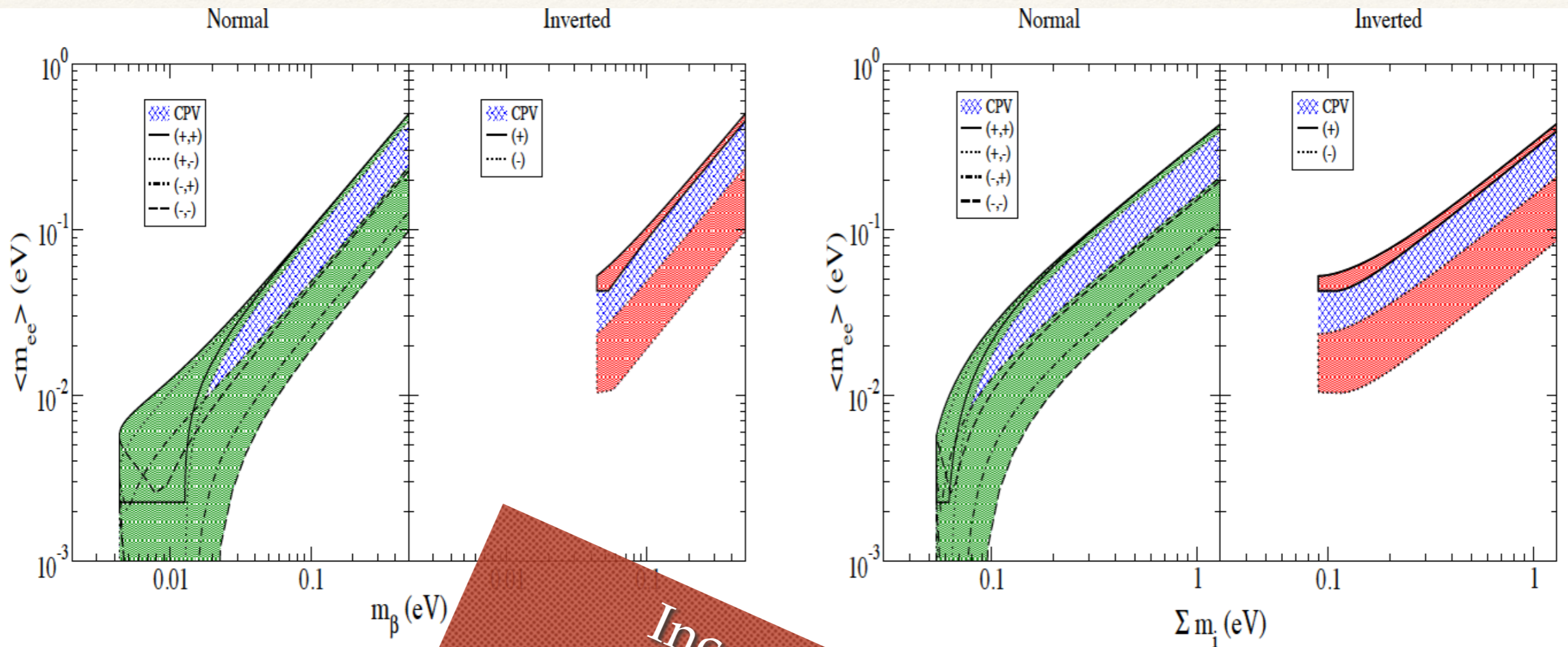
All need to be pursued!

Neutrino Mass Observables



Consistency
would be spectacular
confirmation!

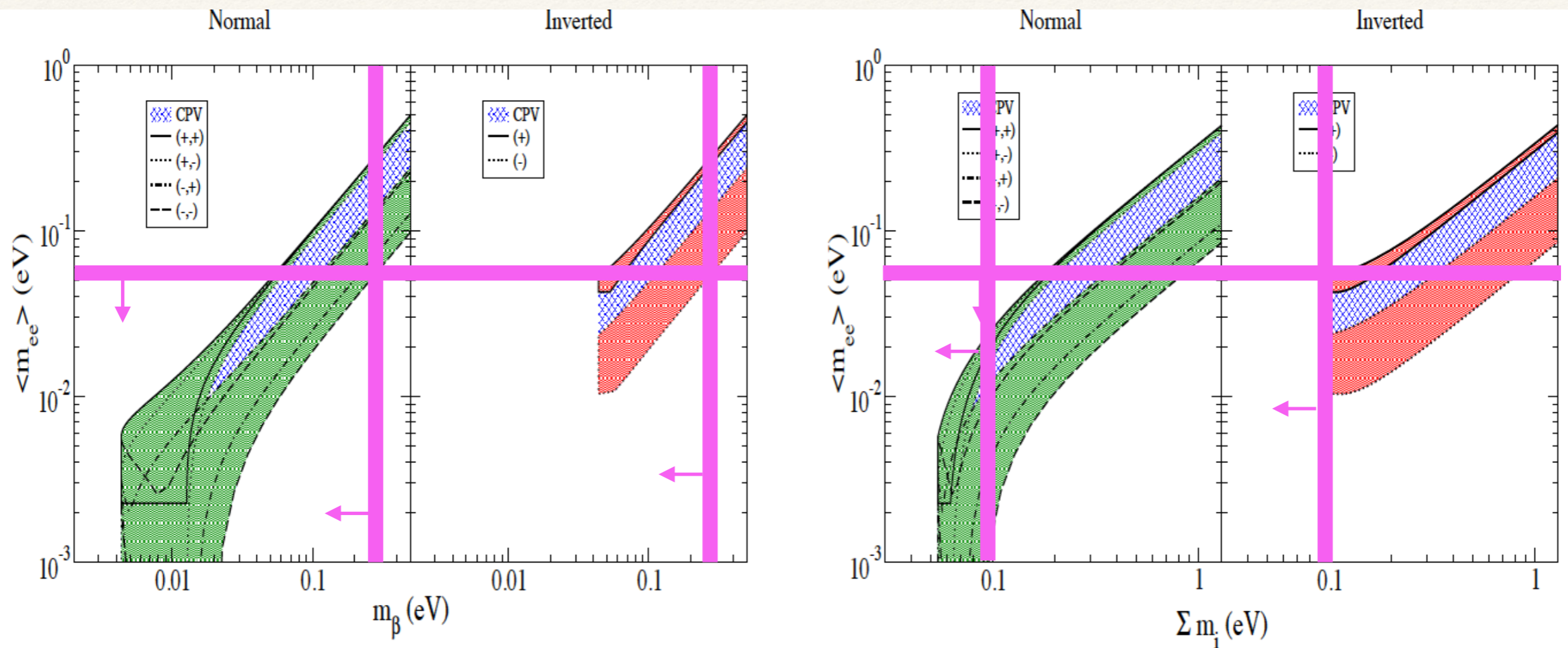
Neutrino Mass Observables



Inconsistencies
would be major
discovery!

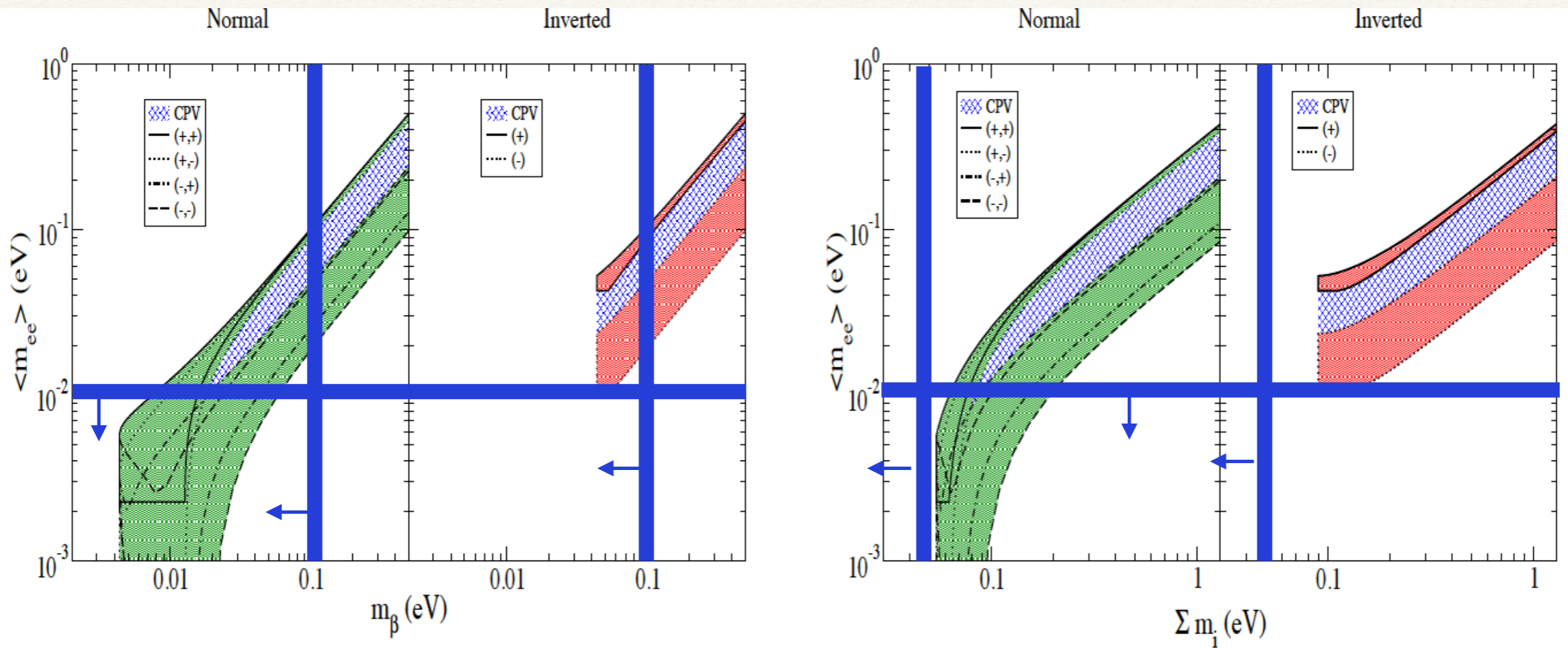
comp
of obs

Neutrino Mass Observables



near future

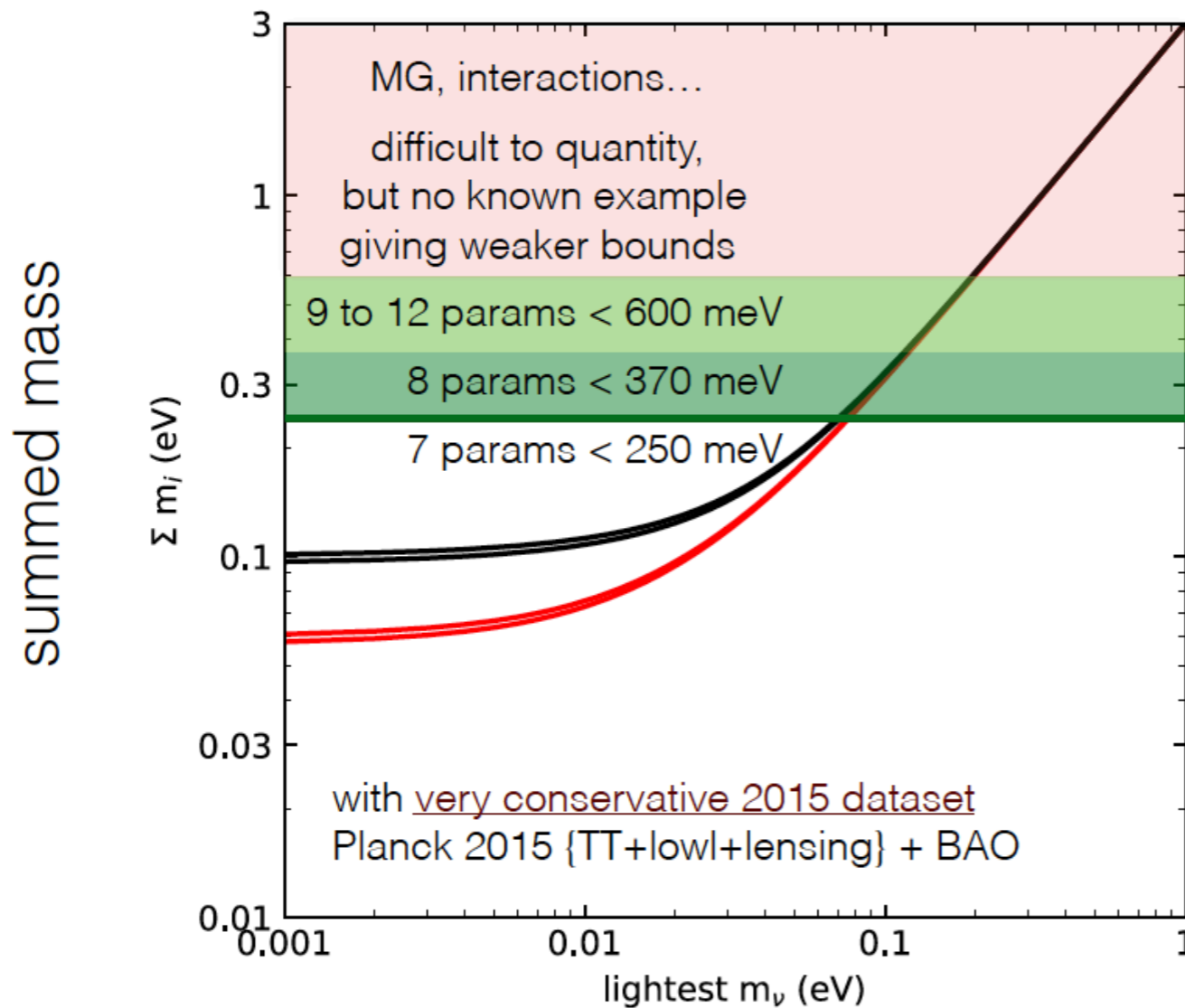
Neutrino Mass Observables



far future

Cosmological Mass Bounds

95%CL upper bounds on $\Sigma_i m_i$ beyond 7 parameters



[Planck col.] 1502.01589; Di Valentino et al. 1507.06646

Usual suspects:

- extra massless relics
- extra light relics
- spatial curvature
- simplest dynamical DE
- primordial GWs
- primordial tilt running

Even more freedom in:

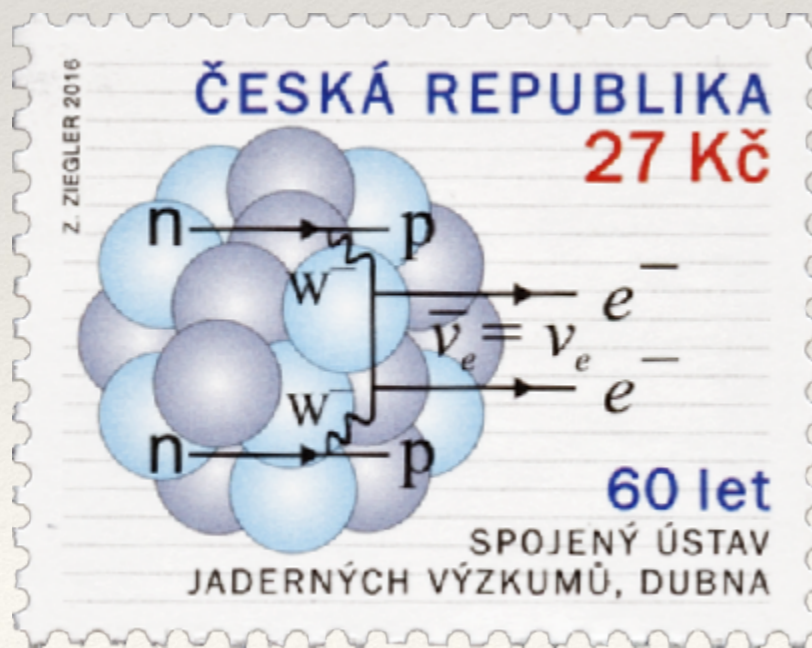
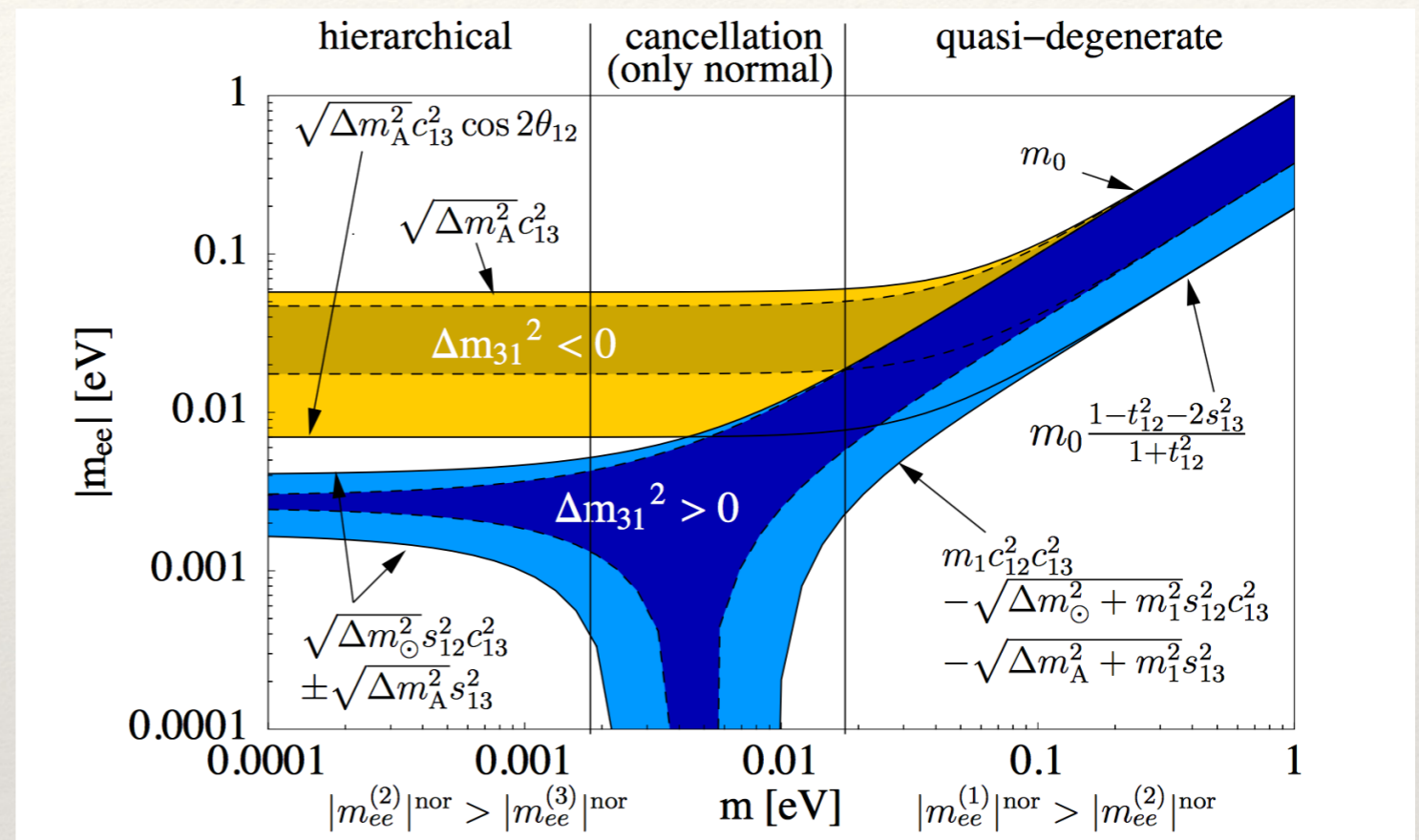
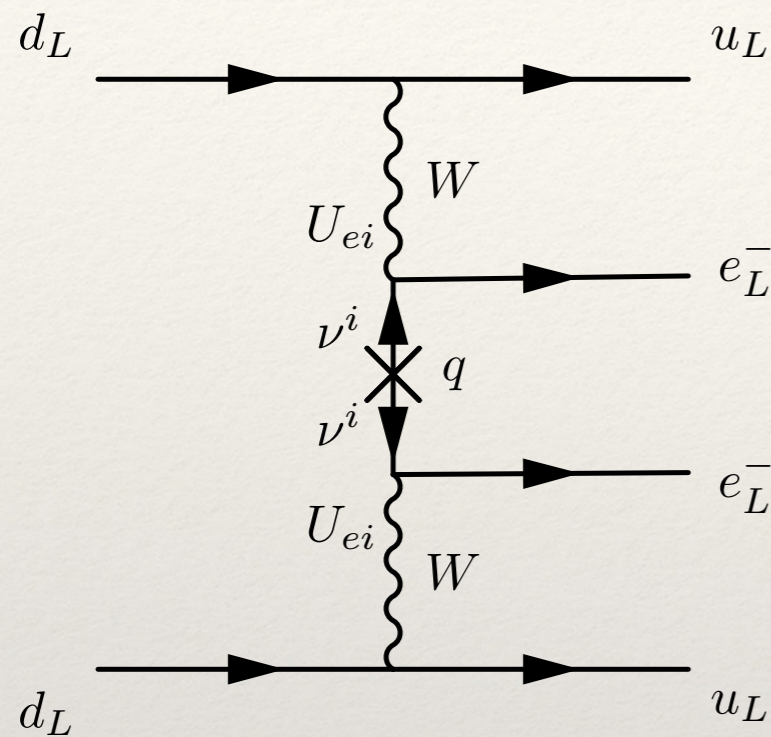
- modified Einstein Gravity
- interactions in DM sector
- ~~primordial perturbations~~

Lesgourgues, talk at Neutrino2018

Plus: future observation will have to see neutrino mass even in modest extensions!

E.g: 5 σ detection when Euclid and SKA are combined!

Double Beta Decay



$$|m_{ee}| = \left| \sum U_{ei}^2 m_i \right| = \left| U_{e1}^2 m_1 + U_{e2}^2 m_2 e^{i\alpha} + U_{e3}^2 m_3 e^{i\beta} \right|$$

$$= f(\theta_{12}, |U_{e3}|, m_i, \text{sgn}(\Delta m_A^2), \alpha, \beta)$$

known

limits

unknown

Experimental Situation

Experiment	Iso.	Iso. Mass [kg _{iso}]	σ [keV]	ROI [σ]	ϵ_{FV} [%]	ϵ_{sig} [%]	\mathcal{E} [$\frac{\text{kg}_{iso} \text{ yr}}{\text{yr}}$]	\mathcal{B} [$\frac{\text{cts}}{\text{kg}_{iso} \text{ ROI yr}}$]	3 σ disc. sens.		Required Improvement		
									$\hat{T}_{1/2}$ [yr]	$\hat{m}_{\beta\beta}$ [meV]	Bkg	σ	Iso. Mass
LEGEND 200 [61, 62]	⁷⁶ Ge	175	1.3	[-2, 2]	93	77	119	$1.7 \cdot 10^{-3}$	$8.4 \cdot 10^{26}$	40–73	3	1	5.7
LEGEND 1k [61, 62]	⁷⁶ Ge	873	1.3	[-2, 2]	93	77	593	$2.8 \cdot 10^{-4}$	$4.5 \cdot 10^{27}$	17–31	18	1	29
SuperNEMO [68, 69]	⁸² Se	100	51	[-4, 2]	100	16	16.5	$4.9 \cdot 10^{-2}$	$6.1 \cdot 10^{25}$	82–138	49	2	14
CUPID [58, 59, 70]	⁸² Se	336	2.1	[-2, 2]	100	69	221	$5.2 \cdot 10^{-4}$	$1.8 \cdot 10^{27}$	15–25	n/a	6	n/a
CUORE [52, 53]	¹³⁰ Te	206	2.1	[-1.4, 1.4]	100	81	141	$3.1 \cdot 10^{-1}$	$5.4 \cdot 10^{25}$	66–164	6	1	19
CUPID [58, 59, 70]	¹³⁰ Te	543	2.1	[-2, 2]	100	81	422	$3.0 \cdot 10^{-4}$	$2.1 \cdot 10^{27}$	11–26	3000	1	50
SNO+ Phase I [66, 71]	¹³⁰ Te	1357	82	[-0.5, 1.5]	20	97	164	$8.2 \cdot 10^{-2}$	$1.1 \cdot 10^{26}$	46–115	n/a	n/a	n/a
SNO+ Phase II [67]	¹³⁰ Te	7960	57	[-0.5, 1.5]	28	97	1326	$3.6 \cdot 10^{-2}$	$4.8 \cdot 10^{26}$	22–54	n/a	n/a	n/a
KamLAND-Zen 800 [60]	¹³⁶ Xe	750	114	[0, 1.4]	64	97	194	$3.9 \cdot 10^{-2}$	$1.6 \cdot 10^{26}$	47–108	1.5	1	2.1
KamLAND2-Zen [60]	¹³⁶ Xe	1000	60	[0, 1.4]	80	97	325	$2.1 \cdot 10^{-3}$	$8.0 \cdot 10^{26}$	21–49	15	2	2.9
nEXO [72]	¹³⁶ Xe	4507	25	[-1.2, 1.2]	60	85	1741	$4.4 \cdot 10^{-4}$	$4.1 \cdot 10^{27}$	9–22	400	1.2	30
NEXT 100 [64, 73]	¹³⁶ Xe	91	7.8	[-1.3, 2.4]	88	37	26.5	$4.4 \cdot 10^{-2}$	$5.3 \cdot 10^{25}$	82–189	n/a	1	20
NEXT 1.5k [74]	¹³⁶ Xe	1367	5.2	[-1.3, 2.4]	88	37	398	$2.9 \cdot 10^{-3}$	$7.9 \cdot 10^{26}$	21–49	n/a	1	300
PandaX-III 200 [65]	¹³⁶ Xe	180	31	[-2, 2]	100	35	60.2	$4.2 \cdot 10^{-2}$	$8.3 \cdot 10^{25}$	65–150	n/a	n/a	n/a
PandaX-III 1k [65]	¹³⁶ Xe	901	10	[-2, 2]	100	35	301	$1.4 \cdot 10^{-3}$	$9.0 \cdot 10^{26}$	20–46	n/a	n/a	n/a

Will enter IH regime soon!

Multi-isotope determination for
mechanism and NMEs!

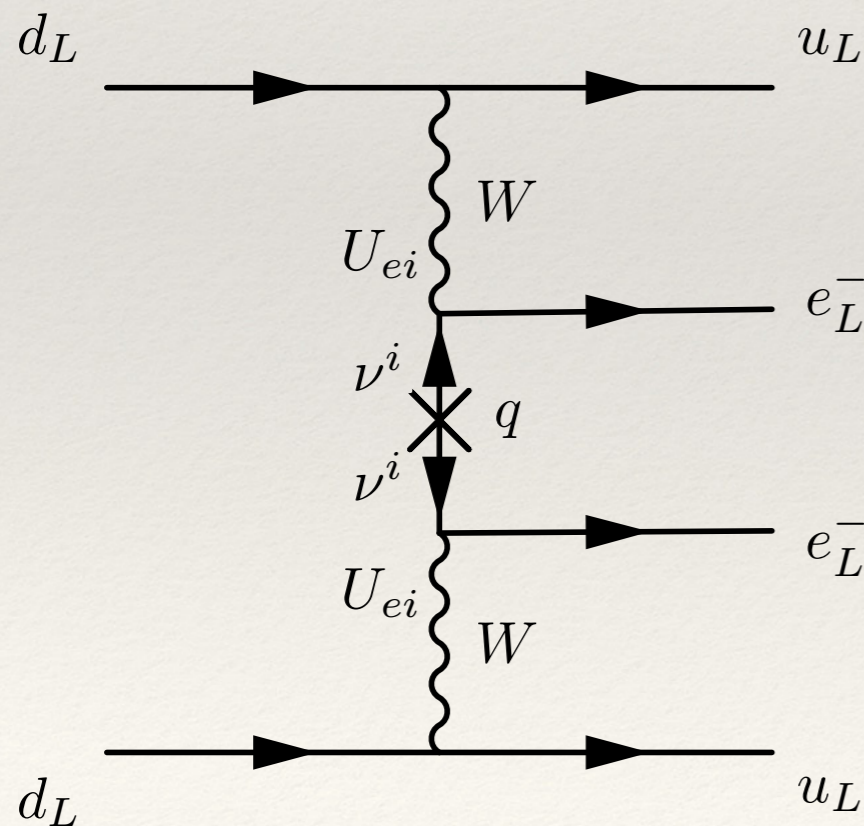
Agostini et al, 1705.02996

New Physics in Double Beta Decay

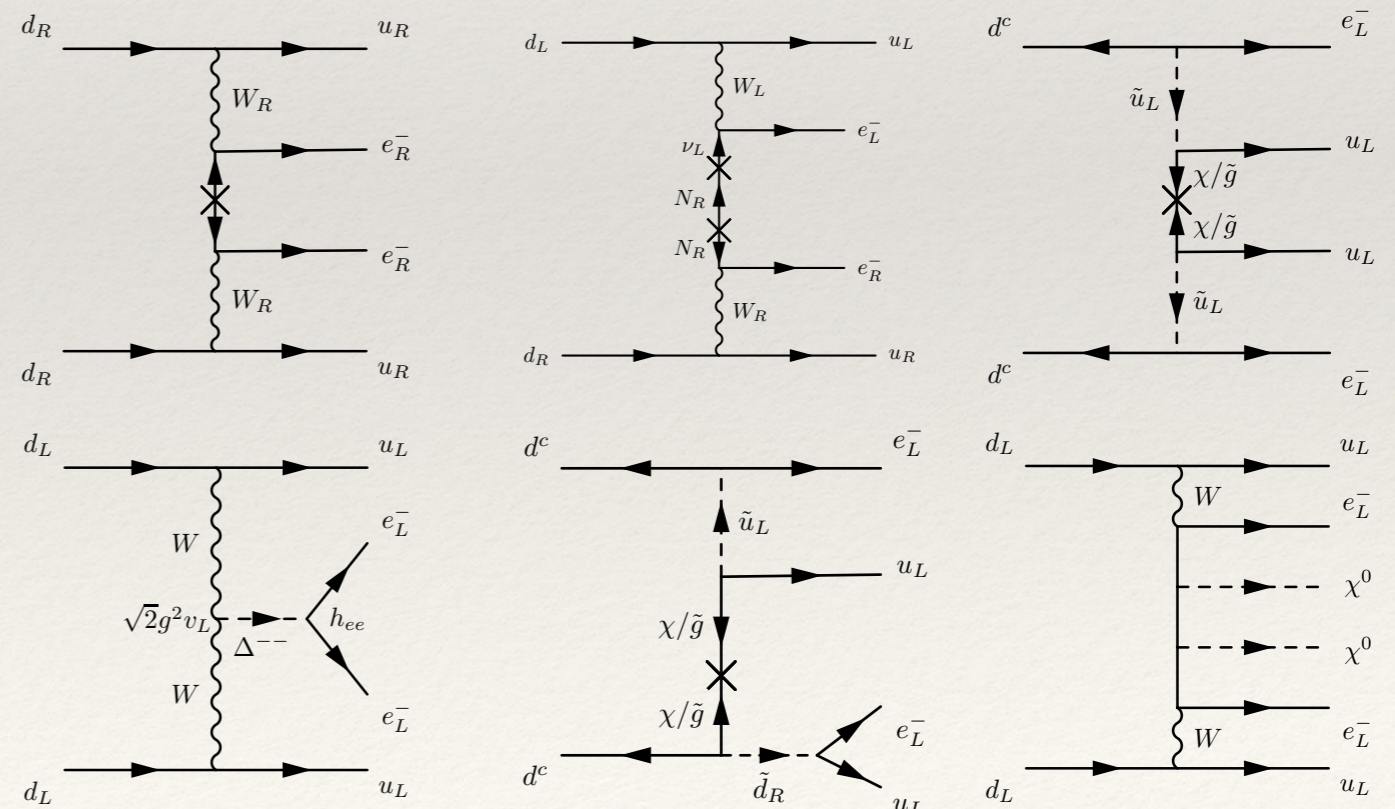
Double Beta Decay is $\Delta L = 2$, not neutrino mass!

Interpretations:

Standard:



Non-Standard:



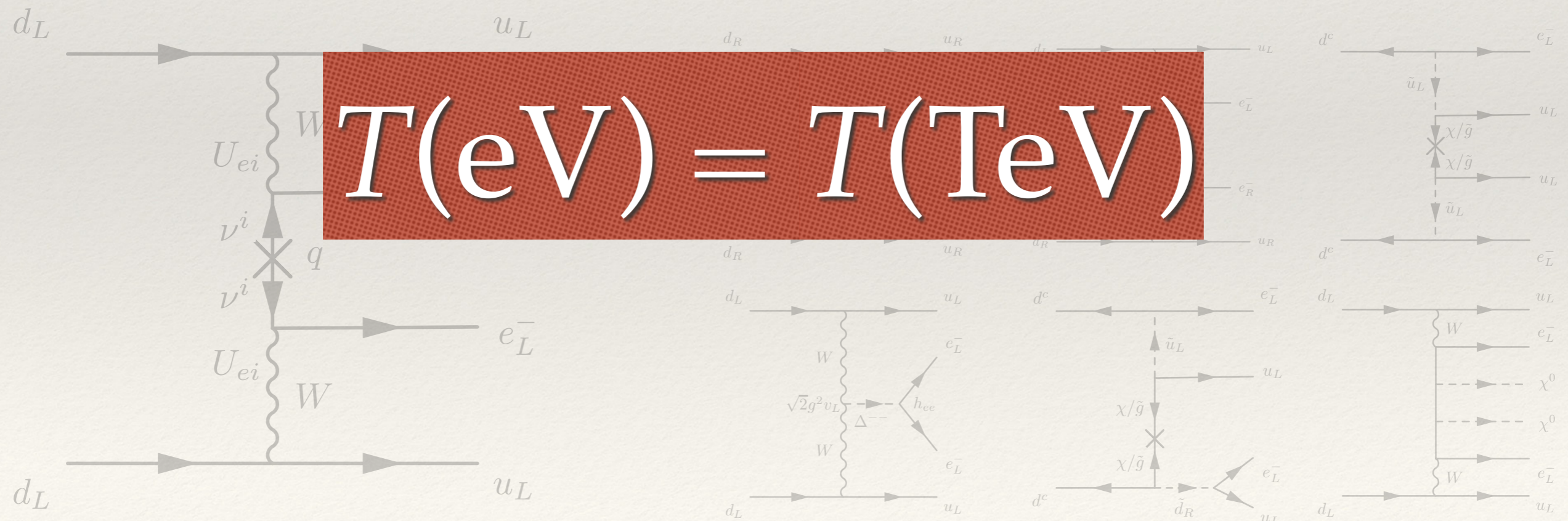
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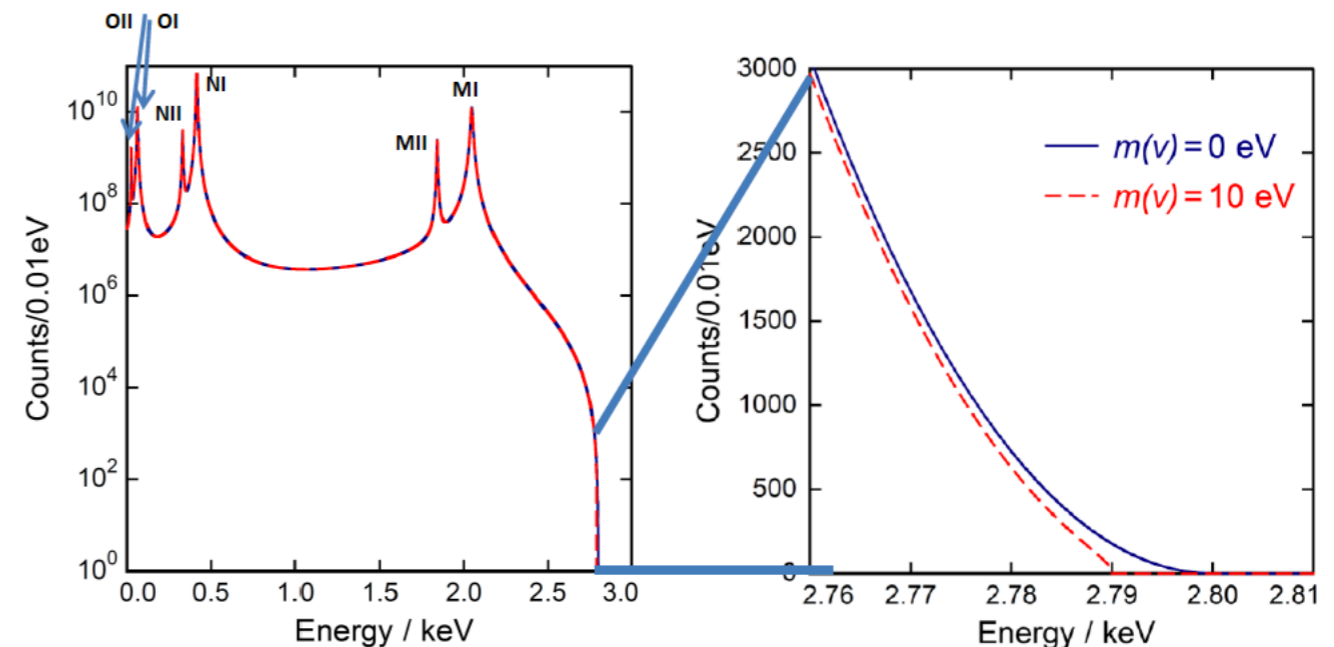
$$T(\text{eV}) = T(\text{TeV})$$

⇒ Tests with LHC, LFV, etc.

see talks by Deppisch, Patra

Direct Neutrino Mass Determination

There are 2 running experiments!!



- ❖ Tritium since May 2018, first ν -mass data in early 2019...
- ❖ (already plans for future versions, aiming at keV- ν , exotic interactions,...)
- ❖ ECHo (EC on ^{163}Ho), spectrum to be measured with low T micro-calorimeters
- ❖ first limit coming soon...

New Physics in Oscillations

- ❖ Various good reasons to expect NP:
 - unitarity violation from new fermions
 - NSIs from new physics
 - new interactions (scalar, tensor, etc.)
 - long-range forces
 - decay, Pseudo-Dirac,...
 - Lorentz / CPT violation: effects $\propto \Lambda/M_{Pl}$ with Λ scale of mass generation (seesaw!), in general growing with ν -energy (IC!)
 - light sterile neutrinos...

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Non-Standard Interactions

$$\mathcal{L} \supset \frac{G_F}{\sqrt{2}} \sum_{q=u,d,e} \bar{\nu}_\alpha \gamma^\mu (1 - \gamma^5) \nu_\beta \left[\varepsilon_{\alpha\beta}^{qV} \bar{q} \gamma^\mu q + \varepsilon_{\alpha\beta}^{qA} \bar{q} \gamma^\mu \gamma^5 q \right]$$

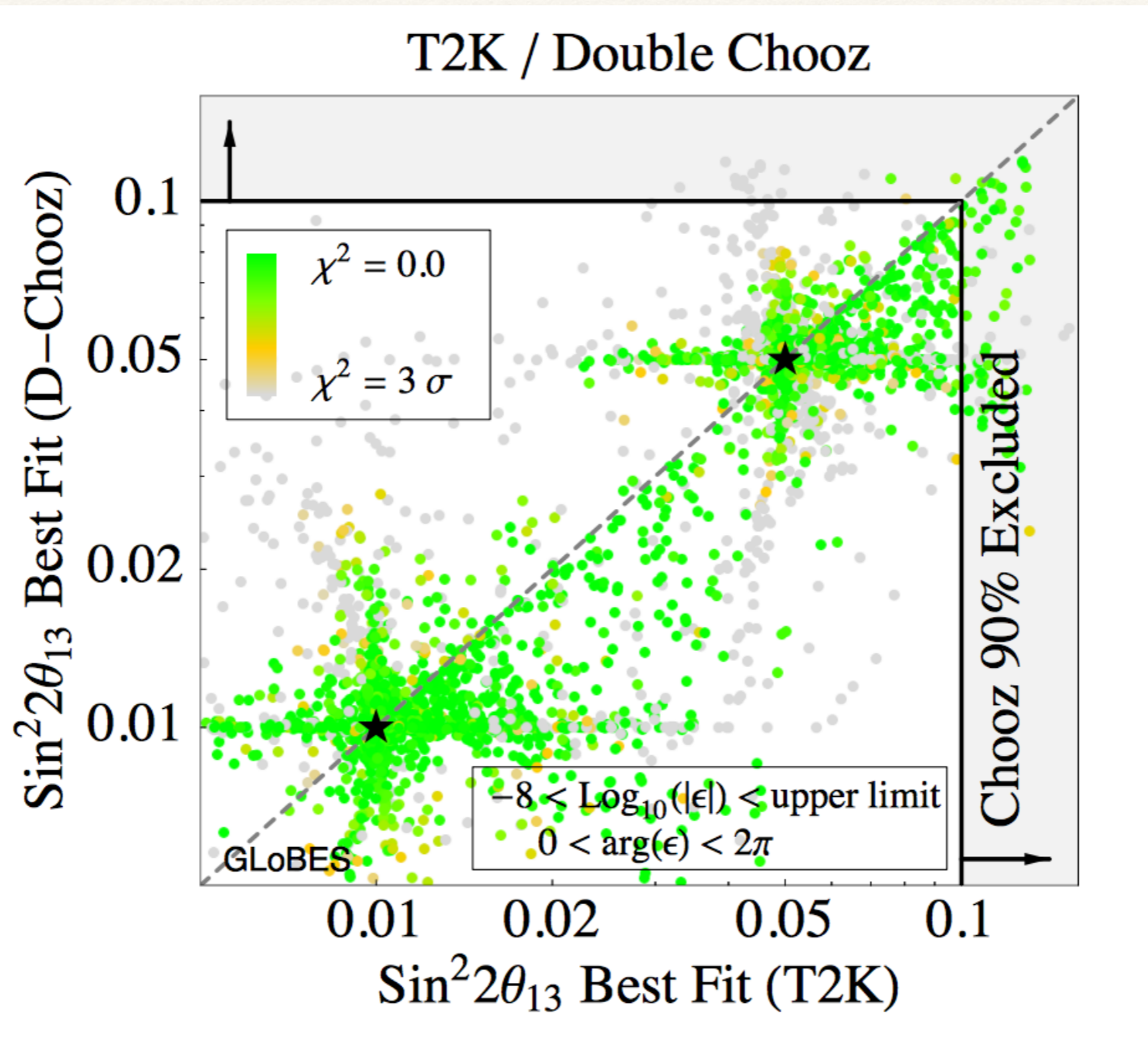
- ❖ $\varepsilon \propto c^2/M_X^2 \Rightarrow \varepsilon = 0.01$ is TeV-scale physics
- ❖ oscillation effect is t -channel forward scattering (q^2 very small), hence c can be very small and M_X MeV-ish
- ❖ can prevent experiments from determining parameters...

Non-Standard Interactions

	LMA	LMA \oplus LMA-D
$\varepsilon_{ee}^u - \varepsilon_{\mu\mu}^u$	$[-0.020, +0.456]$	$\oplus[-1.192, -0.802]$
$\varepsilon_{\tau\tau}^u - \varepsilon_{\mu\mu}^u$	$[-0.005, +0.130]$	$[-0.152, +0.130]$
$\varepsilon_{e\mu}^u$	$[-0.060, +0.049]$	$[-0.060, +0.067]$
$\varepsilon_{e\tau}^u$	$[-0.292, +0.119]$	$[-0.292, +0.336]$
$\varepsilon_{\mu\tau}^u$	$[-0.013, +0.010]$	$[-0.013, +0.014]$
$\varepsilon_{ee}^d - \varepsilon_{\mu\mu}^d$	$[-0.027, +0.474]$	$\oplus[-1.232, -1.111]$
$\varepsilon_{\tau\tau}^d - \varepsilon_{\mu\mu}^d$	$[-0.005, +0.095]$	$[-0.013, +0.095]$
$\varepsilon_{e\mu}^d$	$[-0.061, +0.049]$	$[-0.061, +0.073]$
$\varepsilon_{e\tau}^d$	$[-0.247, +0.119]$	$[-0.247, +0.119]$
$\varepsilon_{\mu\tau}^d$	$[-0.012, +0.009]$	$[-0.012, +0.009]$
$\varepsilon_{ee}^p - \varepsilon_{\mu\mu}^p$	$[-0.041, +1.312]$	$\oplus[-3.328, -1.958]$
$\varepsilon_{\tau\tau}^p - \varepsilon_{\mu\mu}^p$	$[-0.015, +0.426]$	$[-0.424, +0.426]$
$\varepsilon_{e\mu}^p$	$[-0.178, +0.147]$	$[-0.178, +0.178]$
$\varepsilon_{e\tau}^p$	$[-0.954, +0.356]$	$[-0.954, +0.949]$
$\varepsilon_{\mu\tau}^p$	$[-0.035, +0.027]$	$[-0.035, +0.035]$

Esteban et al., 1805.04530

Non-Standard Interactions



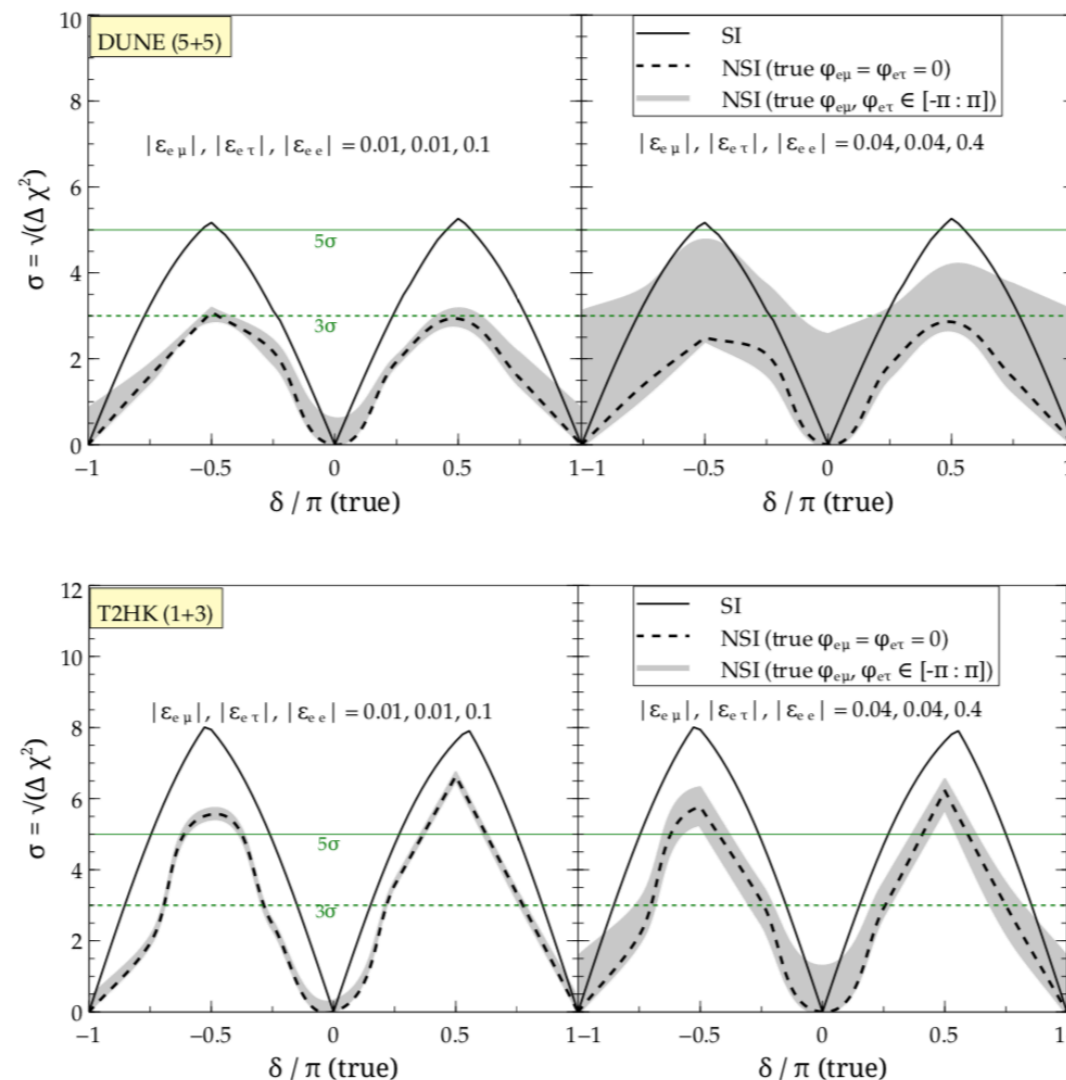
Kopp, Lindner, Ota, Sato, 0708.0152

Non-Standard Interactions

CPV sensitivity - DUNE, T2HK

- DUNE (1300 km)
Runtime = 5 nu + 5 nu bar
35 kton, LArTPC

- T2HK
Runtime = 1 nu + 3 nu bar
560 kton, WC

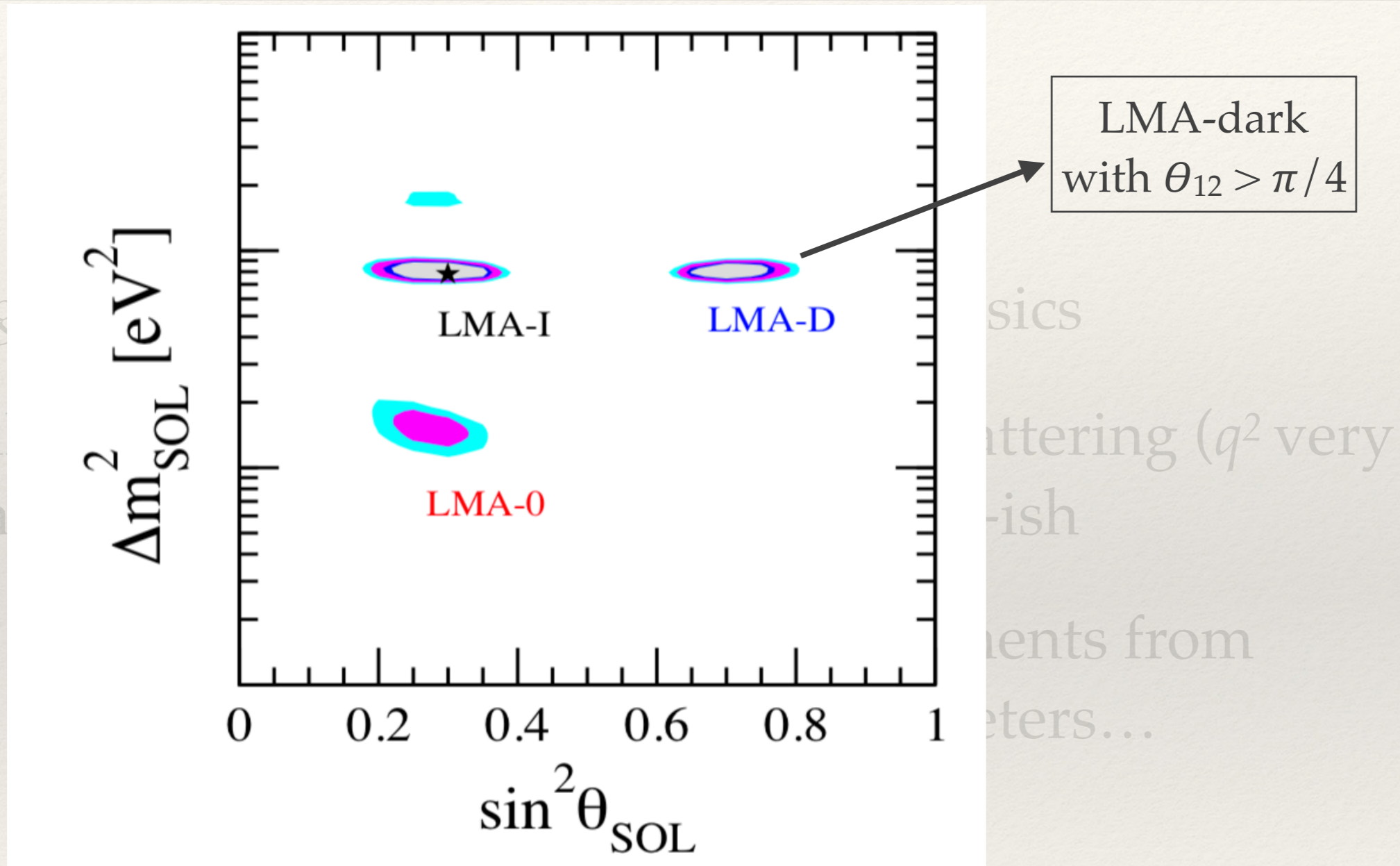


M. Masud and P. Mehta, Phys. Rev. D (2016) [1603.01389]

presence of multiple $\epsilon_{\alpha\beta}$ can make determination of MO, δ and octant of θ_{23} impossible even for DUNE, T2HK and T2HKK

Liao, Marfatia, Whisnant, 1612.01443

Non-Standard Interactions



Miranda, Tortola, Valle, hep-ph/0406280

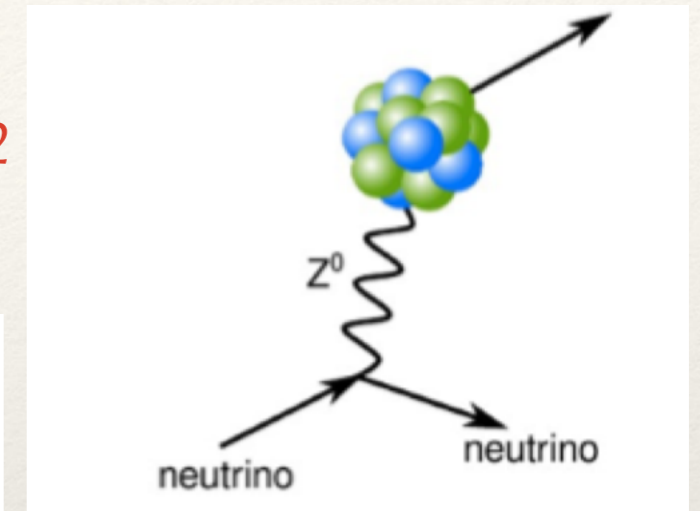
(can also explain small Δm^2 discrepancy in KamLAND/solar and missing upturn of P_{ee})

Coherent Elastic Neutrino-Nucleus Scattering

Freedmann, PRD9, 1974

$$\frac{d\sigma}{dT} = \frac{\sigma_0^{\text{SM}}}{M} \left(1 - \frac{T}{T_{\text{max}}}\right) \propto N^2$$

$$\sigma_0^{\text{SM}} \equiv \frac{G_F^2 [N - (1 - 4s_W^2)Z]^2 F^2(q^2) M^2}{4\pi}$$



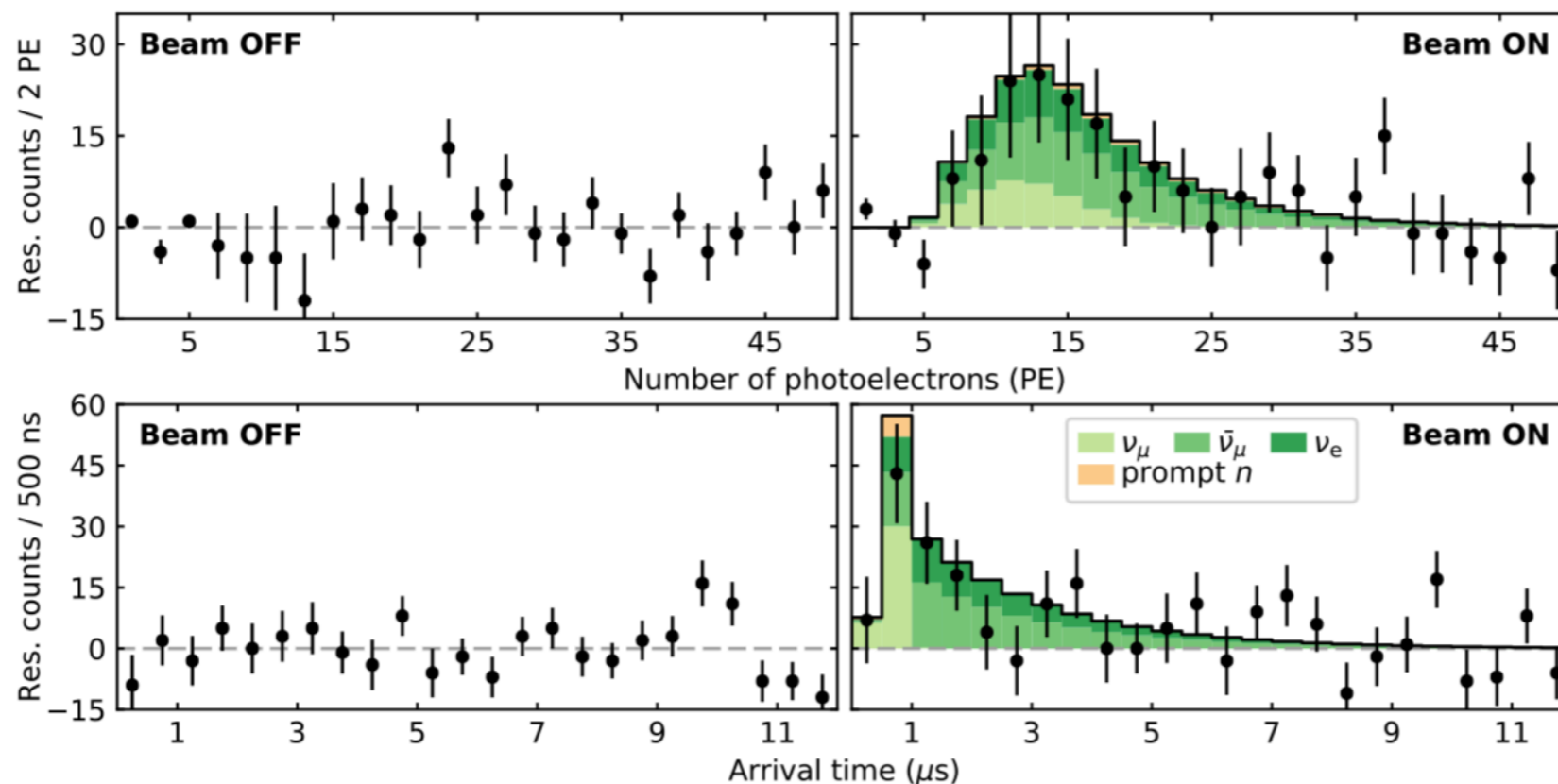
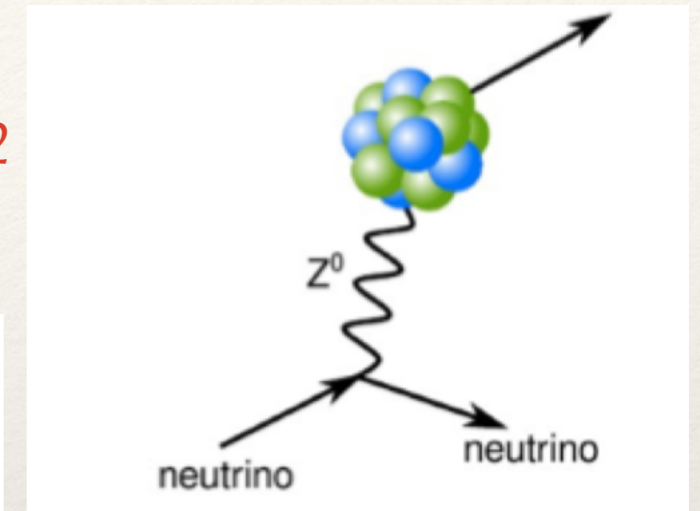
- ❖ needs E_ν below 50 MeV
- ❖ \Rightarrow pion decay or reactors as ν -source
- ❖ \Rightarrow low nuclear recoil below few keV
- ❖ \Rightarrow sensitive detectors and smart shielding

Coherent Elastic Neutrino-Nucleus Scattering

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6.7 σ detection
with pulsed
pion source
at SNS

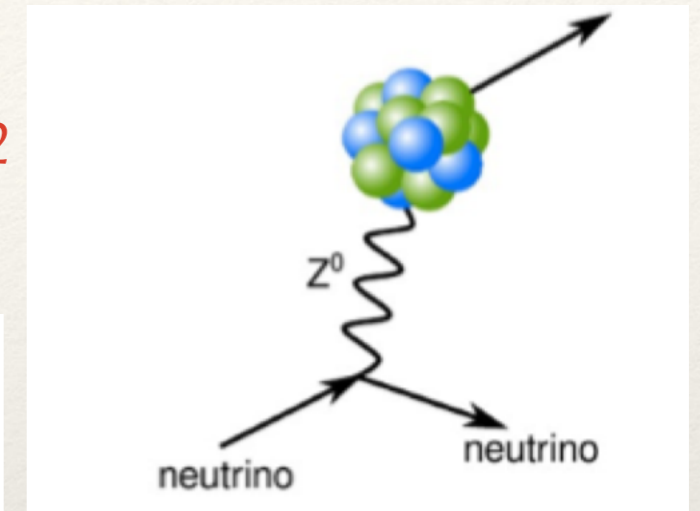
*COHERENT,
Science 357 (2017)*

Coherent Elastic Neutrino-Nucleus Scattering

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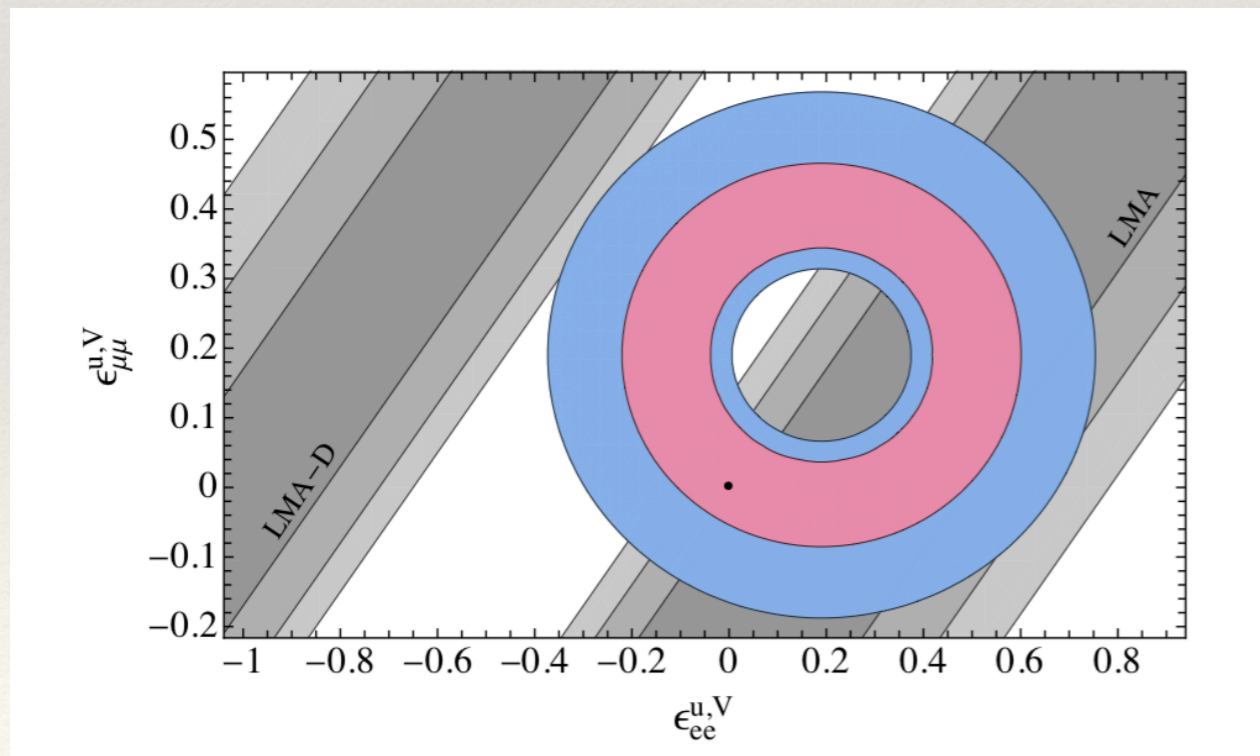


- ❖ last missing ν -cross section in SM (largest one...)
- ❖ helps SN explode
- ❖ neutron charge density \leftrightarrow neutron skin \leftrightarrow NS eos
- ❖ ultimate background for DM direct detection
- ❖ measurement of θ_W at low energies
- ❖ NSIs, exotic NC, Z' , sterile ν ,...

New Physics in Coherent Scattering

if NSIs are present: replace $[N - (1 - 4 s_W^2) Z]^2$ with:

$$Q_{\text{NSI}}^2 \equiv 4 \left[N \left(-\frac{1}{2} + \varepsilon_{ee}^{uV} + 2\varepsilon_{ee}^{dV} \right) + Z \left(\frac{1}{2} - 2s_W^2 + 2\varepsilon_{ee}^{uV} + \varepsilon_{ee}^{dV} \right) \right]^2 + 4 \sum_{\alpha=\mu,\tau} \left[N(\varepsilon_{\alpha e}^{uV} + 2\varepsilon_{\alpha e}^{dV}) + Z(2\varepsilon_{\alpha e}^{uV} + \varepsilon_{\alpha e}^{dV}) \right]^2 .$$

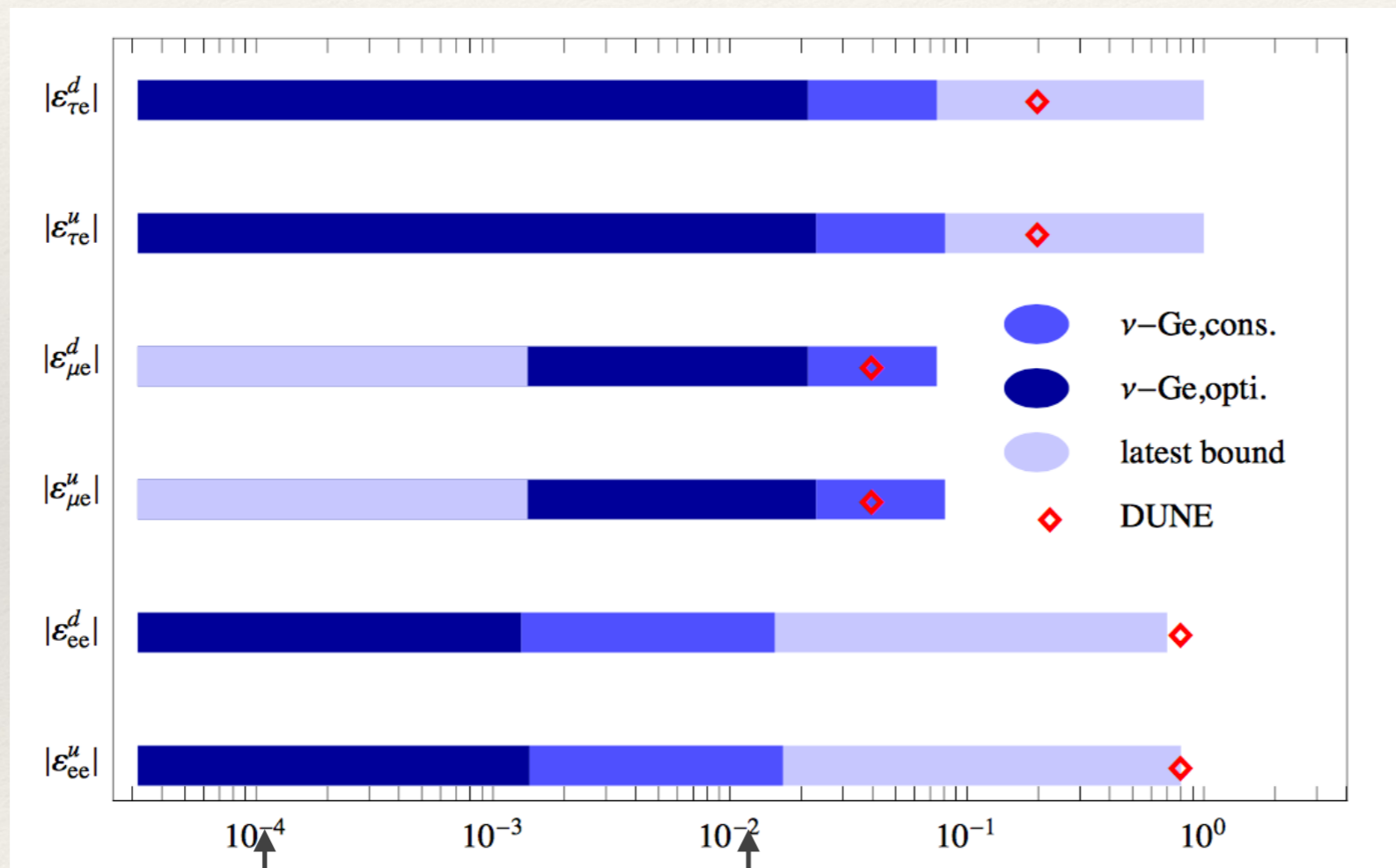


disfavors LMA-dark solution
with more than 3σ

Coloma et al., 1708.02899

New Physics in Coherent Scattering

Example: CONUS-100 like, BG 3 / day / kg / keV,
exposure: 5 kg yr GW m⁻², sys / stat / thresh.:



10 TeV

TeV

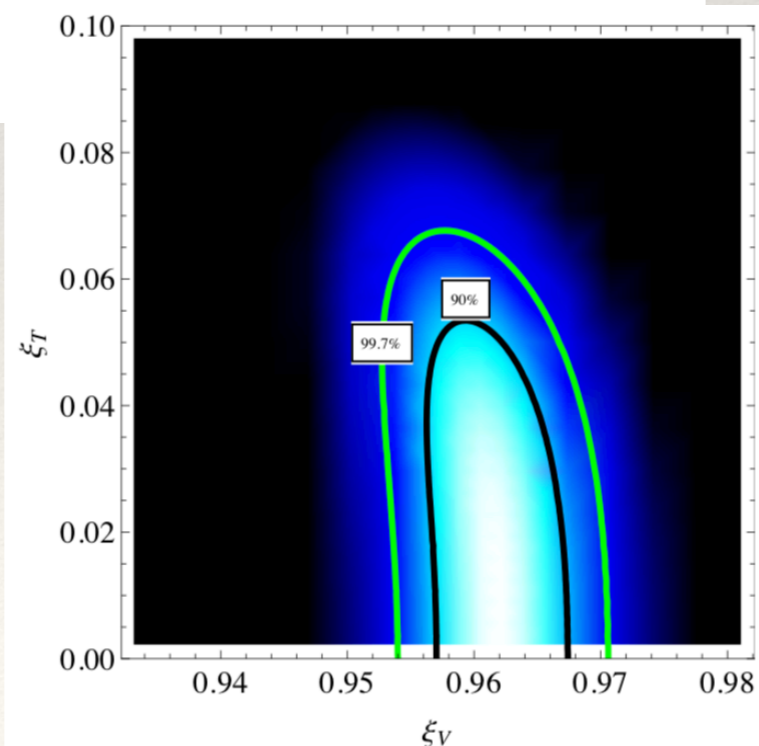
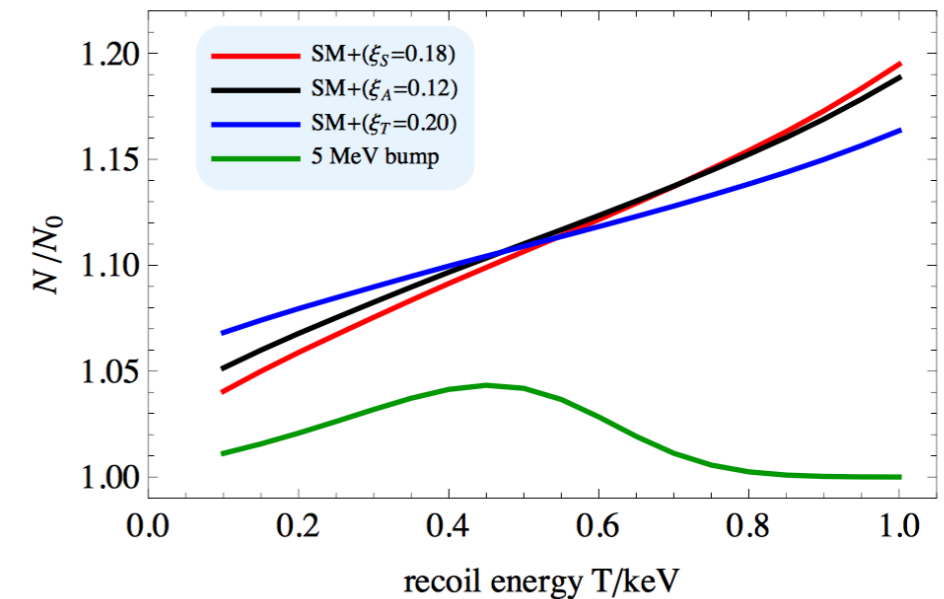
Lindner, WR, Xu, 1612.04150

New Physics in Coherent Scattering

assume exotic neutral currents:

$$\mathcal{L} \supset \frac{G_F}{\sqrt{2}} \sum_{a=S,P,V,A,T} \bar{\nu} \Gamma^a \nu \left[\overline{\psi}_N \Gamma^a (C_a + \overline{D}_a i \gamma^5) \psi_N \right]$$

$$\begin{aligned} \frac{d\sigma}{dT} = & \frac{G_F^2 M}{4\pi} N^2 \left[\xi_S^2 \frac{MT}{2E_\nu^2} \right. \\ & + \xi_V^2 \left(1 - \frac{T}{T_{\max}} \right) - 2\xi_V \xi_A \frac{T}{E_\nu} + \xi_A^2 \left(1 - \frac{T}{T_{\max}} + \frac{MT}{E_\nu^2} \right) \\ & + \xi_T^2 \left(1 - \frac{T}{T_{\max}} + \frac{MT}{4E_\nu^2} \right) \\ & \left. - R \frac{T}{E_\nu} + \mathcal{O} \left(\frac{T^2}{E_\nu^2} \right) \right], \end{aligned}$$



Lindner, WR, Xu, 1612.04150

changes *shape of spectrum*:

Summary

- ❖ Neutrinos still only testable BSM physics
- ❖ PMNS parameters approach CKM-precision
- ❖ first hints on CP, mass ordering
- ❖ standard paradigm tested on all fronts
- ❖ still new windows open up...

how to „predict“ the CP phase: sum-rules

$$U_\nu = \begin{pmatrix} \cos \theta & \sin \theta & 0 \\ -\sin \theta / \sqrt{2} & \cos \theta / \sqrt{2} & \sqrt{\frac{1}{2}} \\ \sin \theta / \sqrt{2} & \cos \theta / \sqrt{2} & \sqrt{\frac{1}{2}} \end{pmatrix} \text{ and } U_\ell \sim \text{CKM}$$

$\Rightarrow \sin^2 \theta_{12} \simeq \sin^2 \theta - |U_{e3}| \sin 2\theta \cos \delta$

*King et al.; Frampton,
Petcov, WR,...*

- ❖ if $\sin^2 \theta = 1/3 = 0.33$ (tri-bimaximal, e.g. A_4, S_4, T')
- ❖ if $\sin^2 \theta = 1/2 = 0.50$ (bimaximal, e.g. D_4)
- ❖ if $\sin^2 \theta = 1/4 = 0.25$ (hexagonal, e.g. D_{12})
- ❖ if $\tan \theta = 1/\phi$ or $\sin^2 \theta = 0.276$ (GRA, e.g. A_5)
- ❖ if $\cos \theta = \phi/2$ or $\sin^2 \theta = 0.346$ (GRB, e.g. D_{10})

\Rightarrow can distinguish only classes of models

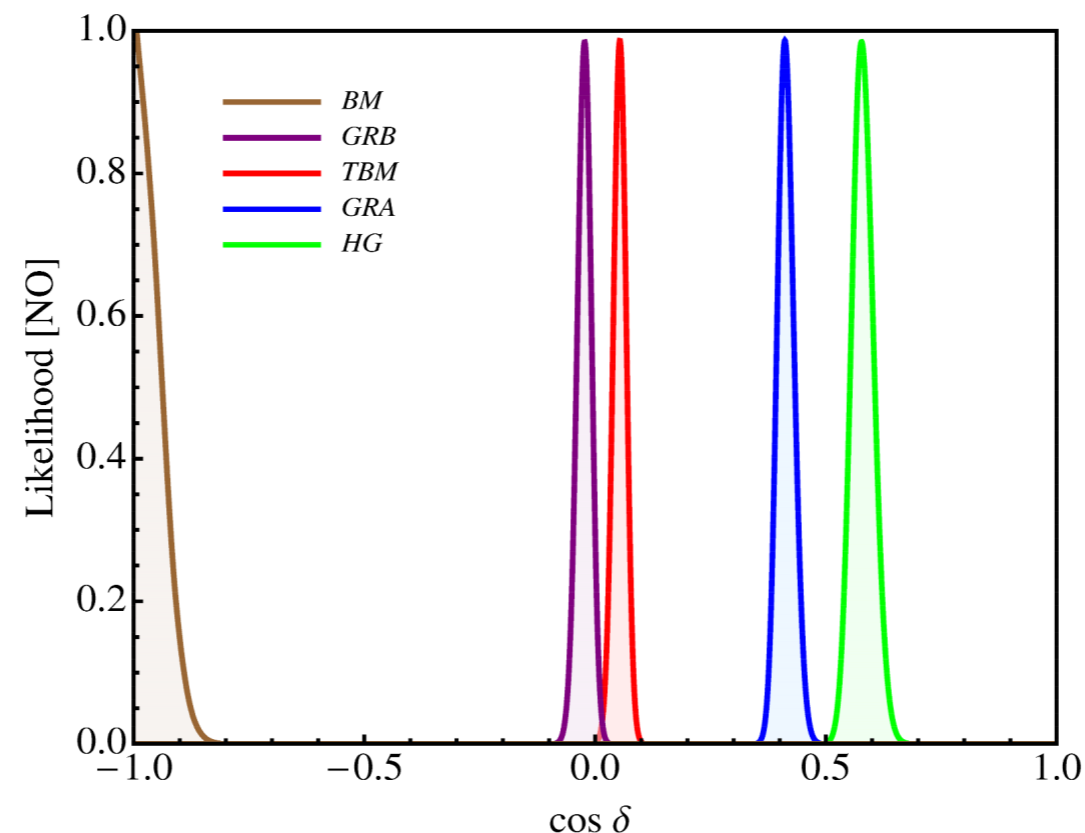
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*King et al.; Frampton,
Petcov, WR,...*

*Girardi, Petcov, Titov,
1410.8056*



\Rightarrow can distinguish only classes of models

Oscillation Parameters

- ❖ Maximal θ_{23} preferred by LBL, slight 1-2 σ shift to $> \pi/4$ by SK
- ❖ LBL prefer $\delta \approx 3\pi/2$, driven by (too many?) ν_e ; also SK due to sub-GeV e -like events
- ❖ normal mass ordering preferred by LBL (tension with reactors) and SK (excess of upward going e -like events), $\approx 2\sigma$ effect each, $\approx 3\sigma$ total

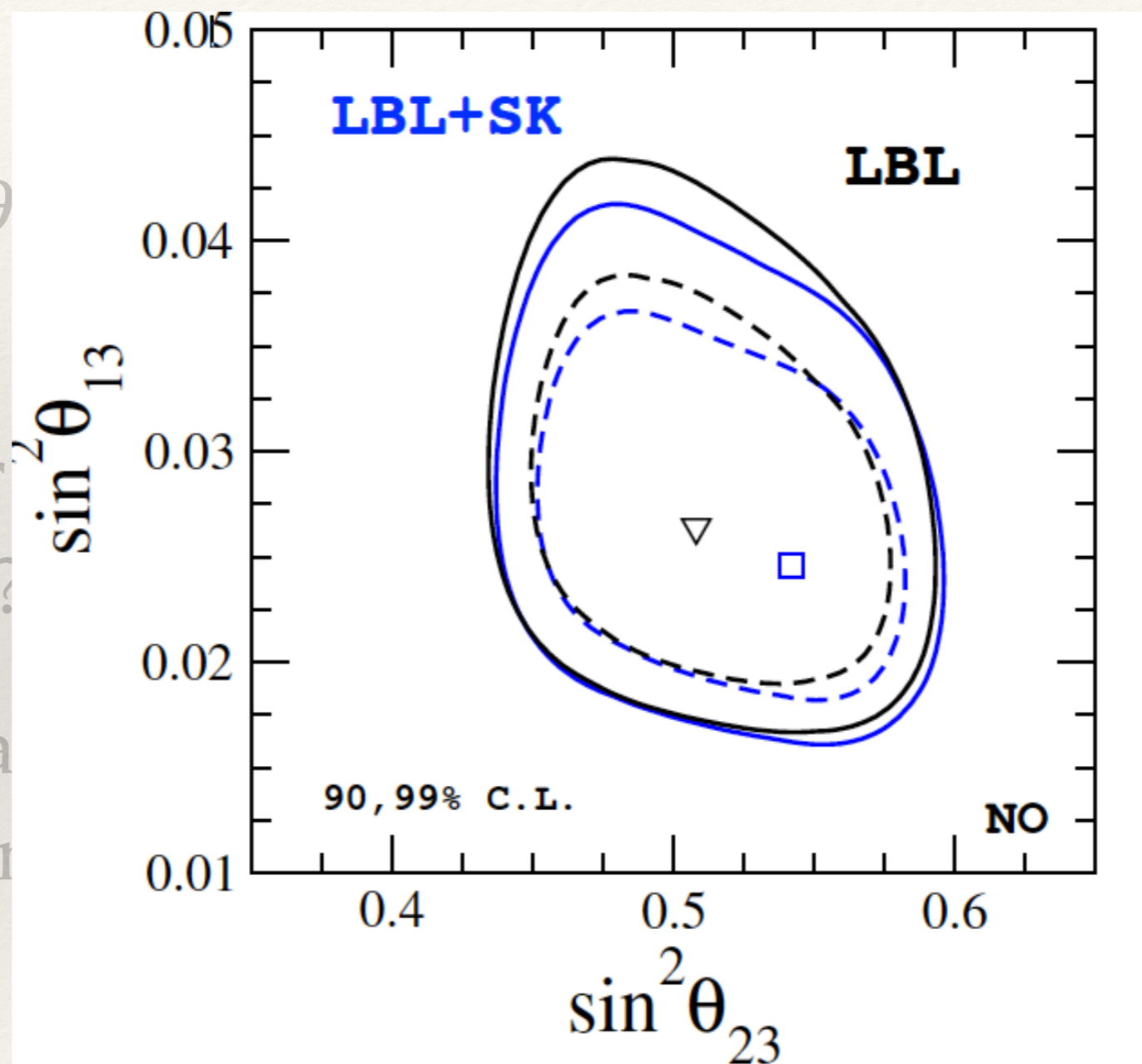
see talks by Sekiguchi, Bhatnagar, Wu, Tanaka

Mass Ordering

- ❖ weak preference for normal ordering
 - tension in the preferred values of θ_{13} in T2K/NO ν A and reactor, found to be stronger for the case of inverted mass ordering
 - tension in the preferred values of Δm^2_{31} in T2K/NO ν A and reactor, found to be stronger for the case of inverted mass ordering
 - e -like multi-GeV events in SK
 - supported by strongest cosmological mass bounds
 - ❖ BUT: depends on sampling with logarithmic or linear prior, using m_i or $m_{sm} + \Delta m^2$ (*Gariazzo et al., 1801.04946, Hannestad and Schwetz, 1606.04691*)

Oscillation Parameters

- ❖ Maximal θ_{13} by SK
- ❖ LBL preferred (too many?)
- ❖ normal mass (reactors) are $\approx 2\sigma$ effect



shift to $> \pi/4$

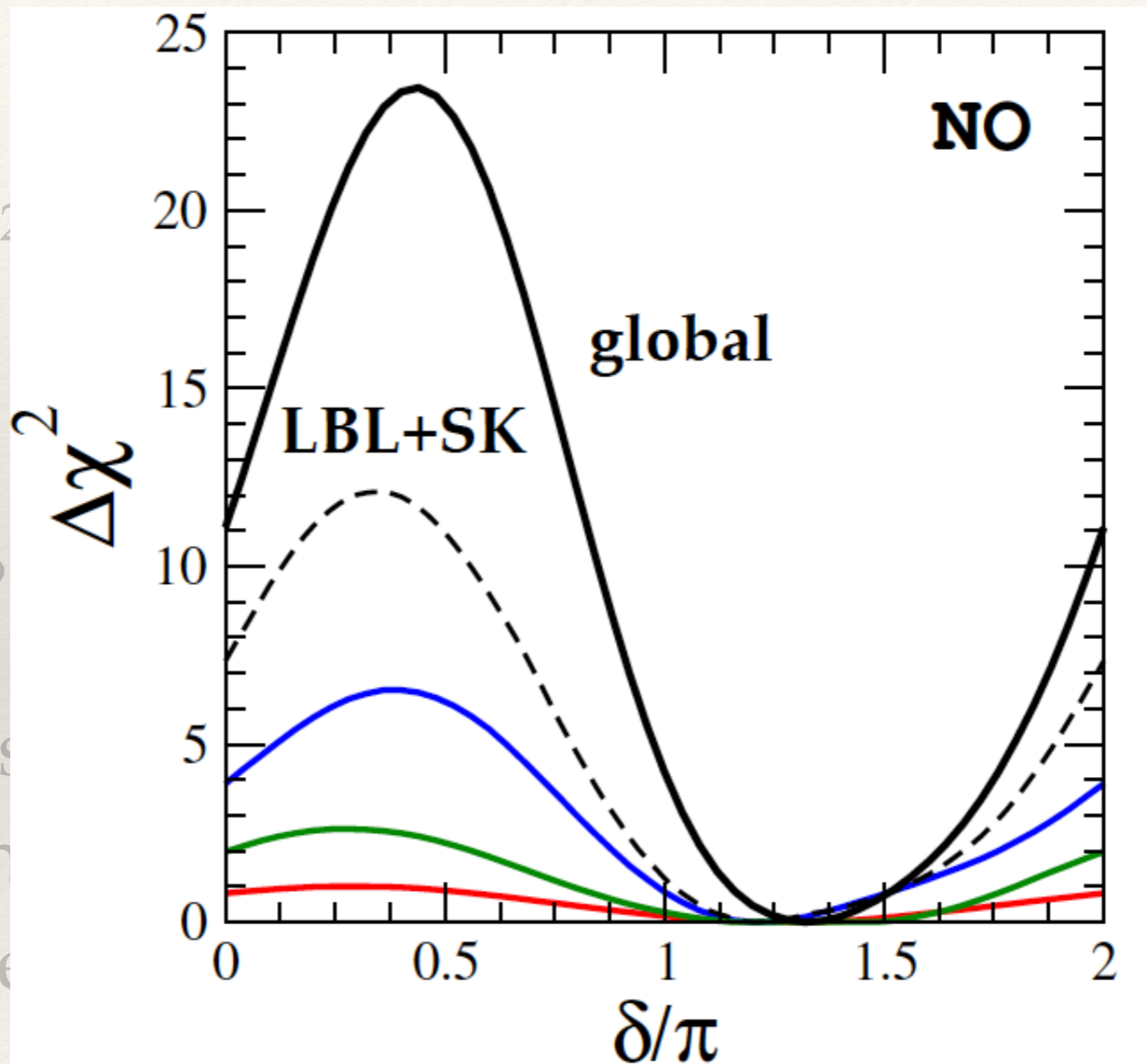
the events

consistency with
like events),

deSalas et al, 1708.01186

Oscillation Parameters

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- ❖ normal mass (reactors) are $\approx 2\sigma$ effect



shift to $> \pi/4$

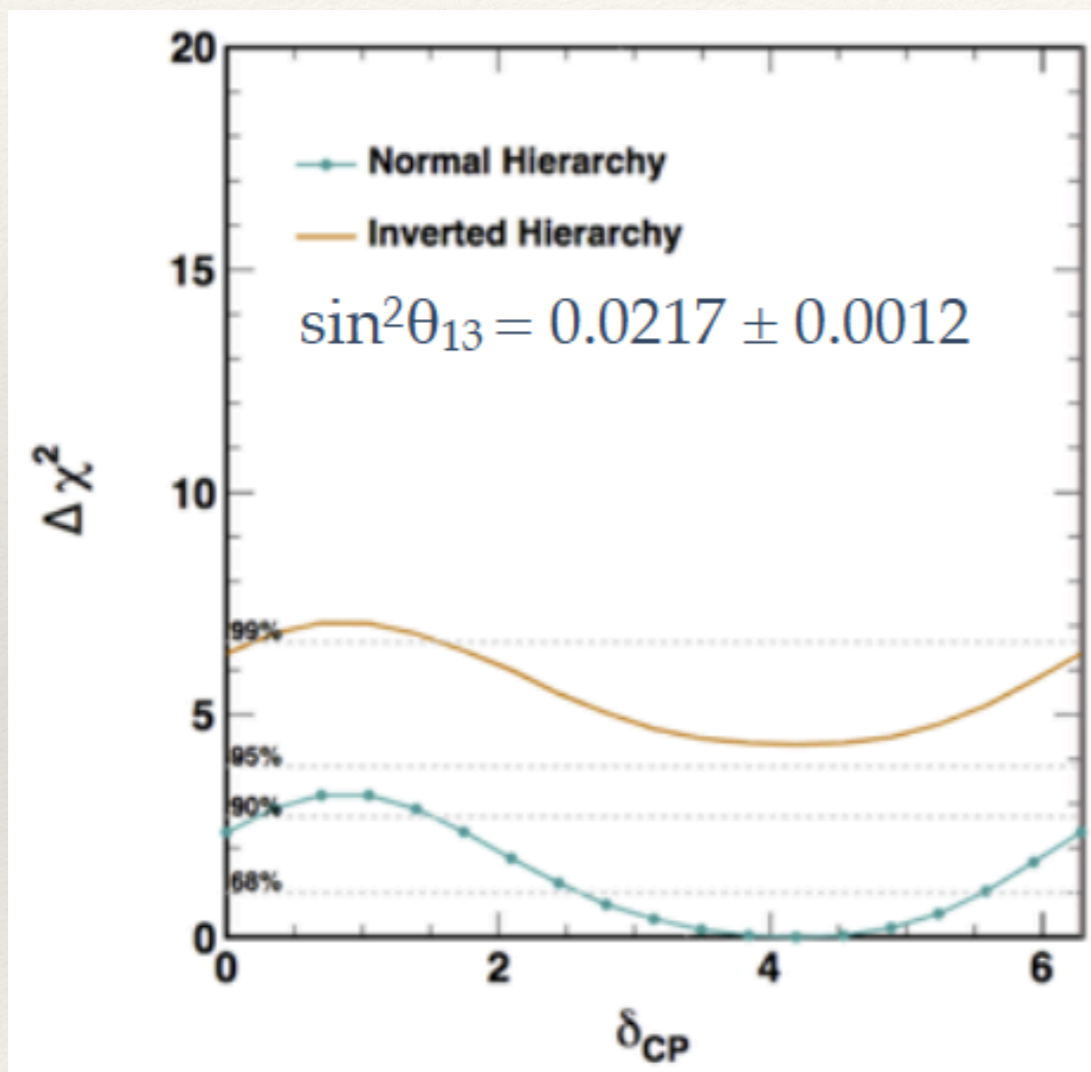
the events

fusion with

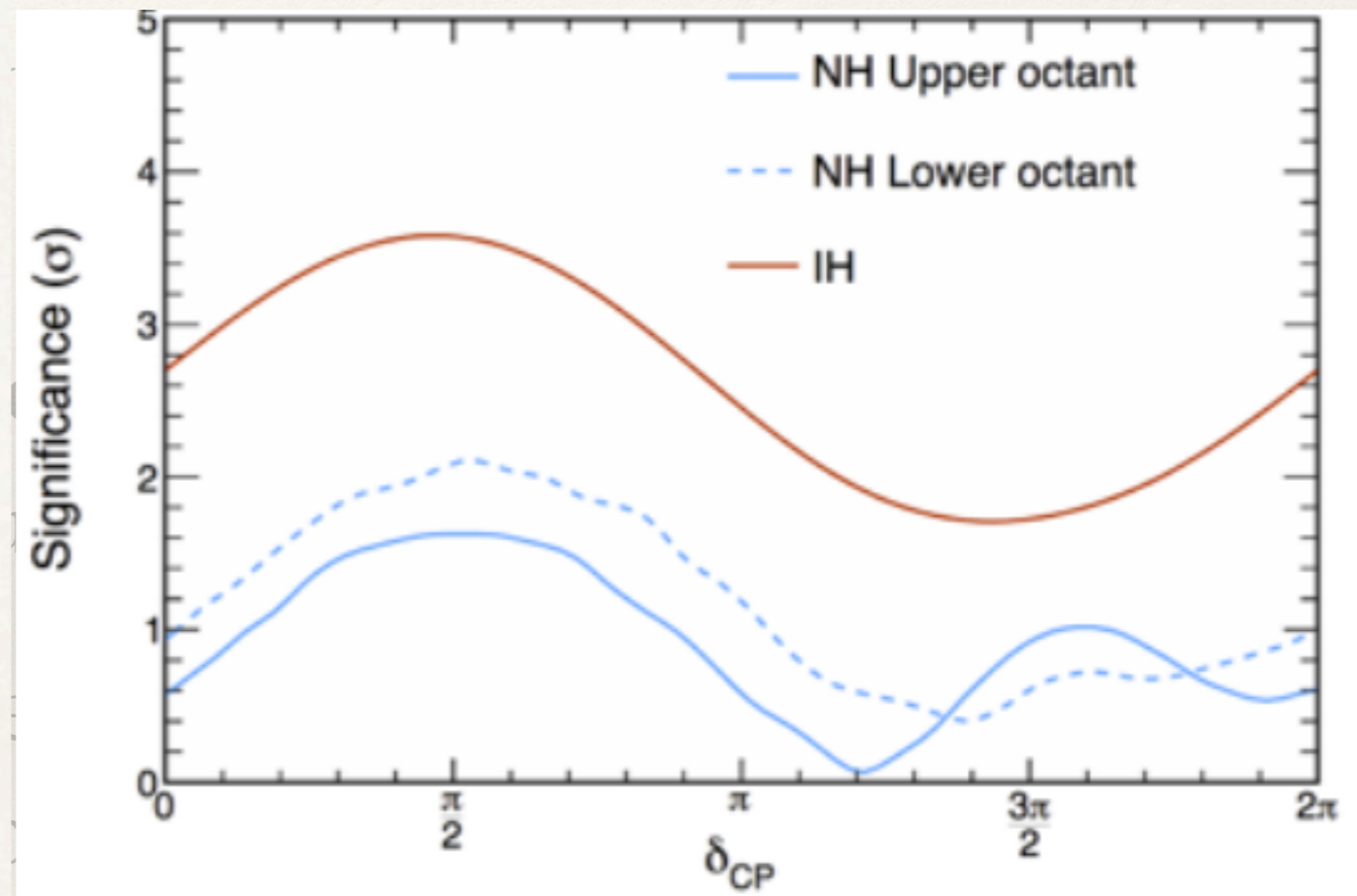
like events),

deSalas et al, 1708.01186

Oscillation Parameters



SK, 1710.09126



NO ν A, 1806.00096

New data presented @Nu2018

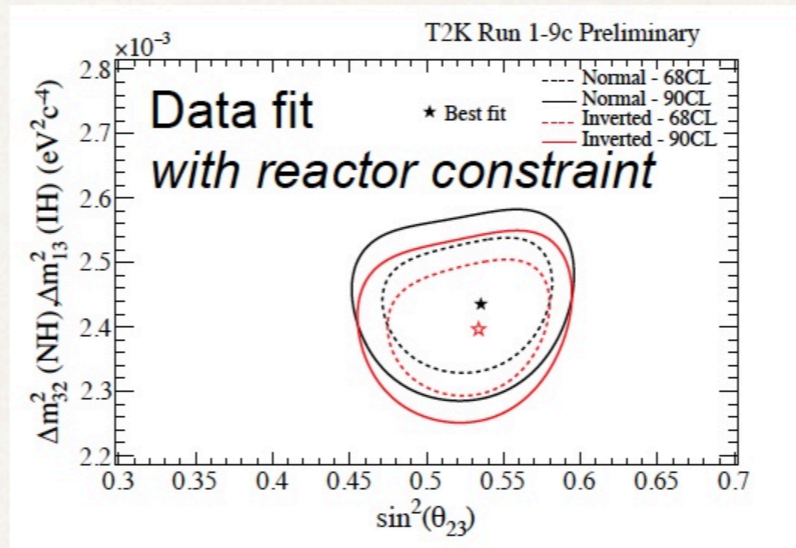
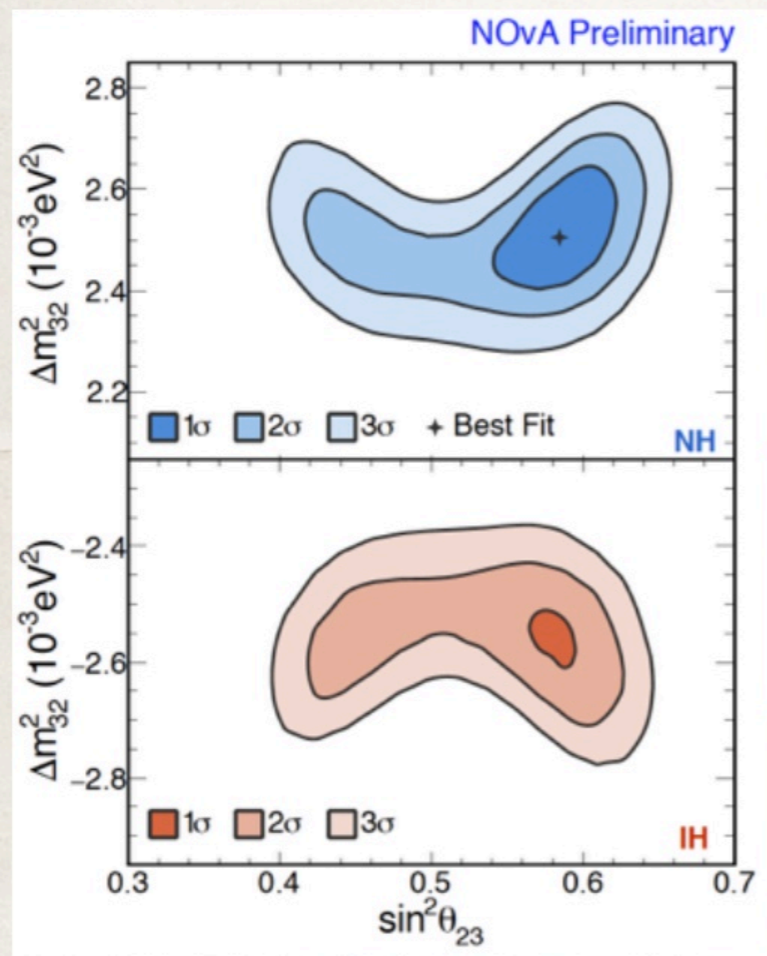
First NOvA antineutrino data

New T2K antineutrino data

Monday
June 4th

Talk by M. Sánchez

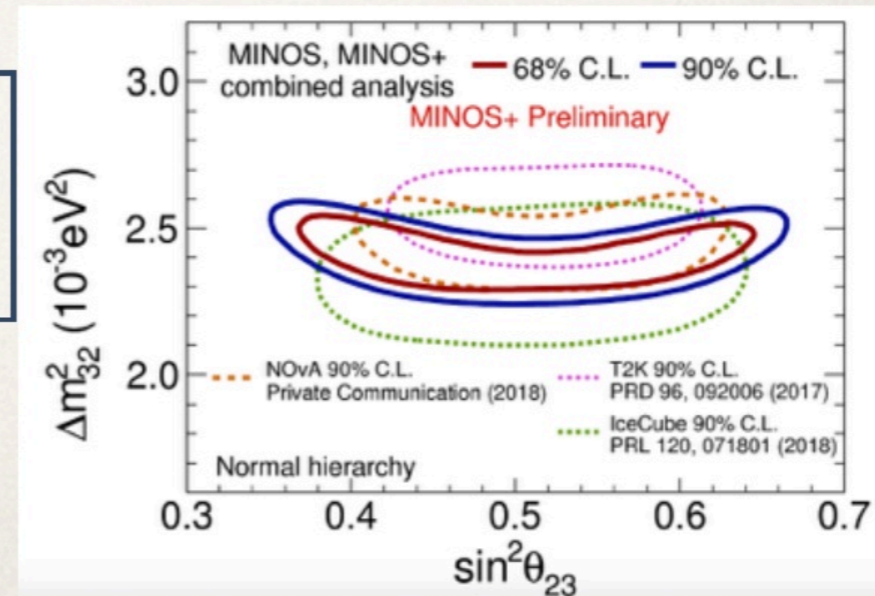
neutrino + antineutrino fit



Talk by
M. Wascko

New combined analysis MINOS/MINOS+

Talk by
A. Aurisano

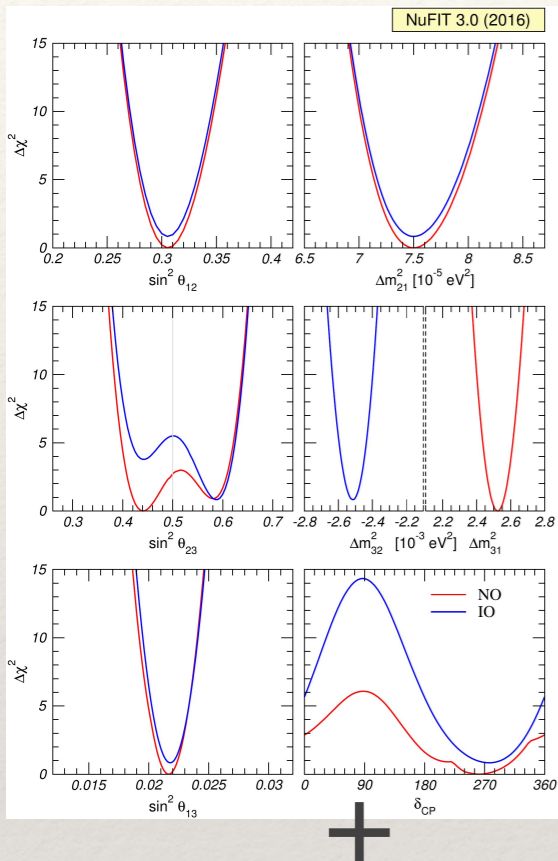


...probably adds another σ to each hint...

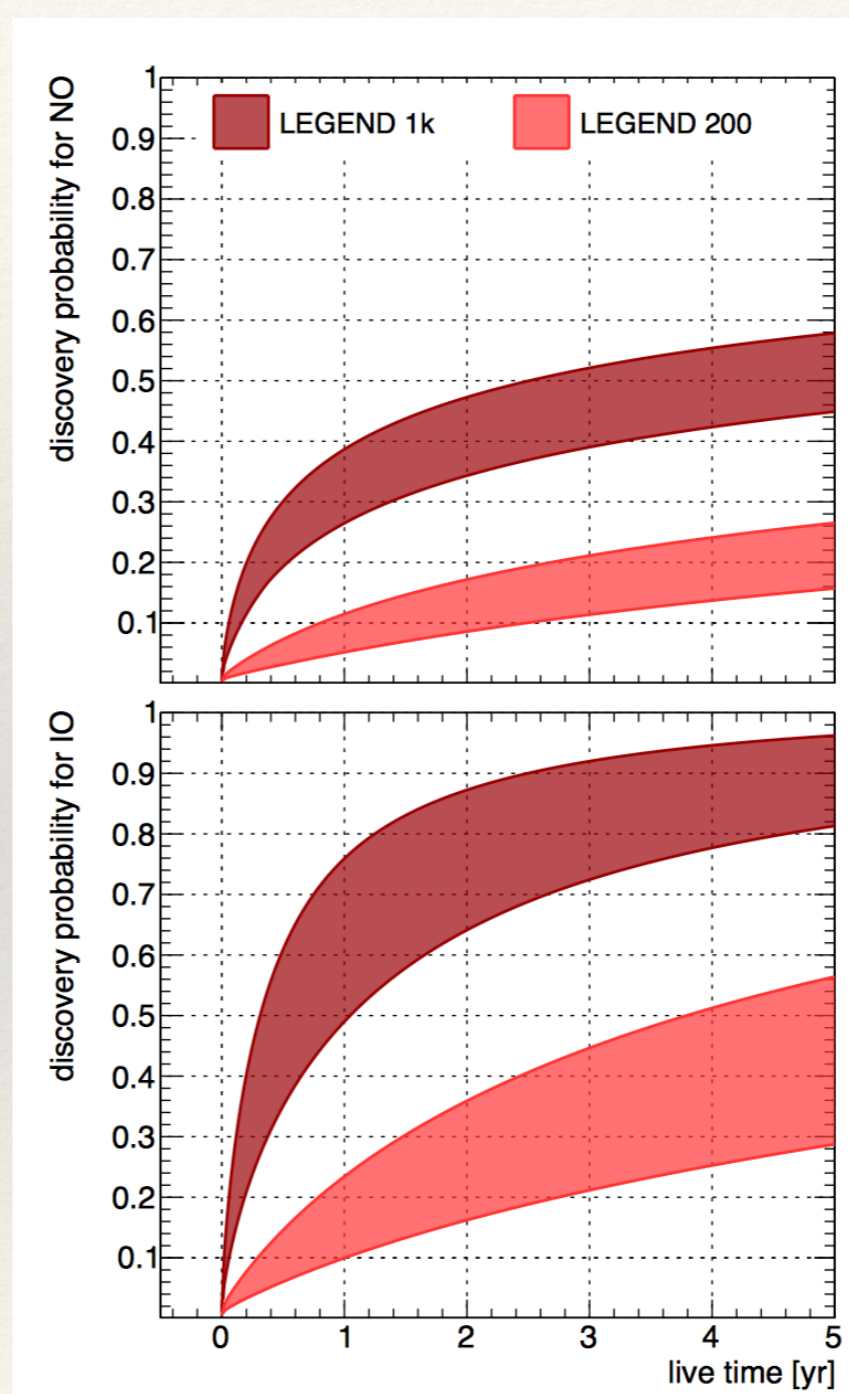
see talks by Sekiguchi, Bhatnagar

Tortola, talk at Neutrino 2018

Expectations of lifetimes



Oscillation fits
 expt. sensitiv.

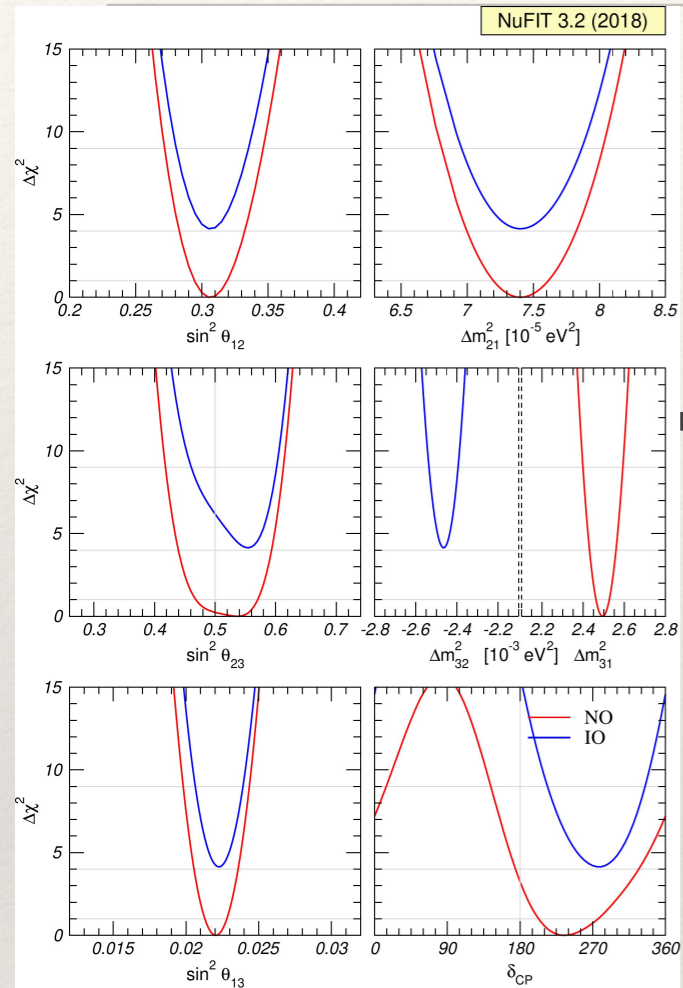


Bayesian discovery probability: discovery sensitivity (value of m_{ee} for which expt. has 50% chance to see it at 3σ) folded with probability distribution of m_{ee}

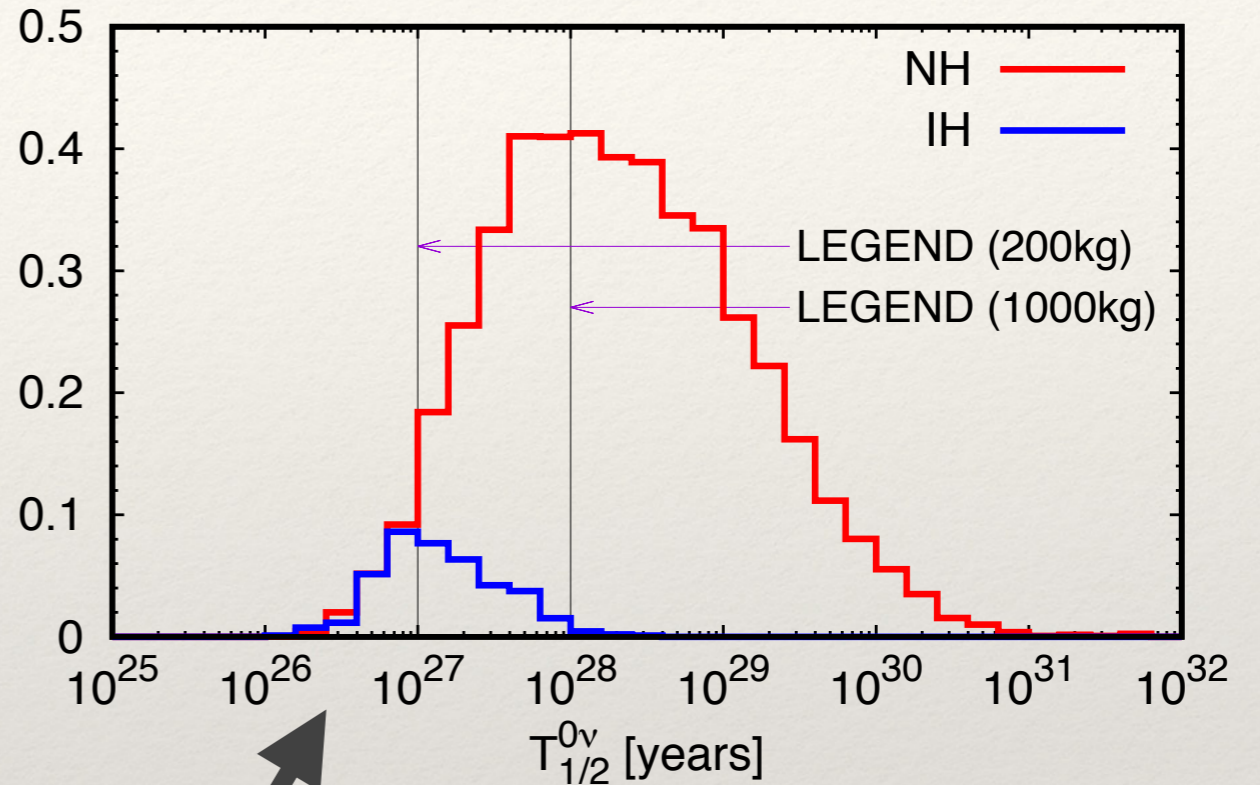
Agostini et al, 1705.02996; also Caldwell et al., 1705.01945; also Zhang, Zhou, 1508.05472

Experiment	Iso.	Iso. Mass [kg _{iso}]	σ [keV]	ROI [σ]	ϵ_{FV} [%]	ϵ_{sig} [%]	\mathcal{E} [kg _{iso} yr / yr]	B [cts / kg _{iso} ROI yr]	3σ disc. sens.		Required Improvement		
									$\hat{T}_{1/2}$ [yr]	$\hat{m}_{\beta\beta}$ [meV]	Bkg	σ	Iso. Mass
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CUPID [58, 59, 70]	⁸² Se	336	2.1	[-2, 2]	100	69	221	$5.2 \cdot 10^{-4}$	$1.8 \cdot 10^{27}$	15-25	n/a	6	n/a
CUORE [52, 53]	¹³⁰ Te	206	2.1	[-1.4, 1.4]	100	81	141	$3.1 \cdot 10^{-1}$	$5.4 \cdot 10^{25}$	66-164	6	1	19
CUPID [58, 59, 70]	¹³⁰ Te	543	2.1	[-2, 2]	100	81	422	$3.0 \cdot 10^{-4}$	$2.1 \cdot 10^{27}$	11-26	3000	1	50
SNO+ Phase I [66, 71]	¹³⁰ Te	1357	82	[-0.5, 1.5]	20	97	164	$8.2 \cdot 10^{-2}$	$1.1 \cdot 10^{26}$	46-115	n/a	n/a	n/a
SNO+ Phase II [67]	¹³⁰ Te	7960	57	[-0.5, 1.5]	28	97	1326	$3.6 \cdot 10^{-2}$	$4.8 \cdot 10^{26}$	22-54	n/a	n/a	n/a
KamLAND-Zen 800 [60]	¹³⁶ Xe	750	114	[0, 1.4]	64	97	194	$3.9 \cdot 10^{-2}$	$1.6 \cdot 10^{26}$	47-108	1.5	1	2.1
KamLAND2-Zen [60]	¹³⁶ Xe	1000	60	[0, 1.4]	80	97	325	$2.1 \cdot 10^{-3}$	$8.0 \cdot 10^{26}$	21-49	15	2	2.9
nEXO [72]	¹³⁶ Xe	4507	25	[-1.2, 1.2]	60	85	1741	$4.4 \cdot 10^{-4}$	$4.1 \cdot 10^{27}$	9-22	400	1.2	30
NEXT 100 [64, 73]	¹³⁶ Xe	91	7.8	[-1.3, 2.4]	88	37	26.5	$4.4 \cdot 10^{-2}$	$5.3 \cdot 10^{25}$	82-189	n/a	1	20
NEXT 1.5k [74]	¹³⁶ Xe	1367	5.2	[-1.3, 2.4]	88	37	398	$2.9 \cdot 10^{-3}$	$7.9 \cdot 10^{26}$	21-49	n/a	1	300
PandaX-III 200 [65]	¹³⁶ Xe	180	31	[-2, 2]	100	35	60.2	$4.2 \cdot 10^{-2}$	$8.3 \cdot 10^{25}$	65-150	n/a	n/a	n/a
PandaX-III 1k [65]	¹³⁶ Xe	901	10	[-2, 2]	100	35	301	$1.4 \cdot 10^{-3}$	$9.0 \cdot 10^{26}$	20-46	n/a	n/a	n/a

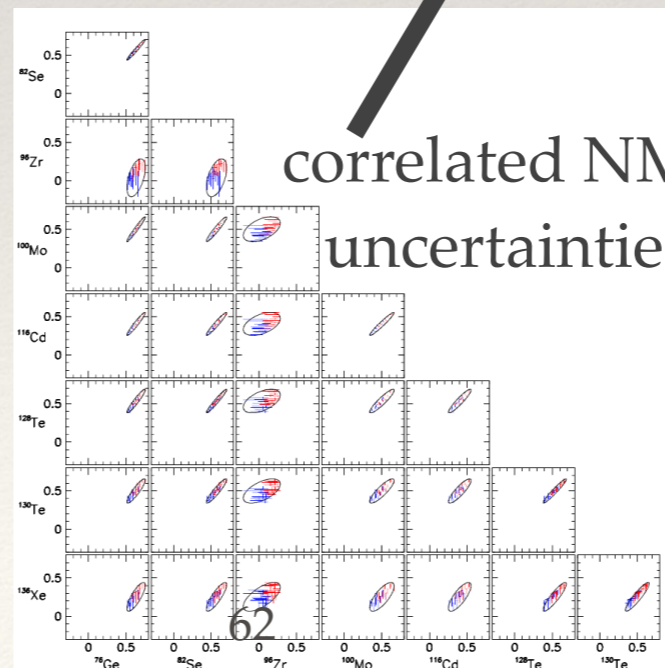
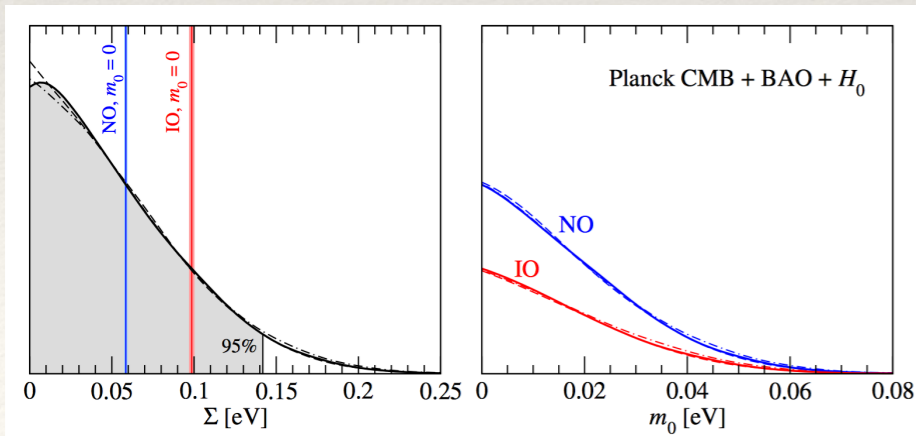
Expectations of lifetimes



Predicted Half-Lifetime for ^{76}Ge



Ge, WR, Zuber, 1707.07904



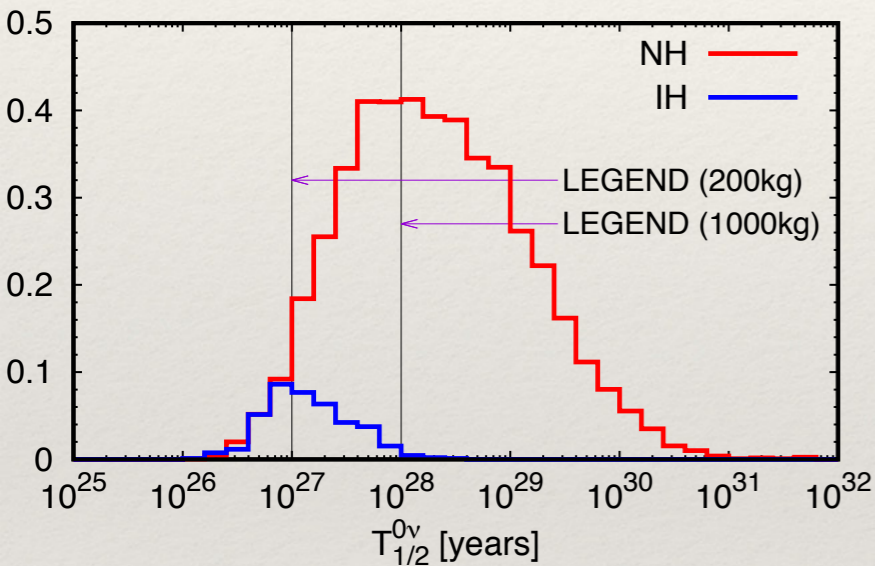
Expectations for half-lives

Standard

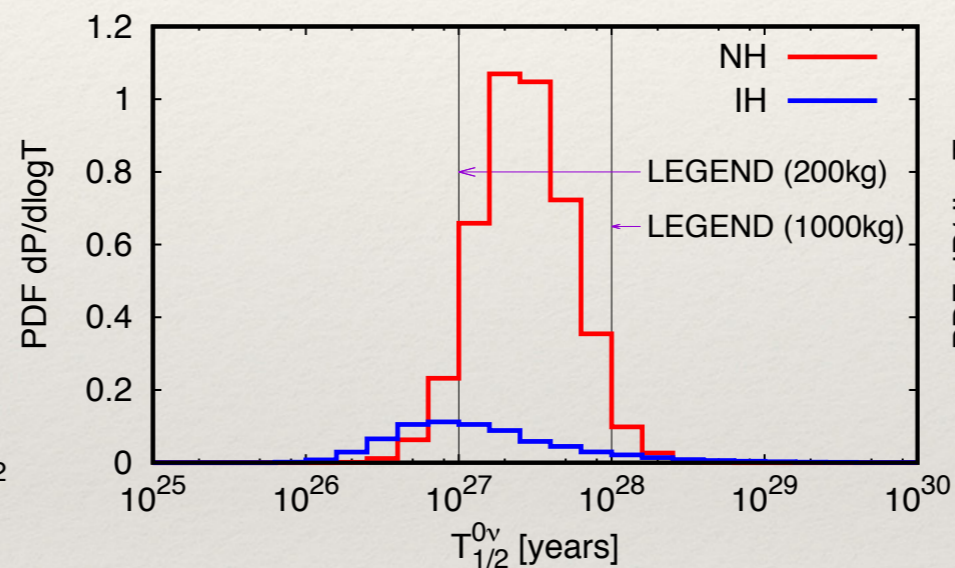
Sterile

Left-right

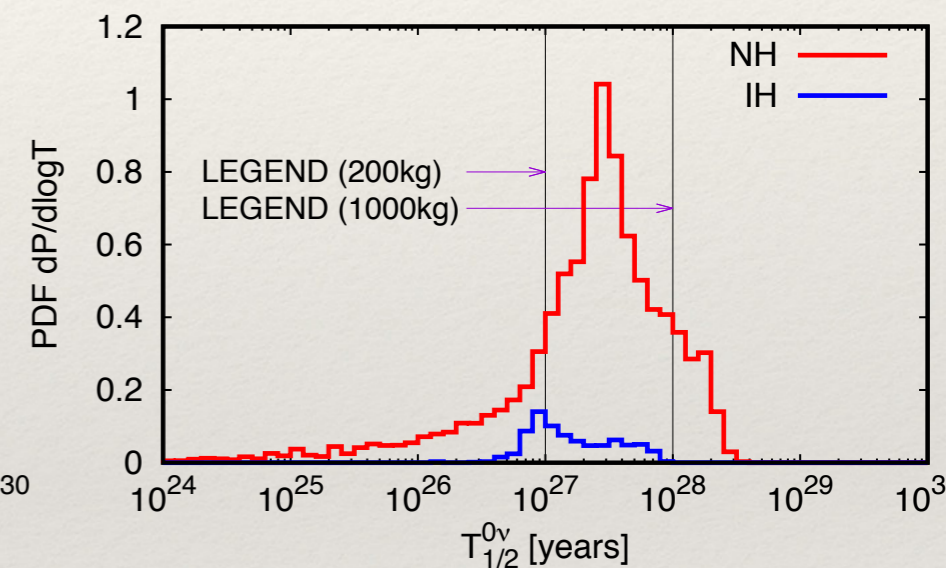
Predicted Half-Lifetime for ^{76}Ge



Predicted Half-Lifetime for ^{76}Ge



Predicted Half-Lifetime for ^{76}Ge [LRSM-typeII]



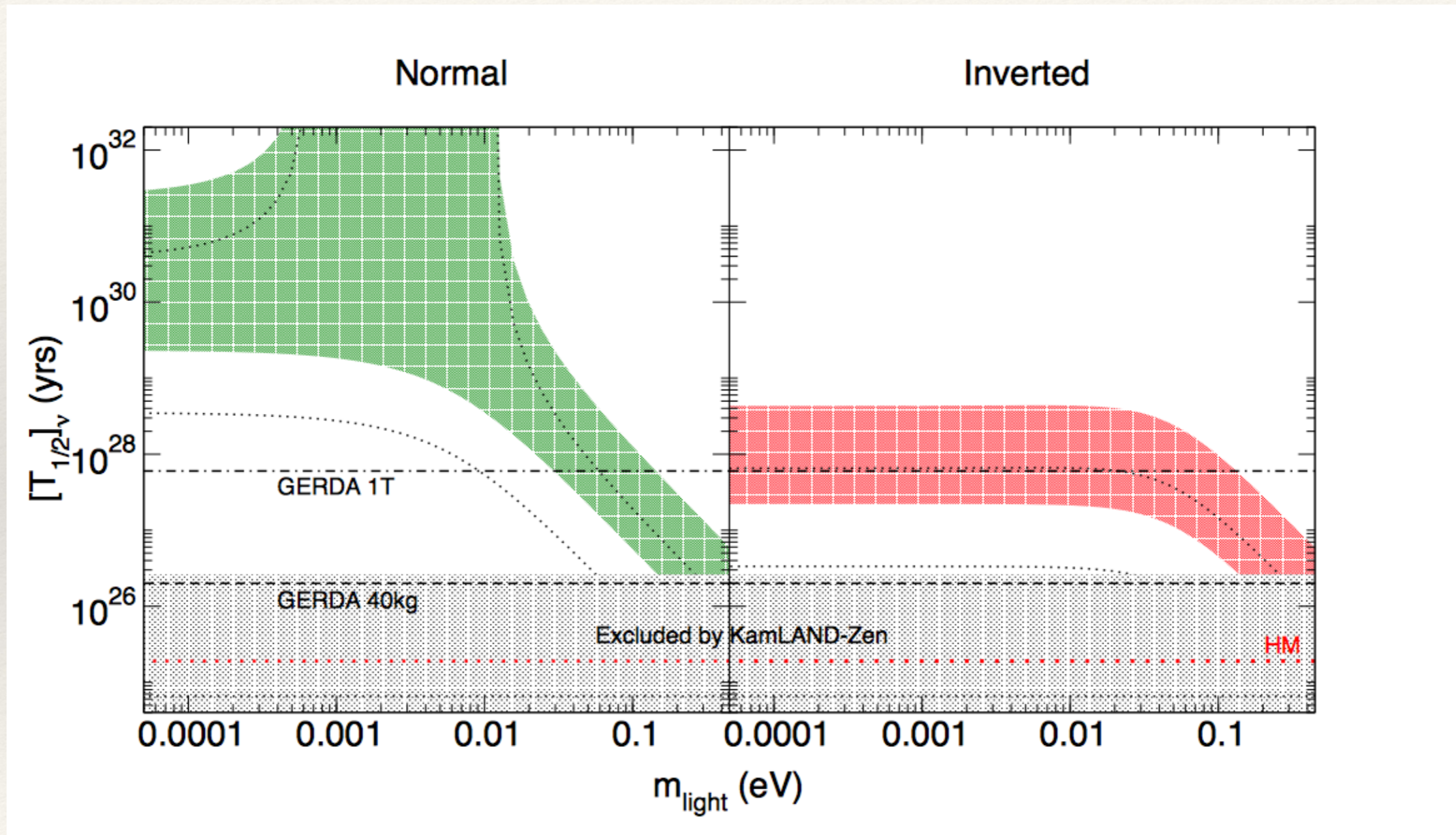
Ge, WR, Zuber, 1707.07904

*For standard scenario, see also Agostini et al, 1705.02996; Caldwell et al., 1705.01945;
Zhang, Zhou, 1508.05472; Benato, 1510.01089*

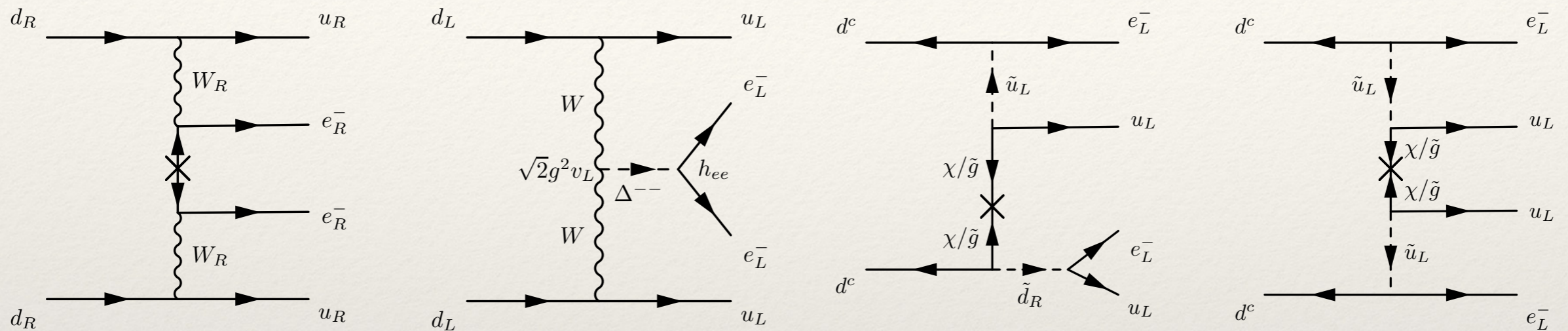
However, most alternative mechanisms unrelated to neutrino parameters...

...thus decoupled from cosmology (and direct experiments)!

The usual plot



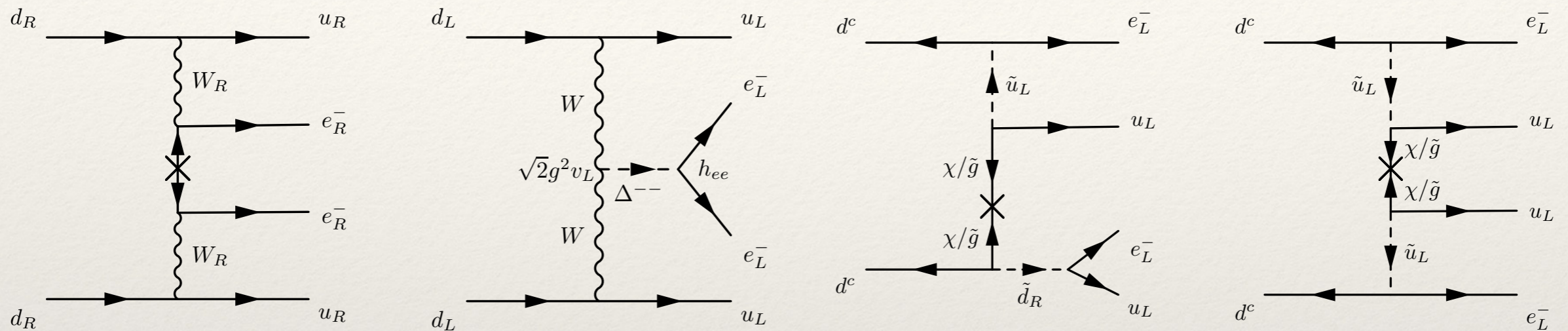
Non-Standard Interpretations



- ❖ typically decouples double beta decay from cosmology and KATRIN

$$\mathcal{A}_{\text{Standard}} = G_F^2 \frac{\langle m \rangle}{q^2} \quad \text{versus} \quad \mathcal{A}_{\text{Non-Standard}} = \frac{c}{M_X^5}$$

Non-Standard Interpretations

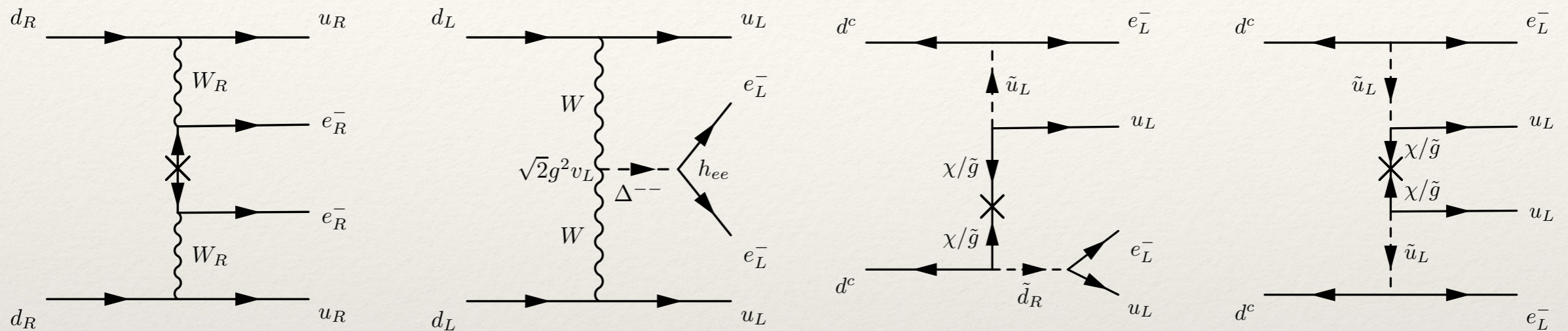


- ❖ typically decouples double beta decay from cosmology and KATRIN

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Therefore:
 $T(\text{eV}) = T(\text{TeV})$

Non-Standard Interpretations



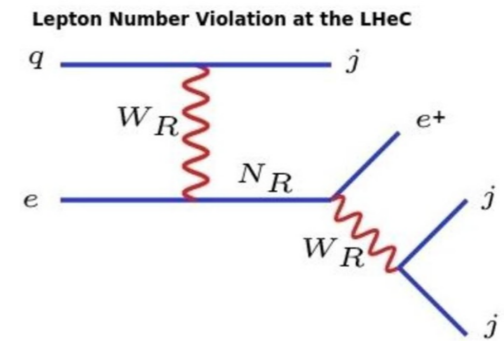
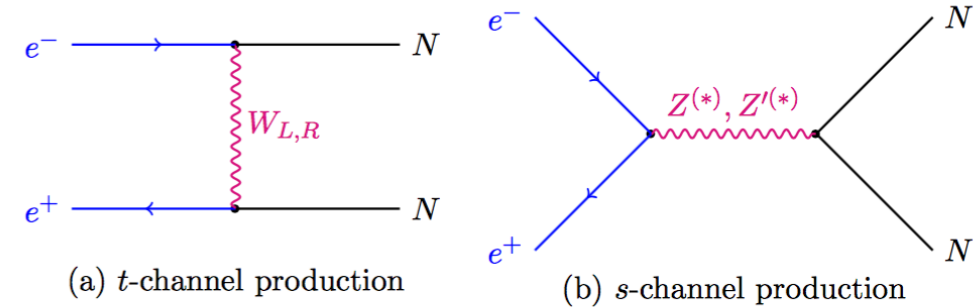
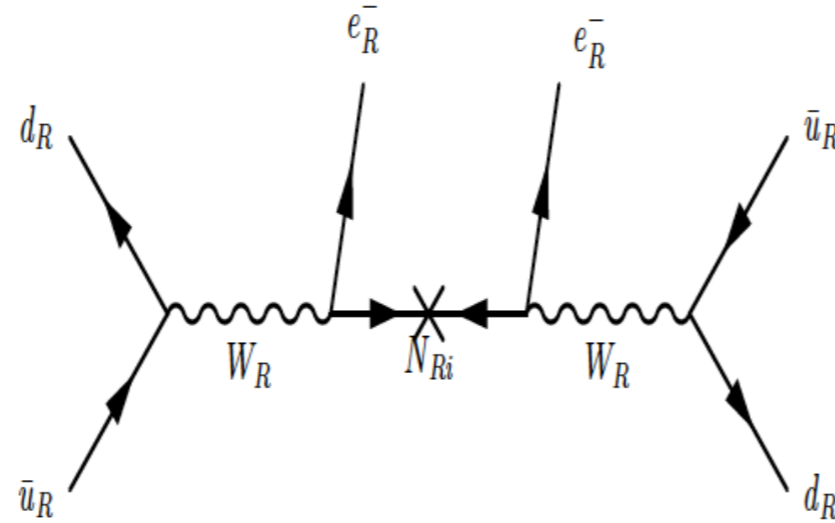
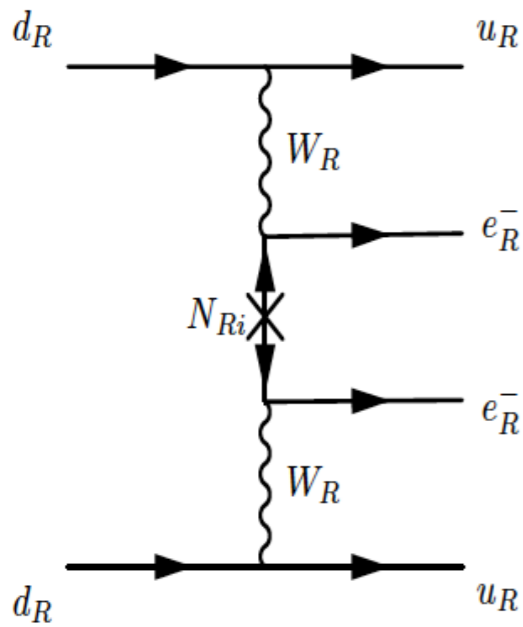
- ❖ typically decouples double beta decay from cosmology and KATRIN

$$\mathcal{A}_{\text{Standard}} = G_F^2 \frac{\langle m \rangle}{q^2} \quad \text{versus} \quad \mathcal{A}_{\text{Non-Standard}} = \frac{c}{M_X^5}$$

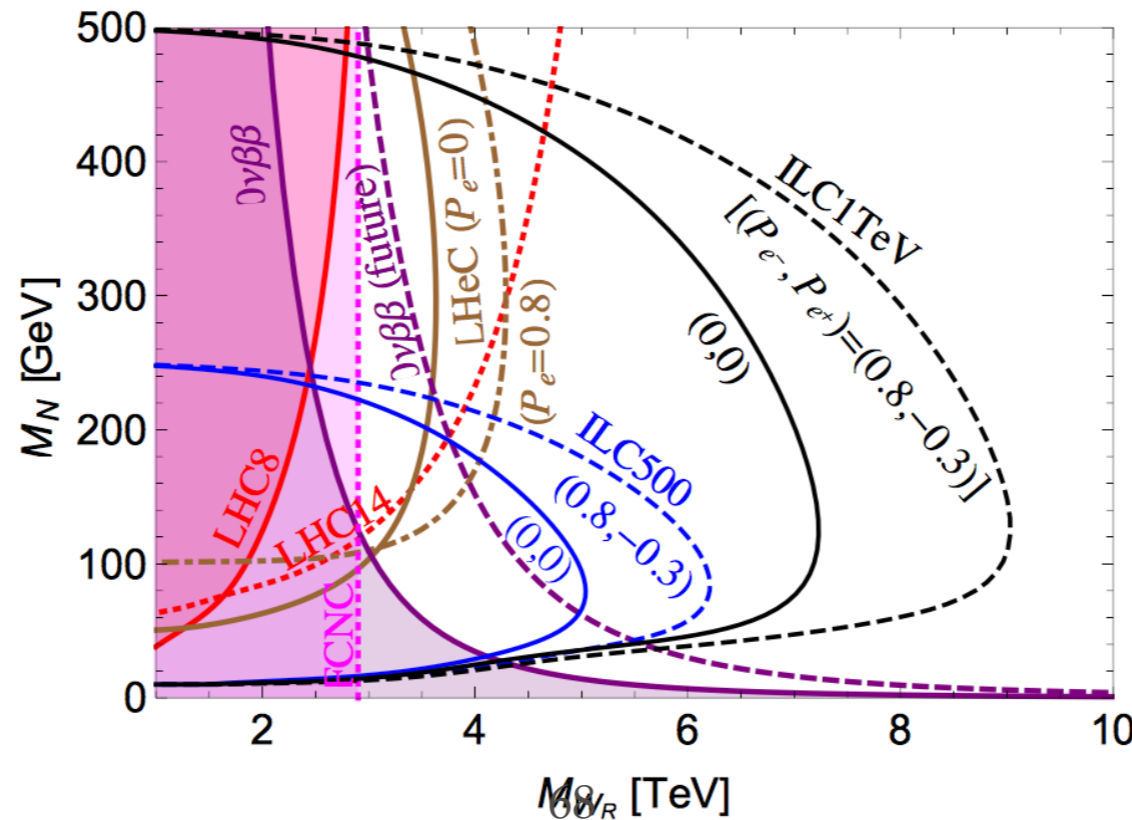
Therefore:
 $T(\text{eV}) = T(\text{TeV})$

⇒ Tests with LHC, LFV, etc.

LHC and Double Beta Decay



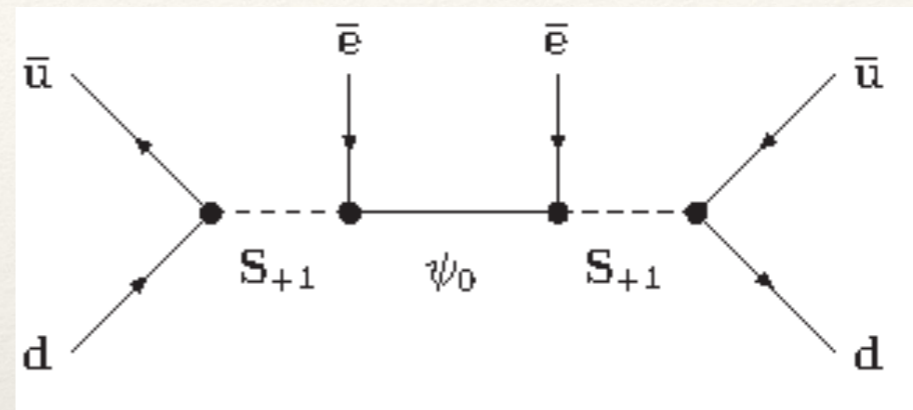
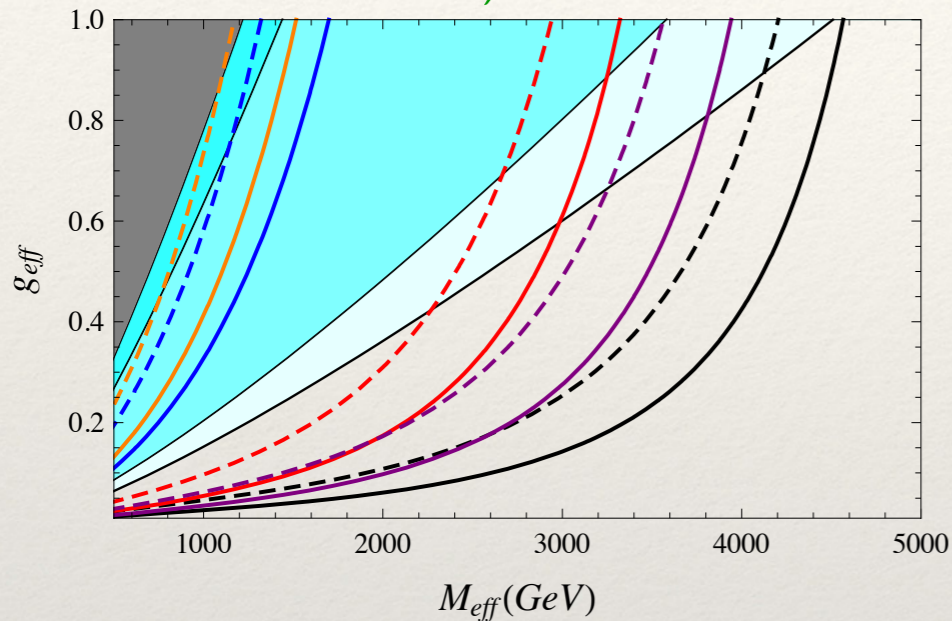
Biwal, Dev, 1701.08751



polarization at
LHeC and ILC

Complementarity of LHC and $0\nu\beta\beta$

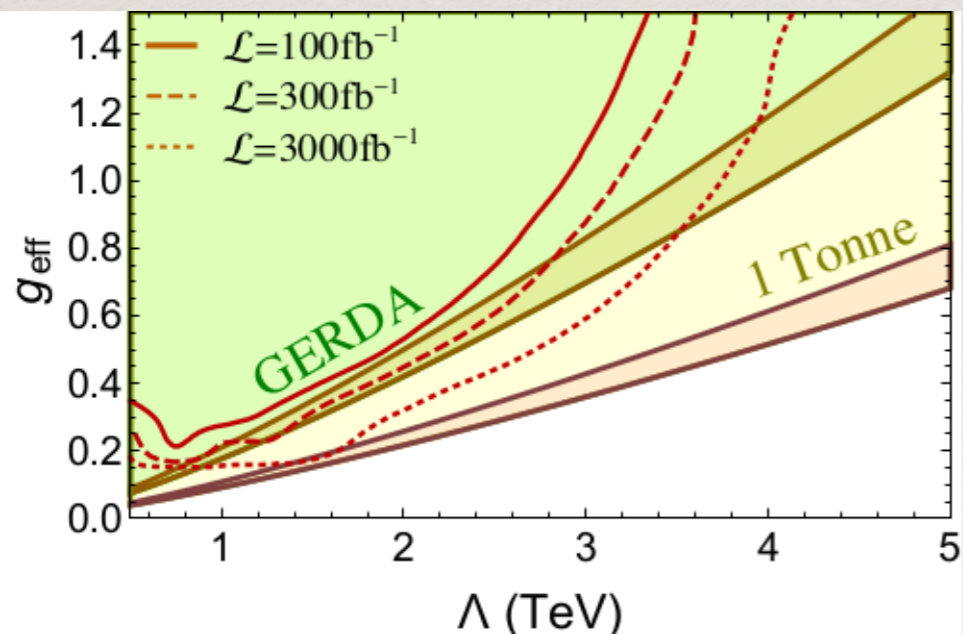
Hirsch et al., 1511.03945



$S \sim (1,2)$

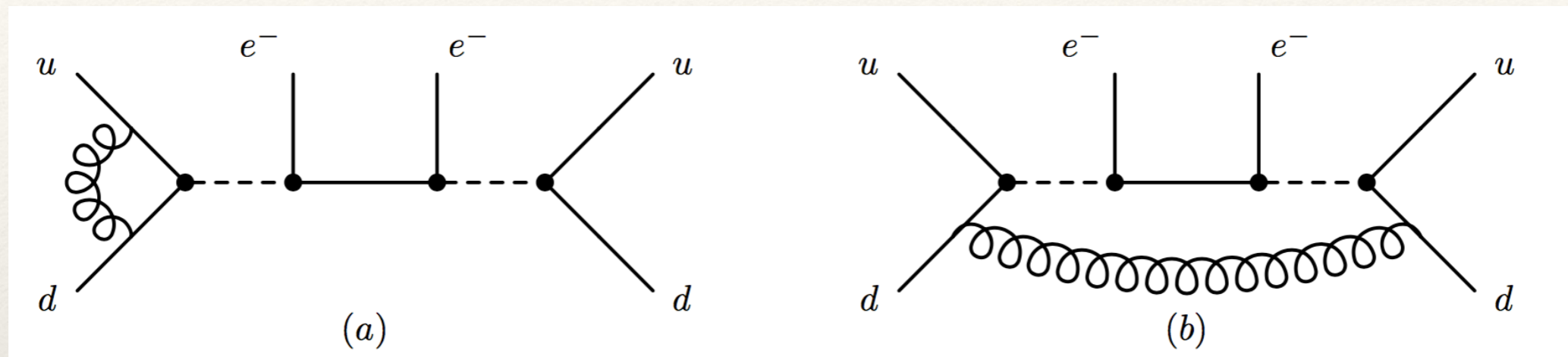
$\psi \sim (1,0)$

- ❖ LHC needs $M_S > M_\psi$
- ❖ LHC has low sensitivity for small M_ψ
- ❖ include jet-fake rate, charge mis-ID, QCD corrections in $0\nu\beta\beta$, etc.
- ❖ \Rightarrow complementary



Ramsey-Musolf et al., 1508.04444

QCD Corrections



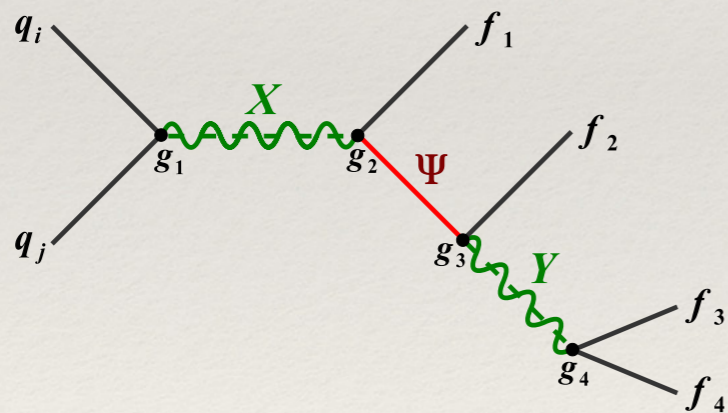
- ❖ naive size $(\alpha_s/4\pi) \ln (M_W/100 \text{ MeV})^2 \approx 10\%$, true for standard diagram
- ❖ creates in non $(V-A) \otimes (V-A)$ short-range mechanisms color non-singlets, Fierzing to singlets gives different operators with vastly different NMEs
- ❖ \Rightarrow can give effect exceeding NME uncertainty...

Mahajan, PRL 112; Gonzalez, Kovalenko, Hirsch, PRD 93;

Peng, Ramsey-Musolf, Winslow, PRD 93

TeV-scale LNV and Baryogenesis

- ❖ Example TeV-scale W_R : leads to washout in early Universe via $e_R e_R \leftrightarrow W_R W_R$ and $e_R W_R \leftrightarrow W_R e_R$; processes stay long in equilibrium (*Frere, Hambye, Vertongen; Bhupal Dev, Mohapatra; Sarkar et al.*)
- ❖ more model-independent (*Deppisch, Harz, Hirsch*):



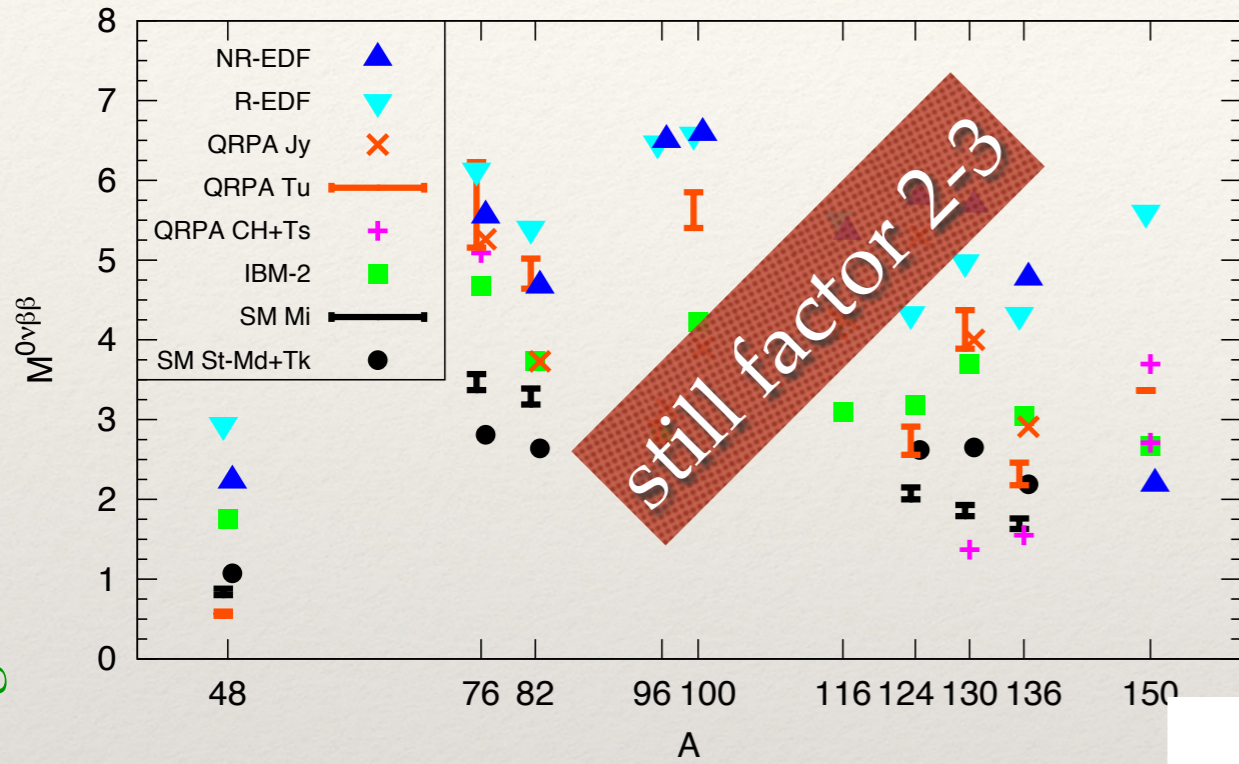
wash-out:

$$\log_{10} \frac{\Gamma_W(qq \rightarrow \ell^+ \ell^+ qq)}{H} \gtrsim 6.9 + 0.6 \left(\frac{M_X}{\text{TeV}} - 1 \right) + \log_{10} \frac{\sigma_{\text{LHC}}}{\text{fb}}$$

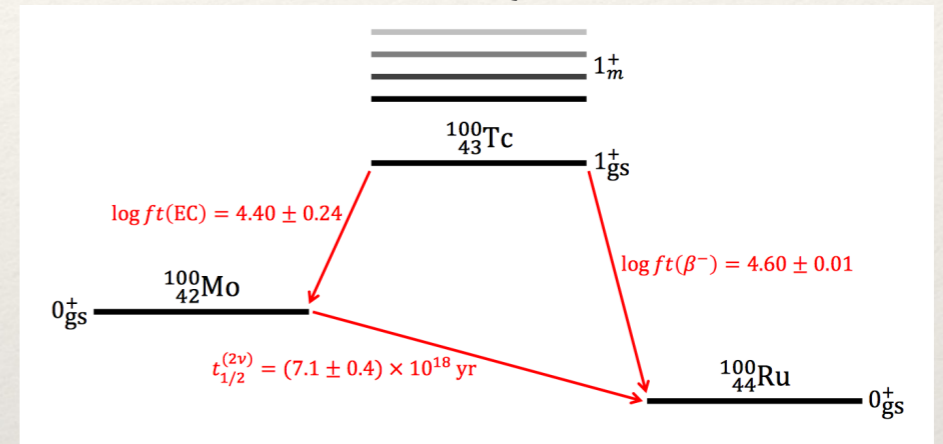
would need electroweak, resonant, ARS, post-sphaleron baryogenesis

Nuclear Matrix Elements

Engel, Menendez, 1610.06548

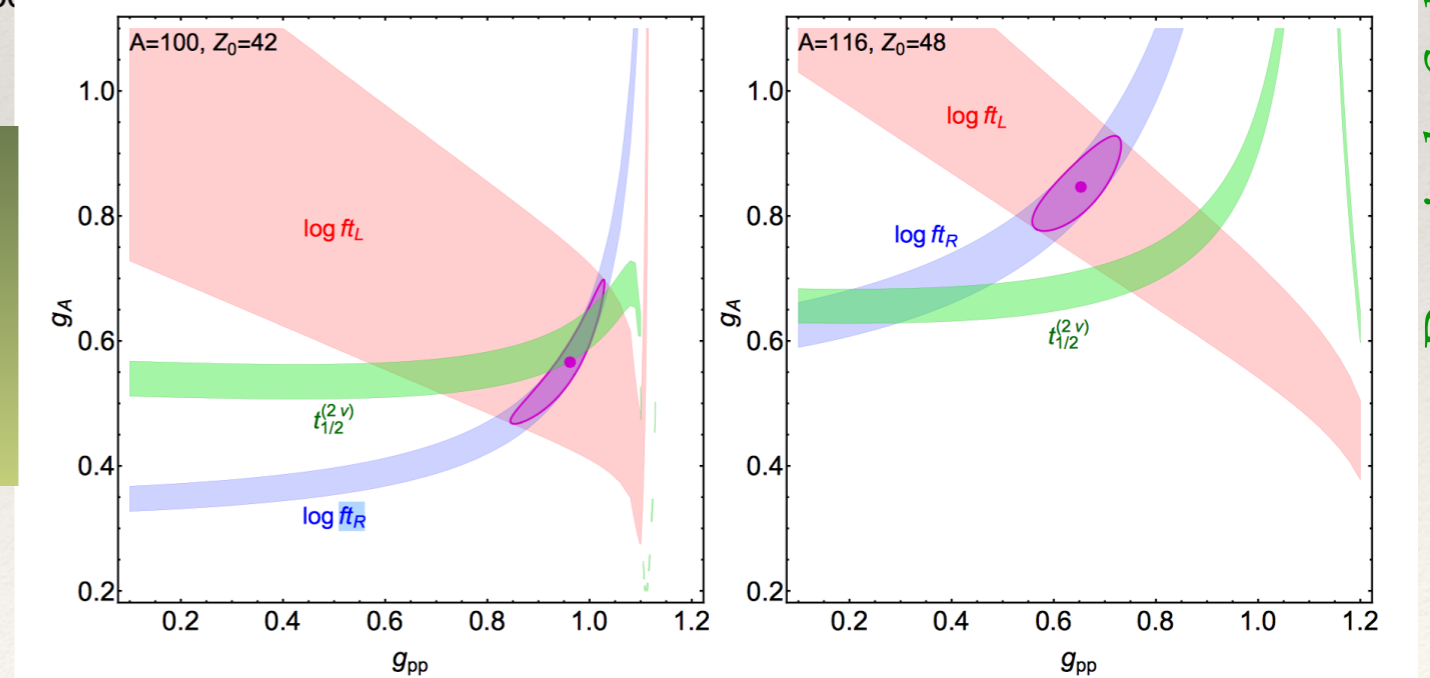


How good are the models?
Example isobaric triplets
within QRPA



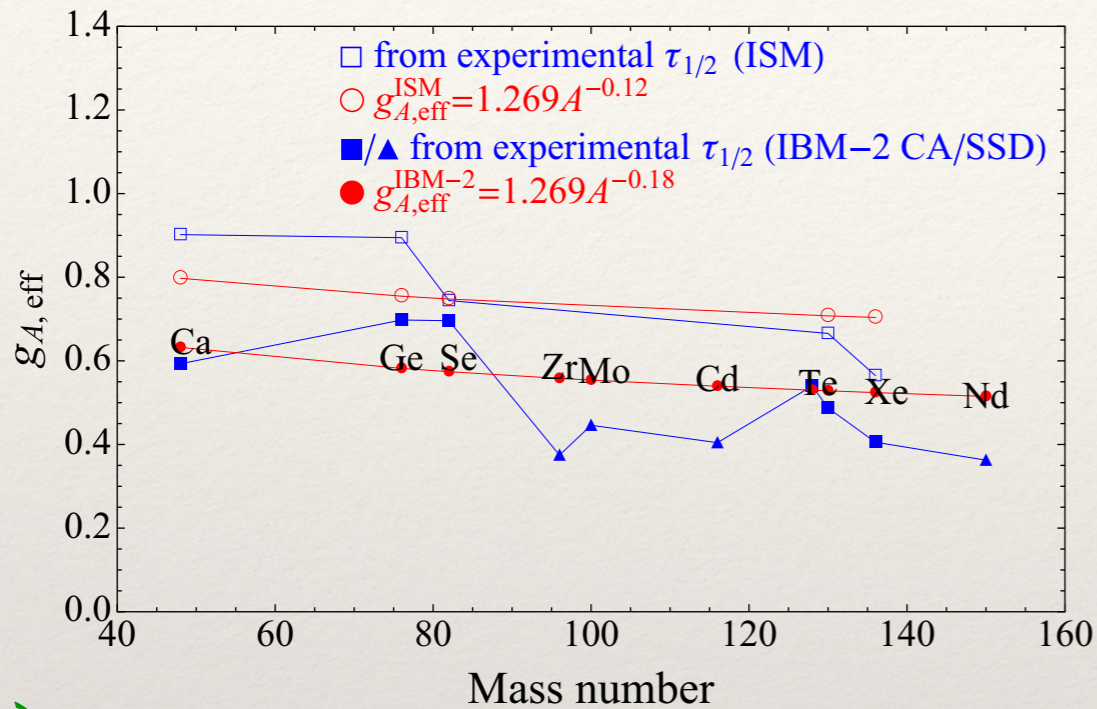
Deppisch, Suhonen, 1606.02908

⇒ Need as much experimental input (e.g. charge exchange) as possible...



Nuclear Matrix Elements

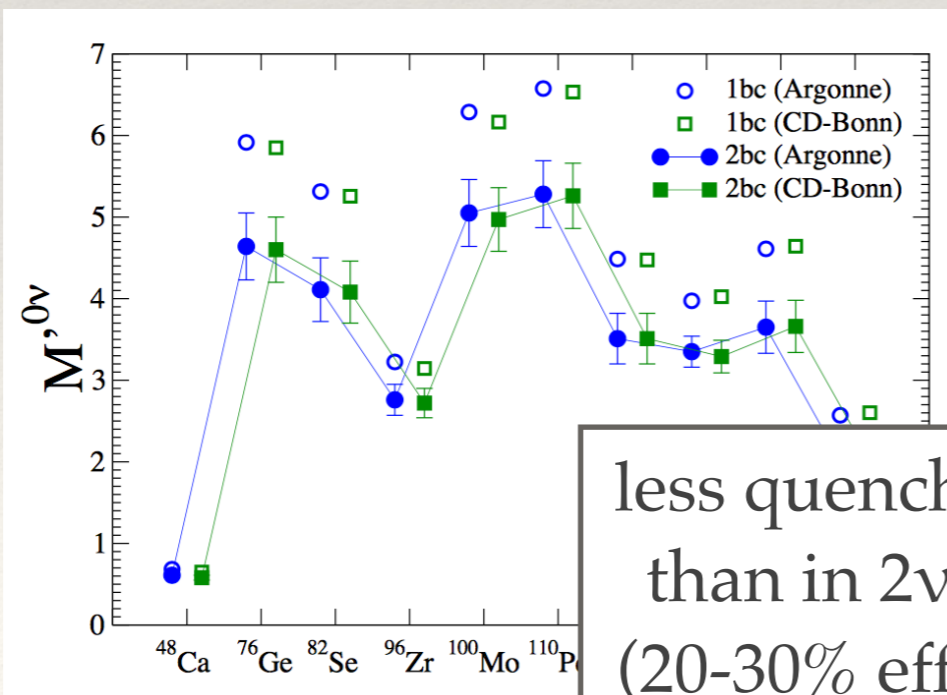
Iachello et al., 1506.08530



QUENCHING??

$$T_{\frac{1}{2}}^{0\nu} \propto g_A^{-4}$$

- ❖ fact in β and $2\nu\beta\beta$
- ❖ truncation of model-space?
- ❖ also in $0\nu\beta\beta$??
 - $q = 10^2$ vs. 10^0 MeV?
 - higher multipolarities?
 - two-body currents?
 - muon capture?
 - SM vs. QRPA



less quenching than in $2\nu\beta\beta$ (20-30% effect)

Menendez, Gazit, Schwenk, 1103.3622; Engel, Simkovic, Vogel, 1403.7860

Why look for Lepton Number Violation?

- ❖ L and B accidentally conserved in SM
- ❖ $\mathcal{L} = \mathcal{L}_{\text{SM}} + 1/\Lambda \mathcal{L}_5 + 1/\Lambda^2 \mathcal{L}_6 + \dots$, with $\mathcal{L}_5 = L^c \phi \phi L \rightarrow m_\nu \nu_{L^c} \nu_L$
- ❖ Baryogenesis: B is violated
- ❖ B, L often connected in GUTs
- ❖ GUTs have seesaw and Majoranas
- ❖ (B and L non-perturbatively violated by 3 units in SM...)

Lepton Number as important as Baryon Number

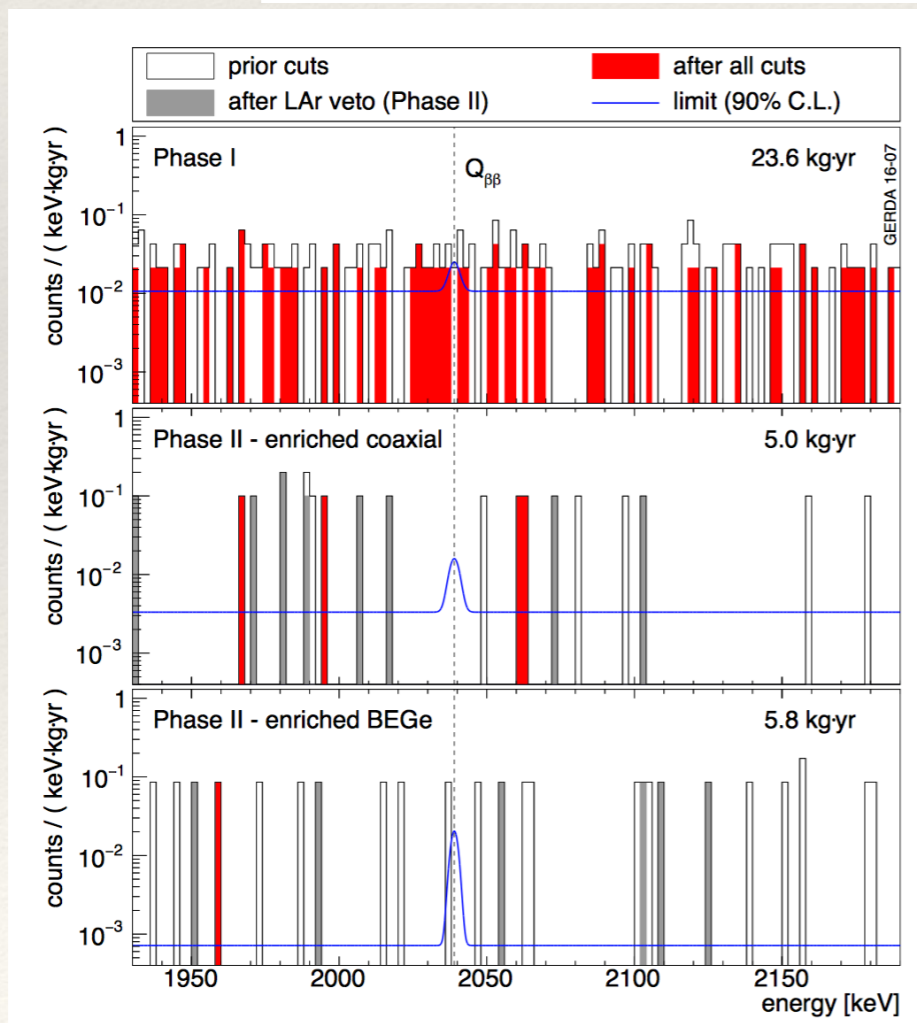
Neutrinoless Double Beta Decay

$$(T_{1/2}^{0\nu})^{-1} \propto \begin{cases} a M \varepsilon t & \text{without background} \\ a \varepsilon \sqrt{\frac{M t}{B \Delta E}} & \text{with background} \end{cases}$$

first background free result

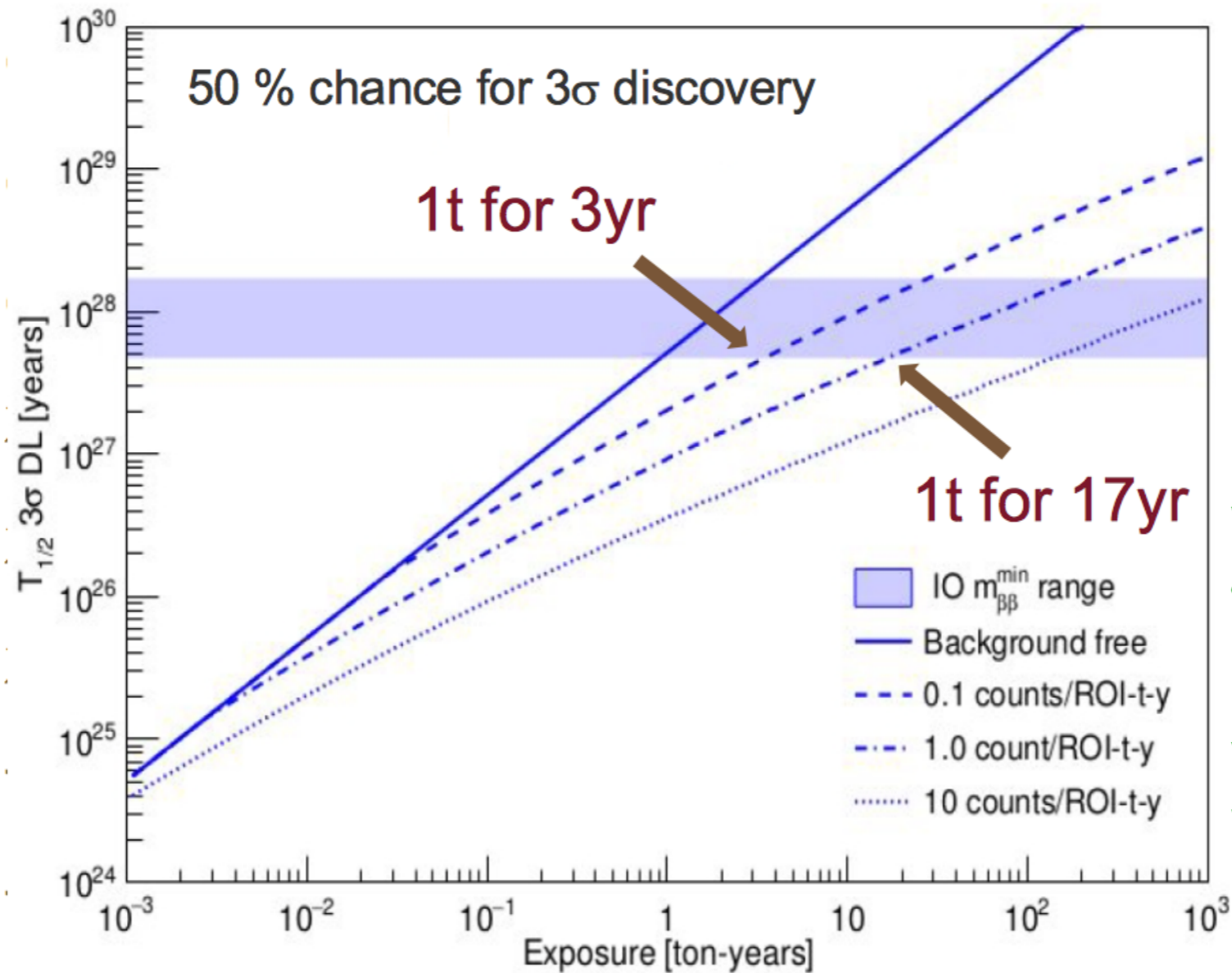
current limits: $T_{1/2} \gtrsim 10^{26}$ years
with exposure of about 100 kg · years

GERDA, 1703.00570, Nature



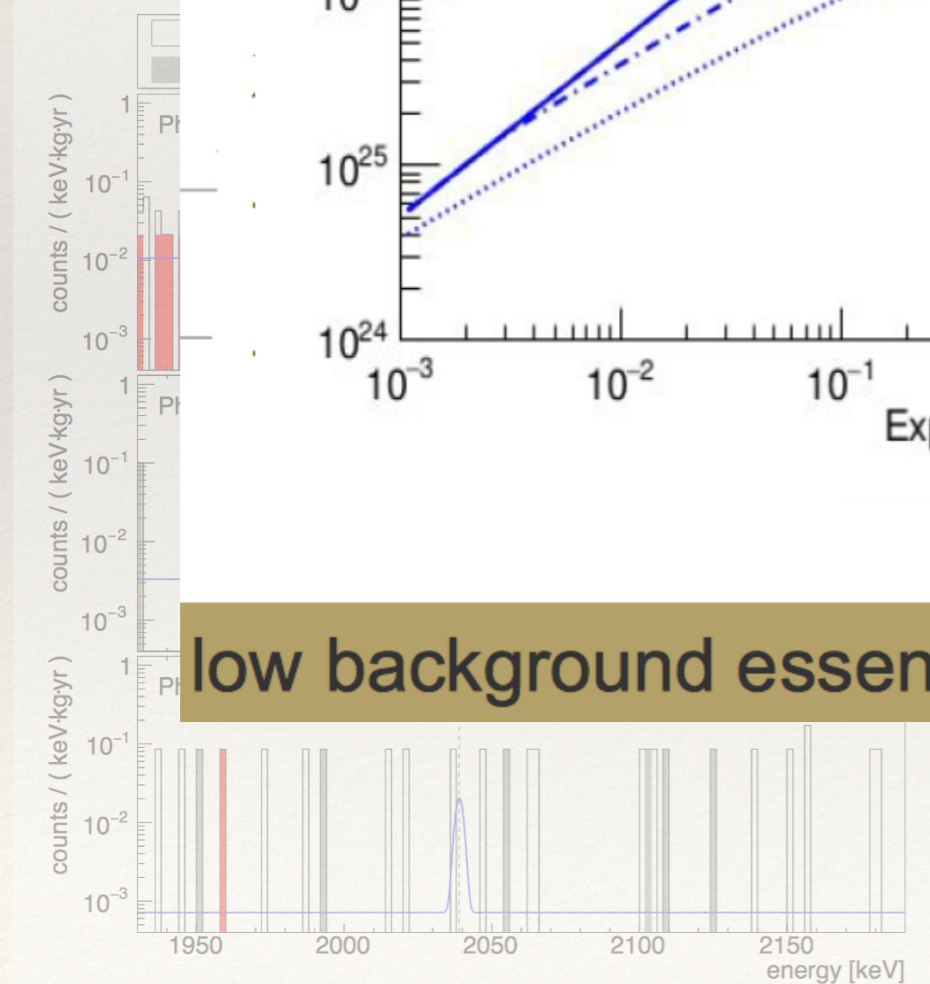
Neutrino

Decay



Plot by Josef Jochum

low background essential for discovery potential



ground
und

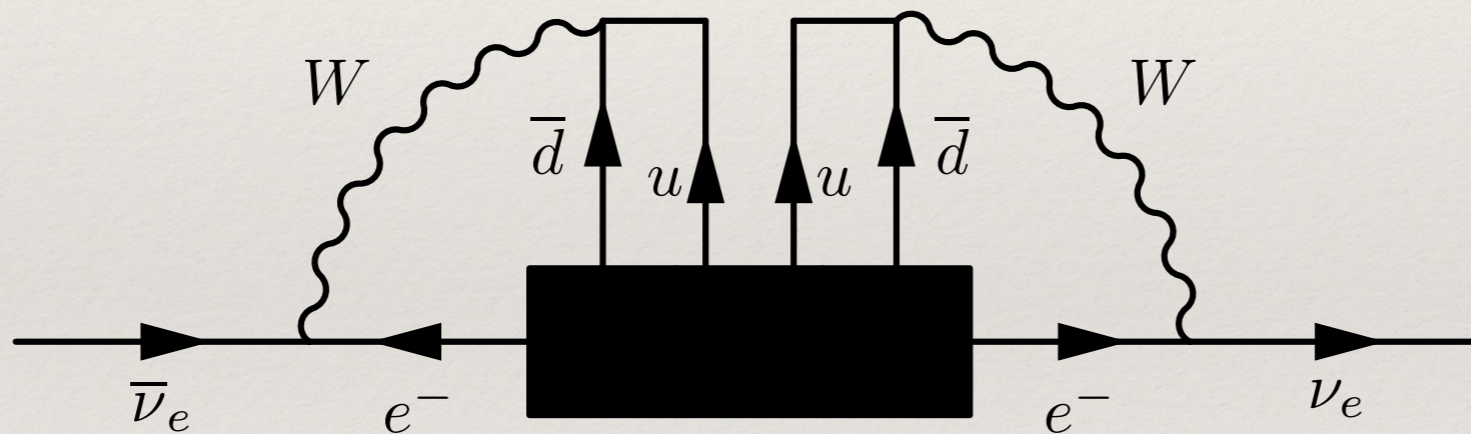
free result

10²⁶ years
100 kg · years

GERDA, 1703.00570, Nature

Black Box Theorem

- ❖ Whatever the mechanism, observation of $0\nu\beta\beta$ implies Majorana neutrinos (*Schechter-Valle, '82*)

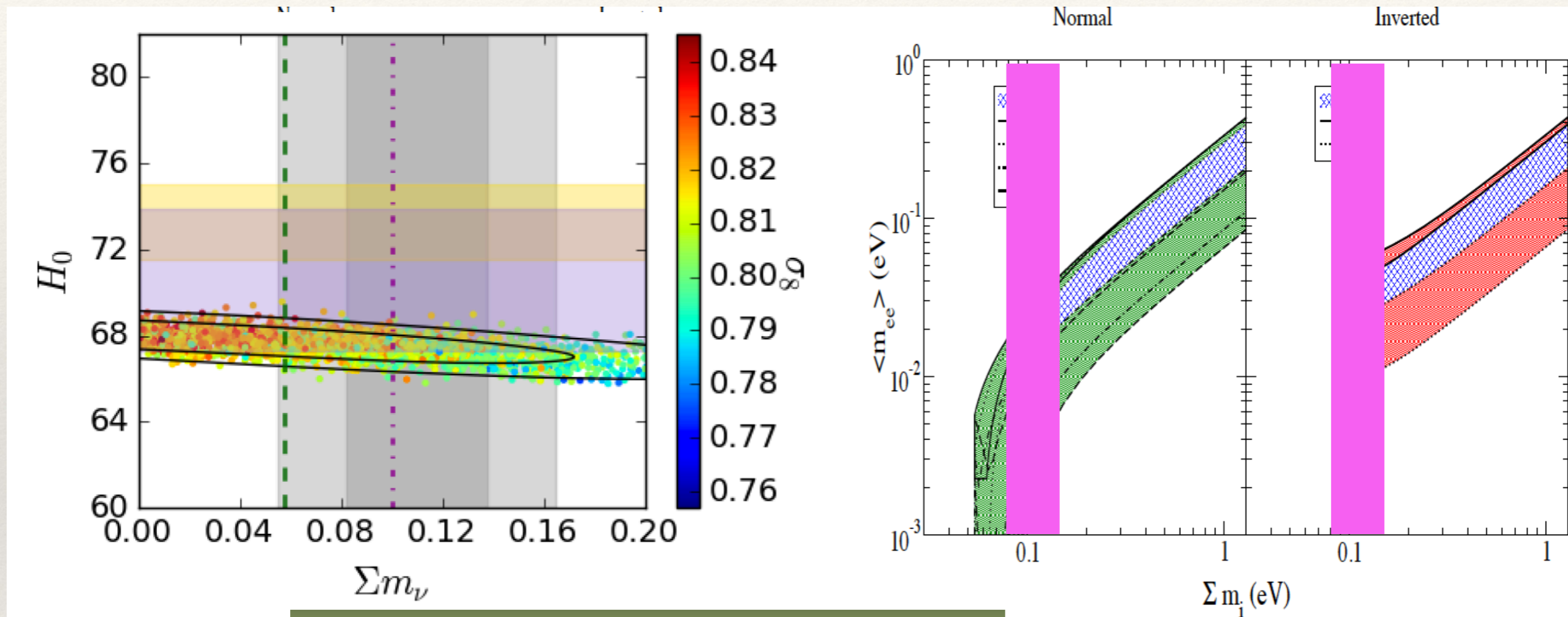


- ❖ is 4-loop diagram \Rightarrow tiny mass (*Dürr, Lindner, Merle, 1105.0901*)

Non-Standard Interpretations

mechanism	physics parameter	current limit	test
light neutrino exchange	$ U_{ei}^2 m_i $	0.2 eV	oscillations, cosmology, neutrino mass
heavy neutrino exchange	$\left \frac{S_{ei}^2}{M_i} \right $	$2 \times 10^{-8} \text{ GeV}^{-1}$	LFV, collider
heavy neutrino and RHC	$\left \frac{V_{ei}^2}{M_i M_{WR}^4} \right $	$4 \times 10^{-16} \text{ GeV}^{-5}$	flavor, collider
Higgs triplet and RHC	$\left \frac{(M_R)_{ee}}{m_{\Delta_R}^2 M_{WR}^4} \right $	$10^{-15} \text{ GeV}^{-1}$	flavor, collider e^- distribution
λ -mechanism with RHC	$\left \frac{U_{ei} \tilde{S}_{ei}}{M_{WR}^2} \right $	$1.4 \times 10^{-10} \text{ GeV}^{-2}$	flavor, collider, e^- distribution
η -mechanism with RHC	$\tan \zeta \left U_{ei} \tilde{S}_{ei} \right $	6×10^{-9}	flavor, collider, e^- distribution
short-range \not{R}	$\frac{ \lambda'_{111} ^2}{\Lambda_{\text{SUSY}}^5}$ $\Lambda_{\text{SUSY}} = f(m_{\tilde{g}}, m_{\tilde{u}_L}, m_{\tilde{d}_R}, m_{\chi_i})$	$7 \times 10^{-18} \text{ GeV}^{-5}$	collider, flavor
long-range \not{R}	$\left \sin 2\theta^b \lambda'_{131} \lambda'_{113} \left(\frac{1}{m_{b_1}^2} - \frac{1}{m_{b_2}^2} \right) \right $ $\sim \frac{G_F}{q} m_b \frac{ \lambda'_{131} \lambda'_{113} }{\Lambda_{\text{SUSY}}^3}$	$2 \times 10^{-13} \text{ GeV}^{-2}$ $1 \times 10^{-14} \text{ GeV}^{-3}$	flavor, collider
Majorons	$ \langle g_\chi \rangle $ or $ \langle g_\chi \rangle ^2$	$10^{-4} \dots 1$	spectrum, cosmology

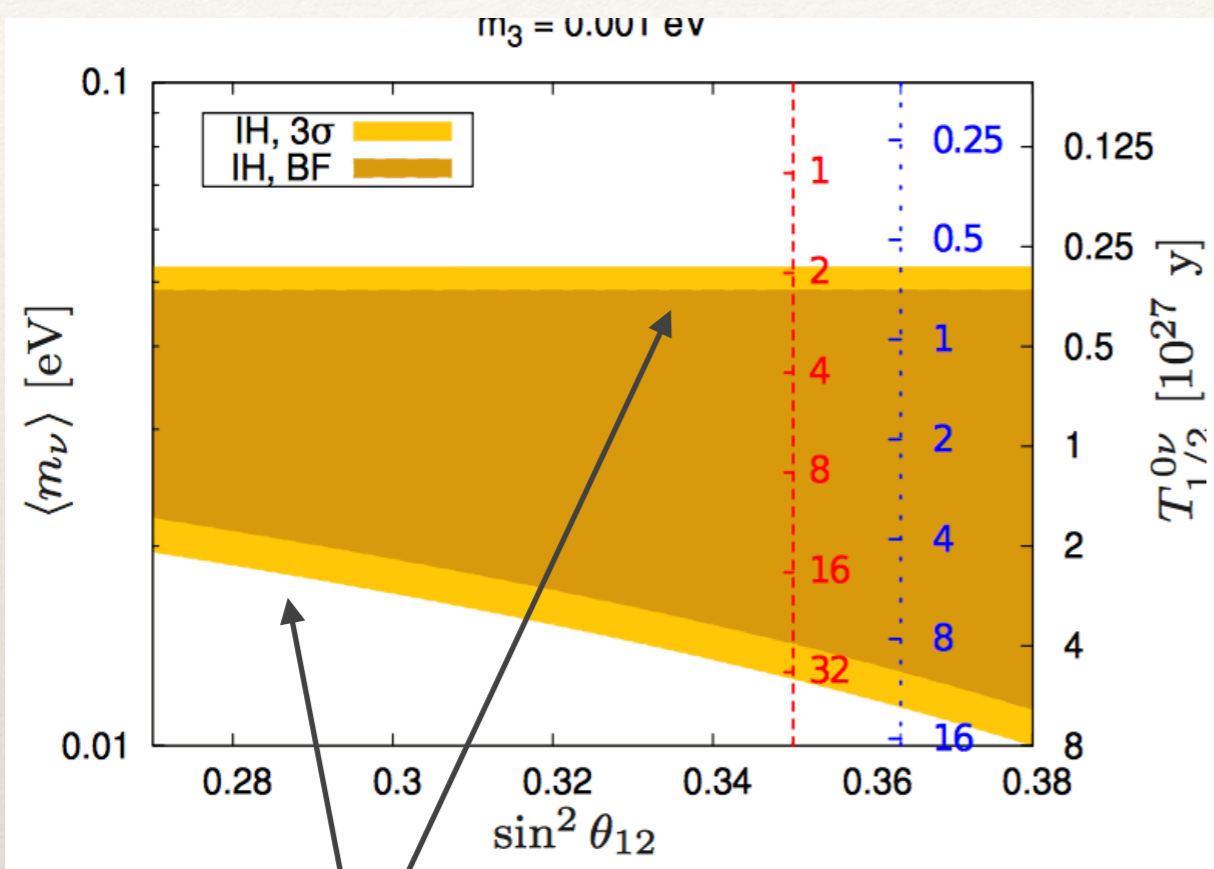
Neutrino Mass Observables



0.11 ± 0.03 eV from 1711.05210

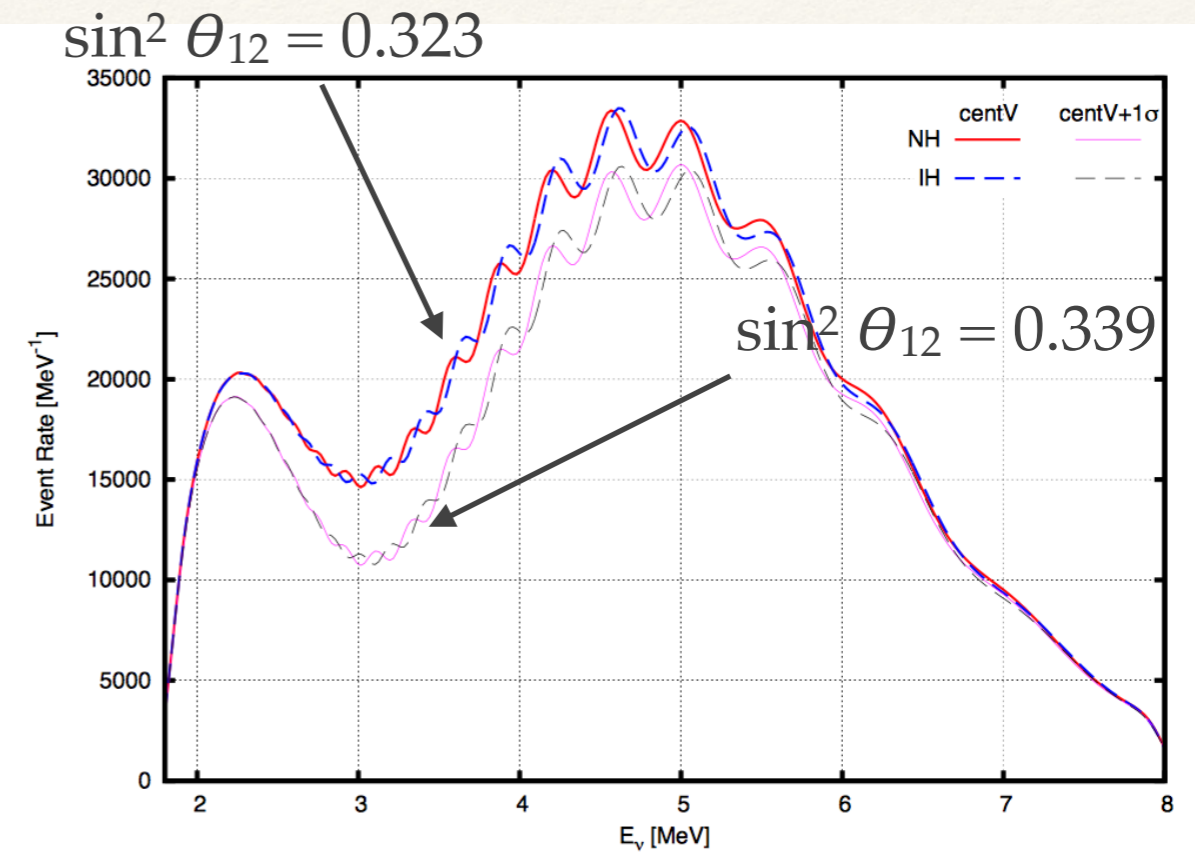
large effect of ν -mass in clustering length of galaxy clusters;
 much larger effect than on power spectrum;
 σ_8 larger locally larger than CMB-value;
 (H_0 still unresolved)

Connections to future Oscillation Experiments



Nature gives us two scales

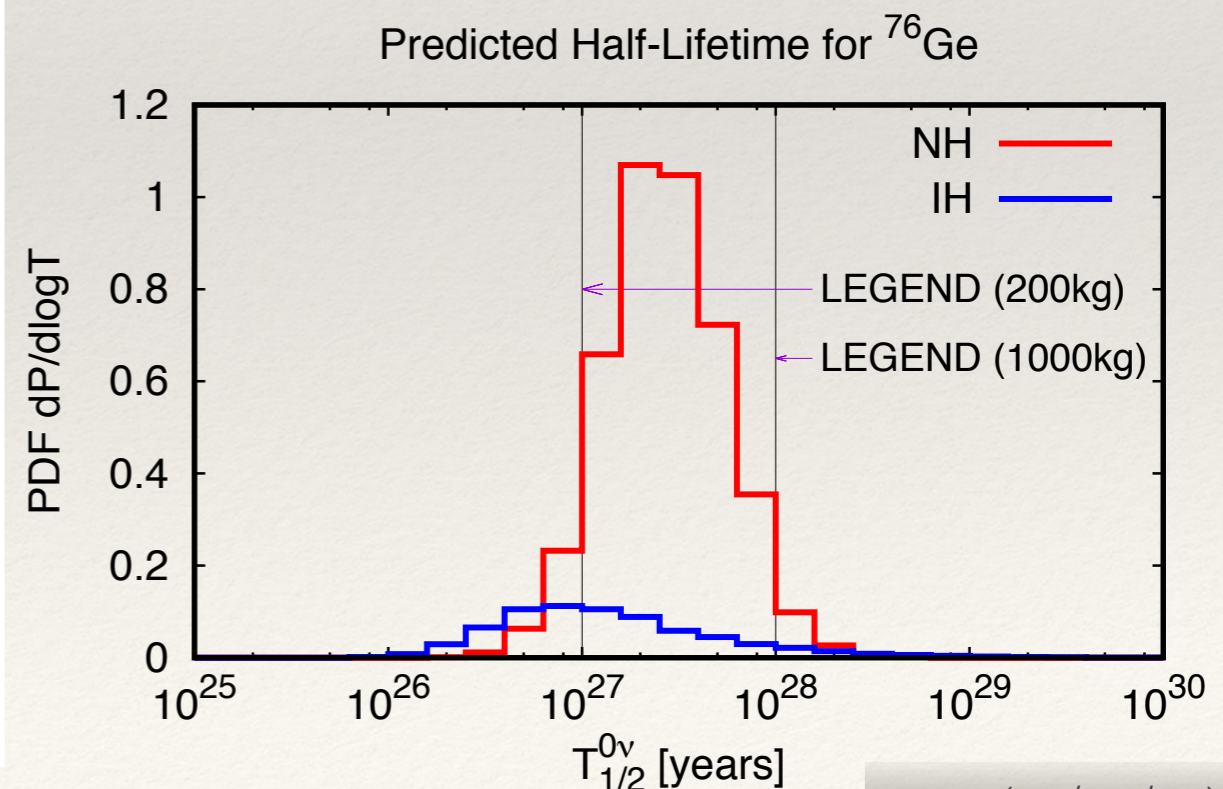
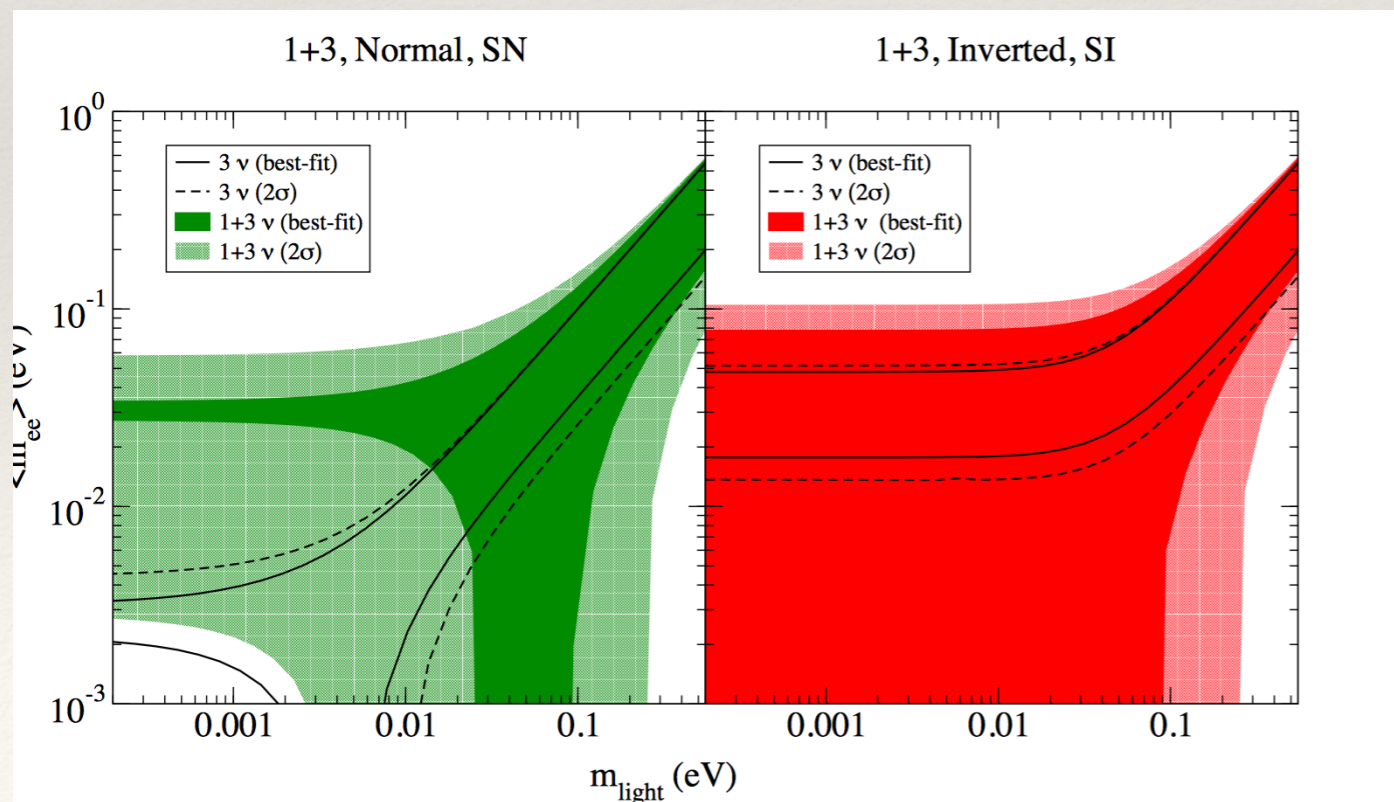
$$\begin{aligned}
 \langle m \rangle_{\text{IH}}^{\text{min}} &\propto \cos 2\theta_{12} \\
 &= 1 - 2 \sin^2 \theta_{12}
 \end{aligned}$$



JUNO fixes θ_{12} and removes uncertainty in value of minimal m_{ee} in IH

Sterile Neutrinos

- ❖ are there sterile states (LSND / reactor / etc.) with mass $\Delta m^2 \approx eV^2$ and mixing $U_{e4} \approx 0.1$?
- ❖ would make m_{ee} sum of 4 terms with sterile contribution $|U_{e4}|^2 \sqrt{\Delta m^2}$ that can cancel almost completely contribution of IH!
- ❖ usual pheno completely turned around!



Origin of Neutrino Mass

- ❖ Most straightforward possibility: add N_R and obtain Dirac mass:

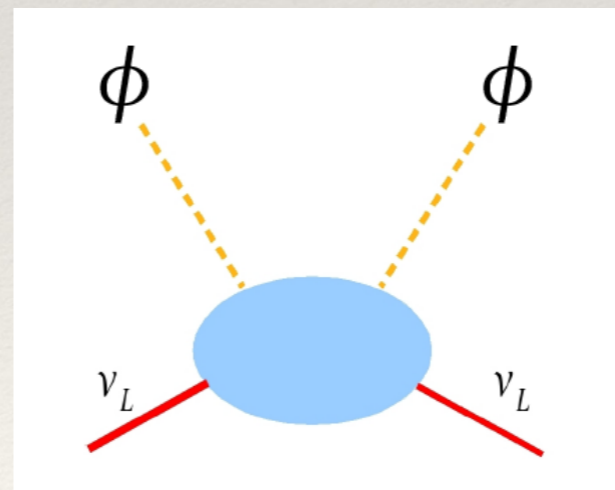
$$L \Phi N_R \rightarrow m_D \nu_L N_R$$

- ❖ Gauge invariance allows Majorana mass:

$$M_R N_R N_R$$

- ❖ in total Majorana mass for SM neutrinos:

$$m_\nu \nu_L^c \nu_L \text{ with } m_\nu = m_D^2 / M_R = m_D \varepsilon \text{ with } \varepsilon = m_D / M_R = m_{SM} / M_R$$



m_ν inverse
proportional to
scale of origin!

Origin of Neutrino Mass

- ❖ Most straightforward possibility:

$$m_D \nu_L N_R \rightarrow m_D \nu_L N_R$$

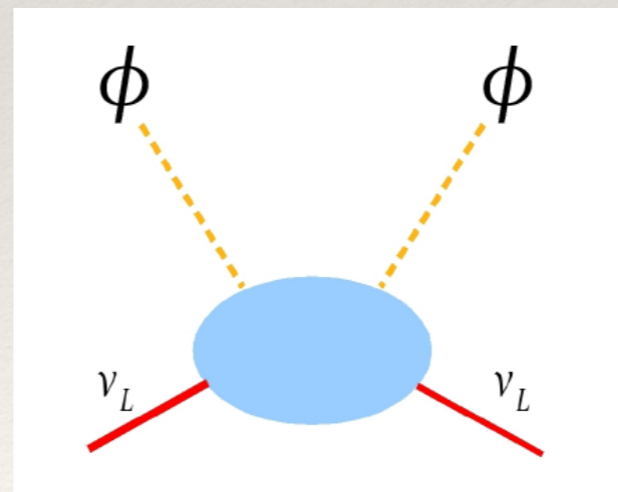
New representation of SM gauge group $N_R \sim (1,0)$ mass

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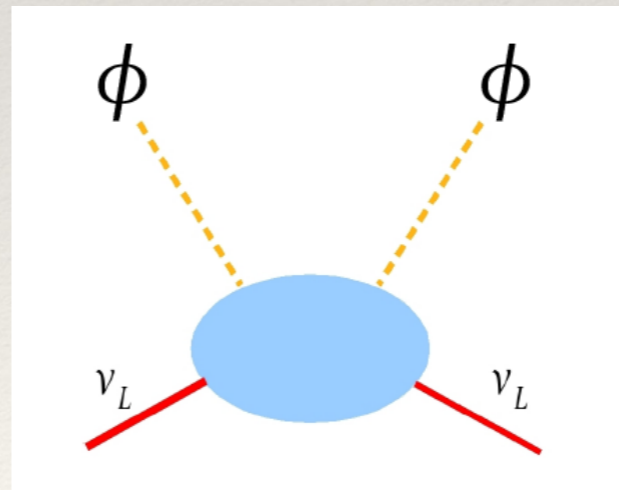
m_ν inverse proportional to scale of origin!

Origin of Neutrino Mass

- ❖ Most straightforward possibility:
 New representation of SM gauge group $N_R \sim (1,0)$ mass
 $\bar{\nu}_L N_R \rightarrow m_D \nu_L N_R$

- ❖ Gauge invariance allows Majorana mass
 New energy scale beyond SM

- ❖ in total Majorana mass for SM neutrinos:
 $m_\nu \nu_L^c \nu_L$ with $m_\nu = m_D^2 / M_R = m_D \epsilon$ with $\epsilon = m_D / M_R = m_{SM} / M_R$



m_ν inverse proportional to scale of origin!

Origin of Neutrino Mass

- ❖ Most straightforward possibility: **New representation of SM gauge group $N_R \sim (1,0)$ mass**

$$m_D \nu_L N_R$$

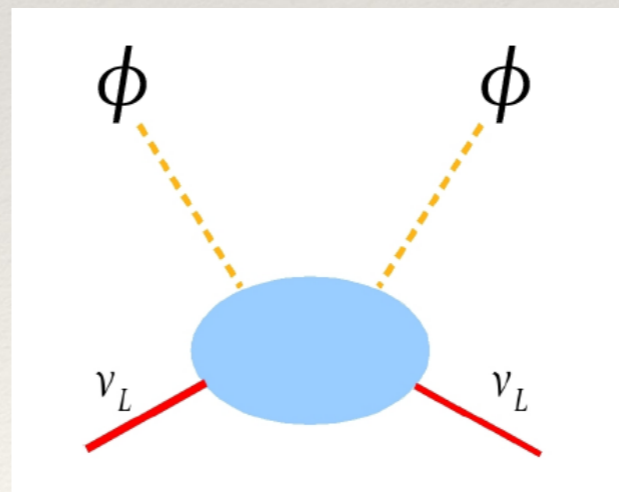
- ❖ Gauge invariance allows Majorana mass

New energy scale beyond SM

- ❖ in total Majorana mass for SM neutrinos:

$$m_\nu \nu_L^c \nu_L \text{ with } m_\nu \sim \frac{m_{SM}}{M_R}$$

New concept: lepton number violation



m_ν inverse proportional to scale of origin!

Origin of Neutrino Mass

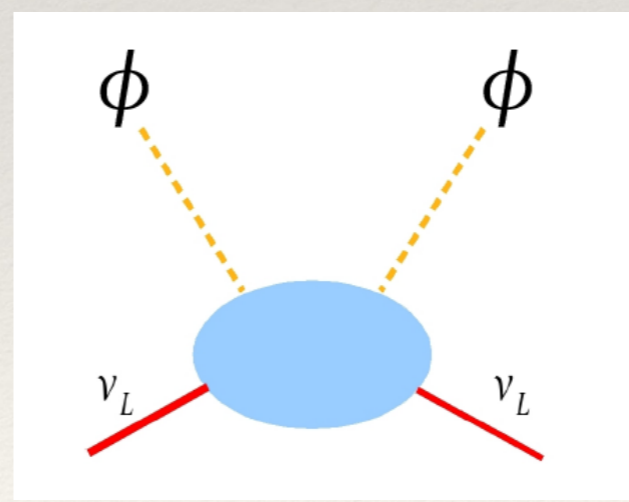
- ❖ Most straightforward possibility: add N_R and obtain Dirac mass

$$L \Phi N_R \rightarrow m_D \nu_L N_R$$
- ❖ Gauge invariance allows Majorana mass

$$M_R N_R N_R$$
- ❖ in total Majorana mass of neutrinos:

$$m_\nu \nu_L^c \nu_L \text{ with } m_\nu/M_R = m_D \epsilon \text{ with } \epsilon = m_D/M_R = m_{SM}/M_R$$

plus possible new interactions of N_R (B-L, LR Symmetry, etc.)



m_ν inverse proportional to scale of origin!

Type I Seesaw $m_\nu = m_D^2 / M_N \propto y^2 / M_N$

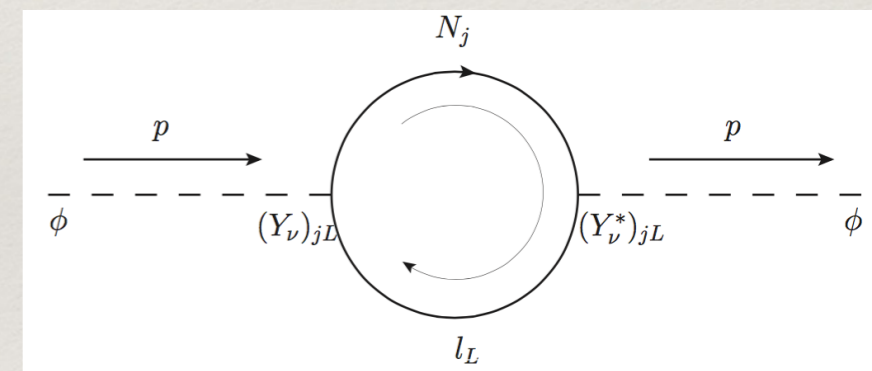
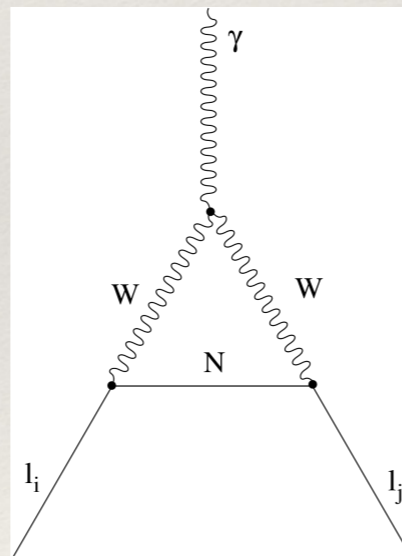
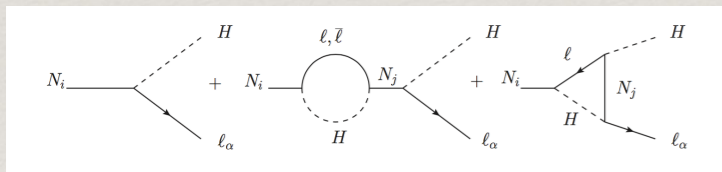
actually, does neither fix m_ν nor m_D nor M_N
needs to be tested or has phenomenology via „seesaw portal“:

Lepton-Higgs-Singlet Vertex: $y L \Phi N_R$

$$N_R \rightarrow L \Phi$$

$$L_\alpha \rightarrow N_R \Phi \rightarrow L_\beta$$

$$\Phi \rightarrow L N_R \rightarrow \Phi$$



Leptogenesis

$$Y_B \propto \text{Im}(y^2)$$

Talks by Deppisch,
Dev

Lepton Flavor Violation

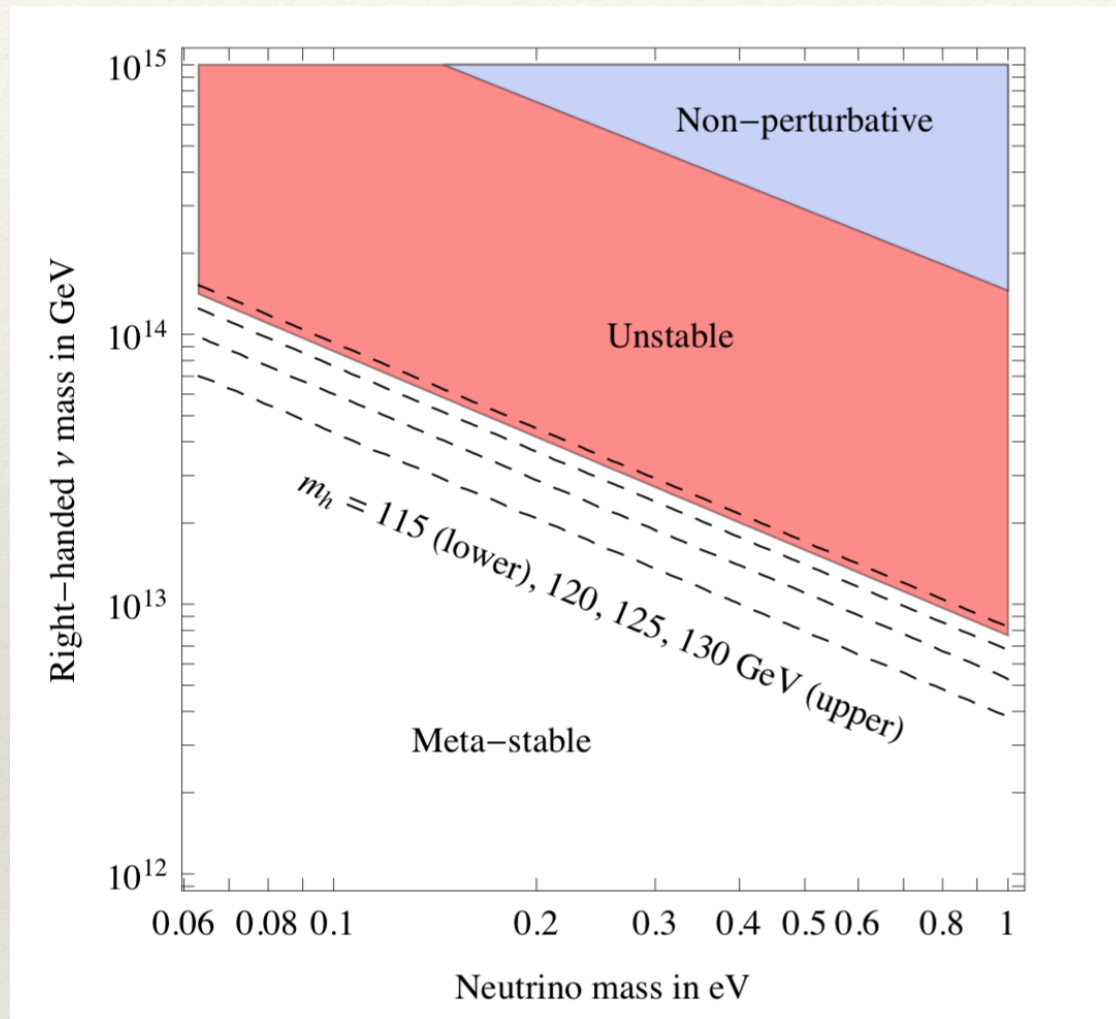
$$\text{BR} \propto y^4 / (M_N^4 \text{ or } M_{\text{SUSY}}^4)$$

**Vacuum stability,
naturalness**

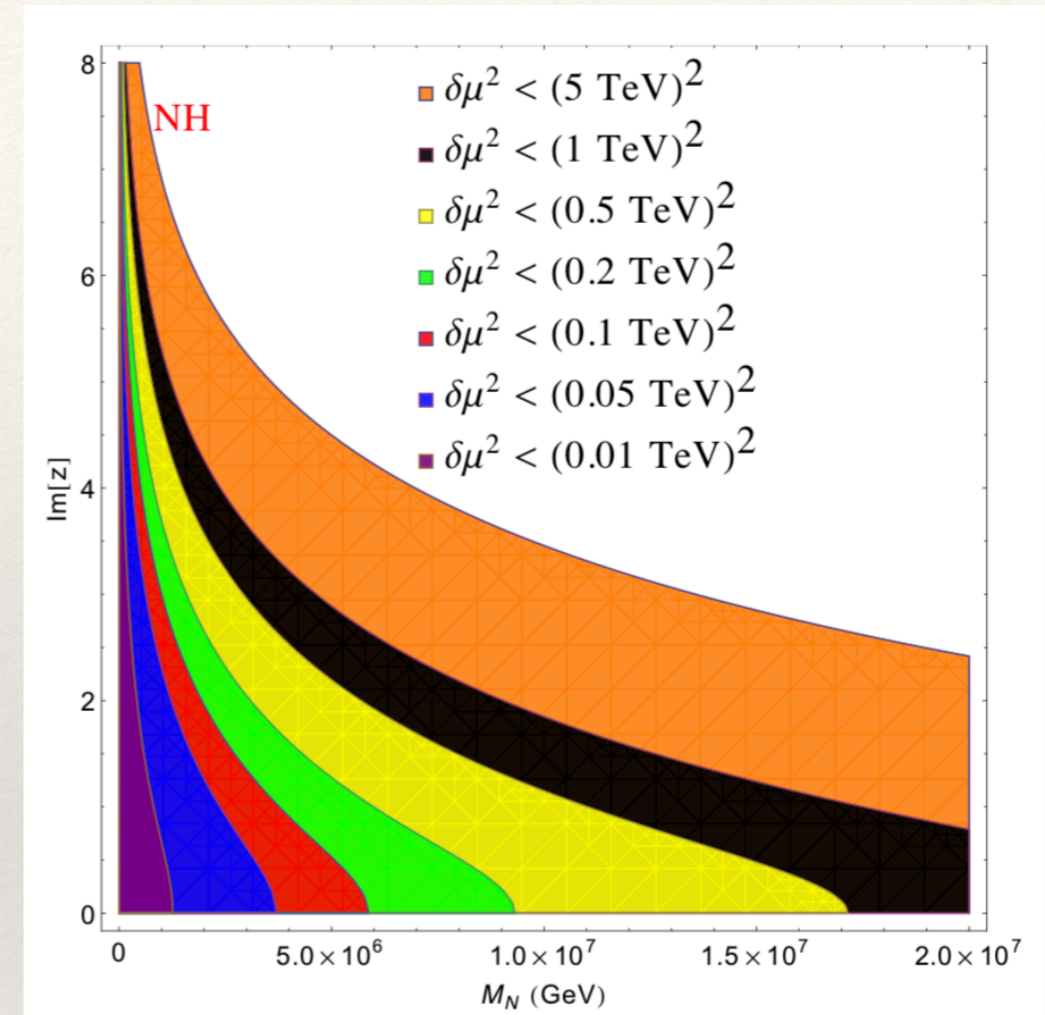
$$d\lambda / dt \propto -y^4$$

$$\delta(m_h^2) \propto y^2 M_N^2$$

Type I Seesaw $m_\nu = m_D^2 / M_R$



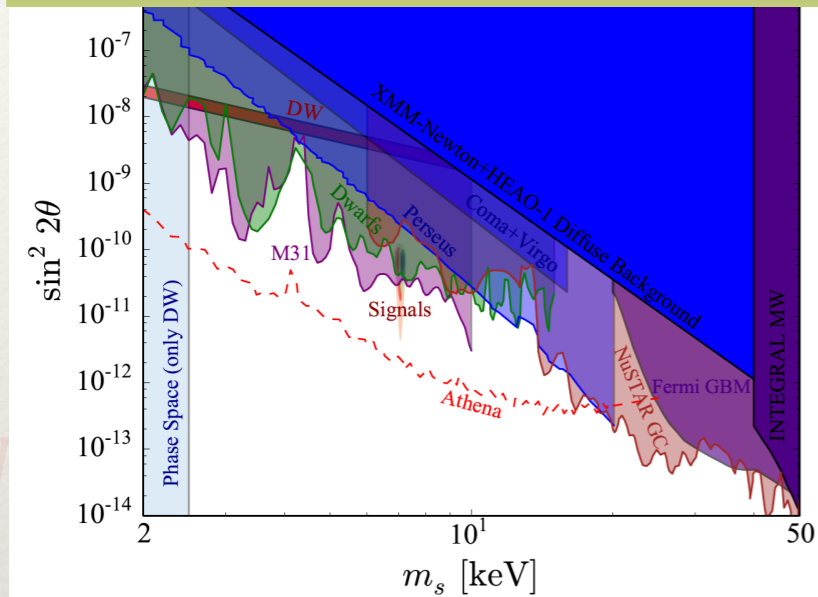
Elias-Miro et al., 1112.3022



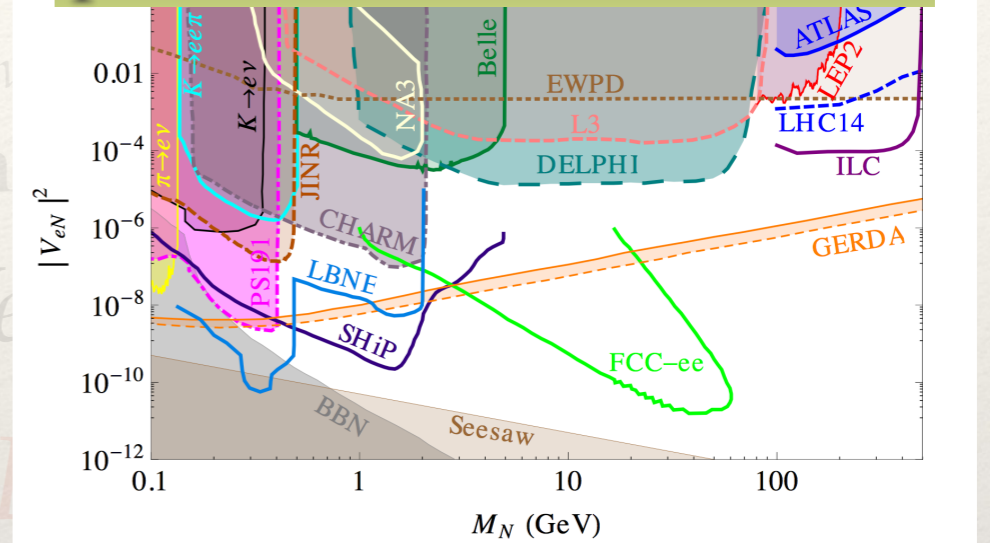
Bambhaniya et al., 1611.03827

Type I Seesaw $m_\nu = m_D^2 / M_R$

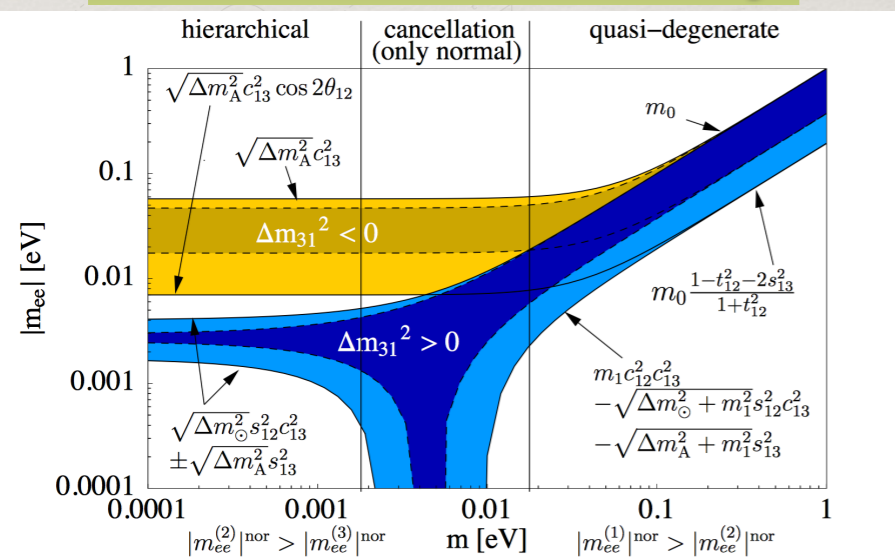
dark matter candidate



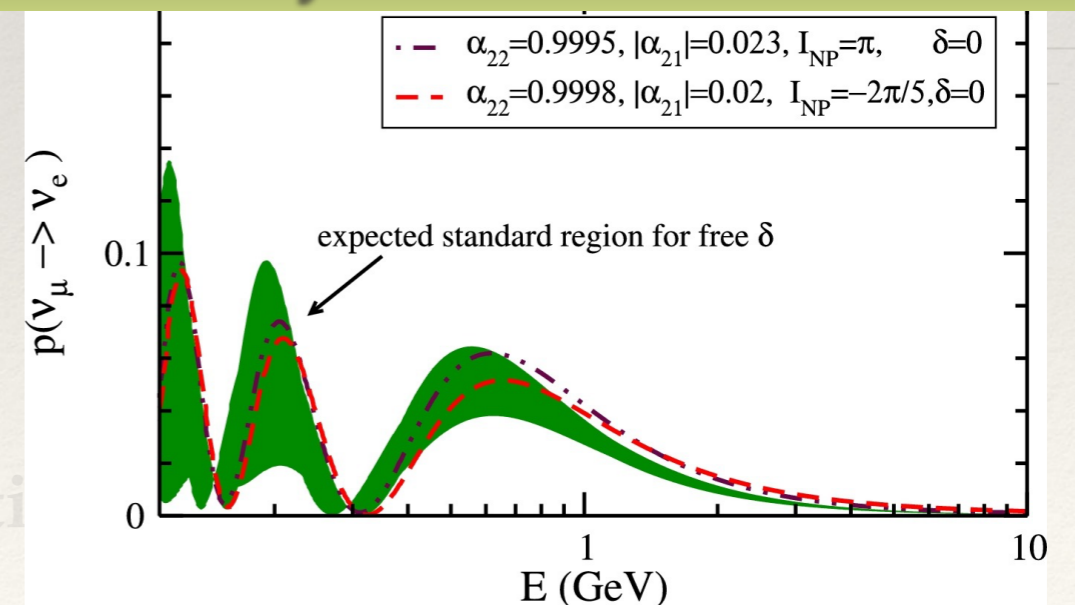
production at colliders



double beta decay



unitarity violation of PMNS



Pathways to Neutrino Mass

similar discussion for all thinkable and unthinkable mass mechanisms

approach	ingredient	quantum number of messenger	\mathcal{L}	m_ν	scale
"SM" (Dirac mass)	RH ν	$N_R \sim (1, 0)$	$h \overline{N}_R \Phi L$	$h v$	$h = \mathcal{O}(10^{-12})$
"effective" (dim 5 operator)	new scale + LNV	–	$h \overline{L^c} \Phi \Phi L$	$\frac{h v^2}{\Lambda}$	$\Lambda = 10^{14}$ GeV
"direct" (type II seesaw)	Higgs triplet + LNV	$\Delta \sim (3, -2)$	$h \overline{L^c} \Delta L + \mu \Phi \Phi \Delta$	$h v_T$	$\Lambda = \frac{1}{h \mu} M_\Delta^2$
"indirect 1" (type I seesaw)	RH ν + LNV	$N_R \sim (1, 0)$	$h \overline{N}_R \Phi L + \overline{N}_R M_R N_R^c$	$\frac{(h v)^2}{M_R}$	$\Lambda = \frac{1}{h} M_R$
"indirect 2" (type III seesaw)	fermion triplets + LNV	$\Sigma \sim (3, 0)$	$h \overline{\Sigma} L \Phi + \text{Tr} \overline{\Sigma} M_\Sigma \Sigma$	$\frac{(h v)^2}{M_\Sigma}$	$\Lambda = \frac{1}{h} M_\Sigma$

plus seesaw variants (linear, inverse, double, singular,...)

plus radiative mechanisms

plus higher dimensional operators

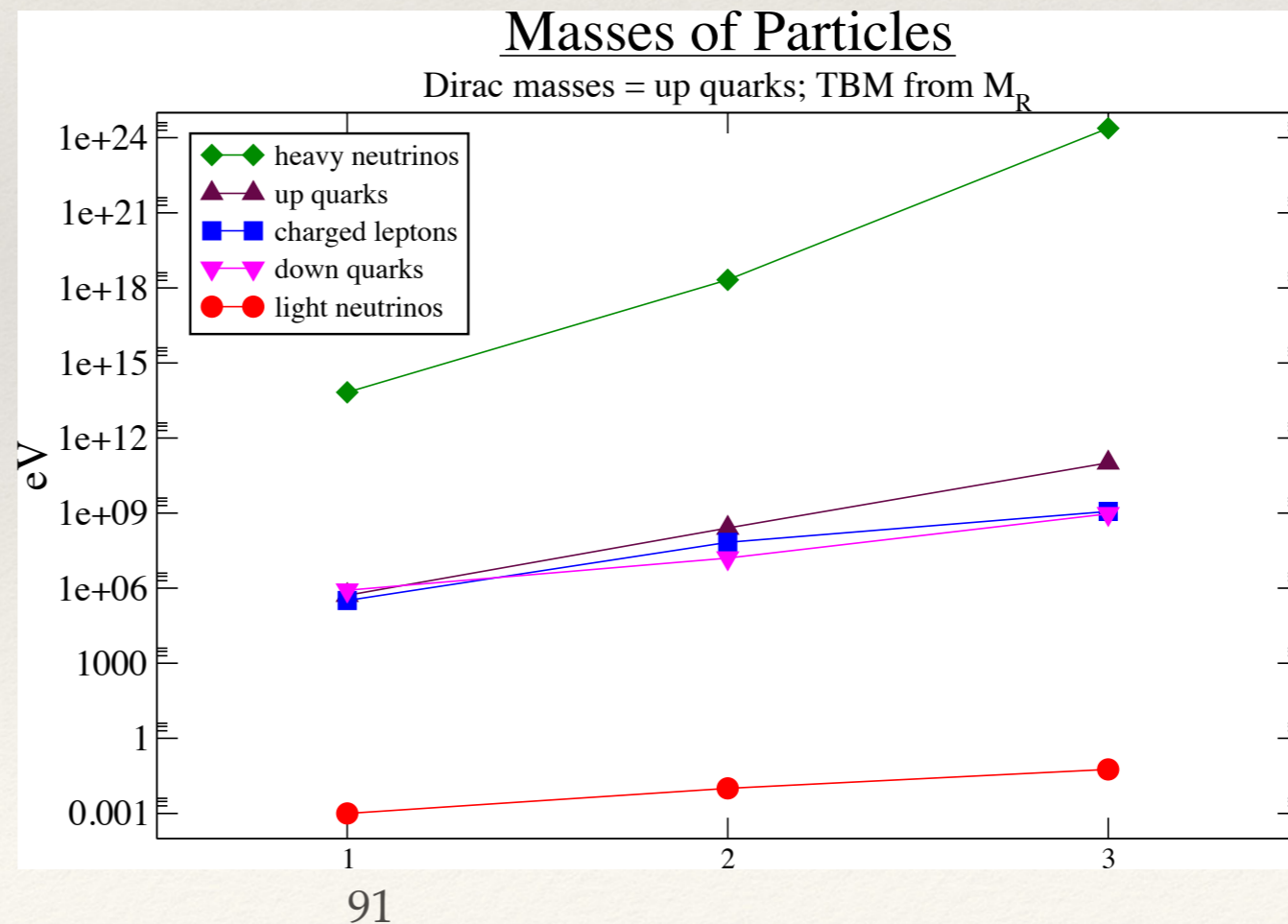
plus extra dimensional

plus plus plus

Seesaw Mechanism

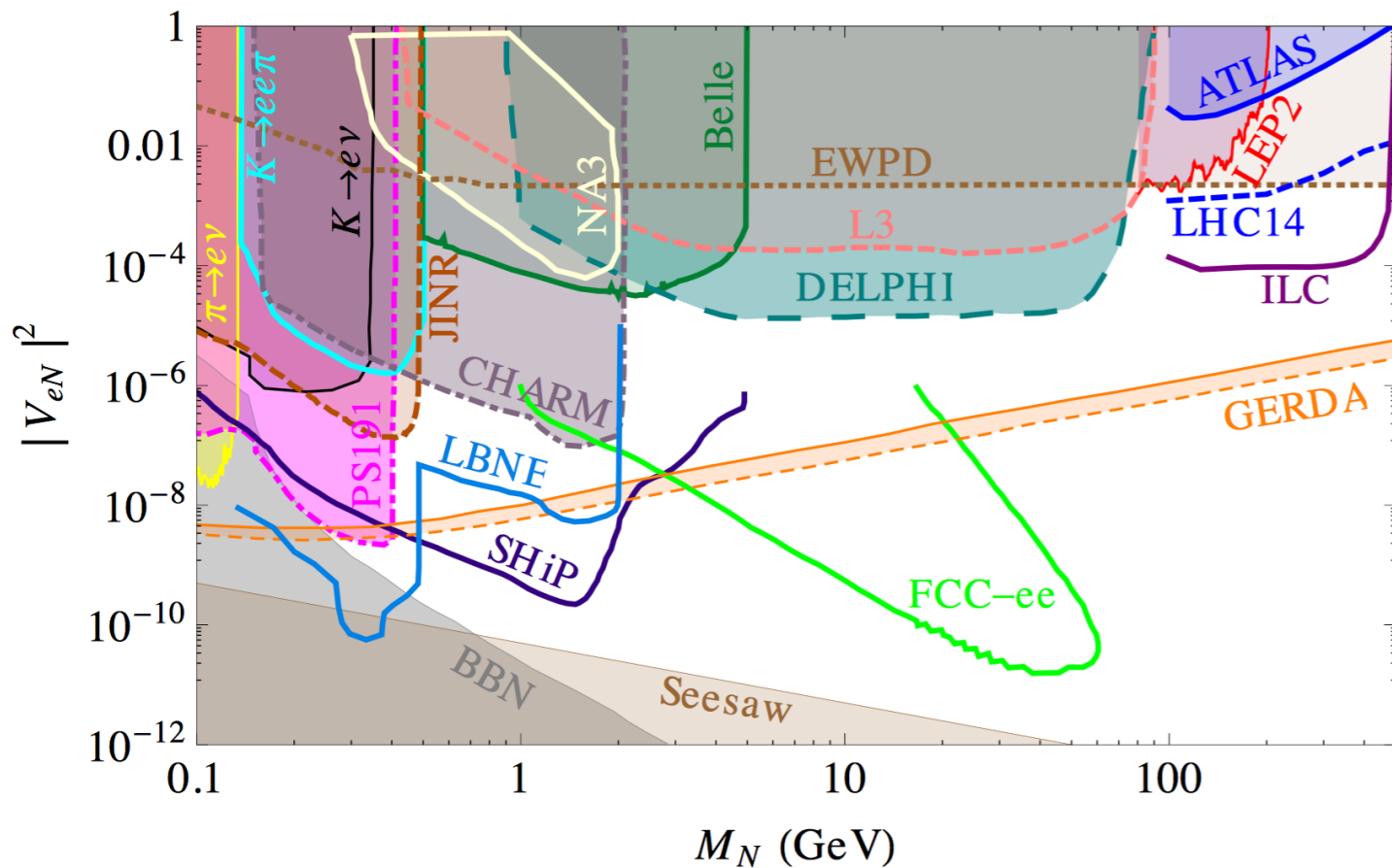
- ❖ suppresses neutrino mass *for each generation*
($m_u \approx m_d$ and $m_b \sim m_t$ vs. $m_{\nu_e} \ll m_e$ and $m_{\nu_\tau} \ll m_\tau$)
- ❖ little hierarchy in m_ν , strong quark-like hierarchy in m_D

⇒ stronger
hierarchy in M_R ?

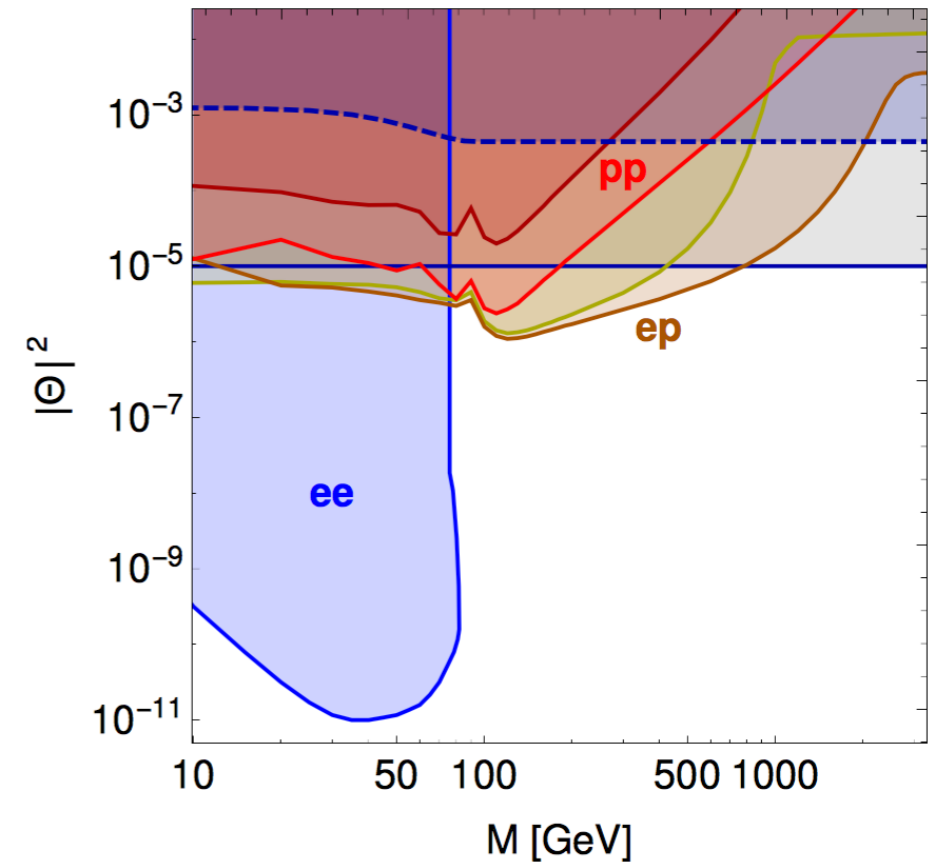


Limits on Heavy Neutrinos

$$M(W_R) \leftrightarrow V_{\alpha N}$$



Deppisch, Dev, Pilaftsis, 1502.06541



Antusch, Cazzato, Fischer, 1612.02728

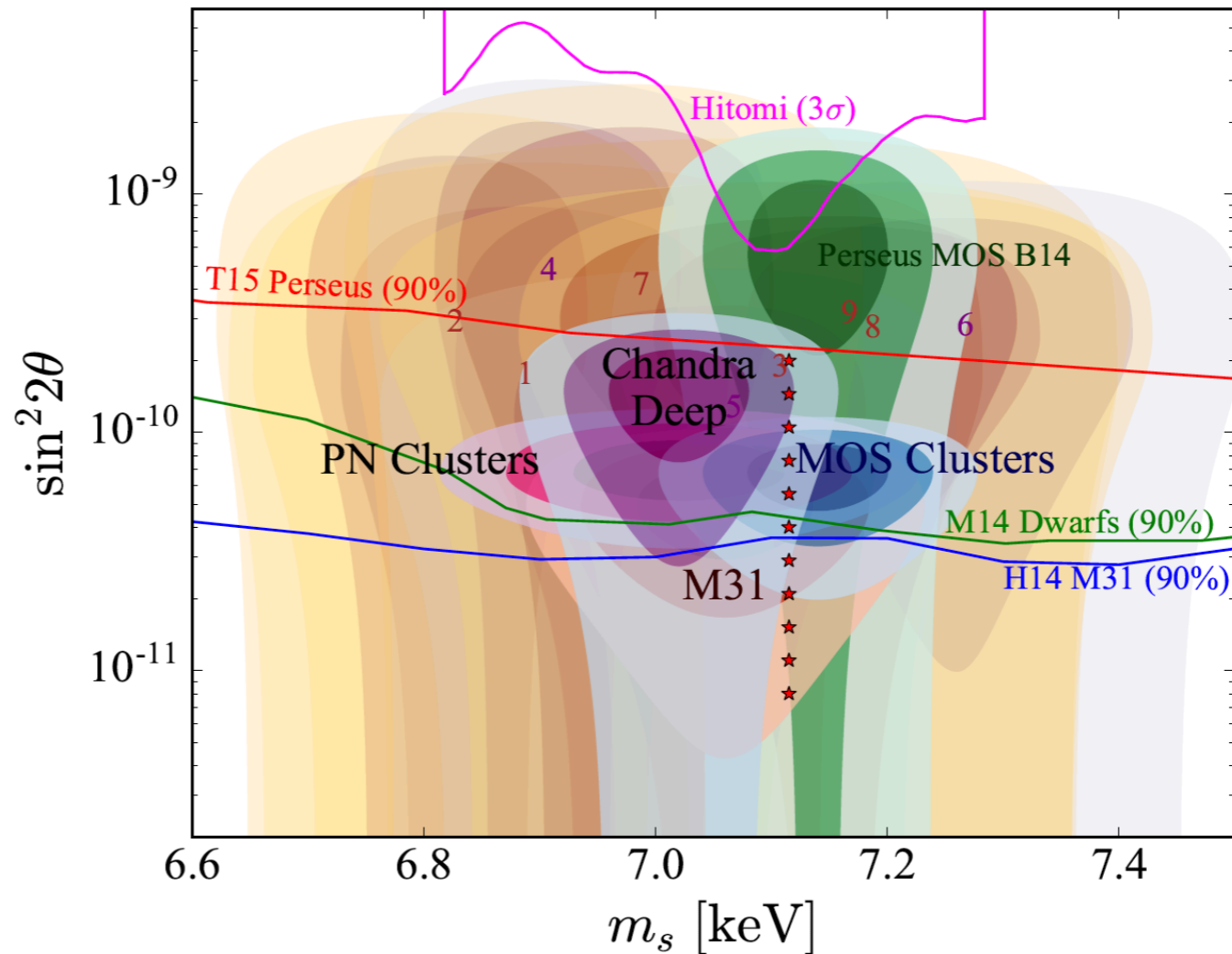
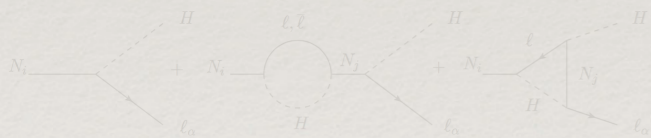
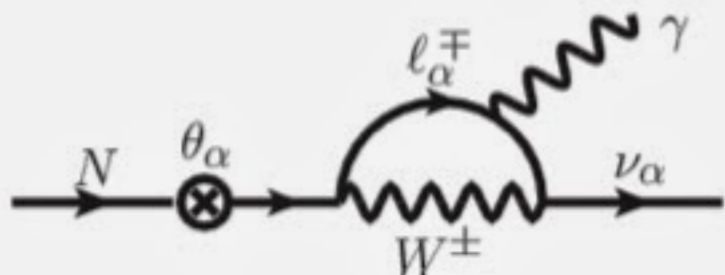
peak searches, kink searches, displaced vertices, LNV decays,...

see also Atre et al., 0901.3589

Type I Seesaw $m_\nu = m_D^2 / M_R$

plus: provides a DM candidate

actually,
needs to be tested



Abazajian, 1705.01837

Leptogenesis

Lepton Flavor Violation

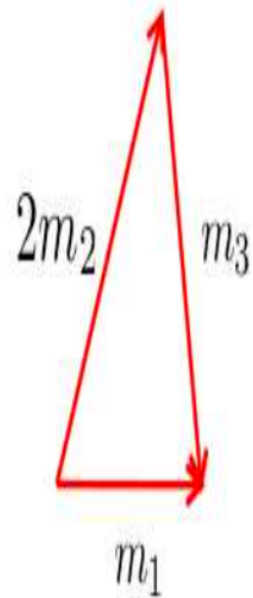
Vacuum stability,
naturalness

Flavor Symmetries

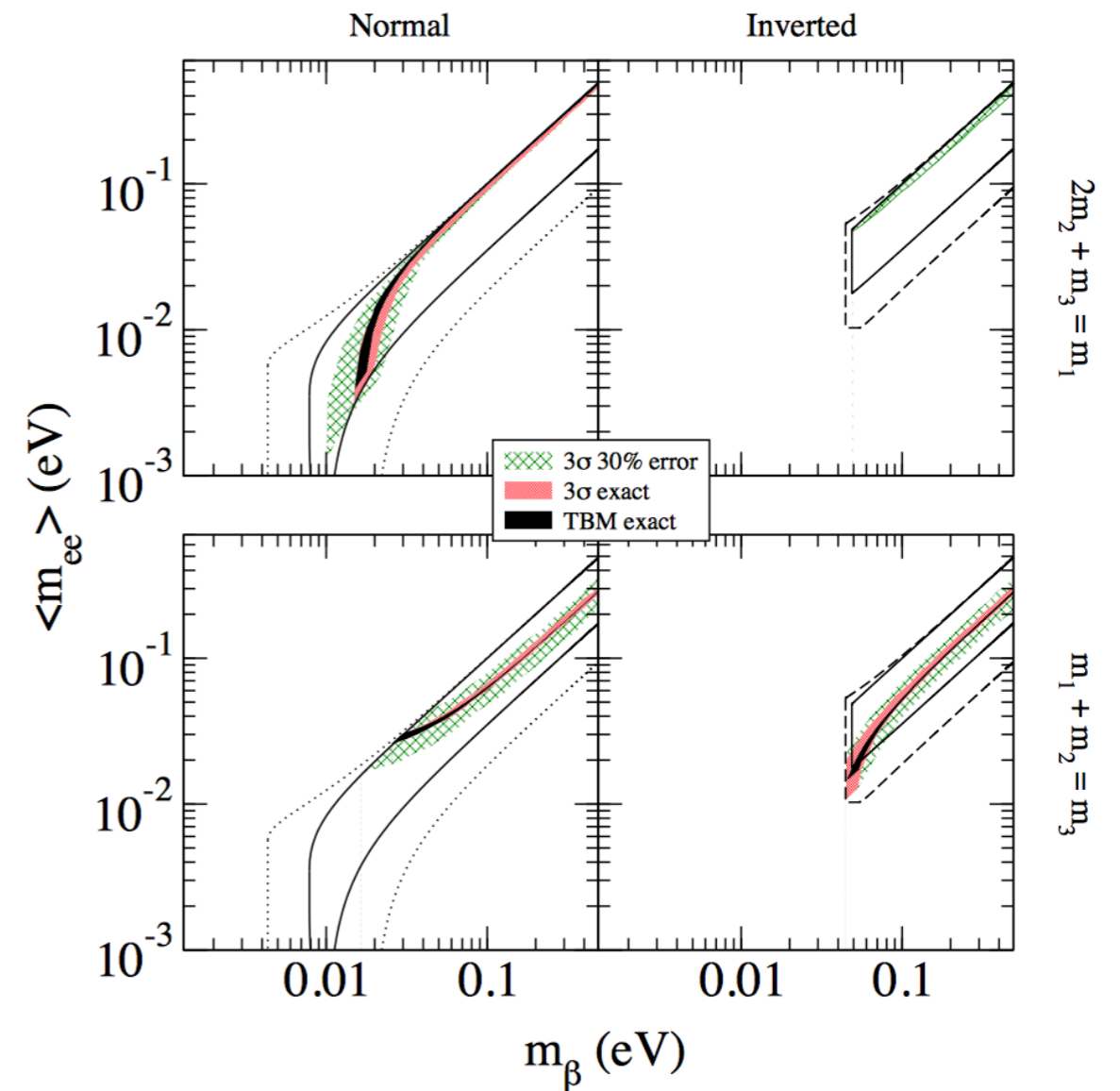
- ❖ Can rule out models by:
 - correlations between angles and phases
 - neutrino mass sum-rules, e.g. $m_1 + m_2 e^{i\alpha} = m_3 e^{i\beta}$
 - LFV if within SUSY or if broken at low scale
 - *minimality*
 - *robustness*
 - *compatibility with larger frameworks (LR symmetry, Pati-Salam, SU(5), SO(10),...)*

Flavor Symmetries

Barry, WR, 1007.5217



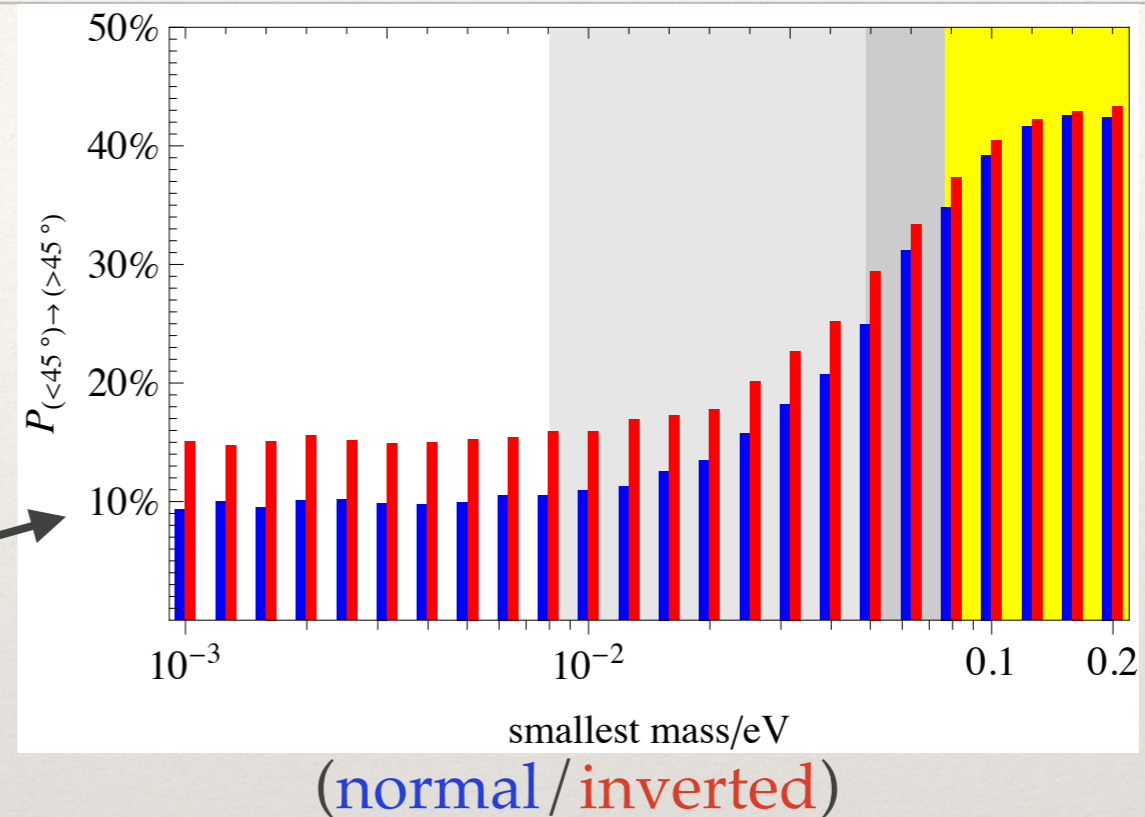
Sum-rule	Flavour symmetry
$2m_2 + m_3 = m_1$	$A_4, T', (S_4)$
$m_1 + m_2 = m_3$	$S_4, (A_4)$
$\frac{2}{m_2} + \frac{1}{m_3} = \frac{1}{m_1}$	A_4, T'
$\frac{1}{m_1} + \frac{1}{m_2} = \frac{1}{m_3}$	S_4



Implications

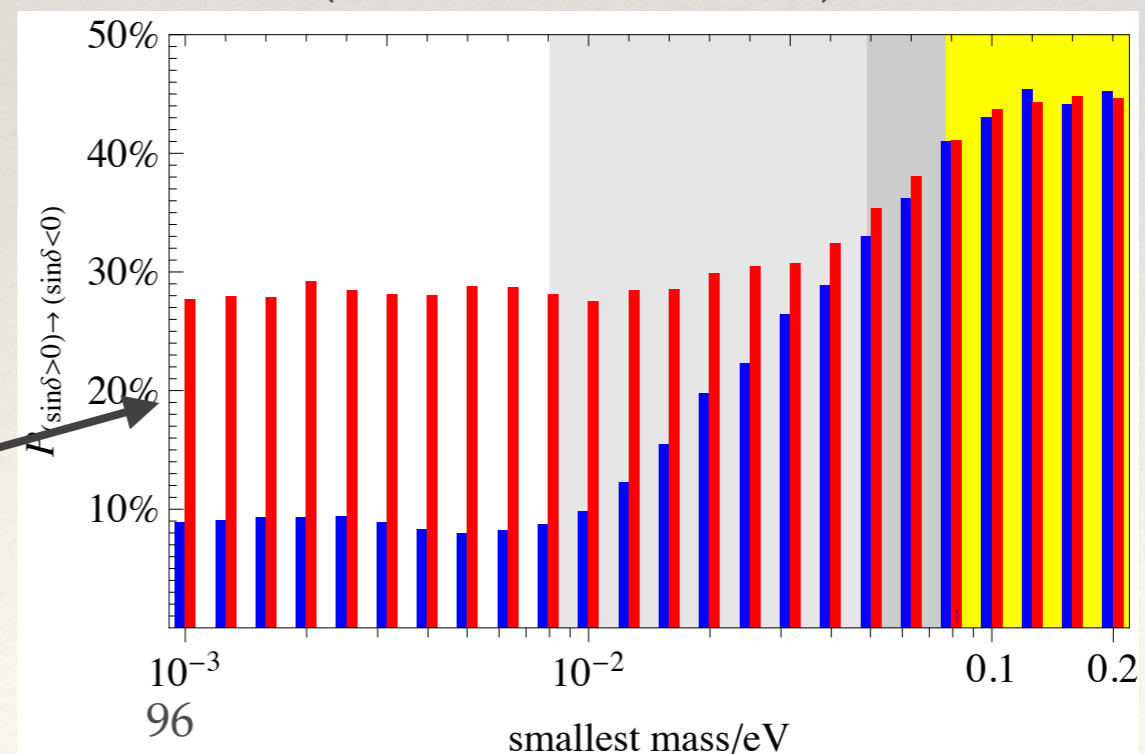
- ❖ Maximal $\theta_{23} = \pi/4$?

probability to change octant of θ_{23}



- ❖ „Maximal“ $\delta = 3\pi/2$?

probability to change sign of $\sin \delta$



WR, Xu, 1508.06063

Implications

❖ Maximal

prob

octant of θ_{23}

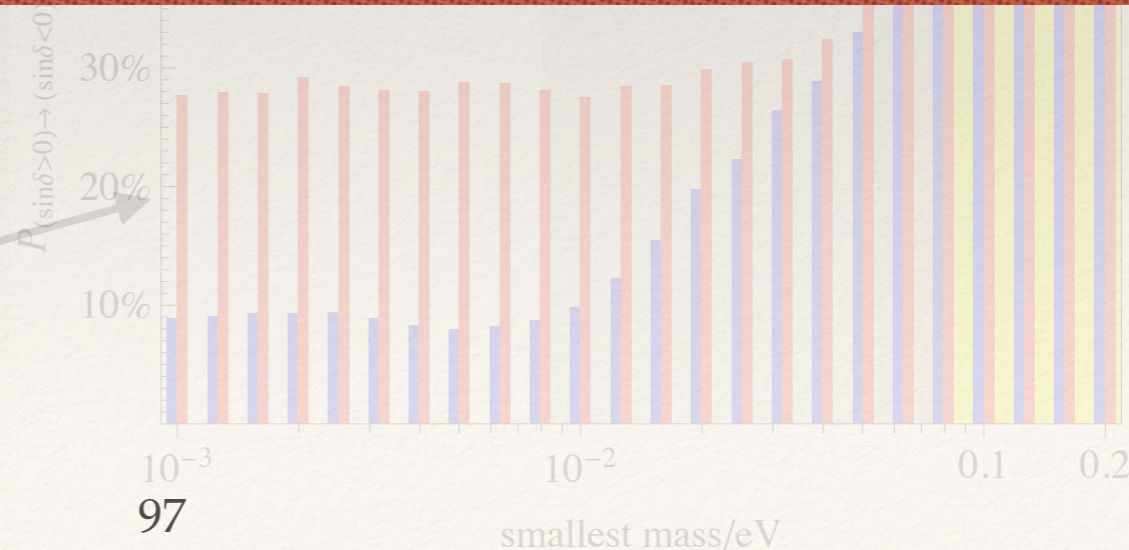
If QD or IH: more need of protection of special values



impact on necessary precision / interpretation of oscillation parameters

❖ „Maximal“ $\delta = 3\pi/2$?

probability to change sign of $\sin \delta$



WR, Xu, 1508.06063

Perturbations

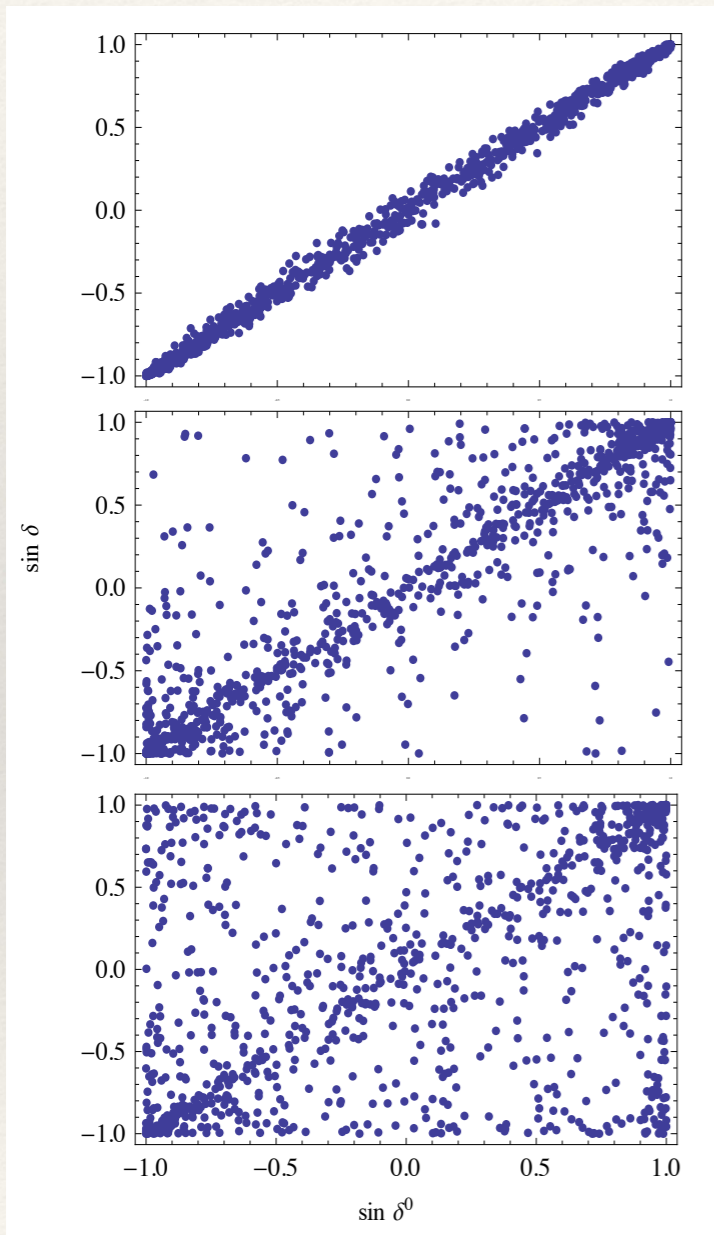
smallest

mass:

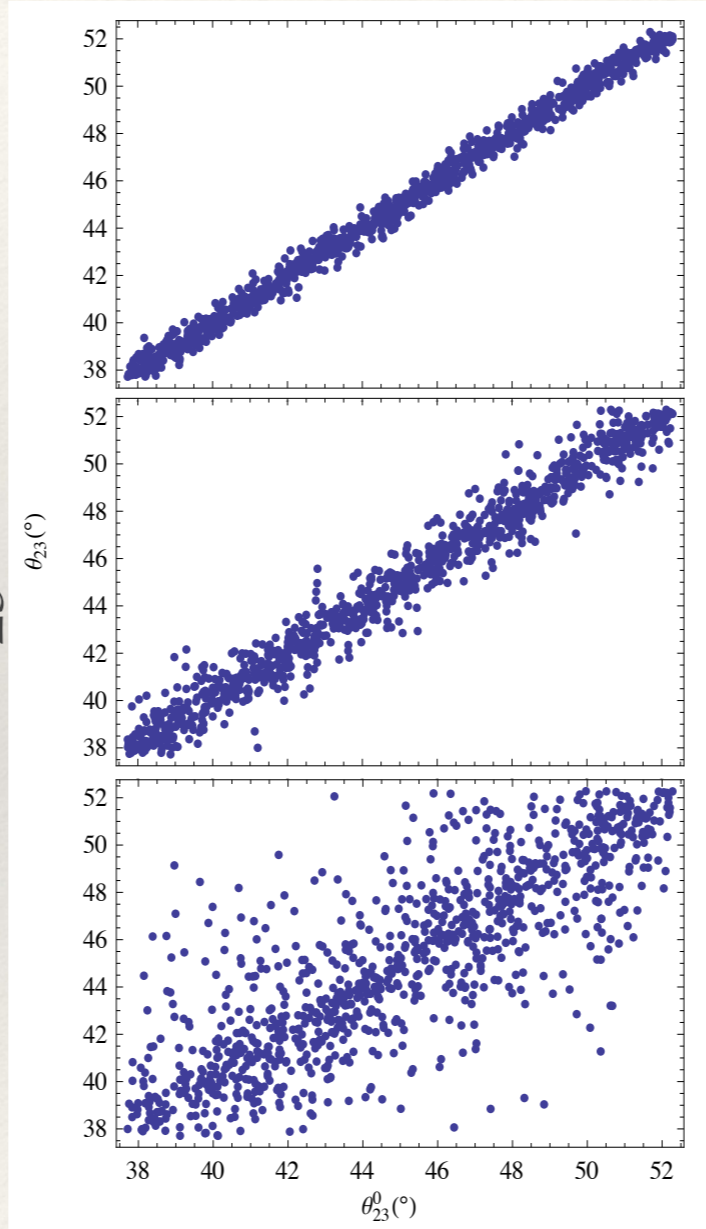
0.001 eV

0.04 eV

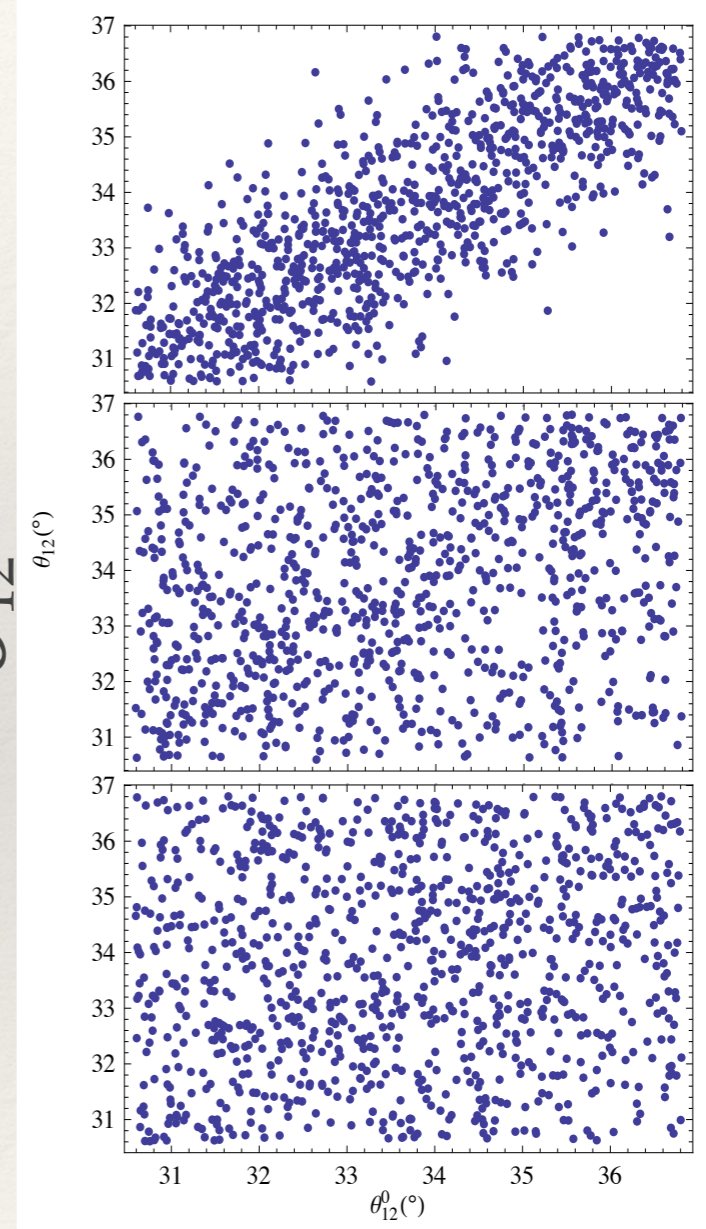
0.1 eV



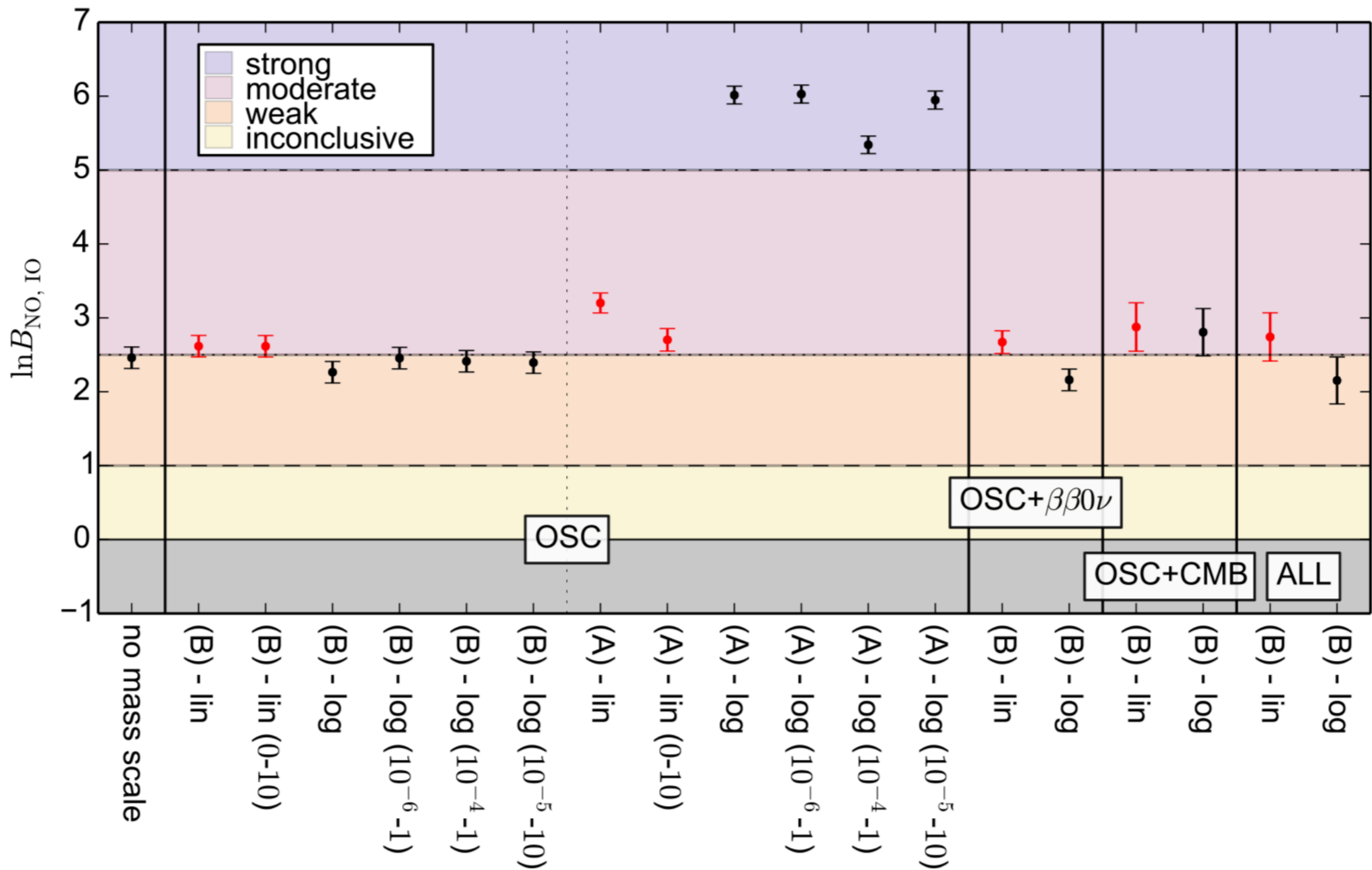
$\sin \delta^0$



θ_{23}^0

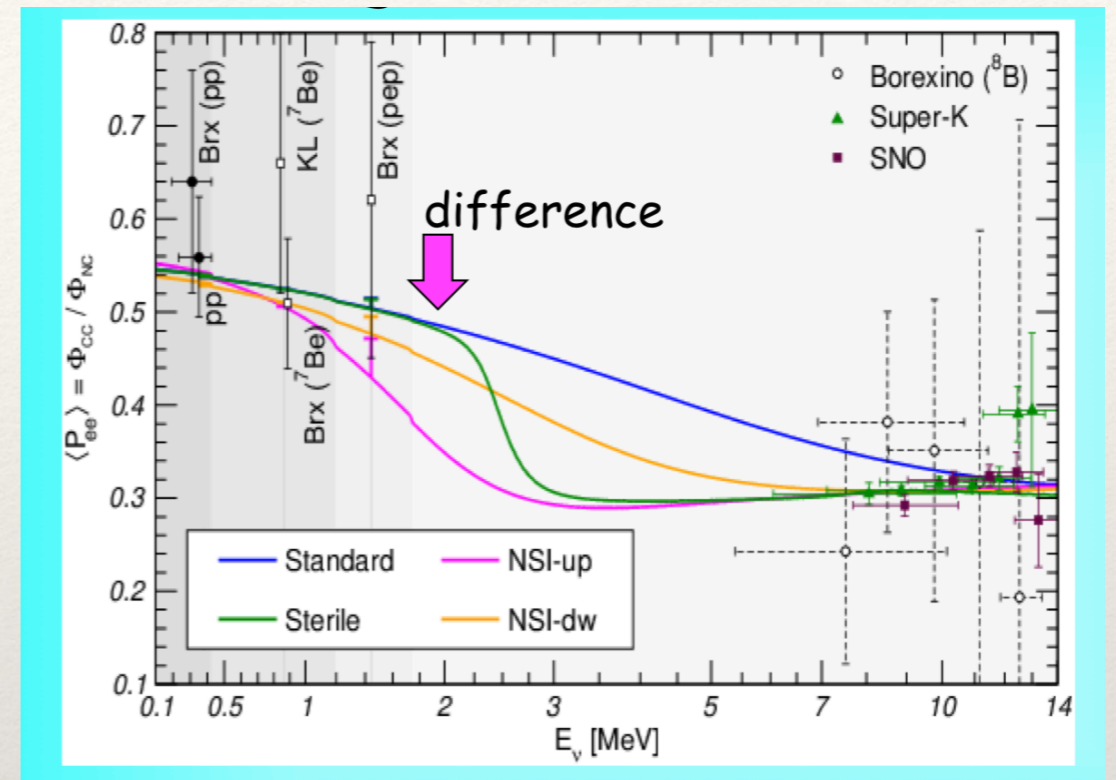
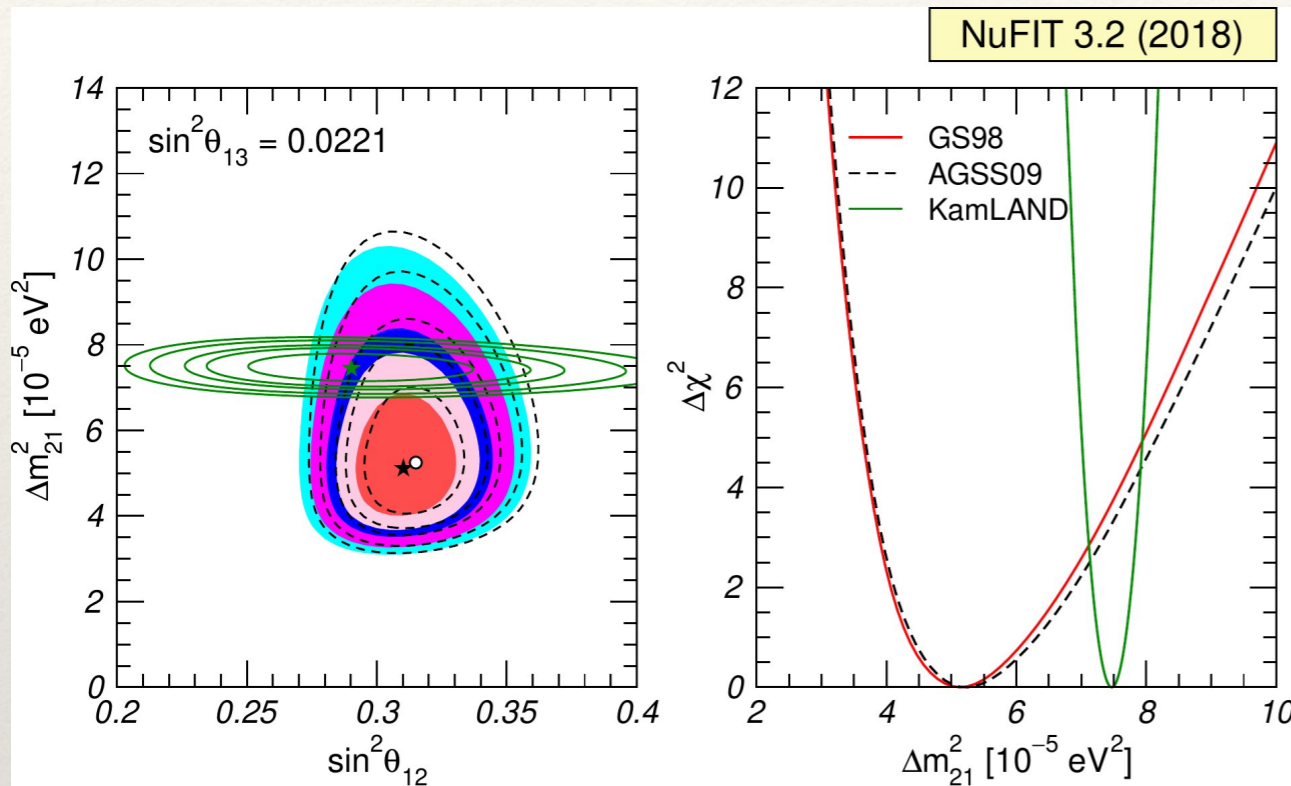


θ_{12}^0



logarithmic priors on masses give more importance to smaller masses, where NO/IO difference is large

Tensions: only in solar sector?



Maltoni, Smirnov, 1507.05287

(plus too large matter effect and too large D/N effect)

Perturbations

- ❖ Various sources:
 - VEV misalignment, NLO terms, RG effects,...
- ❖ Frequent feature: $\delta(\theta_{12}), \delta(\delta) > \delta(\theta_{13}), \delta(\theta_{23})$
- ❖ effects larger for IH and QD

Example RG enhancement:

[in units of $10^{-5} \tan^2 \beta$]

	NH	IH	QD
$\delta(\theta_{12})$	1	$\Delta m_A^2 / \Delta m_\odot^2$	$m_0^2 / \Delta m_\odot^2$
$\delta(\theta_{13})$	$\sqrt{\Delta m_\odot^2 / \Delta m_A^2}$	1	$m_0^2 / \Delta m_A^2$
$\delta(\theta_{23})$	1	1	$m_0^2 / \Delta m_A^2$
$\delta(\delta)$	$\sqrt{\Delta m_\odot^2 / \Delta m_A^2}$	$\Delta m_A^2 / \Delta m_\odot^2$	$m_0^2 / \Delta m_A^2$
$\delta(\alpha, \beta)$	$\sqrt{\Delta m_\odot^2 / \Delta m_A^2}$	$\Delta m_A^2 / \Delta m_\odot^2$	$m_0^2 / \Delta m_A^2$

large

Perturbations

- ❖ Various sources:
 - VEV misalignment, NLO terms, RG effects
- ❖ Frequent feature: $\delta(\theta_{12}), \delta(\delta) > \delta(\theta_{13})$
- ❖ effects larger for IH and QD

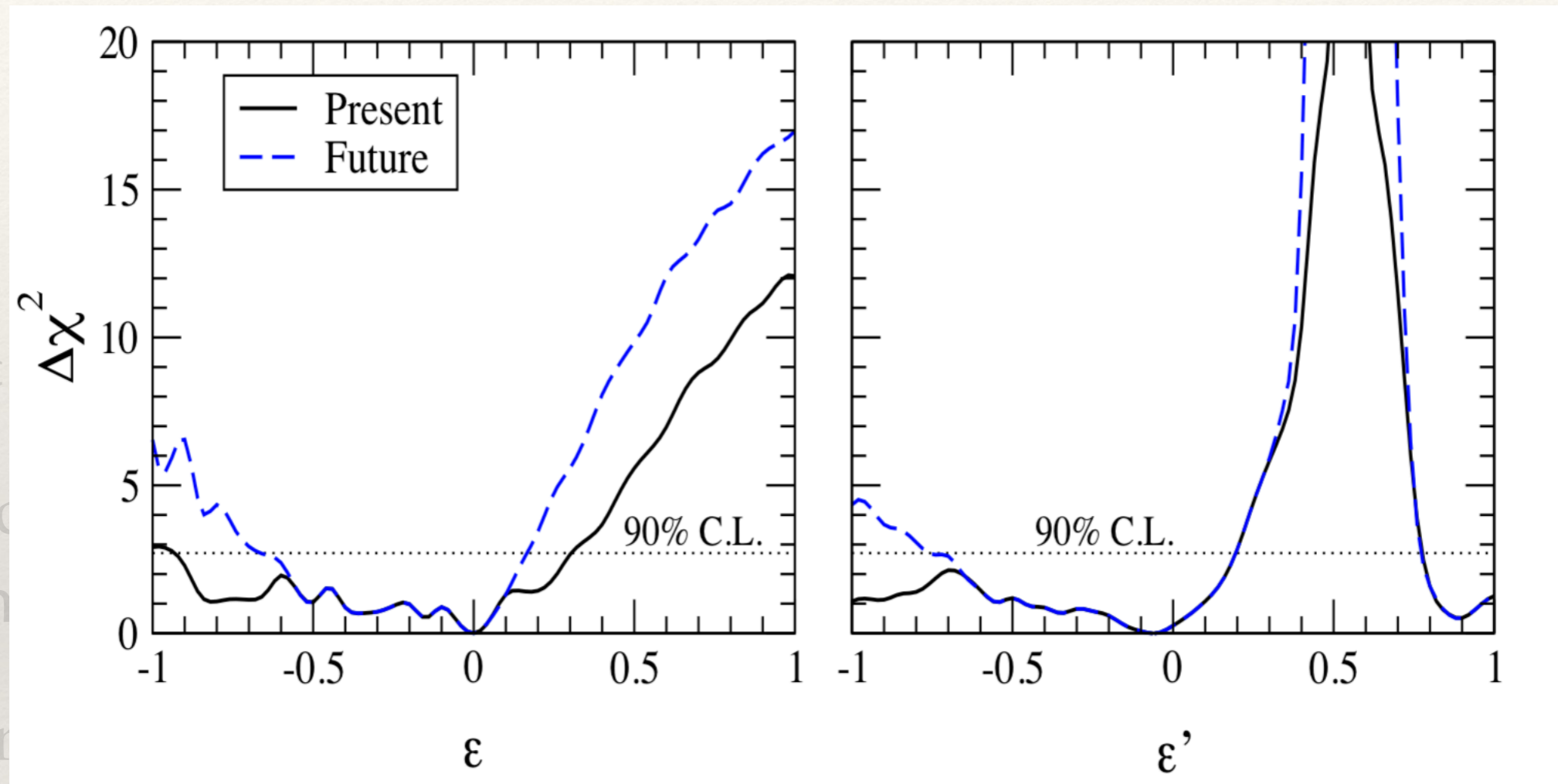
Example RG enhancement:
[in units of $10^{-5} \tan^2 \beta$]

		IH	QD
$\delta(\theta_{12})$	1	$\Delta m_A^2 / \Delta m_\odot^2$	$m_0^2 / \Delta m_\odot^2$
$\delta(\theta_{13})$	$\sqrt{\Delta m_\odot^2 / \Delta m_A^2}$	1	$m_0^2 / \Delta m_A^2$
$\delta(\delta)$	1	1	$m_0^2 / \Delta m_A^2$
$\delta(\alpha, \beta)$	$\sqrt{\Delta m_\odot^2 / \Delta m_A^2}$	$\Delta m_A^2 / \Delta m_\odot^2$	$m_0^2 / \Delta m_A^2$
	$\sqrt{\Delta m_\odot^2 / \Delta m_A^2}$	$\Delta m_A^2 / \Delta m_\odot^2$	$m_0^2 / \Delta m_A^2$

mass scale and ordering helpful

large

Non-Standard Interactions



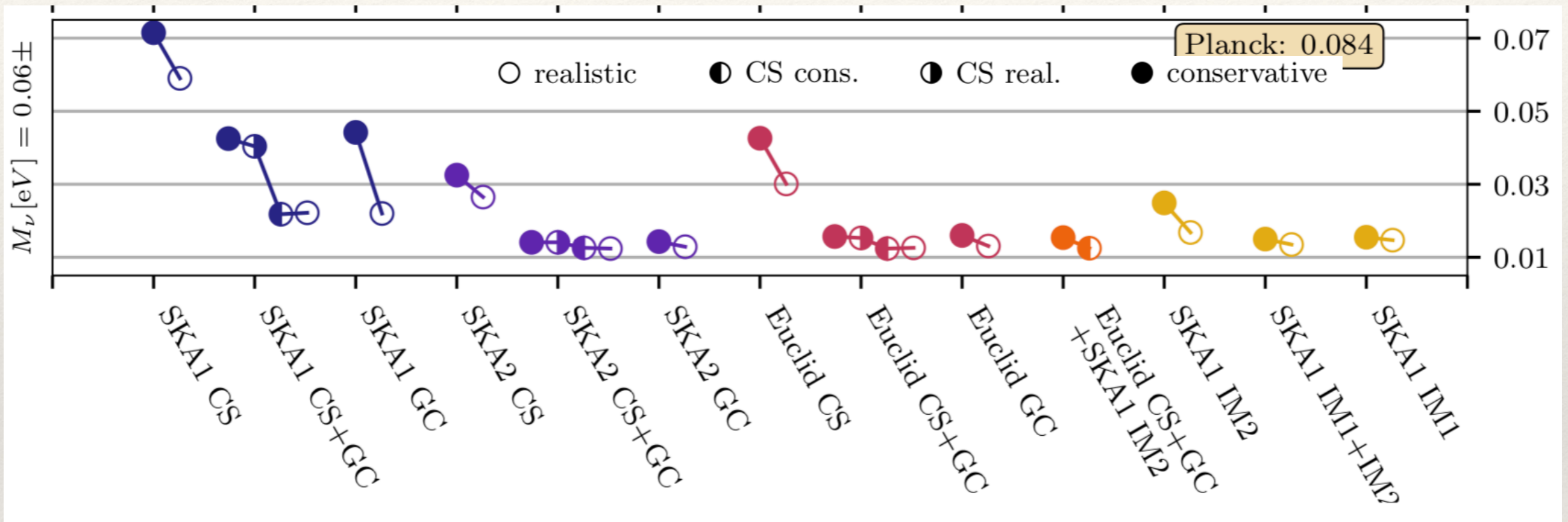
$$\varepsilon = -\sin \theta_{23} \varepsilon_{e\tau}^{dV} \quad \varepsilon' = \sin^2 \theta_{23} \varepsilon_{\tau\tau}^{dV} - \varepsilon_{ee}^{dV}$$

Miranda, Tortola, Valle, hep-ph/0406280

(can also explain small Δm^2 discrepancy in KamLAND/solar and missing upturn of P_{ee})

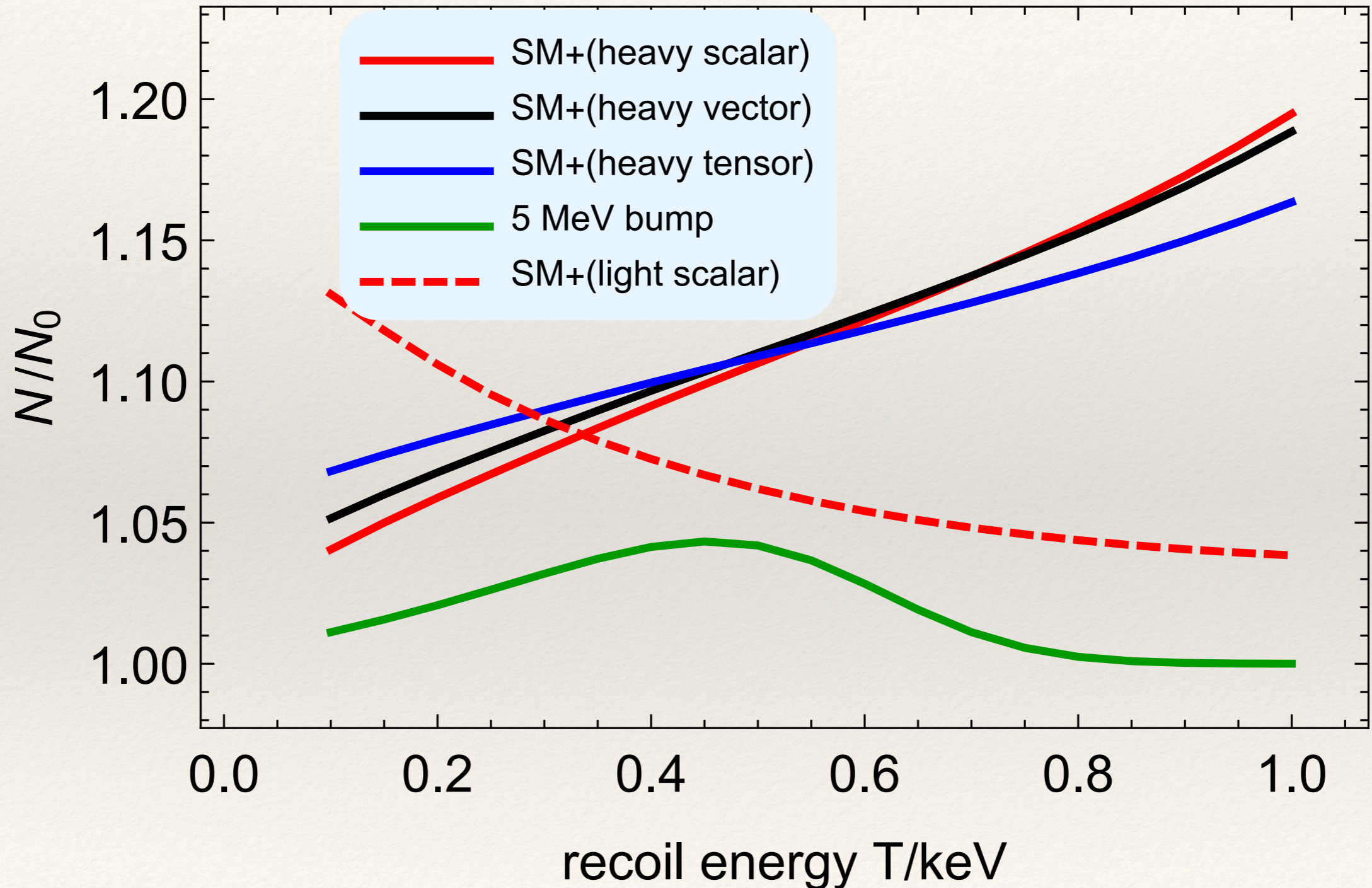
Neutrino Mass guaranteed?

Sprenger et al., 1801.08331



5 σ detection when Euclid and SKA are combined!

New Physics in Coherent Scattering

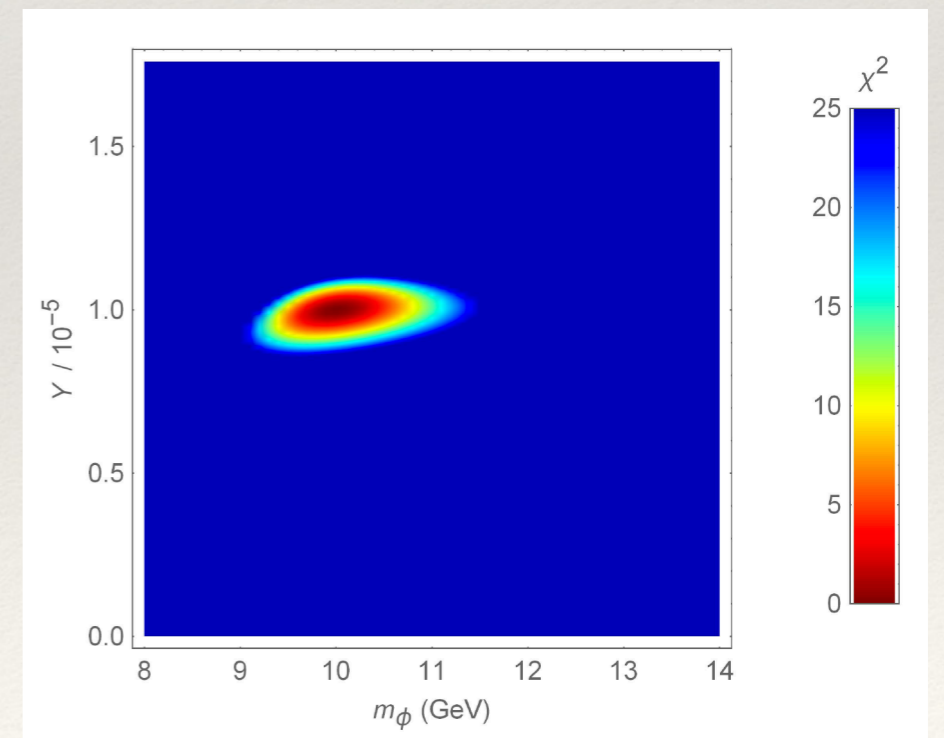
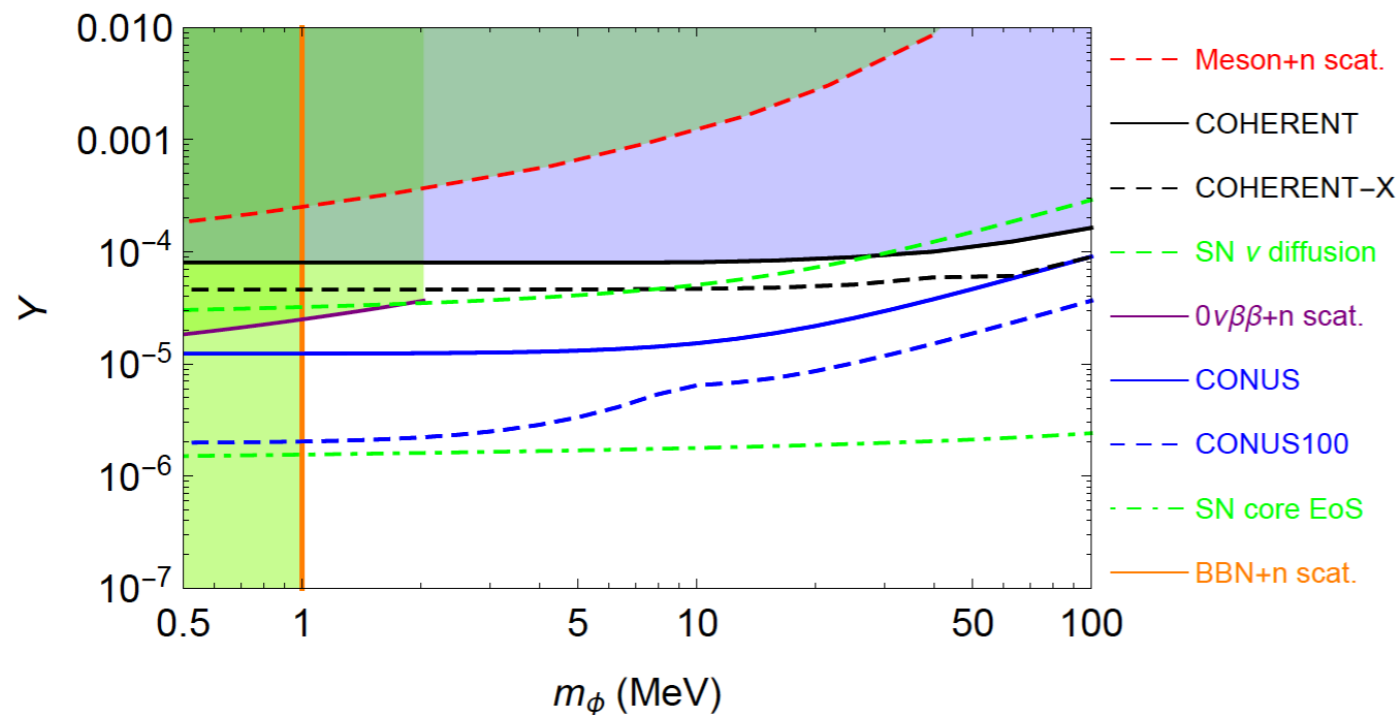
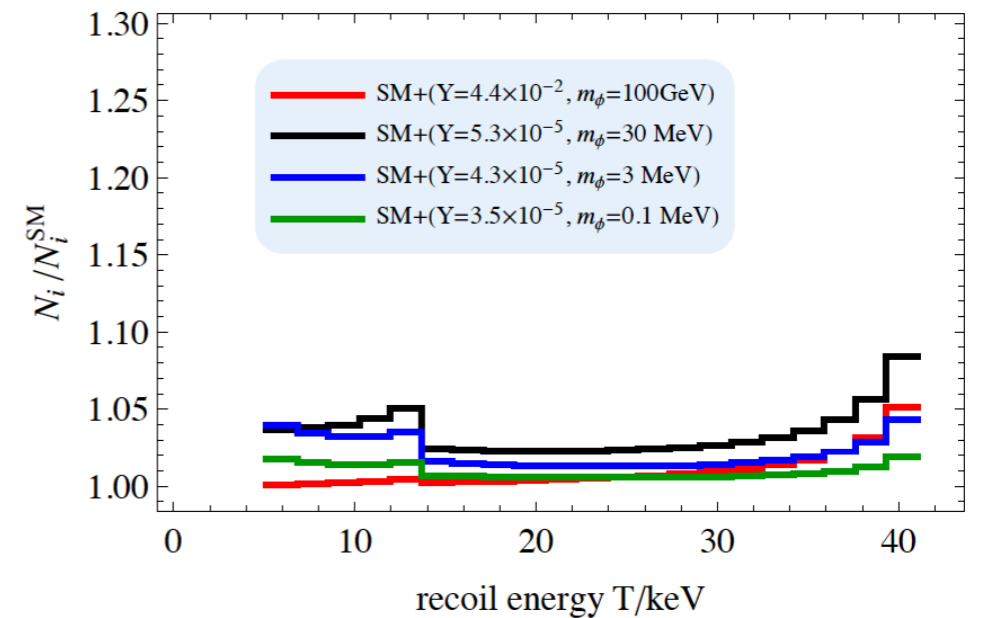
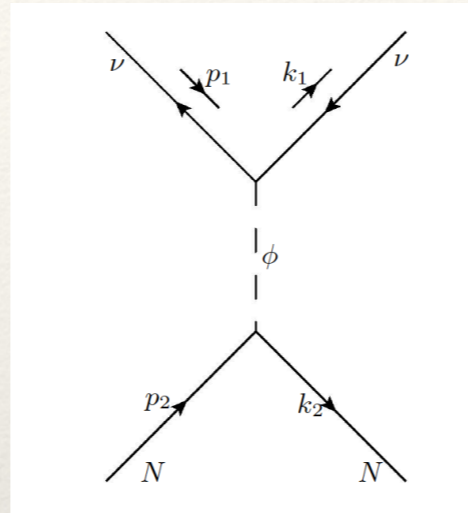


Xun-jie Xu

New Physics in Coherent Scattering

assume light scalar mediator:
(no matter NSI...)

diff. cross section
 $\propto T / (2MT + m_\phi^2)^2 / E_\nu^2$



Farzan, Lindner, WR, Xu, 1802.06171