

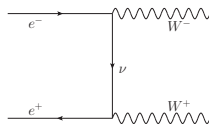
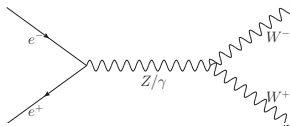
Triple Gauge Couplings (at FCC-ee)

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WG1 & WG2 working meeting
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TGC from $e^+e^- \rightarrow WW$



- ▶ We can measure the triple gauge couplings (TGC) from the diboson process $e^+e^- \rightarrow WW$.
- ▶ Focusing on the anomalous TGC parameters generated by CP-even D6 operators,

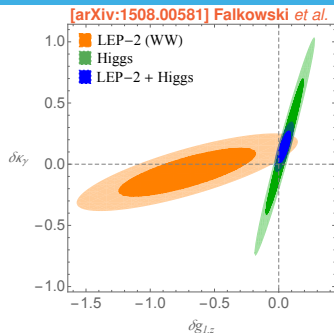
$$\begin{aligned}
 \mathcal{L}_{\text{TGC}} = & \quad ig s_{\theta_W} A^\mu (W^{-\nu} W_{\mu\nu}^+ - W^{+\nu} W_{\mu\nu}^-) \\
 & + ig(1 + \delta g_1^Z) c_{\theta_W} Z^\mu (W^{-\nu} W_{\mu\nu}^+ - W^{+\nu} W_{\mu\nu}^-) \\
 & + ig [(1 + \delta \kappa_Z) c_{\theta_W} Z^{\mu\nu} + (1 + \delta \kappa_\gamma) s_{\theta_W} A^{\mu\nu}] W_\mu^- W_\nu^+ \\
 & + \frac{ig}{m_W^2} (\lambda_Z c_{\theta_W} Z^{\mu\nu} + \lambda_\gamma s_{\theta_W} A^{\mu\nu}) W_\nu^{-\rho} W_{\rho\mu}^+, \tag{1}
 \end{aligned}$$

TGC

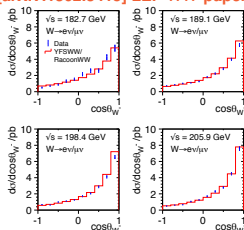
- ▶ Imposing Gauge invariance one obtains $\delta\kappa_Z = \delta g_{1,Z} - t_{\theta_W}^2 \delta\kappa_\gamma$ and $\lambda_Z = \lambda_\gamma$.
- ▶ 3 aTGCs parameters: $\delta g_{1,Z}$, $\delta\kappa_\gamma$ and λ_Z !
- ▶ They correspond to 3 linear combinations of Wilson coefficients in the EFT not constrained by Z-pole measurements.
 - ▶ Even if the Z-pole measurements are infinitely precise, the 3 aTGCs remain unconstrained!
 - ▶ In a convenient basis they correspond to $\mathcal{O}_{HW} = ig(D^\mu H)^\dagger \sigma^a (D^\nu H) W_{\mu\nu}^a$, $\mathcal{O}_{HB} = ig'(D^\mu H)^\dagger (D^\nu H) B_{\mu\nu}$ and $\mathcal{O}_{3W} = \frac{1}{3!} g\epsilon_{abc} W_{\mu\nu}^a W_{\nu\rho}^b W^{c\rho\mu}$.
See e.g. [arXiv:1411.0669] Falkowski & Riva for a general parameterization.
- ▶ $\delta g_{1,Z}$ and $\delta\kappa_\gamma$ are generated by operators that also contribute to Higgs couplings.

Impacts on EFT fits, LHC + LEP

- ▶ $\delta g_{1,Z}$ and $\delta \kappa_\gamma$ are generated by operators that also contribute to Higgs couplings.
- ▶ Higgs better measured \Rightarrow Higgs helps TGC;
- ▶ TGC better measured \Rightarrow TGC helps Higgs.



[arXiv:1302.3415] LEP WW paper



- ▶ Note: LEP bounds should have been better!
 - ▶ LEP did not perform global fits with all 3 aTGCs.
 - ▶ The distributions of W decay angles were not provided.

FCC-ee



W-pairs at FCCee :the OkuW

$$\sqrt{s} \approx 91 \text{ GeV} \sim 10^{12} Z$$

$\sqrt{s}=161$: $L \sim 3 \cdot 10^{35}$ collect 8/ab
30 10^6 WW decays

$\sqrt{s}=240$: $L \sim 0.7 \cdot 10^{35}$ collect 5/ab
80 10^6 WW decays

$\sim 10^6 H$

$\sqrt{s}=350$: $L \sim 10^{34}$ collect 1.5/ab
15 10^6 WW decays

$\sim 2 \times 10^5 H$

FCC-ee



W-pairs at FCCee :the OkuW

$\sqrt{s} \approx 91 \text{ GeV} \sim 10^{12} Z$
(systematics dominated)

$\sqrt{s}=161 : L \sim 3 \cdot 10^{35}$ collect 8/ab
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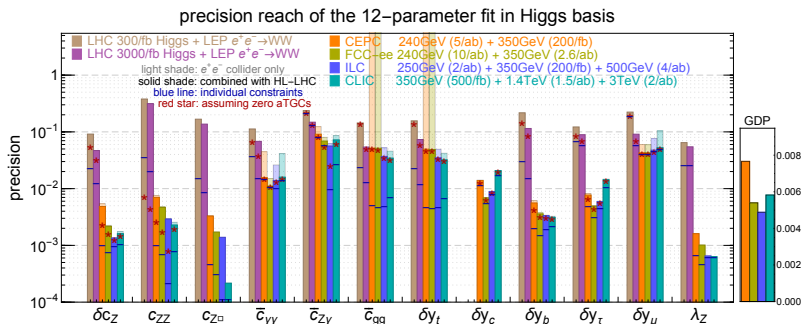
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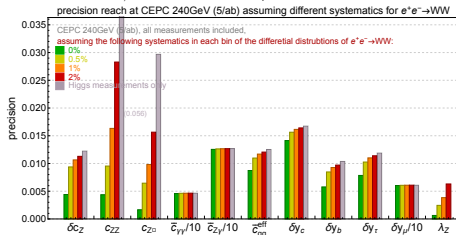
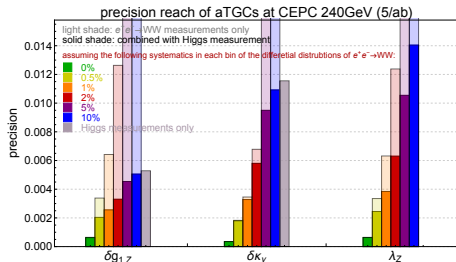
of WW $\sim 10^2 \times$ # of H

EFT fits at future lepton colliders [\[arXiv:1704.02333\]](https://arxiv.org/abs/1704.02333) G. Durieux, C. Grojean, J.G, K. Wang



- ▶ $\delta g_{1,Z}, \delta \kappa_\gamma \rightarrow c_{ZZ}, c_{Z\Box}, c_{\gamma\gamma}, c_{Z\gamma}$
- ▶ see also
 - ▶ [arXiv:1510.04561, 1701.04804] Ellis *et al.*,
 - ▶ [arXiv:1708.08912, 1708.09079] Peskin *et al.*
 - ▶ [arXiv:1711.04046] Wang *et al.*

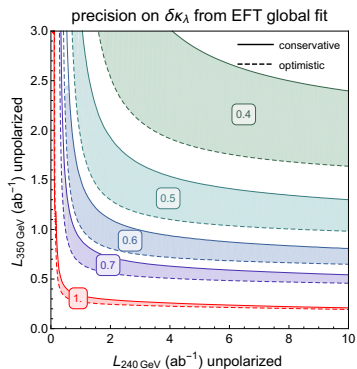
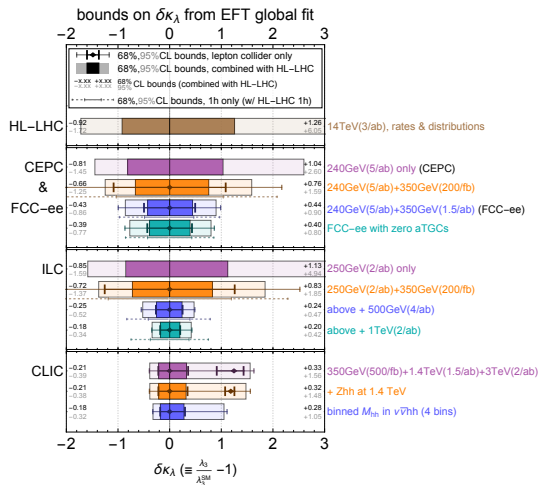
Impact on EFT fits from TGC measurements



- ▶ $\delta g_{1,Z}$, $\delta \kappa_\gamma \leftrightarrow$
 c_{ZZ} , $c_{Z\Box}$, $c_{\gamma\gamma}$, $c_{Z\gamma}$
- ▶ We try different assumptions on the systematic uncertainties (in each bin with the differential distribution divided into 20 bins).
- ▶ Detailed study of $e^+e^- \rightarrow WW$ required to estimate the systematic uncertainties!

Impact on the triple Higgs coupling

[arXiv:1711.03978] S. Di Vita, G. Durieux, C. Grojean, J.G. Z. Liu, G. Panico, M. Riemann, T. Vantalon



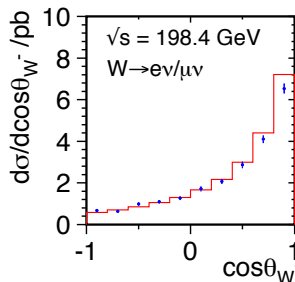
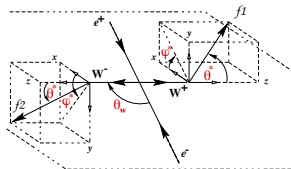
- ▶ Circular colliders can probe the triple Higgs coupling via its loop contributions.

towards a real TGC analysis...

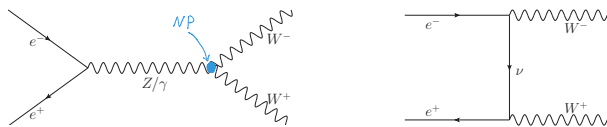
- ▶ We need a real TGC analysis for circular colliders (CEPC or FCC-ee)!
 - ▶ One-sigma bounds & correlations from global fits!
- ▶ Rate measurements are not enough.
 - ▶ At least 3 measurements are needed to independently constrain 3 aTGCs. (OK for FCC-ee)
 - ▶ The aTGCs are sensitive to the angular distributions. (e.g. The region with small θ_{W^-} is dominated by the t -channel diagram which does not contain the TGC vertex.)
- ▶ What's the best way to extract information from the differential (angular) observables?
- ▶ Which channels to use?
 - ▶ **semi-leptonic** ($\sim 42\%$ BR, good reconstruction)
 - ▶ hadronic (hard to reconstruct, angles are folded)
 - ▶ di-leptonic (can solve the neutrino momenta but with ambiguities)

How to extract information from the angles

- ▶ 5 angles (1 production, 2 decays for each W)
- ▶ 1D histogram(s)
 - ▶ easy to do for θ_{W^-}
 - ▶ can also include decay angles, but the correlations among angles are not taken account of.
- ▶ 5D histogram (would need a very large simulated sample)
 - ▶ or 3D histogram focusing on the leptonic decay angles (used in the ILC TGC analysis)
- ▶ fancier methods (see e.g. [hep-ph/9601233](#))
 - ▶ likelihood (unbinned)
 - ▶ optimal observables
- ▶ ??



The TGC dominance assumption



- ▶ Assumption: New physics contributes to $e^+ e^- \rightarrow WW$ only through the TGC vertex.
- ▶ Reality: In principle there can be many other contributions!
- ▶ Other contributions are constrained by Z -pole measurements.
 - ▶ With the Z -pole run, the TGC dominance assumption should be valid at FCC-ee.
 - ▶ Ultimately, a full EFT analysis is desired... (Z -pole + WW + Higgs)

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Conclusion

- ▶ TGC measurements are important (e.g. for the EFT fit of Higgs couplings)!
- ▶ We need to find a good way to extract the information from angular distributions.
- ▶ The TGC dominance assumption is probably OK.
- ▶ A global fit to the $e^+e^- \rightarrow WW$ data with all 3 aTGCs parameters, $\delta g_{1,Z}$, $\delta\kappa_\gamma$ and λ_Z !
 - ▶ One-sigma bounds & correlations!

backup slides

The “12-parameter” framework in the Higgs basis

- ▶ The relevant terms in the EFT Lagrangian are

$$\mathcal{L} \supset \mathcal{L}_{hVV} + \mathcal{L}_{hff} + \mathcal{L}_{\text{tgc}}, \quad (2)$$

- ▶ the Higgs couplings with a pair of gauge bosons

$$\begin{aligned} \mathcal{L}_{hVV} = & \frac{h}{v} \left[(1 + \delta c_W) \frac{g^2 v^2}{2} W_\mu^+ W_\mu^- + (1 + \delta c_Z) \frac{(g^2 + g'^2) v^2}{4} Z_\mu Z_\mu \right. \\ & + c_{WW} \frac{g^2}{2} W_{\mu\nu}^+ W_{\mu\nu}^- + c_{W\Box} g^2 (W_\mu^- \partial_\nu W_{\mu\nu}^+ + \text{h.c.}) \\ & + c_{gg} \frac{g_s^2}{4} G_{\mu\nu}^a G_{\mu\nu}^2 + c_{\gamma\gamma} \frac{e^2}{4} A_{\mu\nu} A_{\mu\nu} + c_{Z\gamma} \frac{e\sqrt{g^2 + g'^2}}{2} Z_{\mu\nu} A_{\mu\nu} \\ & \left. + c_{ZZ} \frac{g^2 + g'^2}{4} Z_{\mu\nu} Z_{\mu\nu} + c_{Z\Box} g^2 Z_\mu \partial_\nu Z_{\mu\nu} + c_{\gamma\Box} gg' Z_\mu \partial_\nu A_{\mu\nu} \right]. \quad (3) \end{aligned}$$

The “12-parameter” framework in the Higgs basis

- ▶ Not all the couplings are independent, for instance one could write the following couplings as

$$\begin{aligned}
 \delta c_W &= \delta c_Z + 4\delta m, \\
 c_{WW} &= c_{ZZ} + 2s_{\theta_W}^2 c_{Z\gamma} + s_{\theta_W}^4 c_{\gamma\gamma}, \\
 c_{W\Box} &= \frac{1}{g^2 - g'^2} \left[g^2 c_{Z\Box} + g'^2 c_{ZZ} - e^2 s_{\theta_W}^2 c_{\gamma\gamma} - (g^2 - g'^2) s_{\theta_W}^2 c_{Z\gamma} \right], \\
 c_{\gamma\Box} &= \frac{1}{g^2 - g'^2} \left[2g^2 c_{Z\Box} + (g^2 + g'^2) c_{ZZ} - e^2 c_{\gamma\gamma} - (g^2 - g'^2) c_{Z\gamma} \right], \quad (4)
 \end{aligned}$$

- ▶ we only consider the diagonal elements in the Yukawa matrices relevant for the measurements considered,

$$\mathcal{L}_{hff} = -\frac{h}{v} \sum_{f=t,c,b,\tau,\mu} m_f (1 + \delta y_f) \bar{f}_R f_L + \text{h.c.} . \quad (5)$$

TGC

$$\begin{aligned}
\mathcal{L}_{\text{TGC}} = & \quad ig s_{\theta_W} A^\mu (W^{-\nu} W_{\mu\nu}^+ - W^{+\nu} W_{\mu\nu}^-) \\
& + ig(1 + \delta g_1^Z) c_{\theta_W} Z^\mu (W^{-\nu} W_{\mu\nu}^+ - W^{+\nu} W_{\mu\nu}^-) \\
& + ig [(1 + \delta \kappa_Z) c_{\theta_W} Z^{\mu\nu} + (1 + \delta \kappa_\gamma) s_{\theta_W} A^{\mu\nu}] W_\mu^- W_\nu^+ \\
& + \frac{ig}{m_W^2} (\lambda_Z c_{\theta_W} Z^{\mu\nu} + \lambda_\gamma s_{\theta_W} A^{\mu\nu}) W_\nu^{-\rho} W_{\rho\mu}^+, \tag{6}
\end{aligned}$$

- ▶ $V_{\mu\nu} \equiv \partial_\mu V_\nu - \partial_\nu V_\mu$ for $V = W^\pm, Z, A$. Imposing Gauge invariance one obtains $\delta \kappa_Z = \delta g_{1,Z} - t_{\theta_W}^2 \delta \kappa_\gamma$ and $\lambda_Z = \lambda_\gamma$.
- ▶ 3 aTGCs parameters $\delta g_{1,Z}, \delta \kappa_\gamma$ and λ_Z , 2 of them related to Higgs observables by

$$\begin{aligned}
\delta g_{1,Z} &= \frac{1}{2(g^2 - g'^2)} \left[-g^2(g^2 + g'^2) c_{Z\Box} - g'^2(g^2 + g'^2) c_{ZZ} + e^2 g'^2 c_{\gamma\gamma} + g'^2(g^2 - g'^2) c_{Z\gamma} \right] \\
\delta \kappa_\gamma &= -\frac{g^2}{2} \left(c_{\gamma\gamma} \frac{e^2}{g^2 + g'^2} + c_{Z\gamma} \frac{g^2 - g'^2}{g^2 + g'^2} - c_{ZZ} \right). \tag{7}
\end{aligned}$$