

# Low-Scale seesaw and the CP violation in neutrino oscillations

## Lepton masses and mixing from modular $S_4$ symmetry

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# Leitmotiv: Symmetry predictions for CPV phases



1<sup>st</sup> part

Froggatt-Nielsen



2<sup>nd</sup> part

Modular Symmetry

# Low-Scale seesaw and the CP violation in neutrino oscillations

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arXiv: 1712.09922

in collaboration with S. T. Petcov and T. T. Yanagida

## SM + 2 right-handed neutrinos

$$\mathcal{L} \supset -\bar{\nu}_R \mathbf{Y}_D^T L H^{c\dagger} - \frac{1}{2} \bar{\nu}_R \mathbf{M}_N \nu_L^C + \text{h.c.}$$

$$\mathbf{Y}_D = \begin{pmatrix} 0 & g_{e2} \\ 0 & g_{\mu 2} \\ 0 & g_{\tau 2} \end{pmatrix}$$

$$\mathbf{M}_N = \begin{pmatrix} 0 & M \\ M & 0 \end{pmatrix}$$

Limit of a conserved **lepton number**, with  $\begin{cases} L(\nu_{1R}) = -1 \\ L(\nu_{2R}) = +1 \end{cases}$

### Symmetry-protected scenario

recall: O. Fischer talk on Monday; See also Shaposhnikov, '07; Kersten, Smirnov, '07

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$$\mathbf{M}_N = \begin{pmatrix} 0 & M \\ M & \mu \end{pmatrix}$$

**“linear” component**

similar to texture in  
Malinsky, Romão, Valle, '05

**subleading “inverse” component**

Wyler, Wolfenstein, '83  
Mohapatra, '86  
Mohapatra, Valle, '86

**Symmetry-protected scenario**

recall: O. Fischer talk on Monday; See also Shaposhnikov, '07; Kersten, Smirnov, '07

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$$\mathbf{M}_N = \begin{pmatrix} 0 & M \\ M & \mu \end{pmatrix}$$

$$\mathbf{m}_\nu \simeq -v^2 \mathbf{Y}_D \mathbf{M}_N^{-1} \mathbf{Y}_D^T$$



$M \sim \text{GeV} - \text{TeV}$   
 Pseudo-Dirac Neutrino  
 Ibarra, Molinaro, Petcov  
 1007.2378, 1103.6217

$$(\mathbf{m}_\nu)_{\ell\ell'} \simeq -\frac{v^2}{M} (g_{\ell1} g_{\ell'2} + g_{\ell2} g_{\ell'1}) \sim \frac{v^2}{M} g_1 g_2$$

$v = 174 \text{ GeV}$

$$|g_{e1}| \sim 10^{-12} - 10^{-8}$$

$$|g_{e2}| \sim 10^{-4} - 10^{-2}$$

# A Froggatt-Nielsen realisation

SUSY setup with a broken  $U(1)_{\text{FN}}$ , flavour scale  $\Lambda$  and charges:

	$\hat{S}$	$\hat{N}_1$	$\hat{N}_2$	$\hat{H}_u$	$\hat{L}_e$	$\hat{L}_\mu$	$\hat{L}_\tau$	$\hat{e}^c$	$\hat{\mu}^c$	$\hat{\tau}^c$
$Q_{\text{FN}}$	-1	$n > 0$	-1	0	2	1	1	4	2	0

“Lopsided”  
Sato, Yanagida, '00

$\epsilon \simeq \lambda_C \simeq 0.2$

$\epsilon \simeq 0.06$  also possible

Buchmuller, Yanagida, '99

$$\langle S \rangle = \epsilon \Lambda$$



$$W_\nu \sim \Lambda (\epsilon^{2n} \hat{N}_1 \hat{N}_1 + \epsilon^{n-1} \hat{N}_1 \hat{N}_2) + (\epsilon \hat{L}_e + \hat{L}_\mu + \hat{L}_\tau) (\epsilon^{n+1} \hat{N}_1 + g_2 \hat{N}_2) \hat{H}_u$$

$$\mu \sim \epsilon^{2n} \Lambda \quad M \sim \epsilon^{n-1} \Lambda$$

$$g_{l1} \sim \epsilon^{n+1}$$

## A Froggatt-Nielsen realisation

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$Q_{\text{FN}}$	-1	$n > 0$	-1	0	2	1	1	4	2	0

### A large charge $n$ mimicks $L$ -conservation

$$\mathbf{Y}_{\mathbf{D}} \sim \begin{pmatrix} g_{e1} & \epsilon g_2 \\ g_{\mu 1} & g_2 \\ g_{\tau 1} & g_2 \end{pmatrix} \sin \beta \quad \left\{ \begin{array}{l} |g_{e2}| : |g_{\mu 2}| : |g_{\tau 2}| \simeq \epsilon : 1 : 1 \quad \star \\ |g_{e1}| : |g_{\mu 1}| : |g_{\tau 1}| \simeq ? : ? : ? \end{array} \right.$$

$$|g_{e1}| \ll |g_{e'2}|$$



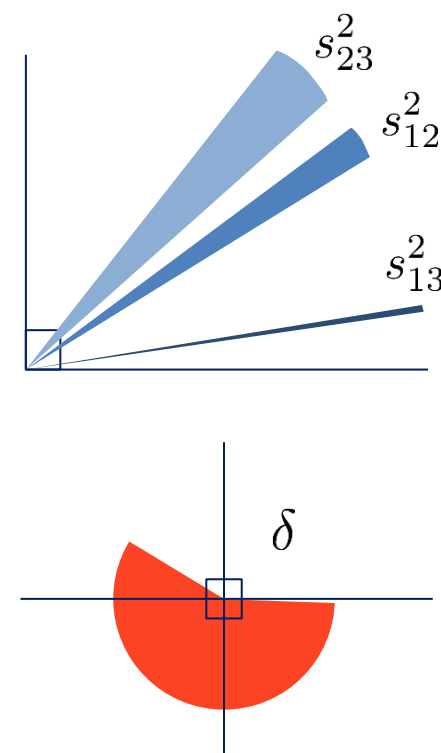
# Interlude: neutrino mixing



For a **NO spectrum**, favoured at the  $3.1\sigma$  CL:

Parameter	Best fit	$1\sigma$ range	$3\sigma$ range
$\Delta m_{\odot}^2 / 10^{-5} \text{ eV}^2$	7.34	7.20 – 7.51	6.92 – 7.91
$ \Delta m_{\text{A}}^2  / 10^{-3} \text{ eV}^2$	2.49	2.46 – 2.53	2.39 – 2.59
$\sin^2 \theta_{12} / 10^{-1}$	3.04	2.91 – 3.18	2.65 – 3.46
$\sin^2 \theta_{13} / 10^{-2}$	2.14	2.07 – 2.23	1.90 – 2.39
$\sin^2 \theta_{23} / 10^{-1}$	5.51	4.81 – 5.70	4.30 – 6.02
$\delta / \pi$	1.32	1.14 – 1.55	0.83 – 1.99

Capozzi, Lisi, Marrone, Palazzo, 1804.09678



Parameterisation:

$$\mathbf{U}_{\text{PMNS}} = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix} \text{diag}(1, e^{i\alpha/2}, 1)$$

a single Majorana phase,  $\alpha_{21} - \alpha_{31} \rightarrow \alpha$

# Predictions for CPV phases

Casas-Ibarra parameterisation (NO):

$$\mathbf{Y}_D = \frac{i}{v} \mathbf{U}_{\text{PMNS}}^* \sqrt{\hat{m}} \mathbf{O} \sqrt{\hat{M}} \mathbf{V}^\dagger$$

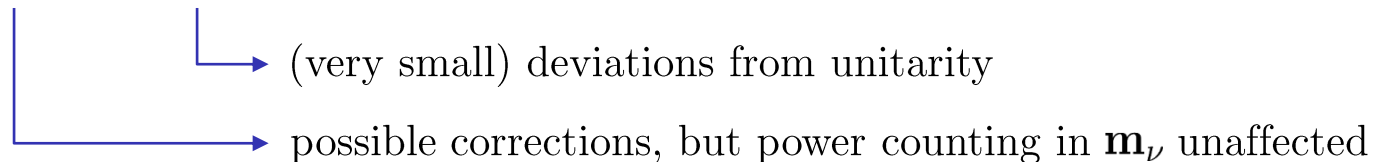
$$\mathbf{O} \equiv \begin{pmatrix} 0 & 0 \\ \cos z & \pm \sin z \\ -\sin z & \pm \cos z \end{pmatrix}, z \in \mathbb{C}$$

$$\mathbf{Y}_D = \begin{pmatrix} g_{e1} & g_{e2} \\ g_{\mu 1} & g_{\mu 2} \\ g_{\tau 1} & g_{\tau 2} \end{pmatrix}$$

small  $|\text{Im } z|$

$$\mathbf{m}_\nu = \mathbf{U}^* \text{diag}(m_1, m_2, m_3) \mathbf{U}^\dagger$$

$$\mathbf{U}_{\text{PMNS}} = \mathbf{U}_l^\dagger (1 + \eta) \mathbf{U} \simeq \mathbf{U}$$


  
 (very small) deviations from unitarity
   
 possible corrections, but power counting in  $\mathbf{m}_\nu$  unaffected

# Predictions for CPV phases

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large  $|\text{Im } z|$  ✓

see also Drewes et al, 1609.09069

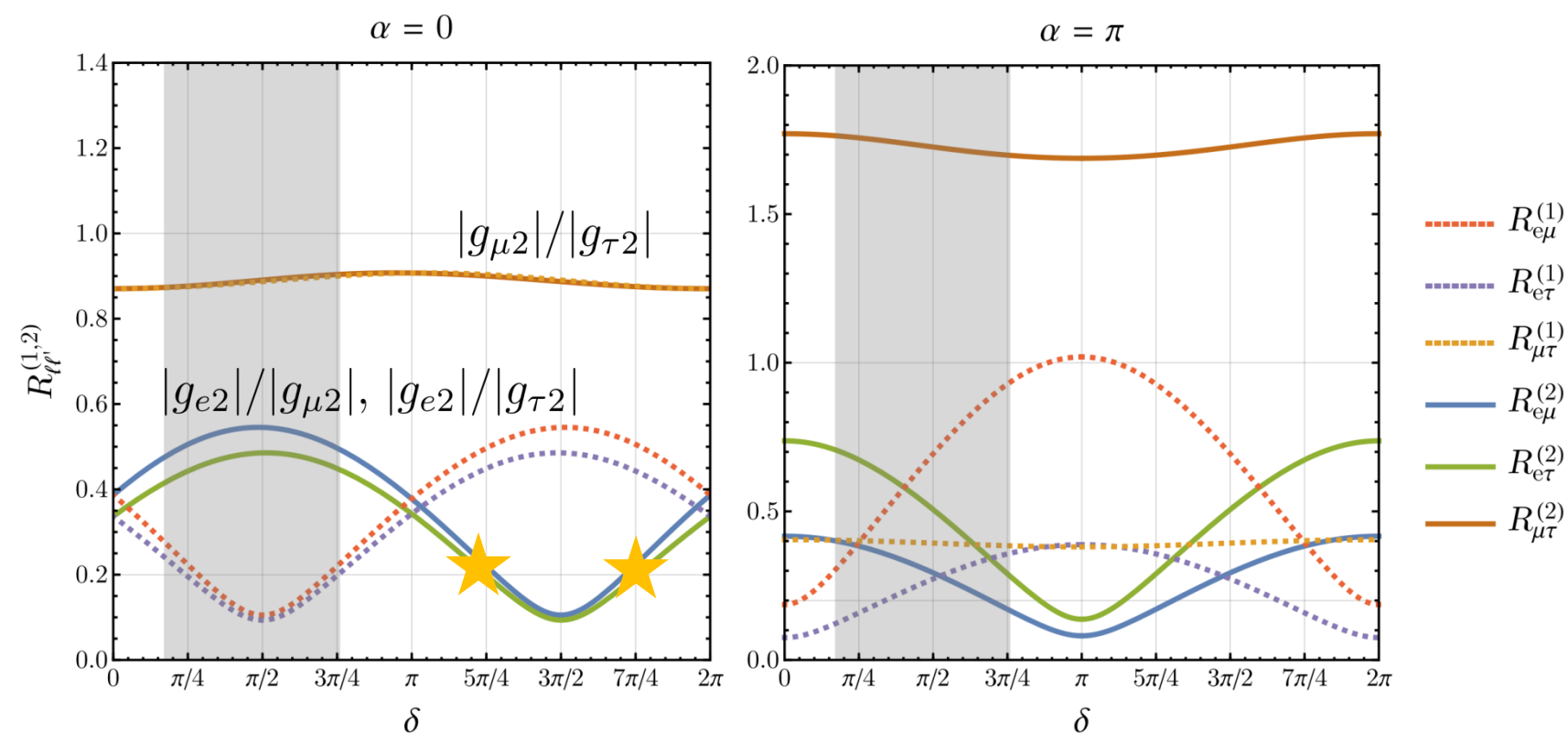
**Ratios of couplings  
independent of the  
Casas-Ibarra parameter**

$$R_{\ell\ell'}^{(1)} \equiv \frac{|g_{\ell 1}|}{|g_{\ell' 1}|} \simeq \frac{|\sqrt{m_2} U_{\ell 2}^* \pm i \sqrt{m_3} U_{\ell 3}^*|}{|\sqrt{m_2} U_{\ell' 2}^* \pm i \sqrt{m_3} U_{\ell' 3}^*|}$$

$$R_{\ell\ell'}^{(2)} \equiv \frac{|g_{\ell 2}|}{|g_{\ell' 2}|} \simeq \frac{|\sqrt{m_2} U_{\ell 2}^* \mp i \sqrt{m_3} U_{\ell 3}^*|}{|\sqrt{m_2} U_{\ell' 2}^* \mp i \sqrt{m_3} U_{\ell' 3}^*|}$$

# Predictions for CPV phases

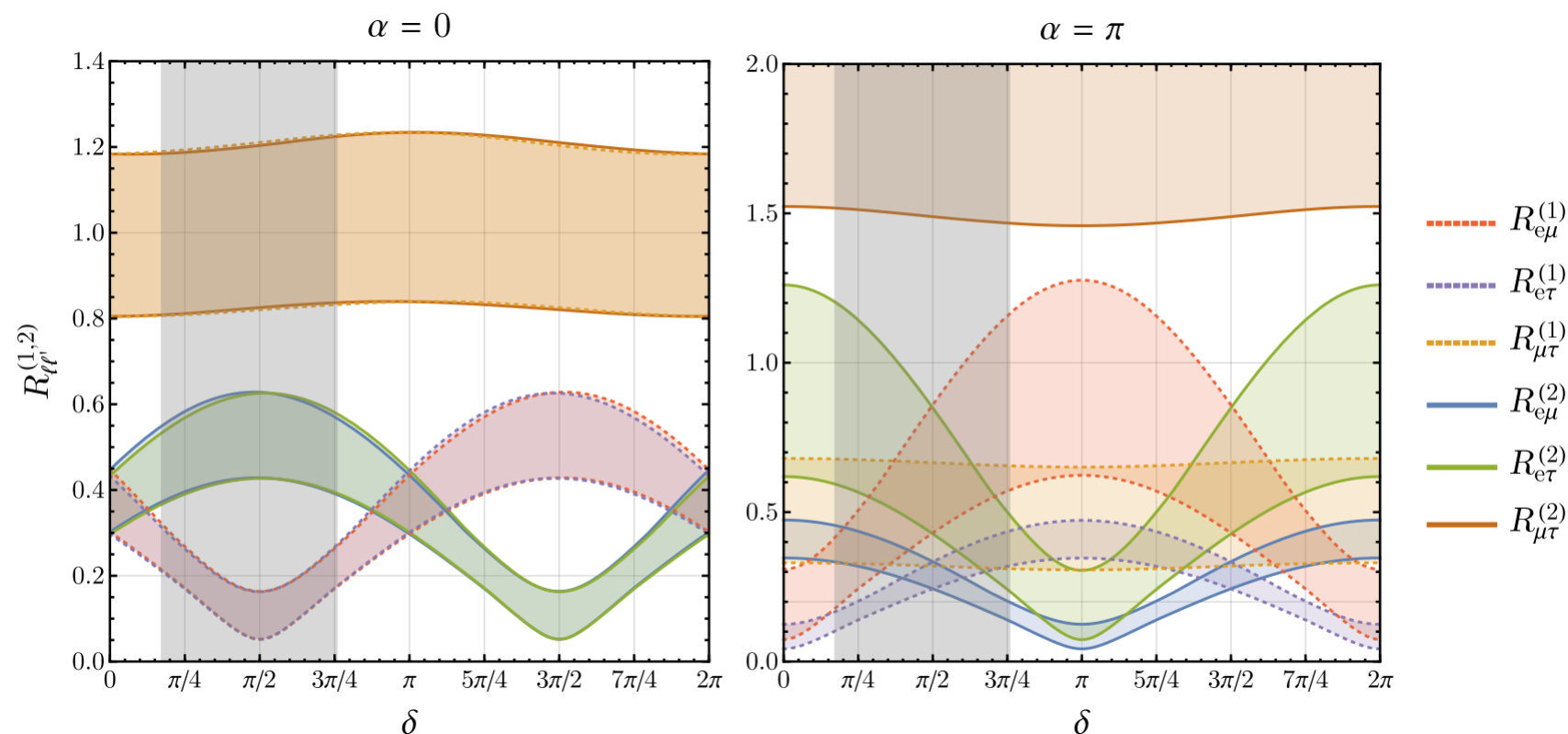
Using best-fit values of mass-squared differences and mixing angles:



$$\epsilon = 0.2 \rightarrow \delta \sim \pm\pi/4, \pm 3\pi/4$$

# Predictions for CPV phases

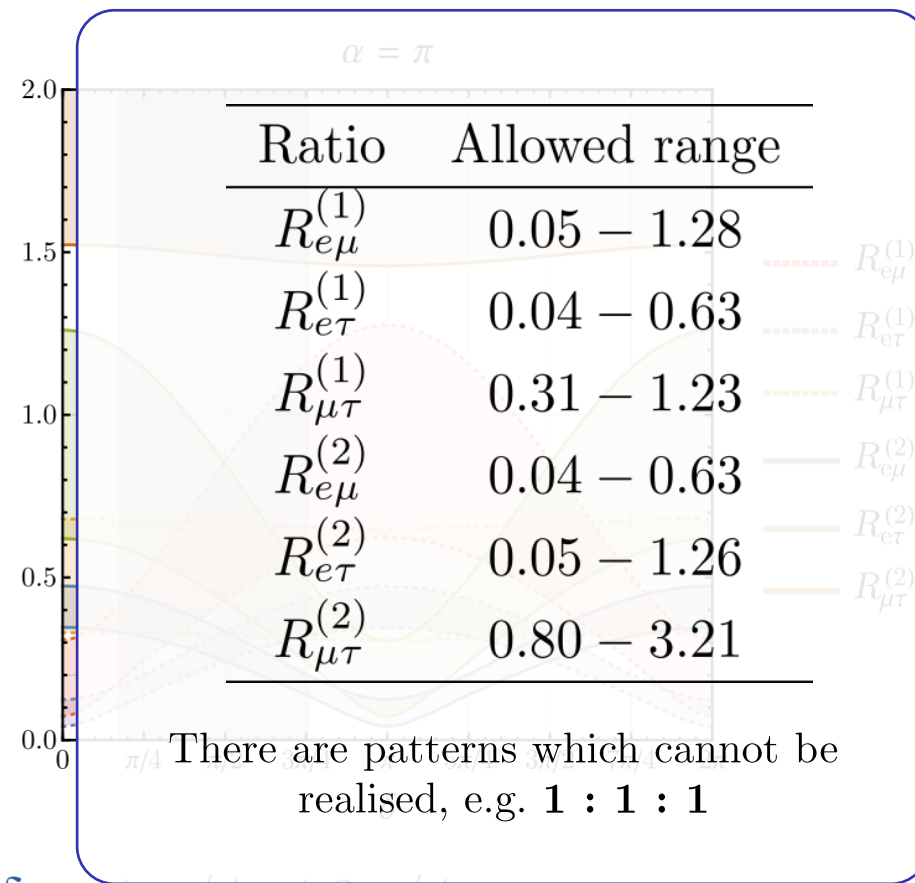
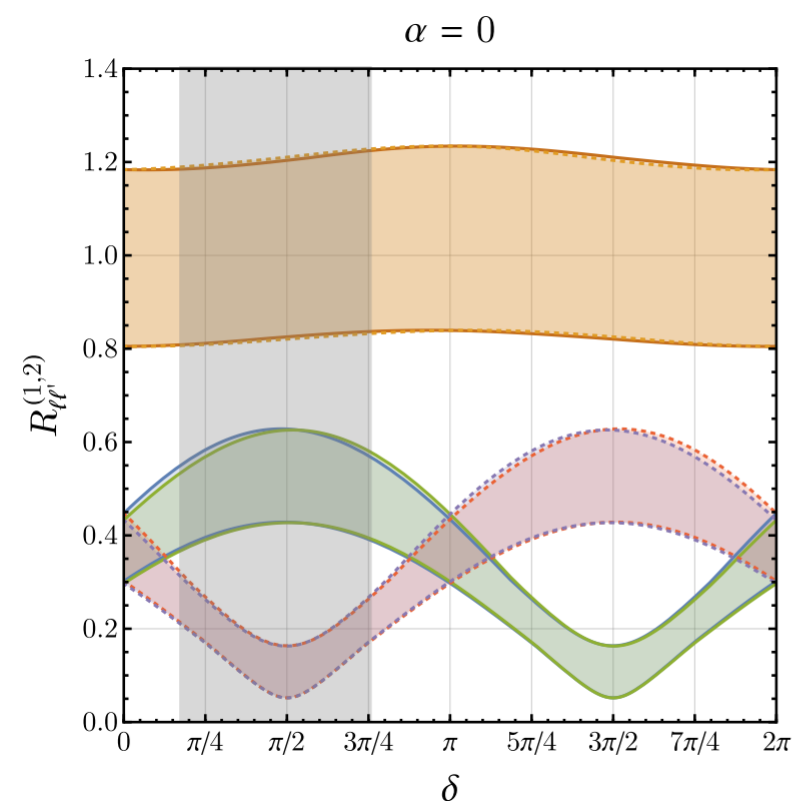
Using  $3\sigma$  ranges for mass-squared differences and mixing angles:



$$\epsilon = 0.2 \rightarrow \delta \sim \pm\pi/4, \pm 3\pi/4$$

# Predictions for CPV phases

Using  $3\sigma$  ranges for mass-squared differences and mixing angles:



$$\epsilon = 0.2 \rightarrow \delta \sim \pm\pi/4, \pm 3\pi/4$$

# Bounds on the couplings $g_1, g_2$

Mass matrix entries:  $|(\mathbf{m}_\nu)_{\ell\ell'}| = |\mathbf{U}_{\ell j}^* m_j \mathbf{U}_{\ell' j}^*|$   $g_1 \ g_2$   
 constrains  $|g_{\ell 1} g_{\ell' 2} + g_{\ell' 1} g_{\ell 2}|$

## Non-unitarity of PMNS:

$\eta = -(\mathbf{R}\mathbf{V})(\mathbf{R}\mathbf{V})^\dagger / 2$  constrains  $|g_{\ell 2} g_{\ell' 2}|$

$$|\eta| < \begin{pmatrix} 1.3 \times 10^{-3} & 1.2 \times 10^{-5} & 1.4 \times 10^{-3} \\ 1.2 \times 10^{-5} & 2.2 \times 10^{-4} & 6.0 \times 10^{-4} \\ 1.4 \times 10^{-3} & 6.0 \times 10^{-4} & 2.8 \times 10^{-3} \end{pmatrix} \quad g_2$$

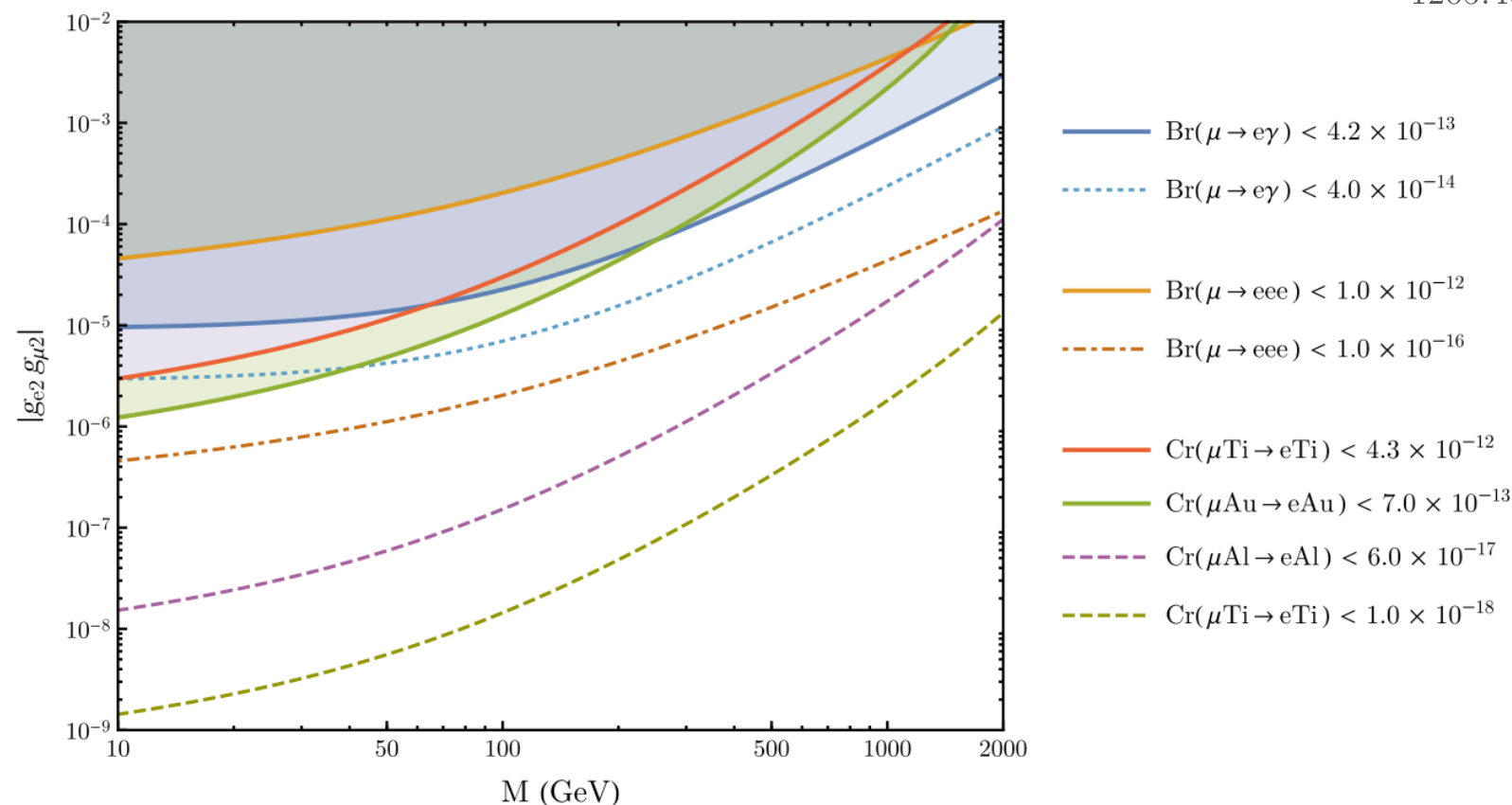
for  $M > v$       Fernandez-Martinez et al, 1605.08774,  
 Blennow et al, 1609.08637

$$\mathbf{R}^* \simeq v \mathbf{Y}_D \mathbf{M}_N^{-1} \quad \mathbf{R}\mathbf{V} \simeq \frac{1}{\sqrt{2}} \frac{v}{M} \begin{pmatrix} g_{e2}^* & -i g_{e2}^* \\ g_{\mu 2}^* & -i g_{\mu 2}^* \\ g_{\tau 2}^* & -i g_{\tau 2}^* \end{pmatrix} \quad \text{controls pheno,}$$

$$\mathbf{M}_N \simeq \mathbf{V}^* \text{diag}(M_1, M_2) \mathbf{V}^\dagger \quad \text{in particular...}$$

# Bounds on the couplings – LFV

see also Ibarra, Molinaro, Petcov, 1103.6217; Dinh, Ibarra, Molinaro, Petcov, 1205.4671



For the  $\epsilon : 1 : 1$  structure, LFV processes like  $\tau \rightarrow e\gamma$ ,  $\tau \rightarrow \mu\gamma$ ,  $\tau \rightarrow \mu\mu\mu$  do not provide stronger constraints on the magnitude of the couplings



## Additional phenomenology

Collider physics: Antusch, Fischer, 1502.05915; Deppisch, Dev, Pilaftsis, 1502.06541; Das, Okada, 1702.04668; Das, Dev, Kim, 1704.00880; and others...

Higgs decays: if new pseudo-Dirac neutrino is light enough, it can affect Higgs **branching ratios**

see also Cely, Ibarra, Molinaro, Petcov, 1208.3654; Antusch, Fischer, 1502.05915

$0\nu\beta\beta$ -decay: contribution from new states **suppressed** with respect to standard 3 neutrino case

(Anti-)leptogenesis: the resonant condition is not satisfied

$$\mu \sim g_1 M \quad \rightarrow \quad \mu \gg \Gamma/2$$

the addition of a third RH neutrino and the suppression of  $\mu$  (e.g. using an extra dim.) may allow to generate BAU

Fukugita, Yanagida, '02

## Summary (1<sup>st</sup> part)

- Froggatt-Nielsen (FN) can be responsible for **approximate  $L$ -conservation**.
- RH neutrinos can have GeV–TeV (low-scale) masses.
- Preferences for CPV phases can be derived, namely  $\delta \sim \pm\pi/4, \pm3\pi/4, \alpha \sim 0$  for  $\epsilon \simeq \lambda_C$ .
- Sizeable LFV signals and collider phenomenology are possible.

# Lepton masses and mixing from modular $S_4$ symmetry

arXiv: 1806.11040

in collaboration with S. T. Petcov



Thu 05/07 (tomorrow)

*9:00*

**F. Feruglio:** “Are neutrino masses modular forms?”

*12:00*

**T. Tatsuishi:** “Neutrino mixing from finite modular groups”

# Modular invariance shaping the leptonic sector

Feruglio, 1706.08749

Principal congruence subgroups, for  $N = 1, 2, \dots$

$$\Gamma(N) \equiv \left\{ \gamma = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \mid a, b, c, d \in \mathbb{Z} \wedge \det \gamma = 1 \wedge \gamma = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \pmod{N} \right\}$$

Modular group:  $\Gamma(1) \simeq \text{SL}(2, \mathbb{Z})$

Groups of linear fractional transformations:

$$\bar{\Gamma}(N) \equiv \Gamma(N)/\{1, -1\}, \text{ for } N = 1, 2$$

$$\bar{\Gamma}(N > 2) \equiv \Gamma(N)$$

$$\bar{\Gamma}(1) : S^2 = (ST)^3 = 1$$

Finite modular groups:

$$\Gamma_N \equiv \bar{\Gamma}(1)/\bar{\Gamma}(N)$$

$$S^2 = (ST)^3 = T^N = 1$$

$$\tau \rightarrow \gamma\tau = \frac{a\tau + b}{c\tau + d},$$

$$\text{with } \gamma = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \bar{\Gamma}(N)$$

$$\{\tau \in \mathbb{C} \mid \text{Im}(\tau) > 0\}$$

$$S : \tau \rightarrow -1/\tau, \quad S = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$$

$$T : \tau \rightarrow \tau + 1, \quad T = \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix}$$

## Modular-invariant SUSY actions

Ferrara et al, '89

$$\mathcal{S} = \int d^4x d^2\theta d^2\bar{\theta} K(\chi_i, \bar{\chi}_i; \tau, \bar{\tau}) + \int d^4x d^2\theta W(\chi_i; \tau) + \text{h.c.}$$

$$W(\chi_i; \tau) = \sum_n \sum_{\{i_1, \dots, i_n\}} (Y_{\{i_1, \dots, i_n\}}(\tau) \chi_{i_1} \cdots \chi_{i_n})_{\mathbf{1}}$$

$$\text{with } \gamma = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \Gamma_N \quad \begin{cases} \tau \rightarrow \gamma\tau = \frac{a\tau + b}{c\tau + d} \\ \chi_i \rightarrow (c\tau + d)^{-k_i} \rho_i(\gamma) \chi_i \\ Y(\tau) \rightarrow Y(\gamma\tau) = (c\tau + d)^{2k_Y} \rho_Y(\gamma) Y(\tau) \end{cases}$$

**Bottom-up approach**

For top-down, see e.g. Kobayashi et al, 1804.06644

# Modular-invariant SUSY actions

$$\text{with } \gamma = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \Gamma_N \quad \begin{cases} \tau \rightarrow \gamma\tau = \frac{a\tau + b}{c\tau + d} \\ \chi_i \rightarrow (c\tau + d)^{-k_i} \rho_i(\gamma) \chi_i \\ Y(\tau) \rightarrow Y(\gamma\tau) = (c\tau + d)^{2k_Y} \rho_Y(\gamma) Y(\tau) \end{cases}$$

Exponents are “weights”

$\tau$  is a dimensionless spurion,  $\langle \tau \rangle$  only source of modular sym. breaking

$$Y(\tau) \text{ are } \mathbf{modular\ forms} \text{ obeying } \begin{cases} 2k_Y - k_{i_1} - \dots - k_{i_n} = 0 \\ \rho_Y \otimes \rho_{i_1} \otimes \dots \otimes \rho_{i_n} \supset \mathbf{1} \end{cases}$$

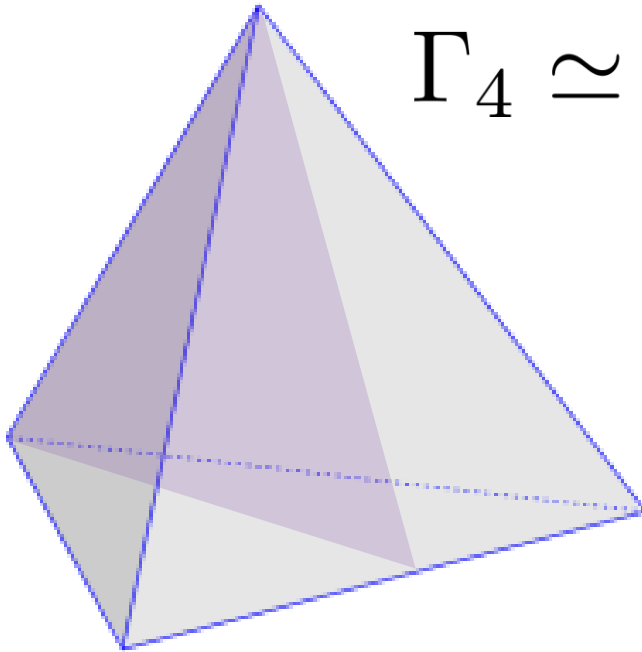
Leading-order predictions potentially modified by  
corrections from Kähler, SUSY breaking

# Modular $S_4$

$S_3$  and  $A_4$  explored in Feruglio, 1706.08749  
and more recently in Kobayashi, Tanaka and  
Tatsuishi, 1803.10391

$$\Gamma_4 \simeq S_4$$

$$\Gamma_2 \simeq S_3, \Gamma_3 \simeq A_4$$



rotations + reflection  
 $\Leftrightarrow$  permutation of vertices

see e.g. Ishhimori et al, 1003.3552

Usual presentation:

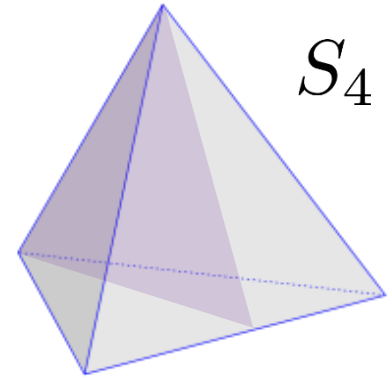
$$\begin{aligned} S^2 &= T^3 = U^2 = (ST)^3 \\ &= (SU)^2 = (TU)^2 = (STU)^4 = 1 \end{aligned}$$

**not convenient**



Modular  $S_4$ 

$$S^2 = (ST)^3 = T^4 = 1$$



$$\mathbf{1} : \rho(S) = 1, \quad \rho(T) = 1$$

$$\mathbf{1}' : \rho(S) = -1, \quad \rho(T) = -1$$

$$\mathbf{2} : \rho(S) = \begin{pmatrix} 0 & \omega \\ \omega^2 & 0 \end{pmatrix}, \quad \rho(T) = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$\mathbf{3} : \rho(S) = \frac{1}{3} \begin{pmatrix} -1 & 2\omega^2 & 2\omega \\ 2\omega & 2 & -\omega^2 \\ 2\omega^2 & -\omega & 2 \end{pmatrix}, \quad \rho(T) = \frac{1}{3} \begin{pmatrix} -1 & 2\omega & 2\omega^2 \\ 2\omega & 2\omega^2 & -1 \\ 2\omega^2 & -1 & 2\omega \end{pmatrix}$$

$$\mathbf{3}' : \rho(S) = -\frac{1}{3} \begin{pmatrix} -1 & 2\omega^2 & 2\omega \\ 2\omega & 2 & -\omega^2 \\ 2\omega^2 & -\omega & 2 \end{pmatrix}, \quad \rho(T) = -\frac{1}{3} \begin{pmatrix} -1 & 2\omega & 2\omega^2 \\ 2\omega & 2\omega^2 & -1 \\ 2\omega^2 & -1 & 2\omega \end{pmatrix}$$

we have adapted group theoretical results from  
Bazzocchi, Merlo, Morisi, 0901.2086

$$\omega = e^{2\pi i/3}$$

## Generators of modular forms – Dedekind eta

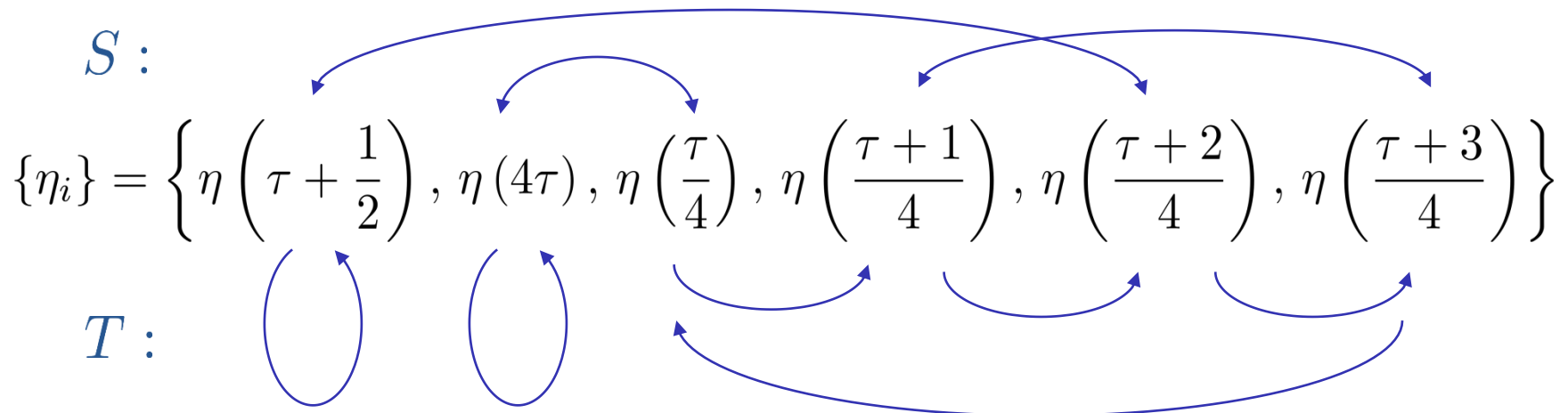
$$\eta(z) \equiv q^{1/24} \prod_{k=1}^{\infty} (1 - q^k)$$

$$q \equiv e^{2\pi i z}$$

$$\eta(z + 1) = e^{i\pi/12} \eta(z)$$

$$\eta(-1/z) = \sqrt{-iz} \eta(z)$$

At level 4, corresponding to S4



up to multiplicative factors

## Generators of modular forms – basis

$$Y(a_1, \dots, a_6 | \tau) \equiv a_1 \frac{\eta'(\tau + 1/2)}{\eta(\tau + 1/2)} + 4a_2 \frac{\eta'(4\tau)}{\eta(4\tau)} + \frac{1}{4} \left[ a_3 \frac{\eta'(\tau/4)}{\eta(\tau/4)} + a_4 \frac{\eta'((\tau + 1)/4)}{\eta((\tau + 1)/4)} + a_5 \frac{\eta'((\tau + 2)/4)}{\eta((\tau + 2)/4)} + a_6 \frac{\eta'((\tau + 3)/4)}{\eta((\tau + 3)/4)} \right]$$

$$\sum_i a_i = 0$$

$$S : Y(a_1, \dots, a_6 | \tau) \rightarrow Y(a_1, a_2, a_3, a_4, a_5, a_6 | -1/\tau) = \tau^2 Y(a_5, a_3, a_2, a_6, a_1, a_4 | \tau)$$

$$T : Y(a_1, \dots, a_6 | \tau) \rightarrow Y(a_1, a_2, a_3, a_4, a_5, a_6 | \tau + 1) = Y(a_1, a_2, a_6, a_3, a_4, a_5 | \tau)$$

Lowest weight forms arrange into:

$$Y_{\mathbf{2}}(\tau) \equiv \begin{pmatrix} Y_1(\tau) \\ Y_2(\tau) \end{pmatrix} \quad \text{doublet}$$

$$Y_{\mathbf{3}' }(\tau) \equiv \begin{pmatrix} Y_3(\tau) \\ Y_4(\tau) \\ Y_5(\tau) \end{pmatrix} \quad \text{triplet'}$$

with

$$Y_1(\tau) \equiv Y(1, 1, \omega, \omega^2, \omega, \omega^2 | \tau)$$

$$Y_2(\tau) \equiv Y(1, 1, \omega^2, \omega, \omega^2, \omega | \tau)$$

$$Y_3(\tau) \equiv Y(1, -1, -1, -1, 1, 1 | \tau)$$

$$Y_4(\tau) \equiv Y(1, -1, -\omega^2, -\omega, \omega^2, \omega | \tau)$$

$$Y_5(\tau) \equiv Y(1, -1, -\omega, -\omega^2, \omega, \omega^2 | \tau)$$

## Constraints



$$\begin{aligned} \frac{1}{3} (Y_3^2 + 2Y_4Y_5) &= Y_1Y_2, & -\frac{1}{\sqrt{3}} (Y_3^2 - Y_4Y_5) &= Y_1Y_4 - Y_2Y_5, \\ \frac{1}{3} (Y_4^2 + 2Y_3Y_5) &= Y_2^2, & -\frac{1}{\sqrt{3}} (Y_5^2 - Y_3Y_4) &= Y_1Y_5 - Y_2Y_3, \\ \frac{1}{3} (Y_5^2 + 2Y_3Y_4) &= Y_1^2, & -\frac{1}{\sqrt{3}} (Y_4^2 - Y_3Y_5) &= Y_1Y_3 - Y_2Y_4. \end{aligned}$$

Related to non-linear realisation of the symmetry as suggested in 1706.08749?

## Modular forms of higher weight

From tensor products of  $Y_2(\tau)$  and  $Y_{3'}(\tau)$ :



$$Y_1^{(4)} = Y_1Y_2 \sim \mathbf{1}$$

$$Y_2^{(4)} = (Y_2^2, Y_1^2)^T \sim \mathbf{2}$$

$$Y_3^{(4)} = (Y_1Y_4 - Y_2Y_5, Y_1Y_5 - Y_2Y_3, Y_1Y_3 - Y_2Y_4)^T \sim \mathbf{3}$$

$$Y_{3'}^{(4)} = (Y_1Y_4 + Y_2Y_5, Y_1Y_5 + Y_2Y_3, Y_1Y_3 + Y_2Y_4)^T \sim \mathbf{3'}$$

Non-zero  $\mathbf{1}'$   
arises only at  
weight 6

# Building minimal models



$$W = \sum_i \alpha_i (E_i^c L H_d Y_2^{a_i} Y_{3'}^{b_i})_1 + \frac{g}{\Lambda} (L H_u L H_u Y_2^c Y_{3'}^d)_1$$

# Building minimal models

## Guidelines



- **No flavons** are introduced,
- Neutrino masses arise from the **Weinberg operator**,
- Higgs multiplets transform trivially,
- Lepton doublets transform as an S4 triplet,
- Lepton singlets transform as S4 singlets, and
- Lowest possible weights are chosen such that a **rank 3 charged-lepton mass matrix** is possible without additional shaping symmetries.

$$W = \sum_i \alpha_i (E_i^c L H_d Y_{\mathbf{2}}^{a_i} Y_{\mathbf{3}'}^{b_i})_{\mathbf{1}} + \frac{g}{\Lambda} (L H_u L H_u Y_{\mathbf{2}}^c Y_{\mathbf{3}'}^d)_{\mathbf{1}}$$

# Model I ( $k_L = 1$ )

Minimality guidelines imply:

	$H_u$	$H_d$	$L$	$E_1^c$	$E_2^c$	$E_3^c$
$\rho_i$	<b>1</b>	<b>1</b>	<b>3</b>	<b>1'</b>	<b>1</b>	<b>1'</b>
			<b>3'</b>	<b>1</b>	<b>1'</b>	<b>1</b>
$k_i$	0	0	1	1	3	3

$$\alpha = 1, \beta = 1.7 \times 10^{-4},$$

$$\gamma = 0.025, \tau = 0.19 + 0.99i$$

## IO spectrum

$$\frac{m_e}{m_\mu} \simeq 0.0049, \sin^2 \theta_{12} \simeq 0.146, \quad \delta \simeq 1.23\pi,$$

$$\frac{m_\mu}{m_\tau} \simeq 0.0533, \sin^2 \theta_{13} \simeq 0.116, \quad \alpha_{21} \simeq 1.92\pi,$$

$$r \simeq 0.0287, \sin^2 \theta_{23} \simeq 0.548, \quad \alpha_{31} \simeq 0.50\pi.$$

$$M_\nu^I = \frac{2g_1 v_u^2}{\Lambda} \begin{pmatrix} 0 & Y_2 & Y_1 \\ Y_2 & Y_1 & 0 \\ Y_1 & 0 & Y_2 \end{pmatrix}$$

$$M_e^I = v_d \begin{pmatrix} \alpha Y_3 & \alpha Y_5 & \alpha Y_4 \\ \beta (Y_1 Y_4 - Y_2 Y_5) & \beta (Y_1 Y_3 - Y_2 Y_4) & \beta (Y_1 Y_5 - Y_2 Y_3) \\ \gamma (Y_1 Y_4 + Y_2 Y_5) & \gamma (Y_1 Y_3 + Y_2 Y_4) & \gamma (Y_1 Y_5 + Y_2 Y_3) \end{pmatrix}^\dagger$$

# Model II ( $k_L = 2$ )

$$\alpha = 0.12, \beta = 4.2 \times 10^{-4},$$

$$\gamma = 1, g/g' = 0.34 - 0.94i,$$

$$g''/g' = 1.93 - 0.53i, \tau = -0.1 + 1.2i$$

Minimality guidelines imply:

	$H_u$	$H_d$	$L$	$E_1^c$	$E_2^c$	$E_3^c$
$\rho_i$	<b>1</b>	<b>1</b>	<b>3</b>	<b>1'</b>	<b>1</b>	<b>1'</b>
			<b>3'</b>	<b>1</b>	<b>1'</b>	<b>1</b>
$k_i$	0	0	2	0	2	2

NO spectrum

$$\frac{m_e}{m_\mu} \simeq 0.0052, \sin^2 \theta_{12} \simeq 0.294, \quad \delta \simeq 0.32\pi,$$

$$\frac{m_\mu}{m_\tau} \simeq 0.0542, \sin^2 \theta_{13} \simeq 0.021, \alpha_{21} \simeq 0.96\pi,$$

$$r \simeq 0.0300, \sin^2 \theta_{23} \simeq 0.574, \alpha_{31} \simeq 1.59\pi.$$

$$M_e^{\text{II}} = M_e^{\text{I}}$$

$$M_\nu^{\text{II}} = \frac{2g'v_u^2}{\Lambda} \left[ \begin{array}{ccc} \left( \begin{array}{ccc} (g/g')Y_1Y_2 & Y_2^2 & Y_1^2 \\ Y_2^2 & Y_1^2 & (g/g')Y_1Y_2 \\ Y_1^2 & (g/g')Y_1Y_2 & Y_2^2 \end{array} \right) \\ + \frac{1}{2} \frac{g''}{g'} \left( \begin{array}{ccc} 2(Y_1Y_4 - Y_2Y_5) & -(Y_1Y_3 - Y_2Y_4) & -(Y_1Y_5 - Y_2Y_3) \\ -(Y_1Y_3 - Y_2Y_4) & 2(Y_1Y_5 - Y_2Y_3) & -(Y_1Y_4 - Y_2Y_5) \\ -(Y_1Y_5 - Y_2Y_3) & -(Y_1Y_4 - Y_2Y_5) & 2(Y_1Y_3 - Y_2Y_4) \end{array} \right) \end{array} \right]$$



# Model II ( $k_L = 2$ )

$$\alpha = 0.12, \beta = 1,$$

$$\gamma = 4.5 \times 10^{-4}, g/g' = 15.1 + 7.7i,$$

$$g''/g' = -0.40 - 1.08i, \tau = 0.35 + 0.85i$$

Minimality guidelines imply:

	$H_u$	$H_d$	$L$	$E_1^c$	$E_2^c$	$E_3^c$
$\rho_i$	<b>1</b>	<b>1</b>	<b>3</b>	<b>1'</b>	<b>1</b>	<b>1'</b>
			<b>3'</b>	<b>1</b>	<b>1'</b>	<b>1</b>
$k_i$	0	0	2	0	2	2

NO spectrum ✓

$$\frac{m_e}{m_\mu} \simeq 0.0045, \sin^2 \theta_{12} \simeq 0.278, \quad \delta \simeq 1.37\pi,$$

$$\frac{m_\mu}{m_\tau} \simeq 0.0557, \sin^2 \theta_{13} \simeq 0.021, \quad \alpha_{21} \simeq 0.25\pi,$$

$$r \simeq 0.0296, \sin^2 \theta_{23} \simeq 0.480, \quad \alpha_{31} \simeq 1.22\pi.$$

$$|\langle m \rangle| \simeq 0.043 \text{ eV}$$

$$M_e^{\text{II}} = M_e^{\text{I}}$$

$$M_\nu^{\text{II}} = \frac{2g'v_u^2}{\Lambda} \left[ \begin{array}{ccc} \left( \begin{array}{ccc} (g/g')Y_1Y_2 & Y_2^2 & Y_1^2 \\ Y_2^2 & Y_1^2 & (g/g')Y_1Y_2 \\ Y_1^2 & (g/g')Y_1Y_2 & Y_2^2 \end{array} \right) \\ + \frac{1}{2} \frac{g''}{g'} \left( \begin{array}{ccc} 2(Y_1Y_4 - Y_2Y_5) & -(Y_1Y_3 - Y_2Y_4) & -(Y_1Y_5 - Y_2Y_3) \\ -(Y_1Y_3 - Y_2Y_4) & 2(Y_1Y_5 - Y_2Y_3) & -(Y_1Y_4 - Y_2Y_5) \\ -(Y_1Y_5 - Y_2Y_3) & -(Y_1Y_4 - Y_2Y_5) & 2(Y_1Y_3 - Y_2Y_4) \end{array} \right) \end{array} \right]$$

## Summary (2<sup>nd</sup> part)

- We have explored the consequences of broken **modular invariance** in the lepton sector. We considered the finite modular subgroup  $\Gamma_4 \simeq S_4$ .
- A minimal model with **no flavons** can accommodate lepton masses and mixing while predicting Dirac and Majorana CPV phases.
- Model-building avenue worthy of **future study**.



Thank you / Merci vilmal !