

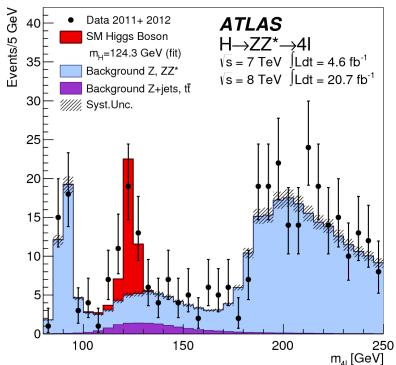
# Contribution of Scalar Singlet and Dimension 5 operator in Higgs Physics

Shekhar Adhikari

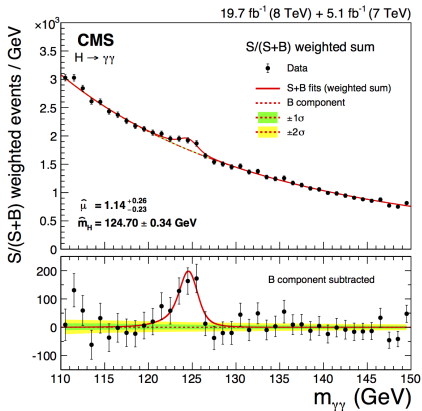
The University of Kansas

May 07, 2018

# Higgs Boson



Higgs boson in 4l channel  
 [www.atlasexperiment.org]

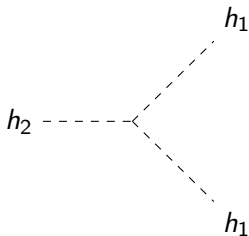


Higgs boson in diphoton channel  
 [cms.web.cern.ch]

- Its discovery verifies well established SM theory.
- Central piece of SM.

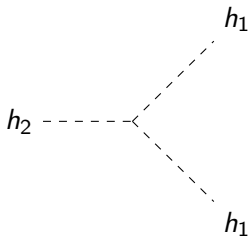
# Motivation for Singlet Extension of SM

- One of the simplest extension of SM is by adding real scalar singlet.
- After EWSB, singlet mixes with SM Higgs as a result, shape of potential as well as Higgs coupling to SM particles are changed.
- New decay signature is observed:  $h_2 \rightarrow h_1 h_1$



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## Further extension of SM

Further potential shape and Higgs coupling are altered with the inclusion of effective Lagrangian.

- The effective Lagrangian approach is a model-independent way to describe new physics (NP) at cut-off scale ( $\Lambda$ ).
- Effective Lagrangian has higher dimension operator suppressed by some power of  $\Lambda$ .

- The model contains SM and new physics(NP) Lagrangian.

$$\mathcal{L} = \mathcal{L}_{SM} + \mathcal{L}_{\leq 5}^S \quad (1)$$

- $\mathcal{L}_{\leq 5}^S$  is NP Lagrangian up to dim 5 built from (S) and SM fields.
- $\mathcal{L}_{\leq 5}^S$  can be split as scalar interaction up to dim 4 and effective interaction( $\mathcal{L}_5^S$ ) of dim 5.

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## Effective Lagrangian

[M. Bauer, A. Butter, J. Gonzalez-Fraile, T. Plehn, and M. Rauch, Phys. Rev., vol. D95, no. 5, p. 055011, 2017]

$$\begin{aligned} \mathcal{L}_5^s = & -\frac{a_3}{2\Lambda} S^3 (\Phi^\dagger \Phi) - \frac{a_4}{2\Lambda} S (\Phi^\dagger \Phi)^2 - \frac{b_5}{5\Lambda} S^5 + g_s^2 \frac{f_{GG}^s}{\Lambda} S G_{\mu\nu}^a G^{a\ \mu\nu} \\ & + \frac{e^2}{\cos^2 \theta_w} \frac{f_{BB}^s}{\Lambda} S B_{\mu\nu} B^{\mu\nu} + \frac{e^2}{\sin^2 \theta_w} \frac{f_{WW}^s}{\Lambda} S W_{\mu\nu}^a W^{a\ \mu\nu} \\ & + \left( -\frac{f_d^s}{\Lambda} S \bar{Q}_L \Phi d_R - \frac{f_u^s}{\Lambda} S \bar{Q}_L \tilde{\Phi} u_R - \frac{f_l^s}{\Lambda} S \bar{L}_L \Phi l_R + h.c \right) \end{aligned} \quad (2)$$

# Procedure to Compute Decay Width

- With our model, we start with scattering amplitude of Higgs ( $h_1$ ) to SM particle.

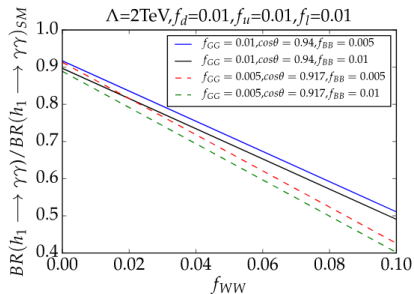
$$\mathcal{M} = \mathcal{M}_{SM} + \mathcal{M}_{dim=5} \quad (3)$$

- The squared amplitude is

$$\begin{aligned} |\mathcal{M}|^2 &= |\mathcal{M}_{SM}|^2 + |\mathcal{M}_{dim=5}|^2 + 2\text{Re}(|\mathcal{M}_{SM}| |\mathcal{M}_{dim=5}|) \\ &\sim \frac{1}{\Lambda^0} \quad \sim \frac{1}{\Lambda^2} \quad \sim \frac{1}{\Lambda^1} \end{aligned} \quad (4)$$

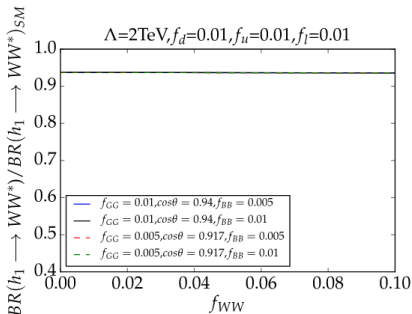
- The decay width has contribution from all three terms.
- We take SM and interference term only and see how branching ratios (BR) of Higgs are effected.

# Result: BR of Higgs to $\gamma\gamma$ and $WW^*$



## Higgs ( $h_1$ ) to $\gamma\gamma$

- BR down by approximately 40-50% compared to SM.
- BR is more sensitive to  $f_{WW}$  in  $h_1 \rightarrow \gamma\gamma$  compared to  $h_1 \rightarrow WW^*$

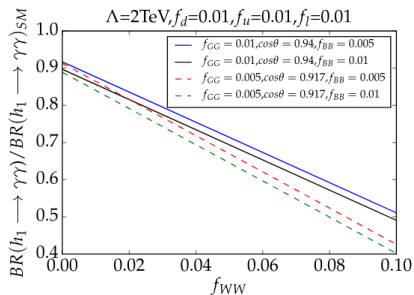


## Higgs ( $h_1$ ) to $WW^*$

- BR down by approximately 5-7% compared to SM.

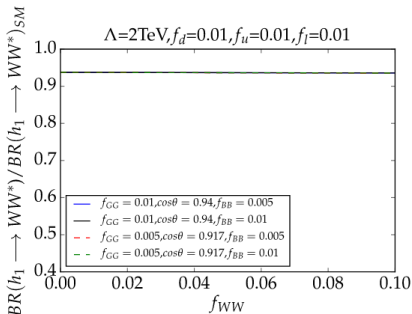


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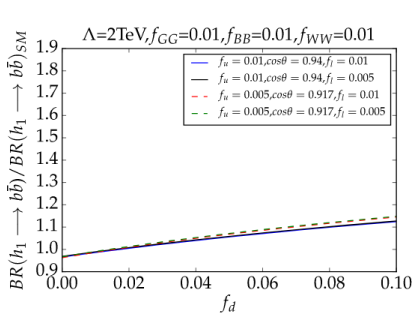


## Higgs ( $h_1$ ) to $WW^*$

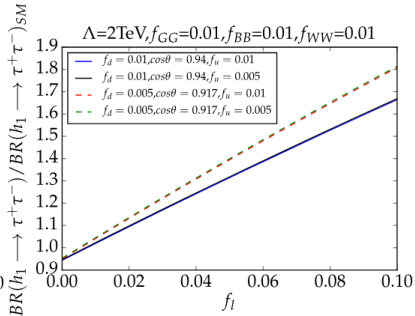
- BR down by approximately 5-7% compared to SM.

[S. Dawson and I. M. Lewis, Phys. Rev., vol. D95, no. 1, p. 015004, 2017]

# BR of Higgs to $b\bar{b}$ and $\tau^+\tau^-$



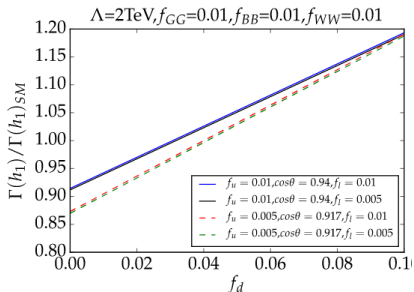
Higgs to  $b\bar{b}$



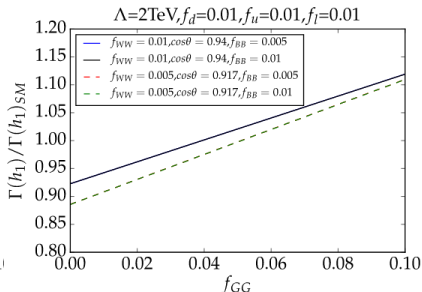
Higgs to  $\tau^+\tau^-$

- BR increases by approximately 10-15% compared to SM.
- BR increases by approximately 65-80% compared to SM.
- BR is more sensitive to effective fermionic coupling coefficient in  $h_1 \rightarrow \tau^+\tau^-$  compared to  $h_1 \rightarrow b\bar{b}$

# Variation of Higgs total width with $f_d$ and $f_{GG}$



Variation of Higgs total width with  $f_d$



Variation of Higgs total width with  $f_{GG}$

- With increase in  $f_d$  and  $f_{GG}$ , the total width also increases.

# Constraints from Higgs Measurements ( $\sqrt{s}=7$ and 8 TeV)

[arxiv:1606.02266]

- The observed signal strength for Higgs production from gluon fusion is:

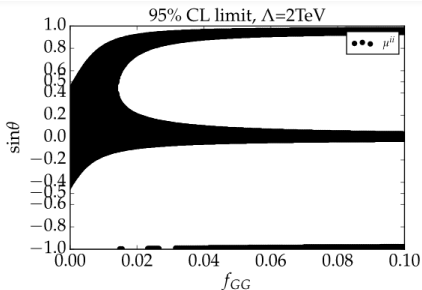
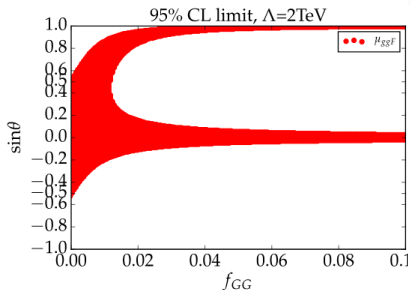
$$\mu_{ggF} = 1.03_{-0.14}^{+0.16} \quad (5)$$

- The signal strength for Higgs decaying to final states are:

$$\begin{aligned} \mu^{\gamma\gamma} &= 1.14_{-0.18}^{+0.19} \\ \mu^{ZZ} &= 1.29_{-0.23}^{+0.26} \\ \mu^{WW} &= 1.09_{-0.16}^{+0.18} \\ \mu^{\tau\tau} &= 1.11_{-0.22}^{+0.24} \\ \mu^{bb} &= 0.70_{-0.27}^{+0.29} \end{aligned} \quad (6)$$

- The observed signal strengths will put bounds on mixing angle and effective coupling coefficients.

# Limits on Mixing angle and $f_{GG}$

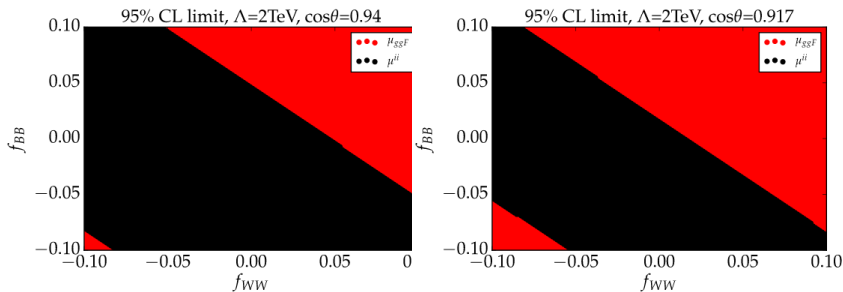


95% CL allowed region of  $\sin\theta$  and  $f_{GG}$

Red is allowed region by a fit to  $\mu_{ggF}$  and black to  $\mu^{ii}$

[S. Dawson and I. M. Lewis, Phys. Rev., vol. D95, no. 1, p. 015004, 2017]

# Limits on Effective Coupling Coefficient ( $f_{BB}$ and $f_{WW}$ )



95% CL allowed region of  $f_{BB}$  and  $f_{WW}$

⇒ Red is allowed region by a fit to  $\mu_{ggF}$  and black to  $\mu^{ii}$ .

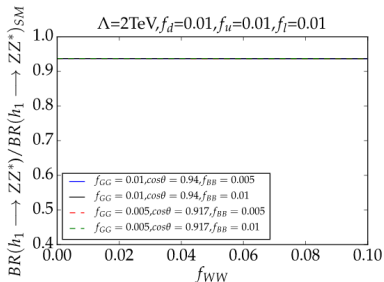
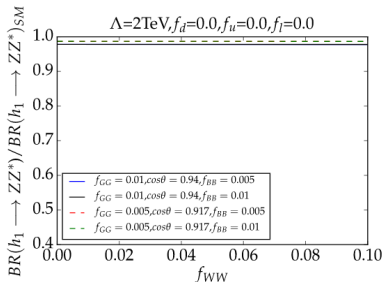
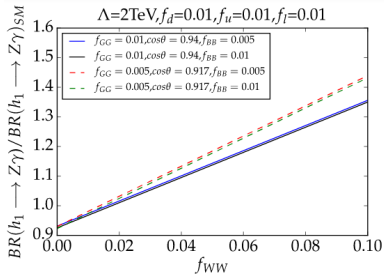
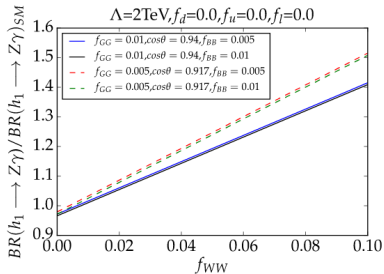
⇒ With increase in mixing angle, the black region shrink.

- Adding real scalar singlet and dimension 5 operator into SM, we found Higgs physics deviated from SM sector.
- Using fit to observed signal strengths, we found the allowed regions of scalar mixing angle and effective coupling coefficients.

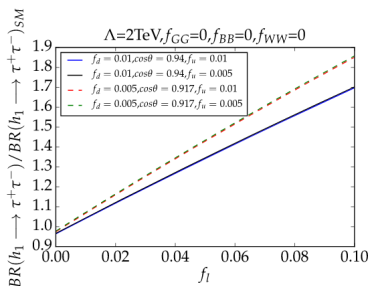
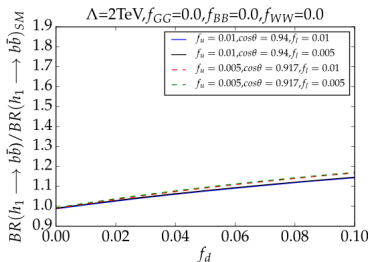
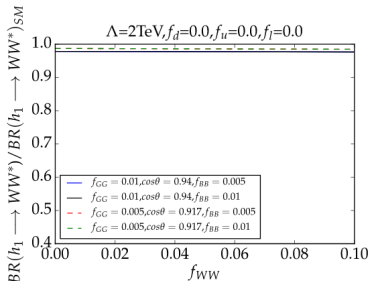
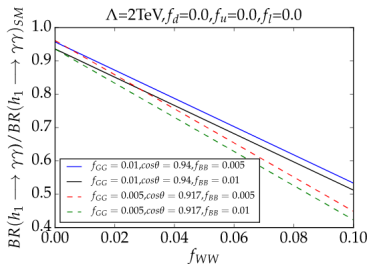
THANK YOU



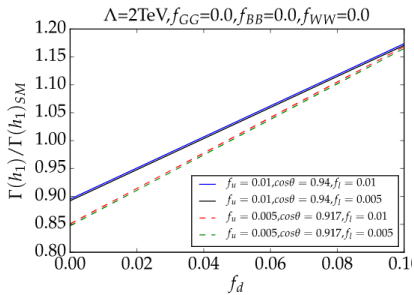
# Higgs to $Z\gamma$ and $ZZ^*$



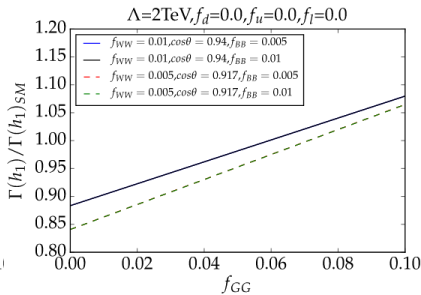
# Higgs to $\gamma\gamma$ , $WW^*$ , $b\bar{b}$ , and $\tau^+\tau^-$



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