Baseline-dependent neutrino oscillations in asymmetrically-warped spacetimes

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European Strategy for Future Neutrino Physics, CERN 01–03 October 2009

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• LSND observed excess of $\bar{\nu}_e$ events in pure $\bar{\nu}_{\mu}$ beam



- with Δm^2_{\odot} & $\Delta m^2_{\rm atm} \Rightarrow$ 3 Δm^2 's \Rightarrow 4 neutrinos
- LSND excess might hint towards deviations from usual oscillation mechanism...
 - oscillations into new "sterile" neutrinos?
 - extra dimensions? [SH, Micu, Päs, Weiler: arXiv:0906.0150 [hep-ph]]
 - CPT & Lorentz violation? [SH, Micu, Päs: arXiv:0906.5072 [hep-ph]]

• altered dispersion relations: "small", effective deviations from common energy-momentum-mass relation

$$E = \sqrt{p^2 + m^2} + V$$

• write Hamiltonian governing oscillations via

$$Ht = H_{\text{diag}}t + \delta(Ht)$$
$$= H_{\text{diag}}t + t\delta H + H\delta t$$

- neutrino oscillations driven by three phase e^{-iHt} differences:
 - diag part H_{diag} : no influence on oscillation probability
 - common phase difference δH : non-degenerate neutrino masses \Rightarrow standard flavour oscillations
 - unconventional phase difference δt : different travel times for different neutrino states

• asymmetrically–warped 5D spacetime

$$ds^{2} = dt^{2} + e^{-2k|u|} dx_{i} dx^{i} - du^{2}$$

- our Minkowskian brane located at $\boldsymbol{u}=\boldsymbol{0}$
- the deeper you dive into the bulk the more your spatial line element gets warped
- greater warp = less travel time



• different travel times for neutrinos on brane & in bulk \Rightarrow phase difference δt in propagation

Neutrino oscillations in a brane-bulk scenario

- two-state active-sterile system with active neutrino confined to brane; sterile neutrino (gauge singlet) may propagate through bulk
- altered dispersion relations encapsulated in potential V prop to "shortcut parameter" $\epsilon \equiv \frac{\delta t}{t}$
- diagonalize Hamiltonian H+V via new mixing angle expressible in terms of vacuum values

$$\tan 2\tilde{\theta} = \frac{\tan 2\theta}{1 - \frac{E^2}{E_{\text{res}}^2}} \quad \text{with} \quad E_{\text{res}} = \sqrt{\frac{\Delta m^2 \cos 2\theta}{2\epsilon}}$$

- resonant oscillations with three energy domains
 - $\begin{array}{ll} E \gg E_{\rm res} & {\rm oscillations \ suppressed} & \tilde{\theta} \to 0 \\ E = E_{\rm res} & {\rm maximal \ mixing} & \tilde{\theta} = \frac{\pi}{4} \\ E \ll E_{\rm res} & {\rm vacuum \ oscillations} & \tilde{\theta} \to \theta \end{array}$

- different modes = sterile's path intersecting the brane an n^{th} time at fixed baseline L \Rightarrow replace l = L/n
- each mode = one possible path; QM: sum over all possible paths \Rightarrow insert e^{iS_n}
- sort out geodesics which do not intersect at L \Rightarrow insert measure Δn
- assume Gaußian distribution of initial momenta p_u about brane \Rightarrow insert distribution with "inverse brane thickness" $\sigma \ge \frac{1}{2\Delta u}$

$$A(\nu_{\rm a} \to \nu_{\rm s}) = \sum_{n=1}^{\infty} \Delta n \ e^{iS_{cl}(n)} \frac{vn}{(n^2 + v^2)^{3/2}} \left[\sqrt{\frac{2}{\pi}} \frac{\beta E}{\sigma} \ e^{-\frac{(\beta E v)^2}{2\sigma^2(n^2 + v^2)}} \right] \sin 2\tilde{\theta}_n \ \sin \frac{L\delta \tilde{H}_n}{2}$$

The Weights

• resonance energy; shortcut parameter; "scaling variable" for mode n

$$\frac{E_{\mathsf{res}}^2(n)}{\Delta m^2} = \frac{\cos 2\theta}{2\epsilon_n} \qquad ; \qquad \epsilon_n = 1 - \left(\frac{n}{v}\right) \operatorname{arcsinh}\left(\frac{v}{n}\right) \qquad ; \qquad v = \frac{kL}{2}$$

• experimentally relevant case: The Near Zone

$$\frac{E_{\rm res}^2(n)}{\Delta m^2} \sim \frac{({\rm MeV})^2}{{\rm eV}^2} \gg 1 \quad \Leftrightarrow \quad v \ll 1 \quad {\rm id \ est} \quad L \ll k^{-1}$$

• Near Zone pheno: a novel "resonant product of L and E"

$$\frac{E^2}{E_{\rm res}^2(n)} = \frac{(LE)^2}{n^2 (LE)_{\rm res}^2} \qquad \text{with} \qquad (LE)_{\rm res} = k^{-1} \sqrt{12 \ \Delta m^2 \cos 2\theta}$$

The Near Zone

• plot $P_{as}(L)$ using Near Zone pheno \Rightarrow resonance peaks at different baselines $L = n(LE)_{res}/E!$



red – 5D osc prob; green – phase averaged 5D osc prob; blue – standard 4D osc prob $\sin^2 2\theta = 0.003$, $k = 5/(10^8 \text{ m})$, E = 15 MeV, $\Delta m^2 = 64 \text{ eV}^2$, and $\sigma = 100 \text{ eV}$; resonances found at multiples $L = n(LE)_{\text{res}}/E = 37n \text{ m}$, n = 1, 2, 3...

- three domains at hand
 - $\begin{array}{ll} \ LE \gg (LE)_{\rm res} & {\rm oscillations\ suppressed} & \tilde{\theta} \to 0 \\ \ LE = (LE)_{\rm res} & {\rm maximal\ mixing} & \tilde{\theta} = \frac{\pi}{4} \\ \ LE \ll (LE)_{\rm res} & {\rm vacuum\ oscillations} & \tilde{\theta} \to \theta \end{array}$
- LSND resonant excess + LSND baseline \Rightarrow resonance condition $LE = (LE)_{res}$ fulfilled \Rightarrow explains resonant excess
- above LSND domain $LE \gg (LE)_{res}$: active & sterile neutrinos decouple \Rightarrow no discernible traces of brane-bulk physics in atmospheric data
- below LSND domain $LE \ll (LE)_{res}$: vacuum oscillations prevail
- Final question:

Is the *LE* resonance phenomenology indeed compatible with the world's oscillation data?