

# Implications of neutrino generalized interactions in CEvNS

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Coherent Elastic  
Neutrino-Nucleus Scattering  
(CEvNS)

- CEvNS
- Relevant neutrino sources
- COHERENT

Sensitivity to new physics

What kind of physics?

Summary

# Coherent Elastic Neutrino-Nucleus Scattering (CEvNS)

CE $\nu$ NS occurs when the neutrino energy  $E_\nu$  is such that nucleon amplitudes sum up coherently  $\Rightarrow$  cross section enhancement

$$\lambda \gtrsim R_N \Rightarrow q \lesssim 200 \text{ MeV}$$

$$E_R = q^2/2m_N \Rightarrow E_\nu \approx \sqrt{E_R^{\max} m_N/2}$$

$$E_\nu \lesssim 100 \text{ MeV}$$

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What kind of physics?

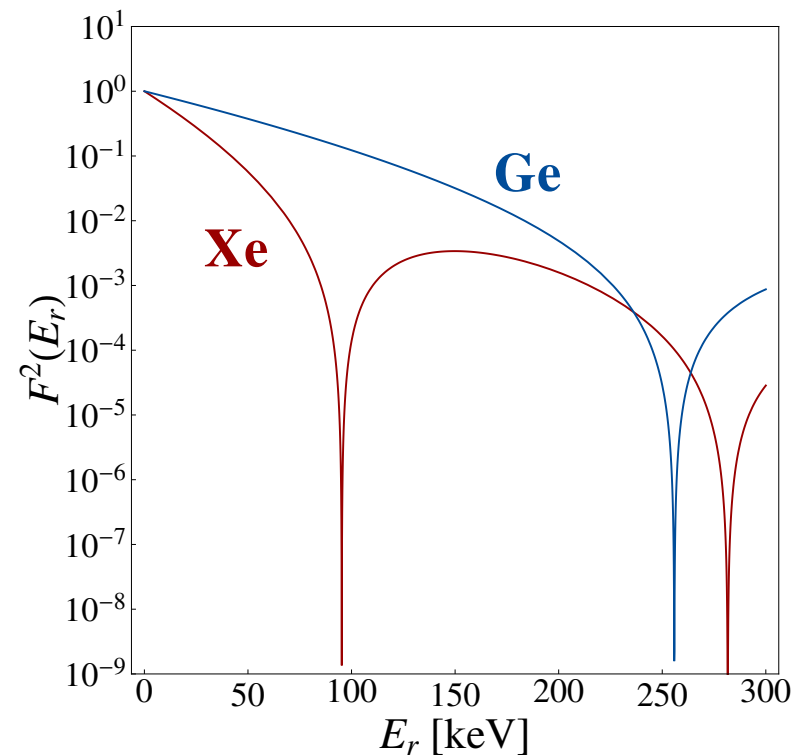
Summary

**Freedman, 1974**

$$\frac{d\sigma_\nu}{dE_R} = \frac{G_F^2}{4\pi} Q_{\text{SM}}^2 m_N \left(1 - \frac{E_R m_N}{2E_\nu^2}\right) \underbrace{F^2(E_R)}_{\text{Form factor}}$$

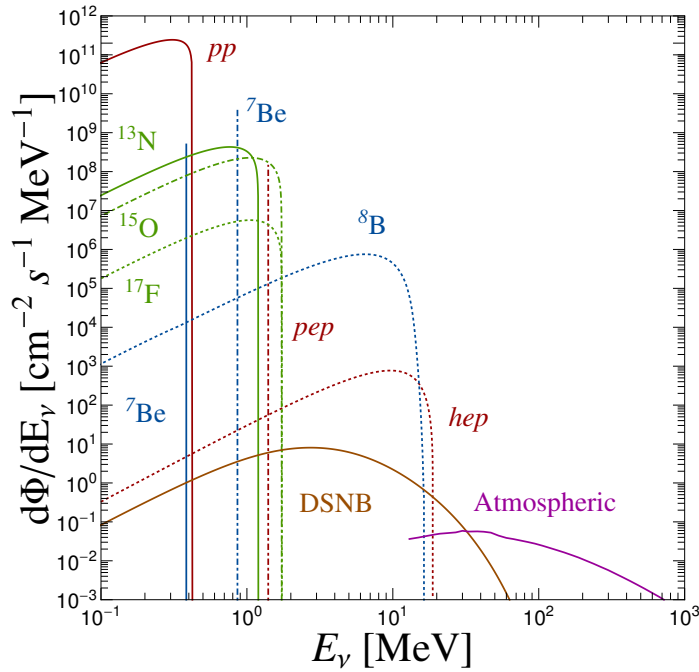
$$Q_{\text{SM}}^2 = [N - (1 - s_W^2)Z]^2 \approx N^2$$

**Helm, 1956**

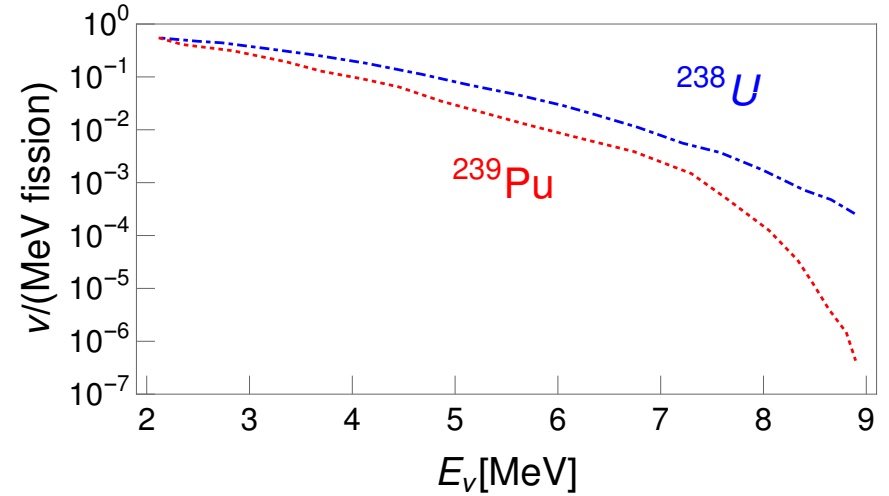


# Relevant neutrino sources

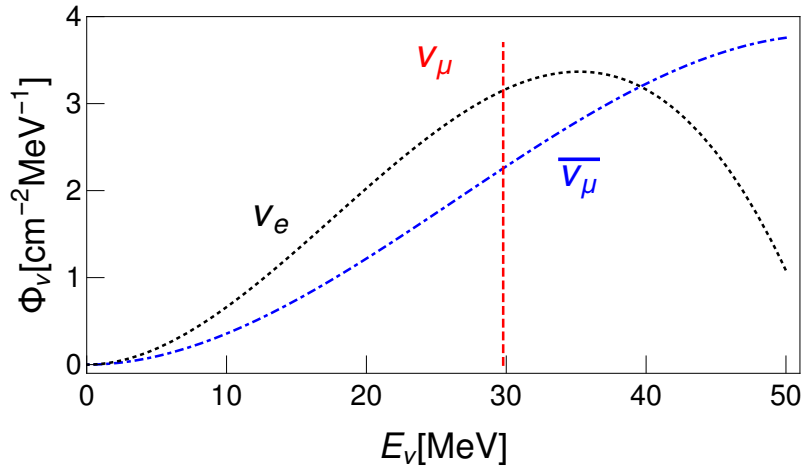
## “Astrophysical” sources



## Reactor Neutrinos



## Fixed target



Solar+Atm:  $\nu$  backgrounds DM detectors

Reactor: Basis for CONUS,  $\nu$ -CLEUS

Fixed target: COHERENT experiment

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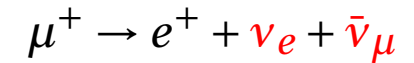
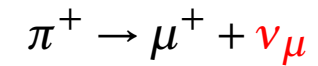
# COHERENT

CE $\nu$ NS observed by COHERENT more than 40 years after its prediction

Akimov et. al. 2017

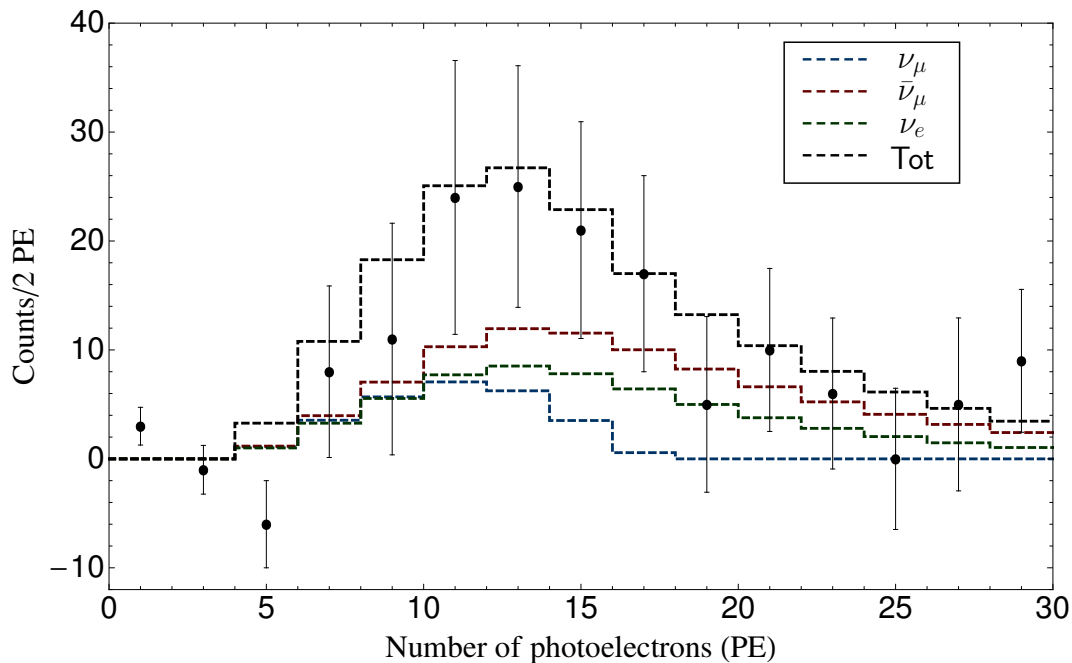
COHERENT uses neutrinos produced in SNS

@ Oak Ridge National Laboratory in the collision  $p - \text{Hg}$



Presence of CE $\nu$ NS favored @ the  $6.7\sigma$  level. Data consistent with SM @ the  $1\sigma$

DAS, De Romeri, Rojas, 2018



$$n_{\text{PE}} = 1.17 (E_R / \text{keV})$$

There is still some room  
for NEW PHYSICS!

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Sensitivity to new physics

- The case of NSI
- Constraints
- The NGI case
- Constraints from oscillations
- Parameter space scenarios
- One-parameter analysis
- Vector -vs- SM+vector

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## Sensitivity to new physics

Non-standard interactions parametrized in a model-independent and phenomenological way

Wolfenstein, 1978

$$\mathcal{L} \sim G_F \sum_{q=u,d} \bar{\nu}_i (1 - \gamma_5) \gamma_\mu \nu_j \bar{q} (\epsilon_{ij}^{qV} - \epsilon_{ij}^{qA} \gamma_5) \gamma^\mu q$$

**Phenomenological constraints from forward coherent scattering  
(matter potentials) DIS and COHERENT data**

Scenarios

Gonzalez-Garcia et. al, 2017

- For  $m_X^2 \ll q^2$  contributions of NSI to DIS are suppressed,  $q_{\text{DIS}}^2 \gtrsim (10\text{GeV})^2$
- Light mediator scenarios:  $M_X \in [10, 10^3] \text{ MeV} \Rightarrow$  DIS constraints evaded
- Heavy mediator scenarios:  $M_X \in [1, 10^3] \text{ GeV}$  all constraints apply

**COHERENT constraints are particularly relevant for light mediators**

# Constraints

COHERENT data has been used to constraint NSI contributions to the  $CE\nu NS$

**Gonzalez-Garcia et. al, 2017**

**J. Liao & D. Marfatia, 2017**

**Kosmas et. al, 2018**

**Billard et. al, 2018**

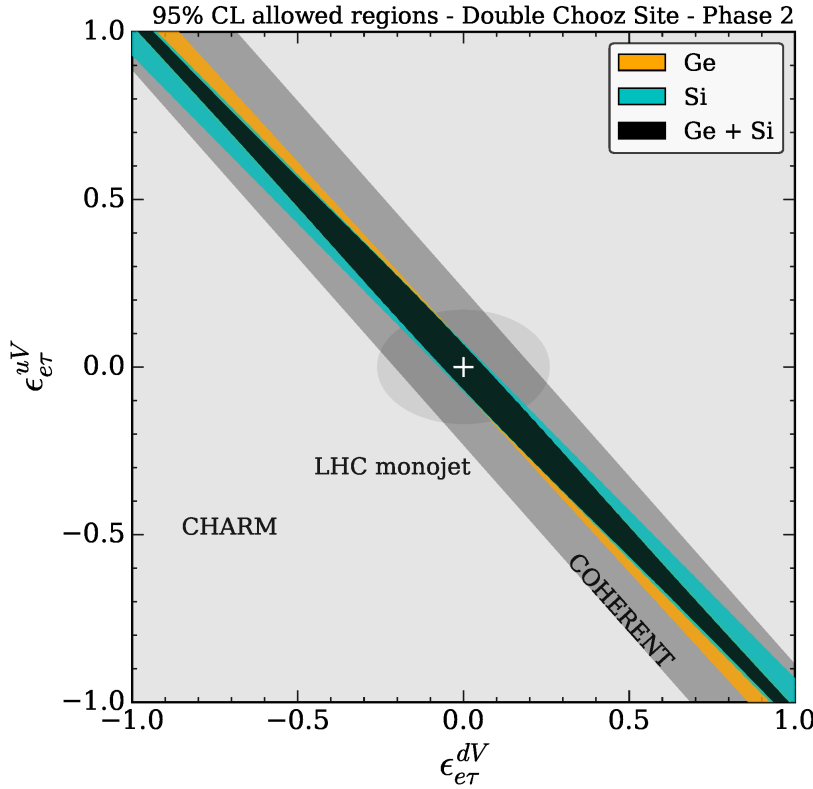
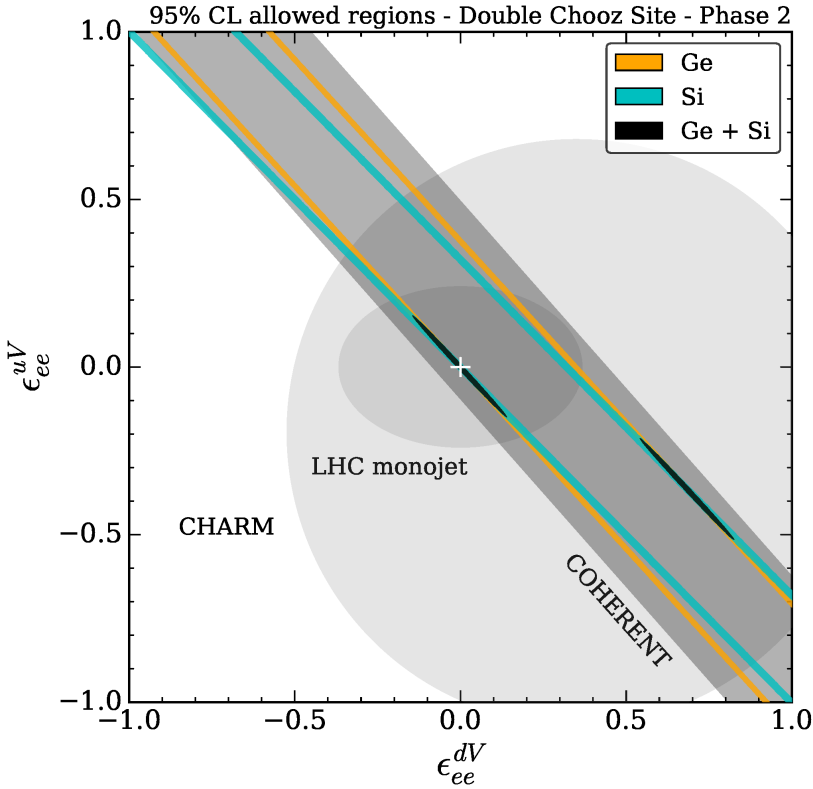
Coherent Elastic Neutrino-Nucleus Scattering (CEvNS)

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NSI are a subset of a larger set of neutrino-quark interactions: Neutrino Generalized Interactions (NGI) Lee & Yang (1957)

$$\mathcal{L} \sim G_F \sum_{q=u,d} (\bar{\nu} \Gamma_A \nu) \left[ \bar{q} \Gamma_A (C_A^q + i D_A^q \gamma_5) q \right]$$

$$\Gamma_A = \{ \mathbb{1}, i\gamma_5, \gamma_\mu, \gamma_5 \gamma_\mu, \sigma_{\mu\nu} \}$$

**Diagonal and non-diagonal LS**

$$\Gamma_P : \mathcal{L} \sim \bar{\nu} \gamma_5 \nu \bar{q} \left( \gamma_5 C_P^q + \mathbb{1} D_P^q \right) q$$

$P$  and  $A$  quark currents are nuclear spin-dependent  $\Rightarrow Z_\uparrow - Z_\downarrow, N_\uparrow - N_\downarrow$

$$\mathcal{L}_S \sim (\bar{\nu} \nu) \left[ \bar{q} \left( C_S^q + i \gamma_5 D_S^q \right) q \right]$$

$$\mathcal{L}_P \sim (\bar{\nu} \gamma_5 \nu) \left[ \bar{q} \left( \gamma_5 C_P^q + i D_P^q \right) q \right]$$

$$\mathcal{L}_V \sim (\bar{\nu} \gamma^\mu \nu) \left[ \bar{q} \left( \gamma_\mu C_V^q + i \gamma_\mu \gamma_5 D_V^q \right) q \right]$$

$$\mathcal{L}_A \sim (\bar{\nu} \gamma^\mu \gamma_5 \nu) \left[ \bar{q} \left( \gamma_\mu \gamma_5 C_A^q + i \gamma_\mu D_A^q \right) q \right]$$

$$\mathcal{L}_T \sim (\bar{\nu} \sigma^{\mu\nu} \nu) \left[ \bar{q} \left( \sigma_{\mu\nu} C_T^q + i \sigma_{\mu\nu} \gamma_5 D_T^q \right) q \right]$$

$$\mathcal{P}_1 = \{ C_S^q, D_P^q, C_V^q, D_A^q, C_T^q \} \quad \checkmark$$

$$\mathcal{P}_2 = \{ C_P^q, D_S^q, C_A^q, D_V^q, D_T^q \} \quad \times$$

**Constraints on  $\mathcal{P}_2$  are weak!**

# Constraints from oscillations

Constraints from forward coherent scattering are only relevant for vector interactions

## Matter potentials

Bergmann, Grossman, Nardi, 1999

$$\mathcal{L}_{\text{int}} \sim \sum_{a,f} (\bar{\nu} \Gamma^a \nu) \underbrace{V_a^f}_{\text{Matter potential}}$$

$$V_{S,P} \sim G_F n_f g_{S,P} \langle \frac{m_f}{E_f} \rangle$$

$$V_V \sim G_F n_f + \dots$$

$$V_{A,T} \sim G_F n_f g_{A,T} \langle \frac{\sigma_f p_f}{E_f} \rangle + \dots$$

Scalar & Pseudoscalar: Helicity suppressed

Axial & Tensor: Relevant only in polarized media

Only vector NGI are constrained by forward coherent scattering

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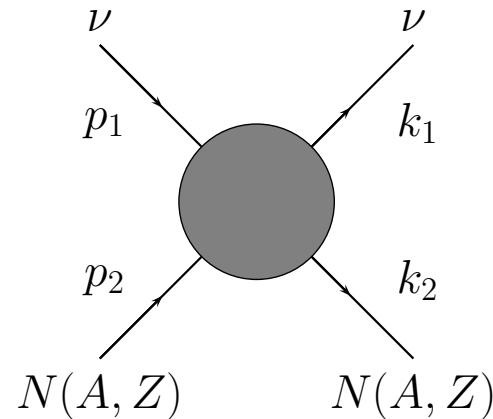
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# Parameter space scenarios



**Cross section parameterized in terms of nuclear currents: Scalar, Vector and Tensor**

Lindner, Rodejohann, Xu, 2016

DAS, De Romeri, Rojas, 2018

$$\frac{d\sigma^a(q^2=0)}{dE_r} = \frac{G_F^2}{4\pi} m_{N_a} N_a^2 \left[ \xi_S^2 \frac{E_r}{E_r^{\max}} + \xi_V^2 \left( 1 - \frac{E_r}{E_r^{\max}} - \frac{E_r}{E_\nu} \right) + \xi_T^2 \left( 1 - \frac{E_r}{2E_r^{\max}} - \frac{E_r}{E_\nu} \right) - R \frac{E_r}{E_\nu} \right]$$

## Scenarios

- Single parameter case: Only one nuclear current present at a time
- Two parameter case: Two nuclear currents are simultaneously present

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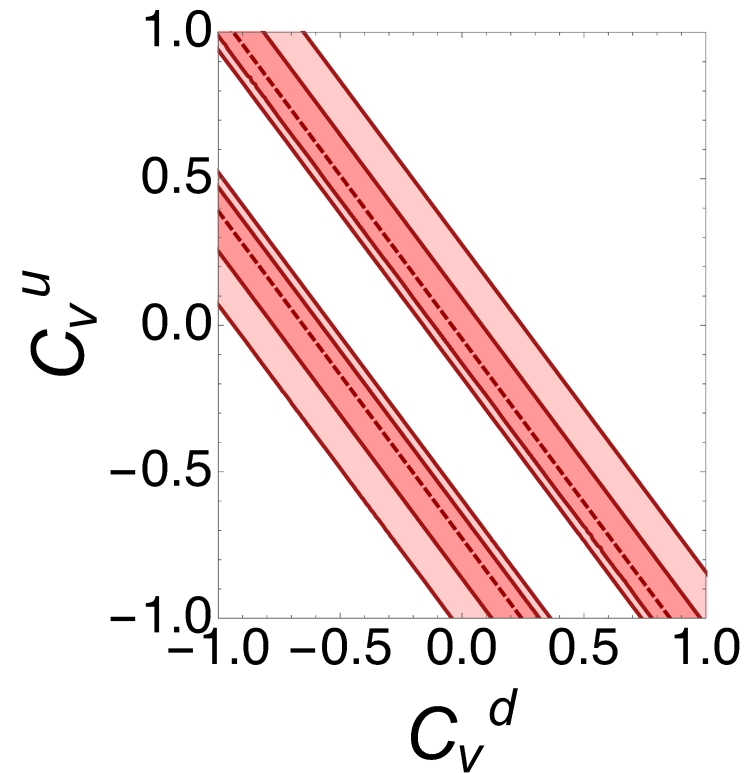
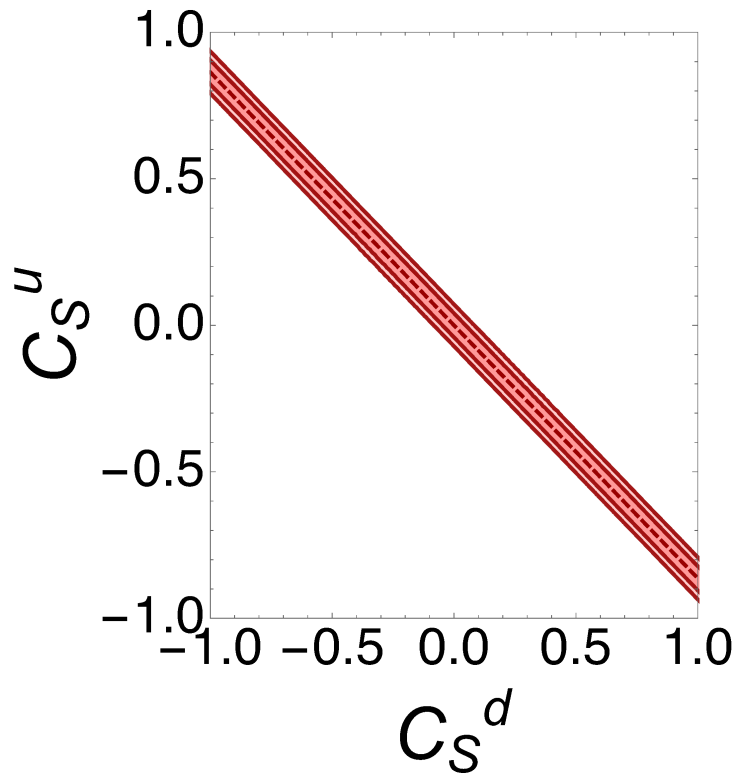
Summary

# One-parameter analysis

Param	BFP value	90% CL	99% CL
$\xi_S$	0	$[-0.62, 0.62]$	$[-1.065, 1.065]$
$\xi_V$	-0.113 -1.764	$[-0.324, 0.224]$ $[-2.102, -1.554]$	$[-0.436, 0.67]$ $[-2.545, -1.442]$
$\xi_T$	0	$[-0.591, 0.591]$	$[-1.071, 1.072]$

$$\xi_S^2 = \frac{C_S^2 + D_P^2}{N^2}$$

$$C_S = Z \sum_{q=u,d} C_S^{(q)} \frac{m_p}{m_q} f_{T_q}^p + (A - Z) \sum_{q=u,d} C_S^{(q)} \frac{m_n}{m_q} f_{T_q}^n$$



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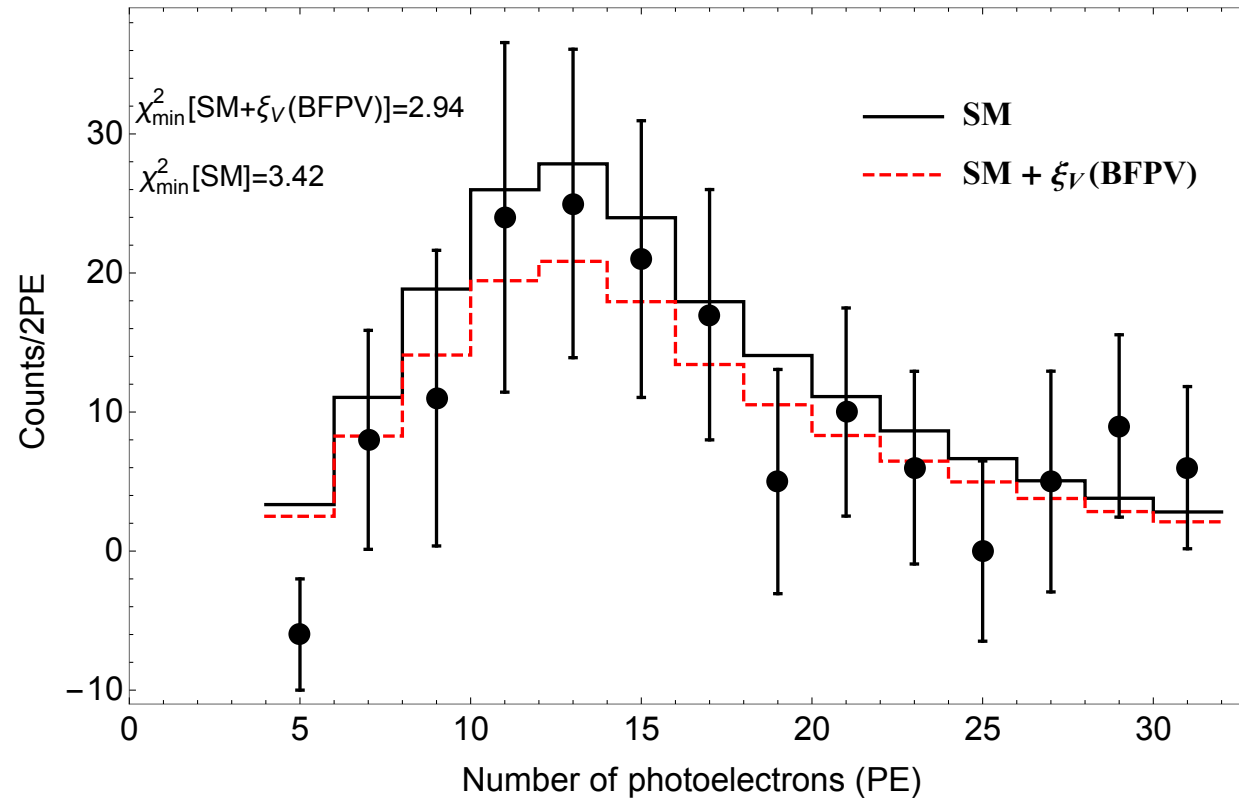
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# Vector -vs- SM+vector

The presence of NGI can indeed improve the data fit... In particular for the vector NGI



If such trend persist with further data... Is there BSM physics hidden in  $\text{CE}_\nu\text{NS}$  [??]

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Sensitivity to new physics

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What kind of physics?

- Ideal world result
- Measure low-energy spectrum

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## What kind of physics?

# Ideal world result

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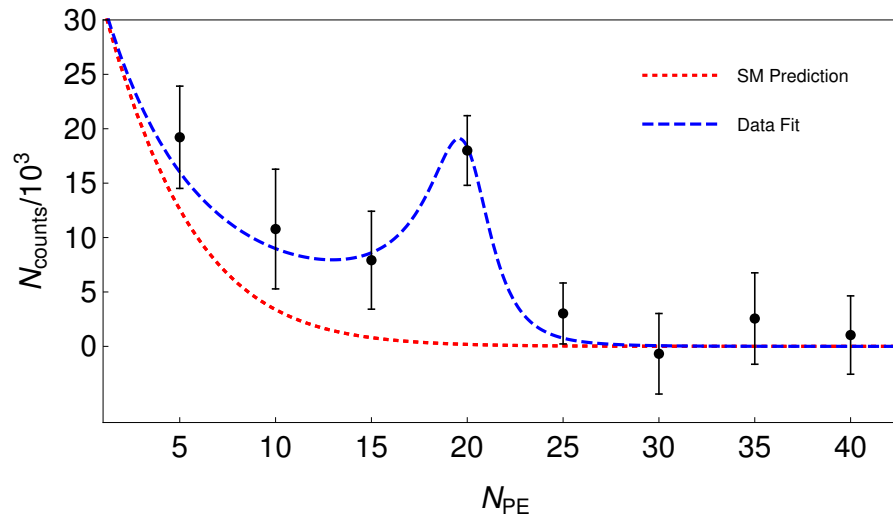
Sensitivity to new physics

What kind of physics?

● Ideal world result

● Measure low-energy  
spectrum

Summary



After the discovery...

What is the nature of the interaction?

What is the BSM physics behind?

Light Physics

Heavy Physics

What Lorentz structure?

# Measure low-energy spectrum

$$\frac{dR}{dE_r} \sim \int \Phi_\nu \left[ 1 - \left(1 - \xi_S^2\right) \frac{E_r}{E_r^{\max}} \right] dE_\nu$$

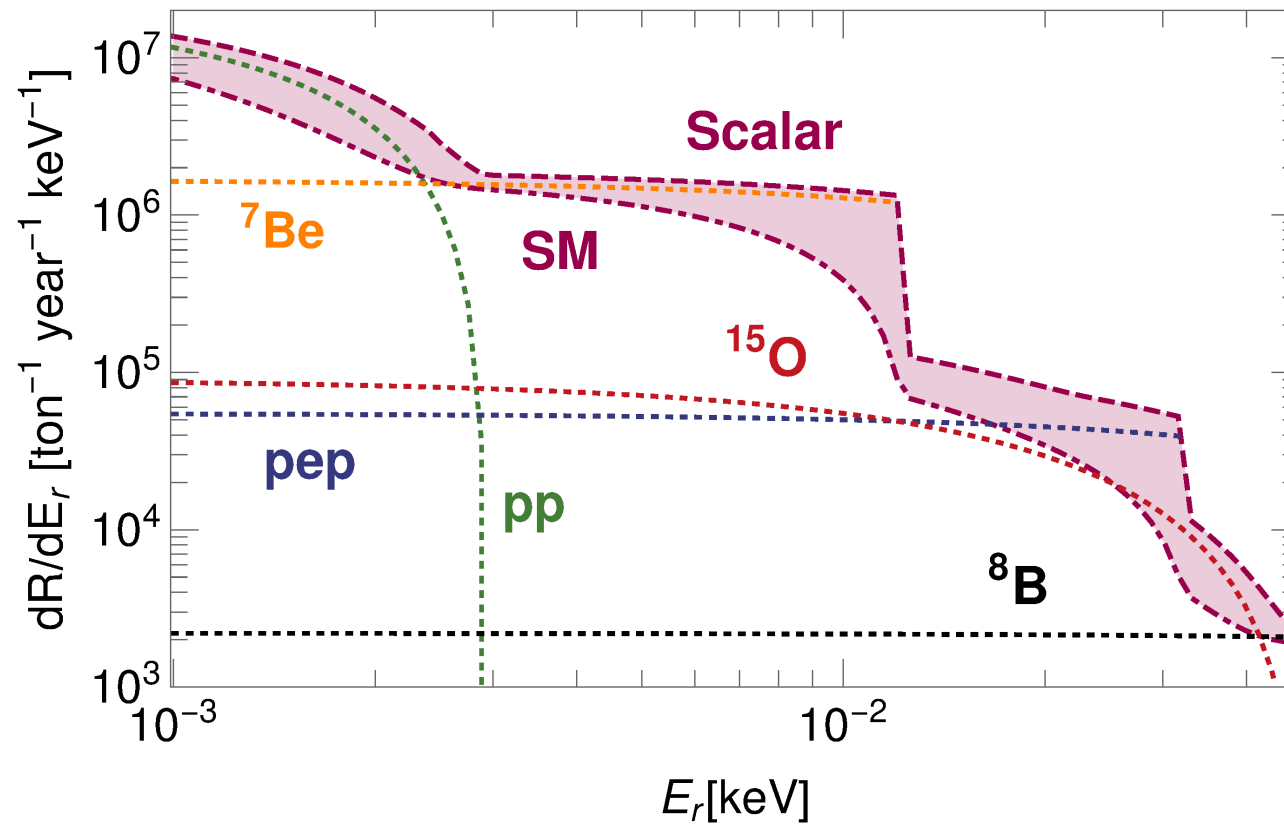
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Sensitivity to new physics

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Summary

● Résumé

# Summary

- COHERENT data and forthcoming data from CONUS and e.g.  $\nu$ -CLEUS will allow studying the presence of new physics at low  $\nu$  energies
- Good understanding of the SM contribution including the axial piece, nuclear physics form factors...
- NGI are the most general set of effective interactions. Using current data we have derived constraints: **NGI can still be fairly large**
- If new interactions are present in the neutrino sector, forthcoming data might allow their discovery