Torsional regularization and finite bare charge

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Problems of general relativity

General relativity describes gravity as curvature of spacetime.

- Singularities: points with infinite density of matter.
- Incompatible with quantum mechanics. We need quantum gravity. It may resolve the singularity problem.
- Field equations contain the conservation of orbital angular momentum, contradicting Dirac equation which gives the conservation of total angular momentum (orbital + spin) and allows spin-orbit exchange in QM.

Simplest extension of GR to include QM spin:

Einstein-Cartan theory. It also eliminates the singularity problem.

Einstein-Cartan-Sciama-Kibble gravity

• Spacetime has curvature and torsion.

$$S^k_{\ ij} = \Gamma^{\ k}_{[i\ j]}$$

- Lagrangian density is proportional to curvature scalar (as in GR).
- Cartan equations:

Torsion is proportional to **spin** density of fermions. ECSK differs significantly from GR at densities $> 10^{45}$ kg/m³; passes all tests.

$$S_{jik} - S_i g_{jk} + S_k g_{ji} = -\frac{1}{2} \kappa s_{ikj}$$

arXiv.org > gr-qc > arXiv:0911.0334

• Einstein equations:

Curvature is proportional to **energy and momentum** density.

$$G^{ik} = \kappa T^{ik} + \frac{1}{2}\kappa^2 \bigg(s^{ij}_{j} s^{kl}_{l} - s^{ij}_{l} s^{kl}_{j} - s^{ijl} s^k_{jl} + \frac{1}{2} s^{jli} s^{k}_{jl} + \frac{1}{4} g^{ik} (2 s^{l}_{m} s^{jm}_{l} - 2 s^{l}_{l} s^{jm}_{m} + s^{jlm} s_{jlm}) \bigg).$$

Universe with spin fluid

Dirac particles can be averaged macroscopically as a spin fluid.

Einstein-Cartan equations for a homogeneous and isotropic Universe become Friedmann equations for the scale factor a.

$$\frac{\dot{a}^2}{c^2} + 1 = \frac{1}{3}\kappa \left(\epsilon - \frac{1}{4}\kappa s^2\right)a^2$$
$$\frac{\dot{a}^2 + 2a\ddot{a}}{c^2} + 1 = -\kappa \left(p - \frac{1}{4}\kappa s^2\right)a^2$$

$$s^2 = \frac{1}{8}(\hbar cn)^2$$

Spin and torsion modify the energy density and pressure with a negative term proportional to the square of the fermion number density *n*, which acts like repulsive gravity and prevents singularities. The Big Bang is replaced by a non-singular Big Bounce.

Problems of quantum field theory

- Ultraviolet divergence: Feynman diagrams involve divergent integrals in the four-momentum space arising from high-energy contributions.
- This unphysical result requires regularization: a mathematical method of turning singular quantities into finite quantities.
 Most common: adding fictitious particles, changing dimensions.
- Renormalization: the original (bare) values of mass and charge absorb divergent terms, giving the measured (dressed) values.
- Dirac was critical about renormalization and expected a realistic regularization based on the principles of physics.
- Solution: torsional regularization, renormalization is finite.

Torsion and noncommutativity of momentum

- Consider two infinitesimal four-vectors dx and dx'.
- In the presence of torsion the parallel transport of dx along dx' and the parallel transport of dx' along dx do not form a closed parallelogram:

$$\delta dx^i = -\Gamma^{\ i}_{jk} dx^j dx'^k \qquad \delta dx'^i = -\Gamma^{\ i}_{jk} dx^j dx'^k \qquad \delta dx'^i - \delta dx^i = -S^i_{\ jk} dx^j dx'^k$$

 Since the momentum is a generator of translation, described by the parallel transport, its operator in quantum mechanics is given by the covariant derivative:

$$p_k = i\hbar \nabla_k$$

• In the presence of torsion, translations do not commute and therefore the four-momentum components do not commute:

$$[p_i, p_j] = 2i\hbar S^k_{\ ij} p_k$$

Integration in momentum space becomes summation over momentum eigenstates

• The classical and quantum partition functions in statistical physics: $\int dq \int dp \, f(H(q,p)) \leftrightarrow 2\pi \sum_{\text{eigenstates}} f(E) \, |[q,p]|$

 One can choose locally a frame of reference in which only the space momentum components do not commute:

$$[p_x, p_y] = iQp_z, \ [p_y, p_z] = iQp_x, \ [p_z, p_x] = iQp_y \qquad Q = -2\hbar A^0$$

$$A^i = \frac{1}{6}\epsilon^{ijkl}S_{jkl}$$

- Einstein-Cartan gravity gives: $Q = Up^3$ (*U* is const ~ M^{-2}_{Pl})
- We obtain a relation analogous to the angular momentum:

$$[n_x, n_y] = in_z, \ [n_y, n_z] = in_x, \ [n_z, n_x] = in_y$$
 $\mathbf{n} = \frac{\mathbf{p}}{Q}$

Integration in momentum space becomes summation over momentum eigenstates

We propose that the integration in n-space satisfying

$$[n_x, n_y] = in_z, \ [n_y, n_z] = in_x, \ [n_z, n_x] = in_y$$

Is replaced with the summation:

$$\int dn_x \int dn_y \int dn_z f(\mathbf{n}^2) \to 4\pi \sum_{\text{eigenstates}} f(\mathbf{n}^2) |n_z|.$$

$$\to 4\pi \sum_{l=1}^{\infty} \sum_{m=-l}^{l} f(\mathbf{n}^2) |m| = 4\pi \sum_{l=1}^{\infty} f(\mathbf{n}^2) l(l+1)$$

Torsional regularization – NP, arXiv:1712.09997

Integration in momentum space becomes summation over momentum eigenstates

Apply TR to a logarithmically divergent integral:

$$\int \frac{d^4p}{(p^2 + \mu^2)^2} = \int \frac{dp_0 d\mathbf{p}}{(p^2 + \mu^2)^2} = \int \frac{dp_0 J d\mathbf{n}}{(p^2 + \mu^2)^2} \to 4\pi \int_{-\infty}^{\infty} dp_0 \sum_{l=1}^{\infty} \frac{J}{(p^2 + \mu^2)^2} l(l+1)$$

$$J = \partial(p_x, p_y, p_z) / \partial(n_x, n_y, n_z)$$
 $p^2 = p_0^2 + U^2 n^2 p^6$

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$$\frac{\partial p}{\partial n_x} = \frac{U^2 p^5 n_x}{1 - 3U^2 n^2 p^4}$$

$$\frac{\partial p_x}{\partial n_x} = \frac{\partial (Qn_x)}{\partial n_x} = Q + 3Un_x p^2 \frac{\partial p}{\partial n_x} \qquad \frac{\partial p_x}{\partial n_y} = \frac{\partial (Qn_x)}{\partial n_y} = 3Un_x p^2 \frac{\partial p}{\partial n_y}$$

$$\frac{\partial p_x}{\partial n_y} = \frac{\partial (Qn_x)}{\partial n_y} = 3Un_x p^2 \frac{\partial p}{\partial n_y}$$

$$dp_0/dp = (1 - 3U^2n^2p^4)/(1 - U^2n^2p^4)^{1/2}$$

$$n = \sqrt{l(l+1)}$$

$$J = \det \begin{pmatrix} \partial p_x / \partial n_x & \partial p_x / \partial n_y & \partial p_x / \partial n_z \\ \partial p_y / \partial n_x & \partial p_y / \partial n_y & \partial p_y / \partial n_z \\ \partial p_z / \partial n_x & \partial p_z / \partial n_y & \partial p_z / \partial n_z \end{pmatrix} = \frac{Q^3}{1 - 3U^2 n^2 p^4}$$

Torsion eliminates ultraviolet divergence

$$\begin{split} &4\pi\int_{-\infty}^{\infty}dp_0\sum_{l=1}^{\infty}\frac{Q^3n^2}{(1-3U^2n^2p^4)(p^2+\mu^2)^2}=4\pi\int dp\frac{dp_0}{dp}\sum_{l=1}^{\infty}\frac{Q^3n^2}{(1-3U^2n^2p^4)(p^2+\mu^2)^2}\\ &=4\pi\int_{-1/\sqrt{Un}}^{1/\sqrt{Un}}dp\sum_{l=1}^{\infty}\frac{Q^3n^2}{(1-U^2n^2p^4)^{1/2}(p^2+\mu^2)^2}=8\pi\int_{0}^{1/\sqrt{Un}}dp\sum_{l=1}^{\infty}\frac{U^3p^9n^2}{(1-U^2n^2p^4)^{1/2}(p^2+\mu^2)^2}\\ &=8\pi\int_{0}^{1}d\xi\sum_{l=1}^{\infty}\frac{U^3\xi^9n^2(Un)^{-5}}{(1-\xi^4)^{1/2}[\xi^2/(Un)+\mu^2]^2}=8\pi\int_{0}^{1}d\xi\sum_{l=1}^{\infty}\frac{\xi^9n^{-1}}{(1-\xi^4)^{1/2}[\xi^2+U\mu^2n]^2}\\ &=4\pi\int_{0}^{1}d\zeta\sum_{l=1}^{\infty}\frac{\zeta^4n^{-1}}{(1-\zeta^2)^{1/2}[\zeta+U\mu^2n]^2}=4\pi\sum_{l=1}^{\infty}\int_{0}^{\pi/2}d\phi\frac{\sin^4\phi\,n^{-1}}{[\sin\phi+U\mu^2n]^2}\\ &=4\pi\sum_{l=1}^{\infty}\int_{0}^{\pi/2}d\phi\frac{\sin^4\phi\,[l(l+1)]^{-1/2}}{[\sin\phi+U\mu^2\sqrt{l(l+1)}]^2}, \qquad \qquad Unp^2=\xi^2=\zeta=\sin\phi \end{split}$$

The logarithmically divergent integral is replaced with a sum that converges as I^{-3} .

Torsion eliminates ultraviolet divergence

This procedure can be generalized to tensor integrals:

$$\begin{split} &\int \frac{d^4p}{(p^2 + \mu^2)^s} \to 8\pi \int_0^{1/\sqrt{Un}} dp \sum_{l=1}^\infty \frac{U^3 p^9 n^2}{(1 - U^2 n^2 p^4)^{1/2} (p^2 + \mu^2)^s} \\ &= 8\pi \int_0^1 d\xi \sum_{l=1}^\infty \frac{U^3 \xi^9 n^2 (Un)^{-5}}{(1 - \xi^4)^{1/2} [\xi^2/(Un) + \mu^2]^s} = 8\pi \int_0^1 d\xi \sum_{l=1}^\infty \frac{U^{s-2} \xi^9 n^{s-3}}{(1 - \xi^4)^{1/2} [\xi^2 + U\mu^2 n]^s} \\ &= 4\pi \int_0^1 d\zeta \sum_{l=1}^\infty \frac{U^{s-2} \zeta^4 n^{s-3}}{(1 - \zeta^2)^{1/2} [\zeta + U\mu^2 n]^s} = 4\pi U^{s-2} \sum_{l=1}^\infty \int_0^{\pi/2} d\phi \frac{\sin^4 \phi \, n^{s-3}}{[\sin \phi + U\mu^2 n]^s} \\ &= 4\pi U^{s-2} \sum_{l=1}^\infty \int_0^{\pi/2} d\phi \frac{\sin^4 \phi \, [l(l+1)]^{(s-3)/2}}{[\sin \phi + U\mu^2 \sqrt{l(l+1)}]^s}. \end{split}$$

$$\int d^4p \frac{\partial}{\partial p_{\nu}} \left(\frac{p^{\mu}}{(p^2 + \Delta)^s} \right) = \int d^4p \frac{\delta^{\mu\nu}}{(p^2 + \Delta)^s} - 2s \int d^4p \frac{p^{\mu}p^{\nu}}{(p^2 + \Delta)^{s+1}}$$

$$\int d^4p \frac{p^{\mu}p^{\nu}}{(p^2 + \Delta)^s} = \frac{\delta^{\mu\nu}}{2(s-1)} \int d^4p \frac{1}{(p^2 + \Delta)^{s-1}}$$

Vacuum polarization

The vacuum polarization tensor is gauge invariant:

$$\begin{split} &\Pi_{\text{bubble}}^{\mu\nu}(q) = -\frac{\alpha_0}{\pi^3} \int d^4p_{\text{E}} \int_0^1 dx \frac{-2p_{\text{E}}^{\mu}p_{\text{E}}^{\nu} + p_{\text{E}}^2\delta^{\mu\nu} + \Delta\delta^{\mu\nu} + 2(q^2g^{\mu\nu} - q^{\mu}q^{\nu})x(1-x)}{(p_{\text{E}}^2 + \Delta)^2} \\ &= -\frac{2\alpha_0}{\pi^3} \int d^4p_{\text{E}} \int_0^1 dx \frac{x(1-x)}{(p_{\text{E}}^2 + \Delta)^2} (q^2g^{\mu\nu} - q^{\mu}q^{\nu}) = \Pi(q^2)q^2 \Big(g^{\mu\nu} - \frac{q^{\mu}q^{\nu}}{q^2}\Big), \end{split}$$

$$\Pi(q^2) = -\frac{2\alpha_0}{\pi^3} \int d^4 p_{\rm E} \int_0^1 dx \frac{x(1-x)}{(p_{\rm E}^2 + \Delta)^2}$$

$$\Delta = m^2 - q^2 x (1 - x)$$

$$\Pi(q^2) \to -\frac{8\alpha_0}{\pi^2} \sum_{l=1}^{\infty} \int_0^1 dx \int_0^{\pi/2} d\phi \frac{\sin^4 \phi \, n^{-1} x (1-x)}{[\sin \phi + U \Delta n]^2}$$
$$= -\frac{8\alpha_0}{\pi^2} \sum_{l=1}^{\infty} \int_0^1 dx \int_0^{\pi/2} d\phi \frac{\sin^4 \phi \, [l(l+1)]^{-1/2} x (1-x)}{[\sin \phi + U \Delta \sqrt{l(l+1)}]^2}$$

The sum-integral in Π is finite.

Torsion makes bare charge finite

Renormalization of the electric charge:

$$\alpha = \frac{\alpha_0}{1 - \Pi(0)}$$

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$$\alpha_{\text{run}} = \frac{\alpha_0}{1 - \Pi(q^2)}$$

Gives the bare electric charge of an electron:

$$e_0 = \frac{e}{(1 + \Pi_R(0))^{1/2}} = e \left[1 - \frac{8\alpha}{\pi^2} \sum_{l=1}^{\infty} \int_0^1 dx \int_0^{\pi/2} d\phi \frac{\sin^4 \phi \, n^{-1} x (1-x)}{[\sin \phi + U m^2 n]^2} \right]^{-1/2}.$$

Including all charged fermions in Π gives the bare charge -1.22 e. The running coupling constant is finite.

Accordingly, the bare fine structure constant is about 1/92.1.

NP, arXiv:1712.09997

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Summary

- The conservation law for total angular momentum (orbital + spin) in curved spacetime, consistent with Dirac equation, requires torsion.
- In the presence of torsion, the four-momentum operator components do not commute. The integration in the momentum space must be replaced with the summation over the momentum eigenvalues.
- The separation between the momentum eigenvalues increases with the magnitude of the momentum as a result of the Einstein-Cartan gravity.
 Consequently, ultraviolet divergent integrals turn into convergent sums.
- Torsion naturally regularizes ultraviolet divergence in QED. Renormalization in QED is finite, leading to a finite bare charge of an electron: -1.22 e.
- Future work: research how torsion affects the electroweak and strong interactions.