

Collider & Astrophysical Phenomenology of Massive Spin-2 Mediated WIMP Dark Matter (Preview!)

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The Standard Model? Dark Matter??

The Universe's Energy Budget

SM Content: $\sim 5\%$	\leftarrow	Good at High E
Dark Energy: $\sim 71\%$	\leftarrow	Vacuum Energy?
Dark Matter: $\sim 24\%$	\leftarrow	???

DM Guess: Non-baryonic thermally-cold neutral massive particle?

The theory DM density must be less than observed DM density (“**relic density/abundance**”), which equals¹

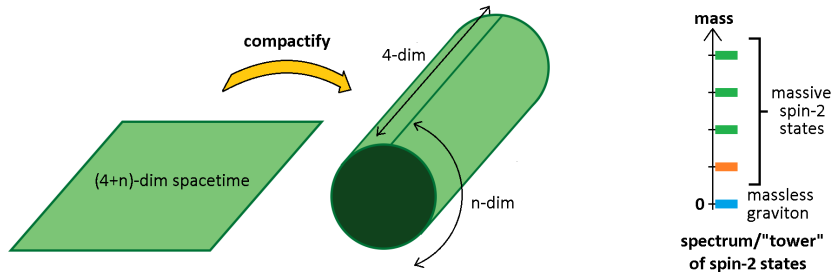
$$(\Omega h^2)_{theory} \leq (\Omega h^2)_{exp} = \left[\frac{\rho_{DM}}{\rho_{critical}} \right] h^2 = 0.11425 \pm 0.00311$$

* **What if DM only interacts gravitationally?** *

¹P. A. R. Ade, et al., Planck 2015 results. XIII. Cosmological parameters, Astron. Astrophys. 594 (2016) A13.

Higher-Dimensional Consequences

Massive spin-2 particles naturally emerge when compactifying higher-dimensional gravity to 4-dimensions (such as in Randall-Sundrum models); **4D spectrum depends on model.**



Q: Can we make any general statements about the phenomenology of these models by focusing on the lowest massive state?

The Effective Massive Spin-2 Theory... and Dark Matter

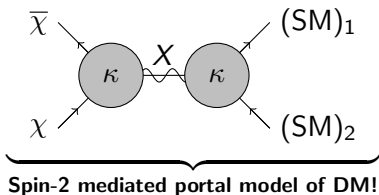
We may write the **simplified massive Spin-2 Lagrangian**,

$$\mathcal{L}_X = \underbrace{\sqrt{-\det g} \left[\frac{2}{\kappa^2} R(g) + \mathcal{L}_{\text{matter}}(g) \right]}_{\text{usual 4D gravity Lagrangian; expand in } \kappa} + \underbrace{m_X^2 (X^{\mu\nu} X_{\mu\nu} - X^\mu{}_\mu X^\nu{}_\nu)}_{\text{spin-2 "Fierz-Pauli" mass term}}$$

where $g_{\mu\nu} \equiv \eta_{\mu\nu} + \kappa X_{\mu\nu}$, and $X_{\mu\nu}$ = massive spin-2 field.

Suppose we also throw in a **dark sector** w/ **WIMP-like** χ ...

$$\mathcal{L}_{\text{matter}} = \mathcal{L}_{\text{SM}} + \mathcal{L}_{\text{DM}} \implies$$



Key Points - The Project in Summary

Model Space: We test viability of this model for...

- We scan a wide range of masses m_χ , $m_{\bar{\chi}}$, and coupling κ and **we include second-order interaction vertices.**
- **DM Nature:** $\chi = \bar{\chi}$ and $\chi \neq \bar{\chi}$
- **DM Spin:** $\chi =$ scalar, fermion, and vector

Constraints: ... where viability is measured against...

- **Collider Limits** (ATLAS & CMS diphoton, dilepton, ...)
- **Relic Density** (Planck)
- **Direct DM Searches** (XENON1T)
- **Indirect DM Searches** (H.E.S.S.-I & II, Fermi-LAT, DAMPE)
- **Tree-Level Partial Wave Unitarity** (b.c. tower is truncated)

Key Points - The Project in Summary

Model Space: We test viability of this model for...

- We scan a wide range of masses m_χ , $m_{\bar{\chi}}$, and coupling κ and we include **second-order interaction vertices**. ($m_\chi = 1 \text{ TeV}$)
- **DM Nature:** $\chi = \bar{\chi}$ and $\chi \neq \bar{\chi}$
- **DM Spin:** $\chi = \text{scalar}$, fermion, and vector

Constraints: ... where viability is measured against...

- **Collider Limits** (ATLAS & CMS diphoton, dilepton, ...)
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Key Complications

Plenty of room for errors, so to ensure correct & confident results,

Guiding Principle: calculate from scratch, and
only then cross-check with community tools

In the process, we've developed programs for...

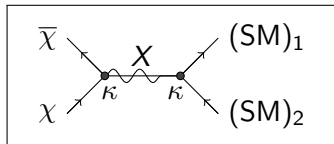
1. Computing **analytic** functions of **helicity amplitudes**
2. Calculating a model's **relic density** given the appropriate σ_{ann}

Furthermore, **we perform a complete $\mathcal{O}(\kappa^2)$ analysis, including multi-mediator vertices.**

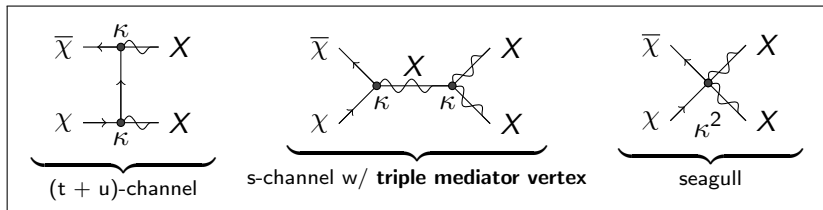
Example Diagrams: DM Annihilation

Many models only include $\mathcal{O}(\kappa)$ matter-mediator vertices:

$$\mathcal{L}_X \supseteq -\frac{\kappa}{2}(T_\chi^{\mu\nu} + T_{SM}^{\mu\nu})X_{\mu\nu} \implies$$



Our simplified model includes additional $\mathcal{O}(\kappa^2)$ diagrams, including **dark matter annihilation to final-state mediator pairs**



Triple Mediator Vertex

So we need the X^3 vertex, which comes from:

$$\mathcal{L}_X \supseteq \frac{\kappa}{8} X_{\mu\nu} \left\{ \eta^{\mu\nu} \left[-(\partial_\rho X)(\partial^\rho X) + 2(\partial_\rho X_{\sigma\tau})(\partial^\rho X^{\sigma\tau}) \right] - 4(\partial^\mu X_{\rho\sigma})(\partial^\nu X^{\rho\sigma}) \right. \\ \left. - 8(\partial_\sigma X^{\nu\rho})(\partial^\sigma X_\rho{}^\mu) + 4(\partial_\rho X^{\mu\nu})(\partial^\rho X) + 8(\partial^\mu X_{\rho\sigma})(\partial^\rho X^{\sigma\nu}) \right\}$$

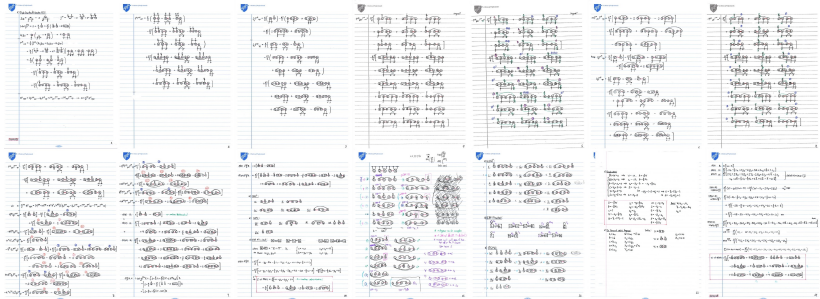
These terms result from expanding \mathcal{L}_{EH} , which naively generates \sim **180 terms**. *Yikes*.

3. We developed a **diagrammatic method** for calculating the $\mathcal{O}(X^3)$ Lagrangian and with it ultimately **we derived the X^3 terms by hand!**

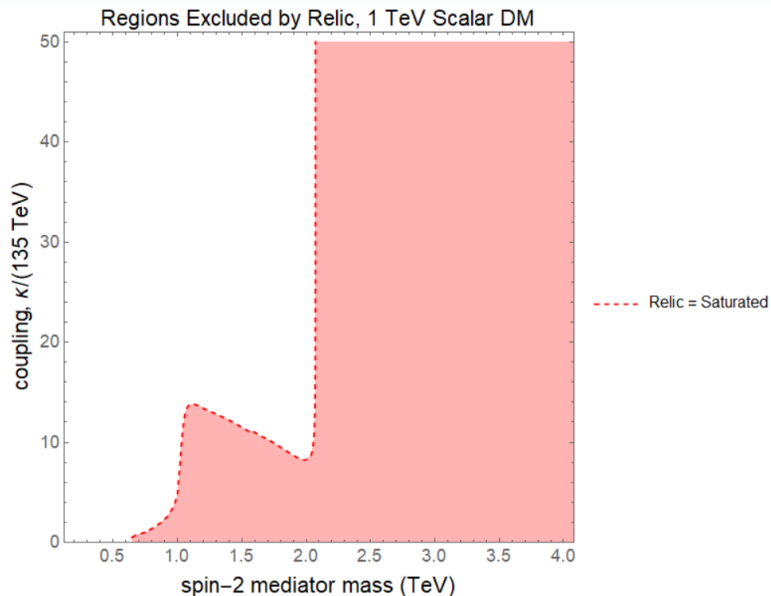
Triple Mediator Vertex

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Example of Relic Density Constraint



Partial Wave Unitarity

$$i\mathcal{M}_{\lambda_a\lambda_b;\lambda_c\lambda_d} = \begin{array}{c} \bar{q}_{\lambda_b} \swarrow \quad \nearrow Q_{\lambda_c} \\ \bullet \\ q_{\lambda_a} \searrow \quad \swarrow Q_{\lambda_d} \end{array}$$

This matrix element decomposes into **partial wave amplitudes**:

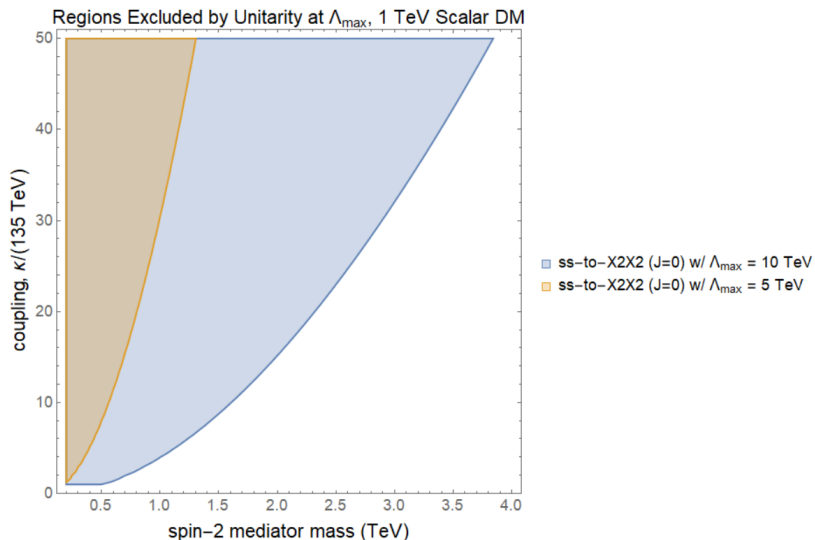
$$\mathcal{M}_{\lambda_a\lambda_b;\lambda_c\lambda_d} \equiv 32\pi^2 \sum_{J=\max\{\lambda_i,\lambda_f\}}^{+\infty} \frac{2J+1}{4\pi} \cdot a_{\lambda_a\lambda_b;\lambda_c\lambda_d}^J(s) \cdot D_{\lambda_a-\lambda_b,\lambda_c-\lambda_d}^{J*}(\theta, \phi)$$

where $D_{\lambda_i,\lambda_f}^J =$ Wigner D -function. The **optical theorem** implies,

$$* \left[\sqrt{1 - \frac{4m_q^2}{s}} \cdot \left| \Re[a_{\lambda_a,\lambda_b;\lambda_c,\lambda_d}^J(s)] \right| \leq 1 \right] *$$

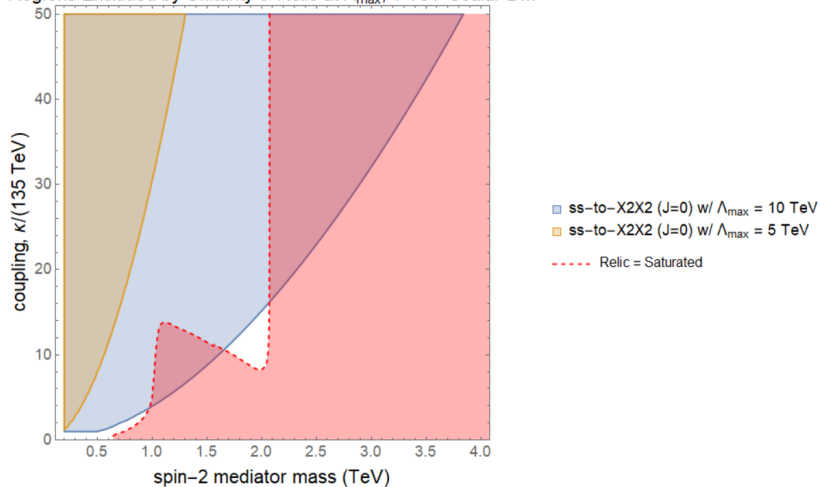
for *all* amplitudes. **We demand satisfaction of this inequality up to some appropriately-chosen scale Λ_{max} .**

Example of Unitarity Constraints



Example of Unitarity + Relic Density Constraints

Regions Excluded by Unitarity & Relic at Λ_{\max} , 1 TeV Scalar DM



How Apply Unitarity To Relic Density?

The relic density depends on the **DM annihilation cross-section**:

$$\sigma_{\text{ann}} = \sum_{\text{SM}} \sigma \left(\begin{array}{c} \bar{\chi} \quad (SM)_1 \text{ or } X \\ \chi \quad (SM)_2 \text{ or } X \end{array} \right)$$

Namely, though the **thermally-averaged cross-section**,

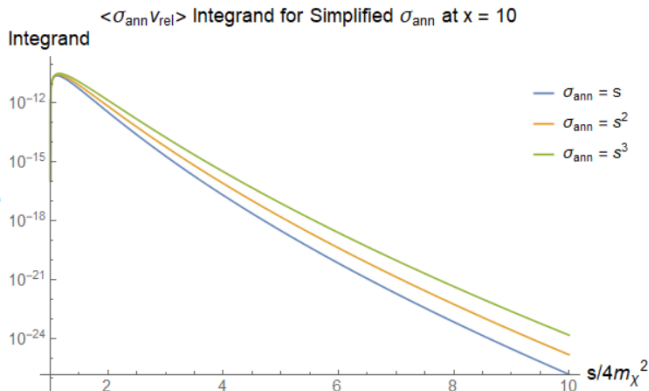
$$\langle \sigma_{\text{ann}} v_{\text{rel}} \rangle(x) = \frac{4x}{K_2(x)^2} \int_1^{+\infty} d\bar{s} \sqrt{\bar{s}} \cdot (\bar{s} - 1) \cdot K_1(2x\sqrt{\bar{s}}) \cdot \sigma_{\text{ann}}$$

where $\bar{s} \equiv s/4m_\chi^2$ and evolution parameter $x = m_\chi/T$.

$\langle \sigma v_{\text{rel}} \rangle$ **needs information about $s = +\infty!$**

Well...

Integrand of Thermally-Averaged Cross Section



Choose pre-freeze-out x -value, demand accuracy ϵ

\implies Integrate over $[1, \bar{s}_{\text{max}}]$ so that $= (1 - \epsilon) \times (\text{full integral})$

\implies Partial-wave analysis w/ $\Lambda_{\text{max}} = 2m_\chi \sqrt{\bar{s}_{\text{max}}}$

Conclusions & Future of this Work

These are elements of a project exploring the viability of **simplified spin-2 mediated dark matter for $s_\chi = 0, \frac{1}{2}, 1$** . Generally...

- **Relic density and partial-wave unitarity** significantly constrain available parameter space in opposite ways.
- **Indirect detection** restricts models on threshold
- **Direct detection** limits are weak.
- **Collider** limits demand very weak coupling or heavy mediator.

*** Expect to see something on the ArXiv soon! ***

Furthermore, following this project's completion, **we plan to optimize & rework our codes for public use.**

Thank you for your attendance and attention!