

WW Production With a Jet Veto Made Simple(r)

In collaboration with: A. Banfi, S. Jaegar & N. Kauer

Luke Arpino

University of Sussex

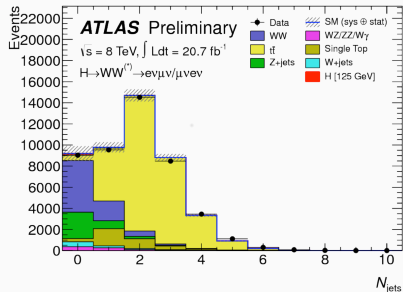
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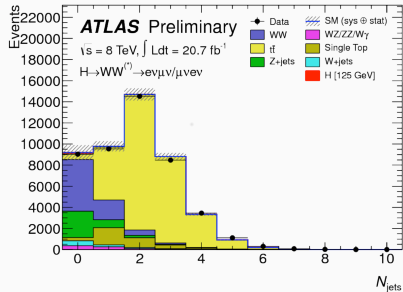
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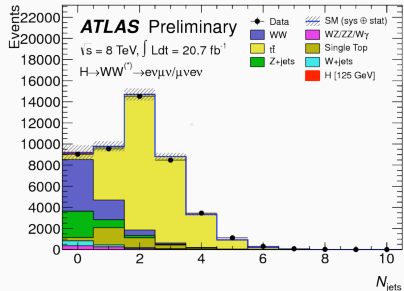
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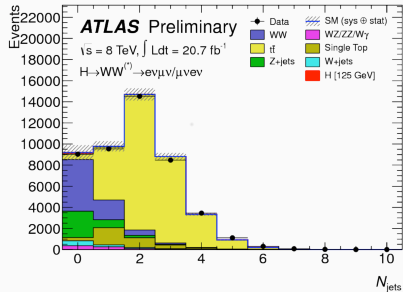
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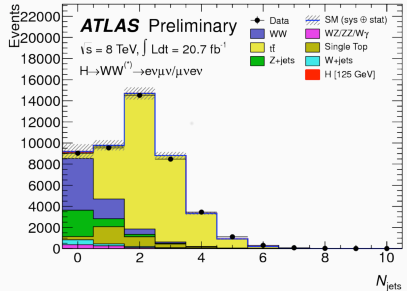
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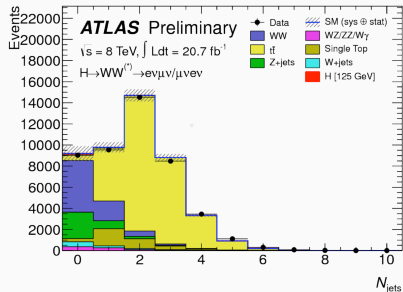
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Higgs + 0-jet cross section becomes the quantity of interest

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Need to quantify how the jet veto reduces the Higgs cross section and the non $t\bar{t}$ cross section

This Talk

1. Jet veto resummation in a nutshell
2. Automation of jet veto resummation
3. A case study: $pp \rightarrow WW$

Jet Veto Resummation in a Nutshell

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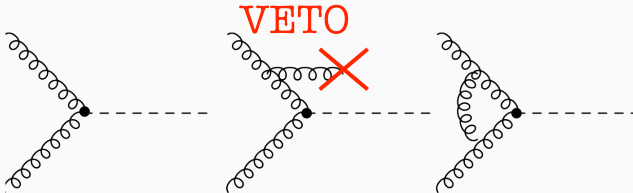
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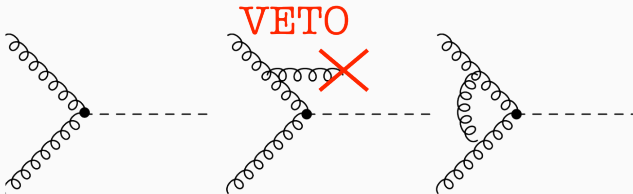
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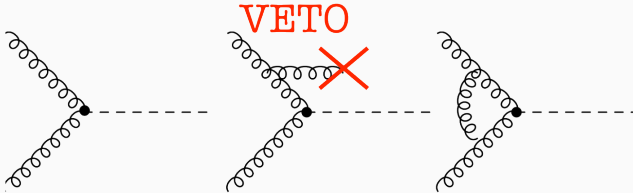
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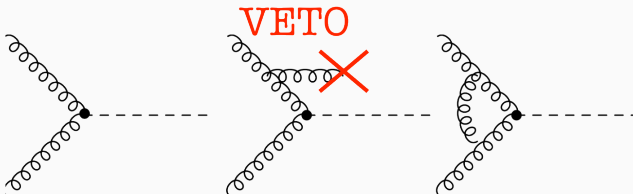
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$$\begin{aligned}\sigma_{0\text{-jet}}(p_{t,\text{veto}}) &\simeq \sigma_0 \left(1 + C \frac{\alpha_s}{\pi} \int \frac{d\omega}{\omega} \frac{d\theta^2}{\theta^2} (\Theta(p_{t,\text{veto}} - \omega\theta) - 1) \right) \\ &\simeq \sigma_0 \left(1 - 2C \frac{\alpha_s}{\pi} \log^2 \frac{M}{p_{t,\text{veto}}} + \dots \right)\end{aligned}$$

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Remnant logs left after cancellation of the real and virtual

The NNLL master formulae

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Banfi, Salam, Monni (2012)

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- Sudakov form factor
- Non-exponentiating multiple emission function

Implementation of Resummation

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Can we do better?

Utilise existing Monte Carlo code

Becher, Frederix, Neubert, Rothen (2014)

The Method

A Quick Review of NLO calculations

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Dipole Subtraction

Cleverly add zero by introducing a subtraction term $d\sigma^A$:

$$\delta\sigma^{NLO} = \int_m d\sigma^V + \int_{m+1} [d\sigma^R - d\sigma^A] + \int_{m+1} d\sigma^A$$

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$$\int_1 d\sigma^A \cong \delta\mathcal{L}$$

Implementation in MCFM

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MCFM + JETVHETO

- Can resum jet vetoes for any colour singlet to NNLL+LL_R
- Keep everything in MCFM—spin correlations, cuts, exclusivity in the leptons, etc.
- Minimal implementation—dipoles and factorisation scale
- Can be brought to other Monte Carlo integrators—POWHEG, AMC@NLO

Completely general!

A case study: $pp \rightarrow WW$

The Setup

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Consider the SMEFT inspired Lagrangian:

$$\mathcal{L} \supseteq -\kappa_t \frac{m_t}{v} H (\bar{t}_R t_L + h.c.) + \kappa_g \frac{\alpha_s}{12\pi} \frac{H}{v} G_{\mu\nu}^a G^{\mu\nu,a}$$

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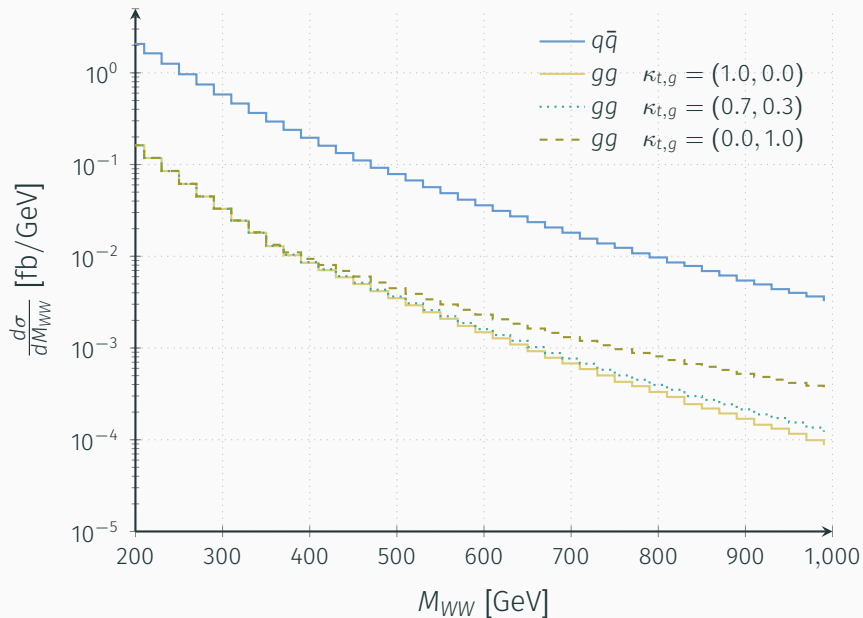
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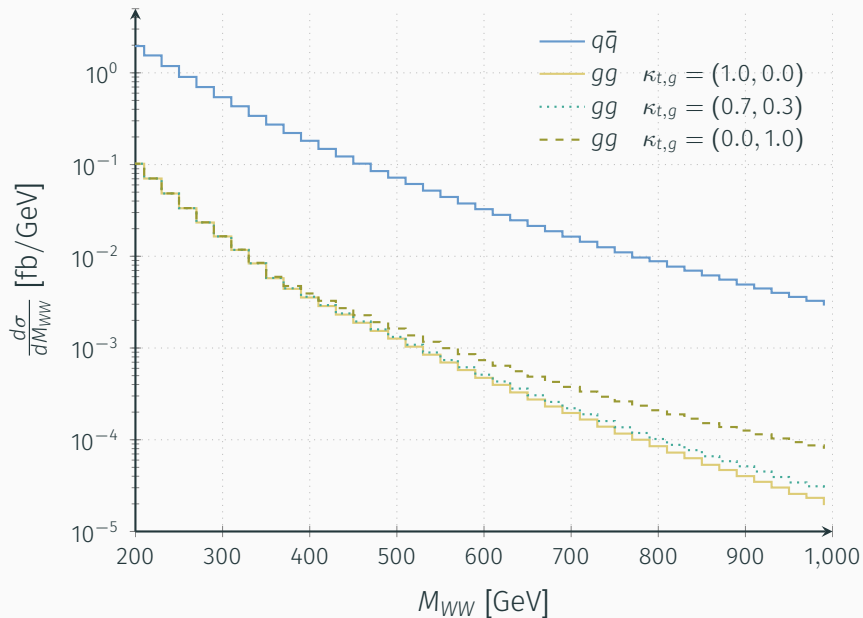
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- Vary ggH and $t\bar{t}H$ couplings together
 - + fix $\kappa_t + \kappa_g = 1$

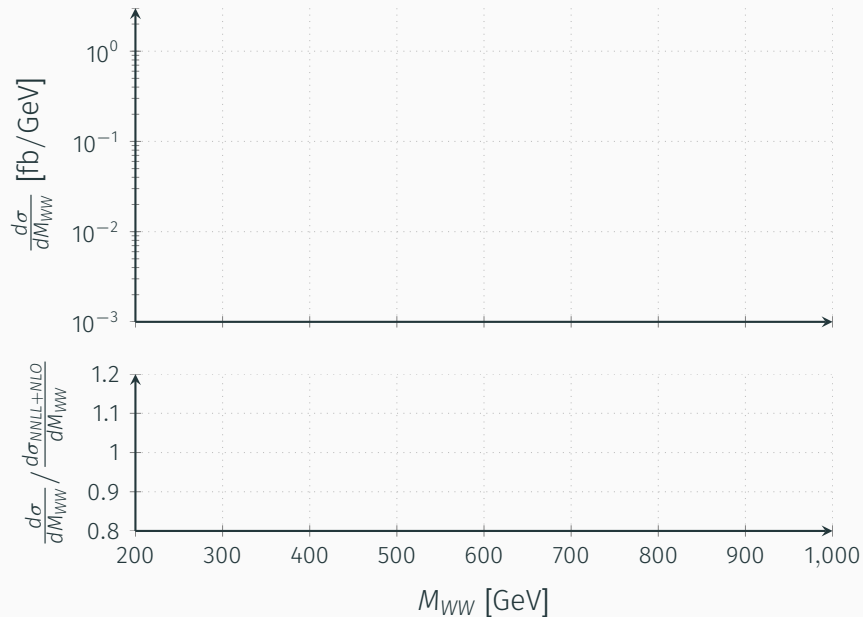
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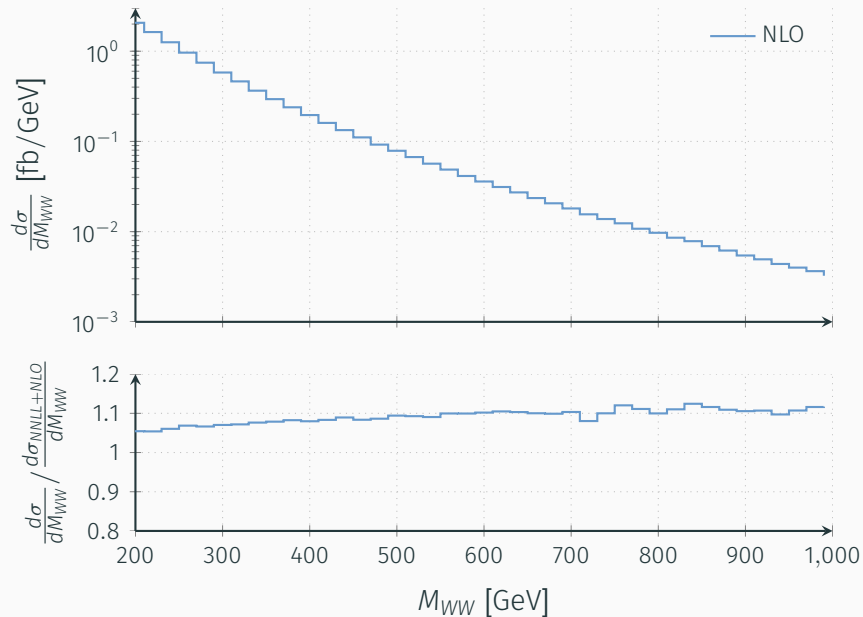
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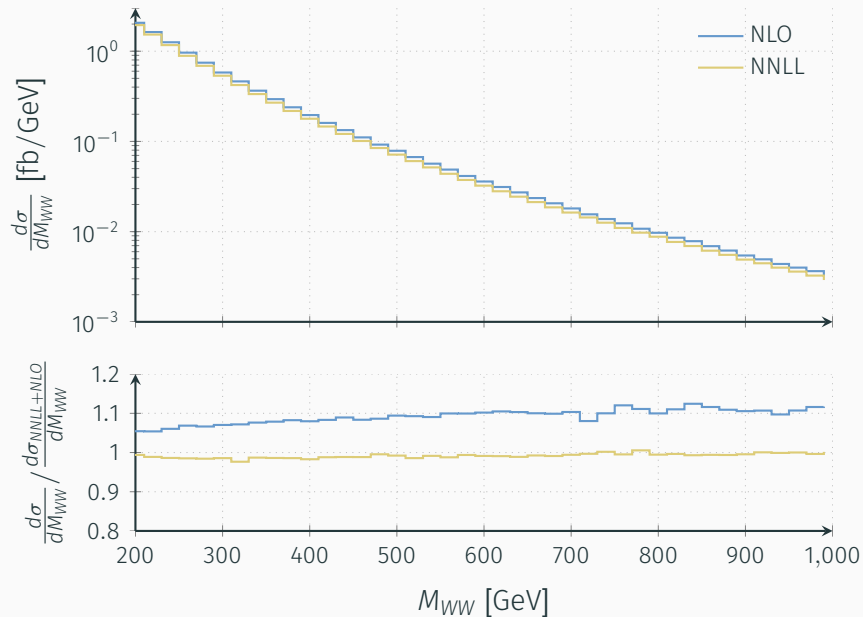
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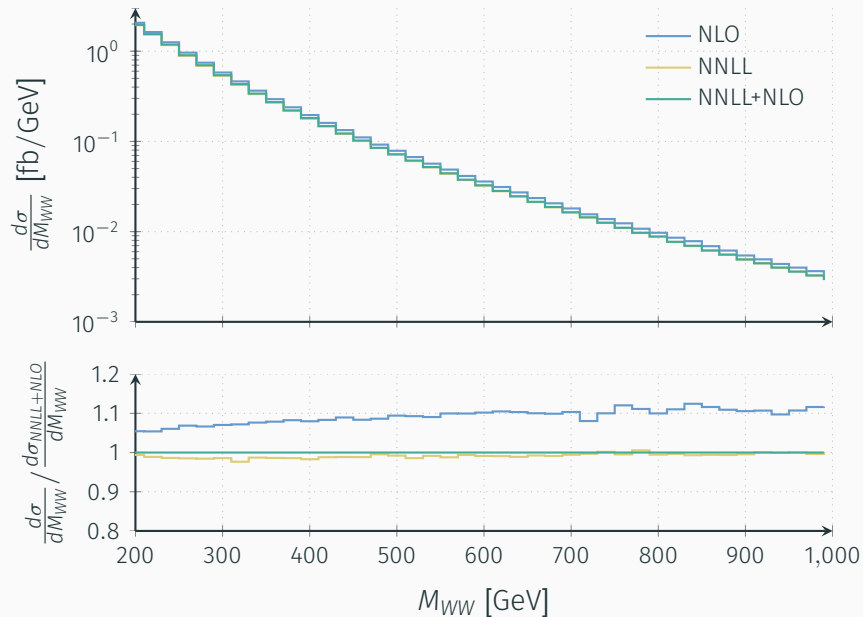
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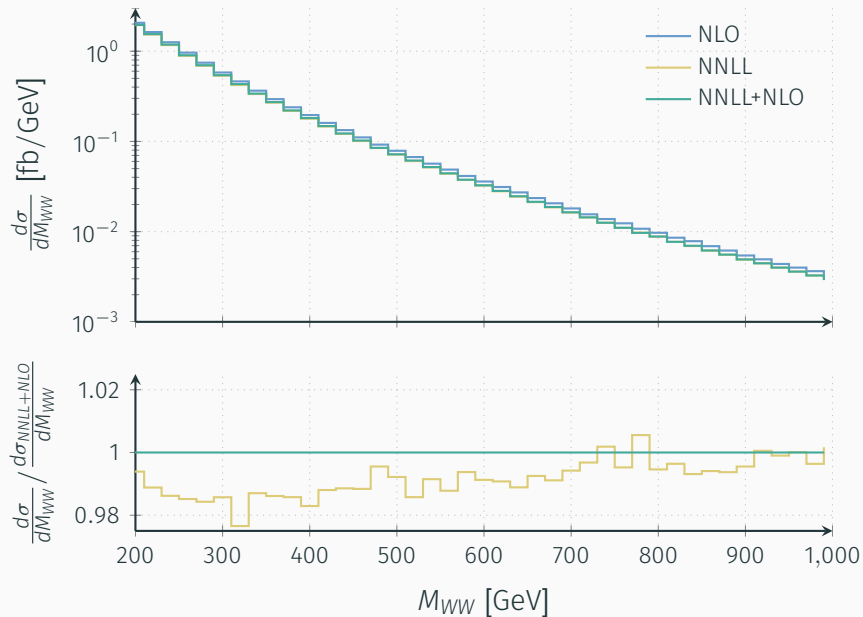
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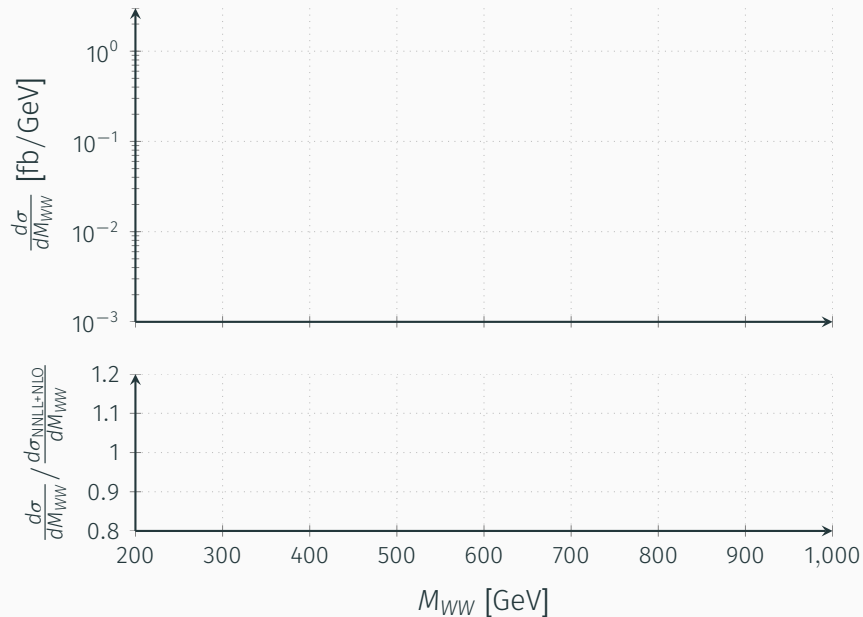
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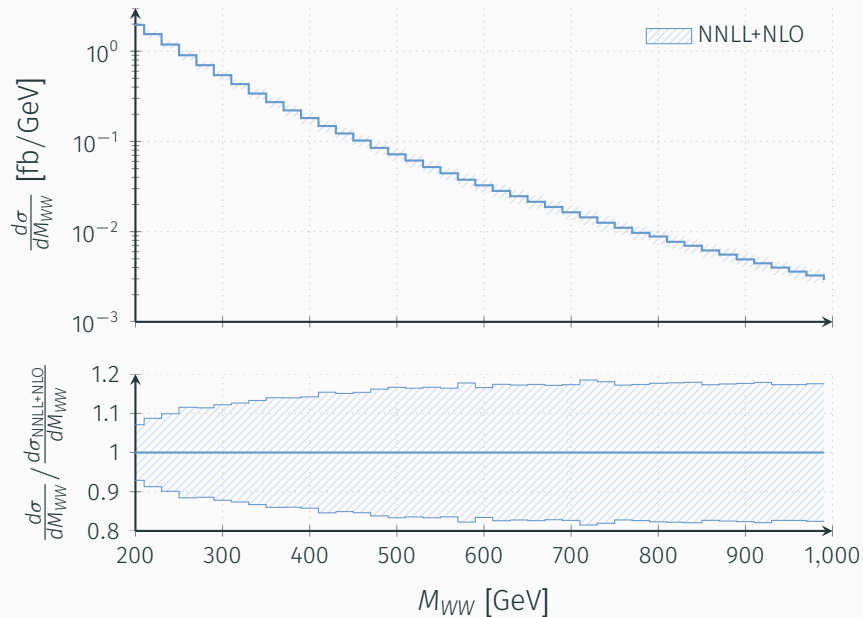
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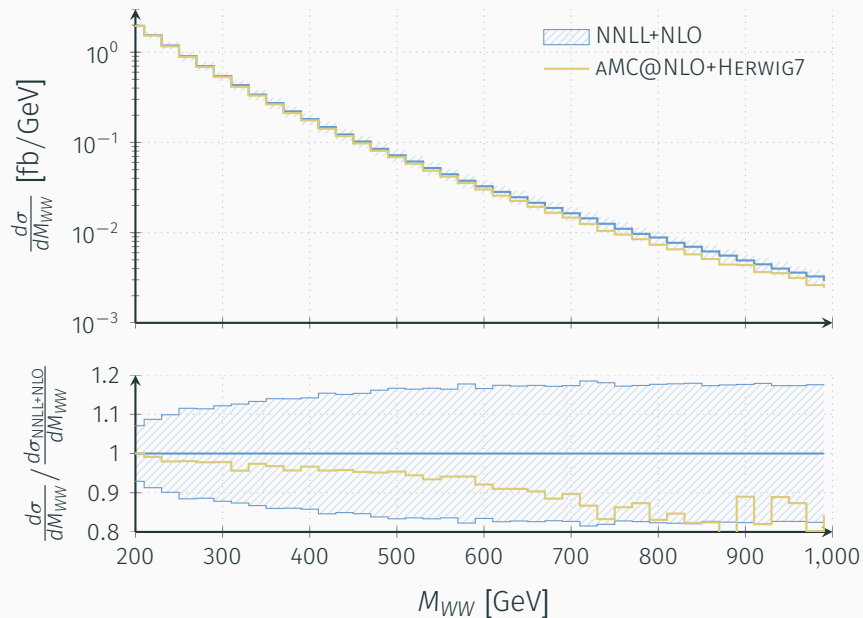
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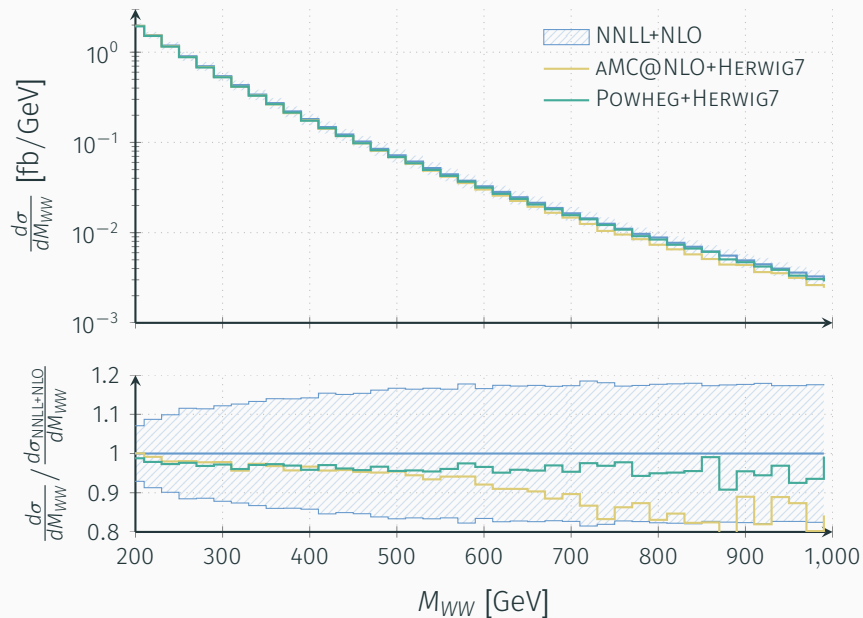
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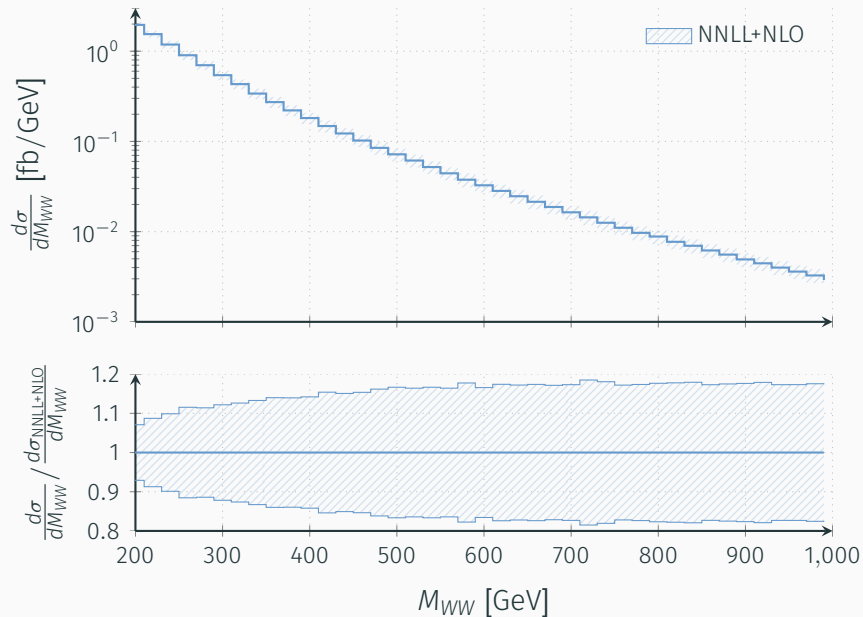
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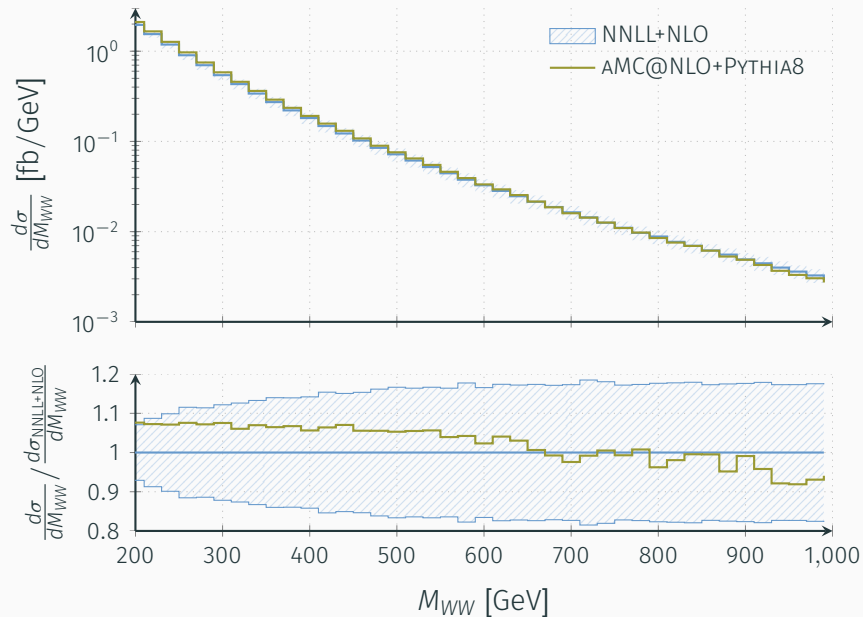
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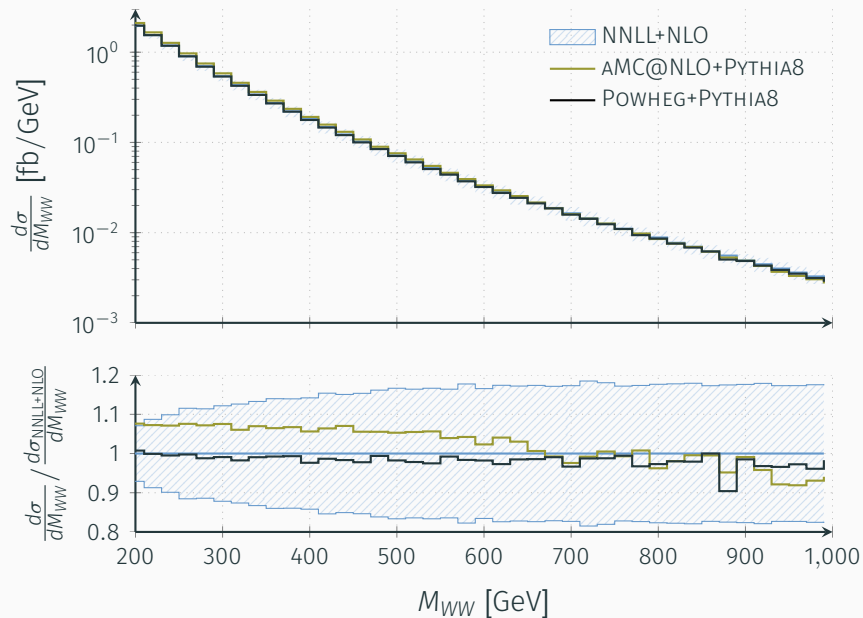
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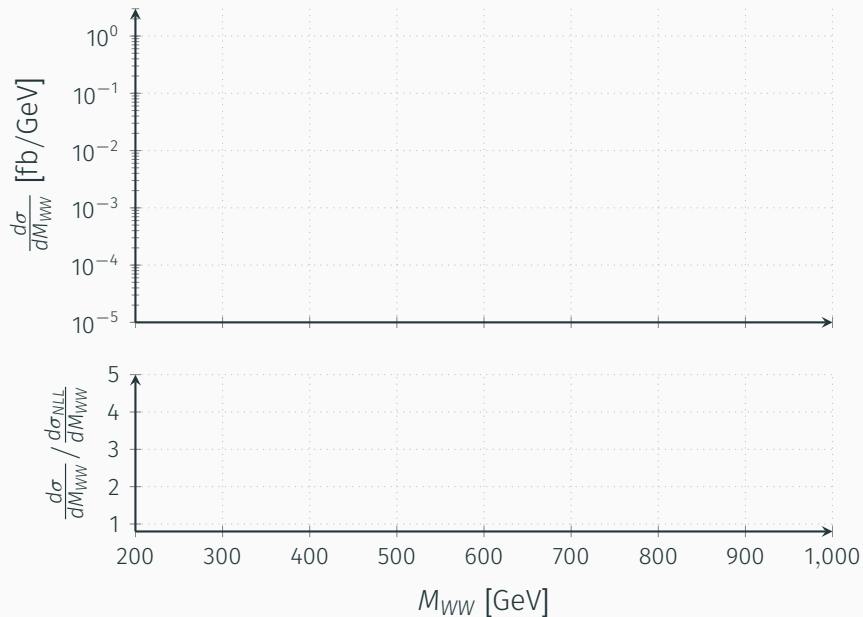
$q\bar{q} \rightarrow WW \rightarrow l^+\nu l^-\bar{\nu}$ (resummed vs. parton shower)



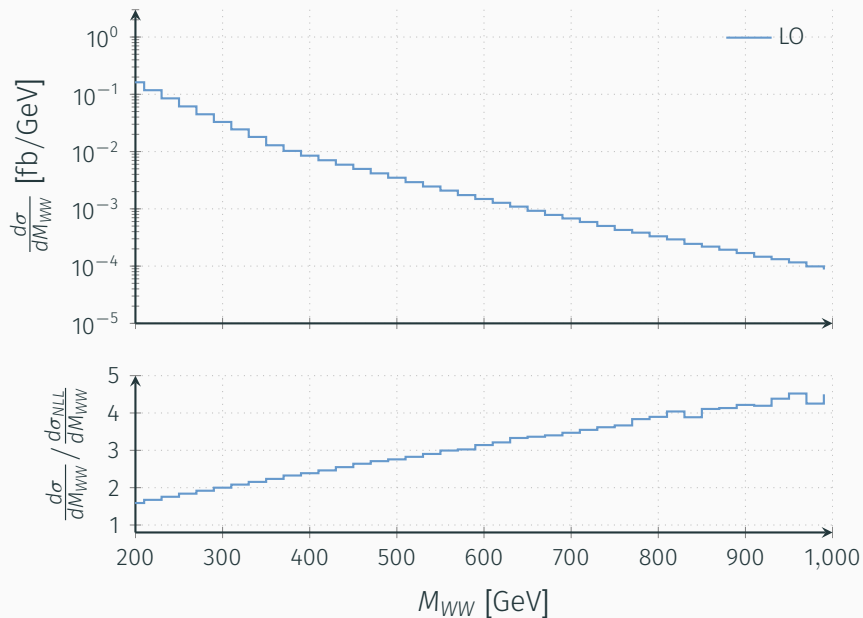
$q\bar{q} \rightarrow WW \rightarrow l^+\nu l^-\bar{\nu}$ (resummed vs. parton shower)



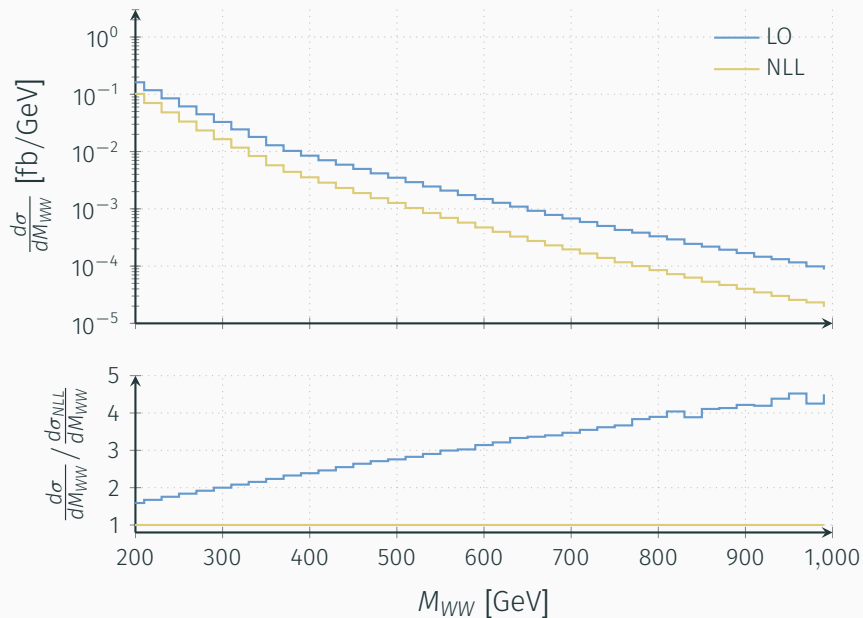
$gg \rightarrow WW \rightarrow l^+ \nu l^- \bar{\nu}$ (fixed-order vs. resummed)



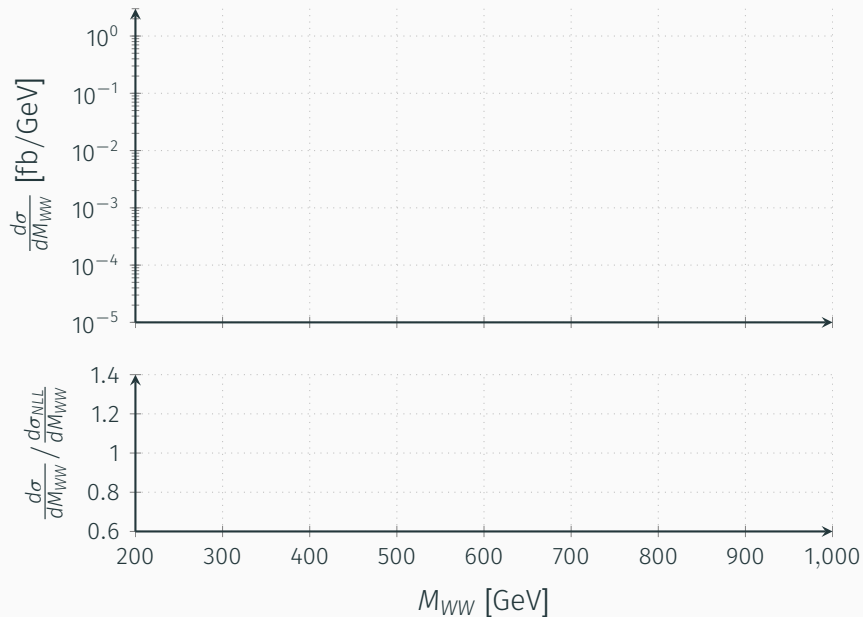
$gg \rightarrow WW \rightarrow l^+ \nu l^- \bar{\nu}$ (fixed-order vs. resummed)



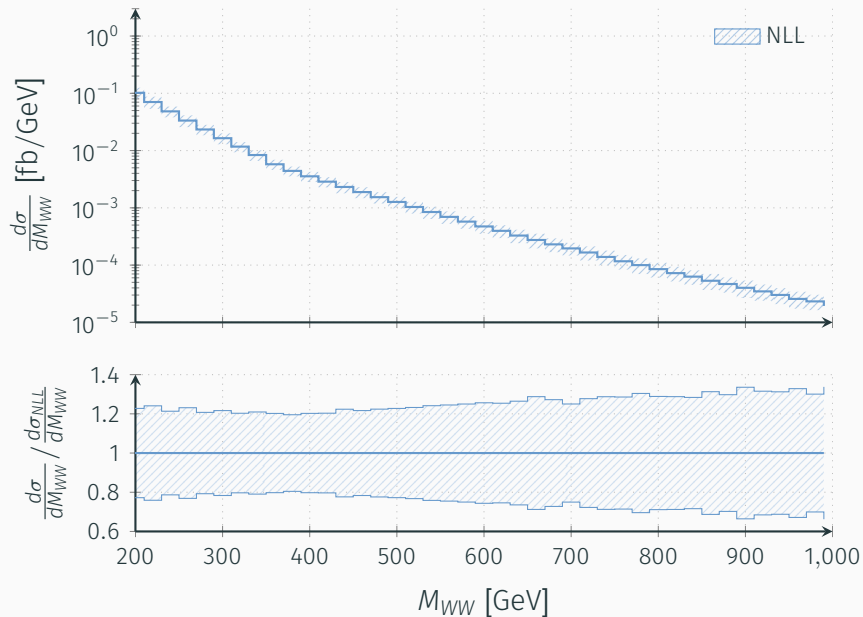
$gg \rightarrow WW \rightarrow l^+ \nu l^- \bar{\nu}$ (fixed-order vs. resummed)



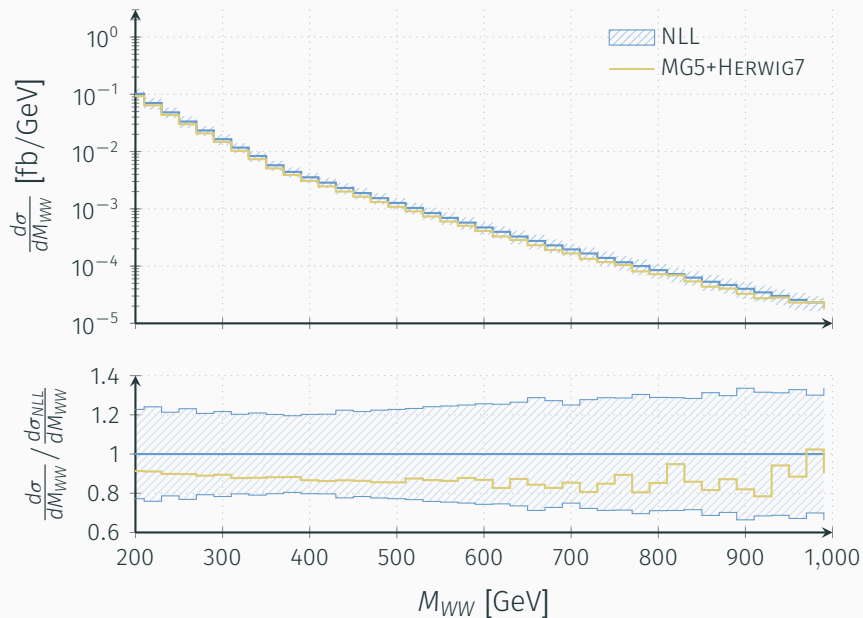
$gg \rightarrow WW \rightarrow l^+ \nu l^- \bar{\nu}$ (resummed vs. parton shower)



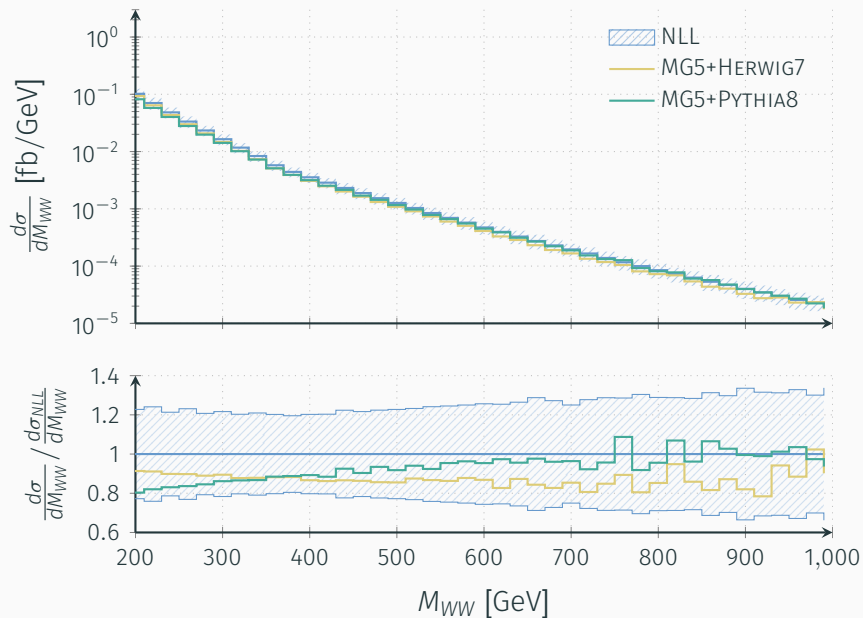
$gg \rightarrow WW \rightarrow l^+ \nu l^- \bar{\nu}$ (resummed vs. parton shower)



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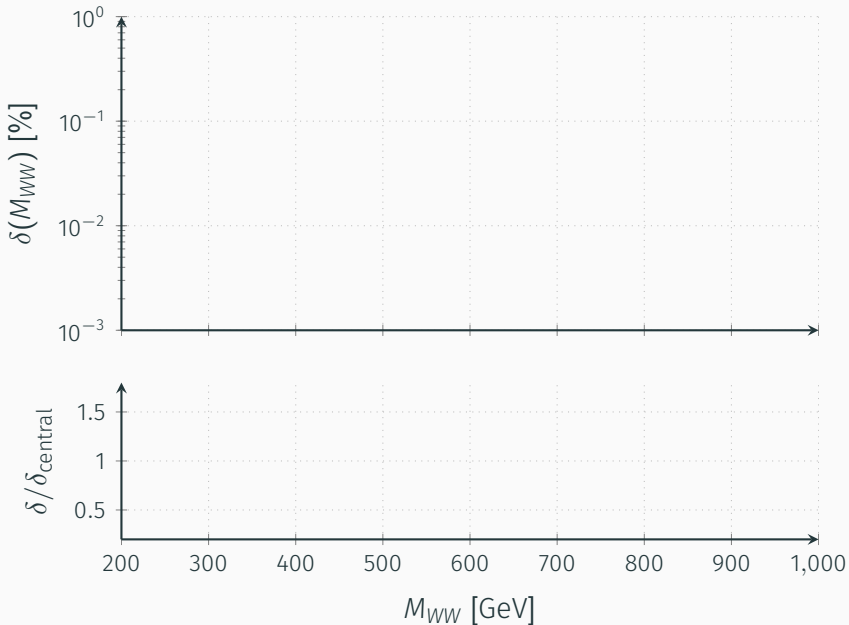


Quantifying these Results

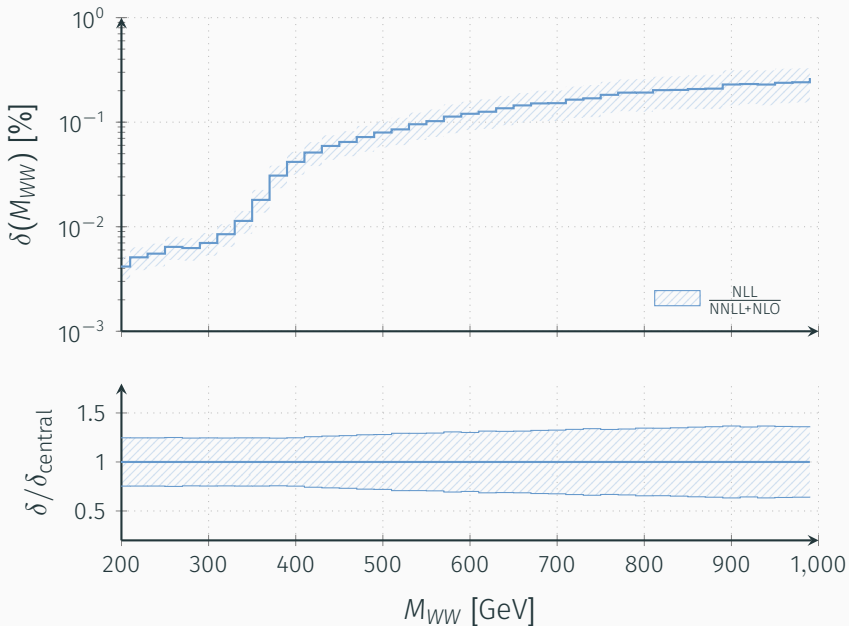
Parameterise discrepancy with respect to the standard model:

$$\delta = \frac{\sigma^{BSM} - \sigma^{SM}}{\sigma^{SM}}$$

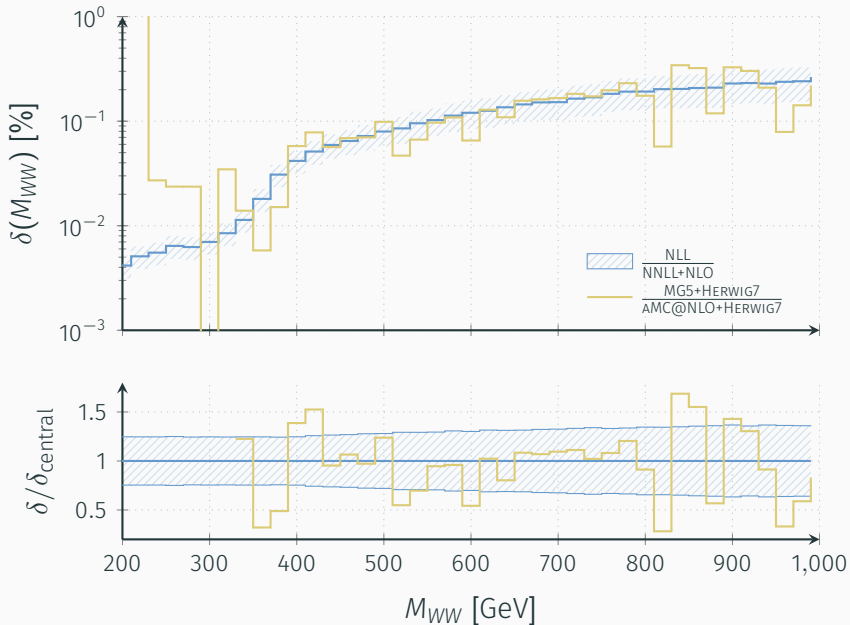
$\delta, \kappa_{t,g} = (0.7, 0.3)$ (resummed vs. parton shower)



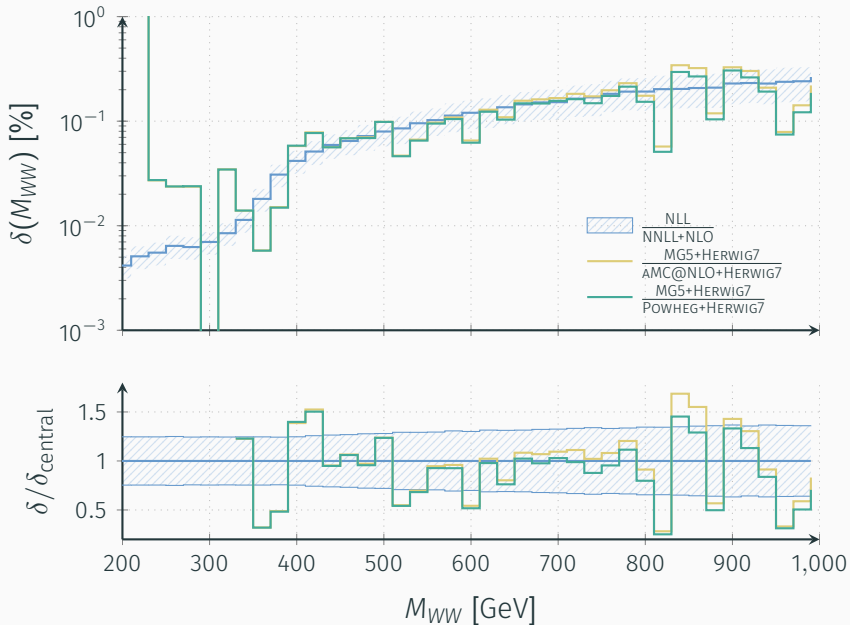
$\delta, \kappa_{t,g} = (0.7, 0.3)$ (resummed vs. parton shower)



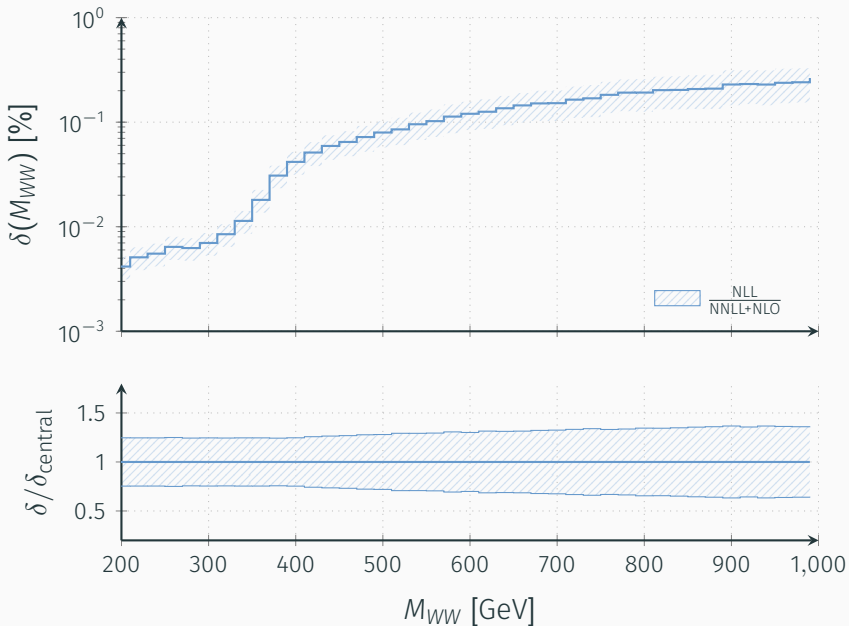
$\delta, \kappa_{t,g} = (0.7, 0.3)$ (resummed vs. parton shower)



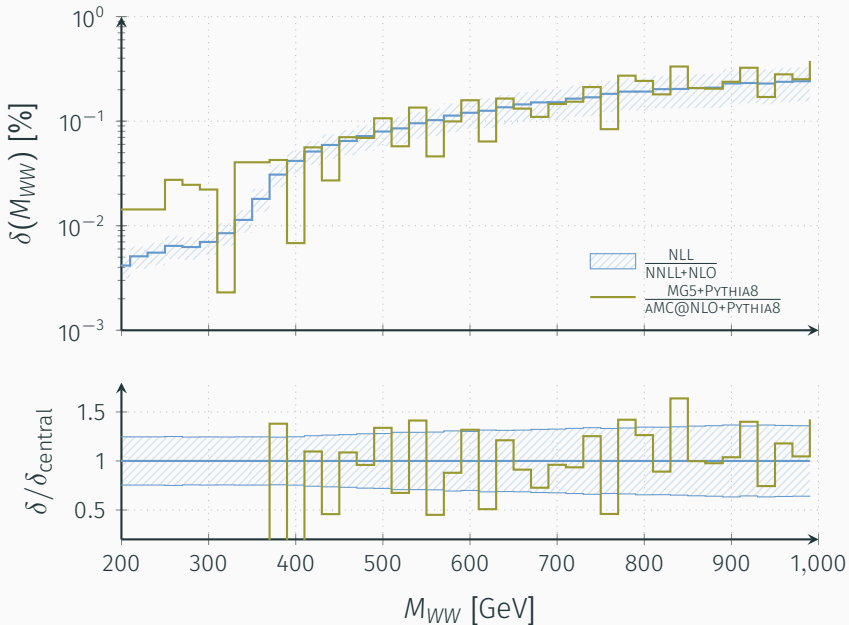
$\delta, \kappa_{t,g} = (0.7, 0.3)$ (resummed vs. parton shower)



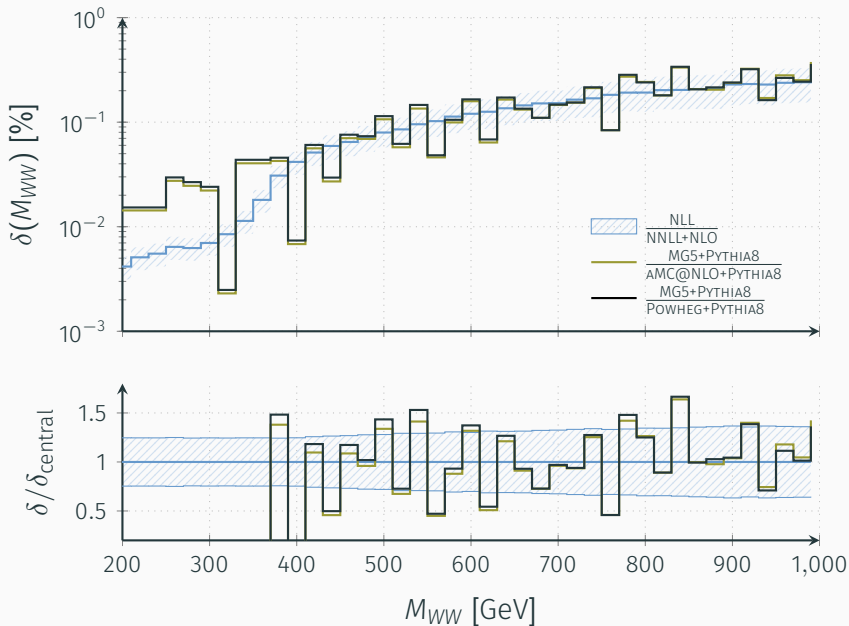
$\delta, \kappa_{t,g} = (0.7, 0.3)$ (resummed vs. parton shower)



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Repurposing Monte Carlo integrators is the way to go for implementing new resummation codes

Any questions?