LU...anomalies, naturalness, new physics & more

Pheno 2018
Univ of Pittsburgh
Amarjit Soni
BNL-HET

05/08/18

Outline

- several looming deviations from SM ...i.e.
 "anomalies"
- For each case:
- briefly mention reservations for expt & for theory/comments
- Model independent collider implications
- Assuming NP is a source: An interesting, minimal setup for a BSM origin
- Summary & Outlook

Anomalies galore!

- RD(*) ~ 46(?)
 RK(*): 2.66(A_K);
- g-2...BNL'06 => FNAL expty 3.66 myn lattie progress y
- E': a personal obsession....for a long^3 time=>'cause of the strong conviction that it is super-sensitive to NP EVER

216[PRL 2015] => ~1200 now => ~1400

[2.1 σ (2.9 σ Buras; Nierste) => ??]few more months to new Inesultang BSM scenarios it is important to keep all these [INCLUDING E] + Higgs nadiative stability in mind

Row

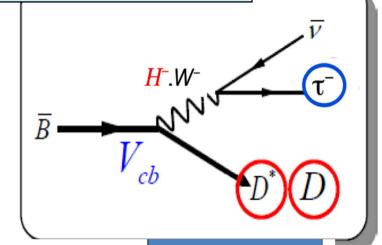
Exclusive $B \to D^{(*)} \tau \nu$

MANNEL FRANCO SEVILLA PLD Thesis

VERA LUTH (BABAR)

FPCP May 2012

(HEFET, CHIMA)



Independent of Vcb!

To test the SM Prediction, we measure

$$R(D) = \frac{\Gamma(\overline{B} \to D\tau \nu)}{\Gamma(\overline{B} \to D\ell \nu)} \qquad R(D^*)$$

$$R(D^*) = \frac{\Gamma(\overline{B} \to D^* \tau \nu)}{\Gamma(\overline{B} \to D^* \ell \nu)}$$

Leptonic τ decays only

Several experimental and theoretical uncertainties cancel in the ratio!

DD accepts and fallic managed accepts.

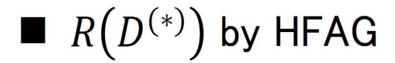
Improving constraints on $\tan \beta/m_H$ using $B \rightarrow D \tau \overline{\nu}$

Ken Kiers* and Amarjit Soni[†]

Department of Physics, Brookhaven National Laboratory, Upton, New York 11973-5000
(Received 12 June 1997)

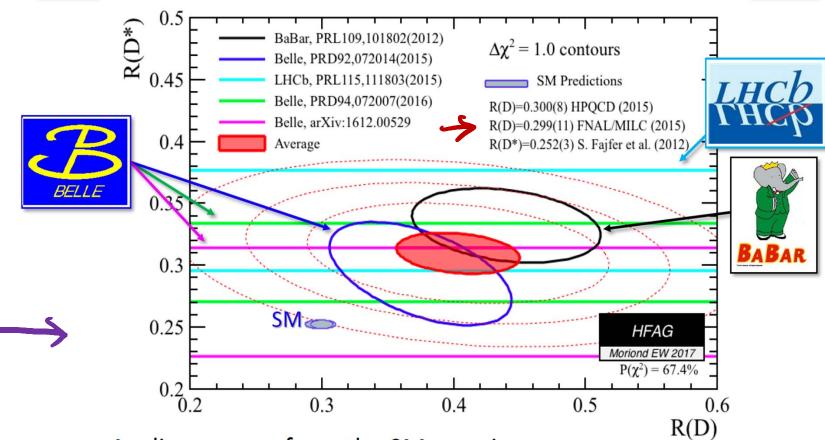
We study the q^2 dependence of the exclusive decay mode $B \to D \tau \overline{\nu}$ in type-II two Higgs doublet models (2HDM's) and show that this mode may be used to put stringent bounds on $\tan \beta/m_H$. There are currently rather large theoretical uncertainties in the q^2 distribution, but these may be significantly reduced by future measurements of the analogous distribution for $B \to D(e,\mu)\overline{\nu}$. We estimate that this reduction in the theoretical uncertainties would eventually (i.e., with sufficient data) allow one to push the upper bound on $\tan \beta/m_H$ down to about 0.06 GeV⁻¹. This would represent an improvement on the current bound by about a factor of 7. We

=) Follover my Vierste et alj fajfer et al 12

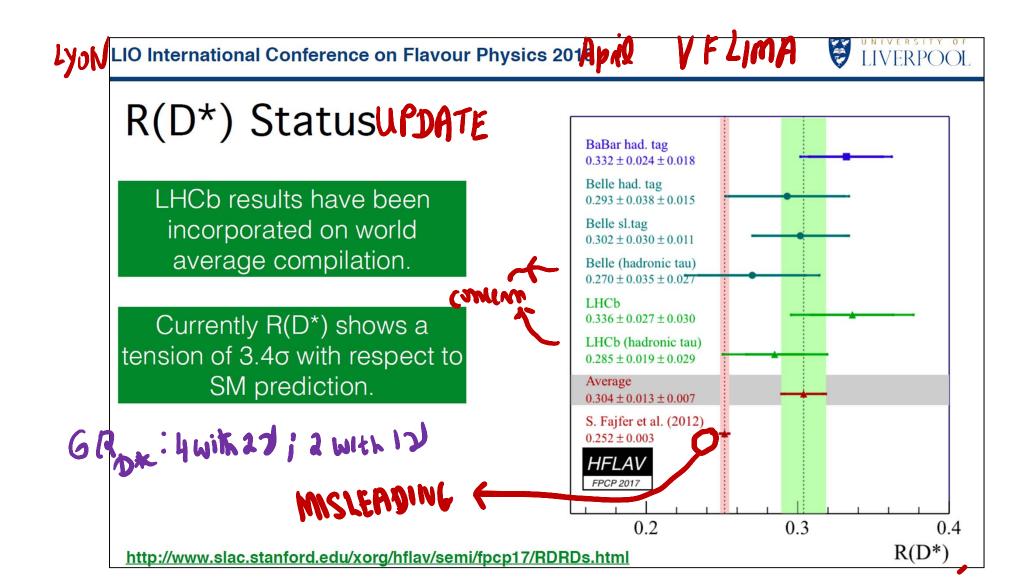


Hirose [BELLE]@EW MORIOND Mar. 2017





- \sim 4 σ discrepancy from the SM remains
 - All the experiments show the larger $Rig(D^{(*)}ig)$ than the SM
- More precise measurements at Belle II and LHCb are essential

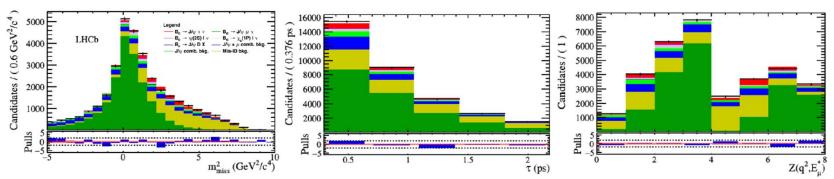


 $B_c
ightarrow J/\psi au
on 2016 Greg Ciezarek,$

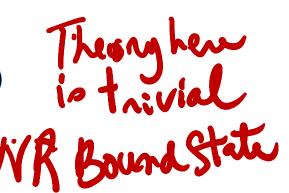
Bc



on behalf of the LHCb collaboration



- $R_{J/\psi} \equiv B_c \rightarrow J/\psi \, \tau \nu/B_c \rightarrow J/\psi \, \mu \nu$
- Measured using very similar techniques to $\mathcal{R}(D^*)$, on run 1 data
- $R_{J/\psi} = 0.71 \pm 0.17 \pm 0.18$
 - $\sim 2\sigma$ from SM
 - But nearly as far from consistency with $\mathcal{R}(D^*)$
- LHCb-PAPER-2017-035(Run 1 data)



Concerns on SM-theory

- Good news is that lattice[FERMIL-MILC] study largely confirms pheno calculations for R_D [our RBC-UKQCD, Witzel et al needs some more time]
- For B=>D* no complete lattice study so far; 4 rather than 2 FF, so , from the lattice perspective, anticipate larger errors than for B=>D...Another ~6 months to complete
- Therefore, O(1%) errors in RD* (and in fact smaller than in RD) are difficult to understand; lattice results should come in some months
- HFAG should update the SM-theory with more realistic errors otherwise their fig is bit misleading つれっているった・のいち
- Meantime recent phenomenological study of Bernlochner, Ligeti, Papucci and Robinson, 1703.05330 [and even more recently...is/are very timely and greatly appreciated. BICI, Cambiant Schackt: 259±.009

 AISMPT NANDI HOURA

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Lepton universality tests

In the SM, ratios

LHCG introduced such v wer



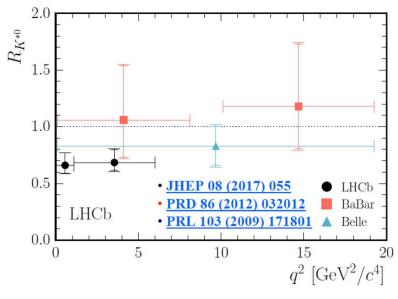
$$R_{\rm K} = \frac{\int d\Gamma[B^+ \to K^+ \mu^+ \mu^-]/dq^2 \cdot dq^2}{\int d\Gamma[B^+ \to K^+ e^+ e^-]/dq^2 \cdot dq^2}$$

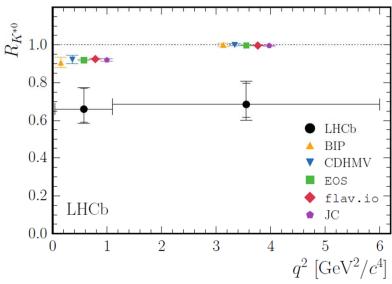


only differ from unity by phase space — the dominant SM processes couple equally to the different lepton flavours.

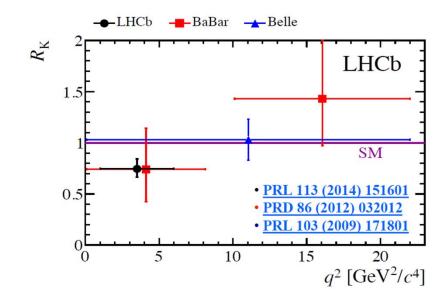
- Theoretically clean since hadronic uncertainties cancel in the ratio.
- Experimentally challenging due to differences in muon/electron reconstruction (in particular Bremsstrahlung from the electrons).
 - → Take double ratios with $B \rightarrow J/\psi X$ decays to cancel possible sources of systematic uncertainty.
 - → Correct for migration of events in q² due to FSR/Bremsstrahlung using MC (with PHOTOS).

Lepton Universality results





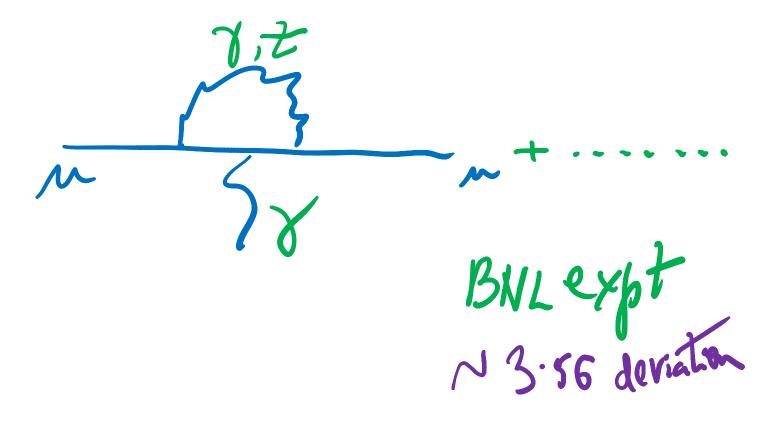
RKX shows similar results as Ru



 R_K : Central- q^2 : 2.6 σ from SM

 R_{K^*} : Low- q^2 : 2.1-2.3 σ from SM

 R_{K^*} : Central- q^2 : 2.4-2.5 σ from SM



POSSIBLE CONNECTION TO G-2

MUON MAY NOT BE JUST A HEAVY ELECTRON: KILE, KOBACH AND AS



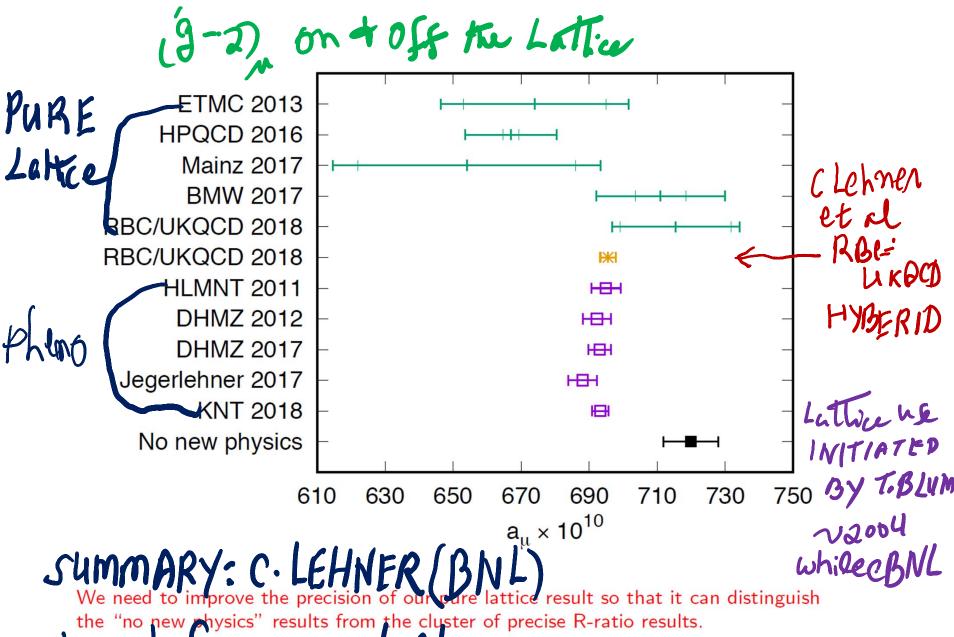
Table 1

Constraints on lepton-flavor violating and conserving processes. For the last four observables, the experimental null results are given in terms of a dimension-6 operator, suppressed by two orders of Λ , which can be interpreted as the nominal scale of new physics.

| | Observable | Limit |
|--------------|-----------------------------------|--|
| | $Br(\mu \rightarrow 3e)$ | $< 1.0 \times 10^{-12} [1]$ |
| 16 P | $Br(\mu \to e \gamma)$ | $< 5.7 \times 10^{-13} [1]$ |
| Istgennot | $Br(\tau \to 3e)$ | $< 2.7 \times 10^{-8} [1]$ |
| sensitive to | $Br(\tau \to e^- \mu^+ \mu^-)$ | $< 2.7 \times 10^{-8} [1]$ |
| anative U | $Br(\tau \to e^+ \mu^- \mu^-)$ | $< 1.7 \times 10^{-8} [1]$ |
| TCW2 In A | $Br(\tau \to \mu^- e^+ e^-)$ | $< 1.8 \times 10^{-8} [1]$ |
| . (P | $Br(\tau \to \mu^+ e^- e^-)$ | $< 1.5 \times 10^{-8} [1]$ |
| NI | $Br(\tau \to 3\mu)$ | $< 2.1 \times 10^{-8} [1]$ |
| 4 | $Br(\tau \to \mu \gamma)$ | $< 4.4 \times 10^{-8}$ [1] |
| | $Br(\tau \to e \gamma)$ | $< 3.3 \times 10^{-8} [1]$ |
| (9-2) | μ – e conversion | $\Lambda \gtrsim 10^3 \text{ TeV [5]}$ |
| Co. M | $e^+e^- \rightarrow e^+e^-$ | $\Lambda \gtrsim 5 \text{ TeV } [3]$ |
| | $e^+e^- \rightarrow \mu^+\mu^-$ | $\Lambda \gtrsim 5 \text{ TeV } [3]$ |
| | $e^+e^- \rightarrow \tau^+\tau^-$ | $\Lambda \gtrsim 4 \text{ TeV } [3]$ |
| WV I WM | podspace 1+0 | |

PKD2015

QUICK UPDATE ON G-2



the "no new physics" results from the cluster of precise R-ratio results.

LUMCh Semimore 03 09 8

Pheno 2018; A Soni (BNL-HET)

Bottom line

- NP or not depends critically not just on precise experiment but also reliable SM prediction from the lattice become mandatory
- Experiment + Lattice M.E. has the last word....[of course should be stressed that the lattice calculations often require sophisticated and demanding input from perturbation theory]
- Experimental results often attained at huge cost can be used effectively, iff commensurate theory predictions are available.....mantra for past several decades

A.S. in Proceedings of Lattice '85 (FSU)..1st Lattice meeting ever attended

The matrix elements of some penguin operators control in the standard model another CP violation parameter, namely ϵ'/ϵ . $^{6,8)}$ Indeed efforts are now underway for an improved measurement of this important parameter. 10) In the absence of a reliable calculation for these parameters, the experimental measurements, often achieved at tremendous effort, cannot be used effectively for constraining the theory. It is therefore clearly important to see how far one can go with MC techniques in alleviating this old but very difficult

With C. Bernard [UCLA]

o 2018; A Soni (BNL-HET)

IV: ε' / ε: Direct CPV EXPERIMENTAL

$$\eta_{+-} = |\eta_{+-}| e^{i\phi_{+-}} = \frac{A(K_L \to \pi^+ \pi^-)}{A(K_S \to \pi^+ \pi^-)}$$
$$\eta_{00} = |\eta_{00}| e^{i\phi_{00}} = \frac{A(K_L \to \pi^0 \pi^0)}{A(K_S \to \pi^0 \pi^0)}$$

$$\eta_{+-} = \epsilon + \epsilon', \quad \eta_{00} = \epsilon - 2\epsilon'$$
 \vdots
 \vdots

Aftenderades of exptal effort CCERN and FNAL ~ 2000 Re E'/e = 16-6+23 × 104

ton Simplicity: 1St Strategy via ChPT

PHYSICAL REVIEW D

VOLUME 32, NUMBER 9

1 NOVEMBER 1985

Application of chiral perturbation theory to $K \rightarrow 2\pi$ decays

Claude Bernard, Terrence Draper,* and A. Soni
Department of Physics, University of California, Los Angeles, California 90024

H. David Politzer and Mark B. Wise

Department of Physics, California Institute of Technology, Pasadena, California 91125 (Received 3 December 1984)

Chiral perturbation theory is applied to the decay $K \rightarrow 2\pi$. It is shown that, to quadratic order in meson masses, the amplitude for $K \rightarrow 2\pi$ can be written in terms of the unphysical amplitudes $K \rightarrow \pi$ and $K \rightarrow 0$, where 0 is the vacuum. One may then hope to calculate these two simpler amplitudes with lattice Monte Carlo techniques, and thereby gain understanding of the $\Delta I = \frac{1}{2}$ rule in K decay. The reason for the presence of the $K \rightarrow 0$ amplitude is explained: it serves to cancel off unwanted renormalization contributions to $K \rightarrow \pi$. We make a rough test of the practicability of these ideas in Monte Carlo studies. We also describe a method for evaluating meson decay constants which does not require a determination of the quark masses.

InspiredI.P. by papers of Shamir [+Furman] + discussions with Creutz

QCD with domain wall quarks

T. Blum* and A. Soni[†]

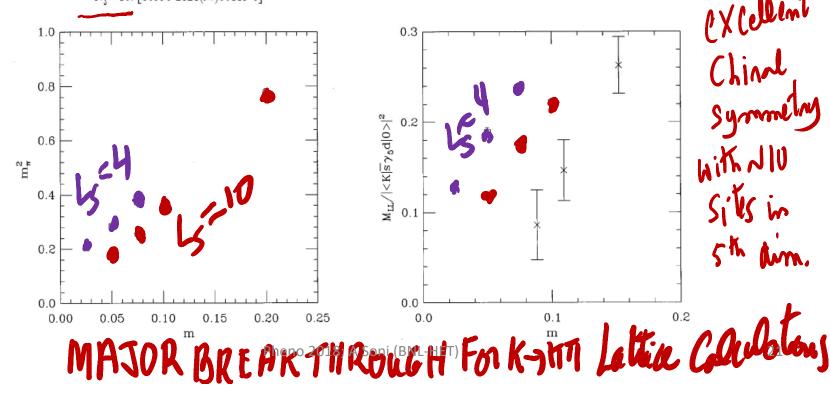
Department of Physics, Brookhaven National Laboratory, Upton, New York 11973

(Received 27 November 1996)

1St Simulation with DUQ

We present lattice calculations in QCD using Shamir's variant of Kaplan fermions which retain the continuum $SU(N)_L \times SU(N)_R$ chiral symmetry on the lattice in the limit of an infinite extra dimension. In particular, we show that the pion mass and the four quark matrix element related to $K_0 - \overline{K_0}$ mixing have the expected behavior in the chiral limit, even on lattices with modest extent in the extra dimension, e.g., $N_c = 10$. [S0556-2821(97)00113-6]

196-47



K->2T ChPT

With DWG in Quench

PHYSICAL REVIEW D 68, 114506 (2003)

121 officials

P BD26m,8

bik DWQ

RBC:

Founding man

Blum, AS

~ 198

Kaon matrix elements and CP violation from quenched lattice QCD: The 3-flavor case

T. Blum, P. Chen, N. Christ, C. Cristian, C. Dawson, G. Fleming, R. Mawhinney, S. Ohta, G. Siegert, A. Soni, P. Vranas, M. Wingate, R. L. Wu, and Y. Zhestkov

¹RIKEN-BNL Research Center, Brookhaven National Laboratory, Upton, New York 11973, USA
 ²Physics Department, Columbia University, New York, New York 10027, USA
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5IBM Research, Yorktown Heights, New York 10598, USA (Received 19 July 2002; published 30 December 2003)

We report the results of a calculation of the $K \rightarrow \pi\pi$ matrix elements relevant for the $\Delta I = 1/2$ rule and ϵ'/ϵ in quenched lattice QCD using domain wall fermions at a fixed lattice spacing a^{-1} ~2 GeV. Working in the three-quark effective theory, where only the u, d, and s quarks enter and which is known perturbatively to next-to-leading order, we calculate the lattice $K \rightarrow \pi$ and $K \rightarrow |0\rangle$ matrix elements of dimension six, fourfermion operators. Through lowest order chiral perturbation theory these yield $K \rightarrow \pi\pi$ matrix elements, which we then normalize to continuum values through a nonperturbative renormalization technique. For the ratio of isospin amplitudes $|A_0|/|A_2|$ we find a value of 25.3 \pm 1.8 (statistical anorom), empared to the value of 22.2, with individual isospin amplitudes 10%-20% below the experimental values. For ϵ'/ϵ , using known central values for standard model parameters, we calculate \ 4.0\pm 2.3\)\times 10\pm (statistical error only compared to the current experimental average of (17.2±1.8)×10⁻⁴. Because we find a large cancellation between the I=0 and I=2 contributions to ϵ'/ϵ , the result may be very sensitive to the approximations employed. Among these are the use of quenched QCD, lowest order chiral perturbation theory, and continuum perturbation theory below 1.3 GeV. We also calculate the kaon B parameter B_K and find $B_{K,MS}(2 \text{ GeV})$ =0.532(11). Although currently unable to give a reliable systematic error, we have control over statistical errors and more simulations will yield information about the effects of the approximations on this firstprinciples determination of these important quantities,

BC Collaboration

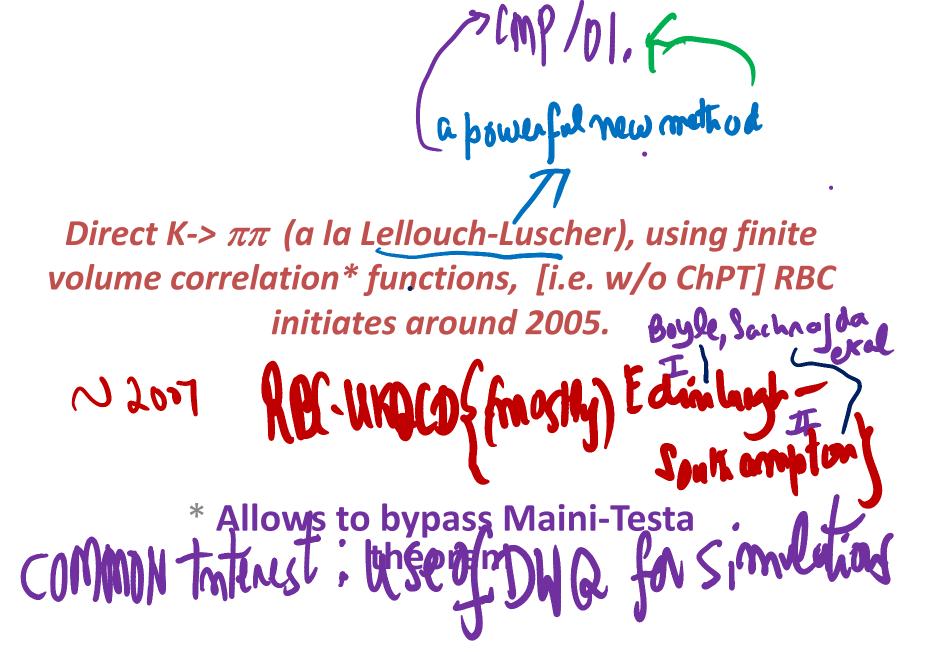
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Now Nachip lagged

Ist laye Sale Simble With DUO

PRO103



DIRECT K-> m

Results for ε'

• Using Re(A) and Re(A) from experiment $Im(A_0)$ and $Im(A_2)$ and the phase shifts,

and our lattice values for

USING 216 independent measurements

 $\operatorname{Re}\left(\frac{\varepsilon'}{\varepsilon}\right) = \operatorname{Re}\left\{\frac{i\omega e^{i(\delta_2 - \delta_0)}}{\sqrt{2}\varepsilon} \left[\frac{\operatorname{Im} A_2}{\operatorname{Re} A_2} - \frac{\operatorname{Im} A_0}{\operatorname{Re} A_0}\right]\right\}$

LARGE CANCELLATION!! (80-85%)

RBC-UKQCD PRL'15 EDITOR'S CHOICE

 $= 1.38(5.15)(4.43) \times 10^{-4},$ (this work) $16.6(2.3) \times 10^{-4}$ (experiment)

Bearing in mind the largish errors in this first calculation, we interpret that our result are consistent with experiment at ~2σ level

w= Ne 1+2 ~ 0.145

Computed ReA2 excellent agreement with expt Computed ReA0 good agreement with expt Offered an "explanation" of the Delta I=1/2 Enhancement [c later]

BURAS et al useow [MF => effect is ~ 2.95" Niente · Pheno 2018; A Soni (BNL-HET)

Generation of New gaye empigs For past 7, 3 years

SUPERCOMPUTERS OVER 3 CONTINENTS!

Progress in the calculation of arepsilon' on the lattice

C. Kelly

| Resource | Million BG/Q equiv core-hours | Independent cfgs. |
|-------------------------------|-------------------------------|-------------------|
| USQCD (BNL 512 BG/Q nodes) | 50 | 220 |
| RBRC/BNL (BNL 512 BG/Q nodes) | 17 | 50 |
| UKQCD (DiRAC 512 BG/Q nodes) | 17 | 50 |
| NCSA (Blue Waters) | 108 | 380 |
| KEK (KEKSC 512 BG/Q nodes) | 74 | 296 |
| Total | 266 | 996 |

Table 1: A breakdown of the various resources we intend to utilize. Note that we require 4 molecular dynamics time units per independent configuration.

4 hilf statems

By mw ~ 1300 m

mersumers

MOTHER of all (lattice) calculations to date: A Personal Perspective

- Calculation K=> $\pi\pi$ & ϵ' were the reasons I went into lattice over 1/3 of a century ago!
- 9 + (3 new) PhD thesis: Terry Draper (UCLA'84), George Hockney(UCLA'86), Cristian Calin (Columbia=CU'01), Jack Laiho(Princeton'04), Sam Li(CU'06), Matthew Lightman(CU'09), Elaine Goode(Southampton'10), Qi Liu(CU'12), Daiqian Zhang(CU'15)+ [new ones starting from CU, U Conn and Southampton] + many PD's & junior facs.. obstacles & challenges (and of course "mistakes"!) ad infinitum.....

•

WHY FOCUS with SUCH intense DETERMINATION All these many many years?

UNDERLYING REALIZATION E': MOST LIKELY A GEM IN SEARCH OF NEW PHENOMENA

Contrarian/Complementary view

- flavor physics is actually hanging by perhaps the weakest link i.e. a single CP-phase endowed by the 3g –SM.
- In many ways this is a contrarian (or complementary)
 point of view, in sharp contrast to the overwhelming
 majority following the naturalness lamp post via Higgs
 radiative stability.

 E' due to its miniscule value, esp because it results from unnatural large cancellations seemed clearly highly vulnerable...The mantra being followed for a very very long time



IF YOU BUILD IT THEY WILL COME

If there is new physics around below ~5 TeV, there is an excellent chance that ε' will find it!

[of course requires accurate theory calculation... RBC-UKQCD plans for X5 in stat and appreciable improvements in systematic in ~2 years]

9/18/2017 LmC2017; SIEGEN; HET-BNL;soni

ALTMANNSHOFFR, Devtas 1704.06659 + Sey WIP

MODEL INDEPENDENT IMPLICATIONS OF RD(*) ANOMALIES FOR [LHC] COLLIDER EXPERIMENTS

- In a nut-shell B-experiments seem to find anomalous behavior in the underlying b=>c tau nu
- This necessarily [by XSym] implies there should be analogous anomaly in g + c => b tau nu...=>pp => b tau nu
- Thus it immediately leads to inescapable search channels for possible NP at the high energy frontier for ATLAS & CMS and these are urgently urged



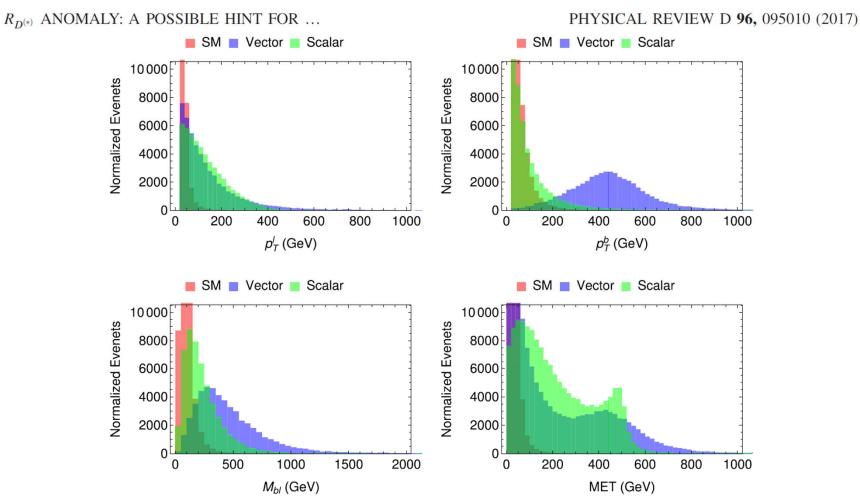


FIG. 1. Normalized kinematic distributions for the $pp \to b\tau\nu \to b\ell + \not\!\!E_T$ signal and background.

EXPECT DISTINCTIVE NP CONTRIBUTIONS IN COLLIDERS

ANOMALY: POSSIBLY A HINT FOR (NATURAL) SUSY-WITH RPV

See Altmannshofer, Dev +AS, arXiv:1704.06659 + WIP

- ASSUMING the anomaly is REAL & HERE TO STAY [BIG ASSUMPTION due to caveats mentioned]
- Anomaly involves simple tree-level semi-leptonic decays
- Also b => tau (3rd family)
- Speculate: May be related to Higgs naturalness
- Seek minimal solution: perhaps 3rd family super-partners(a lot) lighter than other 2 gens > proton decay concerns may not be relevant=> RPV ["natural" SUSY]
- RPV natural setting for LUV ...can accommodate g-2 and eps' if needs be
- Collider signals tend to get a lot harder than (usual-RPC) SUSY

RPV3 preserves gange coupling unification i mespecture of ## of effective gens. 1, 2 003.

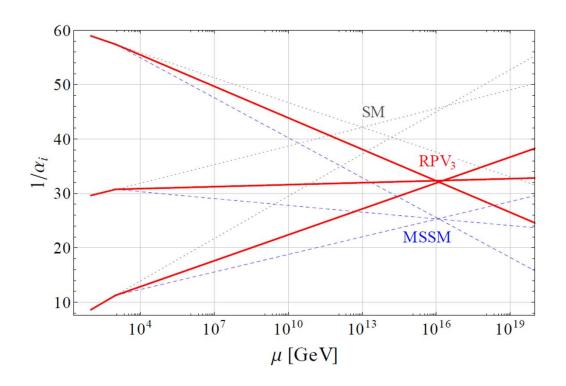


FIG. 2. RG evolution of the gauge couplings in the SM, MSSM and with partial supersymmetrization.

Unification scale astoys some, only value of couplings high

For phono relayant tems:

ADS'PRD 2017

$$\mathcal{L} = \lambda'_{ijk} \left[\tilde{\nu}_{iL} \bar{d}_{kR} d_{jL} + \tilde{d}_{jL} \bar{d}_{kR} \nu_{iL} + \tilde{d}_{kR}^* \bar{\nu}_{iL}^c d_{jL} \right]$$
$$-\tilde{e}_{iL} \bar{d}_{kR} u_{jL} - \tilde{u}_{jL} \bar{d}_{kR} e_{iL} - \tilde{d}_{kR}^* \bar{e}_{iL}^c u_{jL} \right] + \text{H.c.}$$

$$\mathcal{L}_{\text{eff}} \supset \frac{\lambda'_{ijk} \lambda'^*_{mnk}}{2m_{\tilde{d}_{kR}}^2} \left[\bar{\nu}_{mL} \gamma^{\mu} \nu_{iL} \bar{d}_{nL} \gamma_{\mu} d_{jL} - \nu_{mL} \gamma^{\mu} e_{iL} \bar{d}_{nL} \gamma_{\mu} \left(V^{\dagger}_{\text{CKM}} u_L \right)_j + \text{h.c.} \right] - \frac{\lambda'_{ijk} \lambda'^*_{mjn}}{2m_{\tilde{u}_{jL}}^2} \bar{e}_{mL} \gamma^{\mu} e_{iL} \bar{d}_{kR} \gamma_{\mu} d_{nR} ,$$

RPV3 interaction

- DIM-6

For addressing RK(*) in RPV, see e.g. Das et al , 1705.09188

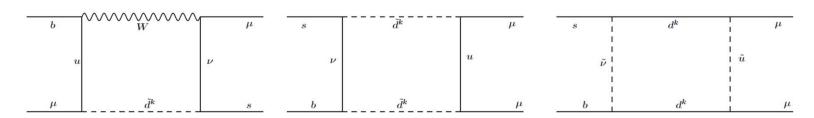


FIG. 1: Representative diagrams for $b \to s\mu^+\mu^-$ transition in R-parity violating interactions.

g-2 with RPV has a long history, see, e.g.Kim, Kyae and Lee, PLB 2001

We [ALTHANNISHOFER+DEV+AS] acexamining+up date in light of current flavor amomalies WORK IN Progress

MULTITUDE OF CONSTRAINTS ON OUR RPV3 [AND ALSO OTHER PROPOSED IDEAS]

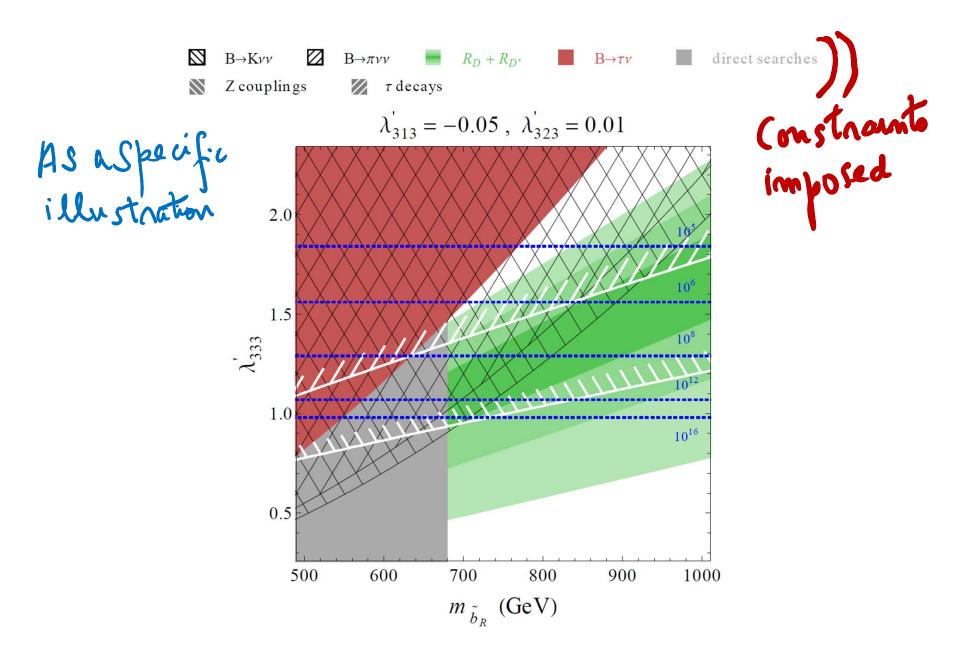
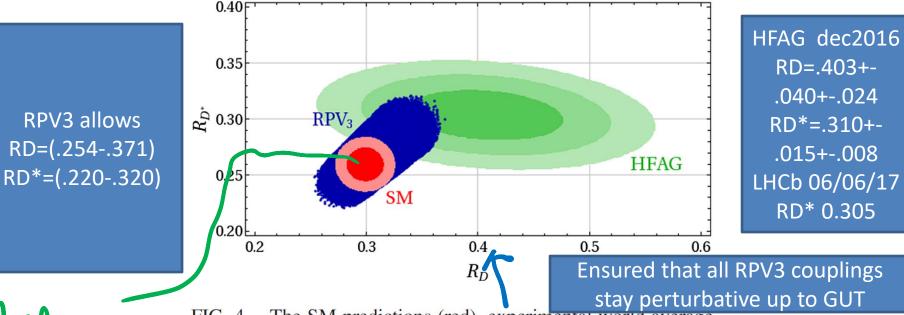


FIG. 3. RPV parameter space satisfying the $R_{D^{(*)}}$ anomaly and other relevant constraints. A Soni (BNL-HET)





Mone Redistic

FIG. 4. The SM predictions (red), experimental world average (green), and accessible values in our RPV-SUSY scenario (blue) in the R_D vs. R_{D^*} plane. For the SM, bearing in mind recent works [17,20,22] we are taking $(R_D^{\rm SM}, R_{D^*}^{\rm SM}) = (0.299 \pm 0.011, 0.260 \pm 0.010)$.

all constraints......RPV(blue) region obtained by scanning with sbottom mass 680-1000Gev, 0<λ333<2;|λ323|<0.1;|λ313|<0.3

••••

Summary and Outlook..[p1 of 2]

- Hints of LUV [from 3 B experiments] in sl B decays claimed to be around 4 sigma and also of FCNC ~ 2.5 sigma in B=K(*) Il are interesting but not yet compelling.
- For RD(*) an important experimental concern is that systematics in tau =>I nu nu may not be in good control; this concern is accentuated as both Belle and LHCb measurements of RD* with tau => hadron + nu appear consistent with SM within ~ 1 sigma.
- From lattice B=>D* semi lep form factors and RD* are urgently needed; likely ~6 months.
- For RK(*) [theory is irrelevant] need confirmation from another expt...Even for Belle II this will take time because of the small Br $^{\sim}O(10^{\circ}-6)$. LHCb needs to study R $_{\phi}$ through Bs=> ϕ $\mu\mu$ AND Bs=>phi e e; also via B-baryons
- More data from LHCb from Run 2 < 1 year should help and further ~2 years down Belle II should start to help more....
- Belle II with X50 Belle lumi and LHCb upgrade have a lot of potential for searches related to these & many others and likely to be a game changer in search of new phenomena.
- It may well be that BNL's observed g-2 signals of possible NP were just a precursor to these observations of LUV in B decays.
- Lattice progress in g-2 by RBC-UKQCD as well as global efforts are impressive ...But needs
 to reduce errors further by ~X4...Expect next reduction X2 in a year or so
- ε': RBC-UKQCD should be able to appreciably improve their 2015 result of ~2.1 sigma tension, in <6 months
- There is now an exciting and may be even a revolutionary possibility that one or more of these avenues will show significant departure from SIVI in the next few years

XTRAS

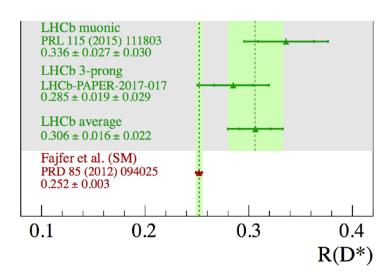
Bernlochner, Ligeti, Papucci and Robinson, 1703.05330

| Scenario | • R(D) | $R(D^*)$ | Correlation |
|-------------------------------------|-------------------|-------------------|-----------------------|
| $\mathcal{L}_{w=1}$ | 0.292 ± 0.005 | 0.255 ± 0.005 | 41% SM hediction |
| $L_{w=1}+SR$ | 0.291 ± 0.005 | 0.255 ± 0.003 | 57% |
| NoL | 0.273 ± 0.016 | 0.250 ± 0.006 | 49% |
| NoL+SR | 0.295 ± 0.007 | 0.255 ± 0.004 | 43% |
| $L_{w\geq 1}$ | 0.298 ± 0.003 | 0.261 ± 0.004 | 19% tok |
| $L_{w\geq 1}+SR$ | 0.299 ± 0.00 a | 0.257 ± 0.003 | 19% We took 44% RDX = |
| $\text{th:L}_{w\geq 1} + \text{SR}$ | 0.306 ± 0.005 | 0.256 ± 0.004 | |
| Data [9] | 0.403 ± 0.047 | 0.310 ± 0.017 | -23% / 0.251T.WS |
| Refs. [48, 52, 54] | 0.300 ± 0.008 | _ | |
| Ref. [53] | 0.299 ± 0.003 | _ | Fajfer, Kamenik, |
| Ref. [34] | _ | 0.252 ± 0.003 | Nisandzic, PRD'12 |

TABLE IV. The R(D) and $R(D^*)$ predictions for our fit scenarios, the world average of the data, and other theory predictions. The fit scenarios are described in the text and in Table I. The bold numbers are our most precise predictions.

Conclusions

- We have measured the ratio $K_{had}(D^*)=BR(B^0\to D^*-\tau v)/BR(B^0\to D^*-3\pi)$ using the $3\pi(\pi^0)$ hadronic decay of the τ lepton.
- The result regarding R(D*) is compatible with all other measurements and with the SM, having the smallest statistical error.
- This analysis was made possible due to the unique LHCb capabilities for separating secondary and tertiary vertices with excellent resolution.



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Using BR(B⁰ \to D* $\mu\nu$) = (4.93 ± 0.11)% [PDG-2016] we measure:

$$R(D^*) = 0.285 \pm 0.019(stat) \pm 0.025(syst) \pm 0.014(ext)$$

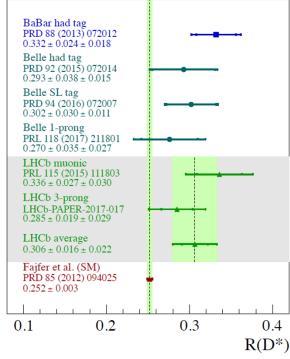
In combination with the muonic LHCb measurement:

$$R(D^*) = 0.336 \pm 0.027 \pm 0.030$$

the LHCb average is:

- $R_{LHCh}(D^*) = 0.306 \pm 0.016 \pm 0.022$
- 2.1σ above the SM.
- Naïve new WA:
 - $R(D^*) = 0.305 \pm 0.015$
 - 3.4σ above the SM.
- Naïve $R(D)/R(D^*)$ combination at 4.1σ from SM.





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Table 13-6. Model-dependent effects of new physics in various processes.

| | CP Viola | tion | | D^0 – \overline{D}^0 |
|----------------------|-----------------------------------|------------------|---|--------------------------|
| Model | $B_d^0 - \overline{B}_d^0$ Mixing | Decay Ampl. | Rare Decays | Mixing |
| MSSM | O(20%) SM | No Effect | $B \to X_s \gamma$ – yes | No Effect |
| | Same Phase | | $B 	o X_s l^+ l^-$ – no | |
| SUSY – Alignment | O(20%) SM | $\mathcal{O}(1)$ | Small Effect | Big Effect |
| | New Phases | | | |
| SUSY - | $\mathcal{O}(20\%)$ SM | $\mathcal{O}(1)$ | No Effect | No Effect |
| Approx. Universality | New Phases | | | |
| R-Parity Violation | Can Do | Everything | Except Make | Coffee |
| MHDM | \sim SM/New Phases | Suppressed | $B \to X_s \gamma, B \to X_s \tau \tau$ | Big Effect |
| 2HDM | \sim SM/Same Phase | Suppressed | $B \to X_s \gamma$ | No Effect |
| Quark Singlets | Yes/New Phases | Yes | Saturates Limits | Q = 2/3 |
| Fourth Generation | ~ SM/New Phases | Yes | Saturates Limits | Big Effect |
| $LRM - V_L = V_R$ | No Effect | No Effect | $B \to X_s \gamma, B \to X_s l^+ l^-$ | No Effect |
| $-V_L \neq V_R$ | Big/New Phases | Yes | $B \to X_s \gamma, B \to X_s l^+ l^-$ | No Effect |
| DEWSB | Big/Same Phase | No Effect | $B \to X_s \ell \ell, B \to X - s \nu \overline{\nu}$ | Big Effect |

though in many cases further data may limit the available parameter space. In the more exciting eventuality that the results are not consistent with Standard Model predictions, the full pattern of the discrepancies both in rare decays and in CP-violating effects will help point to the preferred extension, and possibly rule out others. In either case there is much to be learned.

Results



BELLECEPS July 2017

- ullet $\mathcal{R}(D^*)$ can be calculated as before from extracted yields
- Polarisation from forward/backward asymmetry



$$\frac{\epsilon_{norm}}{\epsilon_{sig}} = 0.97 \pm 0.02 \ (B^{\pm}, \ \tau \to \pi \nu)$$

$$= 1.21 \pm 0.03 \ (B^{0}, \ \tau \to \rho \nu)$$

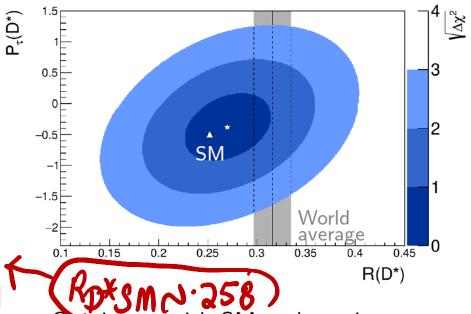
$$= 3.42 \pm 0.07 \ (B^{\pm}, \ \tau \to \rho \nu)$$

$$= 3.83 \pm 0.12 \ (B^{0}, \ \tau \to \rho \nu)$$

Result

$$\mathcal{R}(D^*) = 0.270 \pm 0.035^{+0.028}_{-0.025}$$

 $P_{\tau}(D^*) = -0.38 \pm 0.51^{+0.21}_{-0.16}$



- Consistent with SM and previous measurements!
- Error can be reduced in Belle II

MOTHER of all (lattice) calculations to date: A Personal Perspective

- Calculation $K=>\pi\pi\&\epsilon'$ were the reasons I went into lattice over 1/3 of a century ago!
- 9 + (3 new) PhD thesis: Terry Draper (UCLA'84), George
 Hockney(UCLA'86), Cristian Calin (Columbia=CU'01), Jack
 Laiho(Princeton'04), Sam Li(CU'06), Matthew Lightman(CU'09),
 Elaine Goode(Southampton'10), Qi Liu(CU'12), Daiqian
 Zhang(CU'15)+ [new ones starting from CU, U Conn and
 Southampton] + many PD's & junior facs.. obstacles & challenges
 (and of course "mistakes"!) ad infinitum.....
- Started with CBernard (Wilson F); for this physics Chiral symm on the lattice is a pre-requisite [off-shoot B-physics] => on to DWF (with T Blum)=> RBC with ChPT + quenched => huge quench pathlogies=full QCD is mandatory for this physics; full QCD + ChPT=> large chiral corrections => RBC-UKQCD direct K=> 2 π a la Lellouch- Luscher @ threshold=> @physical kinematics......

Anomalies galore!

- RD(*)
- RK(*)
- •
- g -2
- •
- epsilon': The meaning of life
 216[PRL 2015] => ~720 now => ~1200

$$[2.1\sigma (2.9\sigma?) => ????]$$
some months

Swamonths

eps' and higgs stability are bonus For Delta M_Bs NNLO EW corr may be appricaible?

- Semi-leptonic B-decays r claimed to indicate ~4.1 sigma deviation from SM
- ATLAS, CMS ought to vigorously search for BSM in : b τ v and in t τ
- Expt BG from higher D** etc resonances a concern and should b measured; tau detection via hadronic modes should be given very high priority as its much

less susceptible to D** contaminations

- More independent theory effort on and off lattice for determination of SM value for RD* are urgently needed
- More info from expts on R(D), R(D*), R(π), R(ρ), analogous Bs, B-baryon, B=> τ v are all urgently needed
- Also RD from LHCb as well as Belle would be helpful [since in this case theory is very solid]; BELLE-II and LHCb-upgrades would of
 course help a lot
- RPV-SUSY effectively involving 3rd gen is economical, minimal and natural and may be an interesting origin of the anomaly [if it persists!]
- => classic large missing energy hunt for SUSY not relevant for that scenario
- => many RPV signatures tend to become rather challenging
- => our version gives new interesting avenues in b τ v; t τfinal states
- More studies in progress (inc e,g. RK(*), Bs=>μ μ and much more): see ADS' II

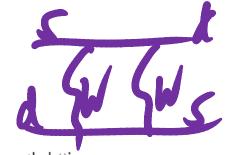
28 39. Statistics



Table 39.1: Area of the tails α outside $\pm \delta$ from the mean of a Gaus distribution.

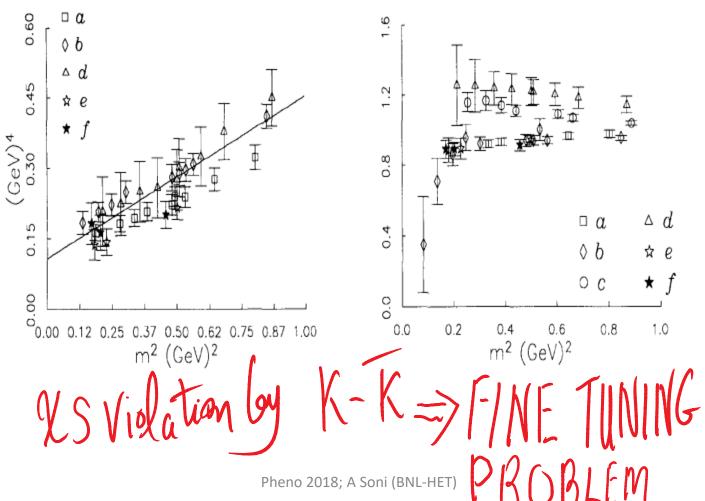
| | α | δ | | α | δ |
|---|-----------------------|-----------|---|-----------|--------------|
| | 0.3173 | 1σ | , | 0.2 | 1.28σ |
| | 4.55×10^{-2} | 2σ | | 0.1 | 1.64σ |
| | 2.7×10^{-3} | 3σ | _ | 0.05 | 1.96σ |
| C | 6.3×10^{-5} | 4σ | | 0.01 | 2.58σ |
| | 5.7×10^{-7} | 5σ | | 0.001 | 3.29σ |
| | 2.0×10^{-9} | 6σ | | 10^{-4} | 3.89σ |

(K||Sm.1)/K)



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C. Bernard, A. Soni / Weak matrix elements on the lattice



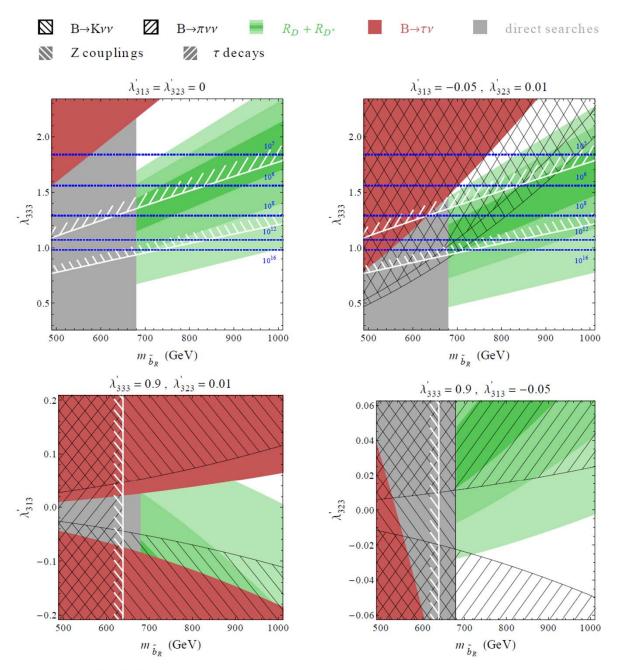


FIG. 3. RPV parameter space satisfying the $R_{D(*)}$ anomaly and other relevant constraints. Pheno 2018; A Soni (BNL-HET)

1 ..

Concern on Experiments

- Leptonic decays: $\tau => \mu \nu \nu ... total 3 v's in event$
- Higher D** etc resonances....use of theo models for subtraction of these backgrounds is fraught with danger.....Backgrounds should be measured experimentally for reliable estimate of errors
- Bearing that in mind, it is striking that LHCb new result june 2017: $B=>D^* \tau \nu; \tau =>3\pi+\nu$ is
- consistent with the SM at ~1- σ => heightens anxiety about D**....contaminations in τ => $\mu\nu\nu$
- Furthermore, new Belle result with hadronic tau decay also consistent with SM well within 1 sigma!
- Claimed ~"4 sigma" probably not that solid

Near future outlook

- LHCb has so far only used Run 1 data
- Plenty more data from Run 2 available but needs to be analyzed...may be will get bit of news on this from EWM in < 1 month
- Lattice calculations for slff for B=>D* in <6 months
- Lattice g-2 improved results will continually come perhaps once/[6 months] for next several years....global effort including a lot from our RBC-UKQCD
- Improved lattice results for eps' from our RBC-UKQCD in a few (O(6))months

 $Q_1 = (\bar{s}_{\alpha} d_{\alpha})_L (\bar{u}_{\beta} u_{\beta})_L,$ $Q_2 = (\bar{s}_{\alpha} d_{\beta})_L (\bar{u}_{\beta} u_{\alpha})_L,$

$$Q_2 = (\bar{s}_{\alpha} d_{\beta})_L (\bar{u}_{\beta} u_{\alpha})_L,$$

$$Q_4 = (\bar{s}_{\alpha} d_{\beta})_L \sum_{q=u,d,s} (\bar{q}_{\beta} q_{\alpha})_L$$

$$Q_{3} = (\bar{s}_{\alpha}d_{\alpha})_{L} \sum_{q=u,d,s} (\bar{q}_{\beta}q_{\beta})_{L},$$

$$Q_{4} = (\bar{s}_{\alpha}d_{\beta})_{L} \sum_{q=u,d,s} (\bar{q}_{\beta}q_{\alpha})_{L},$$

$$Q_{5} = (\bar{s}_{\alpha}d_{\alpha})_{L} \sum_{q=u,d,s} (\bar{q}_{\beta}q_{\beta})_{R},$$

$$Q_{6} = (\bar{s}_{\alpha}d_{\beta})_{L} \sum_{q=u,d,s} (\bar{q}_{\beta}q_{\alpha})_{R},$$

$$Q_6 = (\bar{s}_{\alpha} d_{\beta})_L \sum_{q=u,d,s} (\bar{q}_{\beta} q_{\alpha})_R,$$

 $Q_7 = \frac{3}{2} (\bar{s}_{\alpha} d_{\alpha})_L \sum_{q=u,d,s} e_q (\bar{q}_{\beta} q_{\beta})_R,$ $Q_8 = \frac{3}{2} (\bar{s}_{\alpha} d_{\beta})_L \sum_{q=u,d,s} e_q (\bar{q}_{\beta} q_{\alpha})_R,$ $Q_9 = \frac{3}{2} (\bar{s}_{\alpha} d_{\beta})_L \sum_{q=u,d,s} e_q (\bar{q}_{\beta} q_{\alpha})_R,$

$$Q_8 = \frac{3}{2} (\bar{s}_{\alpha} d_{\beta})_L \sum_{q=u,d,s} e_q (\bar{q}_{\beta} q_{\alpha})_R,$$

$$Q_9 = \frac{3}{2} (\bar{s}_{\alpha} d_{\alpha})_L \sum_{q=u,d,s} e_q (\bar{q}_{\beta} q_{\beta})_L,$$

$$Q_{10} = \frac{3}{2} (\bar{s}_{\alpha} d_{\beta})_L \sum_{q=u,d,s} e_q (\bar{q}_{\beta} q_{\alpha})_L,$$

$$Q_{10} = \frac{3}{2} (\bar{s}_{\alpha} d_{\beta})_L \sum_{q=u,d,s} e_q (\bar{q}_{\beta} q_{\alpha})_L,$$

Pheno 2018; A Soni

Summary of Theo Calculations

R(D)=0.300(8) HPQCD (2015) R(D)=0.299(11) FNAL/MILC (2015) 0.299 ± 0.003 BERNLOCHNER et al 2017 0.299 ± 0.003 D. BIGI etal 2017 R(D*)=0.252(3) S. Fajfer et al. (2012) 0.257 ± 0.003 Bernlochner et al $R(D^*) = 0.258^{+9}$

Lattice computation of the decay constants of B and D mesons

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Amariit Soni

Department of Physics, Brookhaven National Laboratory, Upton, New York 11973 (Received 1 July 1993)

PHYSICAL REVIEW D

VOLUME 45, NUMBER 3

1 FEBRUARY 1992

Semileptonic decays on the lattice: The exclusive 0⁻ to 0⁻ case

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Aida X. El-Khadra

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Amarjit Soni

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(Received 21 December 1990)

PHYSICAL REVIEW D. VOLUME 58, 014501

Lattice study of semileptonic decays of charm mesons into vector mesons

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Aida X. El-Khadra

Theory Group, Fermi National Accelerator Laboratory, P.O. Box 500, Batavia, Illinois 60510

Amarjit Soni

Department of Physics, Brookhaven National Laboratory, Upton, New York 11973
(Received 30 September 1991)

We present our lattice calculation of the semileptonic form factors for the decays $D \to K^*$, $D_s \to \phi$, and $D \to \rho$ using Wilson fermions on a $24^3 \times 39$ lattice at $\beta = 6.0$ with 8 quenched configurations. For $D \to K^*$, we find for the ratio of axial form factors $A_2(0)/A_1(0) = 0.70 \pm 0.16 \pm 0.00$. Results for other form factors and ratios are also given.

form factore and ratios are also given

WORKS LEBACY

Pheno 201

SU(3) flavor breaking in hadronic matrix elements for B-B oscillations

C. Bernard

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Department of Physics, Brookhayen National Laboratory, Upton, New York 11973

Received 28 Juniary 1998; published 5 May 1998

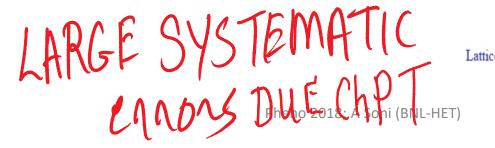
Pheno 2018; A 50111 (DIVE-11E 1)

(Sam)Shu Li, PhD (Sam)shu Li, PhD thesis, Columbia '08

| Quantity | This analysis | Quenched | Experiment |
|---|--------------------------------|------------------------------|-----------------------|
| ReA_0 (GeV) | $4.5(11)(53) \times 10^{-7}$ | $2.96(17) \times 10^{-7}$ | 3.33×10^{-7} |
| ReA_2 (GeV) | $8.57(99)(300) \times 10^{-9}$ | $1.172(53) \times 10^{-8}$ | 1.50×10^{-8} |
| $Im A_0$ (GeV) | $-6.5(18)(77) \times 10^{-11}$ | $-2.35(40) \times 10^{-11}$ | |
| $Im A_2$ (GeV) | $-7.9(16)(39) \times 10^{-13}$ | $-1.264(72) \times 10^{-12}$ | |
| $1/\omega$ | 50(13)(62) | 25.3(1.8) | 22.2 |
| $\operatorname{Re}(\epsilon'/\epsilon)$ | $7.6(68)(256) \times 10^{-4}$ | $-4.0(2.3) \times 10^{-4}$ | 1.65×10^{-3} |



- ChPT approach to $K \to \pi \pi$ faces severe difficulties.
- RBC/UKQCD studying physical $\pi \pi$ final states.
- DWF on coarse lattices and large volumes: 4 → 5 fm?
- Vranas auxiliary determinant (Renfrew talk on Wed.)





N. Christ @LAT08

RBC Sine ~ 98 < Collaboration

- BNL
 - Chulwoo Jung
 - Taku Izubuchi (RBRC)
 - Christoph Lehner
 - Meifeng Lin
- RBR marjit Soni
 - Chris Kelly
 - Tomomi Ishikawa
 - Taichi Kawanai
 - Shigemi Ohta
 - Sergey Syritsyn

- Columbia
 - Ziyuan Bai Xu Feng
 - Norman Christ
 - Luchang Jin
 - Robert Mawhinney
 - Greg McGlynn
 - David Murphy
 - **Daiqian Zhang**
- Connectic

ut

FOUNDING Members of RBC Caled: Blum, Christ, Mawhings

There 2010: A Son: (BALL LIET)

102 Pheno 2018; A Soni (BNL-HET)

UKQCD

Since ~ 07

Collaboration

- Edinburgh
 - Peter Boyle
 - Luigi Del Debbio
 - Julien Frison
 - Jamie Hudspith
 - Richard Kenway
 - Ava Khamseh
 - Brian Pendleton
 - Karthee Sivalingam
 - Oliver Witzel
 - Azusa Yamaguchi
- Plymouth
 - Nicolas Garron
- York (Toronto)
 - Renwick Hudspith

- Southampton
 - Jonathan Flynn
 - Tadeusz Janowski
 - Andreas Juttner
 - Andrew Lawson
 - Edwin Lizarazo
 - Antonin Portelli
 - Chris Sachrajda
 - FrancescoSanfilippo
 - Matthew Spraggs
 - Tobias Tsang
- CERN
 - Marina Marinkovic

While ReA0 and ReA2 and δ 2 agree well with expt a possible difficulty: δ0

The continuum and our lattice determinations of strong phase

difference differs at the ~2
$$\sigma$$
 level:
$$\phi_{\varepsilon'} = \delta_2 - \delta_0 + \frac{\pi}{2} = \begin{cases} (42.3 \pm 1.5)^{\circ} \text{ PDG} & \text{Colongila stal} \\ (54.6 \pm 5.8)^{\circ} & \text{RBC-UKQC} \end{cases}$$
Fortunately, due to the central value of the combination $\delta_2 - \delta_0 + \pi/2 - \phi_{\varepsilon}$ and to the large uncertainties in the determination of the various matrix elements, these two

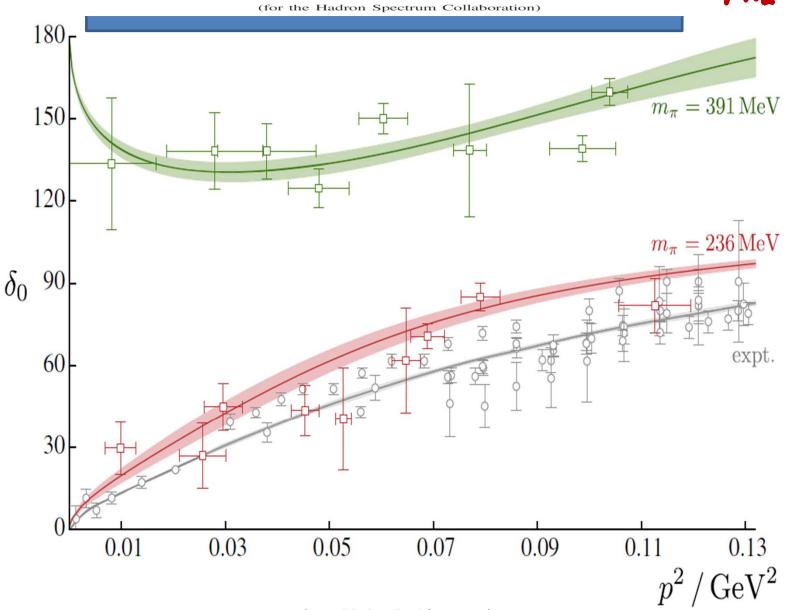
Fortunately, due to the central value of the combination $\delta_2 - \delta_0 + \pi/2 - \phi_\varepsilon$ and to the large uncertainties in the determination of the various matrix elements, these two choices yield almost identical results; for the large uncertainties in the determination of the various matrix elements, these two choices yield almost identical results; for the large uncertainties in the determination of the various matrix elements, these two choices yield almost identical results; for the large uncertainties in the determination of the various matrix elements, these two choices yield almost identical results; for the large uncertainties in the determination of the various matrix elements, these two choices yield almost identical results; for the large uncertainties in the determination of the various matrix elements, these two choices yield almost identical results; for the large uncertainties in the determination of the various matrix elements, the large uncertainties in the determination of the various matrix elements.

Sensitivity of ε' to strong phase(s)

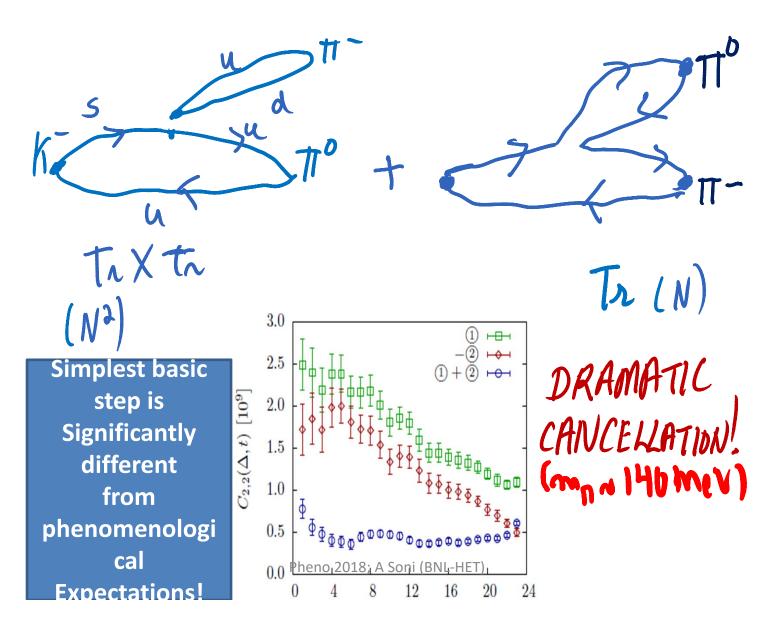
Isoscalar $\pi\pi$ Scattering and the σ Meson Resonance from QCD

Raul A. Briceño, ^{1,*} Jozef J. Dudek, ^{1,2,†} Robert G. Edwards, ^{1,‡} and David J. Wilson ^{3,§}





Dissecting (the much easier) $\Delta I=3/2$ [1=2 $\pi\pi$] Amp on the lattice: 2 contributing topologies only



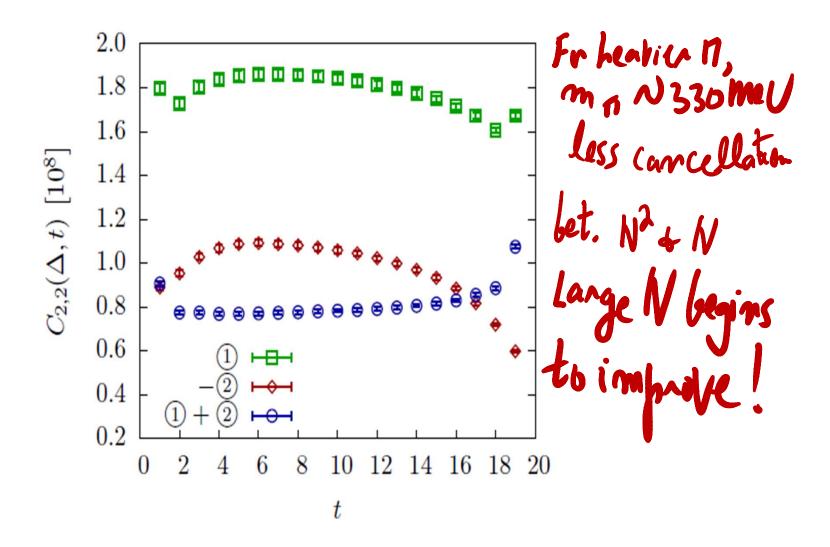


FIG. 3: Contractions ①, -② and ① + ② as functions of t from the simulation at threshold with $m_{\pi} \simeq 330 \,\text{MeV}$ and $\Delta = 20$.

Mass depends of ReA2, A0



| | $a^{-1} [\text{GeV}]$ | $m_{\pi} [\mathrm{MeV}]$ | $m_K[\mathrm{MeV}]$ | $\mathrm{Re}A_{2}\left[10^{\text{-}8}\mathrm{GeV}\right]$ | $\mathrm{Re}A_0[10^{\text{-8}}\mathrm{GeV}]$ | $\frac{\text{Re}A_0}{\text{Re}A_2}$ | notes |
|-------------------------|-----------------------|---------------------------|---------------------|---|--|-------------------------------------|-----------------------|
| 16 ³ Iwasaki | 1.73(3) | 422(7) | 878(15) | 4.911(31) | 45(10) | 9.1(2.1) | threshold calculation |
| 24 ³ Iwasaki | 1.73(3) | 329(6) | 662(11) | 2.668(14) | 32.1(4.6) | 12.0(1.7) | threshold calculation |
| IDSDR | 1.36(1) | 142.9(1.1) | 511.3(3.9) | 1.38(5)(26) | - | | physical kinematics |
| Experiment | - | 135 - 140 | 494 - 498 | 1.479(4) | 33.2(2) | 22.45(6) | |

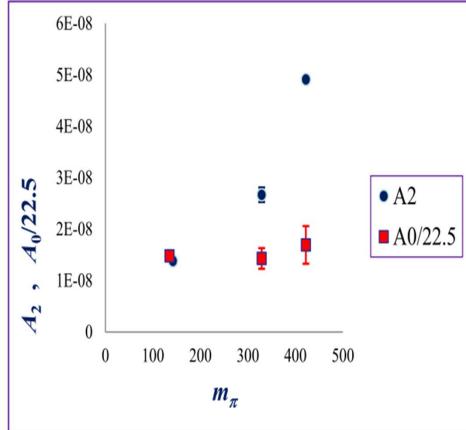
TABLE I: Summary of simulation parameters and results obtained on three DWF ensembles.

Due to the cancellation, 3/2 amplitude decreases significantly as the pion mass is lowered towards its physical value

Delause of mass dependent concellation

Compare A_2 and $A_0/22.5$

ReAz changes with mor dagmataly. For ReAs mass dependence is action mild





$$Q_1 = (\bar{s}_{\alpha} d_{\alpha})_L (\bar{u}_{\beta} u_{\beta})_L,$$

$$Q_2 = (\bar{s}_{\alpha} d_{\beta})_L (\bar{u}_{\beta} u_{\alpha})_L$$

$$Q_3 = (\bar{s}_{\alpha} d_{\alpha})_L \sum_{q=u,d,s} (\bar{q}_{\beta} q_{\beta})_{L,s}$$

$$Q_4 = (\bar{s}_{\alpha} d_{\beta})_L \sum_{q=u,d,s} (\bar{q}_{\beta} q_{\alpha})_L$$

$$Q_5 = (\bar{s}_{\alpha} d_{\alpha})_L \sum_{q=u,d,s} (\bar{q}_{\beta} q_{\beta})_R$$

$$Q_6 = (\bar{s}_{\alpha} d_{\beta})_L \sum_{\alpha = \nu, d} (\bar{q}_{\beta} q_{\alpha})_R,$$

$$Q_{1} = (\bar{s}_{\alpha}d_{\alpha})_{L}(\bar{u}_{\beta}u_{\beta})_{L},$$

$$Q_{2} = (\bar{s}_{\alpha}d_{\beta})_{L}(\bar{u}_{\beta}u_{\alpha})_{L},$$

$$Q_{3} = (\bar{s}_{\alpha}d_{\alpha})_{L}\sum_{q=u,d,s}(\bar{q}_{\beta}q_{\beta})_{L},$$

$$Q_{4} = (\bar{s}_{\alpha}d_{\alpha})_{L}\sum_{q=u,d,s}(\bar{q}_{\beta}q_{\beta})_{L},$$

$$Q_{5} = (\bar{s}_{\alpha}d_{\alpha})_{L}\sum_{q=u,d,s}(\bar{q}_{\beta}q_{\beta})_{L},$$

$$Q_{7} = \frac{3}{2}(\bar{s}_{\alpha}d_{\alpha})_{L}\sum_{q=u,d,s}e_{q}(\bar{q}_{\beta}q_{\alpha})_{R},$$

$$Q_{8} = \frac{3}{2}(\bar{s}_{\alpha}d_{\beta})_{L}\sum_{q=u,d,s}e_{q}(\bar{q}_{\beta}q_{\alpha})_{R},$$

$$Q_{9} = \frac{3}{2}(\bar{s}_{\alpha}d_{\alpha})_{L}\sum_{q=u,d,s}e_{q}(\bar{q}_{\beta}q_{\beta})_{L},$$

$$Q_{9} = \frac{3}{2}(\bar{s}_{\alpha}d_{\alpha})_{L}\sum_{q=u,d,s}e_{q}(\bar{q}_{\beta}q_{\beta})_{L},$$

$$Q_{9} = \frac{3}{2}(\bar{s}_{\alpha}d_{\alpha})_{L}\sum_{q=u,d,s}e_{q}(\bar{q}_{\beta}q_{\beta})_{L},$$

$$Q_8 = \frac{3}{2} (\bar{s}_{\alpha} d_{\beta})_L \sum_{q=u,d,s} e_q (\bar{q}_{\beta} q_{\alpha})_R,$$

$$Q_9 = \frac{3}{2} (\bar{s}_{\alpha} d_{\alpha})_L \sum_{q=u,d,s} e_q (\bar{q}_{\beta} q_{\beta})_L,$$

$$Q_{10} = \frac{3}{2} (\bar{s}_{\alpha} d_{\beta})_L \sum_{q=u,d,s} e_q (\bar{q}_{\beta} q_{\alpha})_L,$$

 $Q_{7} = \frac{1}{2} (\bar{u}_{\beta} u_{\beta})_{L},$ $Q_{7} = \frac{1}{2} (\bar{s}_{\alpha} d_{\beta})_{L} (\bar{u}_{\beta} u_{\alpha})_{L},$ $Q_{8} = \frac{3}{2} (\bar{s}_{\alpha} d_{\beta})_{L} \sum_{q=u,d,s} (\bar{q}_{\beta} q_{\beta})_{L},$ $Q_{9} = \frac{3}{2} (\bar{s}_{\alpha} d_{\alpha})_{L} \sum_{q=u,d,s} \epsilon_{\alpha}$ $Q_{10} = \frac{3}{2} (\bar{s}_{\alpha} d_{\beta})_{L} \sum_{q=u,d,s} e_{q} (\bar{q}_{\beta} q_{\alpha})_{L},$ $Q_{5} = (\bar{s}_{\alpha} d_{\alpha})_{L} \sum_{q=u,d,s} (\bar{q}_{\beta} q_{\alpha})_{R},$ $Q_{6} = (\bar{s}_{\alpha} d_{\beta})_{L} \sum_{q=u,d,s} (\bar{q}_{\beta} q_{\alpha})_{R},$ $Q_{7} = \frac{3}{2} (\bar{s}_{\alpha} d_{\beta})_{L} \sum_{q=u,d,s} \epsilon_{\alpha}$ $Q_{10} = \frac{3}{2} (\bar{s}_{\alpha} d_{\beta})_{L} \sum_{q=u,d,s} e_{q} (\bar{q}_{\beta} q_{\alpha})_{L},$ $Q_{10} = \frac{3}{2} (\bar{s}_{\alpha} d_{\beta})_{L} \sum_{q=u,d,s} e_{q} (\bar{q}_{\beta} q_{\alpha})_{L},$ $Q_{10} = \frac{3}{2} (\bar{s}_{\alpha} d_{\beta})_{L} \sum_{q=u,d,s} e_{q} (\bar{q}_{\beta} q_{\alpha})_{L},$ $Q_{10} = \frac{3}{2} (\bar{s}_{\alpha} d_{\beta})_{L} \sum_{q=u,d,s} e_{q} (\bar{q}_{\beta} q_{\alpha})_{L},$ $Q_{10} = \frac{3}{2} (\bar{s}_{\alpha} d_{\beta})_{L} \sum_{q=u,d,s} e_{q} (\bar{q}_{\beta} q_{\alpha})_{L},$ $Q_{10} = \frac{3}{2} (\bar{s}_{\alpha} d_{\beta})_{L} \sum_{q=u,d,s} e_{q} (\bar{q}_{\beta} q_{\alpha})_{L},$ $Q_{10} = \frac{3}{2} (\bar{s}_{\alpha} d_{\beta})_{L} \sum_{q=u,d,s} e_{q} (\bar{q}_{\beta} q_{\alpha})_{L},$ $Q_{10} = \frac{3}{2} (\bar{s}_{\alpha} d_{\beta})_{L} \sum_{q=u,d,s} e_{q} (\bar{q}_{\beta} q_{\alpha})_{L},$

Pheno 2018; A Soni

Additional Improvements/checks in lattice E' determination underway for past ~2 years

• EM+ isospinfforder extensive study Ciniglians, Neufoly Educat Pich 13 4 Com PhD Stulent Dentrying Completely diff method(s) • 1) excited pipi stateph 1 (Dan Murphy ChiT wash!!
• 11) Revisit ChPT well for physical 11, K moss 4 e NLO!

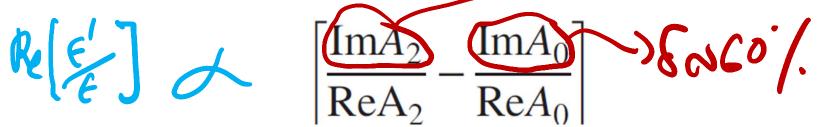
Reexamine BDSPW 84, Lailott AS 04 for E

Guess estimate of reduction of errors

- δ (Im A0) from 65% => 20-25%
- δ (ReA0) from 35% => 15-20% [don't use for ϵ ' for now]
- Uncorrelated fits (due to lack of stat) =>

Very good chance we'll be able to correlated fits now Systematic error from ~ 27% => ~ 20%

- Effect on ε' unclear: 'cause of hefty cancellations



underlying method is systematically improvable =>multitude of successful demonstrations by now

- BK in full QCD with DWF '07 RBC-UKQCD error O(7%)
- Since ~2012 many discretizations, WA error O(1-2%)
- Re A2 from ~25% around 2012 to now ~10% (now no longer due to lattice but only only due to perturbation theory error upto NLO!)
- KI3, A2, fB's , BB's......
- Quark masses; in particular ms no longer anywhere around ~150
 MeV [used to be PDG value] but now ~100 MeV.
- No doubt that A0 and ε' will also go that way for quite sometime to come......to ~10% total in another ~ 3-5years!.

After that EM& isospin effects need to be ascertained quantitatively; WIP