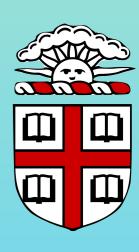
Estimating the local DM content using Gaia DR2

(based on 1807.xxxxx w/ JiJi Fan and John Leung)

Jatan Buch (Brown U.)







Contents:

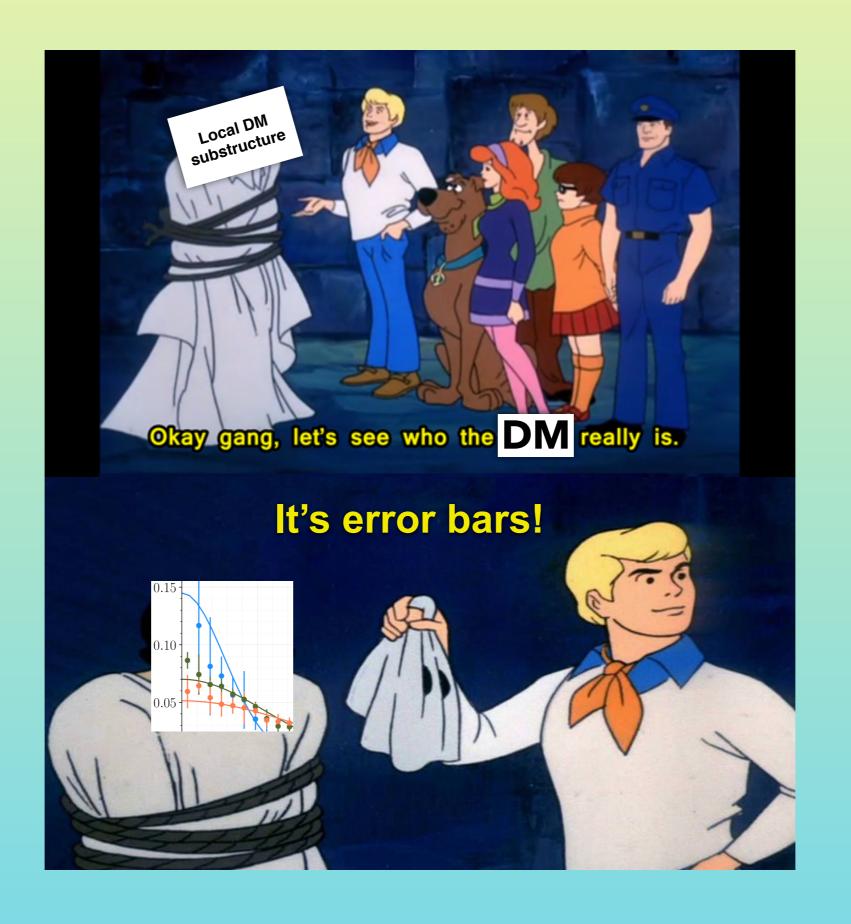
- Overview
- Gaia Data Release 2 (DR2): data, number density, velocity distribution
- Kinematic analysis
- Results: local DM density, dark disks (DD)?!
- Takeaways

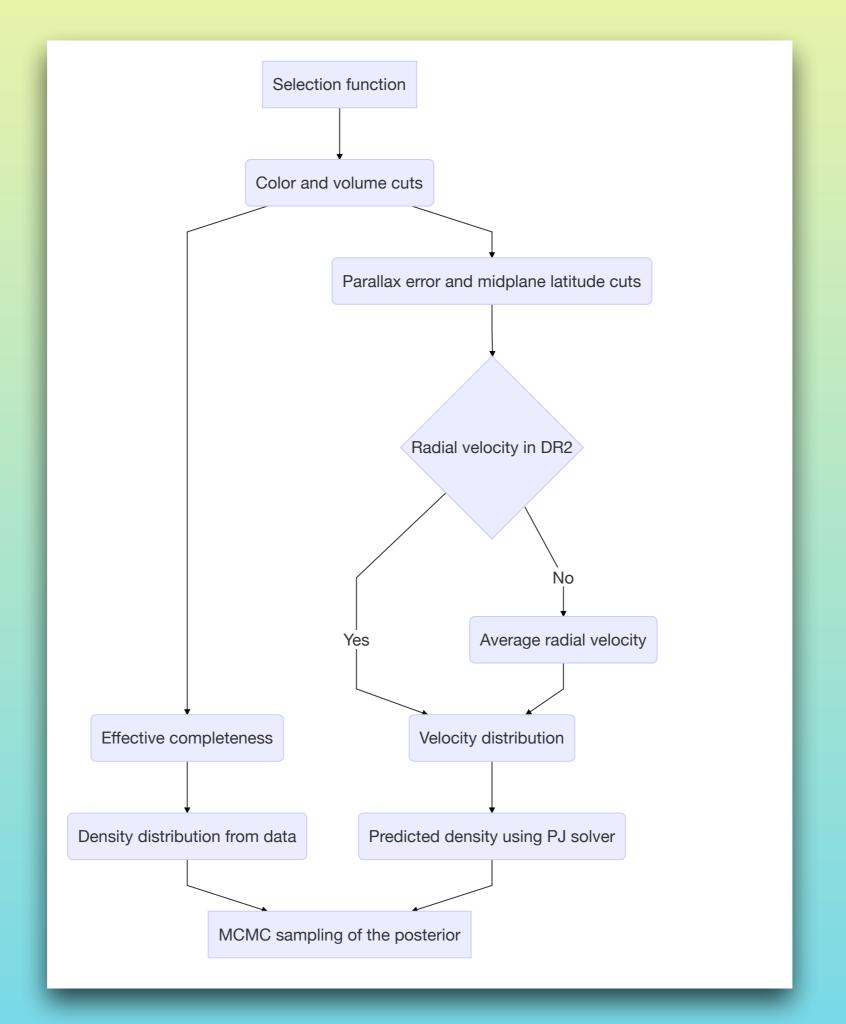
Overview

Motivation

- Density and morphology of local DM density is an important ingredient in direct detection searches on Earth: LZ, Xenon1T, PandaX ...
- Also plays a crucial role in computing the flux of charged cosmic rays produced in DM annihilations. Of course, there's other complicating factors as well [talks by Perez & Cholis on Friday].
- From a theoretical perspective, DM with dissipative self-interactions, for instance U(1)_D, can cool down like baryons and form compact objects (substructure). Depending on the specifics of the astrophysical modeling, various signatures have been proposed: [Fan et al '13; Agrawal & Randall '16; Ghalasi & McQuinn '17; Buckley & DiFranzo '17]

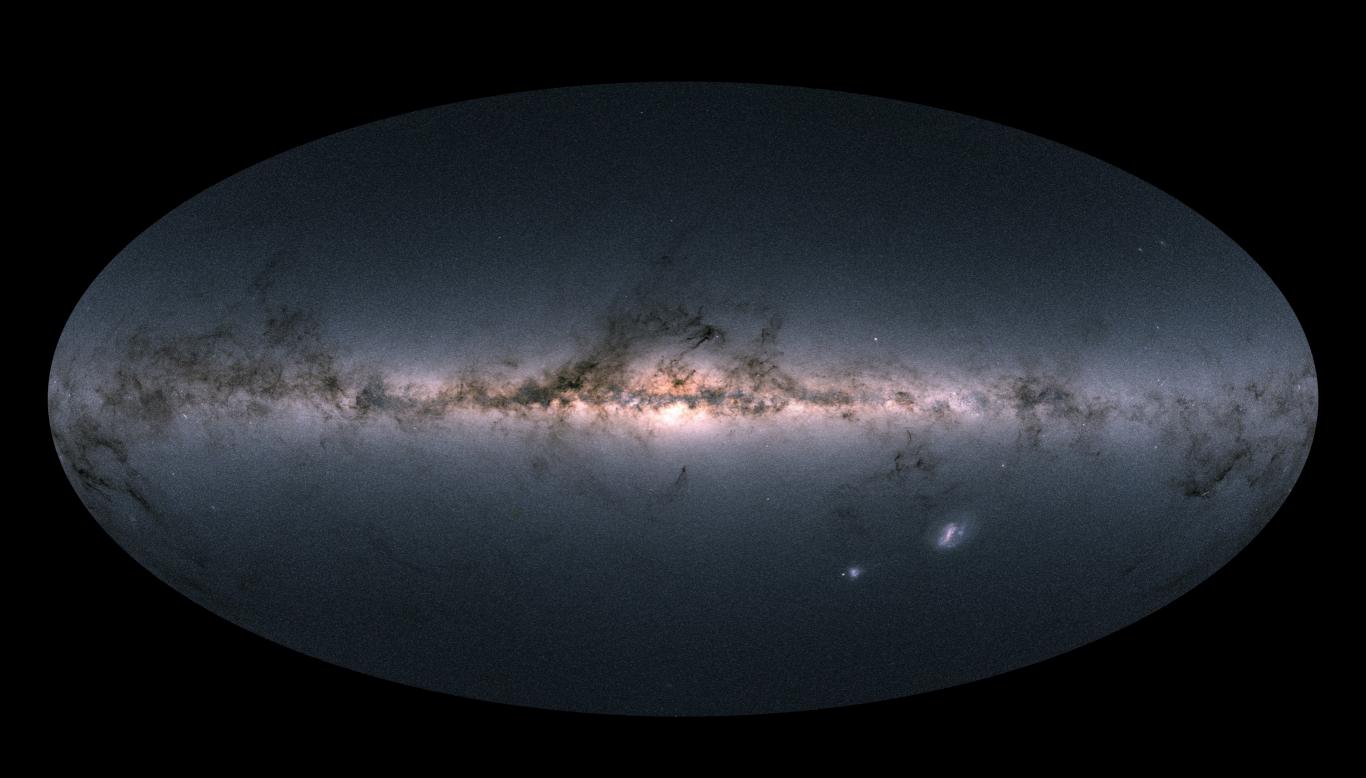
Central question: Can we set *realistic* constraints on the density in DM substructure in the solar neighborhood using current dynamical methods?



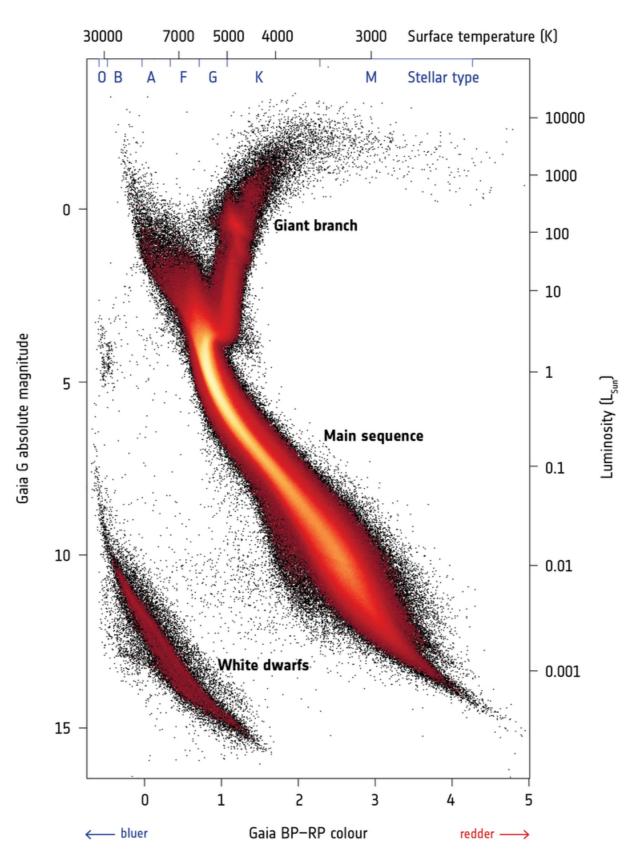


Gaia DR2:

Data, number density, velocity distribution



→ GAIA'S HERTZSPRUNG-RUSSELL DIAGRAM



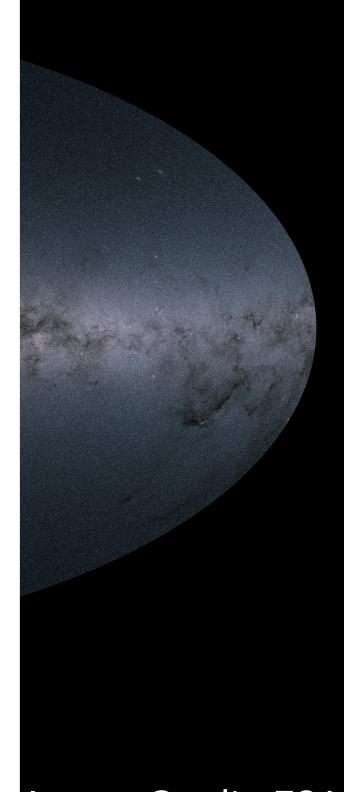
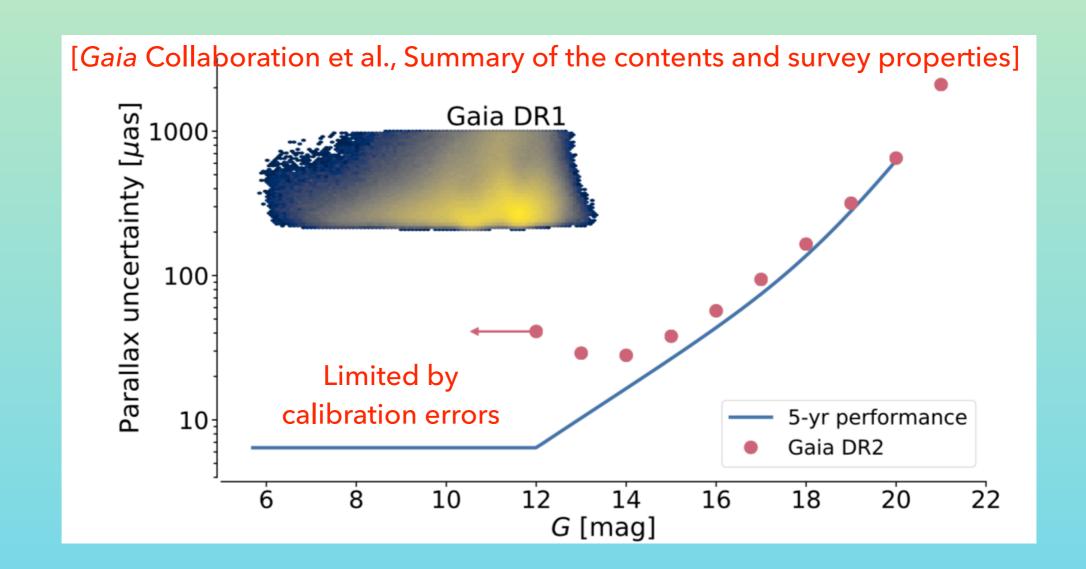


Image Credit: ESA

Some numbers...

- Gaia DR2 provides photometry, and high-precision astrometric data for \approx 1.7 billion sources. The 5-parameter astrometric solution $(\alpha, \delta, \mu_{\tilde{\alpha}}, \mu_{\delta}, \varpi)$ is available for \approx 1.3 billion sources.
- How much has the precision improved over DR1?

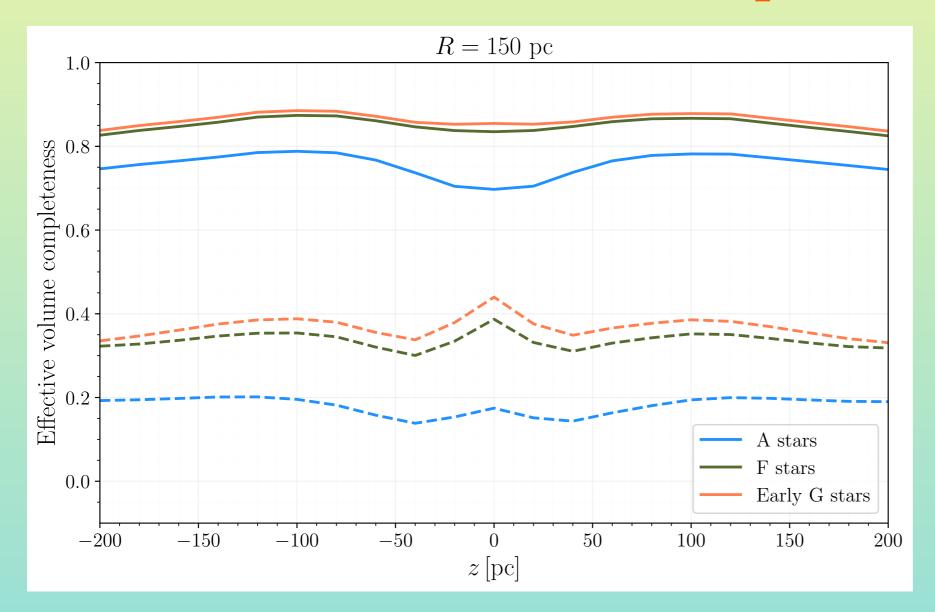


Effective completeness

- The survey has a limiting magnitude $G\approx21$, bright limit $G\approx2$, and is essentially complete between $G\approx12$ and $G\approx17$. While this is definitely an improvement over TGAS, we still need to use an external catalog for constructing a volume complete number density of stars.
- We query the *Gaia* archive for stars in DR2 (full data is ~550 GB*) cross-matched with *2MASS* and apparent magnitude J < 14, and calculate the effective completeness using the gaia_tools_package [Bovy '17].
- 2MASS also provides color information (J, K_s) for DR2 stars, which we use for classifying stars into different spectral types: A, F, early G. An advantage of using (J- K_s) instead of Gaia colors (G_{BP} - G_{RP}) is that these are in the infrared spectrum and only weakly affected by scattering due to interstellar dust.

^{*} If you're interested in working with DR2, I'd be happy to share ideas about handling data.

Effective completeness



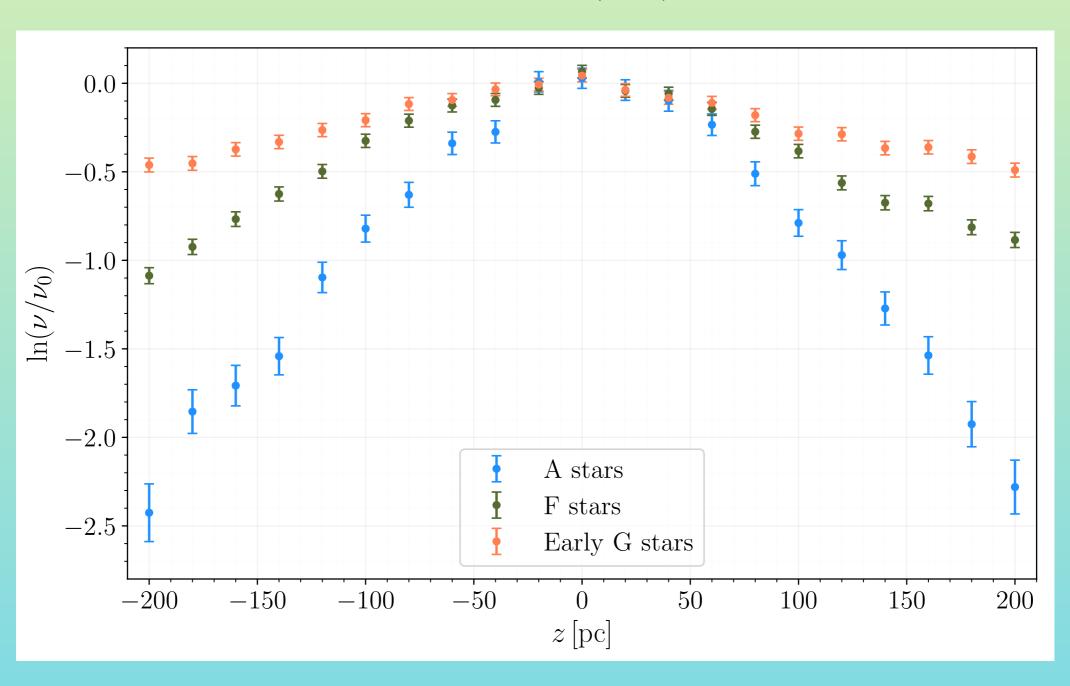
We define the local solar neighborhood as a heliocentric cylinder of radius R=150 pc and half-height z=200 pc.

There's ~2.5x improvement in statistics in the local neighborhood using DR2.

Data set		Gaia DR2		TGAS	
Type	Subtype	Total	Midplane	Total	Midplane
A	A0-A9	4544	310	1729	182
F	F0-F9	38431	2213	16789	1308
Early G	G0-G3	44075	2166	18653	1205

Vertical Number Density

$$z(\text{kpc}) = \frac{\sin b}{\varpi(\text{mas})}$$



Midplane velocity distribution

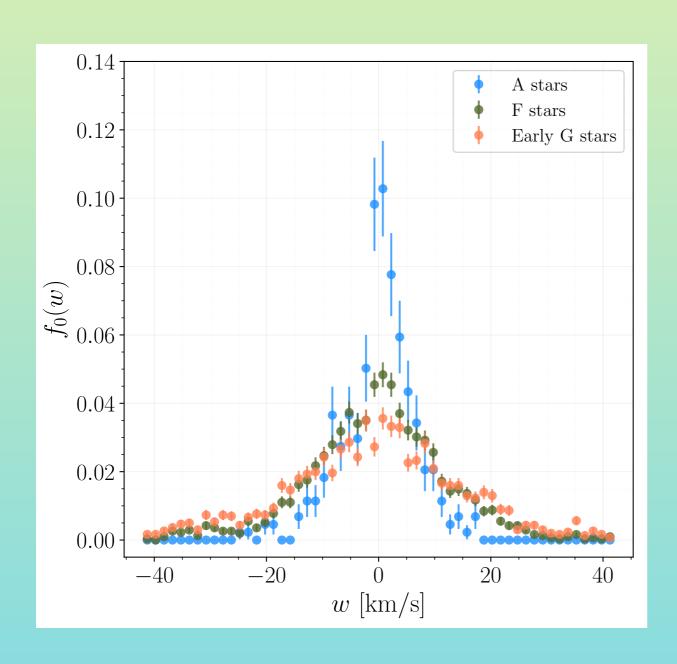
Meanwhile, the vertical velocity of a star is given by,

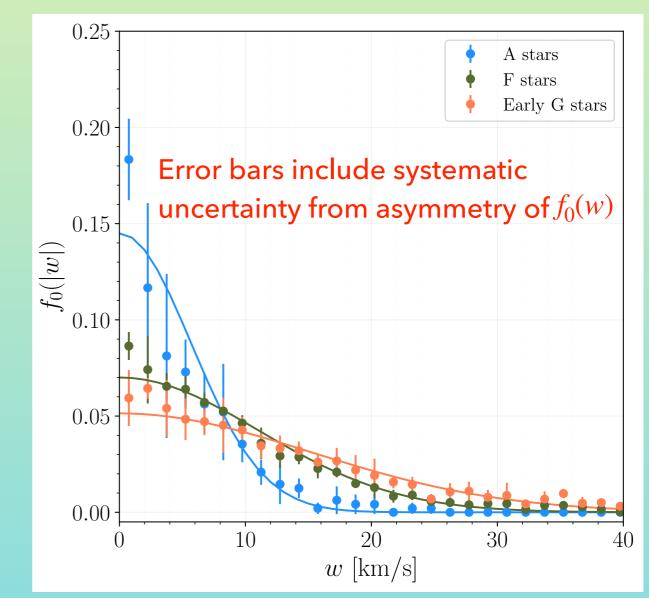
$$w = w_{\odot} + \frac{\kappa \mu_b}{\varpi} \cos b + v_R \sin b,$$

- DR2 also contains line-of-sight radial velocities (RVs) for ≈7.2 million stars measured by its on-board spectrometer. For context, RAVE DR5 presented spectra for ≈450,000 stars.
- Unfortunately, we only have RVs for \approx 2% A, \approx 53% F, and \approx 62% G stars near the midplane (|z| < 20 pc). Thus, we define the midplane using a latitude cut, |b| < 5°, and use an approximation for the mean RV when a star has no RV data in DR2,

$$\langle v_R \rangle = -u_{\odot} \cos l \cos b - v_{\odot} \sin l \sin b - w_{\odot} \sin b,$$

Midplane velocity distribution





Kinematic analysis

Discussion based on: [Flynn & Fuchs '92; Holmberg & Flynn '98; Kramer & Randall '16; Schutz et al. '17]

Poisson-Jeans theory

- The procedure for obtaining the tracer density is straightforward:
 - a) choose a mass model for baryons (gas, stars, and stellar remnants), DM contribution from the halo, and other exotic DM component,
 - b) calculate the local galactic potential of these ingredients, and
 - c) compute the tracer density as a function of the potential.
- To obtain the potential in part b), we solve the Poisson eq.

$$\nabla^2 \Phi = \frac{\partial^2 \Phi}{\partial z^2} + \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial \Phi}{\partial r} \right) = 4\pi G \rho_{\text{tot}},$$

with,

$$\frac{1}{r}\frac{\partial}{\partial r}\left(r\frac{\partial\Phi}{\partial r}\right) \approx (3.4 \pm 0.6) \times 10^{-3} \text{ M}_{\odot}/\text{pc}^3$$
 [Bovy '16]

Poisson-Jeans theory

The total mass density is given by,

$$\rho_{\text{tot}} = \sum_{i=1}^{N_b} \rho_i(0) e^{-\Phi/\sigma_{z;i}^2} + \rho_{\text{DM}} + \rho_{\text{substructure}}(z)$$

where the sum is over N_b components of the Bahcall model that consists of a set of *isothermal* components of baryons characterized by their midplane densities $\rho_i(0)$ and vertical velocity dispersion $\sigma_{z,i}^2$.

The exponential dependence on the potential is due to the vertical Jeans equation, derived by integrating the Boltzmann equation assuming each population is in *equilibrium*,

$$\frac{1}{r\nu_{i}}\frac{\partial}{\partial r}\left(r\nu_{i}\sigma_{rz;i}\right) + \frac{1}{r\nu_{i}}\frac{\partial}{\partial \phi}\left(\nu_{i}\sigma_{\phi z;i}\right) + \frac{1}{\nu_{i}}\frac{d}{dz}\left(\nu_{i}\sigma_{z;i}^{2}\right) = -\frac{d\Phi}{dz}$$
"Tilt" term

"Axial" term

Poisson-Jeans theory

 We can put all these ingredients together by solving the Boltzmann equation in the z direction for each tracer population,

$$w\frac{\partial f_i}{\partial z} - \frac{\partial \Phi}{\partial z} \frac{\partial f_i}{\partial w} = 0.$$

where $f_i(z, w)$ is the distribution function. Assuming separability of phase space, we can integrate over the velocity to obtain the normalized tracer density,

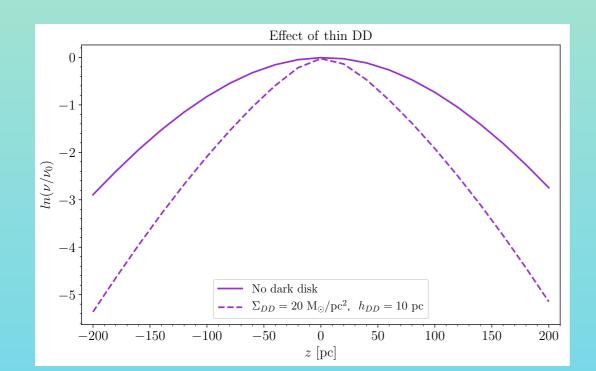
$$\frac{\nu_i(z)}{\nu_i(0)} = 2 \int_0^\infty dw f_{i,z=0}(\sqrt{w^2 + 2\Phi(z)})$$

Results

Data analysis

- We constrain the total matter density by including with the Bahcall model:
 - a) Local DM density $ho_{
 m DM}$,
 - b) Local DM content: $\rho_{\rm DM}$ + thin DD parametrized by,

$$\rho_{DD}(z) = \frac{\Sigma_{DD}}{4h_{DD}} \operatorname{sech}^{2} \left(\frac{z}{2h_{DD}} \right)$$



Parameters	Prior type	Range	Total
$\rho_k(0), \sigma_{z;k}$	Gaussian	Eq. (4.1)	24
$N_{ u_i}$	Uniform	[0.9, 2.0]	3
z_{\odot}	Uniform	[-30.0, 30.0] pc	1
h_{DD}	Uniform	[0.0, 100.0] pc	1
$ ho_{ m DM}$	Uniform	$[0.0, 0.06] \ {\rm M_{\odot}/pc^3}$	1
Σ_{DD}	Uniform	$[0.0, 30.0] \ {\rm M_{\odot}/pc^2}$	1

• Our model \mathcal{M} is characterized by $\theta = \{\psi, \xi\}$, where

 $\psi = \{\rho_{\mathrm{DM}}, \Sigma_{DD}, h_{DD}\}$ are the parameters of interest

ξ are the nuisance parameters, including height of the sun, baryonic uncertainties, ...

Data analysis

 Performing parameter estimation in a Bayesian framework, we sample from the posterior,

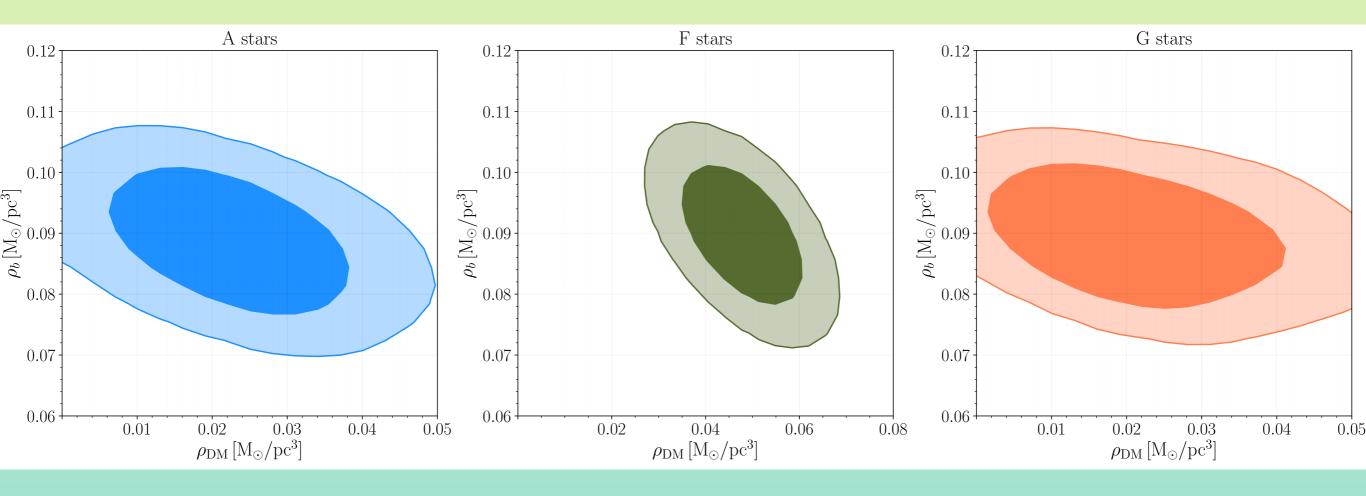
$$p(\boldsymbol{\theta} \mid d) = \frac{p(d \mid \mathcal{M}, \boldsymbol{\theta}) p(\boldsymbol{\theta} \mid \mathcal{M})}{Z^*}$$

using the Markov Chain Monte Carlo (MCMC) sampler emcee [D. Foreman-Mackey et al. '13]

Note that MCMC methods are samplers and not optimizers, so there
is no one 'true' value for each parameter. Instead, results are quoted
using marginalized posteriors of parameters.

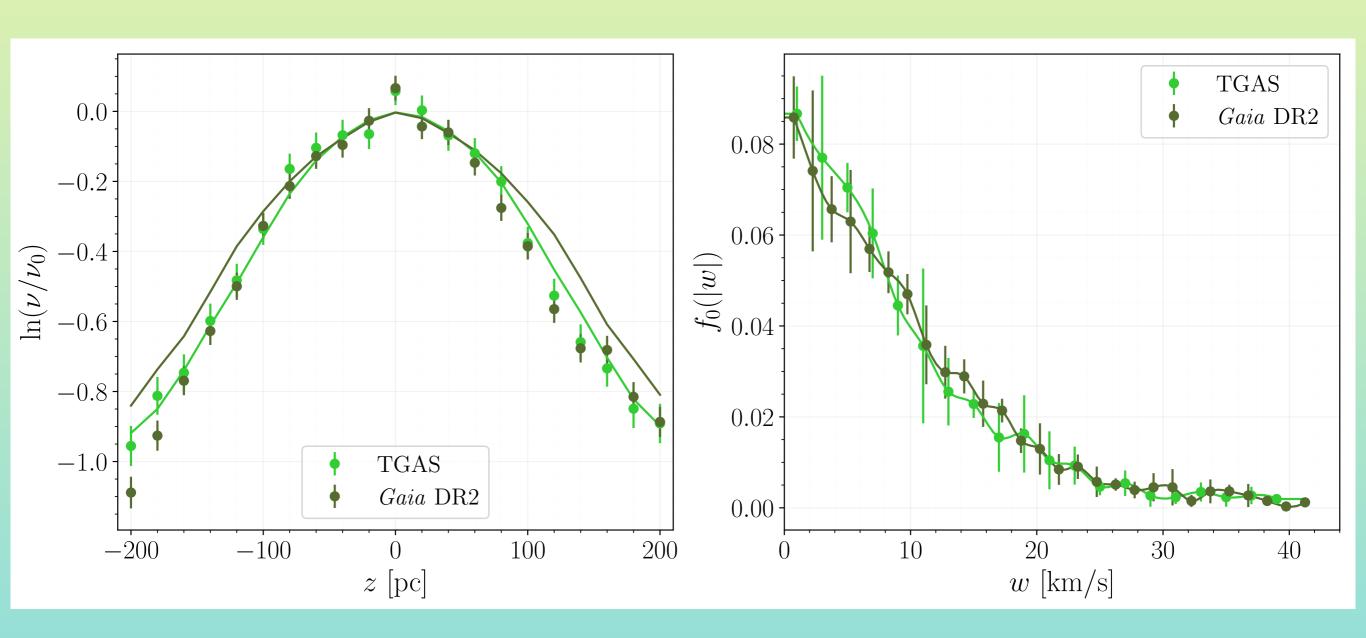
* MCMC samplers don't care about Z, and if you're a Bayesian neither should you! Would love to debate this point in more detail if anyone is interested.

Local DM density



Stellar type	$ ho_{\rm DM}~[{ m M}_{\odot}/{ m pc}^3]$	$\rho_{\mathrm{DM}}~[\mathrm{GeV/cm^3}]$	$\rho_b \; [{\rm M}_{\odot}/{\rm pc}^3]$	$z_{\odot} \; [m pc]$
A stars	$0.023^{+0.010}_{-0.010}$	$0.874^{+0.380}_{-0.380}$	$0.089^{+0.007}_{-0.007}$	$4.95^{+3.78}_{-4.15}$
F stars	$0.047^{+0.006}_{-0.007}$	$1.786^{+0.228}_{-0.266}$	$0.091^{+0.007}_{-0.006}$	$2.52^{+2.58}_{-2.74}$
G stars	$0.021^{+0.014}_{-0.011}$	$0.798^{+0.532}_{-0.418}$	$0.090^{+0.007}_{-0.007}$	$-8.46^{+4.61}_{-4.09}$

Local DM density



- The DR2 midplane velocity distribution has a more gradual falloff as compared to TGAS that results in a broader predicted density. Raises issues regarding the robustness of the method!
- Broader prediction -> accommodates more matter -> weaker constraints

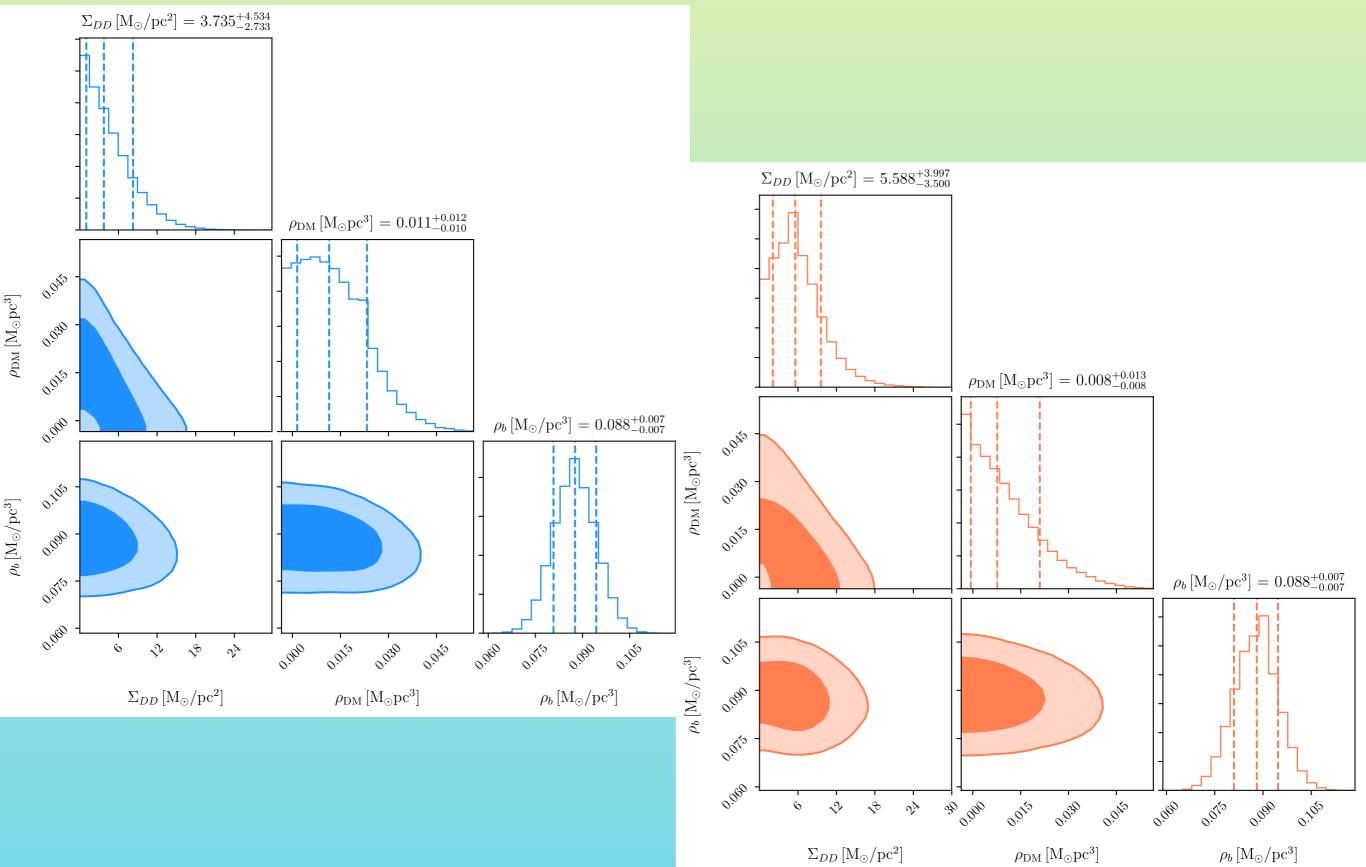
Local DM density

Our results are consistent with previous measurements:

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\rho_{\rm DM} = 0.012^{+0.001}_{-0.002} \, \rm M_{\odot}/pc^3 \quad (within 1\sigma) \quad [Sivertsson et al. '17]
\rho_{\rm DM} = 0.008^{+0.003}_{-0.003} \, \rm M_{\odot}/pc^3 \quad (within 2\sigma) \quad [Bovy \& Tremaine '12]
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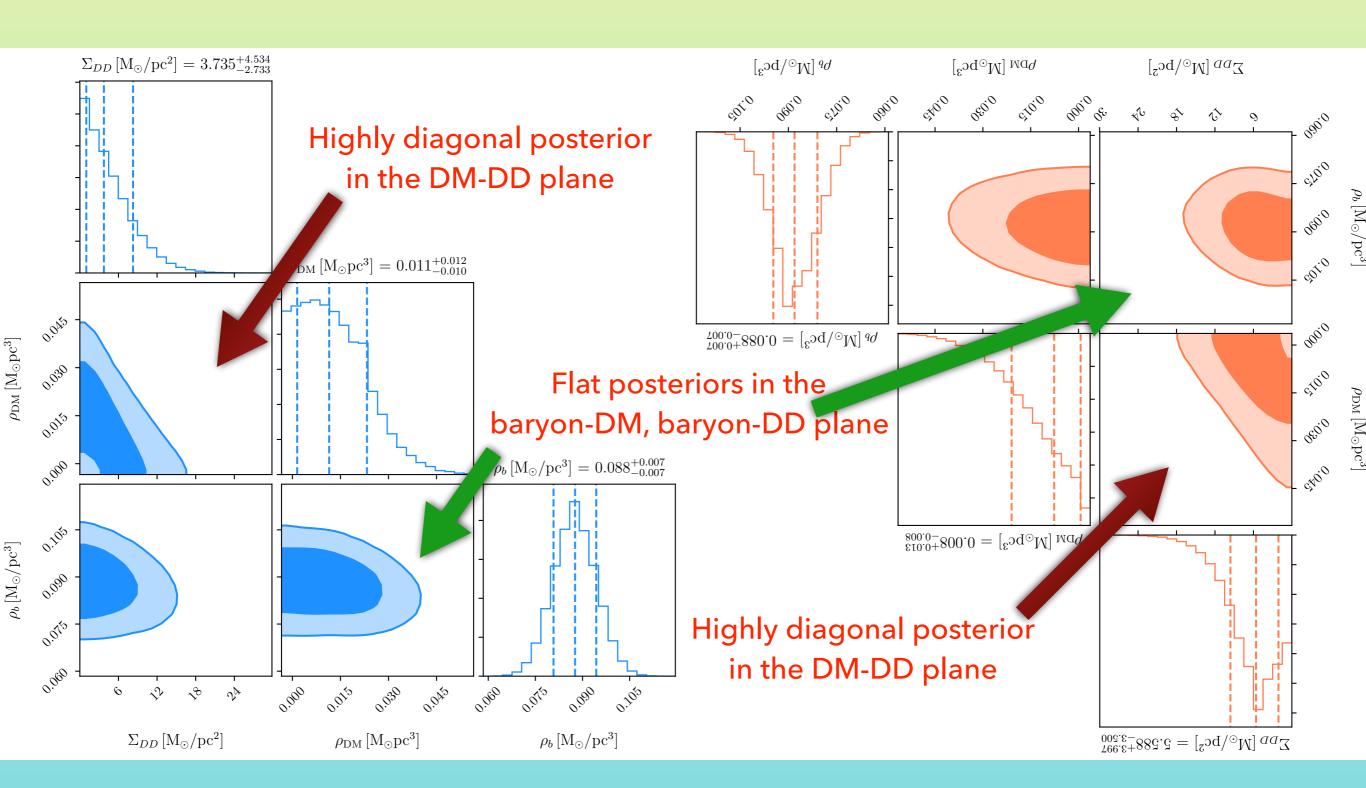
- Notice that the error bars are fairly large in our case. While poorly
 modeled systematics can be a culprit, the posteriors indicate a high level
 of degeneracy between baryons and DM.
- Indeed, as first pointed out by Bahcall (1992) and shown via detailed N-body simulations by Garbari et al. (2011), this degeneracy can only be broken by including the density falloff at z > 1 kpc.
- Since the baryons are mostly confined to the stellar disk with a scale height of ~kpc, any falloff at high z can be attributed to (atleast at leading order) to DM, leading to a more precise measurement with smaller error bars.

Local DM content



Central question: Can we set *realistic* constraints on the density in DM substructure in the solar neighborhood using current dynamical methods?

Local DM content



Answer: Maybe, but ...

Caveats

- Better understanding and modeling of how disequilibria [talk by Necib today] affects dynamics in the solar neighborhood.
- Identifying good tracers that incorporate information about age, and sensitivity to non-equilibrium dynamics: using mono-abundant populations (MAPs) [Lee et al. '11, Bovy et al. '12, Banik et al. '16].
- Need physical observable(s) to break the degeneracy between DM and substructure; ratios, hierarchical modeling?
- Exploring the effect of dissipative DM interactions using cooling prescriptions [Fan & Rosenberg '17] in simulations—semi-analytic or numerical—is still lacking. For discovery or constraining many dissipative DM scenarios listed in the Overview, their input will be key.

Takeaways

(or what you should be able to remember after the cocktail hour tonight!)

- We estimate, using A stars as tracers, the value of local DM density to be $\rho_{\rm DM} = 0.023 \pm 0.01~{\rm M}_{\odot}/{\rm pc}^3$, and exclude a thin DD with Σ_{DD} greater than $(5-15)~{\rm M}_{\odot}/{\rm pc}^2$ at the 95% confidence level.
- Due to the latent degeneracy between baryons and DM (substructure or otherwise) in the solar neighborhood, hard to match the precision (given unknown systematics) of DM density measurements at high z.
 This also leads to only weak upper bounds on the thin DD parameters.
- Studying (sub)structure in phase space could shed light on (sub)structure in theory space: e.g: multicomponent DM sector, tweaks to the CDM paradigm etc. Developing creative and robust dynamical methods for their study will be crucial.
- We've only begun to tap the potential of Gaia; there's a lot more to look forward to in the coming years!

Thank you.

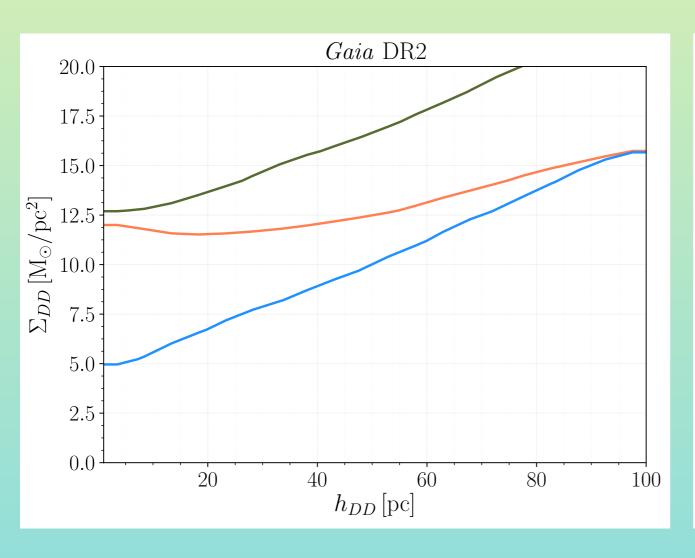
Comments & criticisms welcome!

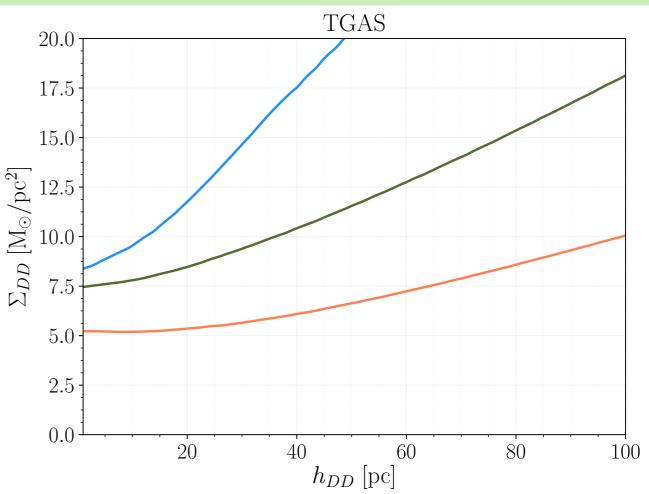
Extra slides

Bahcall model

Baryonic components	$\rho(0) \ [\mathrm{M}_{\odot}/\mathrm{pc}^3]$	$\sigma_z [\mathrm{km/s}]$
Molecular gas (H ₂)	0.0104 ± 0.00312	3.7 ± 0.2
Cold atomic gas $(H_I(1))$	0.0277 ± 0.00554	7.1 ± 0.5
Warm atomic gas $(H_I(2))$	0.0073 ± 0.0007	22.1 ± 2.4
Hot ionized gas $(H_{\rm II})$	0.0005 ± 0.00003	39.0 ± 4.0
Giant stars	0.0006 ± 0.00006	15.5 ± 1.6
$M_V < 3$	0.0018 ± 0.00018	7.5 ± 2.0
$3 < M_V < 4$	0.0018 ± 0.00018	12.0 ± 2.4
$4 < M_V < 5$	0.0029 ± 0.00029	18.0 ± 1.8
$5 < M_V < 8$	0.0072 ± 0.00072	18.5 ± 1.9
$M_V > 8 \text{ (M dwarfs)}$	0.0216 ± 0.0028	18.5 ± 4.0
White dwarfs	0.0056 ± 0.001	20.0 ± 5.0
Brown dwarfs	0.0015 ± 0.0005	20.0 ± 5.0

Consistency with TGAS





 DR2 catalog should be treated as independent from DR1! In particular, there may be significant differences between observations in DR2 and the Tycho-Gaia Astrometric Solution (TGAS) subset of DR1 for some sources.