

A Realistic U(2) Model of Flavor

based on arXiv: 1805.07341 with M. Linster

$$m_{\{u,d,e,\nu\}} \sim \begin{pmatrix} 0 & \varepsilon^2 & 0 \\ \varepsilon^2 & \varepsilon^2 & \{\varepsilon, \varepsilon^2, \varepsilon, \varepsilon^2\} \\ 0 & \{\varepsilon, \varepsilon, \varepsilon^2, \varepsilon^2\} & \{1, \varepsilon, \varepsilon, \varepsilon^2\} \end{pmatrix}$$

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Outline

Realistic extension of original $U(2)$ models by Barbieri & al

R. Barbieri, G. Dvali, L. Hall '96, R. Barbieri, L. Hall, A. Romanino '97, ...

- Differences:**
- viable CKM
 - does not need SUSY
 - Neutrinos

U2

Quarks &
Charged Leptons

Dirac Neutrinos
in $U(2)$

Majorana Neutrinos
in $D_6 \times U(1)$

$U(2)$ Axiflavor

Summary

The Model

Classic *FLASY* Model based on

$$U(2)_F \equiv SU(2)_F \times U(1)_F$$

SM quantum numbers

- compatible with $SU(5)$ GUT
 $10 = Q, U, E \quad \bar{5} = L, E$
- generations transform as $2+1$
 $10_i = 10_a + 10_3$
- need to specify 4 $U(1)_F$ charges
 $X_{10_3} = 0 \quad X_{10_a} = X_{\bar{5}_a} = X_{\bar{5}_3} = 1$

$U(2)$ Breaking with Flavons

- 2 flavons in $2+1$ with charge -1
 $\phi = 2_{-1} \quad \chi = 1_{-1}$
 - VEVs slightly below cutoff Λ
 $\langle \phi \rangle = \begin{pmatrix} \varepsilon_\phi \Lambda \\ 0 \end{pmatrix} \quad \langle \chi \rangle = \varepsilon_\chi \Lambda$
- $\varepsilon_\phi \sim 0.03 \quad \varepsilon_\chi \sim 0.01$

	10_a	$\bar{5}_a$	10_3	$\bar{5}_3$	H	ϕ_a	χ
$SU(2)_F$	2	2	1	1	1	2	1
$U(1)_F$	1	1	0	1	0	-1	-1

[old models: $SO(10)$]

Up-Quark Sector

Invariant Yukawa Lagrangian needs Flavour insertions

$$\begin{aligned} \mathcal{L}_u = & \frac{\lambda_{11}^u}{\Lambda^6} \chi^4 (\phi_a^* Q_a) (\phi_b^* U_b) H + \frac{\lambda_{12}^u}{\Lambda^2} \chi^2 \epsilon_{ab} Q_a U_b H + \frac{\lambda_{13}^u}{\Lambda^3} \chi^2 (\phi_a^* Q_a) U_3 H \\ & + \frac{\lambda_{22}^u}{\Lambda^2} (\epsilon_{ab} \phi_a Q_b) (\epsilon_{cd} \phi_c U_d) H + \frac{\lambda_{23}^u}{\Lambda} (\epsilon_{ab} \phi_a Q_b) U_3 H + \frac{\lambda_{31}^u}{\Lambda^3} \chi^2 Q_3 (\phi_a^* U_a) H \\ & + \frac{\lambda_{32}^u}{\Lambda} Q_3 (\epsilon_{ab} \phi_a U_b) H + \lambda_{33}^u Q_3 U_3 H, \end{aligned}$$

Flavour vevs generate hierarchical Yukawa structure

$$Y_u \approx \begin{pmatrix} \lambda_{11}^u \epsilon_\phi^2 \epsilon_\chi^4 & \lambda_{12}^u \epsilon_\chi^2 & \lambda_{13}^u \epsilon_\phi \epsilon_\chi^2 \\ -\lambda_{12}^u \epsilon_\chi^2 & \lambda_{22}^u \epsilon_\phi^2 & \lambda_{23}^u \epsilon_\phi \\ \lambda_{31}^u \epsilon_\phi \epsilon_\chi^2 & \lambda_{32}^u \epsilon_\phi & \lambda_{33}^u \end{pmatrix} \approx \begin{pmatrix} 0 & \lambda_{12}^u \epsilon_\chi^2 & 0 \\ -\lambda_{12}^u \epsilon_\chi^2 & \lambda_{22}^u \epsilon_\phi^2 & \lambda_{23}^u \epsilon_\phi \\ 0 & \lambda_{32}^u \epsilon_\phi & \lambda_{33}^u \end{pmatrix}$$

drop sub-leading
corrections $\mathcal{O}(\epsilon_\phi^2) \sim 10^{-4}$

Quark and Charged Lepton Sector

Analogous structure in down- and charged lepton sector

$$Y_u \approx \begin{pmatrix} 0 & \lambda_{12}^u \epsilon_\chi^2 & 0 \\ -\lambda_{12}^u \epsilon_\chi^2 & \lambda_{22}^u \epsilon_\phi^2 & \lambda_{23}^u \epsilon_\phi \\ 0 & \lambda_{32}^u \epsilon_\phi & \lambda_{33}^u \end{pmatrix}, \quad Y_d \approx \begin{pmatrix} 0 & \lambda_{12}^d \epsilon_\chi^2 & 0 \\ -\lambda_{12}^d \epsilon_\chi^2 & \lambda_{22}^d \epsilon_\phi^2 & \lambda_{23}^d \epsilon_\phi \epsilon_\chi \\ 0 & \lambda_{32}^d \epsilon_\phi & \lambda_{33}^d \epsilon_\chi \end{pmatrix}$$
$$Y_e \approx \begin{pmatrix} 0 & \lambda_{12}^e \epsilon_\chi^2 & 0 \\ -\lambda_{12}^e \epsilon_\chi^2 & \lambda_{22}^e \epsilon_\phi^2 & \lambda_{23}^e \epsilon_\phi \\ 0 & \lambda_{32}^e \epsilon_\phi \epsilon_\chi & \lambda_{33}^e \epsilon_\chi \end{pmatrix}.$$

I-2 mixing $\sim \epsilon_\chi^2 / \epsilon_\phi^2$

2-3 mixing [d_R & e_L] $\sim \epsilon_\phi / \epsilon_\chi$

2-3 mixing [u_L, u_R, d_L, e_R] $\sim \epsilon_\phi \sim V_{cb}$

I-3 mixing = **(I-2 mixing) × (2-3 mixing)**

Analytics

Texture zeros accurately relate masses and mixing angles

$$|V_{us}| \approx \left| \sqrt{\frac{m_d}{m_s}} \sqrt{c_{23}^{Rd}} - e^{i(\phi_2 - \phi_1)} \sqrt{\frac{m_u}{m_c}} \right|$$

$$|V_{td}| \approx \sqrt{\frac{m_d}{m_s}} \sqrt{c_{23}^{Rd}} |V_{cb}| - e^{i\phi_2} \frac{s_{23}^{Rd}}{c_{23}^{Rd}} \frac{m_s}{m_b}$$

$$|V_{ub}| \approx \left| \sqrt{\frac{m_u}{m_c}} |V_{cb}| - e^{i\phi_1} \sqrt{\frac{m_d}{m_s}} \sqrt{c_{23}^{Rd}} \frac{s_{23}^{Rd}}{c_{23}^{Rd}} \frac{m_s}{m_b} \right|$$

2-3 mixing angle in RH down sector

free phase parameters

In original U(2) models $s_{23}^{Rd} \sim V_{cb}$

$$|V_{ub}/V_{cb}| \approx \sqrt{m_u/m_c}$$

Off by more than 3σ : need $s_{23}^{Rd} \sim 1$ to fix

(but get additional phase dependence)

Numerical Fit

Fit parameters $\{\lambda_{ij}^{u,d,e}, \varepsilon_\phi, \varepsilon_\chi\}$ with SM observables

Use **quality measure for “O(1)-ness”** of $\lambda_{ij}^{u,d,e}$

$$\chi_{O(1)}^2 = \sum_{\lambda_{ij}^p} \frac{(\log(|\lambda_{ij}^p|))^2}{2 \cdot 0.55^2} \longrightarrow \lambda_{ij}^{u,d,e} \text{ with 95\% prob. in } [1/3, 3]$$

e.g. for single parameter $\lambda = \{3, 5, 7, 10, 50, 100\}$
 get contribution $\Delta\chi_{O(1)}^2 = \{2, 4, 6, 9, 25, 35\}$

Fit satisfactory if

$$\chi_{O(1)}^2 \leq \#\text{pars}$$

$$\chi^2 \leq \#\text{obs}$$

Fit	ε_ϕ	ε_χ	min $ \lambda_{ij}^{u,d,\ell} $	max $ \lambda_{ij}^{u,d,\ell} $	χ^2	$\chi_{O(1)}^2$
QL1 _{\mathbb{R}}	0.019	0.008	1/3.1	2.7	1.7	7.8
QL2 _{\mathbb{R}}	0.023	0.008	1/2.7	2.8	12	5.4

Neutrino Sector

Majorana neutrinos don't work because leading order 1-2 entry in Weinberg operator vanishes due to SU(2) antisymmetrization

$$\epsilon_{ab} L_a L_b H H = 0$$

Do Dirac Neutrinos with $N_i = N_a + N_3$ & $X_a^N = X_3^N$



$$m_\nu^D \approx v \epsilon_\chi^{X_a^N - 1} \begin{pmatrix} 0 & \lambda_{12}^\nu \epsilon_\chi^2 & 0 \\ -\lambda_{12}^\nu \epsilon_\chi^2 & \lambda_{22}^\nu \epsilon_\phi^2 & \lambda_{23}^\nu \epsilon_\phi \epsilon_\chi \\ 0 & \lambda_{32}^\nu \epsilon_\phi \epsilon_\chi & \lambda_{33}^\nu \epsilon_\chi^2 \end{pmatrix}$$

get anarchic structure; smallness from largish U(1) charges

Neutrino Fit & Predictions

Good fit just for Normal Ordering

Fit	X_a^N	X_3^N	ε_ϕ	ε_χ	min $ \lambda_{ij}^{u,d,e,\nu} $	max $ \lambda_{ij}^{u,d,e,\nu} $	χ^2	$\chi_{\mathcal{O}(1)}^2$
QL ν_D -1 (NO)	6	6	0.026	0.012	1/2.9	2.6	0.5	10
QL ν_D -2 (NO)	6	6	0.024	0.013	1/2.6	2.2	18	9
QL ν_D -3 (NO)	5	5	0.022	0.006	1/3.1	3.8	1.0	13
QL ν_D -4 (NO)	5	5	0.021	0.006	1/2.5	2.4	18	9

Predict overall mass scale from scanning over successful fits

Quantity	Range [meV]	Preferred values [meV]
$\sum m_i$	58 – 110	60 – 65
m_β	8 – 26	9 – 10

No chance with KATRIN, automatically satisfy PLANCK bound

$$m_\beta = \sqrt{\sum_i m_i^2 |U_{ei}|^2}$$

Majorana Neutrinos from $D_6 \times U(1)$

Consider dihedral $D_6 = D_3 \times Z_2$ instead of $SU(2)$

	10_a	$\bar{5}_a$	10_3	$\bar{5}_3$	H	ϕ_a	χ	
$D_3 \times Z_2$	2_-	2_-	1_+	1_+	1_+	2_-	1_+	$(\psi \otimes \phi)_1 = \psi_1\phi_2 + \psi_2\phi_1$
$U(1)_F$	1	1	0	1	0	-1	-1	$(\psi \otimes \phi \otimes \chi)_1 = \psi_1\phi_1\chi_1 + \psi_2\phi_2\chi_2$

**Mimics $U(2)$ structure except for symmetric 1-2:
Weinberg operator fixed from charged lepton sector**

$$m_u \approx v \begin{pmatrix} 0 & \lambda_{12}^u \varepsilon_\chi^2 & 0 \\ \lambda_{12}^u \varepsilon_\chi^2 & \lambda_{22}^u \varepsilon_\phi^2 & \lambda_{23}^u \varepsilon_\phi \\ 0 & \lambda_{32}^u \varepsilon_\phi & \lambda_{33}^u \end{pmatrix},$$

$$m_d \approx v \begin{pmatrix} 0 & \lambda_{12}^d \varepsilon_\chi^2 & 0 \\ \lambda_{12}^d \varepsilon_\chi^2 & \lambda_{22}^d \varepsilon_\phi^2 & \lambda_{23}^d \varepsilon_\phi \varepsilon_\chi \\ 0 & \lambda_{32}^d \varepsilon_\phi & \lambda_{33}^d \varepsilon_\chi \end{pmatrix}$$

$$m_e \approx v \begin{pmatrix} 0 & \lambda_{12}^e \varepsilon_\chi^2 & 0 \\ \lambda_{12}^e \varepsilon_\chi^2 & \lambda_{22}^e \varepsilon_\phi^2 & \lambda_{23}^e \varepsilon_\phi \\ 0 & \lambda_{32}^e \varepsilon_\phi \varepsilon_\chi & \lambda_{33}^e \varepsilon_\chi \end{pmatrix},$$

$$m_\nu \approx \frac{v^2}{M} \begin{pmatrix} 0 & \lambda_{12}^\nu \varepsilon_\chi^2 & 0 \\ \lambda_{12}^\nu \varepsilon_\chi^2 & \lambda_{22}^\nu \varepsilon_\phi^2 & \lambda_{23}^\nu \varepsilon_\phi \varepsilon_\chi \\ 0 & \lambda_{23}^\nu \varepsilon_\phi \varepsilon_\chi & \lambda_{33}^\nu \varepsilon_\chi^2 \end{pmatrix}$$

Majorana Neutrino Fit

**Neutrino mass matrix automatically anarchic:
better fit with less parameters**

Fit	ε_ϕ	ε_χ	min $ \lambda_{ij}^{u,d,\ell} $	max $ \lambda_{ij}^{u,d,\ell} $	χ^2	$\chi_{O(1)}^2$	M [10^{11} GeV]
QL ν_M -1	0.025	0.009	1/2.8	2.1	0.7	7.9	4.1
QL ν_M -2	0.024	0.009	1/2.6	1.9	18	6.3	3.3

Predict overall mass scale from scanning over successful fits

Quantity	Range [meV]	Preferred values [meV]
$\sum m_i$	59 – 78	60, 70
m_β	8 – 15	9 – 10, 11 – 12
$m_{\beta\beta}^{\max}$	3 – 16	5, 9

Similar to Dirac case, neutrinoless double beta decay hopeless

$$m_{\beta\beta} = |\sum U_{ei}^2 m_i|$$

Testability

- Neutrino sector predictions
- Low-energy UV completion?
- Full-fledged SU(5) model?
- **The U(2) Axiflavor**

$U(1)_F$ spontaneously broken and has QCD anomaly:
Goldstone is QCD axion solving Strong CP [“axiflavor”]

Wilczek '82; Ema, Hamaguchi, Moroi, Nakayama '16; Calibbi, Goertz, Redigolo, RZ, Zupan '16

Predict flavor-violating axion couplings, fix cutoff by Axion DM

The U(2) Axiflavor

→ see yesterday's talk by Chun

$$\mathcal{L}_a = \frac{\partial_\mu a}{2f_a} \bar{f}_i \gamma^\mu \left[C_{f_i f_j}^V + C_{f_i f_j}^A \gamma_5 \right] f_j + \frac{E}{N} \frac{a(x)}{f_a} \frac{\alpha_{\text{em}}}{8\pi} F_{\mu\nu} \tilde{F}^{\mu\nu}$$

→ U(1)_F charges give anomaly coefficients

$$N_{\text{DW}} = 2N = 9$$

$$E/N = 8/3 \quad \{\text{SU}(5)\}$$

→ Flavon breaking scale (-cutoff) sets mass scale

$$f_a \sim \sqrt{\epsilon_\chi^2 + \epsilon_\phi^2} \Lambda$$

$$m_a = 5.7 \mu\text{eV} \left(\frac{10^{12} \text{ GeV}}{f_a} \right)$$

→ Yukawa structure gives axion-fermion couplings

$$C_{f_i f_j}^V = (V_{fL})_{ki} X_{f_k} (V_{fL})_{kj}^*$$

$$C_{f_i f_j}^A = (V_{fR})_{ki}^* X_{f_k^c} (V_{fR})_{kj}$$

The U(2) Axiflavor

$$\begin{aligned}
 C_{u_i u_j}^V &= \frac{\varepsilon_{L,ij}^u - \varepsilon_{R,ij}^u}{9}, & C_{u_i u_j}^A &= \frac{2\delta_{ij} - \varepsilon_{L,ij}^u - \varepsilon_{R,ij}^u}{9}, \\
 C_{d_i d_j}^V &= \frac{\varepsilon_{L,ij}^d}{9}, & C_{d_i d_j}^A &= \frac{2\delta_{ij} - \varepsilon_{L,ij}^d}{9}, \\
 C_{e_i e_j}^V &= -\frac{\varepsilon_{R,ij}^e}{9}, & C_{e_i e_j}^A &= \frac{2\delta_{ij} - \varepsilon_{R,ij}^e}{9},
 \end{aligned}$$

$$\varepsilon_L^u \sim \varepsilon_R^u \sim \varepsilon_L^d \sim \varepsilon_R^e \sim \begin{pmatrix} \lambda^6 & \lambda^5 & \lambda^3 \\ \lambda^5 & \lambda^4 & \lambda^2 \\ \lambda^3 & \lambda^2 & 1 \end{pmatrix}$$

SU(2)/D₆ protects s-d transitions, strongest bound from SN

Coupling	m_a^{\max}/C [eV]	$m_a^{\max, U(2)}$ [eV]	$f_a^{\min, U(2)}$ [GeV]	Constraint
C_{bs}^V	$9.1 \cdot 10^{-2}$	16	$3.6 \cdot 10^5$	$B^+ \rightarrow K^+ a$ [28]
C_{sd}^V	$1.7 \cdot 10^{-5}$	0.58	$9.8 \cdot 10^6$	$K^+ \rightarrow \pi^+ a$ [29]
C_{ee}^A	$3.1 \cdot 10^{-3}$	0.014	$4.1 \cdot 10^8$	WD Cooling [30]
C_N	$3.5 \cdot 10^{-3}$	0.0092	$6.2 \cdot 10^8$	SN1987A [31]

U(2) Axiflavor Phenomenology

U(1)_F broken before inflation:

$$\Omega_{\text{DM}} h^2 \approx 0.12 \left(\frac{6 \mu\text{eV}}{m_a} \right)^{1.165} \theta^2$$

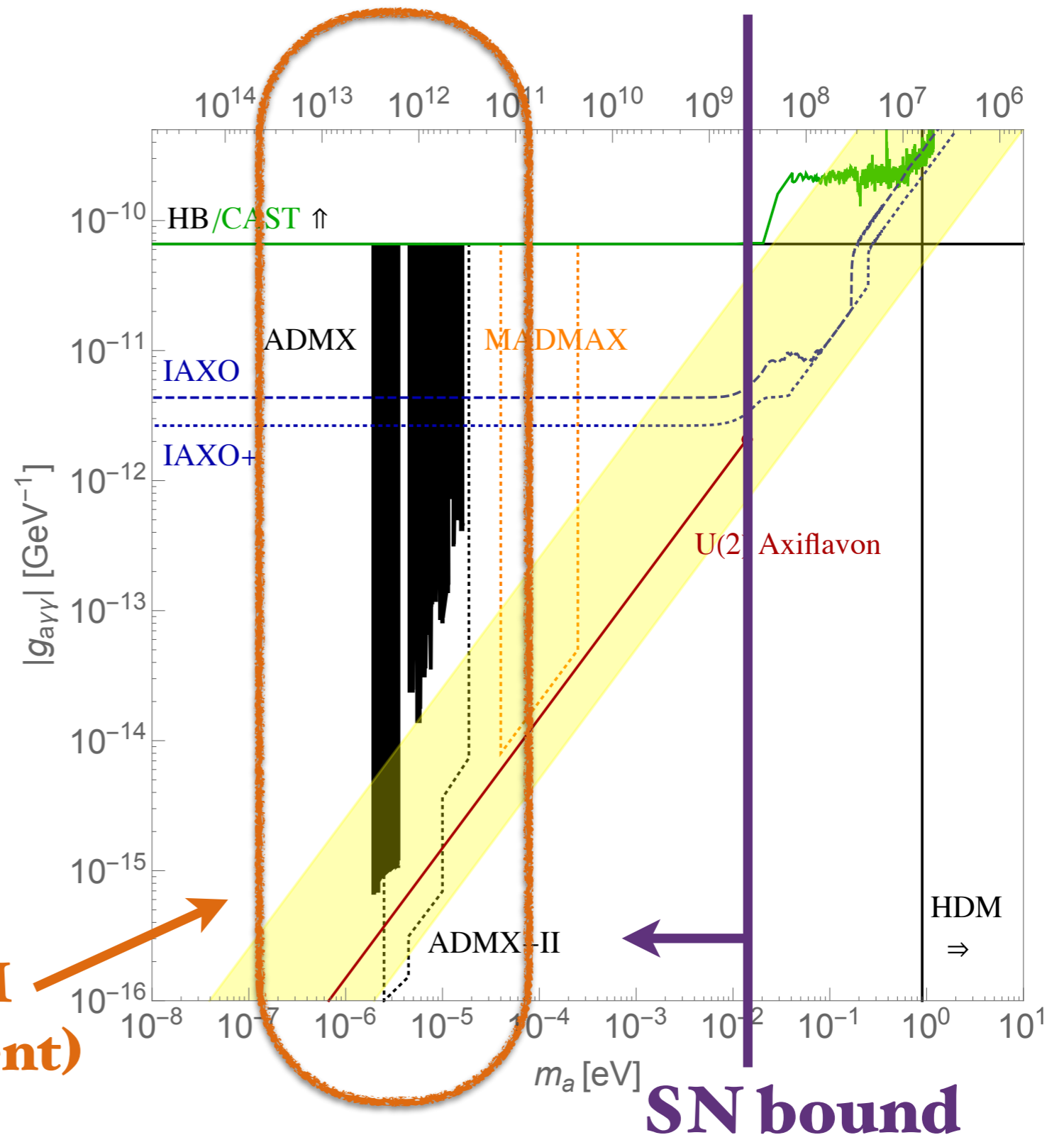
with misalignment angle

$$\theta \in [-\pi, \pi]$$

Axion DM gives preferred range for cutoff

$$\Lambda \sim (10^{13} \div 10^{15}) \text{ GeV}$$

→ see talk by Ringwald



Natural Axion DM window (misalignment)

SN bound

Summary

- Simple, realistic Flavor Model from $U(2)$, works without SUSY, compatible with $SU(5)$
- Dirac Neutrinos in $U(2)$; Majorana Neutrinos in $D_6 \times U(1)_F$, which are automatically anarchic; can predict absolute neutrino mass scale
- Naturally get QCD Axion as DM from $U(1)_F$; fixes breaking /cutoff scale; $SU(2)$ strongly suppresses flavor violation

Backup

Fit Results

Parameter	QL ν_D -1	QL ν_D -2	QL ν_D -3	QL ν_D -4	QL ν_M -1	QL ν_M -2
λ_{12}^u	0.902	0.843	3.831	1.162	-1.633	-1.176
λ_{22}^u	1.187	-1.047	1.859	1.148	1.339	1.112
λ_{23}^u	2.222	-2.175	-2.138	-1.799	2.127	1.925
λ_{32}^u	-1.103	-1.419	1.511	2.422	1.196	1.615
λ_{33}^u	0.787	0.779	-0.787	0.786	0.787	0.785
δ_{33}	-0.640	-0.720	-3.948	-1.097	-3.837	-3.988
λ_{12}^d	0.479	-0.479	2.165	2.173	-0.888	0.976
λ_{22}^d	-1.000	-1.156	-1.075	-0.972	-0.973	0.976
λ_{23}^d	0.913	-0.786	-1.304	-1.155	1.073	0.985
λ_{32}^d	-0.355	0.401	0.414	0.423	0.365	-0.394
λ_{33}^d	0.665	0.651	1.394	1.497	-0.902	-0.948
λ_{12}^ℓ	0.402	-0.376	-1.752	-1.758	-0.801	0.856
λ_{22}^ℓ	0.987	-1.134	1.821	2.052	1.306	1.497
λ_{23}^ℓ	0.343	0.381	0.393	-0.414	-0.368	0.391
λ_{32}^ℓ	-0.992	-1.132	1.175	1.193	-1.198	1.294
λ_{33}^ℓ	0.432	-0.399	-0.945	0.992	-0.503	-0.536
λ_{12}^ν	0.882	-1.416	0.938	1.006	2.130	-1.873
λ_{22}^ν	-0.994	-1.303	0.325	0.398	-0.844	-0.760
λ_{23}^ν	-2.588	-1.074	-1.505	1.681	1.137	-1.078
λ_{32}^ν	1.065	-0.704	0.601	0.680		
λ_{33}^ν	0.952	-1.572	-0.890	0.891	-0.489	-0.655
X_a^N	6	6	5	5		
X_3^N	6	6	5	5		
$v/M \times 10^9$					-0.421	-0.520
ε_ϕ	0.026	0.024	0.022	0.021	0.025	0.024
ε_χ	0.012	0.013	0.006	0.006	0.009	0.009

Flavon Potential

$$\phi = 2_-[-1],$$

$$\chi = 1_+[-1],$$

$$\psi = 1_-[+1],$$

$$\begin{aligned}
 V_{\text{scal}} = & m_\chi^2 |\chi|^2 + (m_\phi^2 + \kappa_\chi |\chi|^2 + \kappa_\psi |\psi|^2) (\tilde{\phi} \cdot \phi) + m_\psi^2 |\psi|^2 \\
 & + \frac{\lambda_1}{4} (\tilde{\phi} \cdot \tilde{\phi})(\phi \cdot \phi) + \frac{\lambda_2}{2} (\tilde{\phi} \cdot \tilde{\phi} \cdot \phi \cdot \phi) + \lambda_3 |\chi|^2 |\psi|^2 + \frac{\lambda_\chi}{2} |\chi|^4 + \lambda_\psi |\psi|^4 \\
 & + \left[\frac{\kappa_1}{2} \psi \psi (\phi \cdot \phi) + \frac{\kappa_2}{2} \chi^* \chi^* (\phi \cdot \phi) + \frac{1}{2} \lambda_{\chi\psi} \psi \psi \chi \chi + \rho \psi (\tilde{\phi} \cdot \phi \cdot \phi) + \text{h.c.} \right]
 \end{aligned}$$

$$\begin{aligned}
 m_\phi^2 = -2m^2, m_\chi^2 = -3/10m^2, m_\psi^2 = 2m^2, \lambda_1 = 1, \lambda_2 = 1/9, \lambda_\chi = 1, \kappa_\chi = -1/8 \\
 \lambda_3 = 2/3, \lambda_{\chi\psi} = -1/20, \kappa_1 = -1/3, \kappa_\psi = 7/10, \rho = -1/20, \lambda_\psi = 9/10, \kappa_2 = 1/2000
 \end{aligned}$$

$$v_1 = 4.6m, \quad v_2 = -3.6 \cdot 10^{-4} m, \quad v_\chi = 1.7m, \quad v_\psi = -2.1 \cdot 10^{-5} m$$