

Quark and lepton mixing from flavor and CP symmetries

Gui-Jun Ding

University of Science and Technology of China



FLASY18, July 2nd-July 5th, 2018, University of Basel, Switzerland

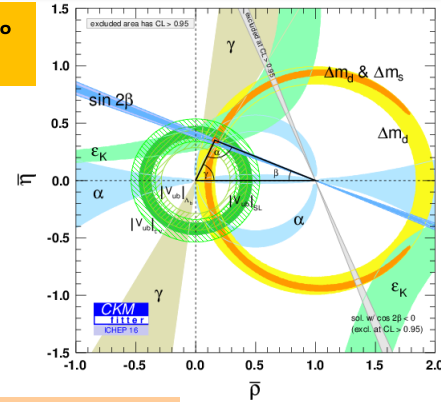
“Old” flavor mixing puzzle in SM

[Particle Data Group 2018]

$$\alpha = (88.8 \pm 2.3)^\circ$$

Quarks:

$$\|V_{CKM}\| \approx \begin{pmatrix} 0.97434 & 0.22506 & 0.00357 \\ 0.22492 & 0.97351 & 0.0414 \\ 0.00875 & 0.0403 & 0.99915 \end{pmatrix}$$



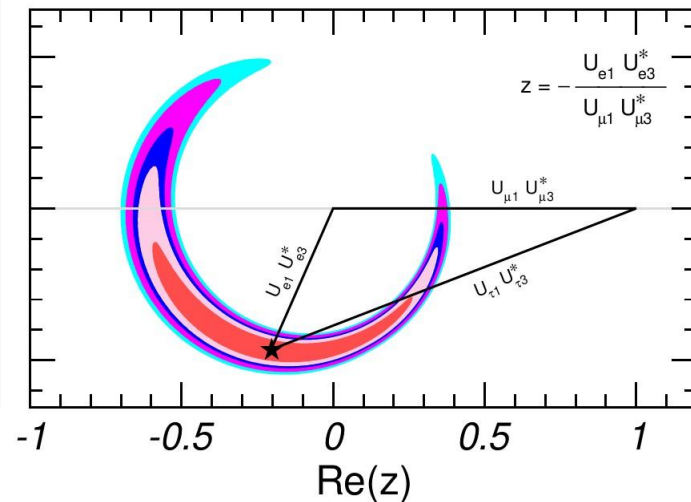
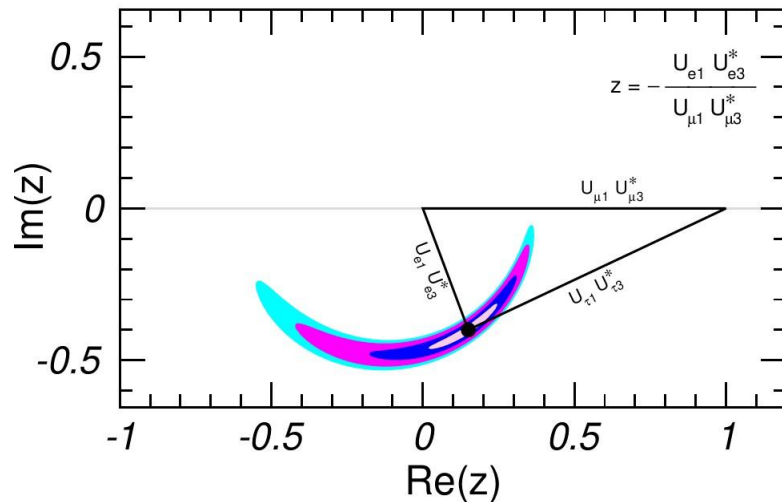
Leptons:

$$\|U_{PMNS}\| = \begin{pmatrix} 0.799 \sim 0.844 & 0.516 \sim 0.582 & 0.141 \sim 0.156 \\ 0.242 \sim 0.494 & 0.467 \sim 0.678 & 0.639 \sim 0.774 \\ 0.284 \sim 0.521 & 0.490 \sim 0.695 & 0.615 \sim 0.754 \end{pmatrix}$$

[Gonzalez-Garcia et al., NuFIT3.2 (2018)]

NuFIT 3.2 (2018)

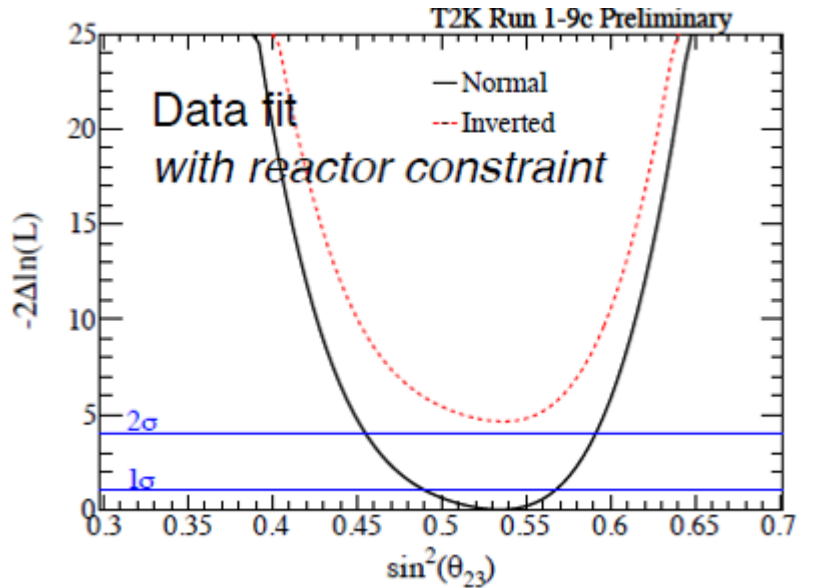
IO | NO



Latest results on θ_{23} octant

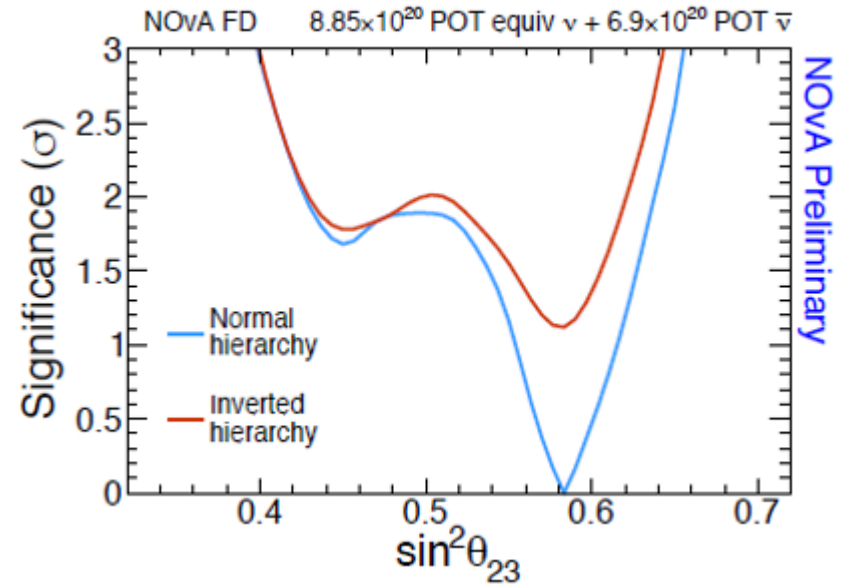
T2K

Wascko@ Neutrino 18



NOvA

Sanchez @ Neutrino 18



- Best fit values:

	NO	IO
$\sin^2\theta_{23}$	$0.536^{+0.031}_{-0.046}$	$0.536^{+0.031}_{-0.041}$
$ \Delta m^2 $	2.434 ± 0.064	$2.410^{+0.062}_{-0.063}$

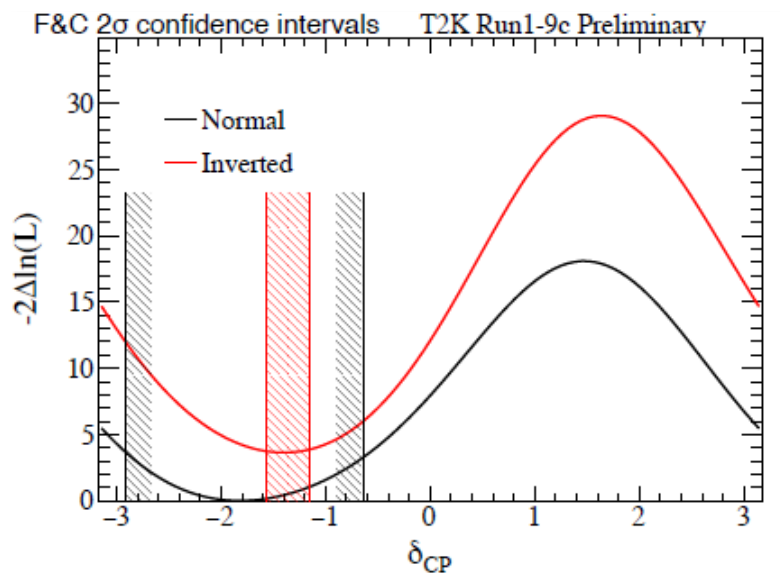
- Best fit: $\sin^2\theta_{23} = 0.58 \pm 0.03$ for NO
- Prefer non-maximal at 1.8σ , exclude lower octant at similar level

more disfavored lower octant?

Latest results on δ_{CP}

T2K

Wascko@ Neutrino 18



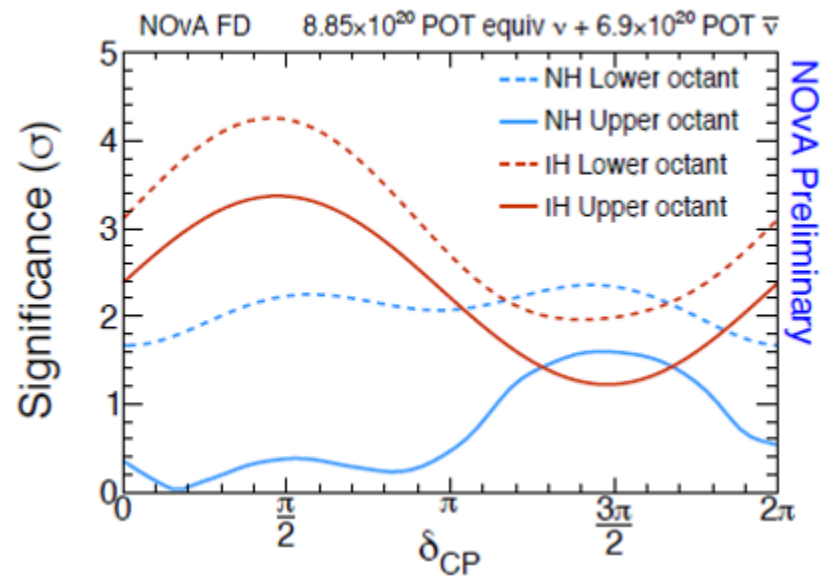
- CP conserving values $\delta_{CP} = 0, \pi$ outside 2σ region for NO & IO

- $-\pi < \delta_{CP} < 0$ is favored



NOvA

Sanchez @ Neutrino 18



- Best fit: $\delta_{CP} \approx 0.17\pi$ for NO, $\delta_{CP} \approx 1.5\pi$ for IO
- Prefer NO by 1.8σ , exclude $\delta_{CP} = \pi/2$ in the IO at $> 3\sigma$

a simple predictive CP symmetry: $\mu\tau$ reflection

$\mu\tau$ reflection = $\mu\tau$ exchange + canonical CP

See talk by Celso Nishi

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} \xrightarrow{\text{CP}} \begin{pmatrix} \nu_e^c \\ \nu_\tau^c \\ \nu_\mu^c \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} \begin{pmatrix} \nu_e^c \\ \nu_\mu^c \\ \nu_\tau^c \end{pmatrix}$$

This CP transformation is **not** a unit matrix.

If the neutrino mass matrix is **invariant** under the $\mu\tau$ reflection

$$m_\nu = \begin{pmatrix} a & b & b^* \\ b & c & d \\ b^* & d & c^* \end{pmatrix}$$

$\nu_e \quad \nu_\mu \leftrightarrow \nu_\tau^c$



$$\theta_{23} = \frac{\pi}{4}, \quad \delta_{CP} = \pm \frac{\pi}{2}$$

[Harrison, Scott, 2002; Grimus, Lavoura, 2004]

A general residual CP transformation

Definition of generalized CP transformation [Chen, Li, Ding, 2014]

$$\psi \xrightarrow{CP} iX\gamma^0\mathcal{C}\bar{\psi}^T, \quad X \text{ is unitary and symmetric}$$

Invariance of the ψ mass matrix under X

$$\begin{cases} X^T m_\psi X = m_\psi^*, & \text{for Majorana } \psi \\ X^\dagger m_\psi^\dagger m_\psi X = (m_\psi^\dagger m_\psi)^*, & \text{for Dirac } \psi \end{cases}$$

Performing Takagi factorization of X

$$X = \Sigma \Sigma^T \longrightarrow \begin{cases} \Sigma^T m_\psi \Sigma = (\Sigma^T m_\psi \Sigma)^*, & \text{for Majorana } \psi \\ \Sigma^\dagger m_\psi^\dagger m_\psi \Sigma = (\Sigma^\dagger m_\psi^\dagger m_\psi \Sigma)^*, & \text{for Dirac } \psi \end{cases}$$

The mass matrix m_ψ is diagonalized by

$$U_\psi = \Sigma O_{3 \times 3} P, \quad P \equiv \text{diag}(e^{i\alpha_1}, e^{i\alpha_2}, e^{i\alpha_3})$$

with $\alpha_{1,2,3} = 0, \pm\pi/2, \pi$ for Majorana ψ

Determined up to a real orthogonal matrix $O_{3 \times 3}$

Generalized $\mu\tau$ reflection on neutrinos

the generalized $\mu\tau$ reflection of neutrino

[Chen, Ding, Gonzalez-Canales, Valle, 2015]

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} \rightarrow \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos \Theta & i \sin \Theta \\ 0 & i \sin \Theta & \cos \Theta \end{pmatrix} \begin{pmatrix} \nu_e^c \\ \nu_\mu^c \\ \nu_\tau^c \end{pmatrix},$$

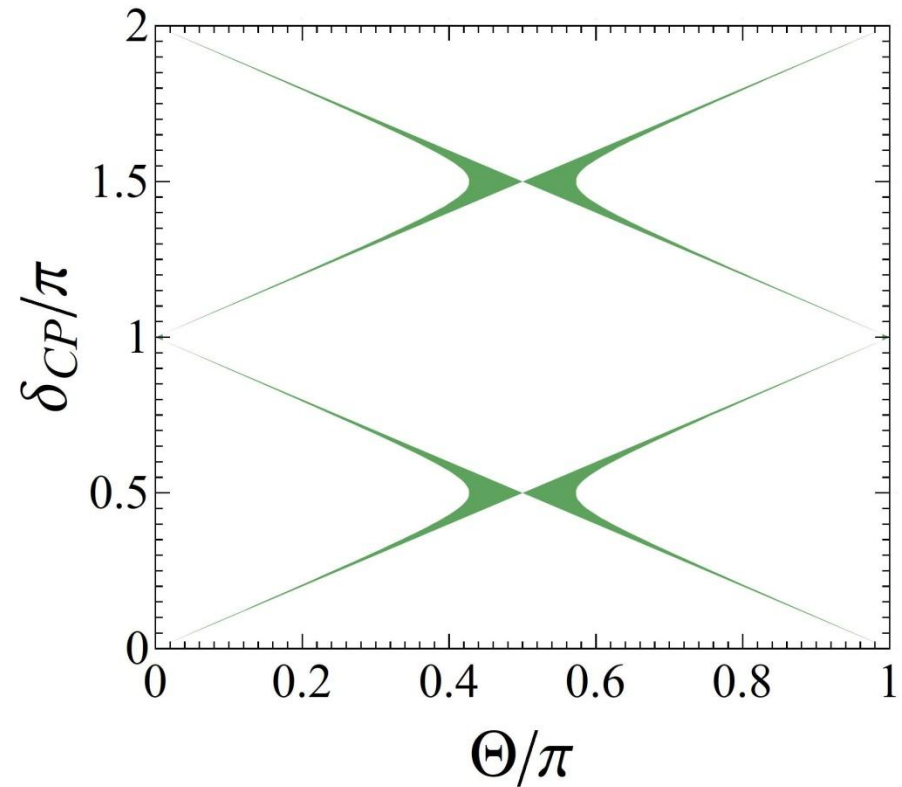
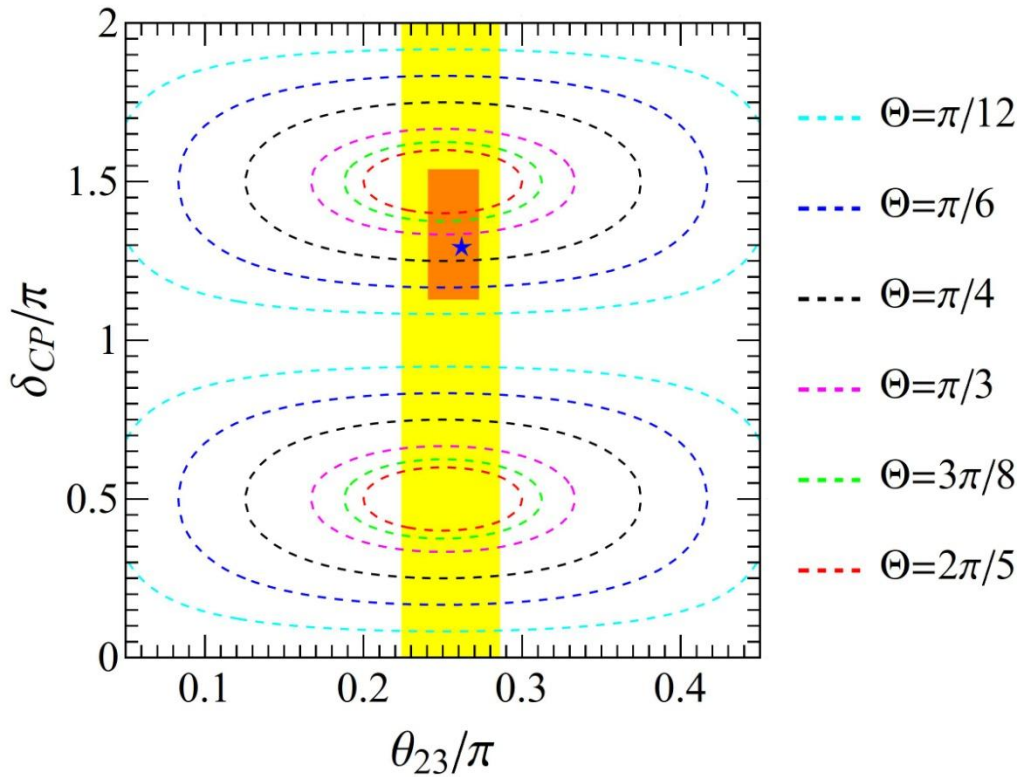
$\mu\tau$ reflection: $\Theta = \pi/2$

In the charged lepton diagonal basis, the lepton mixing matrix is

$$U = \Sigma_{\mu\tau} O_{3\times 3} \text{diag}(1, i^{k_1}, i^{k_2}), \quad \Sigma_{\mu\tau} \equiv \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos \frac{\Theta}{2} & i \sin \frac{\Theta}{2} \\ 0 & i \sin \frac{\Theta}{2} & \cos \frac{\Theta}{2} \end{pmatrix}$$

Predictions for lepton mixing parameters:

$$\sin^2 \delta_{CP} \sin^2 2\theta_{23} = \sin^2 \Theta, \quad \sin \alpha_{21} = \sin \alpha'_{31} = 0$$

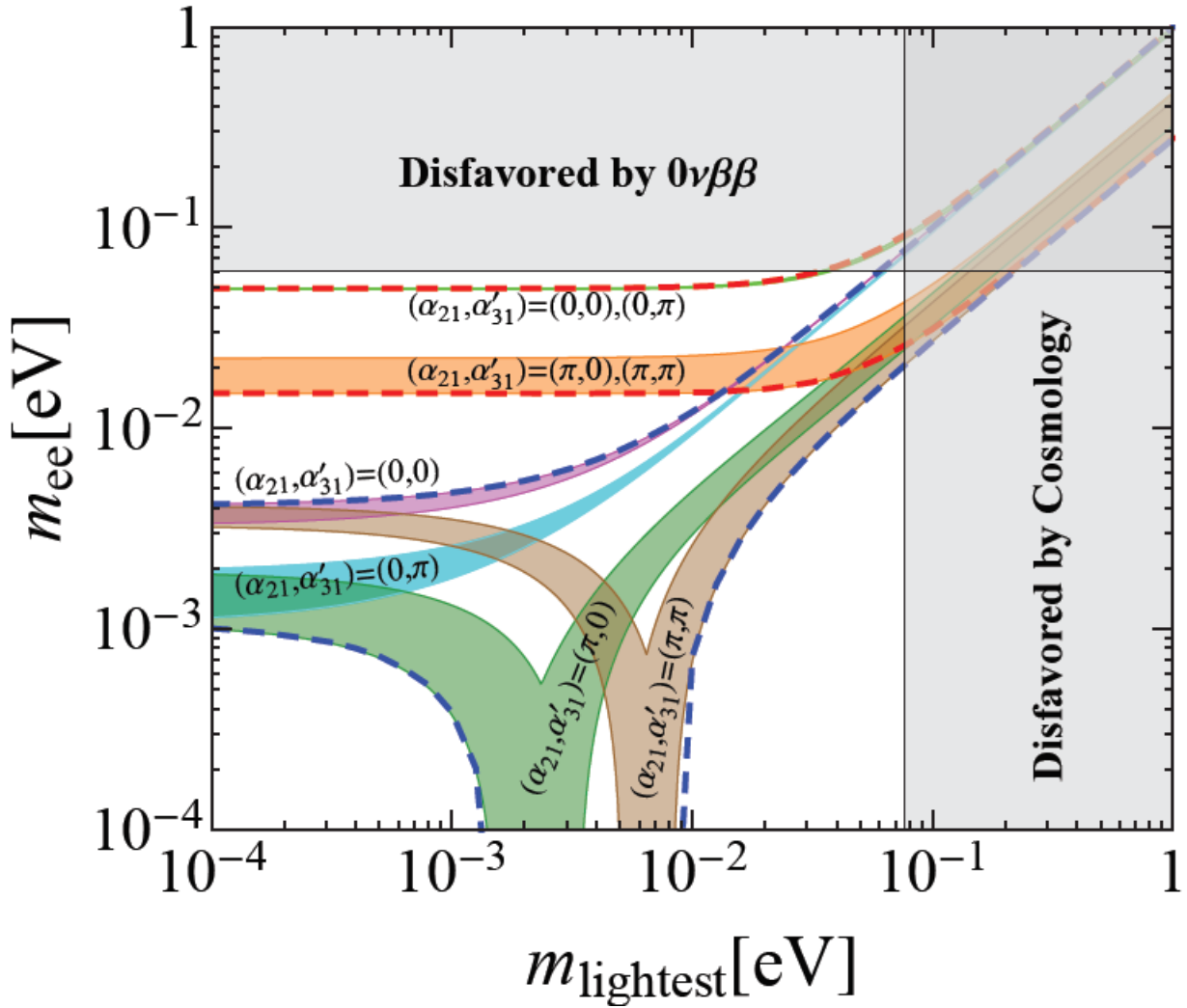


Numerical example:

$$\Theta = \frac{\pi}{3}, \quad \theta_1 = 0.274\pi, \quad \theta_2 = 0.0475\pi, \quad \theta_3 = 0.813\pi,$$

$$\sin^2 \theta_{12} = 0.307, \quad \sin^2 \theta_{13} = 0.0221, \quad \sin^2 \theta_{23} = 0.538, \quad \delta_{CP} = 1.335\pi$$

Predictions for the effective mass of $0\nu 2\beta$ decay



Generalized $e\mu$ reflection on charged leptons

the generalized $e\mu$ reflection of charged lepton

$$\begin{pmatrix} e \\ \mu \\ \tau \end{pmatrix} \rightarrow \begin{pmatrix} \cos \Theta & i \sin \Theta & 0 \\ i \sin \Theta & \cos \Theta & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} e^c \\ \mu^c \\ \tau^c \end{pmatrix}$$

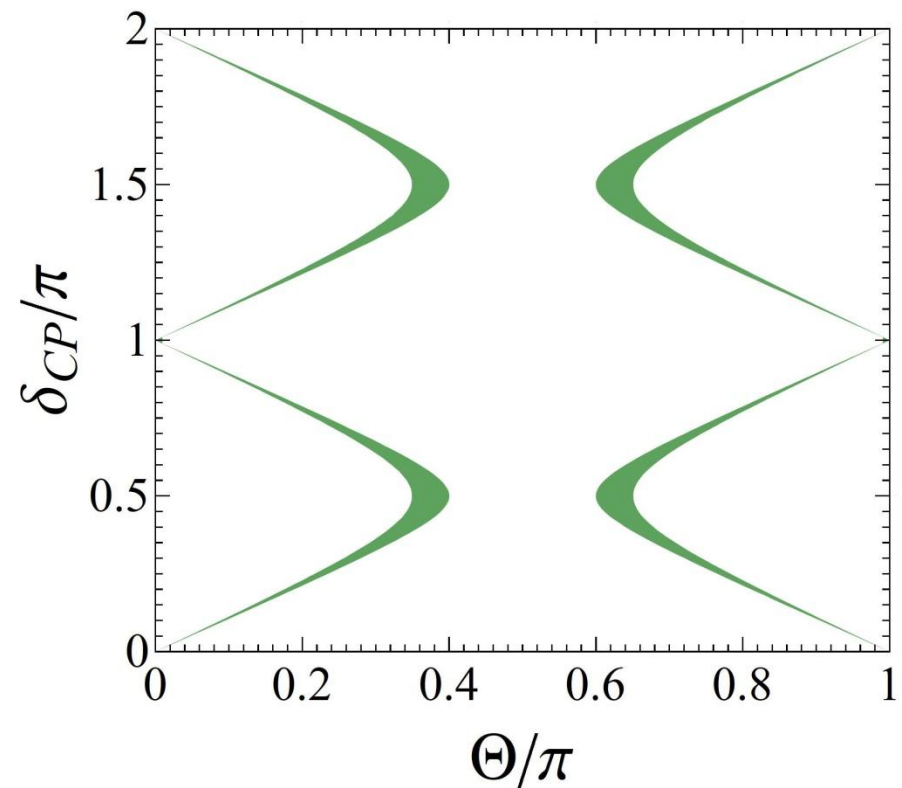
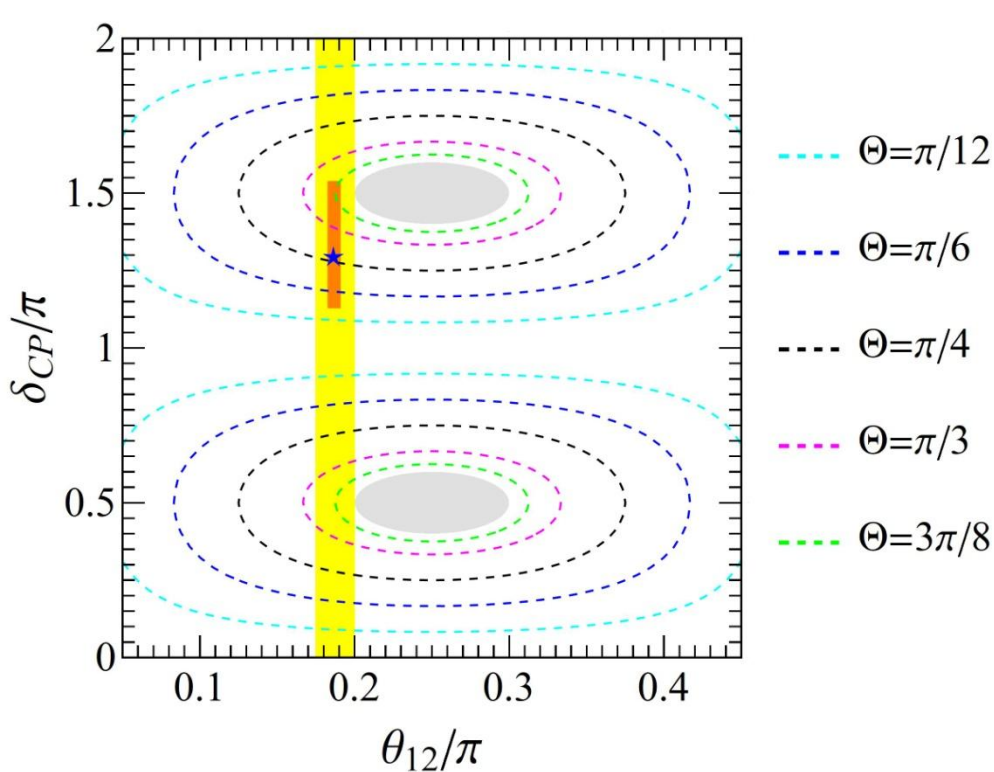
[Chen,Chulia,Ding,Srivastava, Valle,2018]

In the diagonal neutrino basis, the lepton mixing matrix reads

$$U = \text{diag}(e^{i\alpha_1}, e^{i\alpha_2}, e^{i\alpha_3}) O_{3 \times 3} \Sigma_{e\mu}^\dagger, \quad \Sigma_{e\mu} = \begin{pmatrix} \cos \frac{\Theta}{2} & i \sin \frac{\Theta}{2} & 0 \\ i \sin \frac{\Theta}{2} & \cos \frac{\Theta}{2} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

Sum rule for lepton mixing parameters:

$$\sin^2 \delta_{CP} \sin^2 2\theta_{12} = \sin^2 \Theta$$



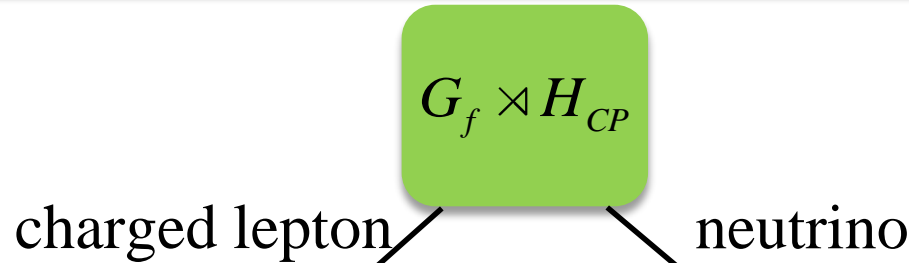
Numerical example:

$$\Theta = \frac{\pi}{3}, \quad \theta_1 = 0.262\pi, \quad \theta_2 = 0.0475\pi, \quad \theta_3 = 0.890\pi,$$

$$\sin^2 \theta_{12} = 0.307, \quad \sin^2 \theta_{13} = 0.0221, \quad \sin^2 \theta_{23} = 0.538, \quad \delta_{CP} = 1.388\pi$$

Universal flavor symmetry for quark and lepton mixing

[Lu, Ding, 2016,2018;Li, Lu, Ding, 2017]



$$X_l = \Sigma_l \Sigma_l^T,$$

$$\Sigma_l^\dagger \rho_3(g_l) \Sigma_l = \pm \text{diag}(1, -1, -1)$$

$$Z_2^{g_l} \times X_l$$

$$Z_2^{g_\nu} \times X_\nu$$

$$X_\nu = \Sigma_\nu \Sigma_\nu^T,$$

$$\Sigma_\nu^\dagger \rho_3(g_\nu) \Sigma_\nu = \pm \text{diag}(1, -1, -1)$$

$$U_l = \Sigma_l R_{23}(\theta_l) P_l$$

$$U_\nu = \Sigma_\nu R_{23}(\theta_\nu) P_\nu$$

$$U_{PMNS}(\theta_l, \theta_\nu) = R_{23}^T(\theta_l) \Sigma_l^\dagger \Sigma_\nu R_{23}(\theta_\nu) P_\nu$$

- All mixing angles and CP phases are expressed in terms of **two free parameters** $\theta_{l,\nu} \in [0, \pi)$
- This scheme can be extended to quark sector, and the CKM mixing matrix is of similar form

CKM matrix from $\Delta(6n^2)$ and CP

➤ $\Delta(6n^2)$ is isomorphic to $(Z_n \times Z_n) \rtimes S_3$ with the multiplication rules

$$a^3 = b^2 = (ab)^2 = 1,$$

$$S_3 = \langle a, b \rangle$$

$$c^n = d^n = 1, \quad cd = dc,$$

$$Z_n \times Z_n = \langle c, d \rangle$$

$$aca^{-1} = c^{-1}d^{-1}, \quad ada^{-1} = c, \quad bcb^{-1} = d^{-1}, \quad bdb^{-1} = c^{-1} \quad \boxtimes$$

for triplet representation $\eta \equiv e^{2\pi i/n}$ [Escobar,Luhn,2008]

$$a = \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & 0 & 0 \end{pmatrix}, \quad b = -\begin{pmatrix} 0 & 0 & 1 \\ 0 & 1 & 0 \\ 1 & 0 & 0 \end{pmatrix}, \quad c = \begin{pmatrix} \eta & 0 & 0 \\ 0 & \eta^{-1} & 0 \\ 0 & 0 & 1 \end{pmatrix}, \quad d = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \eta & 0 \\ 0 & 0 & \eta^{-1} \end{pmatrix}$$

Only one viable quark mixing pattern is obtained up to row and column permutations

Residual symmetry:

$$Z_2^{g_u} = Z_2^{bc^x d^x}, \quad X_u = \left\{ c^\gamma d^{-2x-\gamma}, bc^{x+\gamma} d^{-x-\gamma} \right\},$$

$$Z_2^{g_d} = Z_2^{bc^y d^y}, \quad X_d = \left\{ c^\delta d^{-2y-\delta}, bc^{y+\delta} d^{-y-\delta} \right\}$$

➤ The CKM matrix is determined to be

$$V_{CKM} = \begin{pmatrix} s_d \sin \varphi_1 & \boxed{\cos \varphi_1} & -c_d \sin \varphi_1 \\ c_u c_d e^{i\varphi_2} + s_u s_d \cos \varphi_1 & -s_u \sin \varphi_1 & c_u s_d e^{i\varphi_2} - c_d s_u \cos \varphi_1 \\ c_d s_u e^{i\varphi_2} - c_u s_d \cos \varphi_1 & c_u \sin \varphi_1 & s_u s_d e^{i\varphi_2} + c_u c_d \cos \varphi_1 \end{pmatrix}$$

with

$$\varphi_1 = \frac{x-y}{n} \pi, \quad \varphi_2 = \frac{3(x-y+\gamma-\delta)}{n} \pi \longleftarrow \text{fixed by residual symmetry}$$

➤ Sum rules

$$\begin{cases} \cos^2 \theta_{13}^q \sin^2 \theta_{12}^q = \cos^2 \varphi_1, \\ J_{CP}^q \approx \frac{1}{2} \sin 2\varphi_1 \sin \varphi_2 \sin \theta_{13}^q \sin \theta_{23}^q \end{cases}$$

➤ Example: $\varphi_1 = \varphi_2 = 3\pi/7$ which can be achieved in $\Delta(6 \cdot 7^2) = \Delta(294)$ group

	θ_u^{bf}/π	θ_d^{bf}/π	$\sin \theta_{12}^q$	$\sin \theta_{23}^q$	$\sin \theta_{13}^q$	J_{CP}^q
Our	0.4867	0.4988	0.22252	0.04204	0.00359	3.202×10^{-5}
Data	—	—	0.22497 ± 0.00069	0.04229 ± 0.00057	0.00368 ± 0.00010	$(3.115 \pm 0.093) \times 10^{-5}$

The measured quark mixing angles and CP violation phases can be accommodated.

➤ Lepton sector : $g_l=b$, $X_l=cd^{-1}$, $g_\nu=abc^{-2}$, $X_\nu=cd^{-2}$ in $\Delta(6 \cdot 7^2) = \Delta(294)$

$$U_{PMNS} = R_{13}^T(\theta_l) \frac{1}{2} \begin{pmatrix} -\sqrt{2}e^{i\pi/7} & -1 & 1 \\ -\sqrt{2}e^{i\pi/7} & 1 & -1 \\ 0 & \sqrt{2}e^{-i\pi/7} & \sqrt{2}e^{-i\pi/7} \end{pmatrix} R_{13}(\theta_\nu) Q_\nu$$

Sum rule between $\cos\delta_{CP}$ and mixing angles :

$$\cos\delta_{CP} = \frac{4\cos^2\theta_{12}\cos^2\theta_{23} + 4\sin^2\theta_{12}\sin^2\theta_{23}\sin^2\theta_{13} - 1}{2\sin 2\theta_{12}\sin 2\theta_{23}\sin\theta_{13}}$$

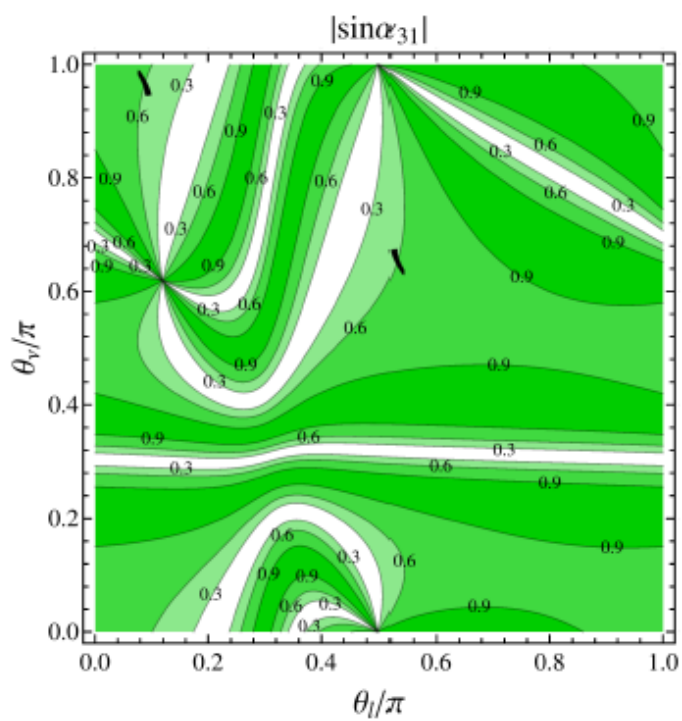
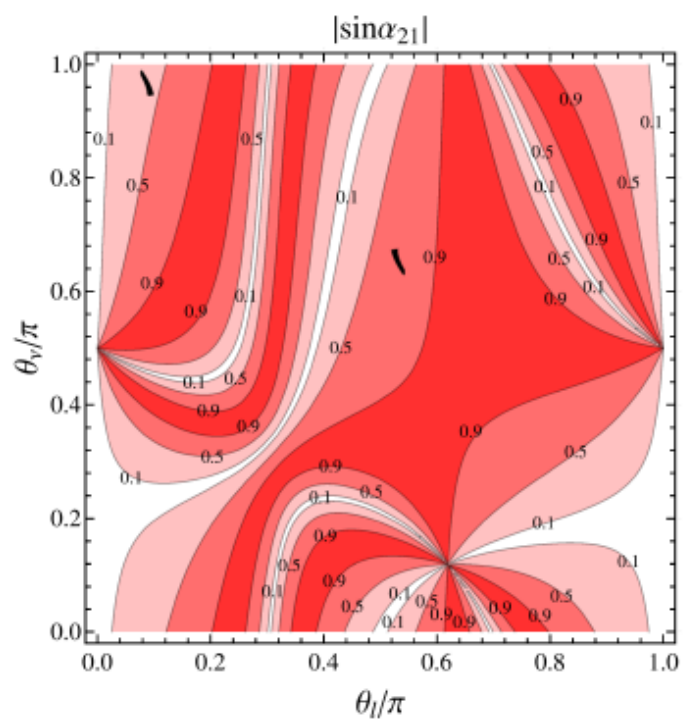
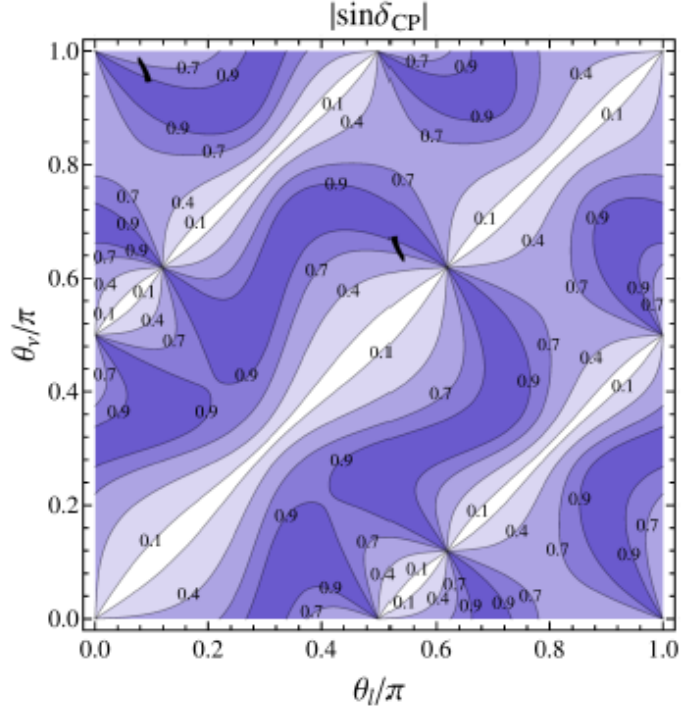
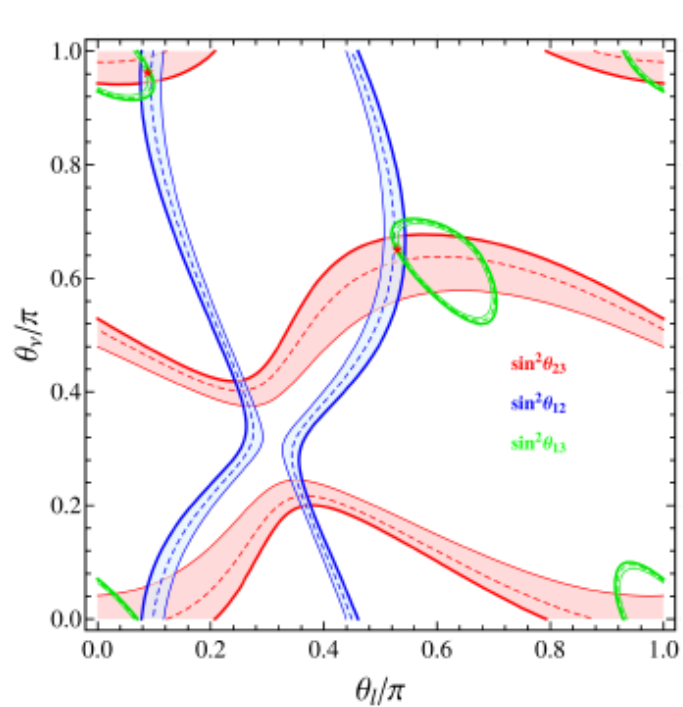
Best fit values:

$$\theta_l = 0.089\pi, \quad \theta_\nu = 0.965\pi, \quad \chi_{\min}^2 = 2.416,$$

$$\sin^2\theta_{13} = 0.0222, \quad \sin^2\theta_{12} = 0.319, \quad \sin^2\theta_{23} = 0.579,$$

$$\delta_{CP} = 1.204\pi, \quad |\sin\alpha_{21}| = 0.391, \quad |\sin\alpha_{31}| = 0.596$$

[Li, Lu, Ding, 2017]



A model building paradigm: Littlest Seesaw

- Three lepton doublets transform as triplet of flavor symmetry $L \sim 3$

$$\mathcal{L} = -y_{\text{atm}} (L \cdot \Phi_{\text{atm}}) N_{\text{atm}}^c - y_{\text{sol}} (L \cdot \Phi_{\text{sol}}) N_{\text{sol}}^c - \frac{1}{2} M_{\text{atm}} N_{\text{atm}}^c N_{\text{atm}}^c - \frac{1}{2} M_{\text{sol}} N_{\text{sol}}^c N_{\text{sol}}^c$$

- Vacuum alignment **CSD(n)** [King, 2015; Luhn, King, 2016]

$$\langle \Phi_{\text{atm}} \rangle = v_{\text{atm}} \begin{pmatrix} 0 \\ 1 \\ 1 \end{pmatrix}, \quad \langle \Phi_{\text{sol}} \rangle = v_{\text{sol}} \begin{pmatrix} 1 \\ n \\ n-2 \end{pmatrix}$$

See talk by Steve King

- Neutrino mass matrix

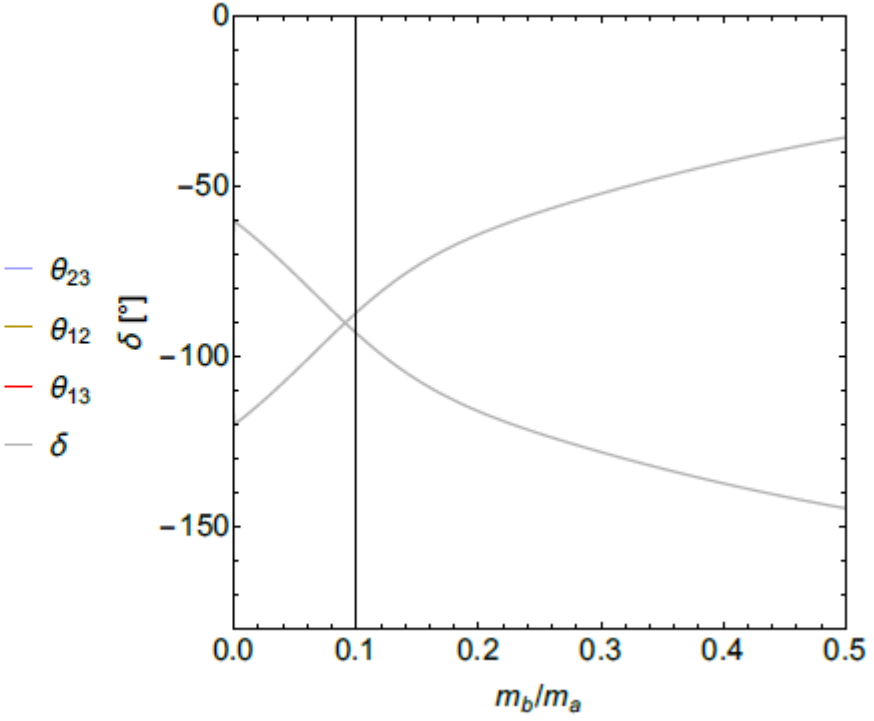
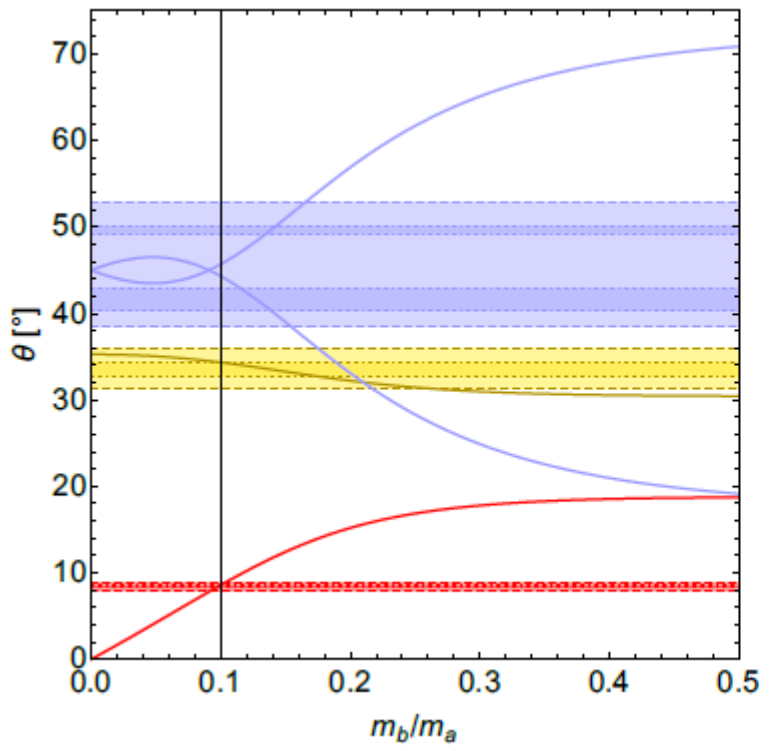
Only one phase

$$\vec{m}_v = m_a \begin{pmatrix} 0 & 0 & 0 \\ 0 & 1 & 1 \\ 0 & 1 & 1 \end{pmatrix} + m_b e^{i\eta} \begin{pmatrix} 1 & n & n-2 \\ n & n^2 & n(n-2) \\ n-2 & n(n-2) & (n-2)^2 \end{pmatrix}$$

m_a and m_b are fixed by the neutrino mass squared differences

➤ Numerical benchmark: $n=3, \eta=2\pi/3$ [Ballett, King, Pascoli et al, 2016]

m_a (meV)	m_b (meV)	η (rad)	θ_{12} ($^\circ$)	θ_{13} ($^\circ$)	θ_{23} ($^\circ$)	δ_{CP} ($^\circ$)	m_1 (meV)	m_2 (meV)	m_3 (meV)
26.57	2.684	$\frac{2\pi}{3}$	34.3	8.67	45.8	-86.7	0	8.59	49.8
Value	from	[25]	$33.48^{+0.78}_{-0.75}$	$8.50^{+0.20}_{-0.21}$	$42.3^{+3.0}_{-1.6}$	-54^{+39}_{-70}	0	8.66 ± 0.10	49.57 ± 0.47

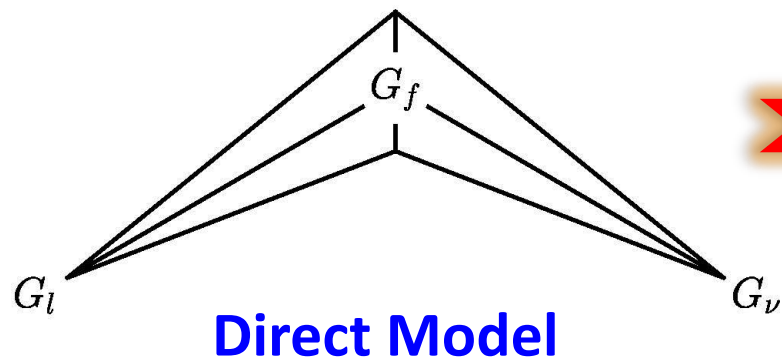


Littlest Seesaw from flavor and CP symmetries

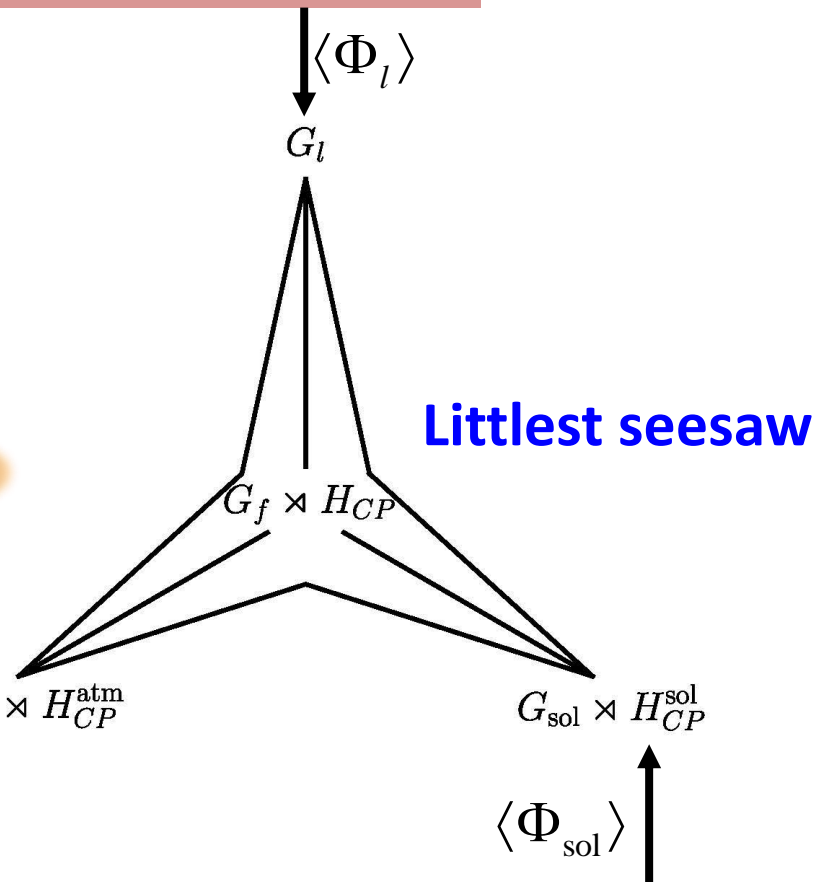
$$m_D = (y_{\text{atm}} \langle \Phi_{\text{atm}} \rangle, y_{\text{sol}} \langle \Phi_{\text{sol}} \rangle),$$

$$m_N = \begin{pmatrix} M_{\text{atm}} & 0 \\ 0 & M_{\text{sol}} \end{pmatrix}$$

$$\mathcal{L}_c = -y_l (L \cdot \Phi_l) l^c$$



Direct Model



Littlest seesaw

$$\mathcal{L}_{\text{atm}} = -\frac{1}{2} M_{\text{atm}} N_{\text{atm}}^c N_{\text{atm}}^c - y_{\text{atm}} (L \cdot \Phi_{\text{atm}}) N_{\text{atm}}^c$$

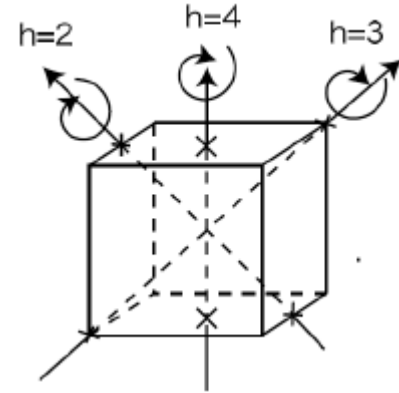
$$\mathcal{L}_{\text{sol}} = -\frac{1}{2} M_{\text{sol}} N_{\text{sol}}^c N_{\text{sol}}^c - y_{\text{sol}} (L \cdot \Phi_{\text{sol}}) N_{\text{sol}}^c$$

Example of new Littlest seesaw in S_4

- Flavor symmetry is S_4 , its generators in the triplet representation 3 and $3'$ are:

$$\omega \equiv e^{2\pi i/3}$$

$$S = \frac{1}{3} \begin{pmatrix} -1 & 2 & 2 \\ 2 & -1 & 2 \\ 2 & 2 & -1 \end{pmatrix}, \quad T = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \omega^2 & 0 \\ 0 & 0 & \omega \end{pmatrix}, \quad U = \mp \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix},$$



Symmetry of a cube

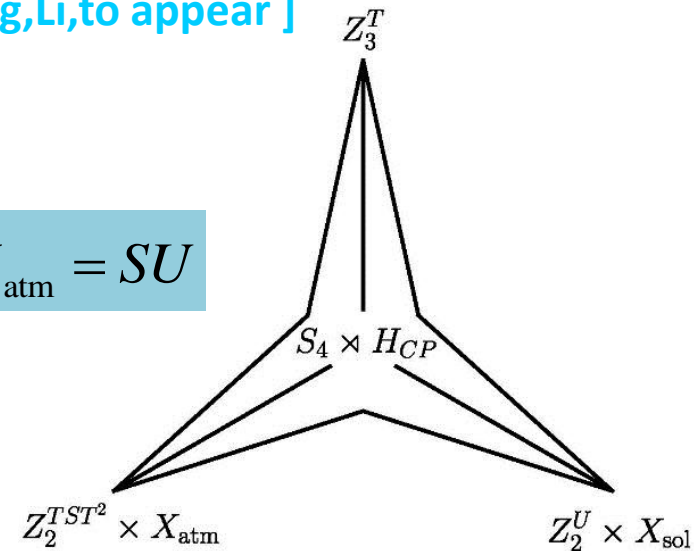
- assignment $L \sim 3, N_{\text{atm}}^c \sim 1, N_{\text{sol}}^c \sim 1', \Phi_{\text{atm}} \sim 3, \Phi_{\text{sol}} \sim 3'$

- S_4 is broken to different subgroups [Ding, King, Li, to appear]

Charged lepton sector: $G_l = Z_3^T$

Atmospheric neutrino sector: $G_{\text{atm}} = Z_2^{TST^2}, X_{\text{atm}} = SU$

Solar neutrino sector: $G_{\text{sol}} = Z_2^U, X_{\text{sol}} = U$



➤ Vacuum alignment

$$\begin{cases} G_{\text{atm}} \langle \Phi_{\text{atm}} \rangle = \langle \Phi_{\text{atm}} \rangle \\ X_{\text{atm}} \langle \Phi_{\text{atm}} \rangle = \langle \Phi_{\text{atm}} \rangle^* \end{cases} \longrightarrow \langle \Phi_{\text{atm}} \rangle \propto (1, \omega^2, \omega)^T$$

$$\begin{cases} G_{\text{sol}} \langle \Phi_{\text{sol}} \rangle = \langle \Phi_{\text{sol}} \rangle \\ X_{\text{sol}} \langle \Phi_{\text{sol}} \rangle = \langle \Phi_{\text{sol}} \rangle^* \end{cases} \longrightarrow \langle \Phi_{\text{sol}} \rangle \propto (1, n, n)^T$$

➤ Mass matrices

$$m_D = \begin{pmatrix} a & b \\ \omega a & nb \\ \omega^2 a & nb \end{pmatrix}, \quad m_N = \begin{pmatrix} M_{\text{atm}} & 0 \\ 0 & M_{\text{sol}} \end{pmatrix}$$

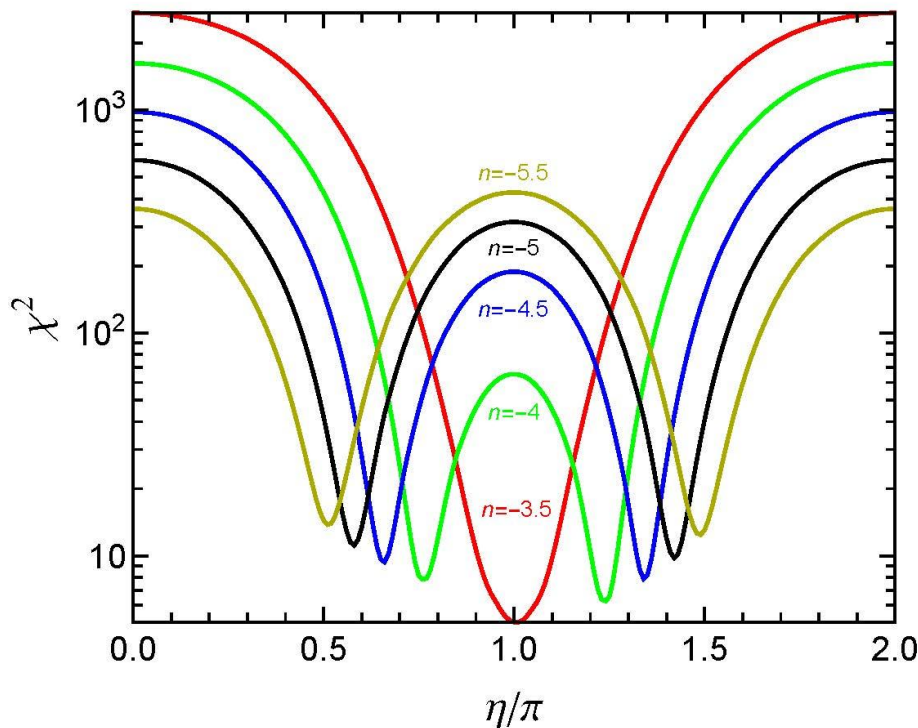
After seesaw

$$m_\nu = m_a \begin{pmatrix} 1 & \omega & \omega^2 \\ \omega & \omega^2 & 1 \\ \omega^2 & 1 & \omega \end{pmatrix} + m_s e^{i\eta} \begin{pmatrix} 1 & n & n \\ n & n^2 & n^2 \\ n & n^2 & n^2 \end{pmatrix}$$

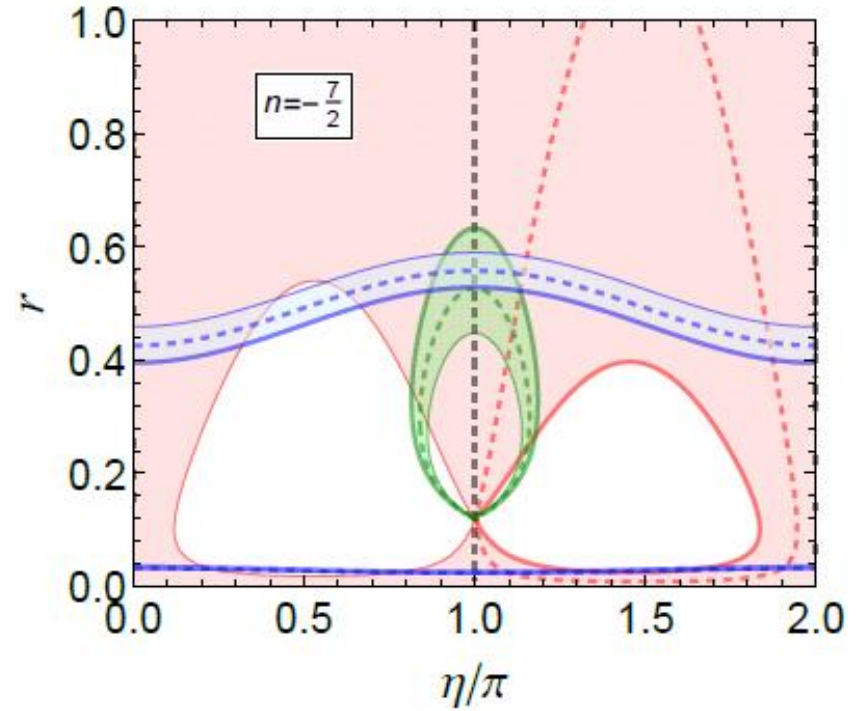
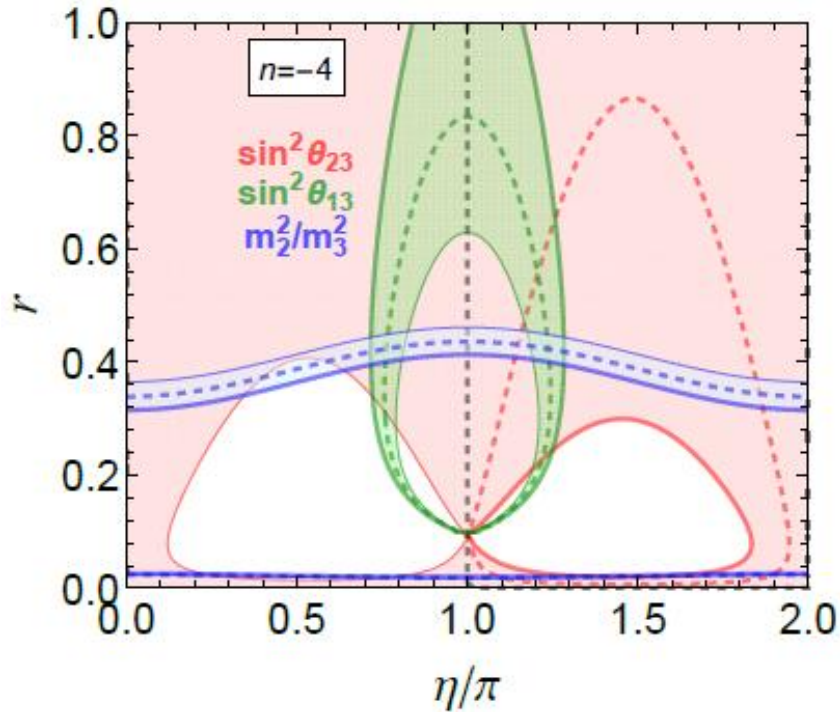
Four parameters m_a, m_s, η, n explain nine observables (3 neutrino masses+ 3 mixing angles+3 CP phases).

$$m_\nu \begin{pmatrix} -i\sqrt{3}n \\ n - \omega^2 \\ -n + \omega \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix} \quad \longrightarrow \quad U_{PMNS} = \begin{pmatrix} \frac{\sqrt{3}n}{\sqrt{5n^2 + 2n + 2}} \times & \times \\ \sqrt{\frac{n^2 + n + 1}{5n^2 + 2n + 2}} \times & \times \\ \sqrt{\frac{n^2 + n + 1}{5n^2 + 2n + 2}} \times & \times \end{pmatrix}$$

The neutrino mass spectrum is **normal ordering** and the lightest one is massless $m_1=0$



n	$\langle \Phi_{\text{sol}} \rangle$	η	m_s/m_a	χ_{min}^2	$\sin^2 \theta_{13}$	$\sin^2 \theta_{12}$	$\sin^2 \theta_{23}$	δ_{CP}/π	β/π	m_2^2/m_3^2
-4	$(1, -4, -4)^T$	$\frac{5\pi}{4}$	0.422	6.140	0.0227	0.336	0.560	-0.412	0.264	0.0292
		$\frac{9\pi}{5}$	0.428	10.489	0.0204	0.338	0.549	-0.424	0.211	0.0290
-5	$(1, -5, -5)^T$	$\frac{7\pi}{5}$	0.266	11.880	0.0210	0.345	0.576	-0.379	0.422	0.0294
		$\frac{10\pi}{7}$	0.264	10.047	0.0225	0.344	0.578	-0.381	0.452	0.0291
$-\frac{7}{2}$	$(1, -\frac{7}{2}, -\frac{7}{2})^T$	π	0.557	4.711	0.0227	0.331	0.5	-0.5	0	0.0294
		$\frac{10\pi}{9}$	0.550	11.708	0.0242	0.330	0.530	-0.459	0.117	0.0297
$-\frac{9}{2}$	$(1, -\frac{9}{2}, -\frac{9}{2})^T$	$\frac{4\pi}{3}$	0.331	8.237	0.0216	0.341	0.571	-0.391	0.352	0.0293
		$\frac{11\pi}{8}$	0.329	13.713	0.0239	0.340	0.575	-0.390	0.396	0.0289



Model of new Littlest seesaw

➤ Field content

[Ding, King, Li, to appear]

matter	L	e^c	μ^c	τ^c	ν_{atm}^c	ν_{sol}^c	$H_{u,d}$
S_4	3	1	1	1	1	1'	1
Z_5	ω_5^4	ω_5	ω_5	ω_5	1	1	1
Z_8	ω_8^7	ω_8^6	ω_8^7	1	ω_8^5	ω_8	1
Z'_8	ω_8	ω_8^7	ω_8^7	ω_8^7	ω_8^5	1	1

driving	ξ_l^0	ϕ_l^0	ϕ_a^0	σ^0	ρ^0	η^0	χ^0	φ^0	Δ^0	κ^0	$\zeta_{1,2,3,4}^0$
S_4	1	3'	3'	2	2	2	3'	3'	3'	1	1
Z_5	1	1	ω_5^3	ω_5^2	ω_5	ω_5^4	1	1	ω_5^2	ω_5^3	1
Z_8	ω_8^6	ω_8^6	1	1	1	1	1	1	1	1	1
Z'_8	1	1	ω_8^4	ω_8^6	ω_8^6	ω_8^6	ω_8^6	ω_8^7	ω_8^4	ω_8^2	1

flavon	η_l	ϕ_l	ξ_a	ϕ_a	ξ_s	η_s	χ_s	φ_s	Δ_s	ϕ_s	ψ_s
S_4	2	3	1	3	1	2	3'	3'	3'	3'	3'
Z_5	1	1	1	ω_5	1	ω_5^3	ω_5^2	1	ω_5^3	ω_5	ω_5^4
Z_8	ω_8	ω_8	ω_8^6	-1	ω_8^6	1	1	1	1	1	1
Z'_8	1	1	ω_8^6	ω_8^2	1	ω_8^5	ω_8^5	ω_8	ω_8^4	ω_8^7	ω_8^5

$$\langle \eta_l \rangle = (0, v_{\eta_l})^T, \quad \langle \phi_l \rangle = (0, v_{\phi_l}, 0)^T \Rightarrow G_l = Z_3^T$$

$$\langle \xi_a \rangle = v_{\xi_a}, \quad \langle \phi_a \rangle = v_{\phi_a} (1, \omega^2, \omega)^T \Rightarrow G_{\text{atm}} = Z_2^{TST^2}$$

$$\begin{aligned} \langle \chi_s \rangle &= v_{\chi_s} (1, 0, 0)^T, & \langle \psi_s \rangle &= v_{\psi_s} (1, -2, -2)^T \\ \langle \varphi_s \rangle &= v_{\varphi_s} (1, -1, -1)^T, & \langle \Delta_s \rangle &= v_{\Delta_s} (1, 2, 2)^T, \\ \langle \eta_s \rangle &= v_{\eta_s} (1, 1)^T, & \langle \xi_s \rangle &= v_{\xi_s}, & \langle \phi_s \rangle &= v_{\phi_s} (1, -4, -4)^T \end{aligned} \Rightarrow G_{\text{sol}} = Z_2^U$$

➤ Charged lepton sector

$$w_l = \frac{y_\tau}{\Lambda} (L\phi_l)_1 \tau^c H_d + \frac{y_\mu}{\Lambda^2} (L(\eta_l\phi_l)_3)_1 \mu^c H_d + \frac{y_{e,1}}{\Lambda^3} (L\eta_l\eta_l\phi_l) e^c H_d \\ + \frac{y_{e,2}}{\Lambda^3} (L\eta_l\phi_l\phi_l) e^c H_d + \frac{y_{e,3}}{\Lambda^3} (L\phi_l\phi_l\phi_l) e^c H_d$$

✓ Charged lepton mass matrix is diagonal

✓ Hierarchical charged lepton masses

$$m_\tau = \left| y_\tau \frac{v_{\phi_l}}{\Lambda} \right| v_d, \quad m_\mu = \left| y_\mu \frac{v_{\eta_l} v_{\phi_l}}{\Lambda^2} \right| v_d, \\ m_e = \left| (y_{e,1} v_{\eta_l}^2 - 2y_{e,2} v_{\eta_l} v_{\phi_l} + y_{e,3} v_{\phi_l}^2) \frac{v_{\phi_l}}{\Lambda^3} \right| v_d$$

✓ Higher order corrections appear at relative order $1/\Lambda^3$

➤ Neutrino sector

$$w_\nu = \frac{y_a}{\Lambda} (L\phi_a) \nu_{\text{atm}}^c H_u + \frac{y_s}{\Lambda} (L\phi_s) \nu_{\text{sol}}^c H_u + x_a \xi_a \nu_{\text{atm}}^c \nu_{\text{atm}}^c + x_s \xi_s \nu_{\text{sol}}^c \nu_{\text{sol}}^c$$

After flavor and EW symmetry breaking, the neutrino Dirac and Majorana mass matrices read

$$m_D = \begin{pmatrix} a & b \\ \omega a & -4b \\ \omega^2 a & -4b \end{pmatrix}, \quad m_N = \begin{pmatrix} M_{\text{atm}} & 0 \\ 0 & M_{\text{sol}} \end{pmatrix}$$

with

$$a = \frac{y_a v_{\phi_a}}{\Lambda}, \quad b = \frac{y_s v_{\phi_s}}{\Lambda}, \quad M_{\text{atm}} = x_a v_{\xi_a}, \quad M_{\text{sol}} = x_s v_{\xi_s}$$

new Littlest seesaw texture is reproduced exactly.

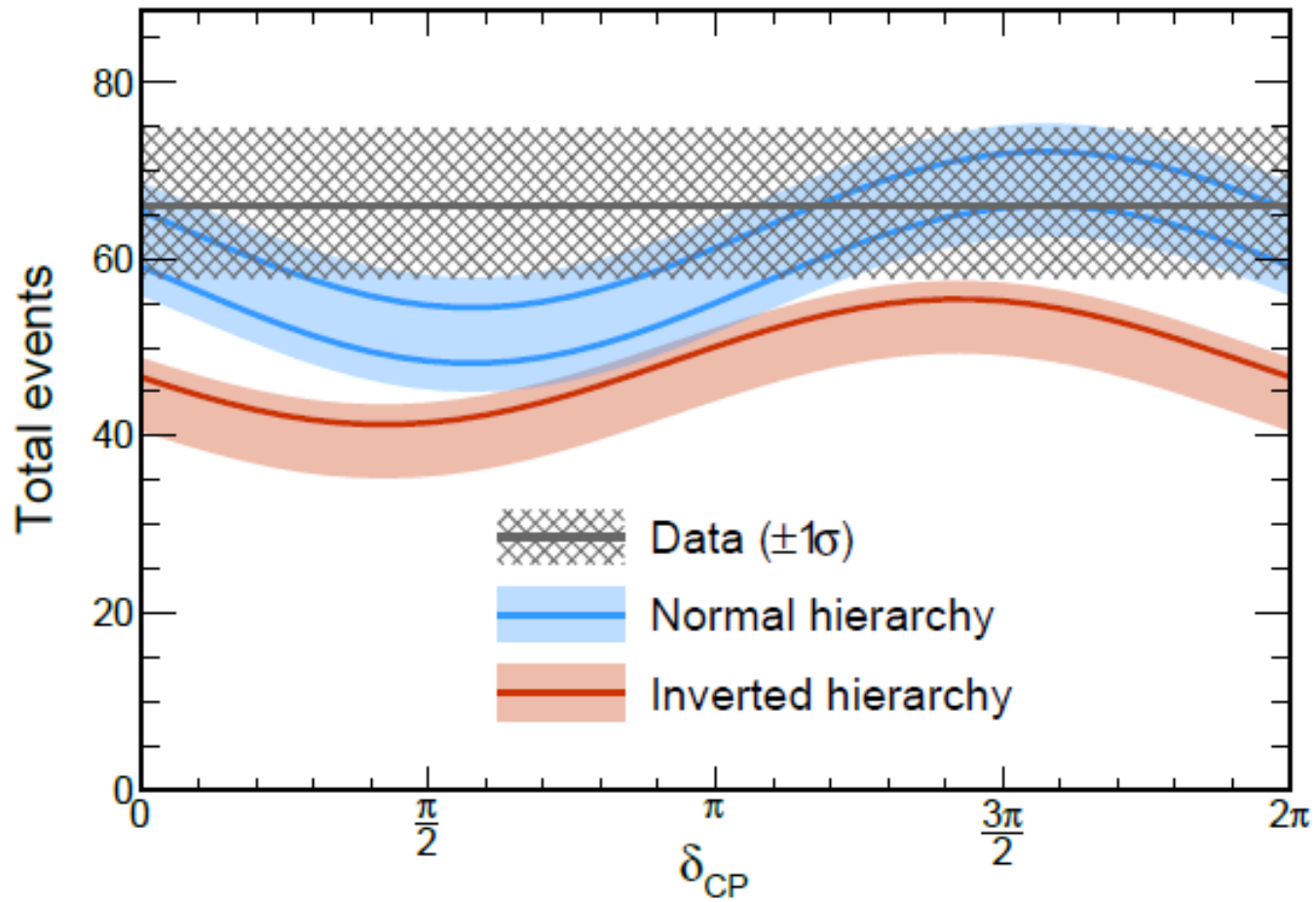
Summary

- ✚ Generalized $\mu\tau$ reflection is a simple alternative to understand deviations of θ_{23} and δ_{CP} from maximal values.
- ✚ The drastically different quark and lepton mixing patterns can be explained from the same flavor symmetry combined with CP, if the residual symmetry in quark sectors is $Z_2 \times CP$. $\Delta(6 \cdot 7^2) = \Delta(294)$ is a good candidate for such flavor symmetry.
- ✚ Littlest seesaw inspired by flavor and CP symmetries (or indirect CP approach) provides new model building opportunity. New mixing patterns can be found, and an explicit model with S_4 and CP is constructed.

Thank you for your attention!

Backup

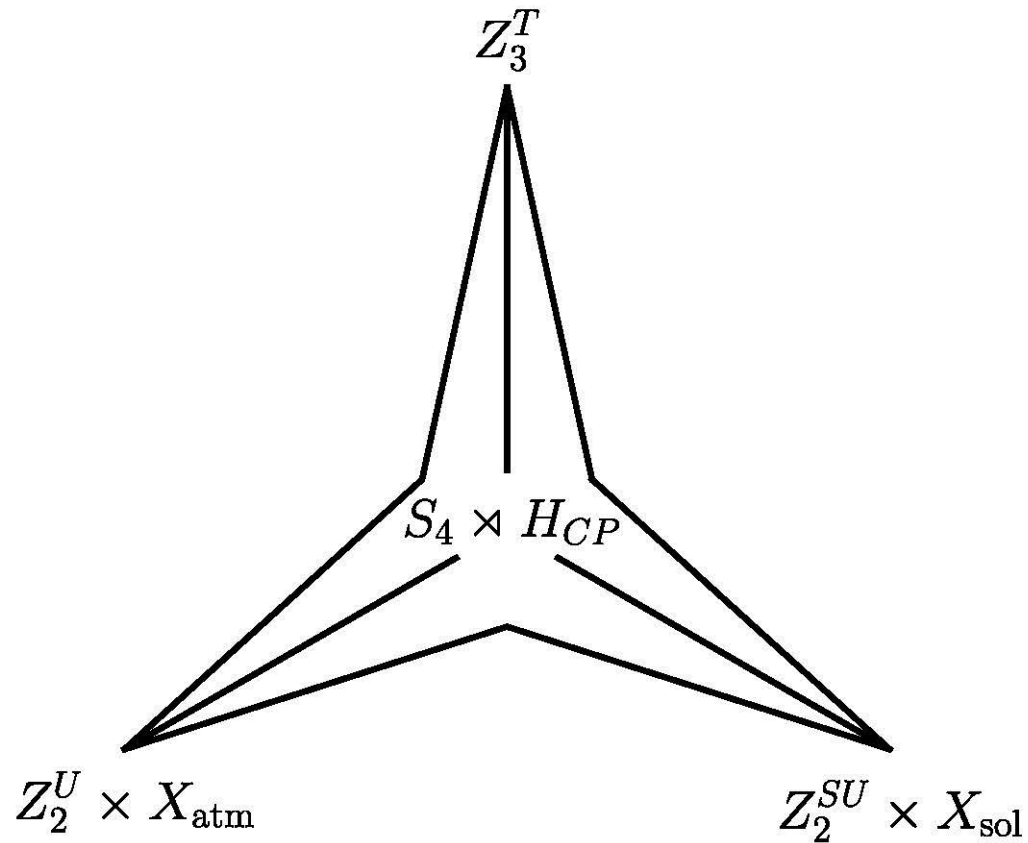
NOvA, Acero et al, 1806.00096



➤ Littlest seesaw

$$\langle \Phi_{\text{atm}} \rangle = v_{\text{atm}} \begin{pmatrix} 0 \\ 1 \\ 1 \end{pmatrix}, \quad \langle \Phi_{\text{sol}} \rangle = v_{\text{sol}} \begin{pmatrix} 1 \\ n \\ n-2 \end{pmatrix}$$

$$X_{\text{atm}} = 1, \quad X_{\text{sol}} = 1$$



The overall phase of the vacuum

$$w_d^{\text{phase}} = M_{\xi_s}^2 \zeta_1^0 + \frac{g_1}{\Lambda^2} \zeta_1^0 \xi_s^4 + M_{\xi_a}^2 \zeta_2^0 + \frac{g_2}{\Lambda^2} \zeta_2^0 \xi_a^4 + M_1^2 \zeta_3^0 + \frac{1}{\Lambda^2} \zeta_3^0 [g_3 (\eta_s \eta_s \varphi_s \psi_s)_1 \\ + g_4 (\chi_s \varphi_s \psi_s \psi_s)_1 + g_5 (\phi_s \phi_s \psi_s \psi_s)_1] + M_2^2 \zeta_4^0 + \frac{g_6}{\Lambda^2} \zeta_4^0 (\chi_s \phi_s \phi_a \phi_a)_1 ,$$