

Acknowledgements: The E158, MOLLER, EXO-200 and nEXO collaborations
and many other theoretical and experimental colleagues

Low Energy Probes of the Standard Model (I)



Parity-Violating Electron Scattering and Neutrinoless Double Beta Decay Experiments

Krishna Kumar

Stony Brook University

The Standard Model at 50: Success and Challenges

SLAC Summer Institute

July 31, 2018

Outline

$$Q^2 \ll M_Z^2$$

◆ Low Energy Weak Neutral Current Measurements

- ★ Historical Perspective on parity-violating electron scattering
- ★ Motivation for Modern Low Q^2 Measurements
- ★ Current Experimental Status

◆ Neutrinoless Double Beta Decay

- ★ Physics Motivation Quick repeat of key points from V. Cirigliano lectures
- ★ Current Experimental Status

◆ Conclusion and Outlook

*Parity Violating
Electron Scattering:
Historical
Perspective*

The Weak Neutral Current

$\tan \theta_W = \frac{g'}{g} \quad e = g \sin \theta_W$ *One free parameter: weak mixing angle θ_W*

The Z boson incorporated

	Left-	Right-
W “Charge”	$T = \pm \frac{1}{2}$	zero
Z “Charge”	$T - q \sin^2 \theta_W$	$-q \sin^2 \theta_W$

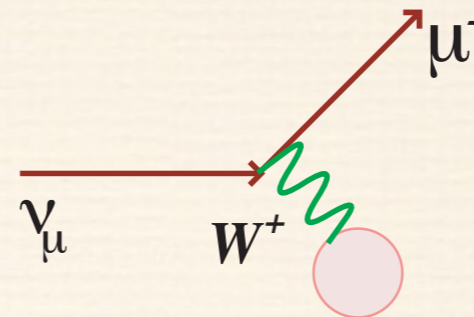
The Weak Neutral Current

$$\tan \theta_W = \frac{g'}{g} \quad e = g \sin \theta_W$$

One free parameter: weak mixing angle θ_W

The Z boson incorporated

	Left- g_L	Right- g_R
W "Charge"	$T = \pm \frac{1}{2}$	zero
Z "Charge"	$T - q \sin^2 \theta_W$	$-q \sin^2 \theta_W$



Charged Current

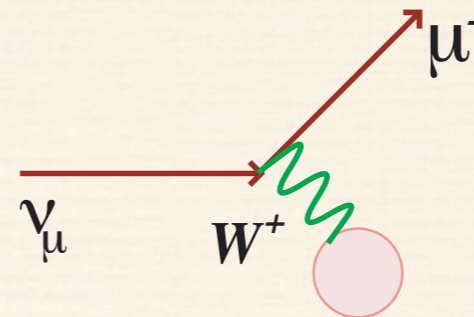
The Weak Neutral Current

$$\tan \theta_W = \frac{g'}{g} \quad e = g \sin \theta_W$$

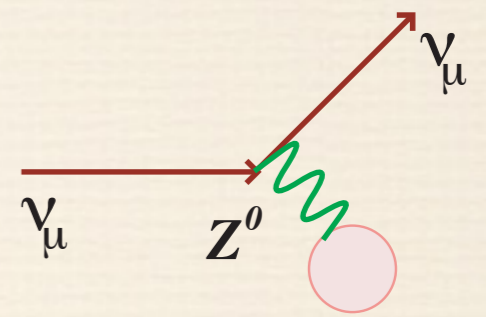
One free parameter: weak mixing angle θ_W

The Z boson incorporated

	Left- g_L	Right- g_R
W “Charge”	$T = \pm \frac{1}{2}$	zero
Z “Charge”	$T - q \sin^2 \theta_W$	$-q \sin^2 \theta_W$



Charged Current



Neutral Current

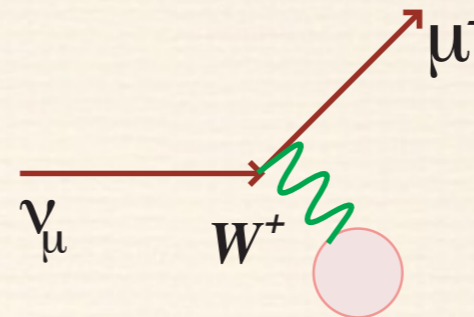
The Weak Neutral Current

$$\tan \theta_W = \frac{g'}{g} \quad e = g \sin \theta_W$$

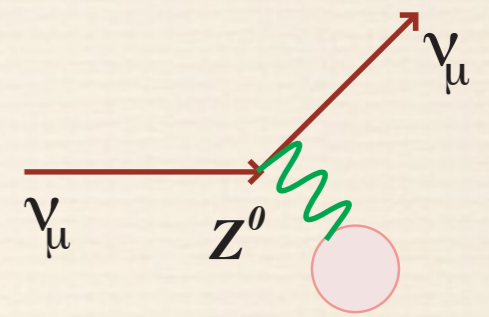
One free parameter: weak mixing angle θ_W

The Z boson incorporated

	Left- g_L	Right- g_R
W "Charge"	$T = \pm \frac{1}{2}$	zero
Z "Charge"	$T - q \sin^2 \theta_W$	$-q \sin^2 \theta_W$



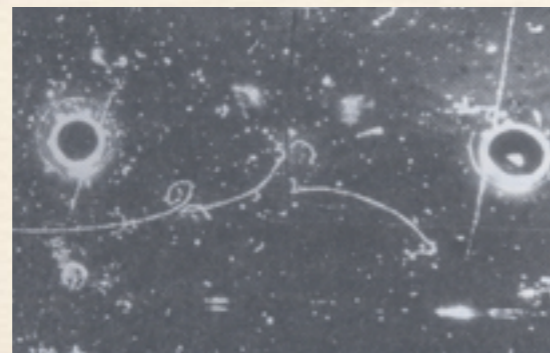
Charged Current



Neutral Current

- Gargamelle observes one $\nu_\mu e^-$ event
- First measurement of weak mixing angle

*early 1970s:
antineutrino-
electron
scattering*



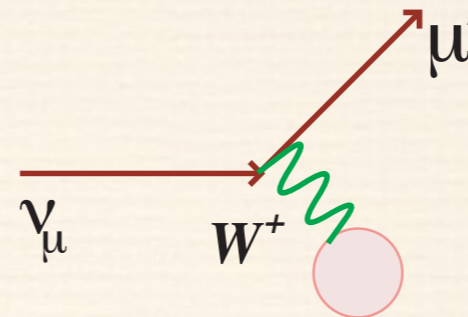
The Weak Neutral Current

$$\tan \theta_W = \frac{g'}{g} \quad e = g \sin \theta_W$$

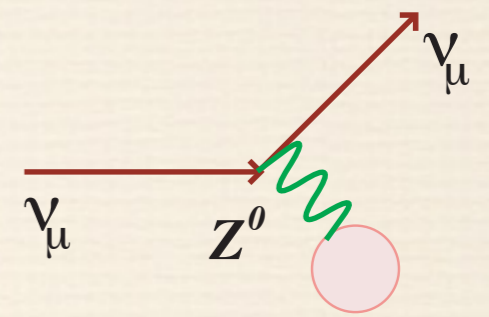
One free parameter: weak mixing angle θ_W

The Z boson incorporated

	Left- g_L	Right- g_R
W "Charge"	$T = \pm \frac{1}{2}$	zero
Z "Charge"	$T - q \sin^2 \theta_W$	$-q \sin^2 \theta_W$



Charged Current



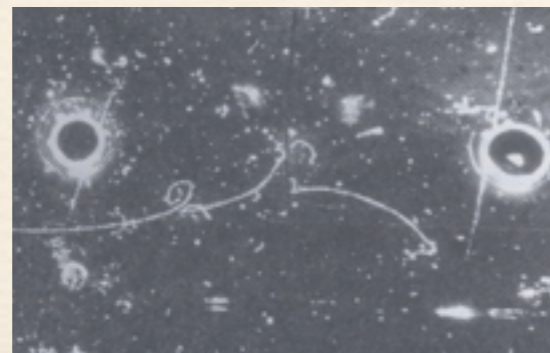
Neutral Current

- **Does the Weak Neutral Current interfere with the Electromagnetic Current?**
 - Central to establishing $SU(2)_L \times U(1)_Y$

mid-1970s

- **Gargamelle observes one $\nu_\mu e^-$ event**
- **First measurement of weak mixing angle**

*early 1970s:
antineutrino-
electron
scattering*



Glashow, Weinberg and Salam: $SU(2)_L \times U(1)_Y$

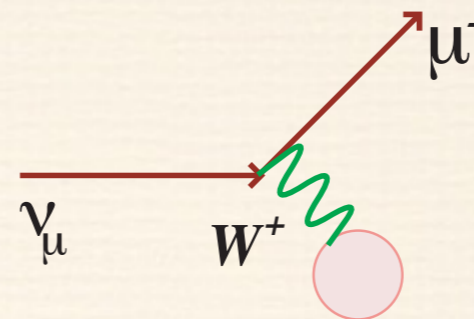
The Weak Neutral Current

$$\tan \theta_W = \frac{g'}{g} \quad e = g \sin \theta_W$$

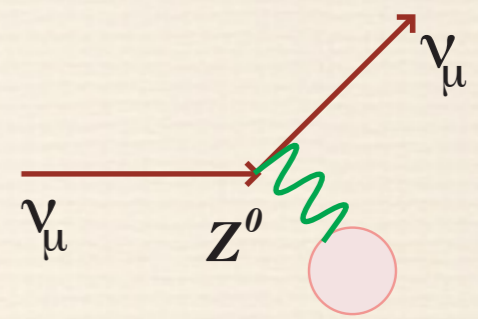
One free parameter: weak mixing angle θ_W

The Z boson incorporated

	Left- g_L	Right- g_R
W “Charge”	$T = \pm \frac{1}{2}$	zero
Z “Charge”	$T - q \sin^2 \theta_W$	$-q \sin^2 \theta_W$



Charged Current



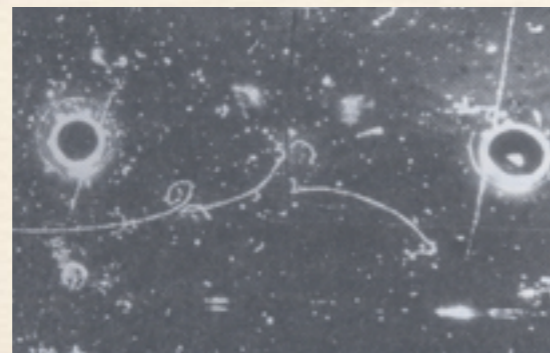
Neutral Current

- **Does the Weak Neutral Current interfere with the Electromagnetic Current?**
 - Central to establishing $SU(2)_L \times U(1)_Y$

mid-1970s

- **Gargamelle observes one $\nu_\mu e^-$ event**
- **First measurement of weak mixing angle**

*early 1970s:
antineutrino-
electron
scattering*



Consider fixed target electron scattering

Glashow, Weinberg and Salam: $SU(2)_L \times U(1)_Y$

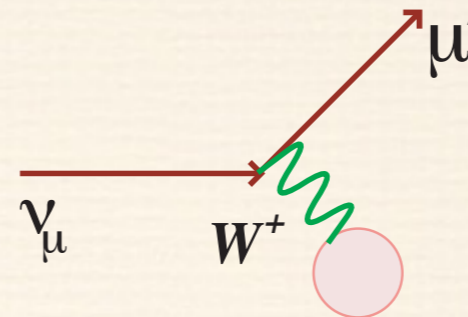
The Weak Neutral Current

$$\tan \theta_W = \frac{g'}{g} \quad e = g \sin \theta_W$$

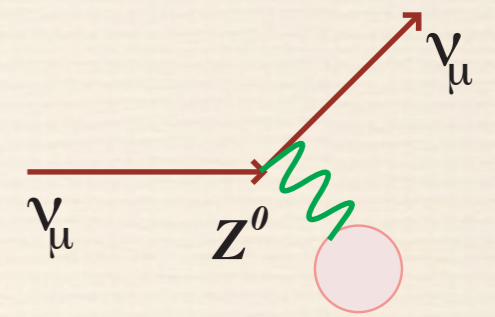
One free parameter: weak mixing angle θ_W

The Z boson incorporated

	Left- g_L	Right- g_R
W “Charge”	$T = \pm \frac{1}{2}$	zero
Z “Charge”	$T - q \sin^2 \theta_W$	$-q \sin^2 \theta_W$



Charged Current



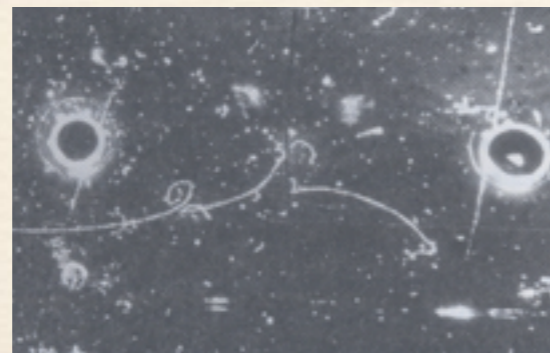
Neutral Current

- **Does the Weak Neutral Current interfere with the Electromagnetic Current?**
 - Central to establishing $SU(2)_L \times U(1)_Y$

mid-1970s

- **Gargamelle observes one $\nu_\mu e^-$ event**
- **First measurement of weak mixing angle**

*early 1970s:
antineutrino-
electron
scattering*



$$\begin{pmatrix} \nu \\ e \end{pmatrix}_l \quad \begin{pmatrix} E^0 \\ e \end{pmatrix}_r$$

Parity is conserved

Glashow, Weinberg and Salam: $SU(2)_L \times U(1)_Y$

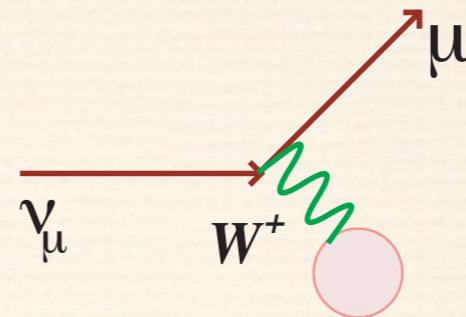
The Weak Neutral Current

$$\tan \theta_W = \frac{g'}{g} \quad e = g \sin \theta_W$$

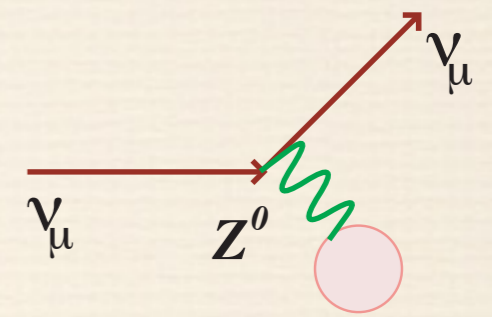
One free parameter: weak mixing angle θ_W

The Z boson incorporated

	Left- g_L	Right- g_R
W "Charge"	$T = \pm \frac{1}{2}$	zero
Z "Charge"	$T - q \sin^2 \theta_W$	$-q \sin^2 \theta_W$



Charged Current

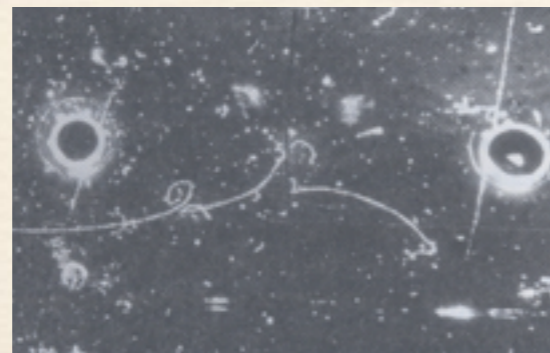


Neutral Current

- **Does the Weak Neutral Current interfere with the Electromagnetic Current?**
 - Central to establishing $SU(2)_L \times U(1)_Y$ mid-1970s

- **Gargamelle observes one $\nu_\mu e^-$ event**
- **First measurement of weak mixing angle**

*early 1970s:
antineutrino-
electron
scattering*



Consider fixed target electron scattering

$$\begin{pmatrix} \nu \\ e \end{pmatrix}_l \quad \begin{pmatrix} E^0 \\ e \end{pmatrix}_r$$

Parity is conserved

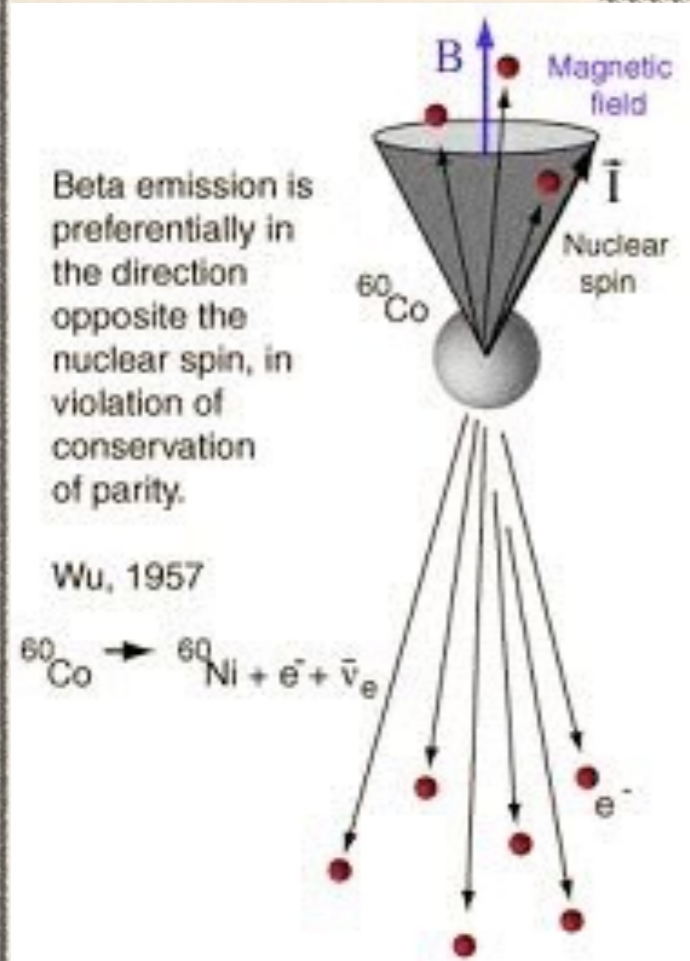
$$\begin{pmatrix} \nu \\ e \end{pmatrix}_l \quad (e)_r$$

Parity is violated

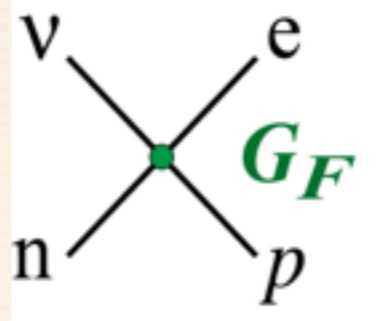
Zel'dovich speculation: Is Electron Scattering Parity-Violating?

Electroweak Scattering

JETP 36, pp 964-66 (1959)



Nuclear β Decay

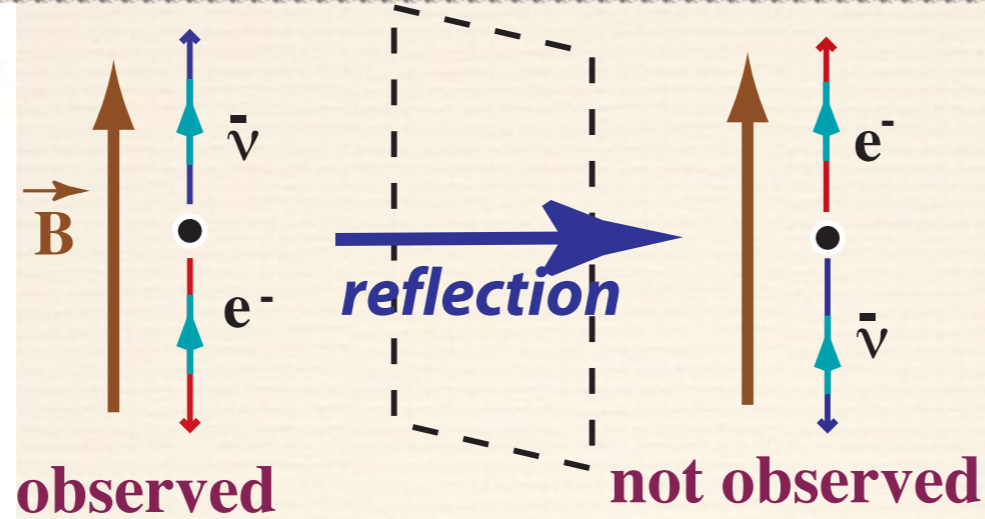
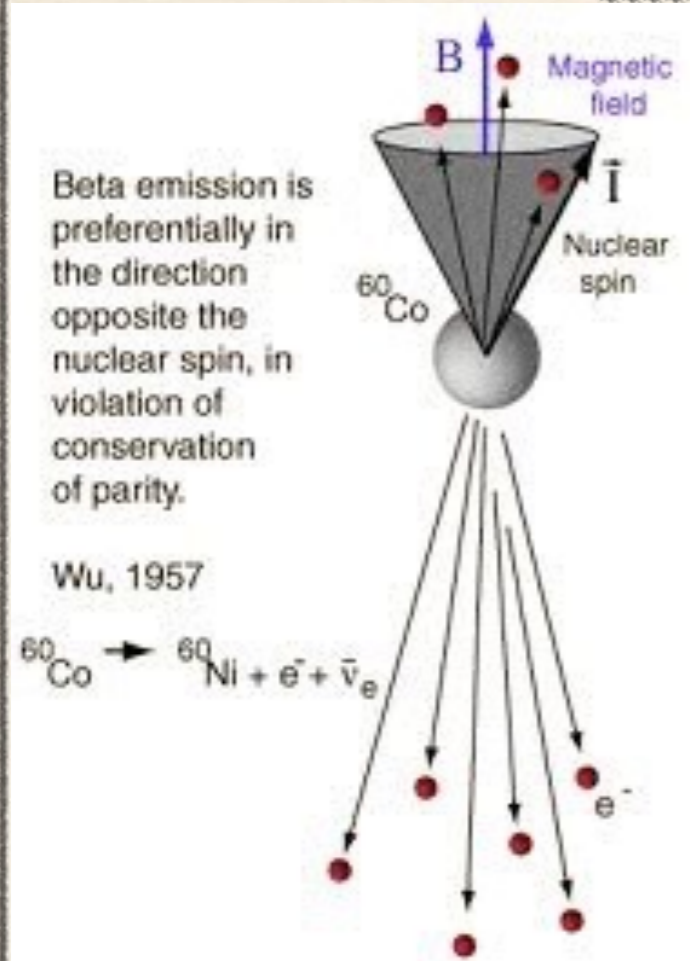


charge and flavor-changing

Zel'dovich speculation: Is Electron Scattering Parity-Violating?

Electroweak Scattering

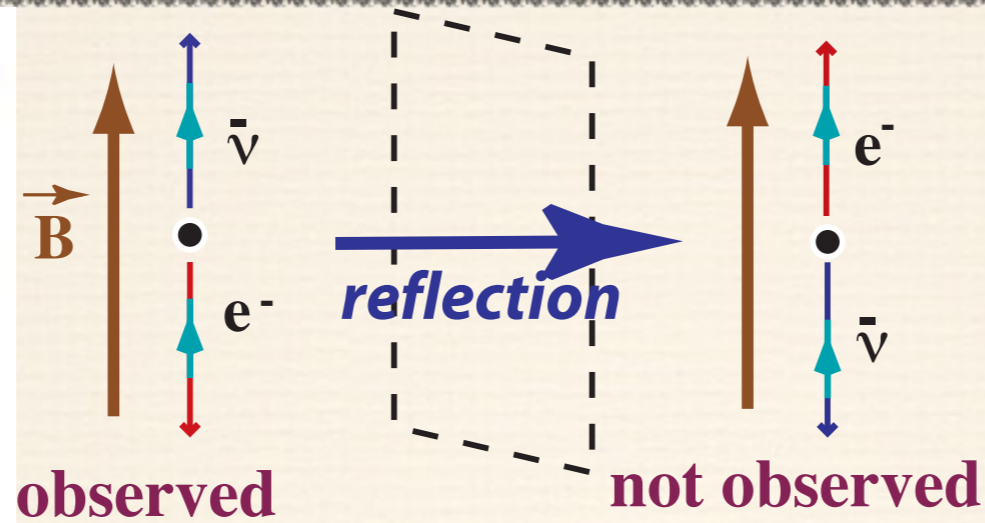
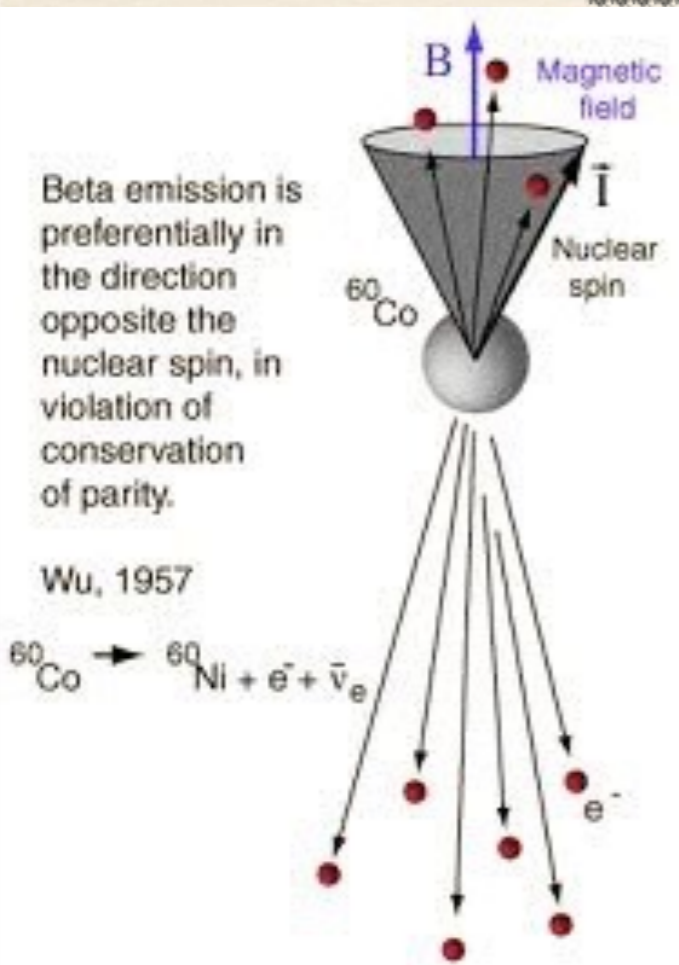
JETP 36, pp 964-66 (1959)



Zel'dovich speculation: Is Electron Scattering Parity-Violating?

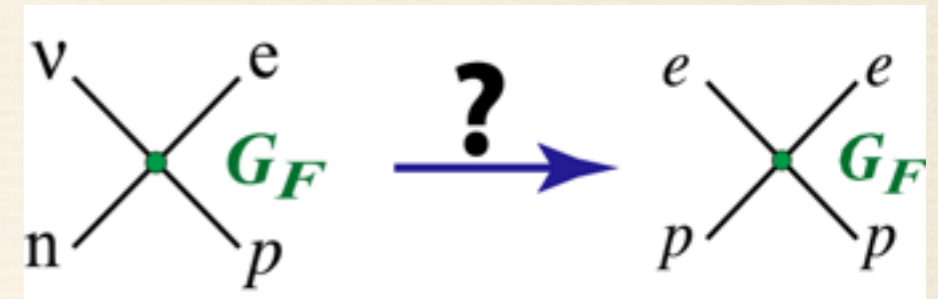
Electroweak Scattering

JETP 36, pp 964-66 (1959)



Neutron β Decay

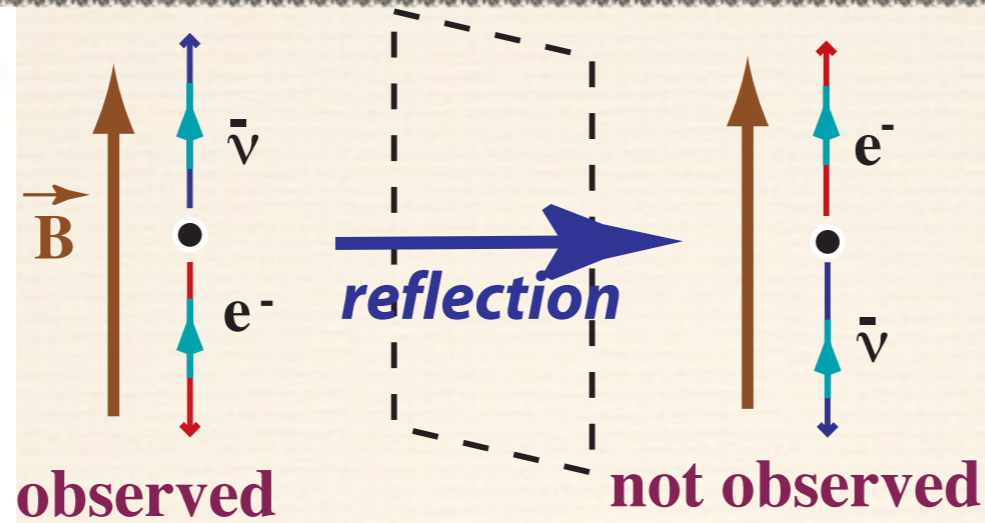
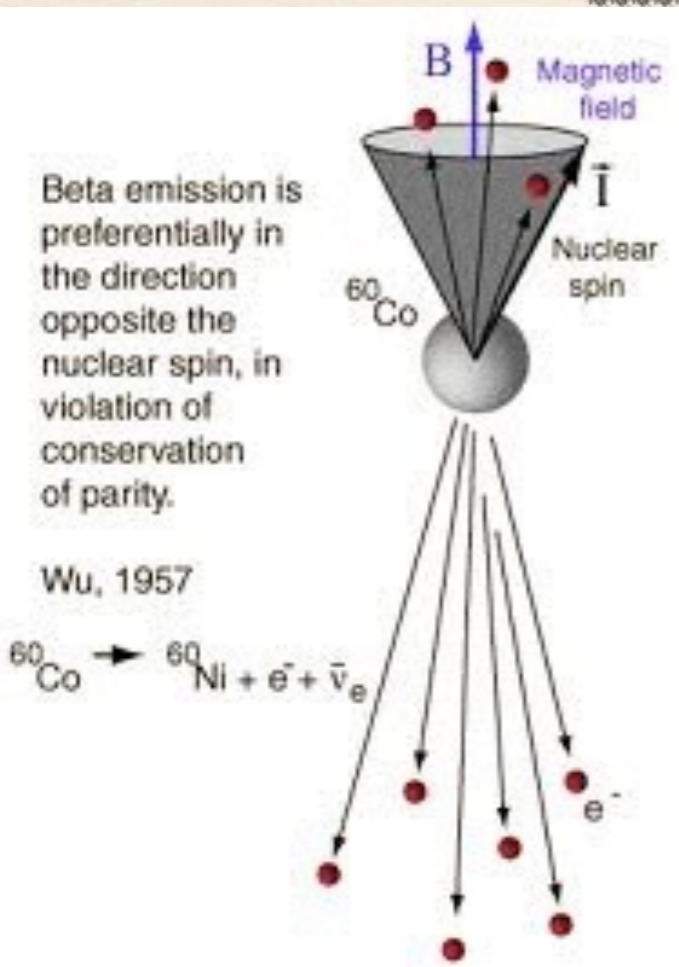
Electron-proton
Weak Scattering



Zel'dovich speculation: Is Electron Scattering Parity-Violating?

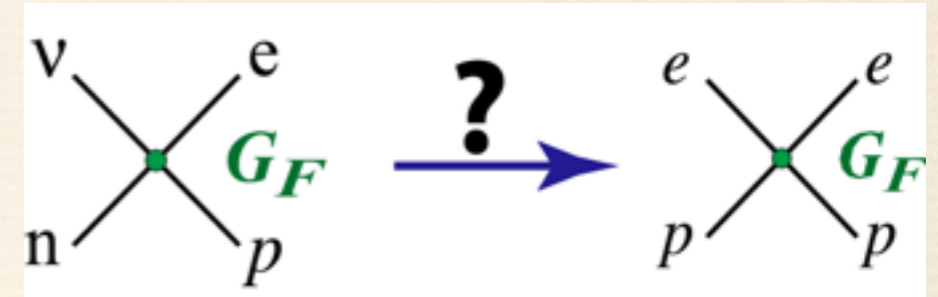
Electroweak Scattering

JETP 36, pp 964-66 (1959)



Neutron β Decay

Electron-proton
Weak Scattering



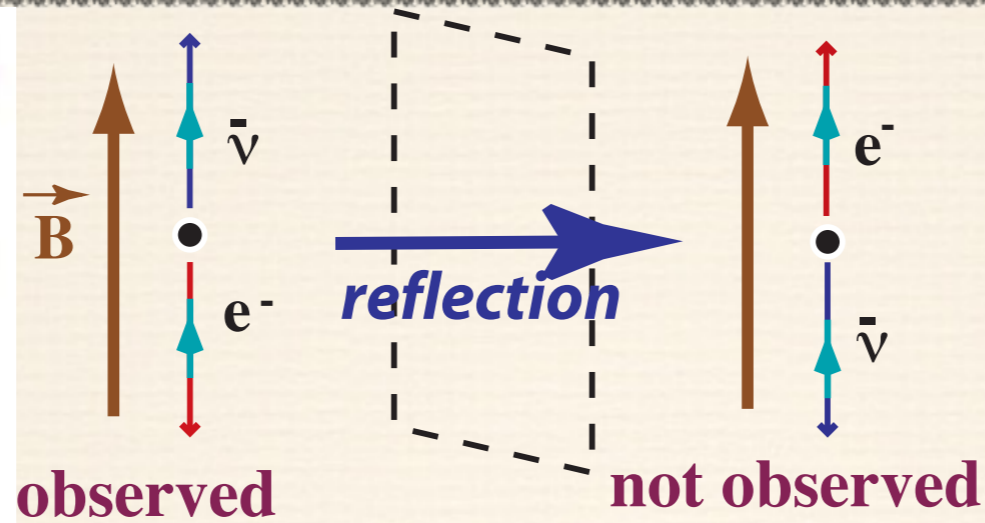
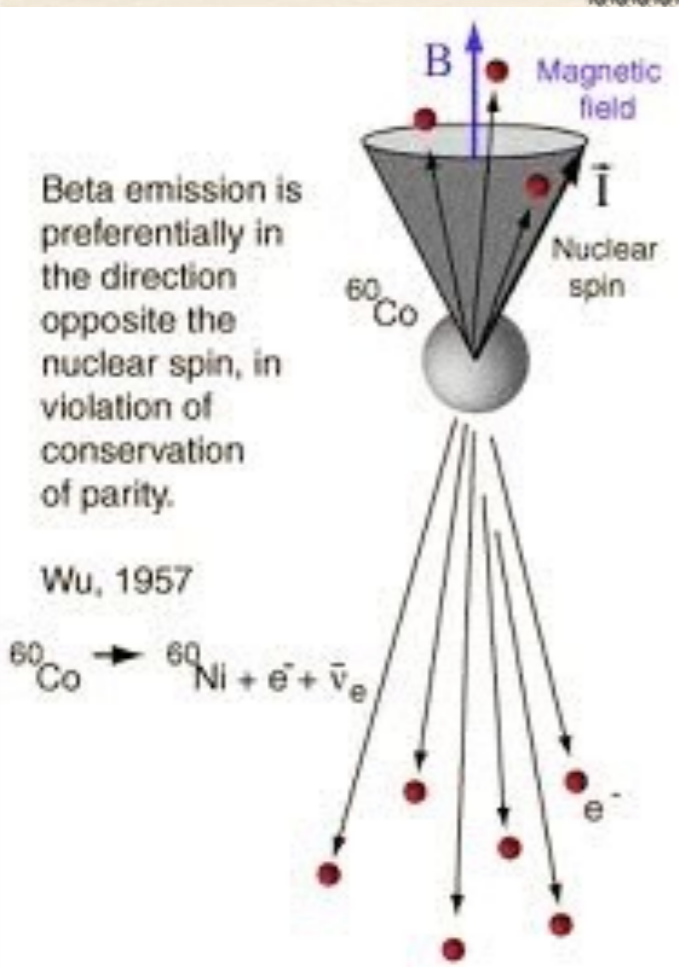
$$\sigma \propto |A_{EM} + A_{weak}|^2$$

$$\sim |A_{EM}|^2 + 2A_{EM}A_{weak}^* + \dots$$

Zel'dovich speculation: Is Electron Scattering Parity-Violating?

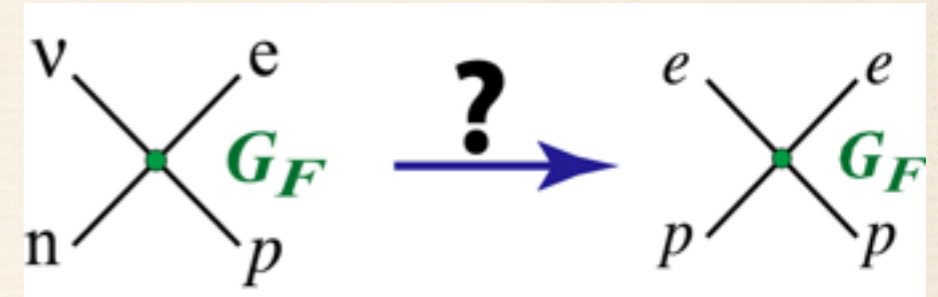
Electroweak Scattering

JETP 36, pp 964-66 (1959)



Neutron β Decay

Electron-proton
Weak Scattering



$$\sigma \propto |A_{EM} + A_{weak}|^2$$

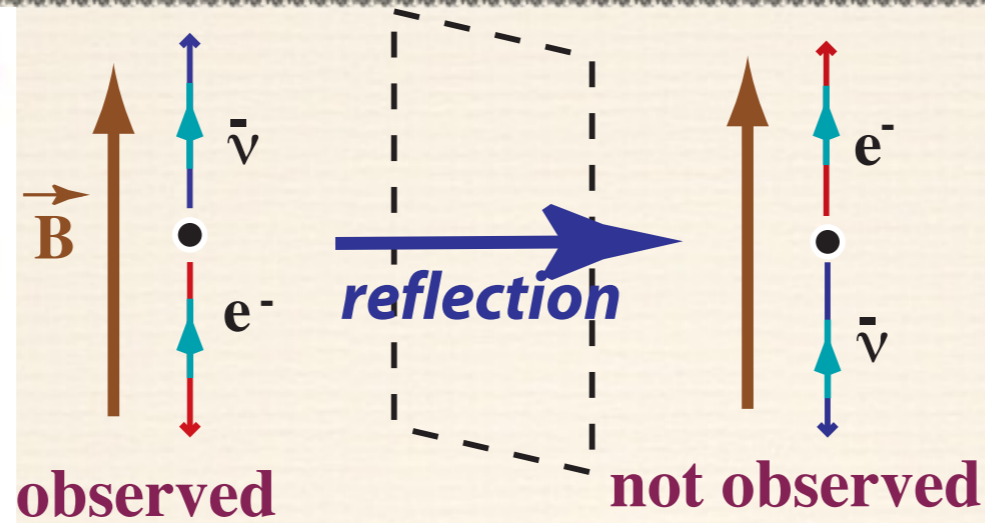
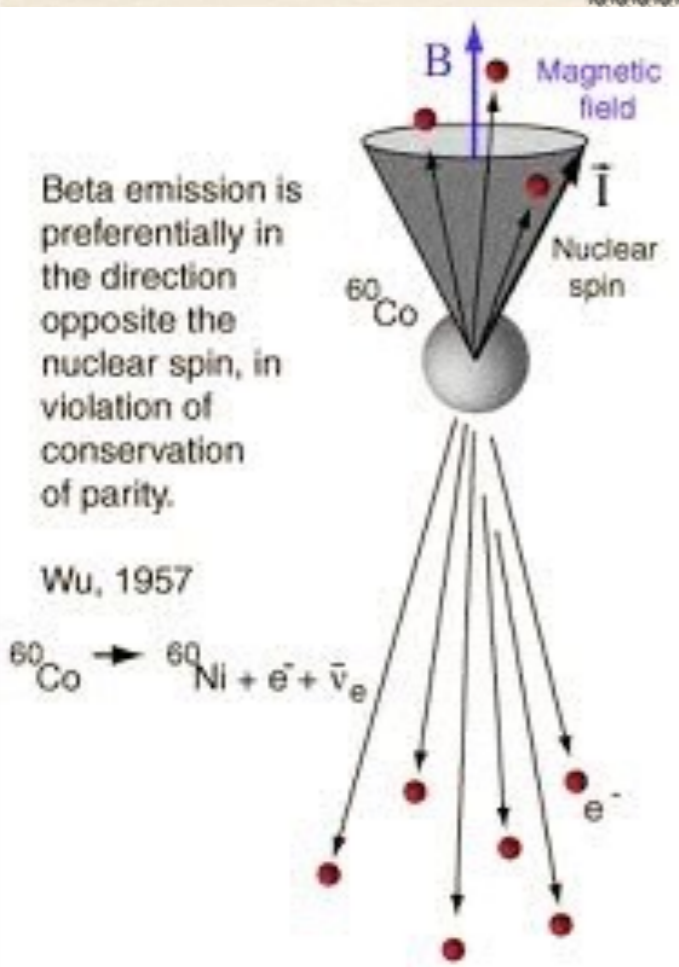
$$\sim |A_{EM}|^2 + 2A_{EM}A_{weak}^* + \dots$$

Parity-violating

Zel'dovich speculation: Is Electron Scattering Parity-Violating?

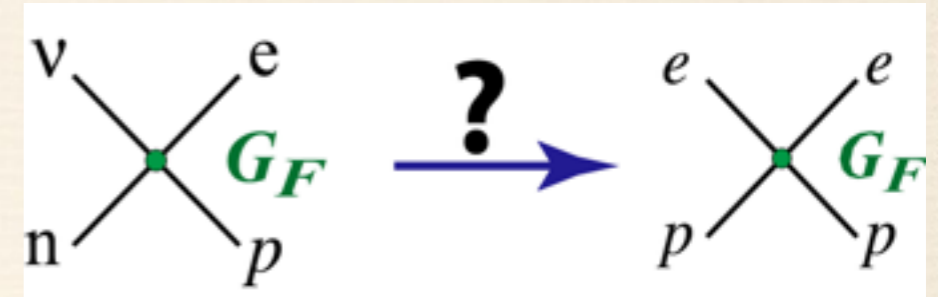
Electroweak Scattering

JETP 36, pp 964-66 (1959)



Neutron β Decay

Electron-proton
Weak Scattering



$$A_{PV} = \frac{\sigma_{\uparrow} - \sigma_{\downarrow}}{\sigma_{\uparrow} + \sigma_{\downarrow}} = -A$$

$$\sigma \propto |A_{EM} + A_{weak}|^2$$

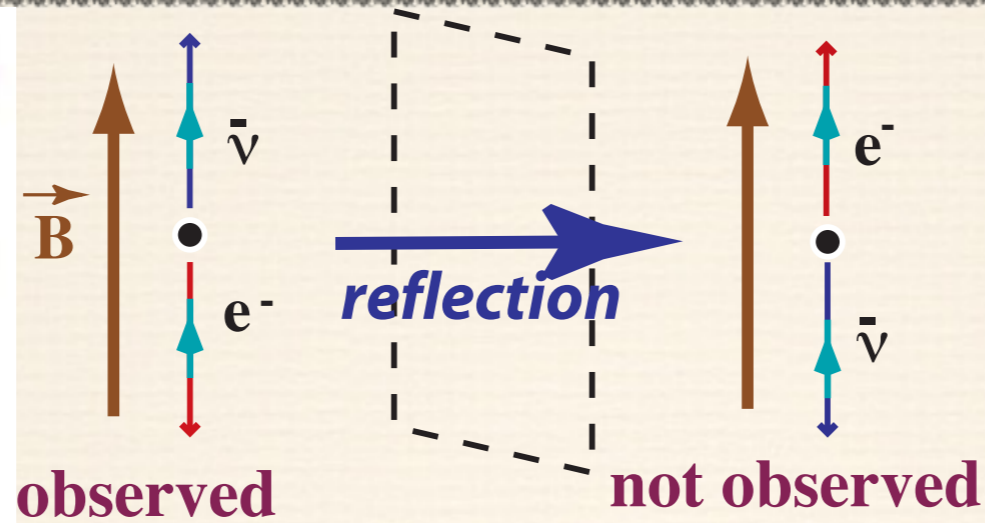
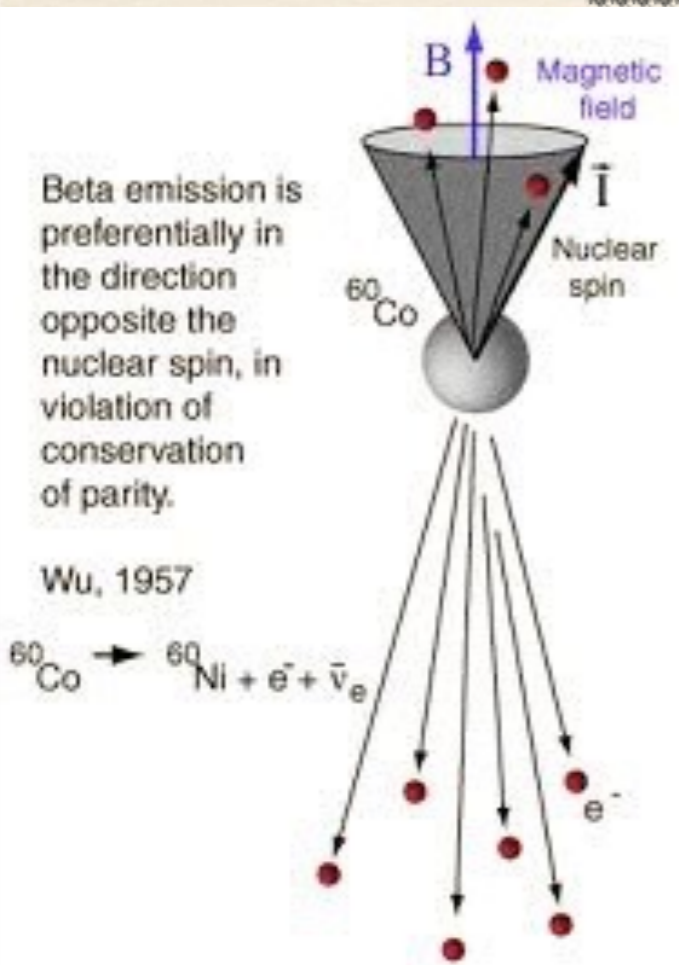
$$\sim |A_{EM}|^2 + \boxed{2A_{EM}A_{weak}^*} + \dots$$

Parity-violating

Zel'dovich speculation: Is Electron Scattering Parity-Violating?

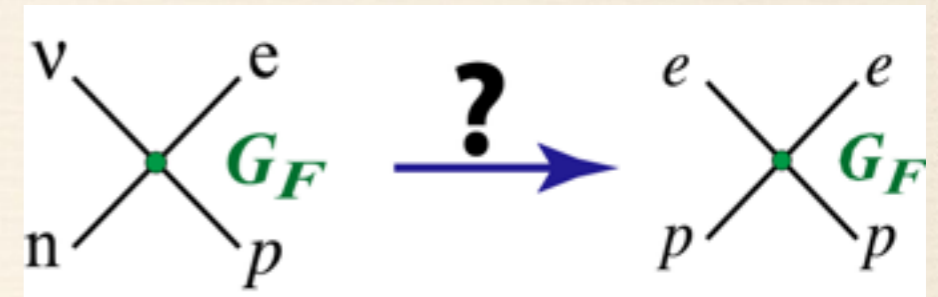
Electroweak Scattering

JETP 36, pp 964-66 (1959)



Neutron β Decay

Electron-proton
Weak Scattering



$$A_{PV} = \frac{\sigma_{\uparrow} - \sigma_{\downarrow}}{\sigma_{\uparrow} + \sigma_{\downarrow}} = -A$$

$$\sim \frac{A_{\text{weak}}}{A_{\text{EM}}} \sim \frac{G_F Q^2}{4\pi\alpha}$$

$$\sigma \propto |A_{\text{EM}} + A_{\text{weak}}|^2$$

$$\sim |A_{\text{EM}}|^2 + \boxed{2A_{\text{EM}}A_{\text{weak}}^* + \dots}$$

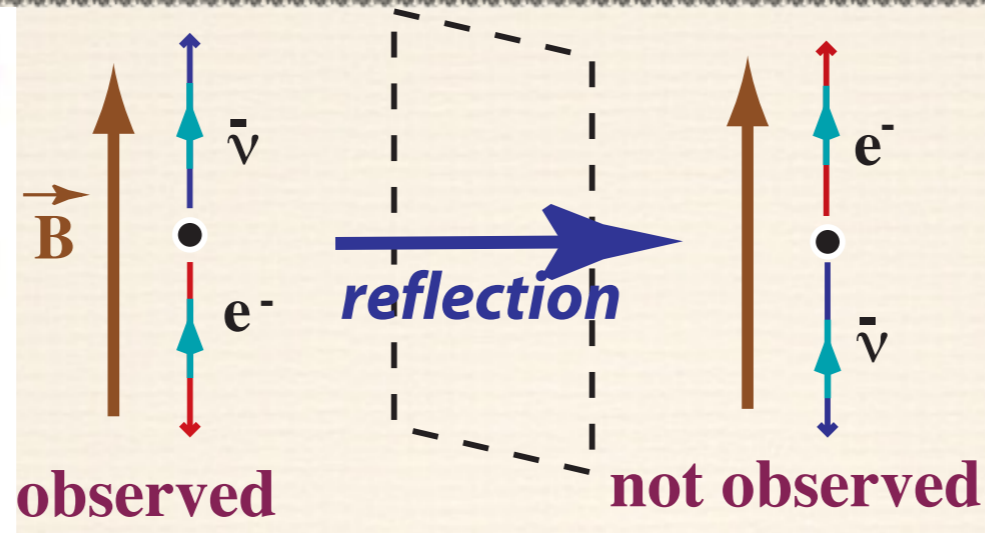
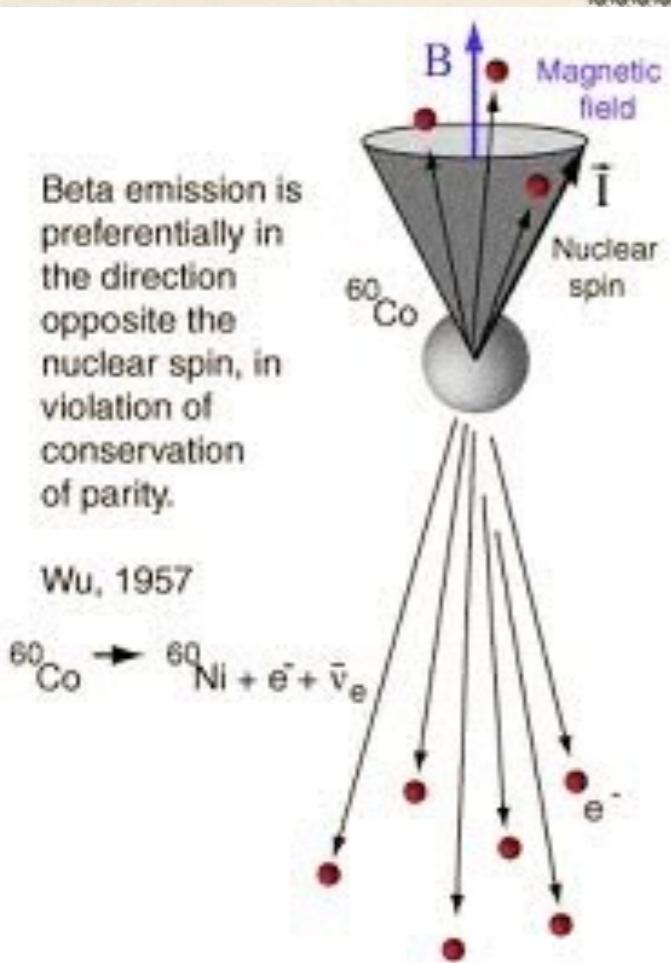
Parity-violating

$$A_{PV} \sim 10^{-4} \cdot Q^2 \text{ (GeV}^2\text{)}$$

Zel'dovich speculation: Is Electron Scattering Parity-Violating?

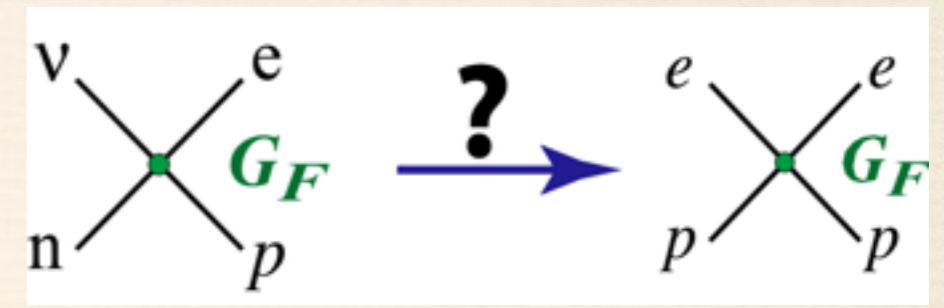
Electroweak Scattering

JETP 36, pp 964-66 (1959)



Neutron β Decay

Electron-proton Weak Scattering



$$A_{PV} = \frac{\sigma_{\uparrow} - \sigma_{\downarrow}}{\sigma_{\uparrow} + \sigma_{\downarrow}} = -A$$

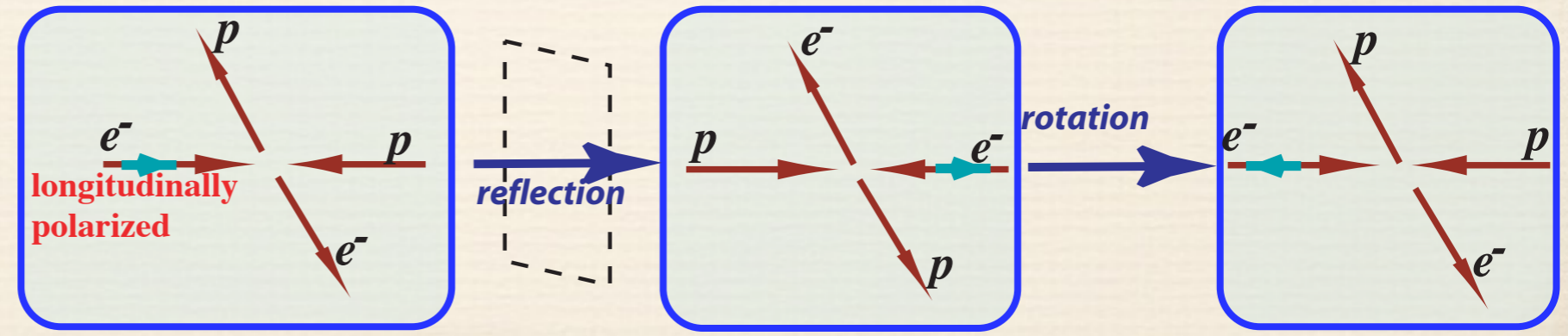
$$\sigma \propto |A_{EM} + A_{weak}|^2$$

$$\sim |A_{EM}|^2 + \boxed{2A_{EM}A_{weak}^* + \dots}$$

$$\sim \frac{A_{weak}}{A_{EM}} \sim \frac{G_F Q^2}{4\pi\alpha}$$

Parity-violating

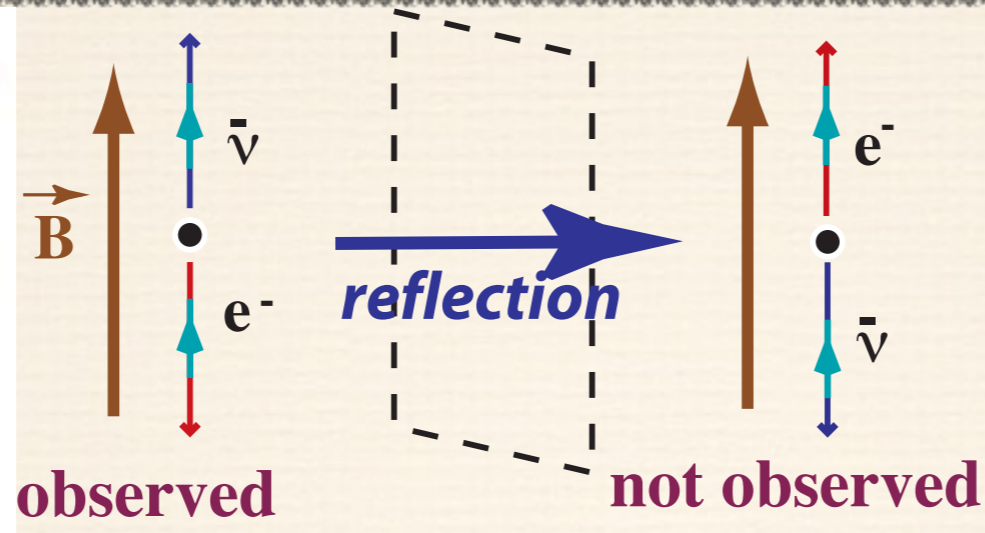
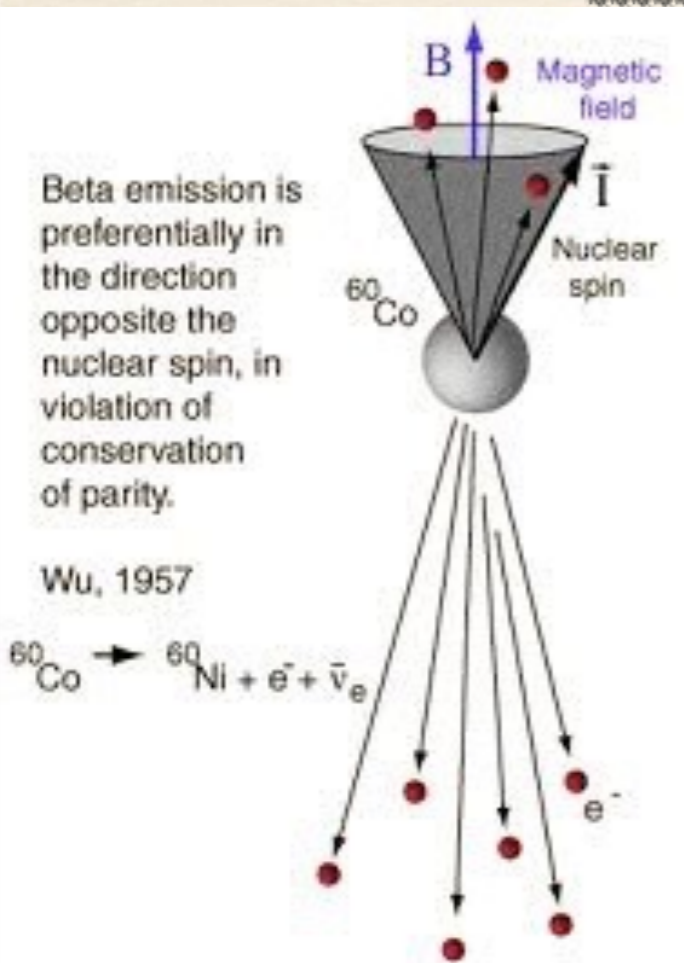
$$A_{PV} \sim 10^{-4} \cdot Q^2 \text{ (GeV}^2\text{)}$$



Zel'dovich speculation: Is Electron Scattering Parity-Violating?

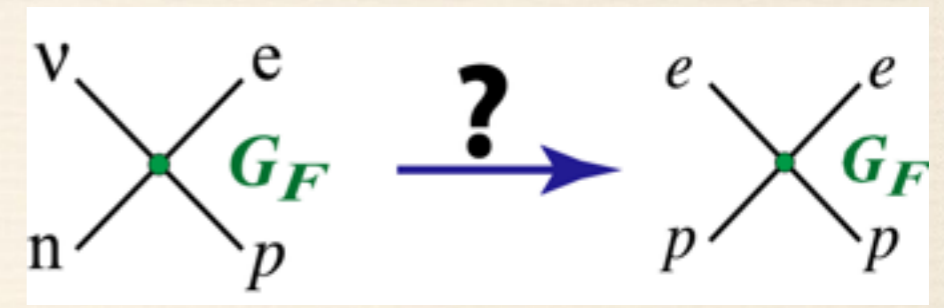
Electroweak Scattering

JETP 36, pp 964-66 (1959)



Neutron β Decay

Electron-proton Weak Scattering



$$A_{PV} = \frac{\sigma_{\uparrow} - \sigma_{\downarrow}}{\sigma_{\uparrow} + \sigma_{\downarrow}} = -A$$

$$\sigma \propto |A_{EM} + A_{weak}|^2$$

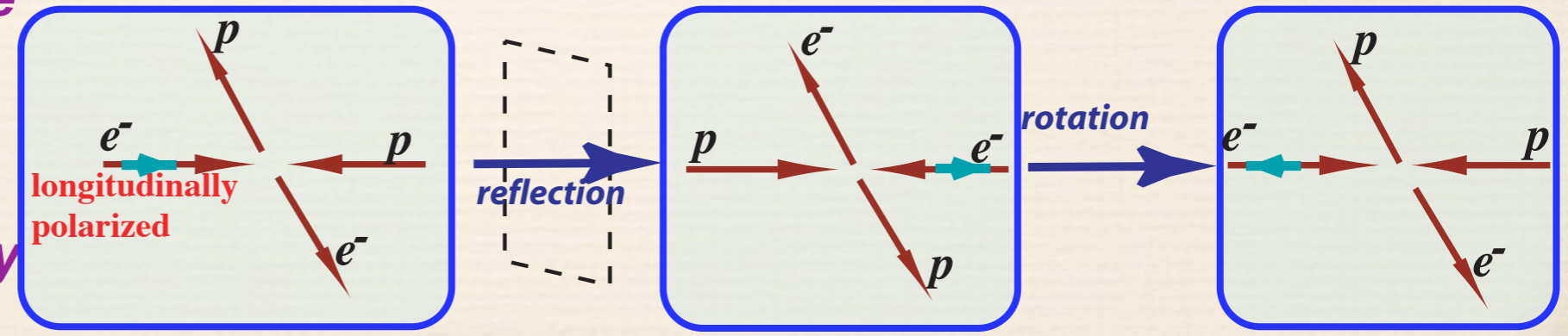
$$\sim |A_{EM}|^2 + \boxed{2A_{EM}A_{weak}^* + \dots}$$

$$\sim \frac{A_{weak}}{A_{EM}} \sim \frac{G_F Q^2}{4\pi\alpha}$$

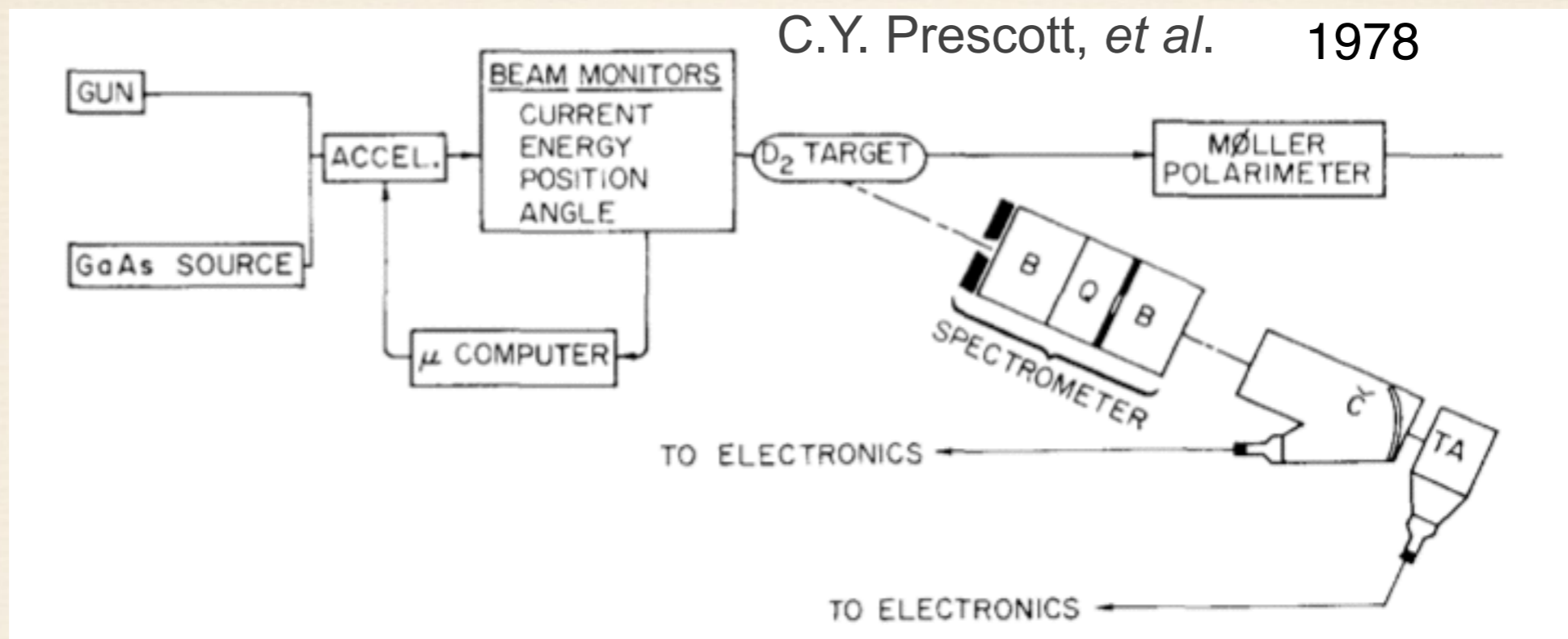
Parity-violating

$$A_{PV} \sim 10^{-4} \cdot Q^2 \text{ (GeV}^2\text{)}$$

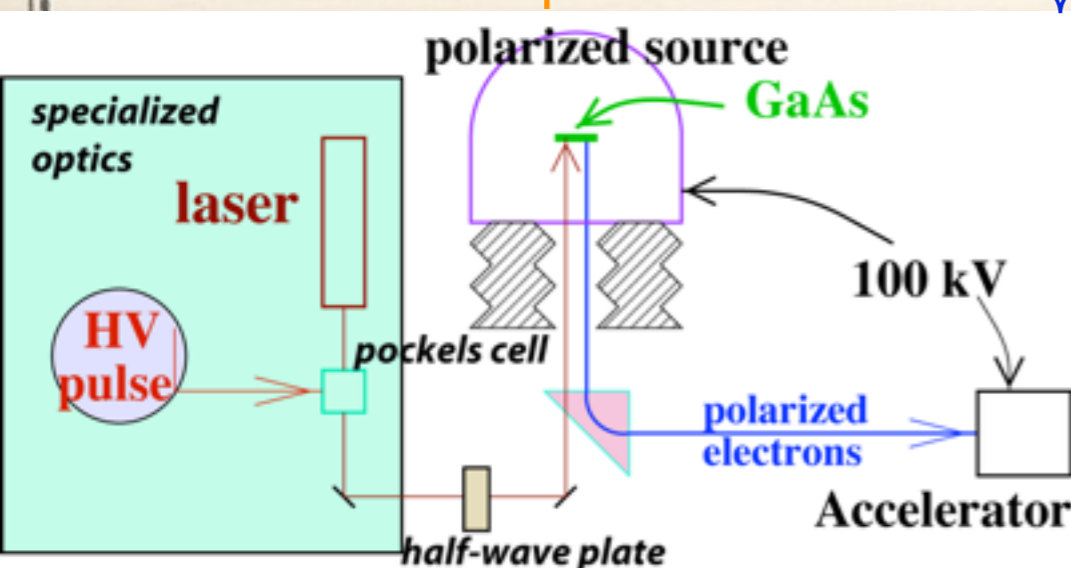
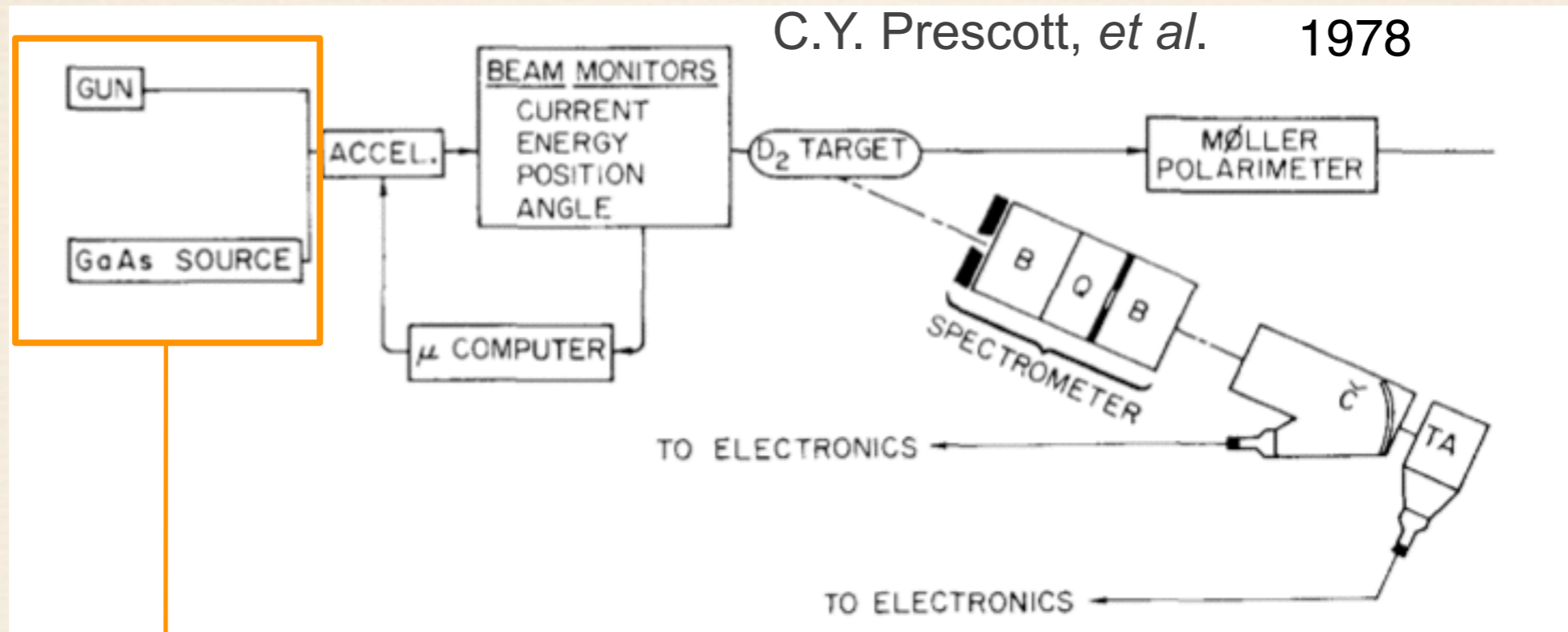
- longitudinally polarize one beam with the ability to change its sign
- Measure fractional rate difference with a sensitivity of a part in 10,000



The first Parity-Violating Electron Scattering (PVES) Experiment Anatomy of E122 at SLAC

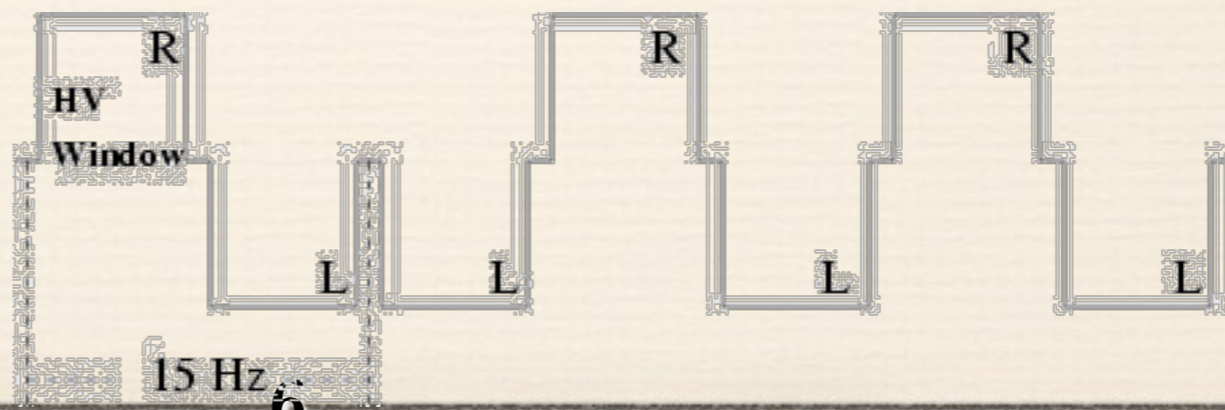


The first Parity-Violating Electron Scattering (PVES) Experiment Anatomy of E122 at SLAC

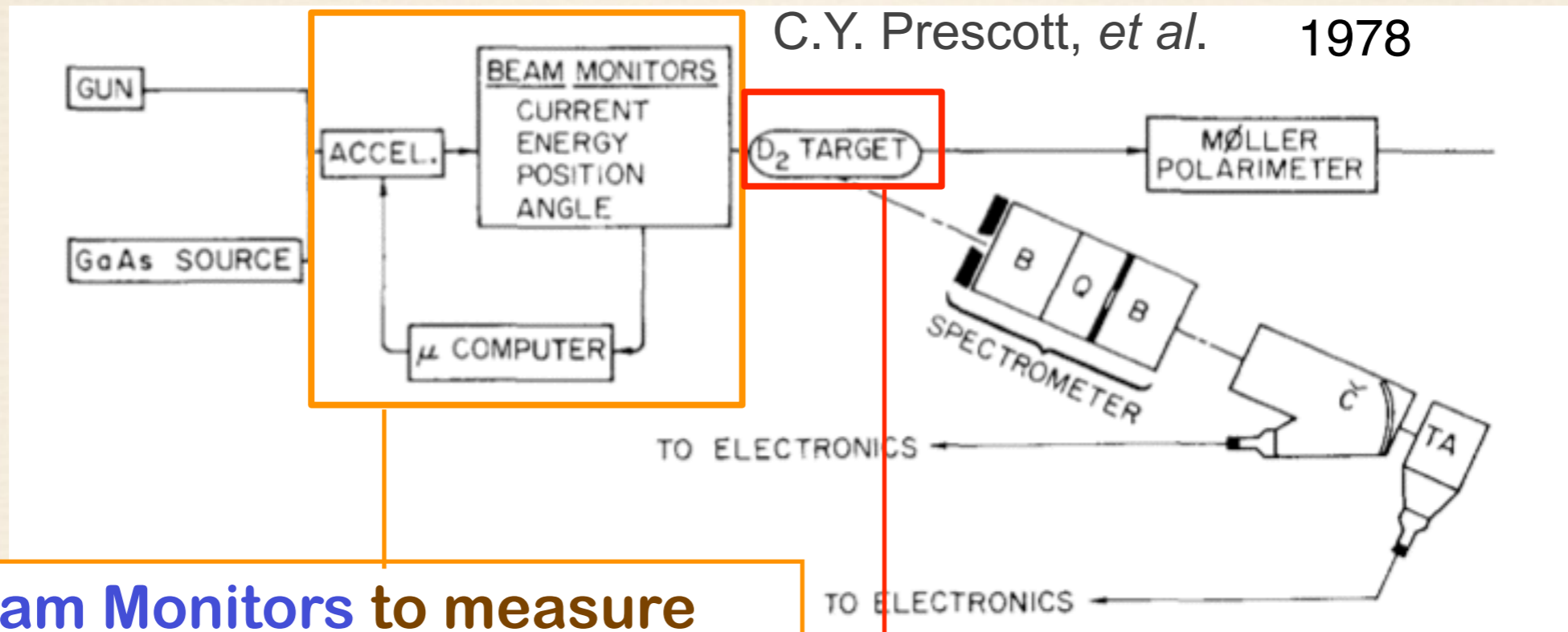


✧ Beam helicity sequence is chosen pseudo-randomly

- Helicity state, followed by its complement
- Data analyzed as "pulse-pairs"



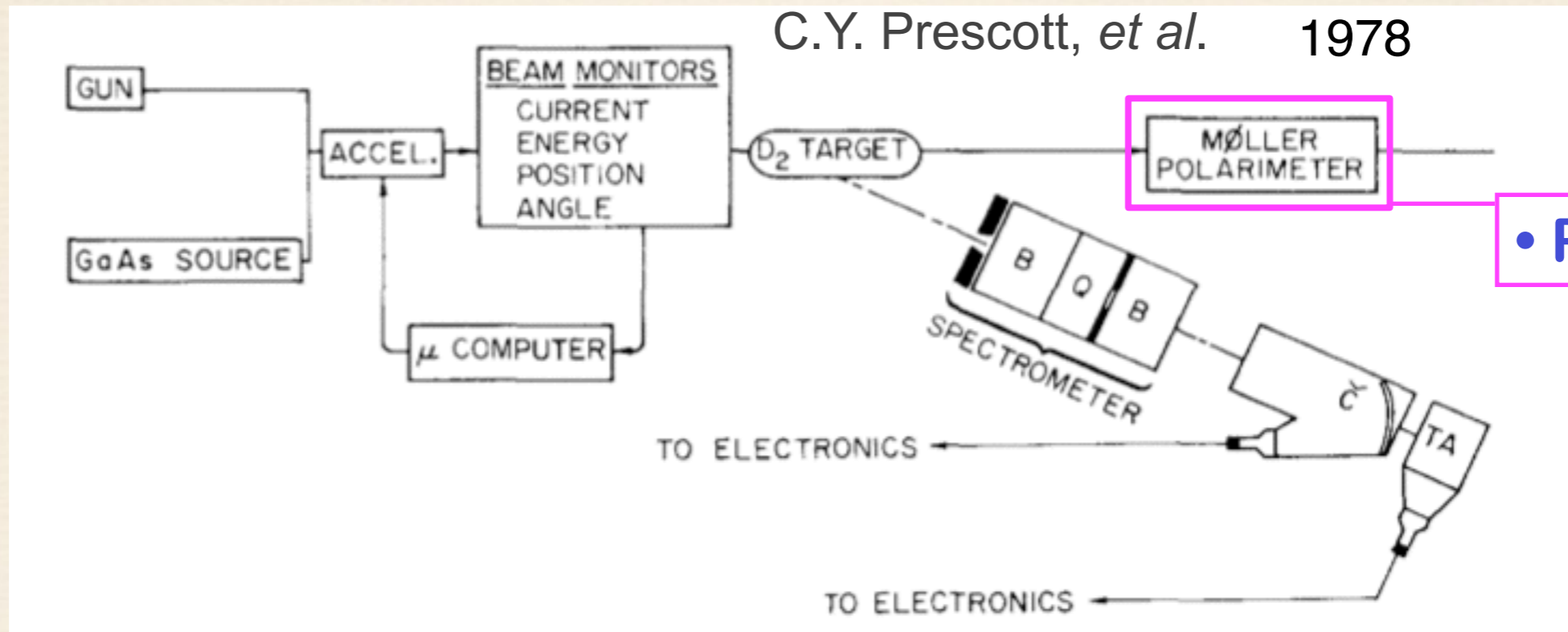
The first Parity-Violating Electron Scattering (PVES) Experiment Anatomy of E122 at SLAC



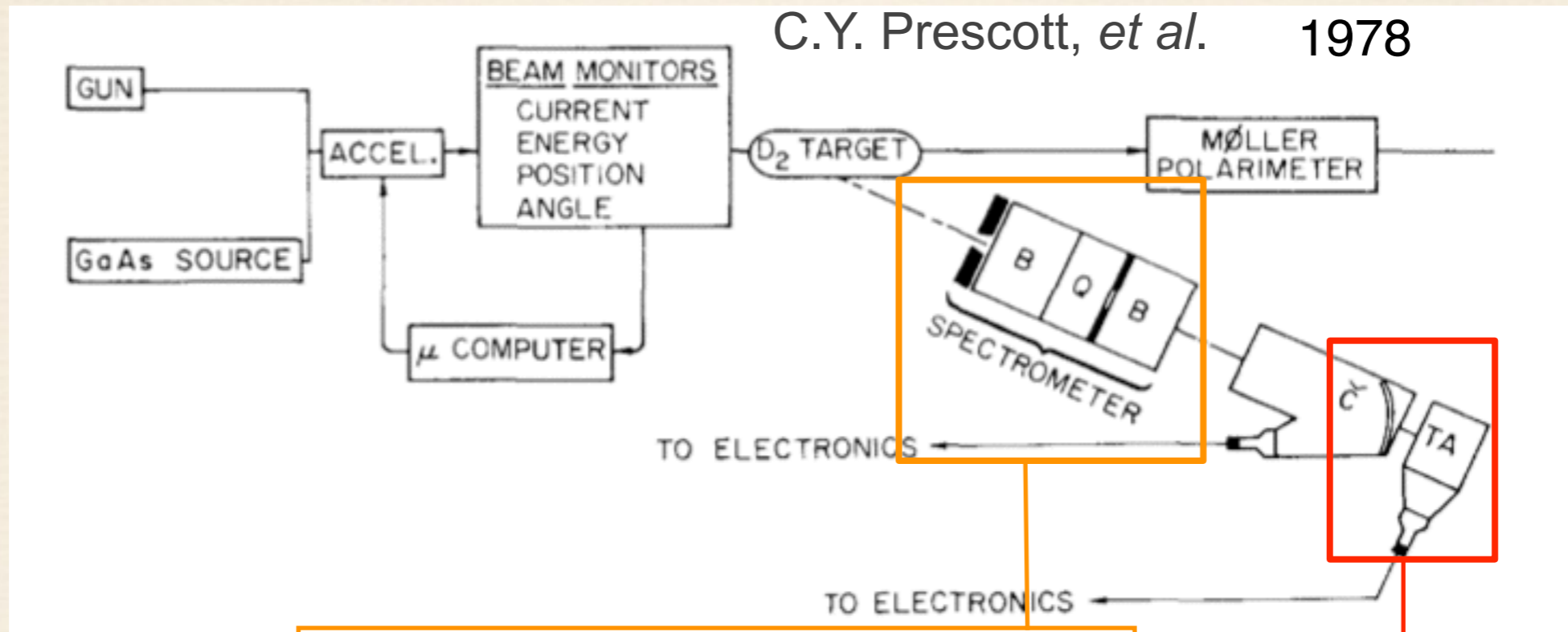
- **Beam Monitors** to measure helicity-correlated changes in beam parameters

- **High-power cryotarget** 30 cm long for high luminosity

The first Parity-Violating Electron Scattering (PVES) Experiment Anatomy of E122 at SLAC

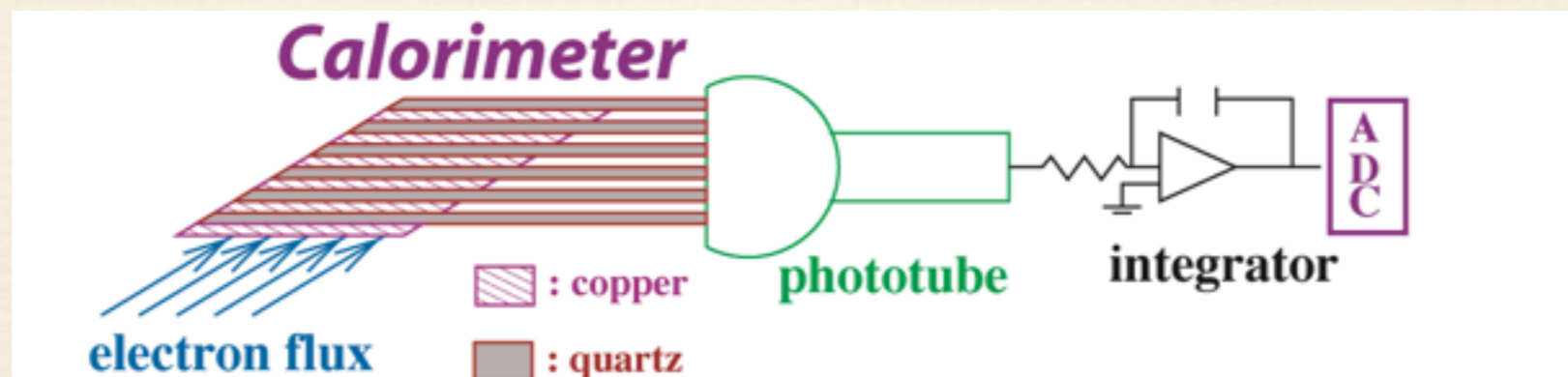


The first Parity-Violating Electron Scattering (PVES) Experiment Anatomy of E122 at SLAC



• **Magnetic spectrometer** directs flux to background-free region

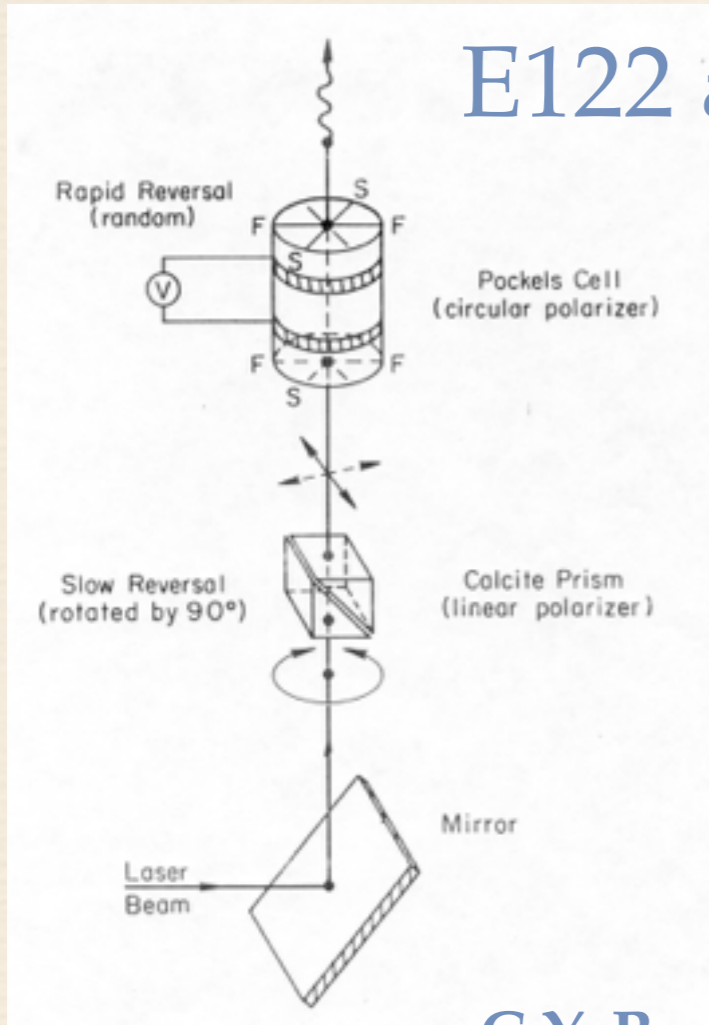
• **Flux Integration** measures high rate without deadtime



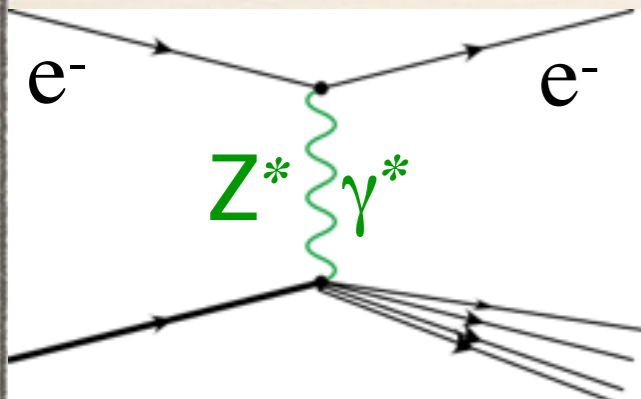
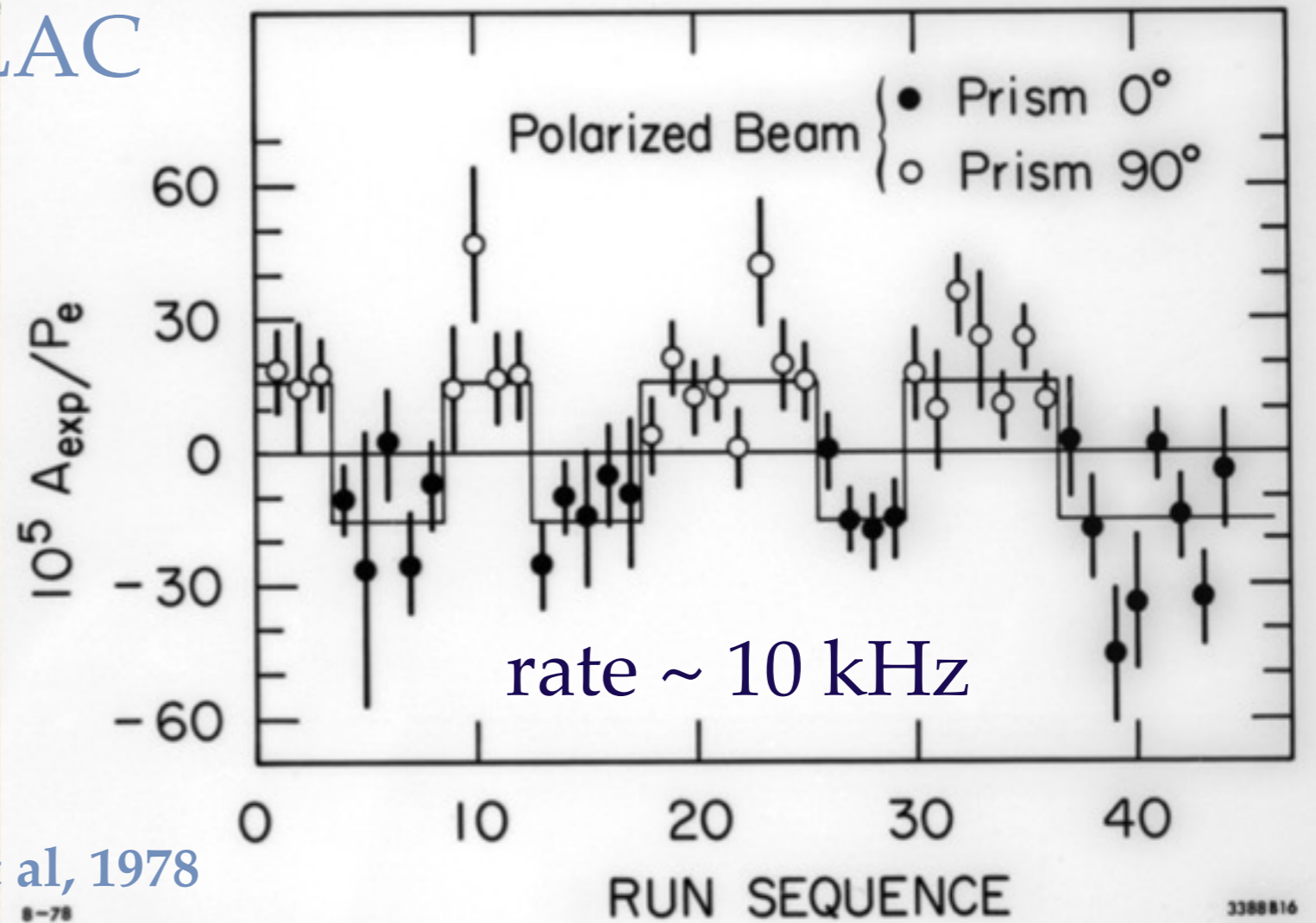
Does the weak neutral current amplitude interfere with the electromagnetic amplitude?

SLAC E122 Result

E122 at SLAC



C.Y. Prescott et al, 1978

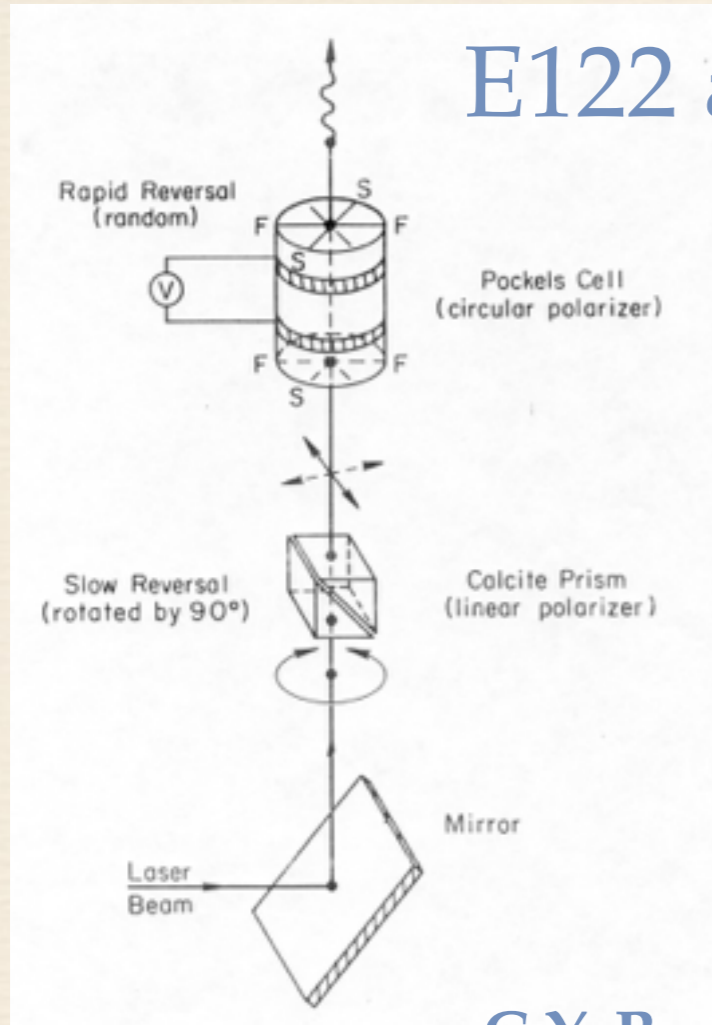


$$A_{PV} \sim 10^{-4}$$

$$\delta(A_{PV}) \sim 10^{-5}$$

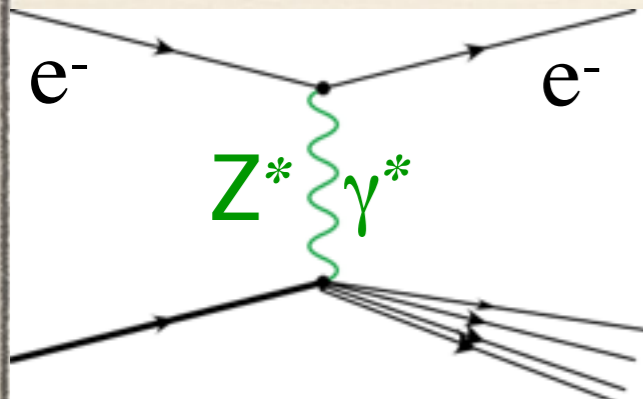
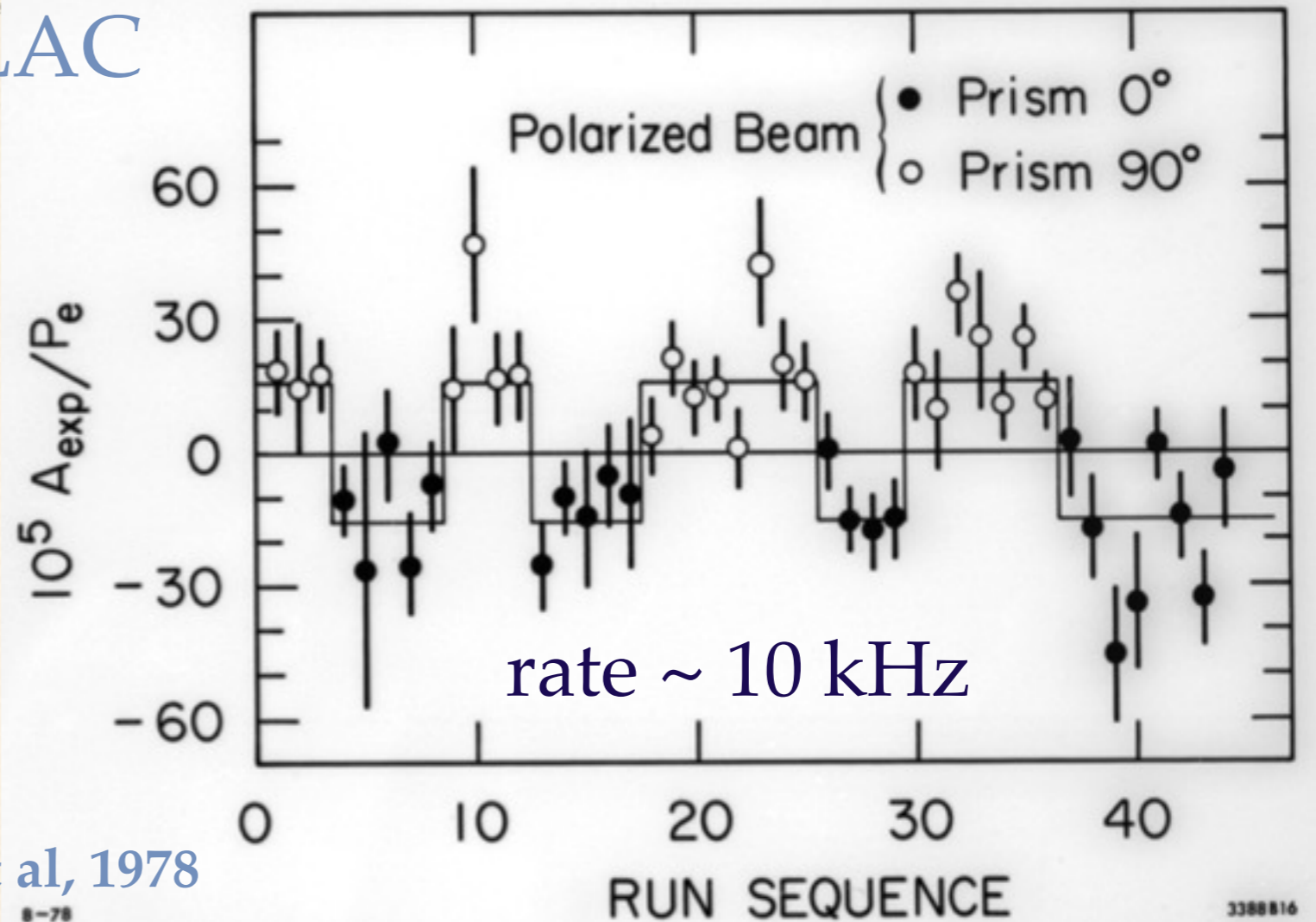
Does the weak neutral current amplitude interfere with the electromagnetic amplitude?

SLAC E122 Result



E122 at SLAC

C.Y. Prescott et al, 1978



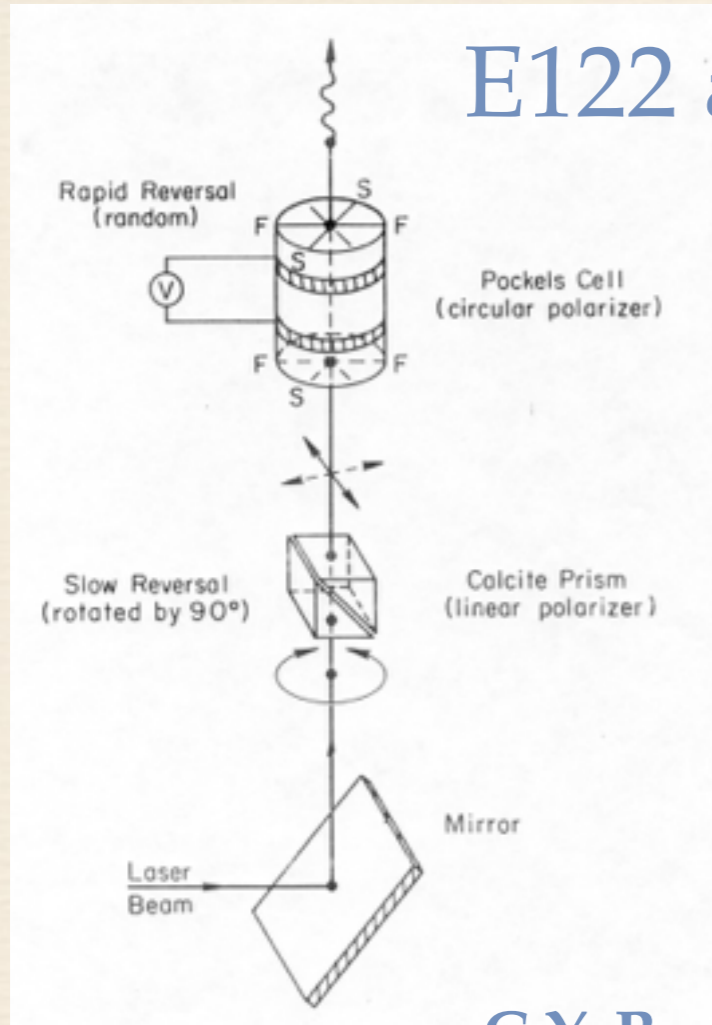
- **Parity Violation in Weak Neutral Current Interactions**
- $\sin^2\theta_W = 0.224 \pm 0.020$: same as in neutrino scattering

$$A_{PV} \sim 10^{-4}$$

$$\delta(A_{PV}) \sim 10^{-5}$$

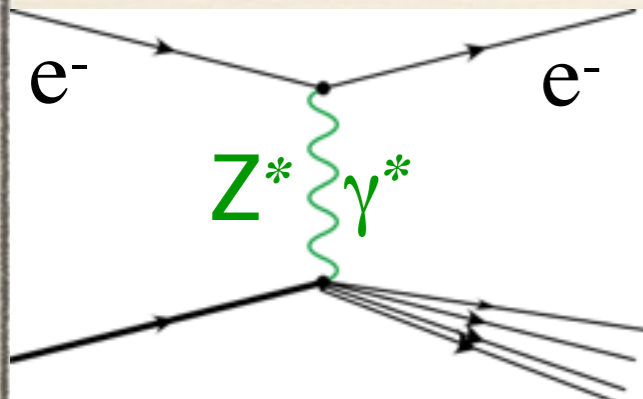
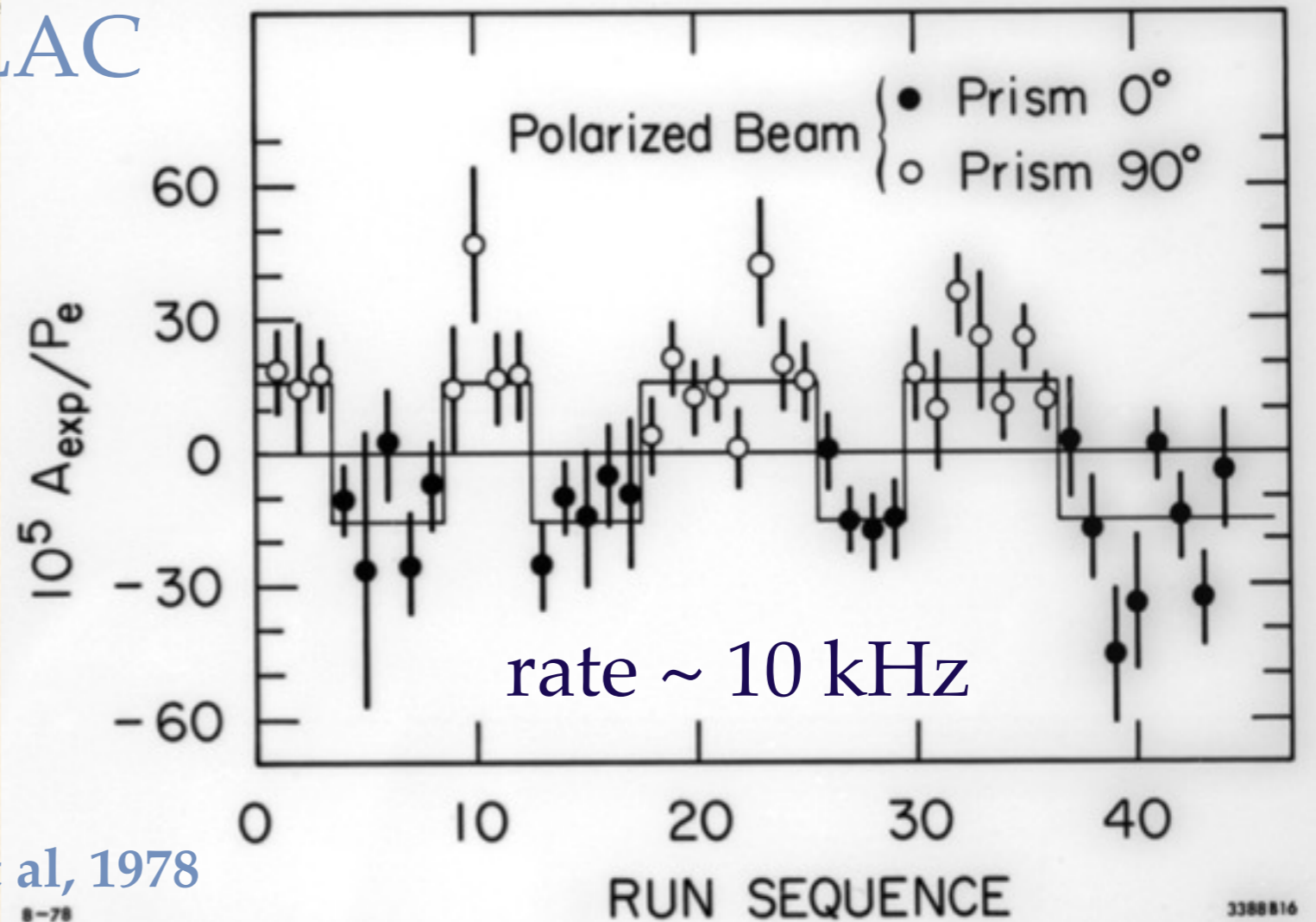
Does the weak neutral current amplitude interfere with the electromagnetic amplitude?

SLAC E122 Result



E122 at SLAC

C.Y. Prescott et al, 1978



- **Parity Violation in Weak Neutral Current Interactions**
- $\sin^2\theta_W = 0.224 \pm 0.020$: same as in neutrino scattering

$$A_{PV} \sim 10^{-4}$$

$$\delta(A_{PV}) \sim 10^{-5}$$

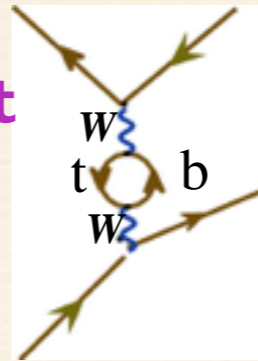
Glashow, Weinberg, Salam Nobel Prize awarded in 1979

Electroweak Theory at 1-Loop

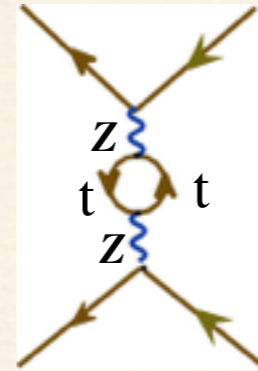
Lecture by A. Freitas

For electroweak interactions, 3 input parameters needed:

1. Rb-87 mass + Ry constant
2. The muon lifetime
3. The Z line shape



Muon decay



Z production

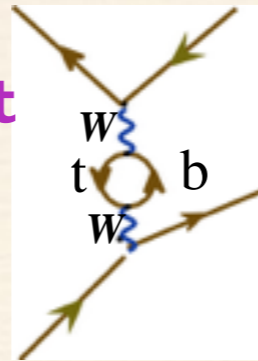
α_{QED} G_F M_Z

Electroweak Theory at 1-Loop

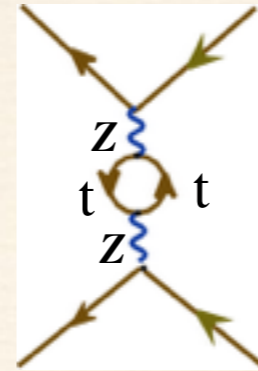
Lecture by A. Freitas

For electroweak interactions, 3 input parameters needed:

1. Rb-87 mass + Ry constant
2. The muon lifetime
3. The Z line shape



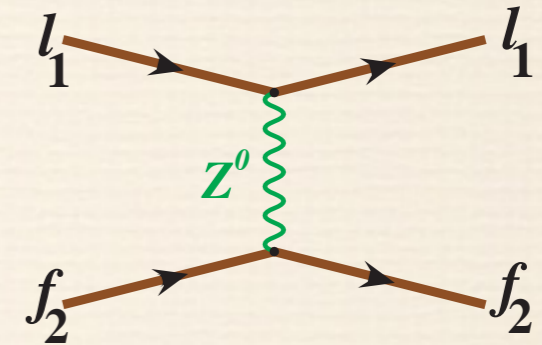
Muon decay



Z production

Weak Neutral Current interactions

4th and 5th best measured parameters:
 M_W and $\sin^2\theta_W$



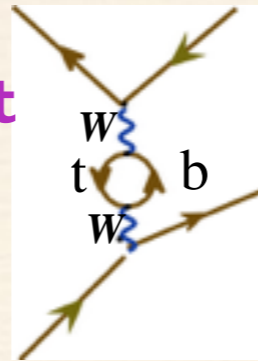
α_{QED} G_F M_Z

Electroweak Theory at 1-Loop

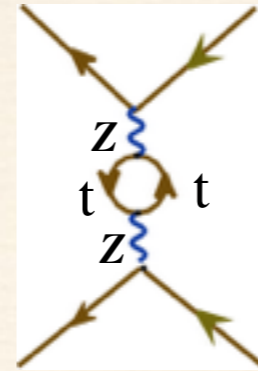
Lecture by A. Freitas

For electroweak interactions, 3 input parameters needed:

1. Rb-87 mass + Ry constant
2. The muon lifetime
3. The Z line shape



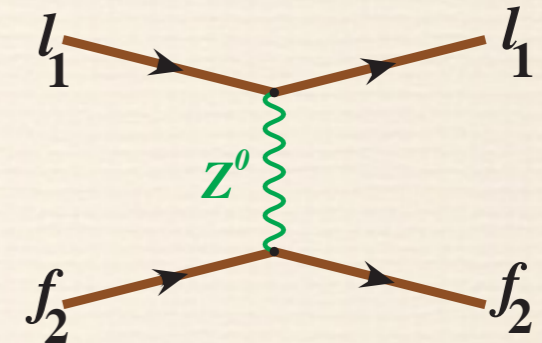
Muon decay



Z production

Weak Neutral Current interactions

4th and 5th best measured parameters:
 M_W and $\sin^2\theta_W$



simple definition; disfavored due to heavy m_t

$$\alpha_{QED} \quad G_F \quad M_Z$$

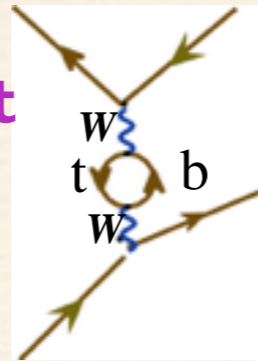
$$\sin^2 \theta_W \equiv 1 - m_W^2 / m_Z^2$$

Electroweak Theory at 1-Loop

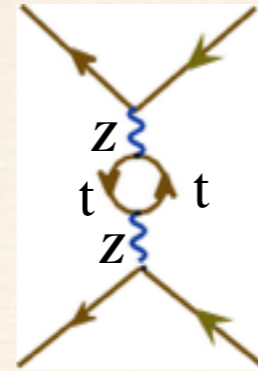
Lecture by A. Freitas

For electroweak interactions, 3 input parameters needed:

1. **Rb-87 mass + Ry constant**
2. **The muon lifetime**
3. **The Z line shape**

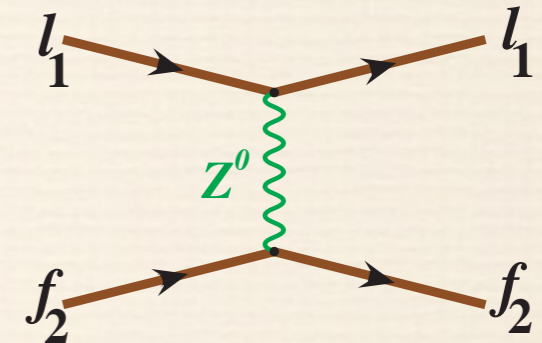


Muon decay



Z production

**4th and 5th best measured parameters:
M_W and sin²θ_W**



Weak Neutral Current interactions

$$\sin^2 \theta_W \equiv 1 - m_W^2 / m_Z^2$$

simple definition; disfavored due to heavy m_t

$$\sin^2 \theta_W^{eff} \equiv (1 - g_{\mu\mu Z}) / 4$$

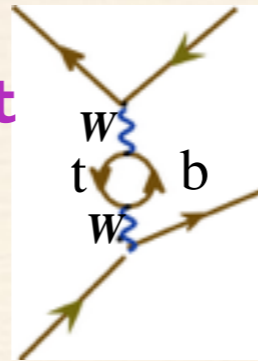
good at Z-pole; nasty counterterms at other scales

Electroweak Theory at 1-Loop

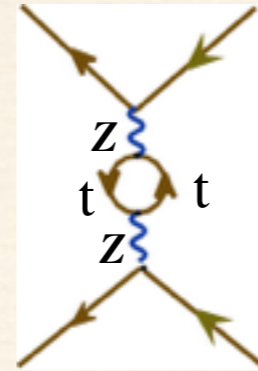
Lecture by A. Freitas

For electroweak interactions, 3 input parameters needed:

1. **Rb-87 mass + Ry constant**
2. **The muon lifetime**
3. **The Z line shape**

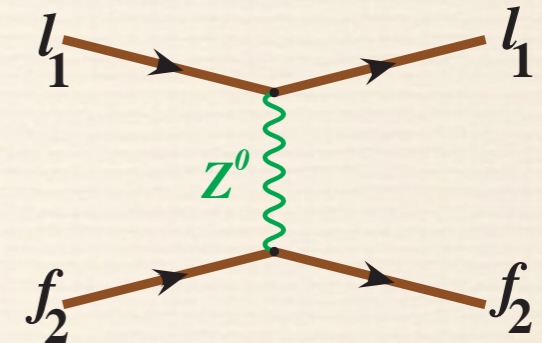


Muon decay



Z production

**4th and 5th best measured parameters:
M_W and sin²θ_W**



Weak Neutral Current interactions

$$\sin^2 \theta_W \equiv 1 - m_W^2 / m_Z^2$$

simple definition; disfavored due to heavy m_t

$$\sin^2 \theta_W^{eff} \equiv (1 - g_{\mu\mu Z}) / 4$$

good at Z-pole; nasty counterterms at other scales

$$\sin^2 \theta_W(\mu)_{\overline{MS}} \equiv e^2(\mu)_{\overline{MS}} / g^2(\mu)_{\overline{MS}}$$

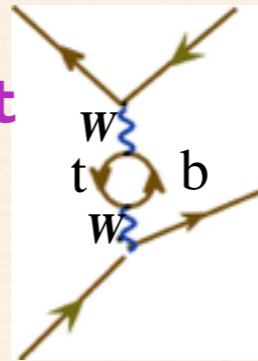
theoretically motivated; but not physical

Electroweak Theory at 1-Loop

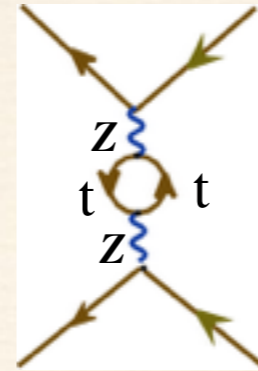
Lecture by A. Freitas

For electroweak interactions, 3 input parameters needed:

1. **Rb-87 mass + Ry constant**
2. **The muon lifetime**
3. **The Z line shape**

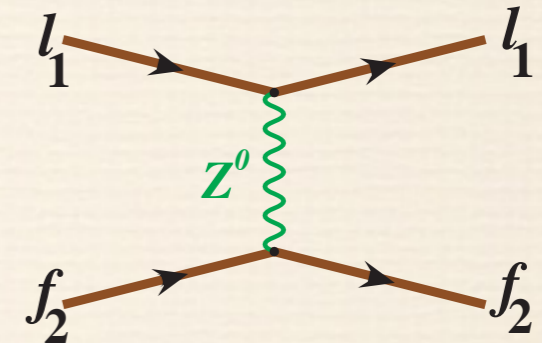


Muon decay



Z production

**4th and 5th best measured parameters:
M_W and sin²θ_W**



Weak Neutral Current interactions

$$\sin^2 \theta_W \equiv 1 - m_W^2 / m_Z^2$$

simple definition; disfavored due to heavy m_t

$$\sin^2 \theta_W^{eff} \equiv (1 - g_{\mu\mu Z}) / 4$$

good at Z-pole; nasty counterterms at other scales

$$\sin^2 \theta_W(M_Z)_{\overline{MS}} = \sin^2 \theta_W^{eff} - 0.00028$$

$$\sin^2 \theta_W(\mu)_{\overline{MS}} \equiv e^2(\mu)_{\overline{MS}} / g^2(\mu)_{\overline{MS}}$$

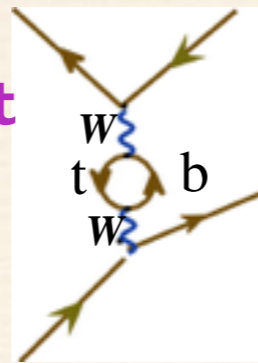
theoretically motivated; but not physical

Electroweak Theory at 1-Loop

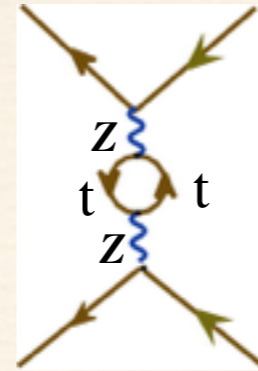
Lecture by A. Freitas

For electroweak interactions, 3 input parameters needed:

1. **Rb-87 mass + Ry constant**
2. **The muon lifetime**
3. **The Z line shape**

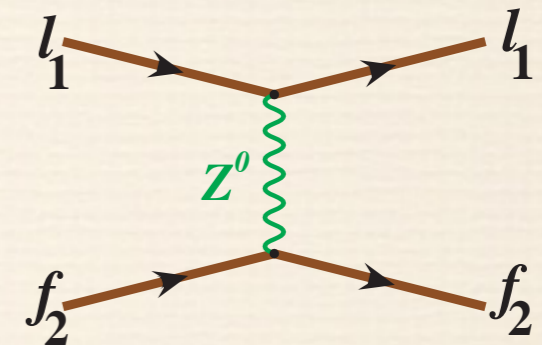


Muon decay



Z production

**4th and 5th best measured parameters:
M_W and sin²θ_W**



Weak Neutral Current interactions

$$\sin^2 \theta_W \equiv 1 - m_W^2 / m_Z^2$$

simple definition; disfavored due to heavy m_t

$$\sin^2 \theta_W^{eff} \equiv (1 - g_{\mu\mu Z}) / 4$$

good at Z-pole; nasty counterterms at other scales

$$\sin^2 \theta_W(M_Z)_{\overline{MS}} = \sin^2 \theta_W^{eff} - 0.00028$$

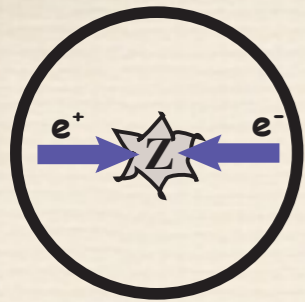
$$\sin^2 \theta_W(\mu)_{\overline{MS}} \equiv e^2(\mu)_{\overline{MS}} / g^2(\mu)_{\overline{MS}}$$

theoretically motivated; but not physical

Theory Predictions

$$\sin^2 \theta_W(m_Z)_{\overline{MS}} = 0.23122 (2)$$

$$M_W = 80.358 (4) \text{ GeV}$$

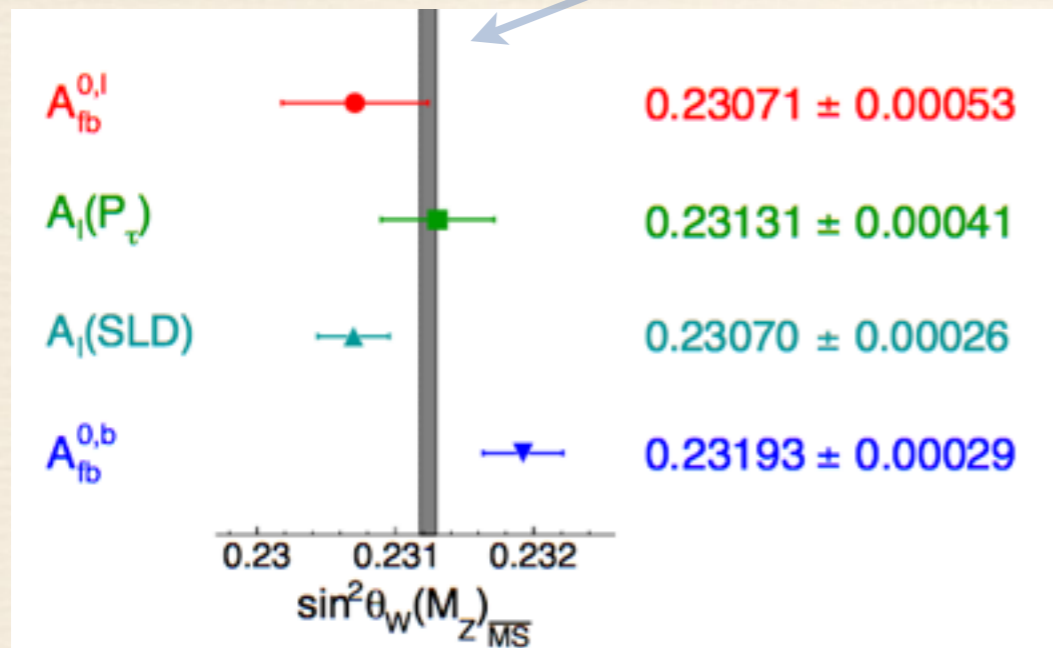


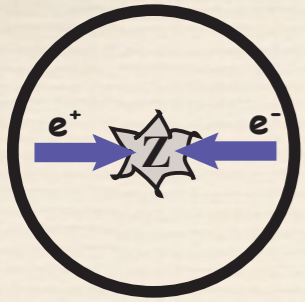
Theory vs Experiment

The most precise measurements at LEP/SLC

colliders:
LEP, SLC

Prediction for 125 GeV Higgs



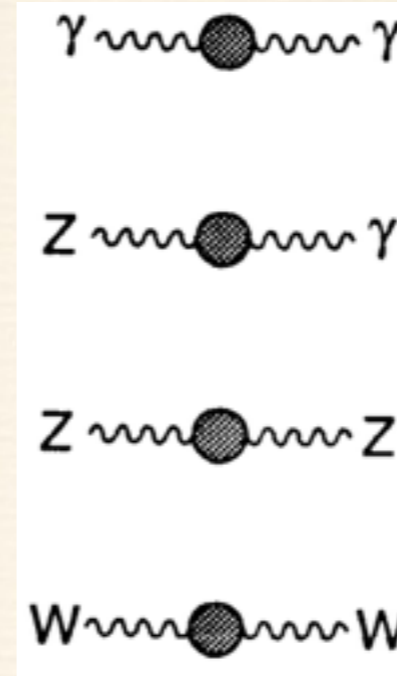
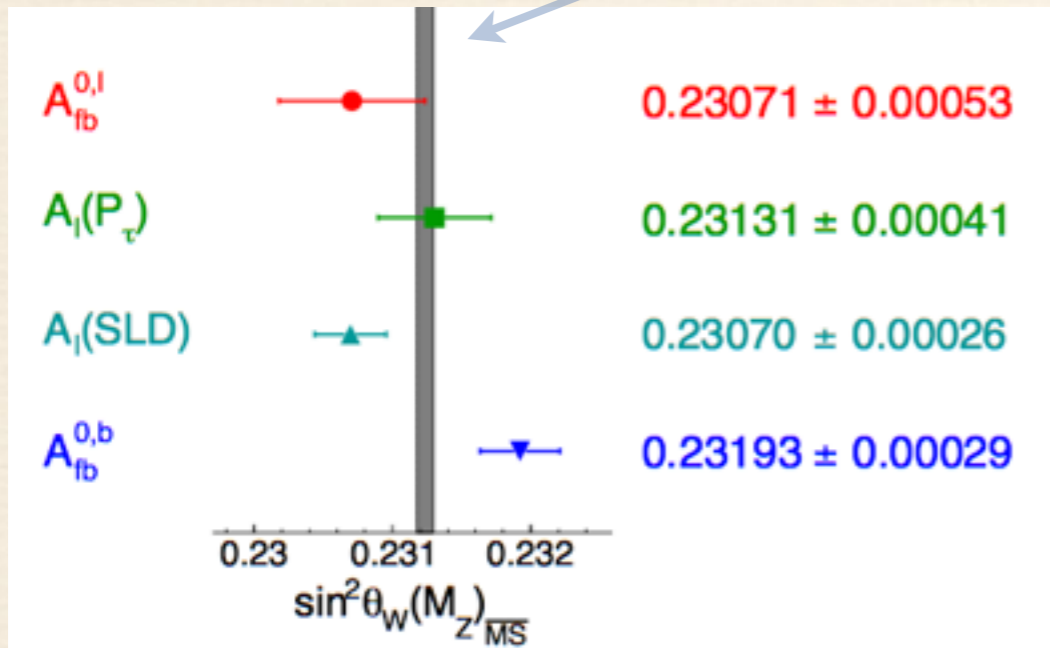


Theory vs Experiment

The most precise measurements at LEP/SLC

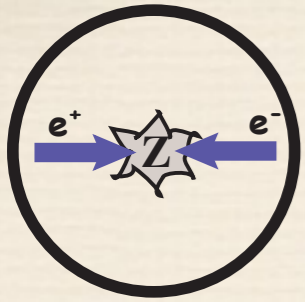
colliders:
LEP, SLC

Prediction for 125 GeV Higgs



S, T, U
parameters

Stringent constraints
on large classes of
new physics models

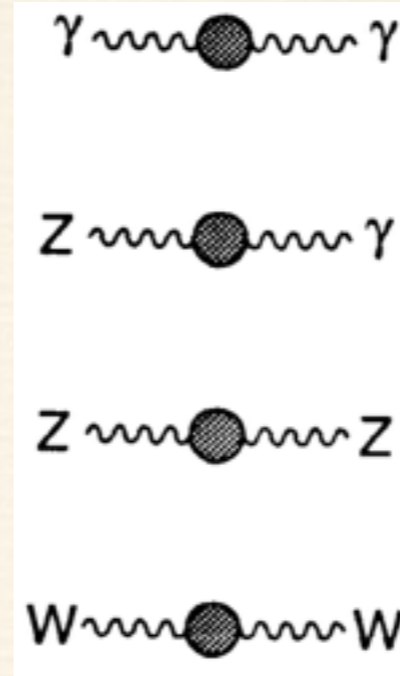
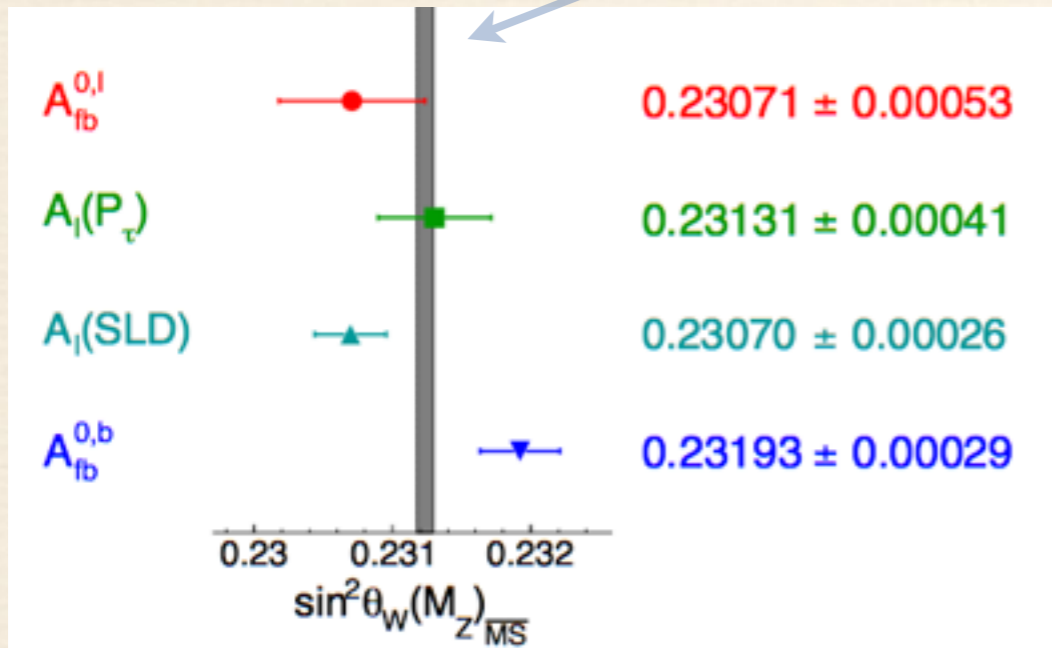


Theory vs Experiment

The most precise measurements at LEP/SLC

colliders:
LEP, SLC

Prediction for 125 GeV Higgs



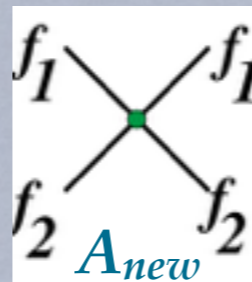
S, T, U
parameters

Stringent constraints
on large classes of
new physics models

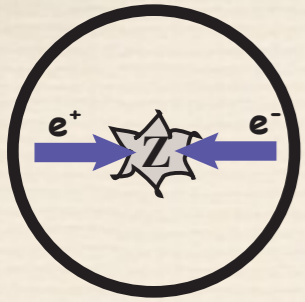
Flavor Diagonal Contact Interactions

Consider $f_1 \bar{f}_1 \rightarrow f_2 \bar{f}_2$ or $f_1 f_2 \rightarrow f_1 f_2$

$$L_{f_1 f_2} = \sum_{i,j=L,R} \frac{4\pi}{\Lambda_{ij}^2} \eta_{ij} \bar{f}_{1i} \gamma_\mu f_{1i} \bar{f}_{2j} \gamma^\mu f_{2j}$$



New heavy physics that does not couple directly to SM gauge bosons

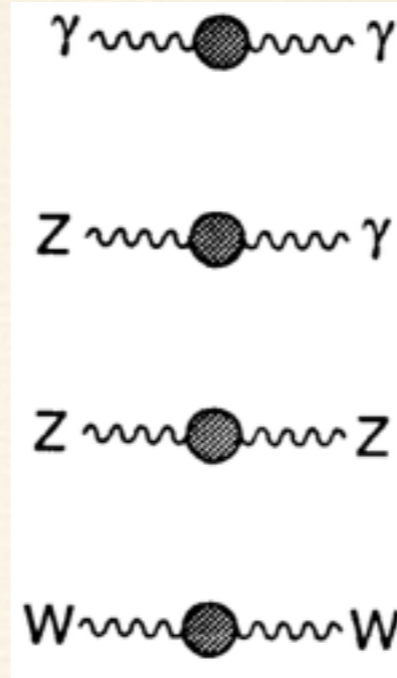
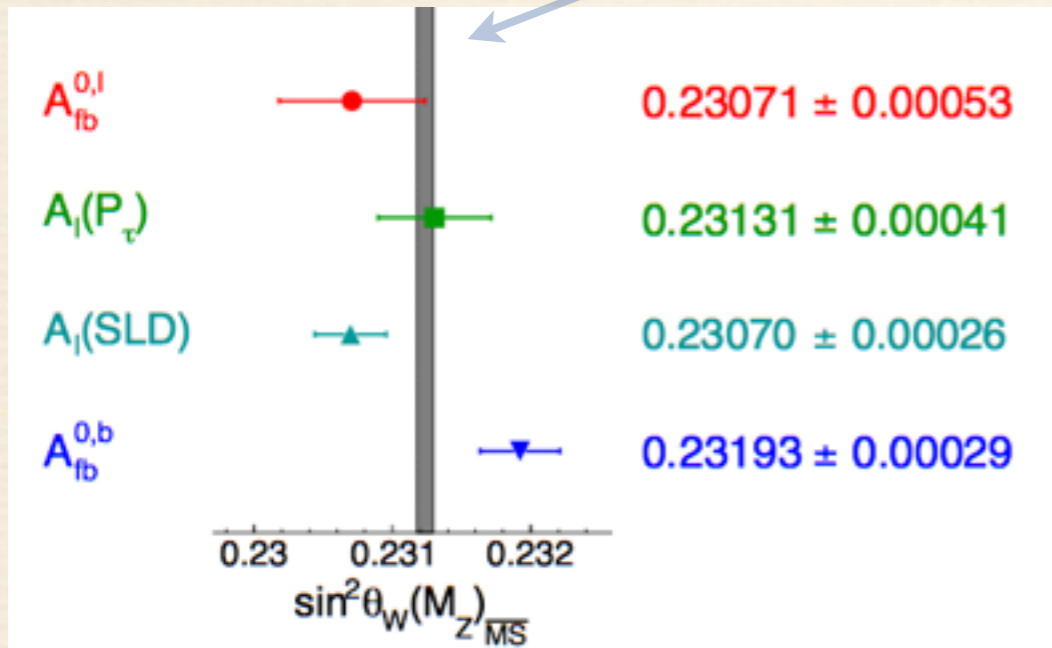


Theory vs Experiment

The most precise measurements at LEP/SLC

colliders:
LEP, SLC

Prediction for 125 GeV Higgs



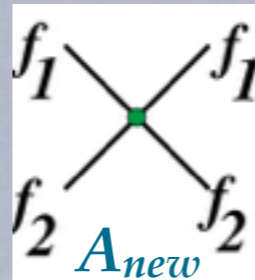
S, T, U
parameters

Stringent constraints
on large classes of
new physics models

Flavor Diagonal Contact Interactions

Consider $f_1 \bar{f}_1 \rightarrow f_2 \bar{f}_2$ or $f_1 f_2 \rightarrow f_1 f_2$

$$L_{f_1 f_2} = \sum_{i,j=L,R} \frac{4\pi}{\Lambda_{ij}^2} \eta_{ij} \bar{f}_{1i} \gamma_\mu f_{1i} \bar{f}_{2j} \gamma^\mu f_{2j}$$

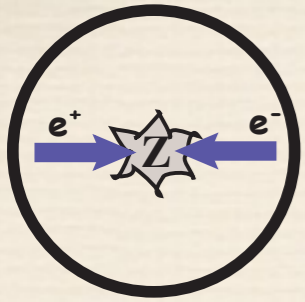


New heavy physics that does not couple directly to SM gauge bosons

on resonance: A_Z is imaginary

$$\left| A_Z + A_{\text{new}} \right|^2 \rightarrow A_Z^2 \left[1 + \left(\frac{A_{\text{new}}}{A_Z} \right)^2 \right]$$

no interference!

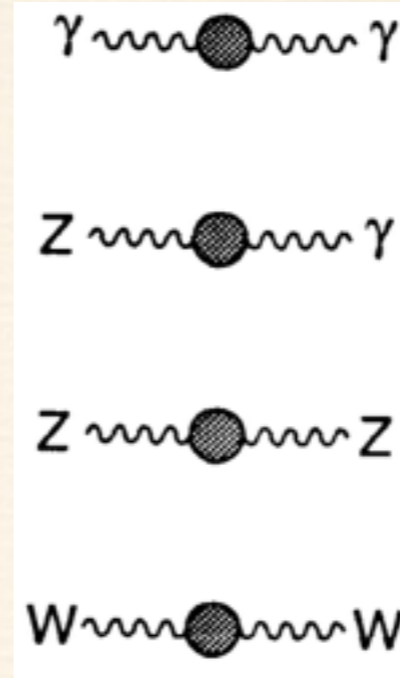
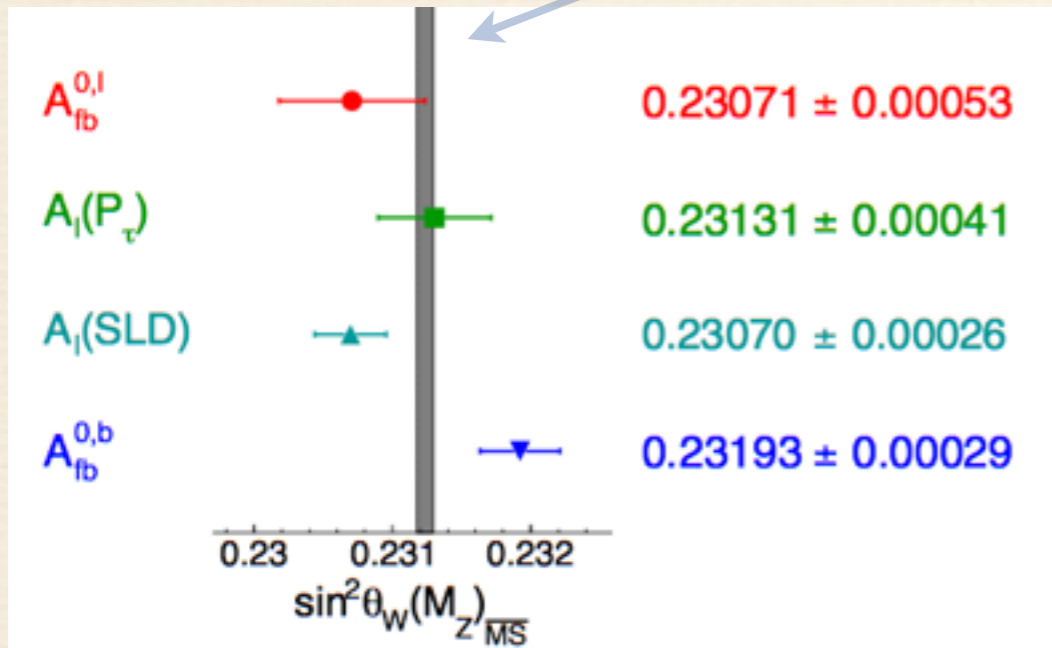


Theory vs Experiment

The most precise measurements at LEP/SLC

colliders:
LEP, SLC

Prediction for 125 GeV Higgs



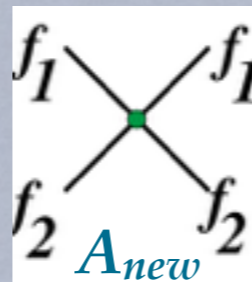
S, T, U
parameters

Stringent constraints
on large classes of
new physics models

Flavor Diagonal Contact Interactions

Consider $f_1 \bar{f}_1 \rightarrow f_2 \bar{f}_2$ or $f_1 f_2 \rightarrow f_1 f_2$

$$L_{f_1 f_2} = \sum_{i,j=L,R} \frac{4\pi}{\Lambda_{ij}^2} \eta_{ij} \bar{f}_{1i} \gamma_\mu f_{1i} \bar{f}_{2j} \gamma^\mu f_{2j}$$



on resonance: A_Z is imaginary

$$\left| A_Z + A_{\text{new}} \right|^2 \rightarrow A_Z^2 \left[1 + \left(\frac{A_{\text{new}}}{A_Z} \right)^2 \right]$$

no interference!

New heavy physics that does not couple directly to SM gauge bosons

Unique role for $\sin^2 \theta_w$ measurements at $Q^2 \ll M_Z^2$

*Modern Low Q^2 Weak
Neutral Current
Measurements*

Physics down to a length scale of 10^{-19} m well understood but....

Modern Electroweak Physics

Many questions still unanswered....

The High Energy Frontier: Collider Physics

The Cosmic Frontier: Particle, Nuclear and Gravitational Astrophysics

A comprehensive search for clues requires, in addition:

The Intensity/Precision Frontier

Physics down to a length scale of 10^{-19} m well understood but....

Modern Electroweak Physics

Many questions still unanswered....

The High Energy Frontier: Collider Physics

The Cosmic Frontier: Particle, Nuclear and Gravitational Astrophysics

A comprehensive search for clues requires, in addition:

The Intensity/Precision Frontier

◆ **Violation of Approximate (?) Symmetries**

★ **Neutrinoless Double-Beta Decay, EDMs, CLFV,...**

◆ **Direct Detection of Dark Matter**

◆ **Measurements of Neutrino Masses and Mixing**

◆ **Precise Measurements of SM observables**

*Intense beams, ultra-high precision, exotic nuclei,
table-top experiments, rare processes....*

BSM Indirect Searches

courtesy
V. Cirigliano,
H. Maruyama,
M. Pospelov

High Energy Dynamics

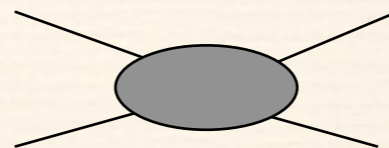
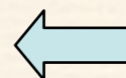
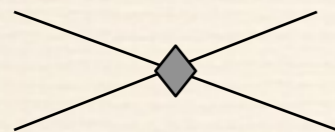
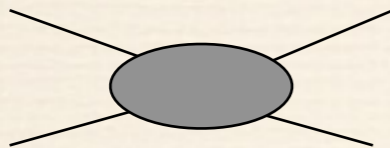
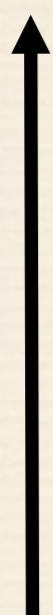
$$\mathcal{L} = \mathcal{L}_{SM} + \frac{1}{\Lambda} \mathcal{L}_5 + \frac{1}{\Lambda^2} \mathcal{L}_6 + \dots$$

higher dimensional operators
can be systematically classified

Λ (\sim TeV)

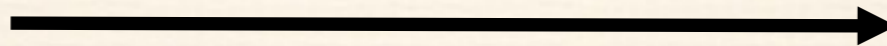
$M_{W,Z}$
(100 GeV)

E



Dark Sector

$(\text{coupling})^{-1}$



Heavy Z's, light (dark) Z's, L-R models, compositeness, extra dimensions, SUSY...

Weak Neutral Current Interactions (WNC) at $Q^2 \ll M_Z^2$

BSM Indirect Searches

courtesy
V. Cirigliano,
H. Maruyama,
M. Pospelov

High Energy Dynamics

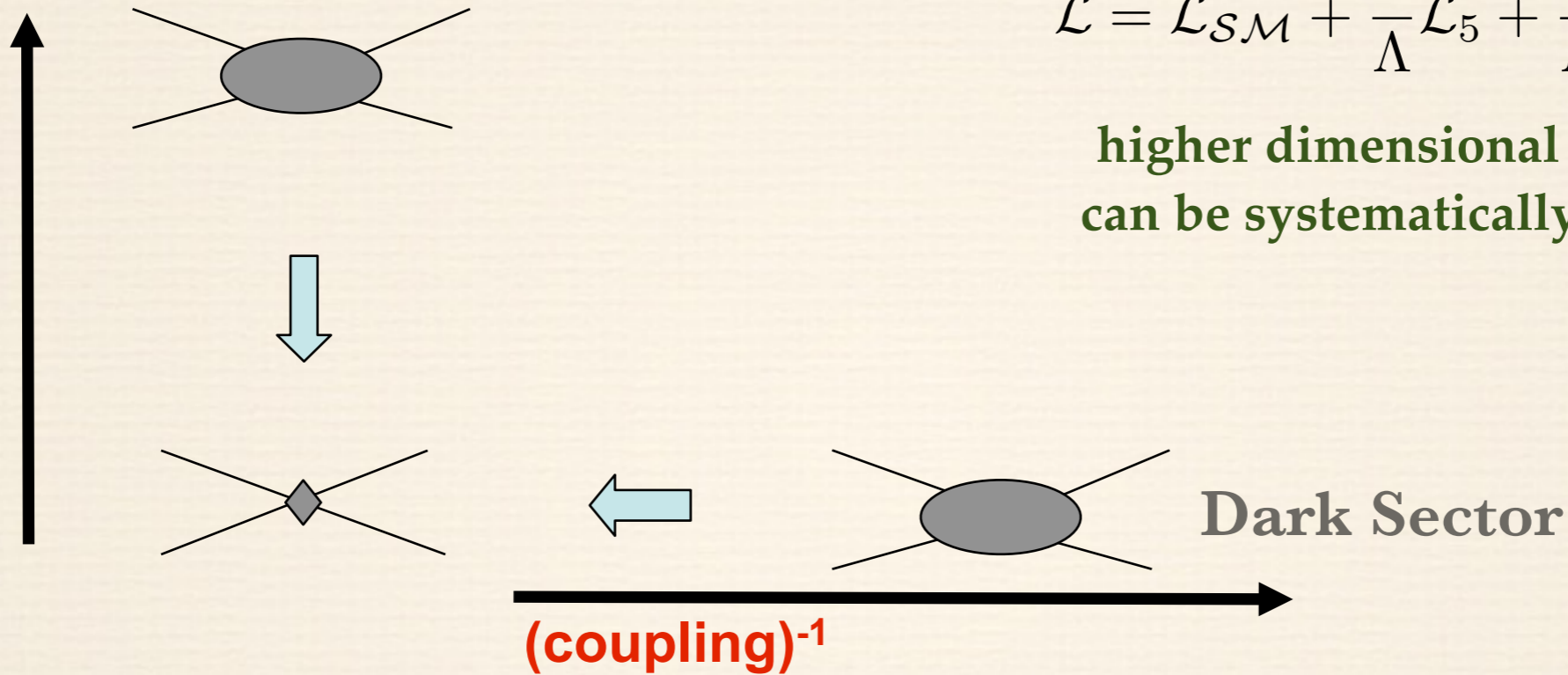
$$\mathcal{L} = \mathcal{L}_{SM} + \frac{1}{\Lambda} \mathcal{L}_5 + \frac{1}{\Lambda^2} \mathcal{L}_6 + \dots$$

higher dimensional operators
can be systematically classified

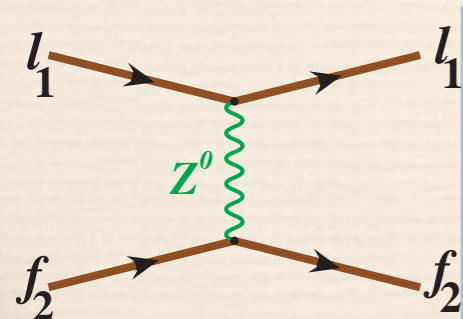
Λ (\sim TeV)

$M_{W,Z}$
(100 GeV)

E



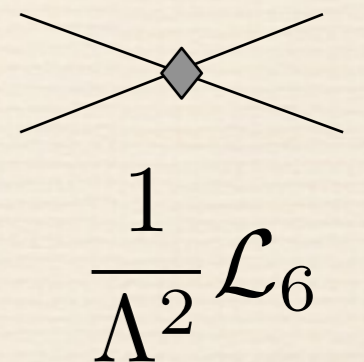
Heavy Z's, light (dark) Z's, L-R models, compositeness, extra dimensions, SUSY...



Search for new flavor diagonal CP-conserving
neutral currents

Tiny yet measurable deviations from
SM processes with precise predictions

must reach $\Lambda \sim 10$ TeV



Thumb Rule: Weak mixing angle must be measured to sub-1% precision

WNC “Bookkeeping”



◆ Atomic Parity Violation: Cs-133

◆ future measurements and theory challenging

◆ Neutrino Deep Inelastic Scattering: NuTeV

◆ future measurements and theory challenging

◆ PV Møller Scattering: E158 at SLAC

◆ statistics limited, theory robust

◆ next generation: MOLLER (factor of 5 better)

◆ PV elastic e-p scattering: Qweak

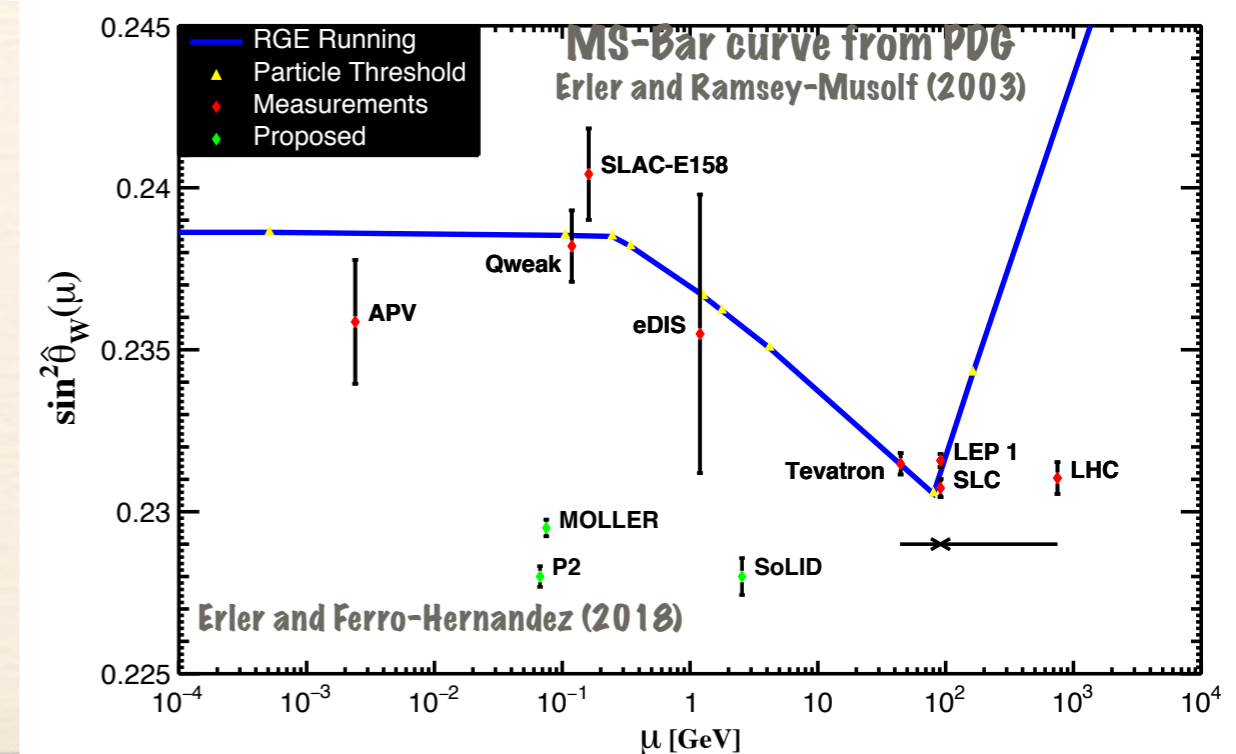
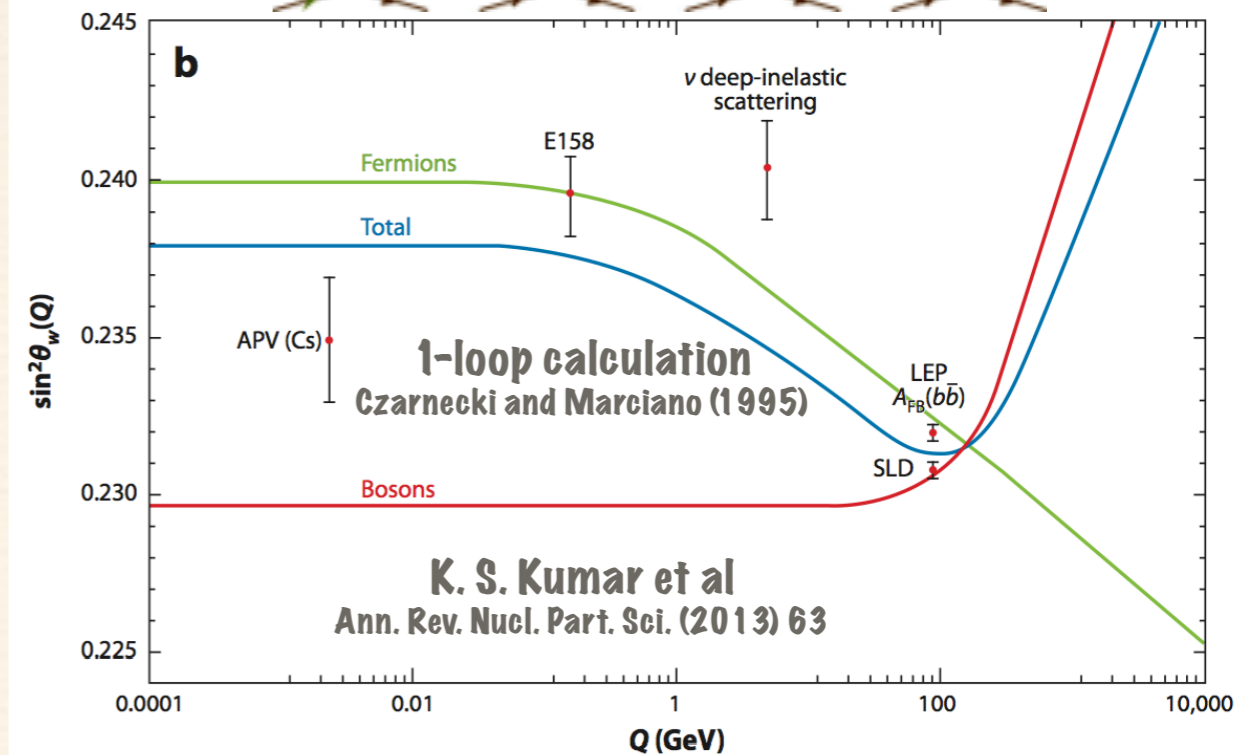
◆ theory robust at low beam energy

◆ next generation: P2 (factor of 3 better)

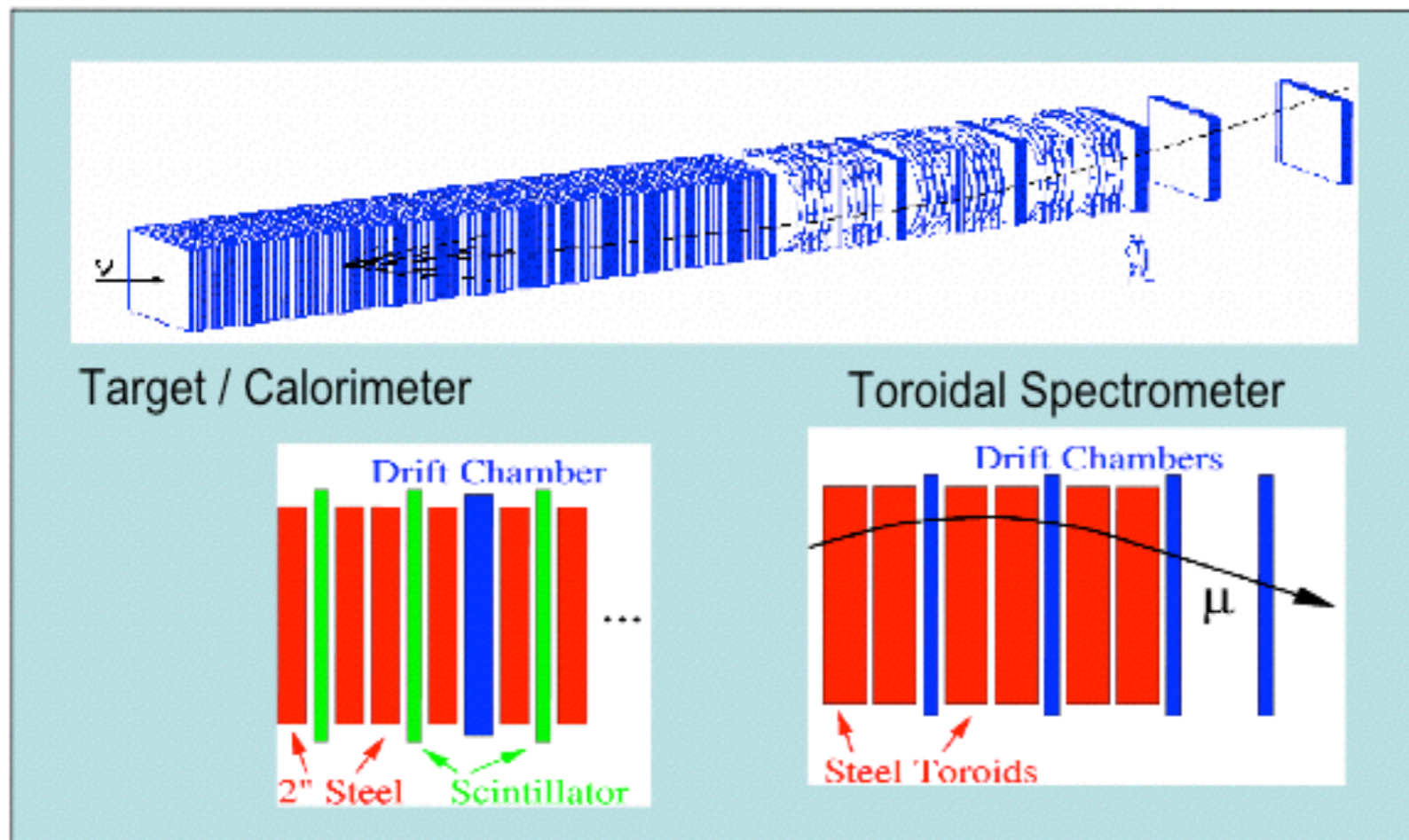
◆ PV Deep Inelastic Scattering: PVDIS

◆ theory robust for ^2H in valence quark region

◆ factor of 5 to 8 improvement possible: SOLID



NuTeV



Most precise measurement of neutrino-quark coupling

Subtle parton physics effects can affect the result

generated great interest in both nuclear and particle phenomenology

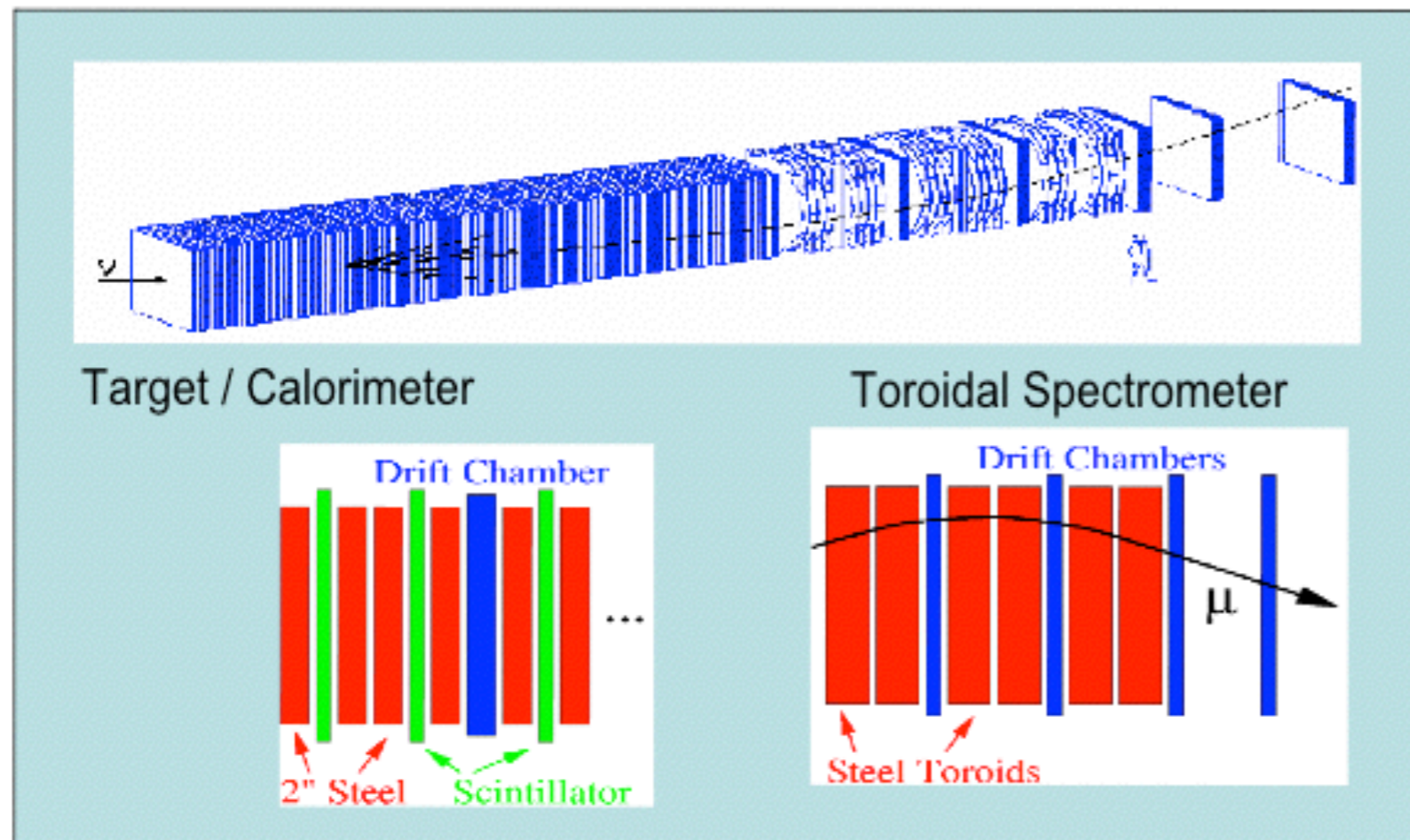
$$R^- = \frac{\sigma_{\nu N}^{NC} - \sigma_{\bar{\nu} N}^{NC}}{\sigma_{\nu N}^{CC} - \sigma_{\bar{\nu} N}^{CC}} \approx \rho^2 \left(\frac{1}{2} - \sin^2 \theta_W \right)$$

$$\sin^2 \theta_W^{(on-shell)} = 0.2277 \pm 0.0013(stat.) \pm 0.0009(syst.)$$

Standard Model prediction is 0.2227 (3σ deviation)

NuTeV

Significant discovery: but EW or QCD Physics?



Most precise measurement of neutrino-quark coupling

Subtle parton physics effects can affect the result

generated great interest in both nuclear and particle phenomenology

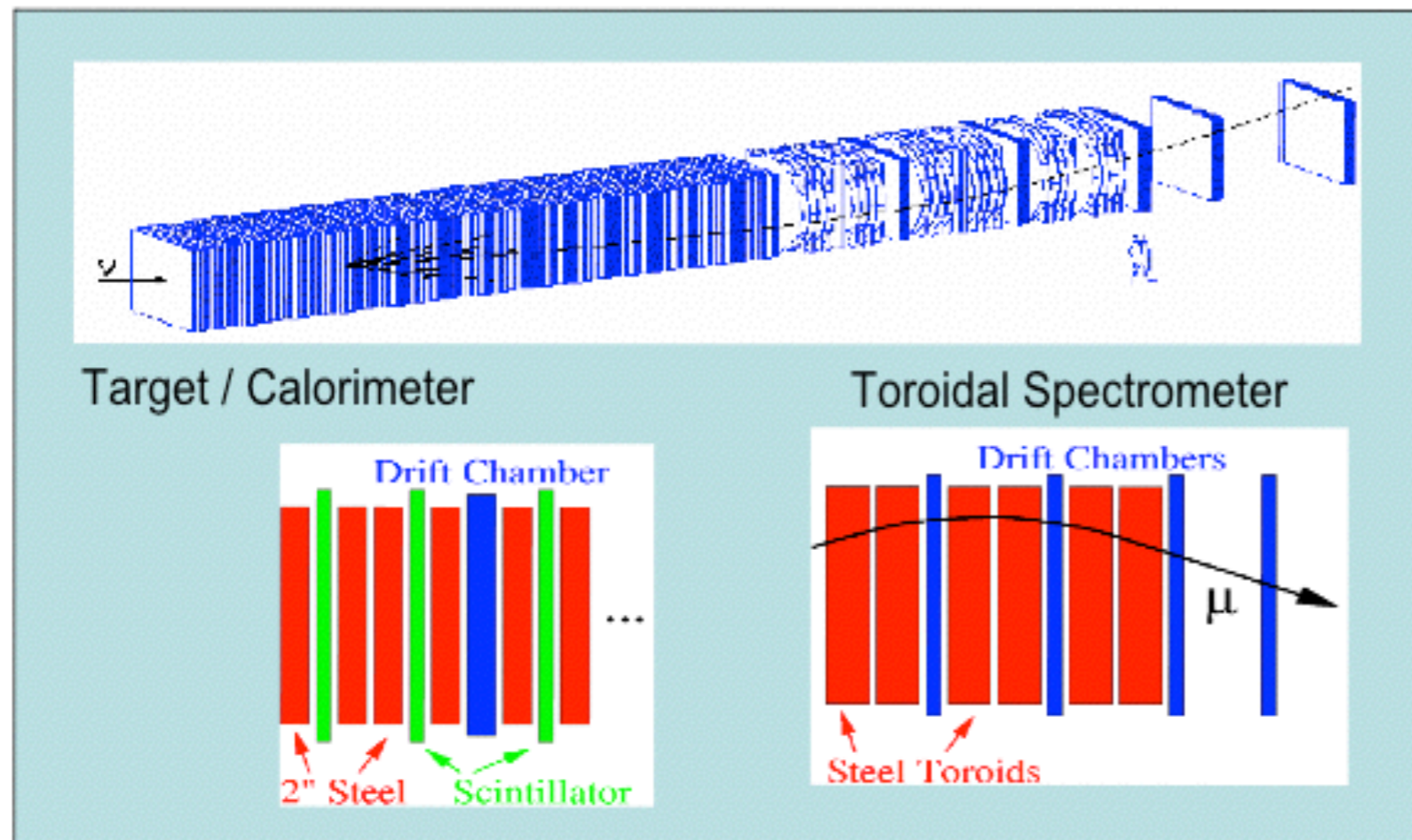
$$R^- = \frac{\sigma_{\nu N}^{NC} - \sigma_{\bar{\nu} N}^{NC}}{\sigma_{\nu N}^{CC} - \sigma_{\bar{\nu} N}^{CC}} \approx \rho^2 \left(\frac{1}{2} - \sin^2 \theta_W \right)$$

$$\sin^2 \theta_W^{(on-shell)} = 0.2277 \pm 0.0013(stat.) \pm 0.0009(syst.)$$

Standard Model prediction is 0.2227 (3σ deviation)

NuTeV

Significant discovery: but EW or QCD Physics?



Most precise measurement of neutrino-quark coupling

Subtle parton physics effects can affect the result

generated great interest in both nuclear and particle phenomenology

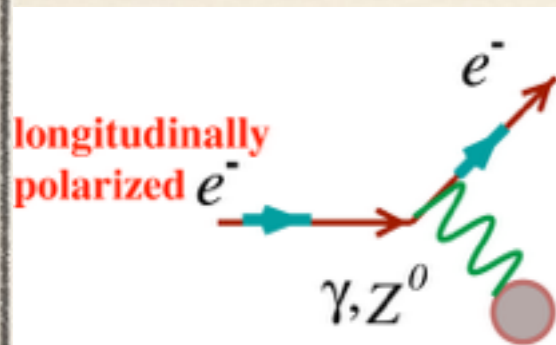
$$R^- = \frac{\sigma_{\nu N}^{NC} - \sigma_{\bar{\nu} N}^{NC}}{\sigma_{\nu N}^{CC} - \sigma_{\bar{\nu} N}^{CC}} \approx \rho^2 \left(\frac{1}{2} - \sin^2 \theta_W \right)$$

$$\sin^2 \theta_W^{(on-shell)} = 0.2277 \pm 0.0013(stat.) \pm 0.0009(syst.)$$

Standard Model prediction is 0.2227 (3 σ deviation)

4 decades of measurements: emergence as a **precision tool**

Parity-Violating Electron Scattering (PVES)



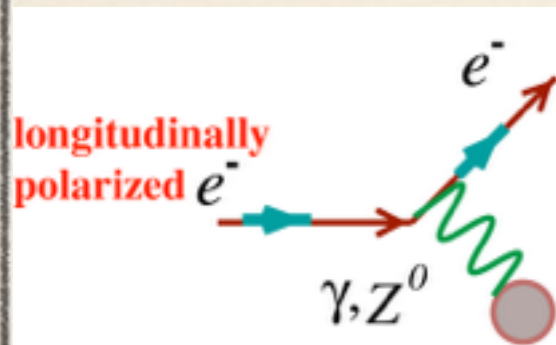
$$-A_{LR} = A_{PV} = \frac{\sigma_{\uparrow}^- - \sigma_{\downarrow}^-}{\sigma_{\uparrow}^- + \sigma_{\downarrow}^-} \sim \frac{A_{\text{weak}}}{A_{\gamma}} \sim \frac{G_F Q^2}{4\pi\alpha} (g_A^e g_V^T + \beta g_V^e g_A^T)$$

g_V is a function of $\sin^2\theta_W$

Weak Charge Q_w

4 decades of measurements: emergence as a **precision tool**

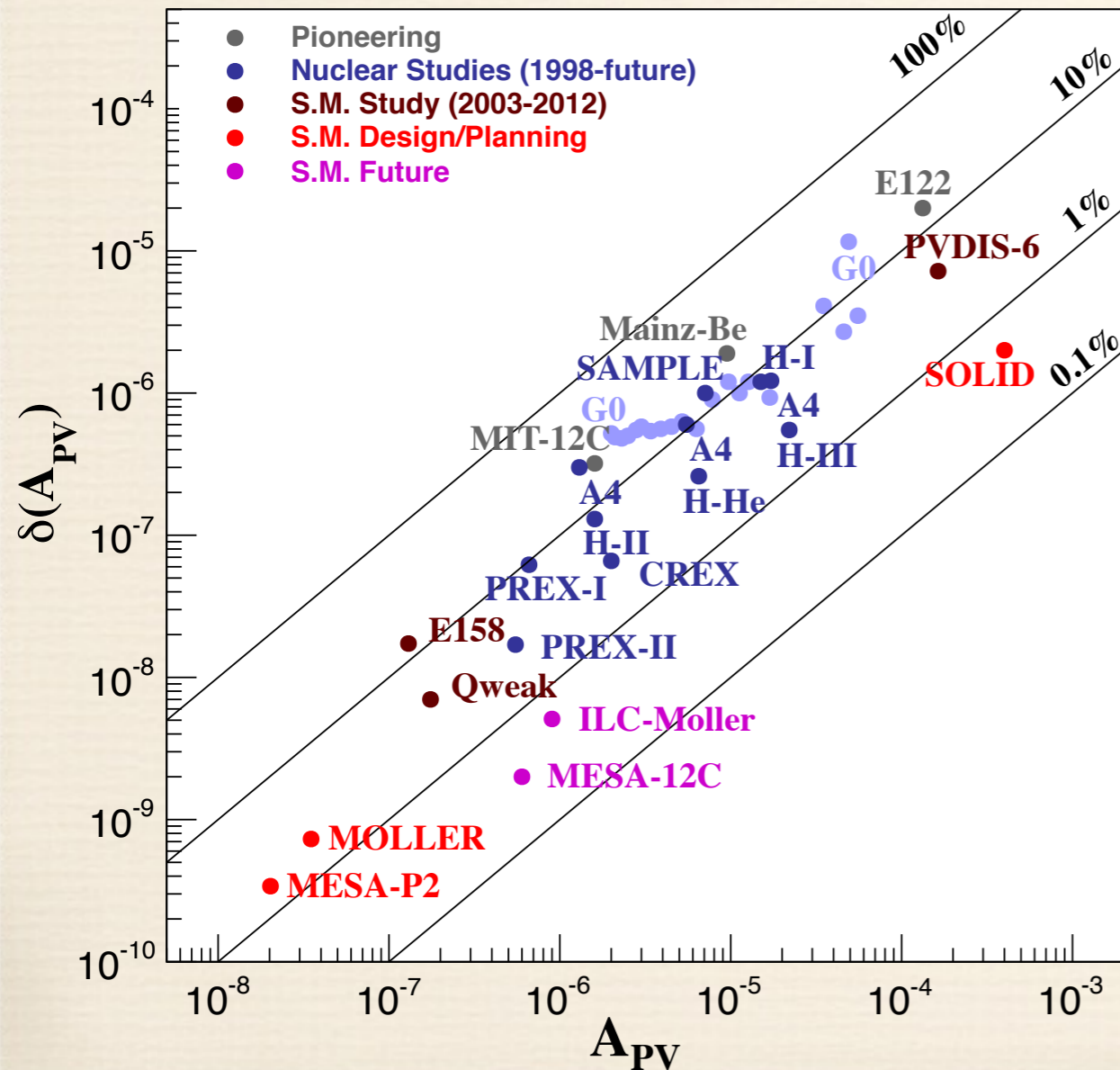
Parity-Violating Electron Scattering (PVES)



$$-A_{LR} = A_{PV} = \frac{\sigma_{\uparrow} - \sigma_{\downarrow}}{\sigma_{\uparrow} + \sigma_{\downarrow}} \sim \frac{A_{\text{weak}}}{A_{\gamma}} \sim \frac{G_F Q^2}{4\pi\alpha} (g_A^e g_V^T + \beta g_V^e g_A^T)$$

g_V is a function of sin²θ_w

Weak Charge Q_w



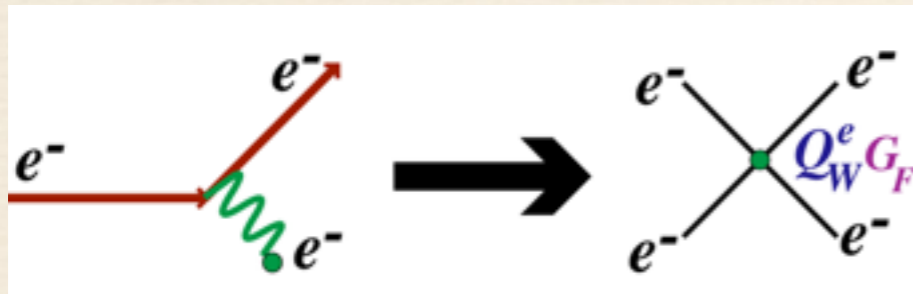
Variety of Physics Topics:
continuous interplay between
hadron physics and electroweak
physics

*Steady improvements in
accelerator and detector
technology*

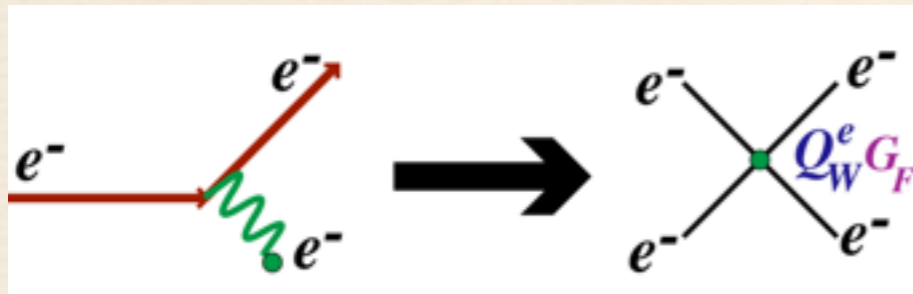
State of the Art

- sub-part per billion statistical reach and systematic control
- **sub-1% normalization control**

PV Electron-Electron Scattering



PV Electron-Electron Scattering

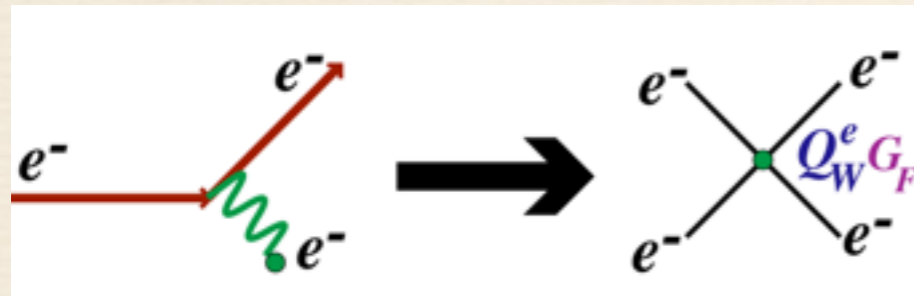


electron target:

$$Q_W = 1 - 4 \sin^2 \theta_W$$

$$\frac{\delta(Q_W)}{Q_W} \sim 10\% \implies \frac{\delta(\sin^2 \theta_W)}{\sin^2 \theta_W} \sim 0.5\%$$

PV Electron-Electron Scattering



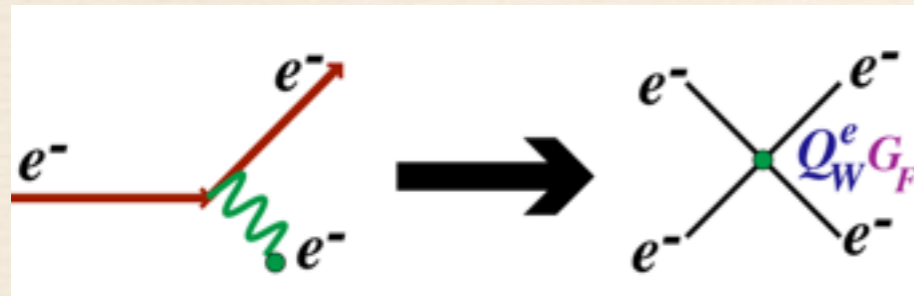
$$+ \text{[Crossed lines diagram]} \frac{1}{\Lambda^2} \mathcal{L}_6$$

electron target:

$$Q_W = 1 - 4 \sin^2 \theta_W$$

$$\frac{\delta(Q_W)}{Q_W} \sim 10\% \implies \frac{\delta(\sin^2 \theta_W)}{\sin^2 \theta_W} \sim 0.5\%$$

PV Electron-Electron Scattering

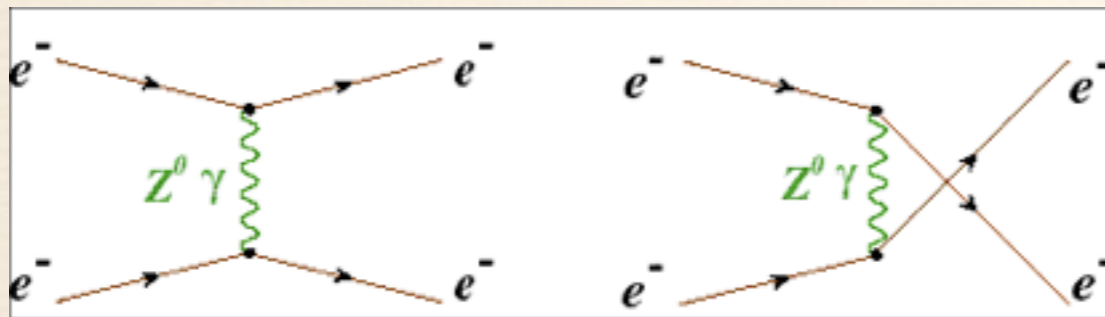


$$+ \text{[Crossed diagram]} \frac{1}{\Lambda^2} \mathcal{L}_6$$

electron target:

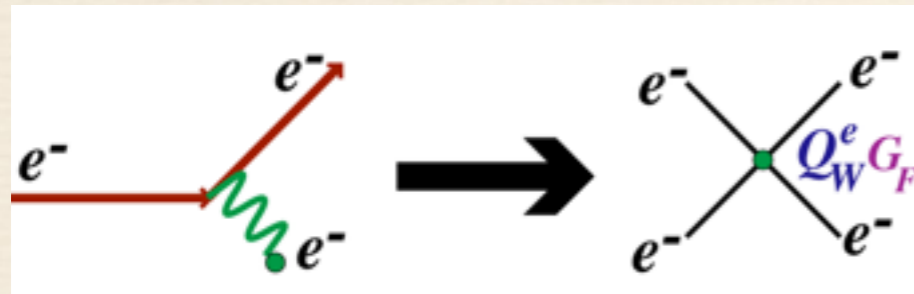
$$Q_W = 1 - 4 \sin^2 \theta_W$$

$$\frac{\delta(Q_W)}{Q_W} \sim 10\% \implies \frac{\delta(\sin^2 \theta_W)}{\sin^2 \theta_W} \sim 0.5\%$$



$$|A_\gamma + A_Z + A_{\text{new}}|^2 \rightarrow A_\gamma^2 \left[1 + 2 \left(\frac{A_Z}{A_\gamma} \right) + 2 \left(\frac{A_{\text{new}}}{A_\gamma} \right) \right]$$

PV Electron-Electron Scattering

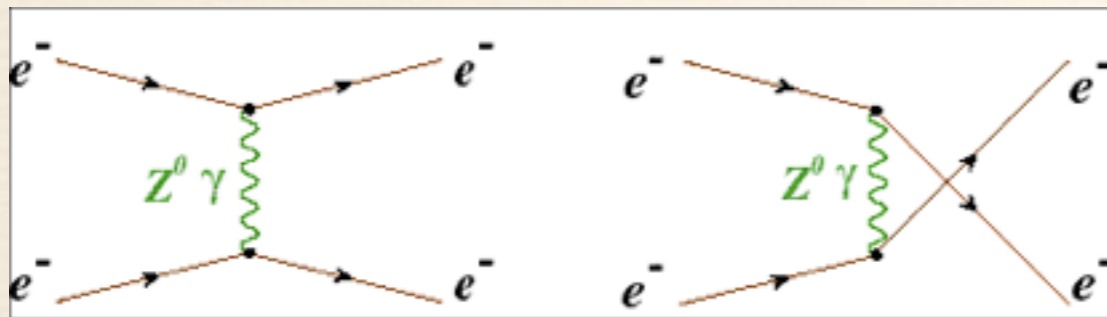


$$+ \text{[Crossed lines diagram]} \frac{1}{\Lambda^2} \mathcal{L}_6$$

electron target:

$$Q_W = 1 - 4 \sin^2 \theta_W$$

$$\frac{\delta(Q_W)}{Q_W} \sim 10\% \implies \frac{\delta(\sin^2 \theta_W)}{\sin^2 \theta_W} \sim 0.5\%$$

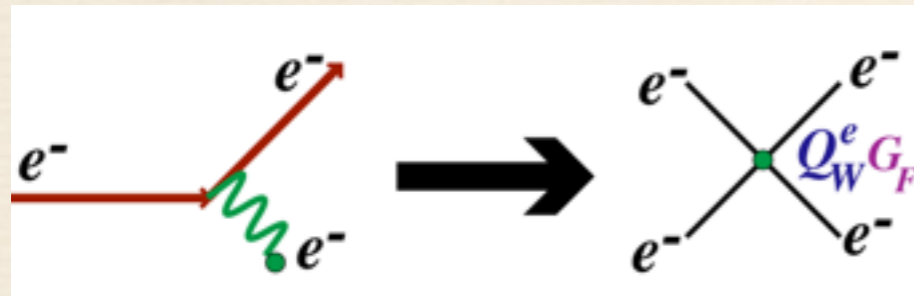


$$|A_\gamma + A_Z + A_{\text{new}}|^2 \rightarrow A_\gamma^2 \left[1 + 2 \left(\frac{A_Z}{A_\gamma} \right) + 2 \left(\frac{A_{\text{new}}}{A_\gamma} \right) \right]$$

$$A_{PV} \approx 8 \times 10^{-8} E_{\text{beam}} (1 - 4 \sin^2 \theta_W)$$

➡ **Tiny!**

PV Electron-Electron Scattering

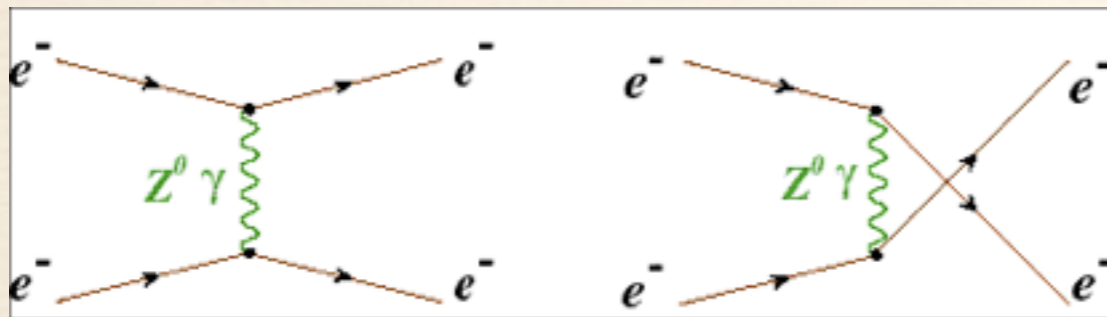


$$+ \text{[t-channel diagram]} \frac{1}{\Lambda^2} \mathcal{L}_6$$

electron target:

$$Q_W = 1 - 4 \sin^2 \theta_W$$

$$\frac{\delta(Q_W)}{Q_W} \sim 10\% \implies \frac{\delta(\sin^2 \theta_W)}{\sin^2 \theta_W} \sim 0.5\%$$



$$|A_\gamma + A_Z + A_{\text{new}}|^2 \rightarrow A_\gamma^2 \left[1 + 2 \left(\frac{A_Z}{A_\gamma} \right) + 2 \left(\frac{A_{\text{new}}}{A_\gamma} \right) \right]$$

$$A_{PV} \approx 8 \times 10^{-8} E_{\text{beam}} (1 - 4 \sin^2 \theta_W)$$

Tiny!



45 & 48 GeV Beam
85% longitudinal polarization

4-7 mrad

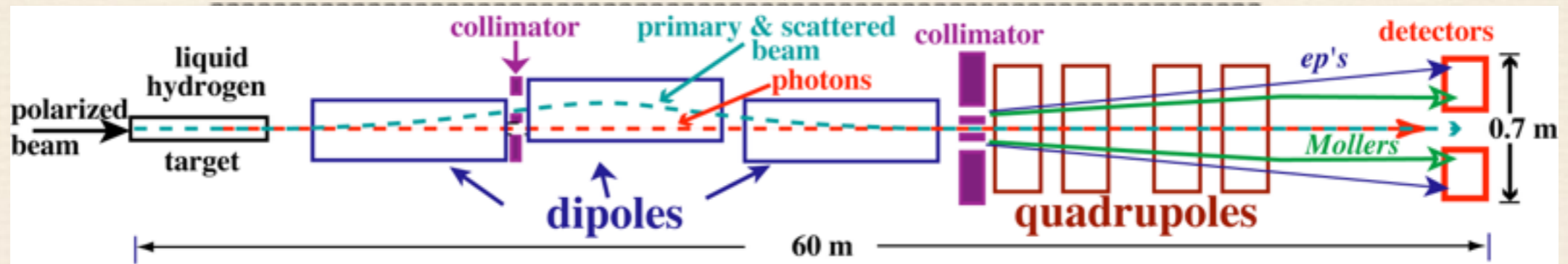
SLAC E158: 1997-2004

End Station A at SLAC

Goal: error small enough to probe TeV scale physics

SLAC E158

~ 10 ppb statistical error at highest E_{beam} , ~ 0.5% error on weak mixing angle



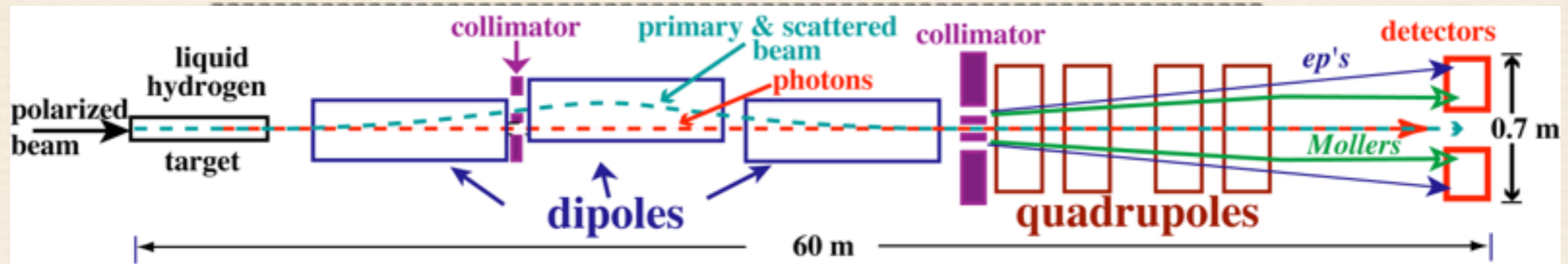
A large number of technical challenges



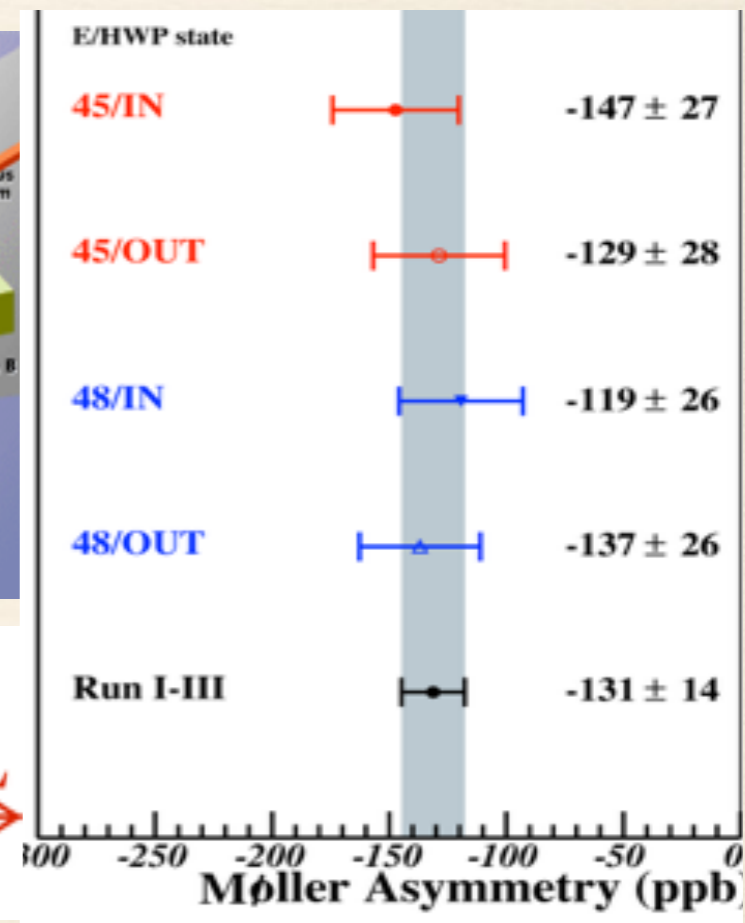
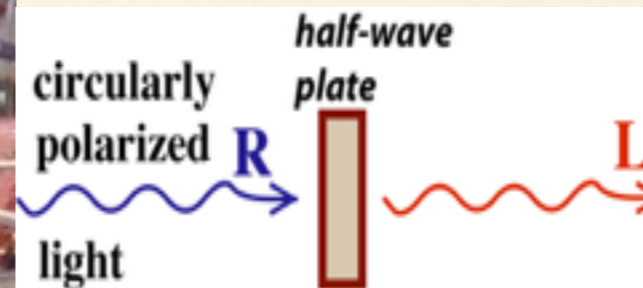
Goal: error small enough to probe TeV scale physics

SLAC E158

~ 10 ppb statistical error at highest E_{beam} , $\sim 0.5\%$ error on weak mixing angle



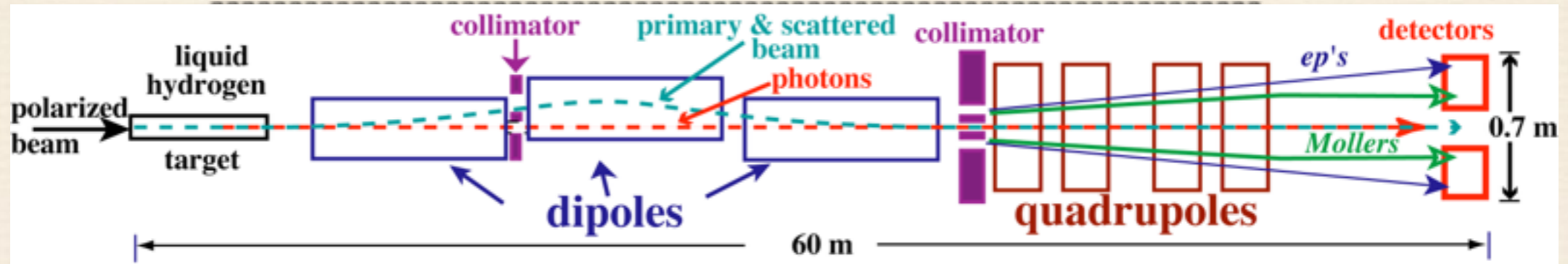
A large number of technical challenges



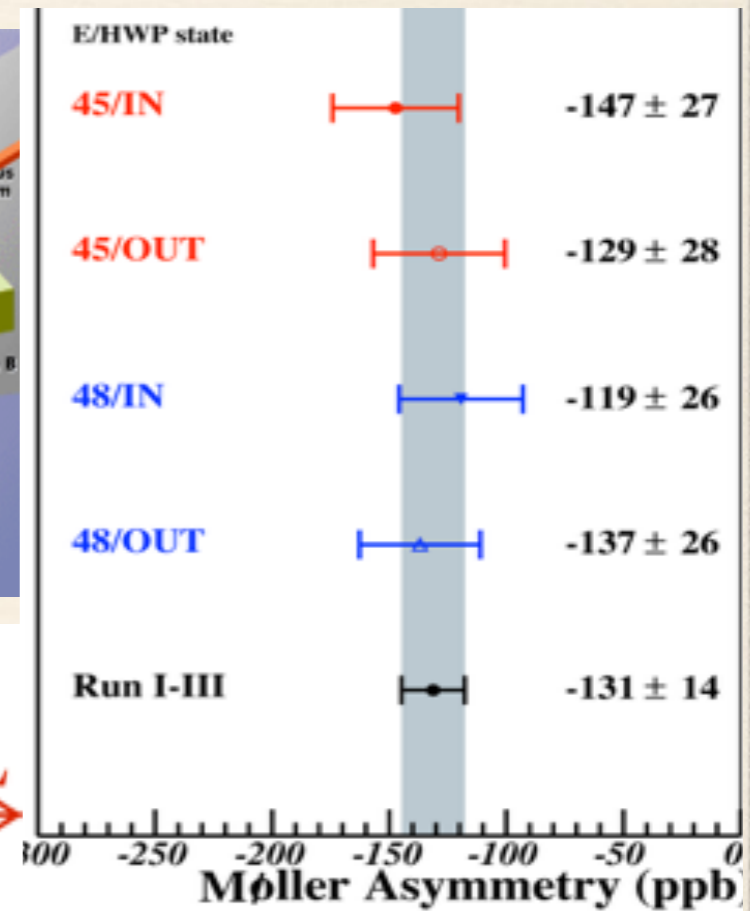
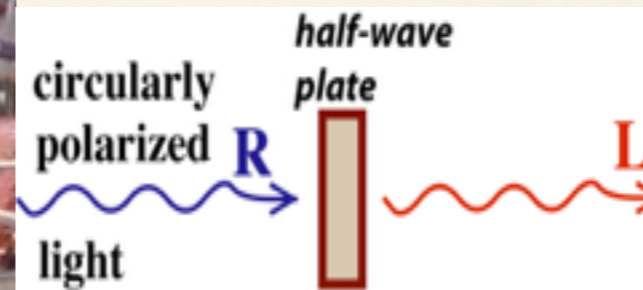
Goal: error small enough to probe TeV scale physics

SLAC E158

~ 10 ppb statistical error at highest E_{beam} , ~ 0.5% error on weak mixing angle



A large number of technical challenges



$$A_{PV} = (-131 \pm 14 \pm 10) \times 10^{-9}$$

Phys. Rev. Lett. **95** 081601 (2005)

Tree-level prediction: ~ 250 ppb

E158 Implications

Tree-level prediction: ~ 250 ppb

$$A_{PV} = (-131 \pm 14 \pm 10) \times 10^{-9}$$

E158 Implications

Final E158 Result

Phys. Rev. Lett. **95** 081601 (2005)

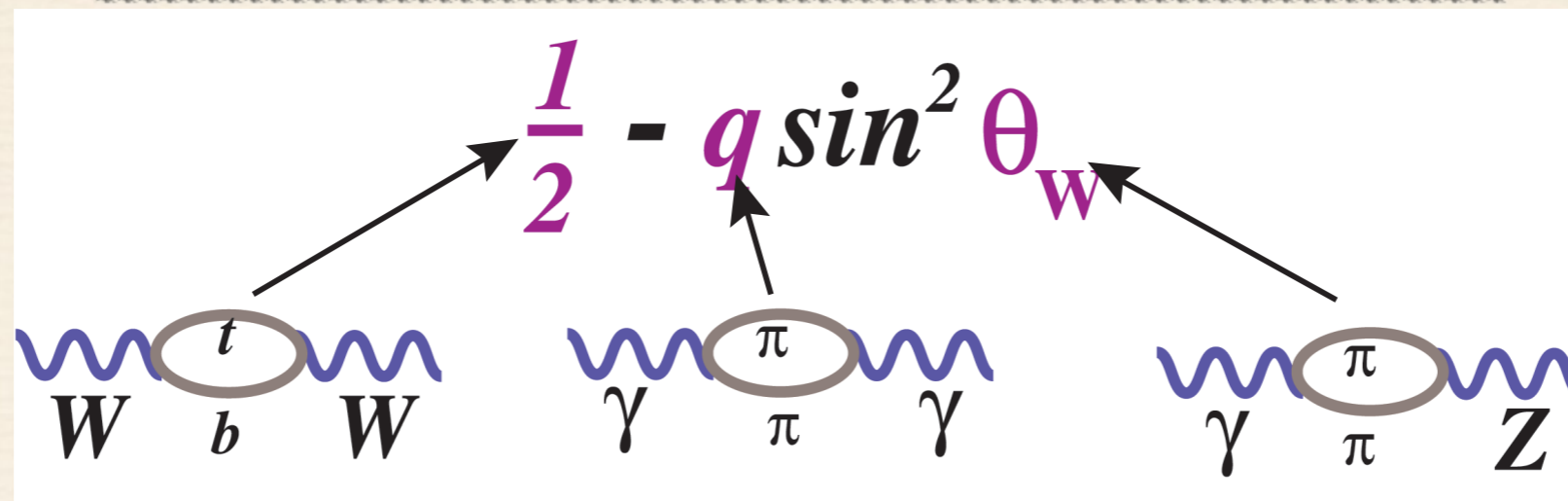
Tree-level prediction: ~ 250 ppb

$$A_{PV} = (-131 \pm 14 \pm 10) \times 10^{-9}$$

E158 Implications

Final E158 Result

Phys. Rev. Lett. **95** 081601 (2005)



Tree-level prediction: ~ 250 ppb

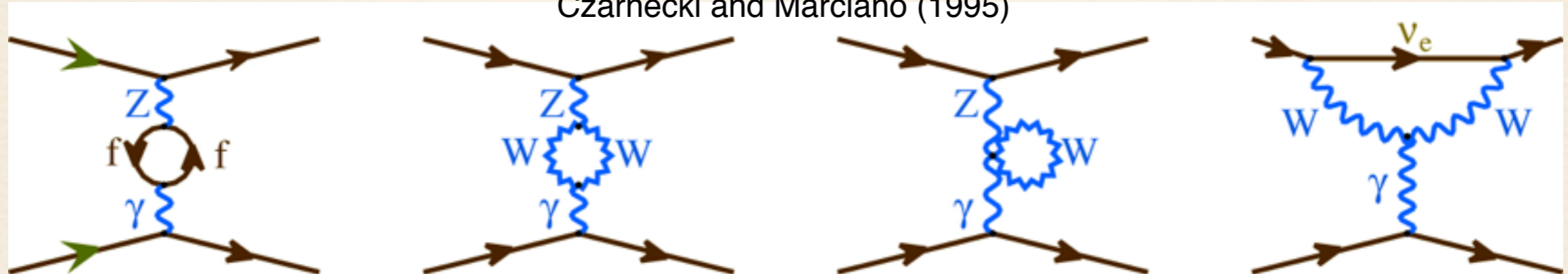
$$A_{PV} = (-131 \pm 14 \pm 10) \times 10^{-9}$$

E158 Implications

Final E158 Result

Phys. Rev. Lett. **95** 081601 (2005)

Czarnecki and Marciano (1995)



Tree-level prediction: ~ 250 ppb

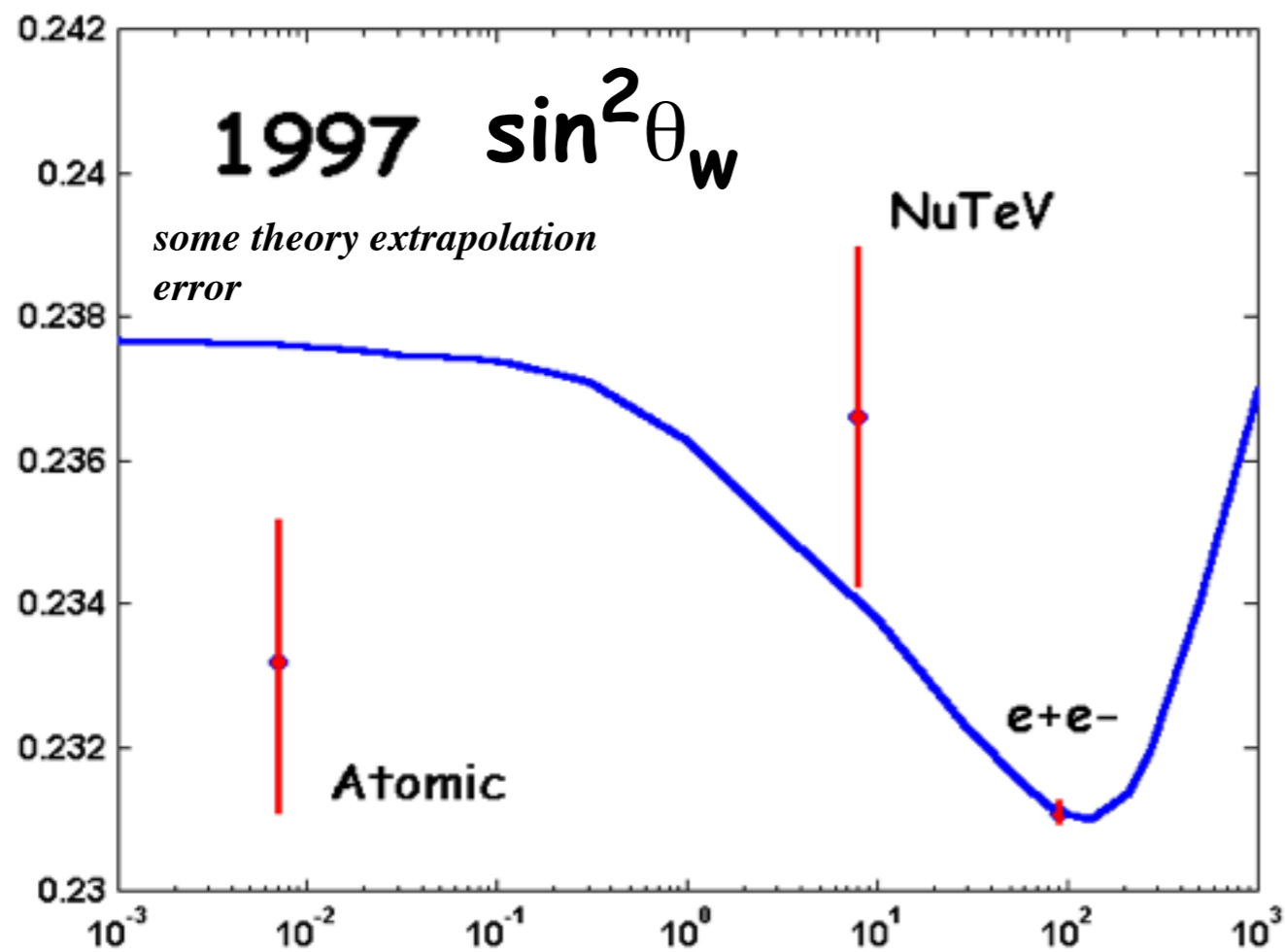
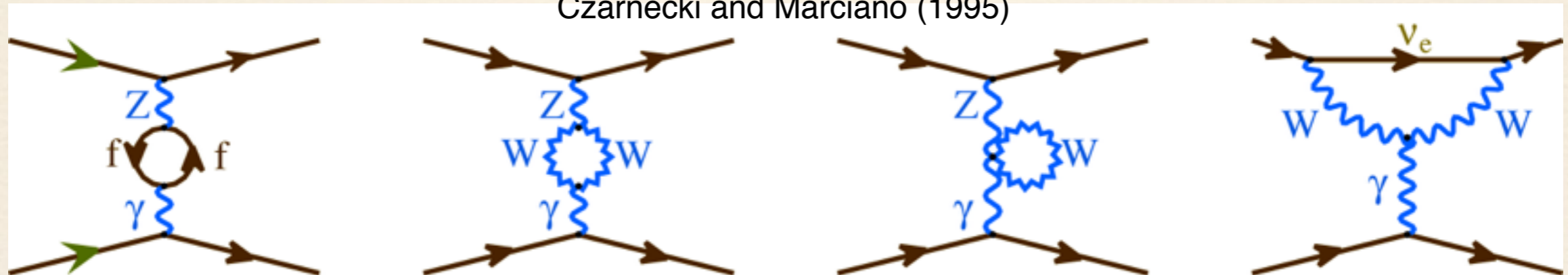
$$A_{PV} = (-131 \pm 14 \pm 10) \times 10^{-9}$$

E158 Implications

Final E158 Result

Phys. Rev. Lett. **95** 081601 (2005)

Czarnecki and Marciano (1995)



Tree-level prediction: ~ 250 ppb

$$A_{PV} = (-131 \pm 14 \pm 10) \times 10^{-9}$$

E158 Implications

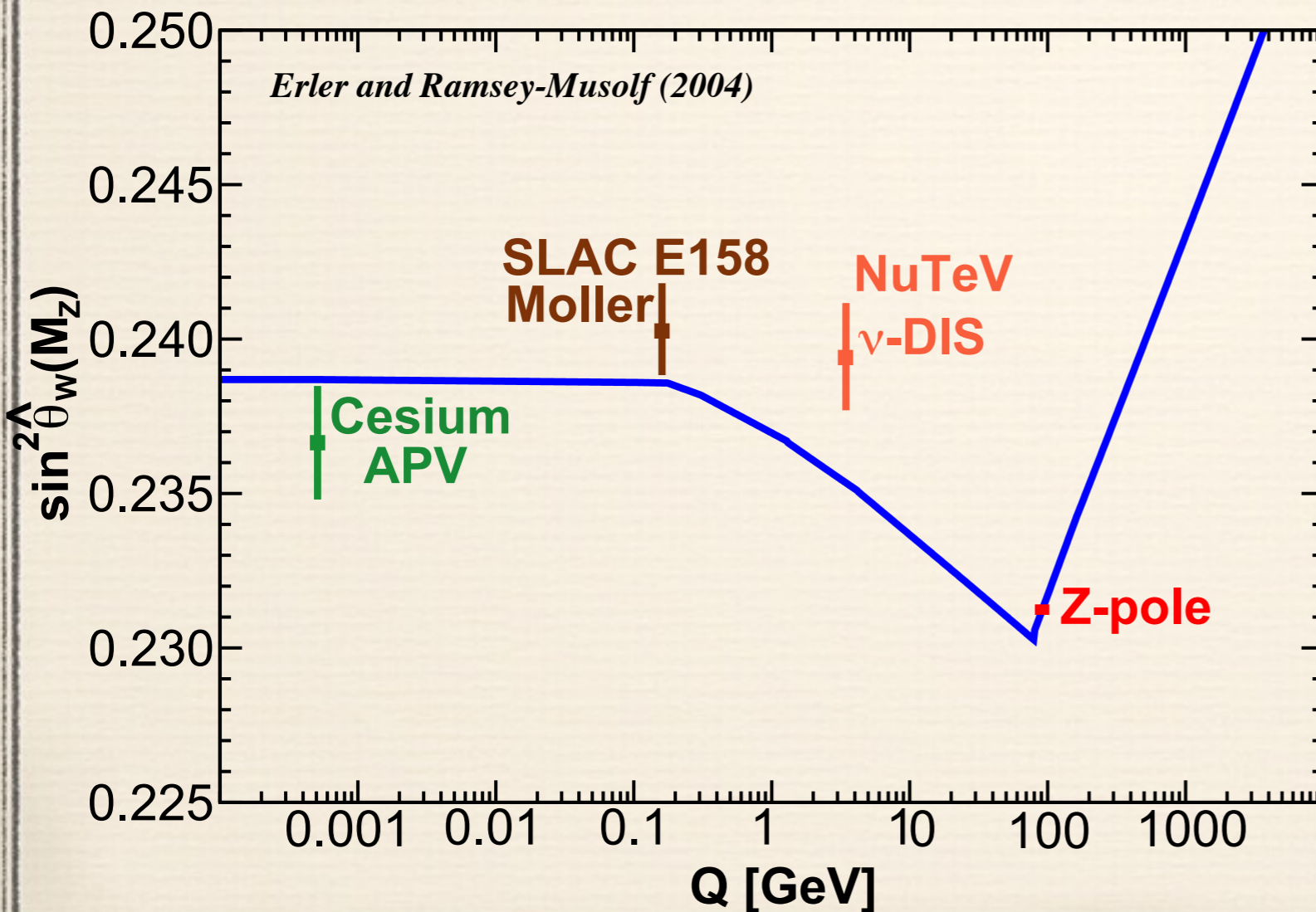
Final E158 Result

Phys. Rev. Lett. **95** 081601 (2005)

Czarnecki and Marciano (1995)



Erler and Ramsey-Musolf (2004)



Tree-level prediction: ~ 250 ppb

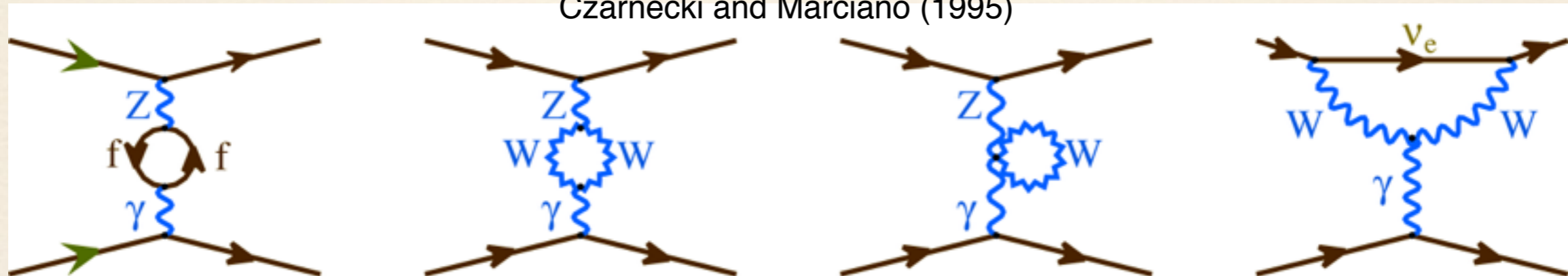
$$A_{PV} = (-131 \pm 14 \pm 10) \times 10^{-9}$$

E158 Implications

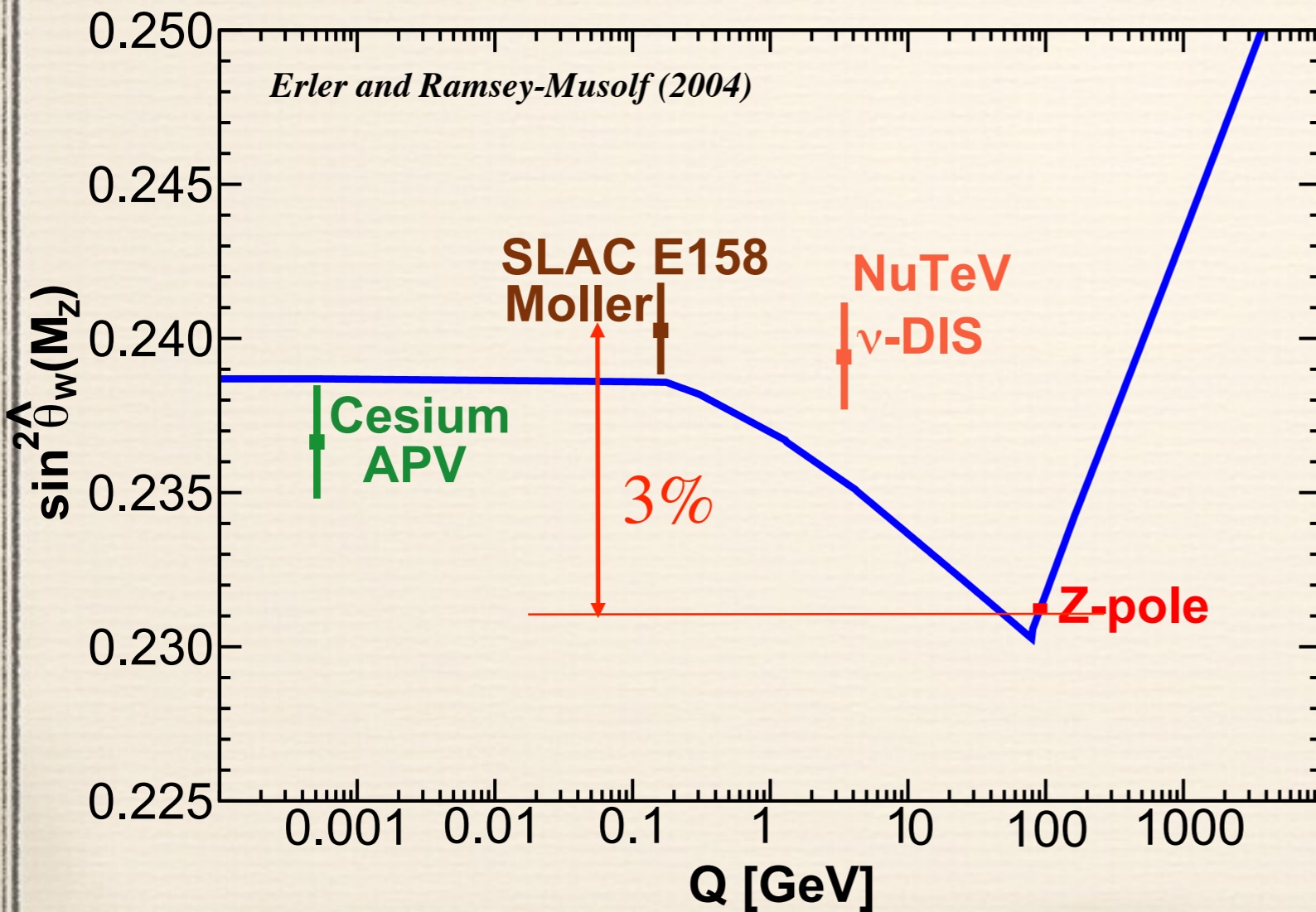
Final E158 Result

Phys. Rev. Lett. **95** 081601 (2005)

Czarnecki and Marciano (1995)



Erlar and Ramsey-Musolf (2004)



Tree-level prediction: ~ 250 ppb

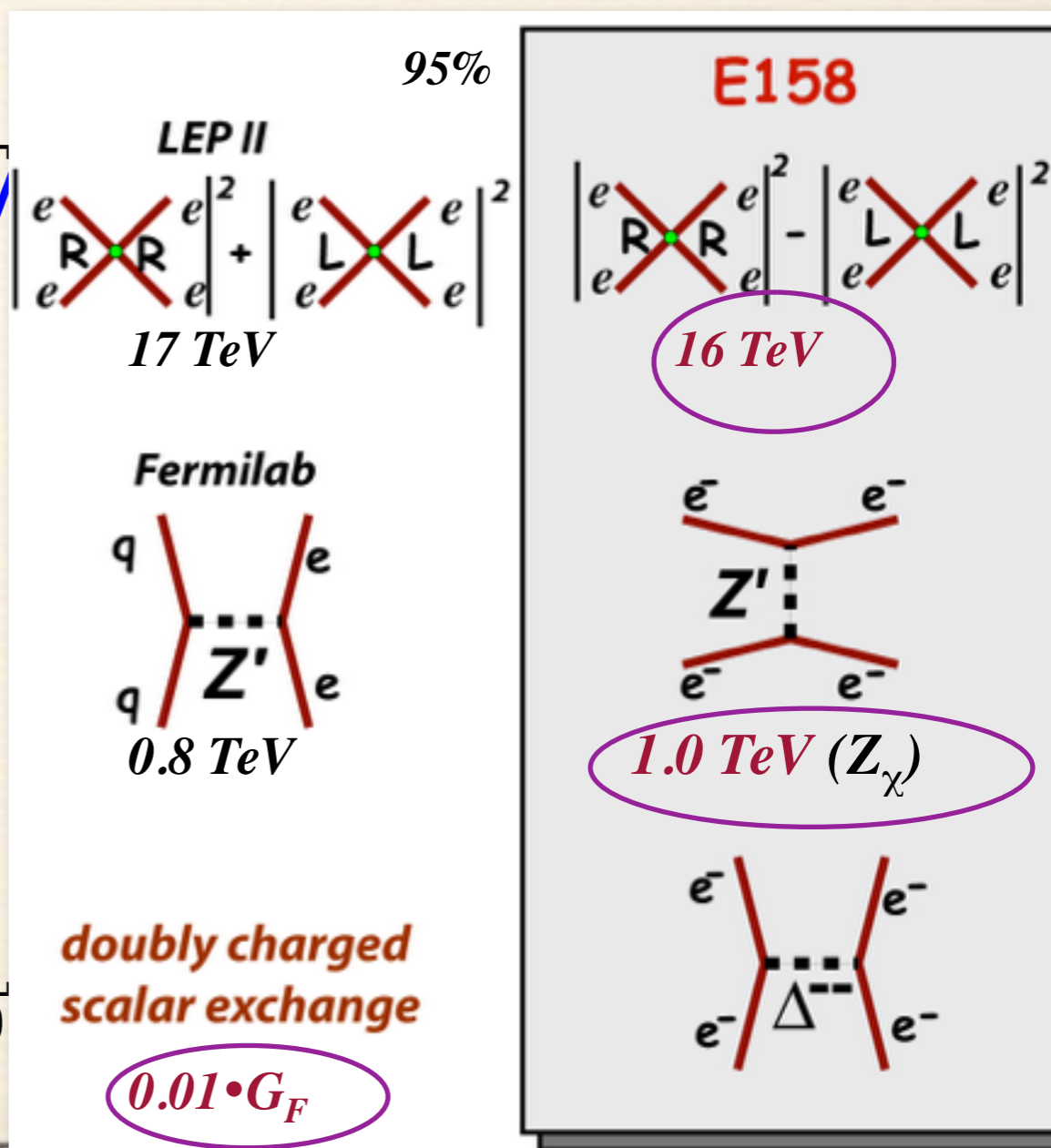
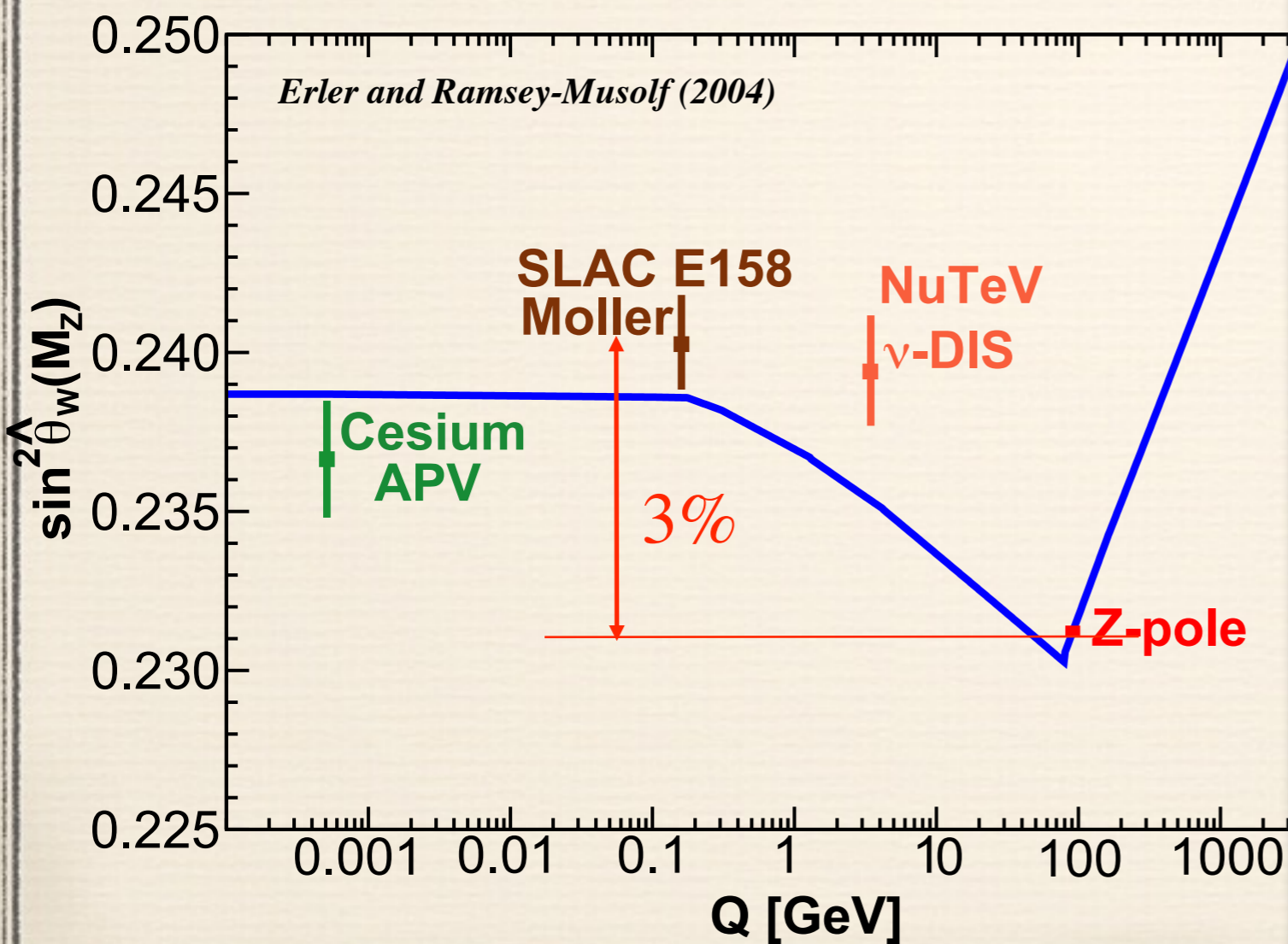
$$A_{PV} = (-131 \pm 14 \pm 10) \times 10^{-9}$$

Final E158 Result

E158 Implications

Phys. Rev. Lett. **95** 081601 (2005)

Limits on "New" Physics



$$A_{PV} = (-131 \pm 14 \pm 10) \times 10^{-9}$$

Nature

Vol 435 26 May 2005

E158 Implications

Final E158 Result

NEWS AND VIEWS

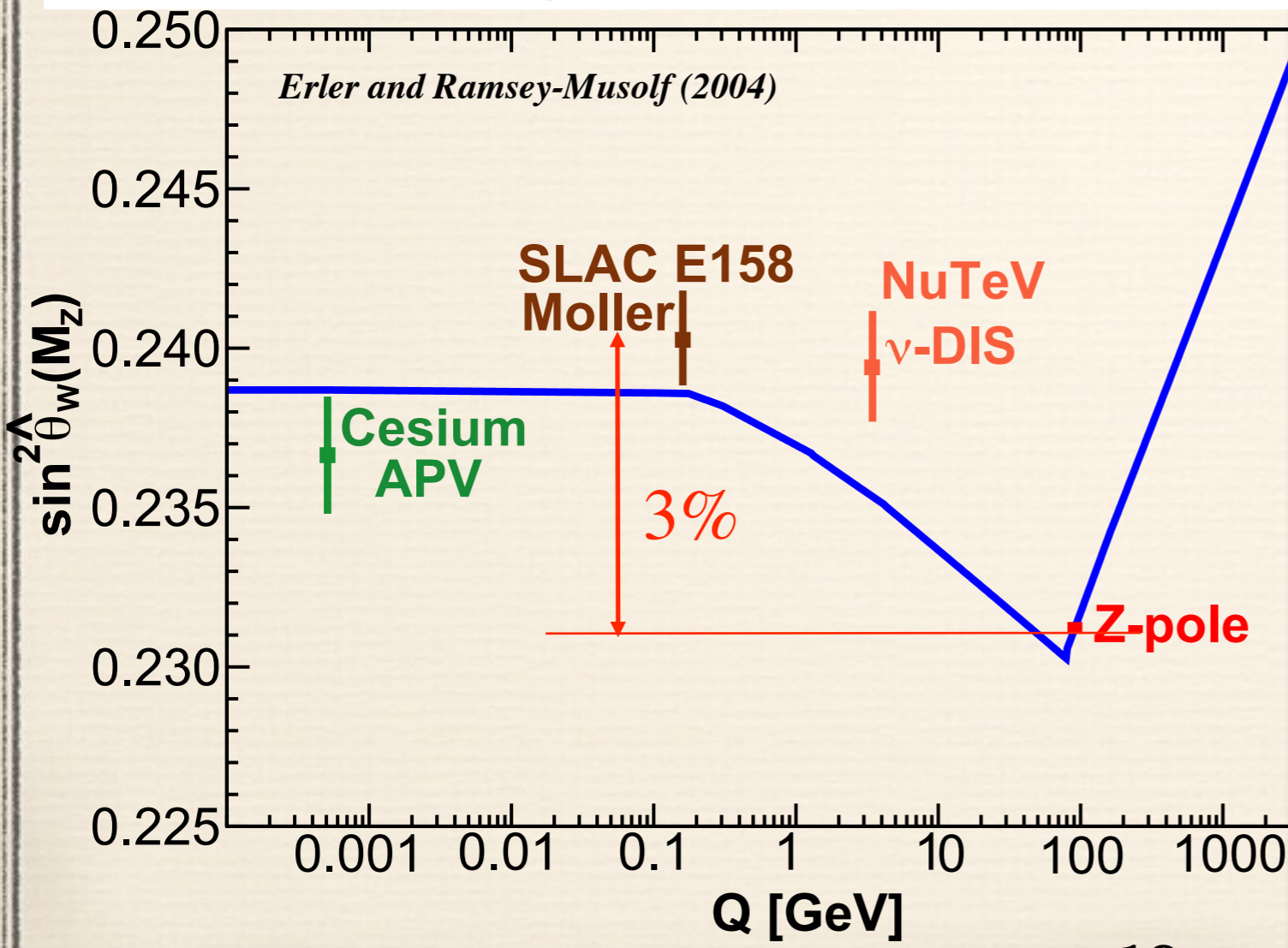
Phys. Rev. Lett. 95 081601 (2005)

PARTICLE PHYSICS

Electrons are not ambidextrous

Andrzej Czarnecki and William J. Marciano

The best low-energy measurement yet obtained of the electroweak mixing angle — a central parameter of the standard model of particle physics — is the last hurrah for Stanford's powerful two-mile linear accelerator.



Limits on "New" Physics

95%

LEP II
 $|e_R e_R|^2 + |e_L e_L|^2$
 17 TeV

Fermilab
 $q q \rightarrow Z' e e$
 0.8 TeV

E158
 $|e_R e_R|^2 - |e_L e_L|^2$
 16 TeV

$e^- e^- \rightarrow Z' e e$
 1.0 TeV (Z_χ)

$e^- e^- \rightarrow \Delta e e$
 doubly charged scalar exchange
 0.01 $\cdot G_F$

$$A_{PV} = (-131 \pm 14 \pm 10) \times 10^{-9}$$

Nature

Vol 435 26 May 2005

E158 Implications

Final E158 Result

NEWS AND VIEWS

Phys. Rev. Lett. 95 081601 (2005)

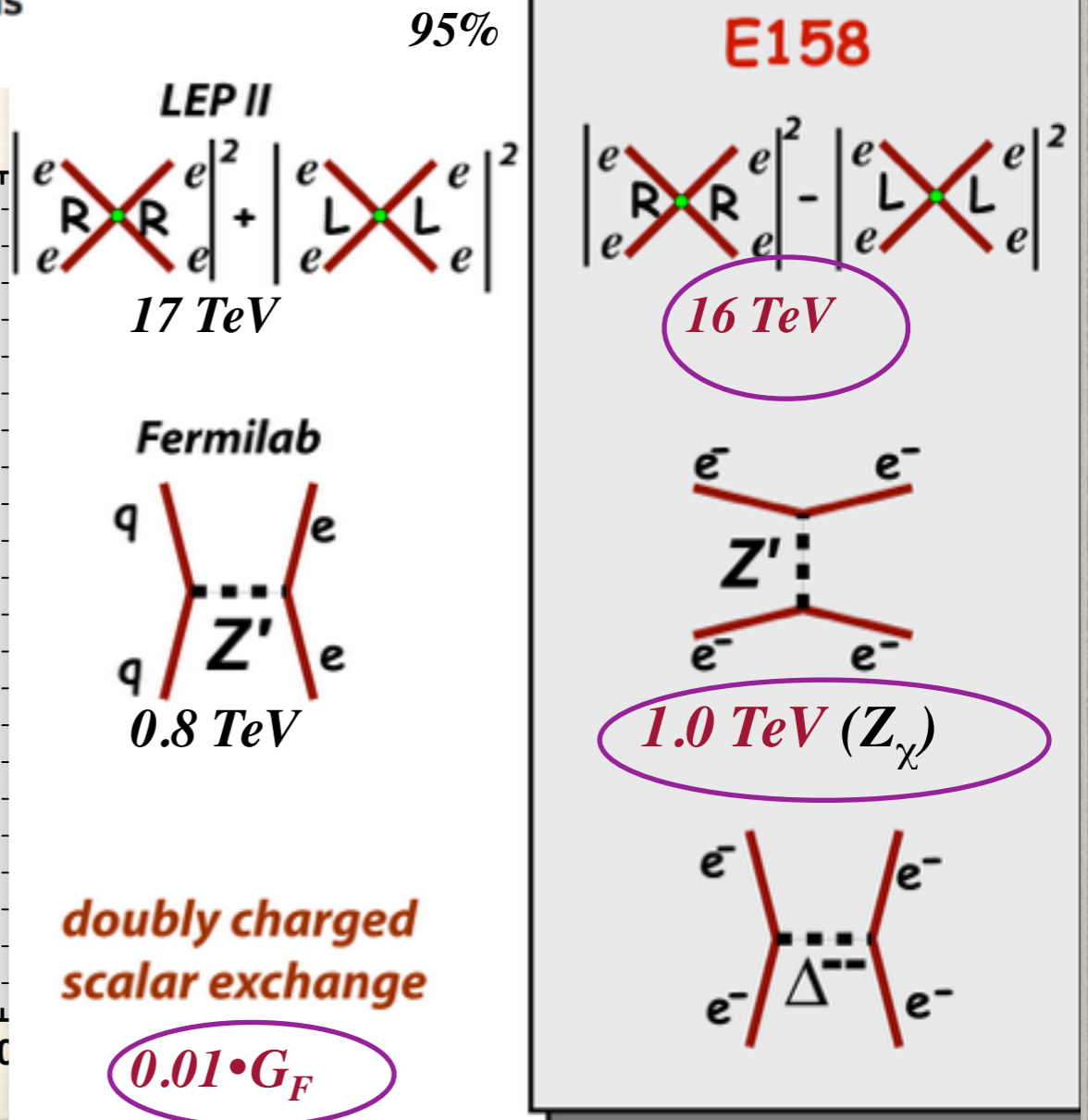
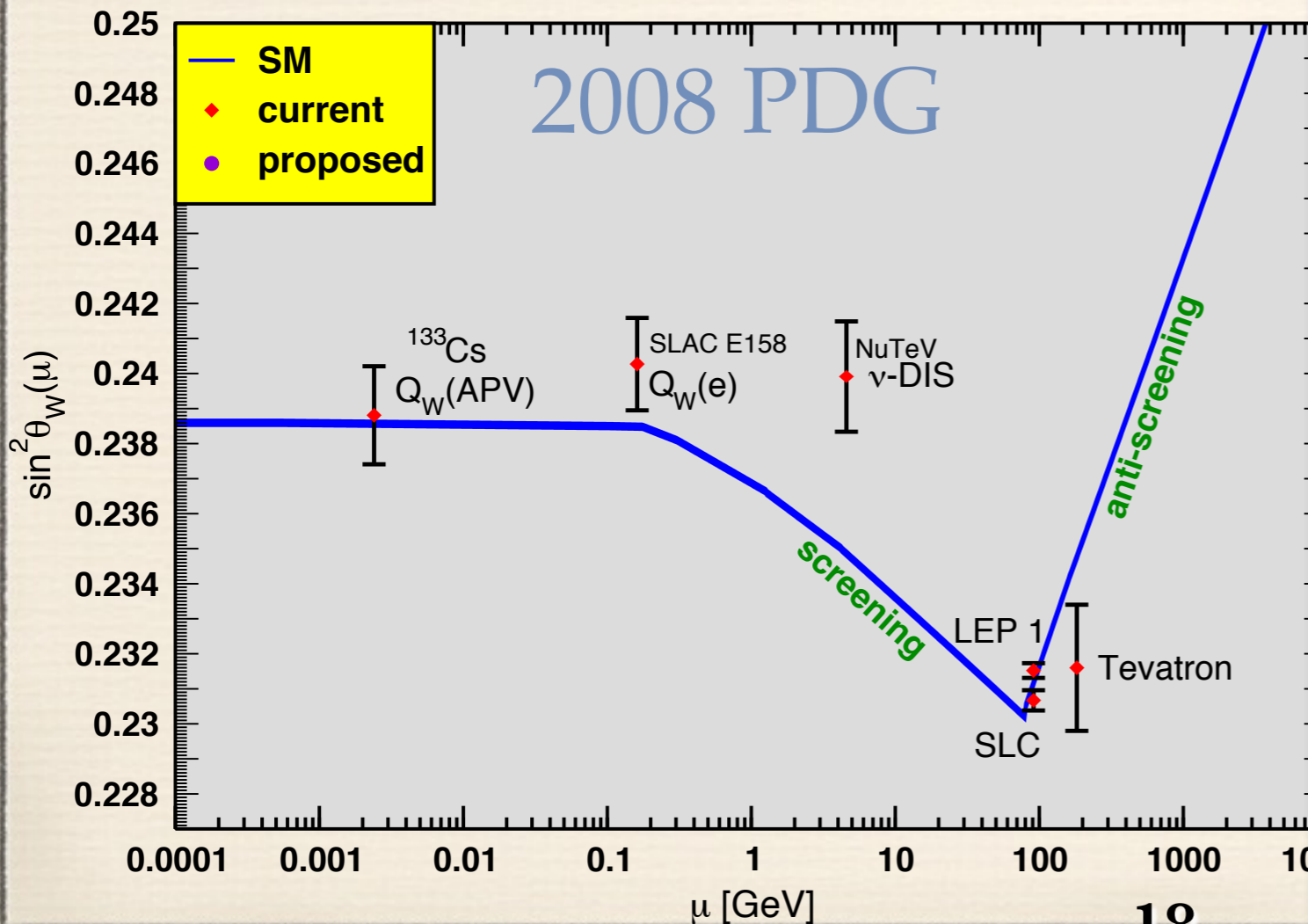
PARTICLE PHYSICS

Electrons are not ambidextrous

Andrzej Czarnecki and William J. Marciano

The best low-energy measurement yet obtained of the electroweak mixing angle — a central parameter of the standard model of particle physics — is the last hurrah for Stanford's powerful two-mile linear accelerator.

Limits on "New" Physics



$$A_{PV} = (-131 \pm 14 \pm 10) \times 10^{-9}$$

Nature

Vol 435 26 May 2005

E158 Implications

Final E158 Result

NEWS AND VIEWS

Phys. Rev. Lett. 95 081601 (2005)

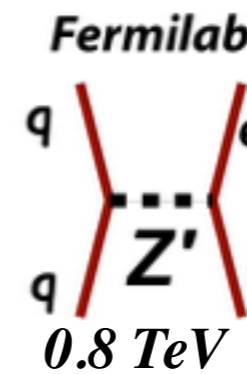
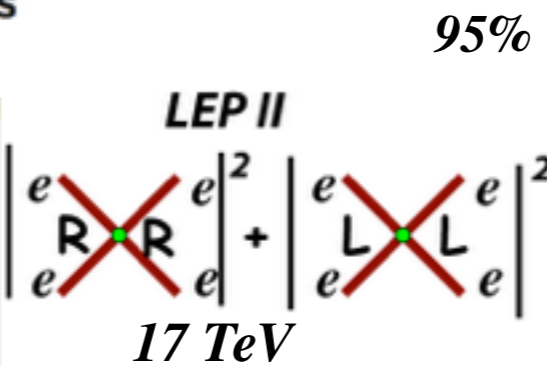
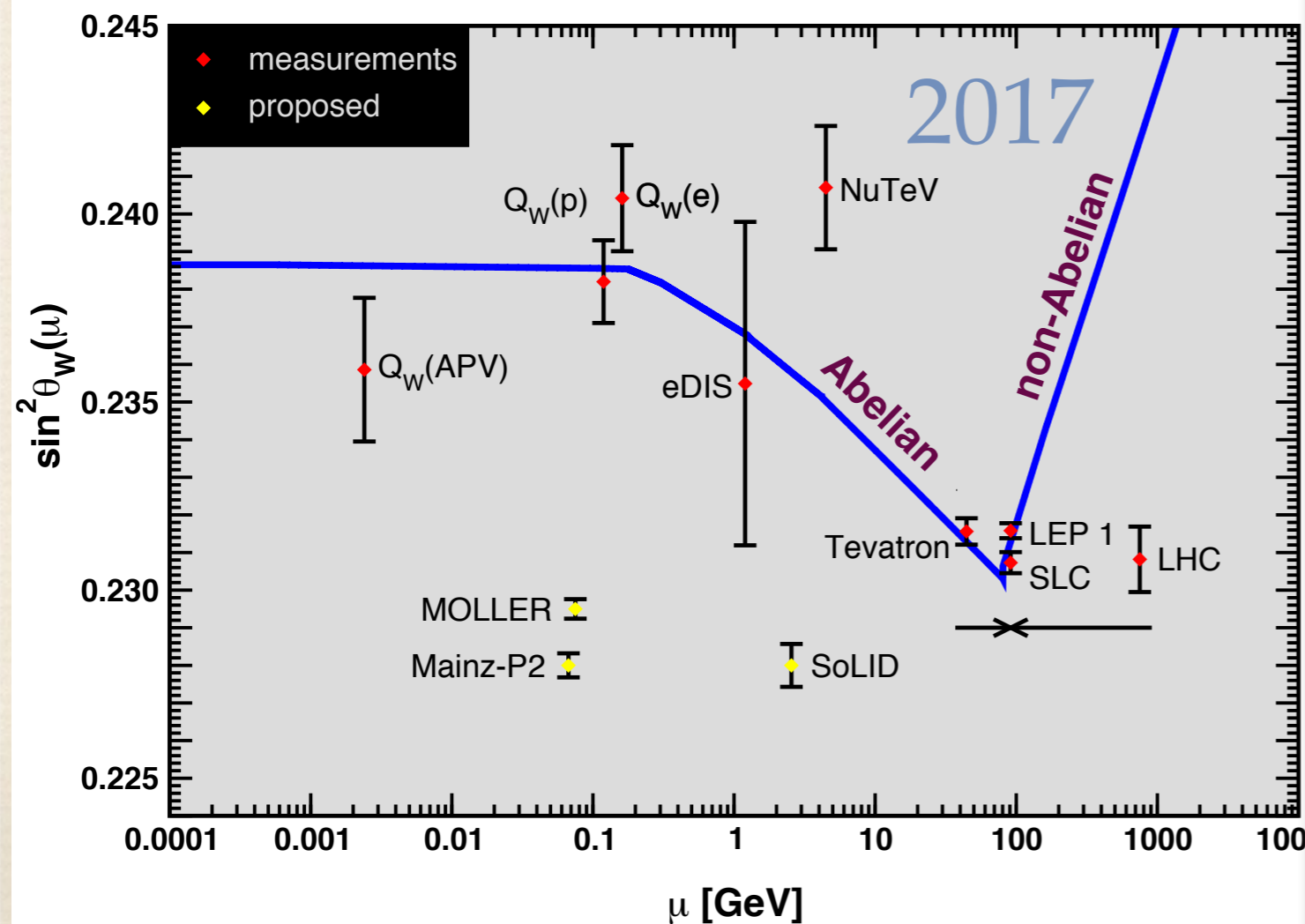
PARTICLE PHYSICS

Electrons are not ambidextrous

Andrzej Czarnecki and William J. Marciano

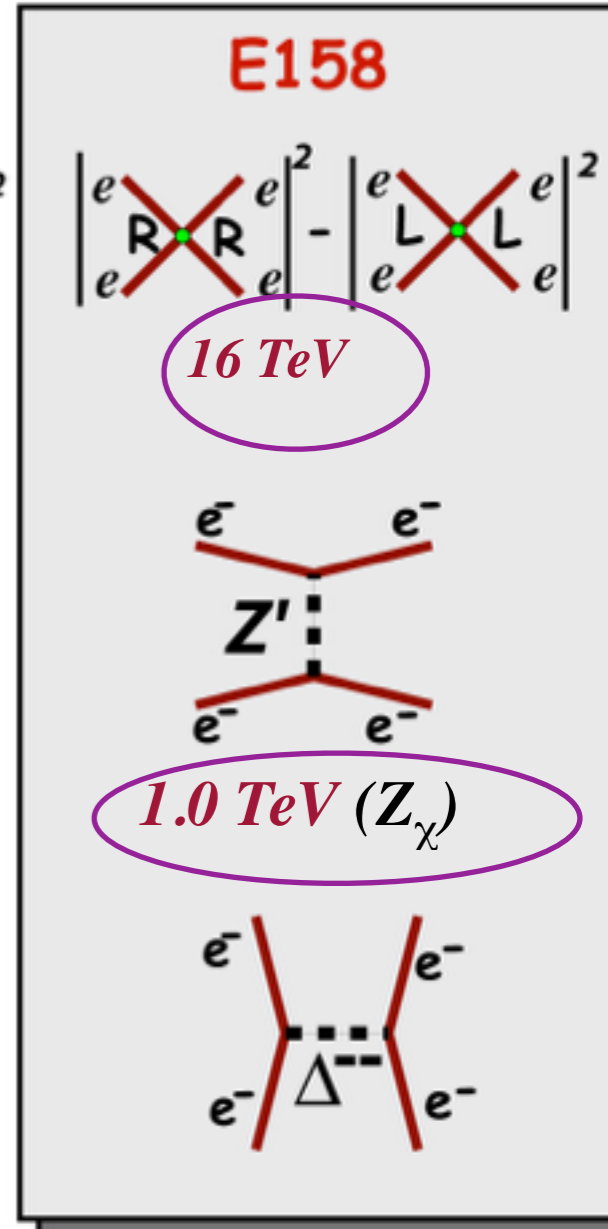
The best low-energy measurement yet obtained of the electroweak mixing angle — a central parameter of the standard model of particle physics — is the last hurrah for Stanford's powerful two-mile linear accelerator.

Limits on "New" Physics



doubly charged scalar exchange

$$0.01 \cdot G_F$$



$$A_{PV} = (-131 \pm 14 \pm 10) \times 10^{-9}$$

Nature

Vol 435 26 May 2005

E158 Implications

Final E158 Result

NEWS AND VIEWS

Phys. Rev. Lett. 95 081601 (2005)

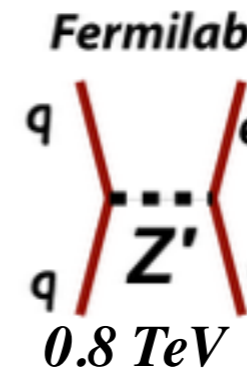
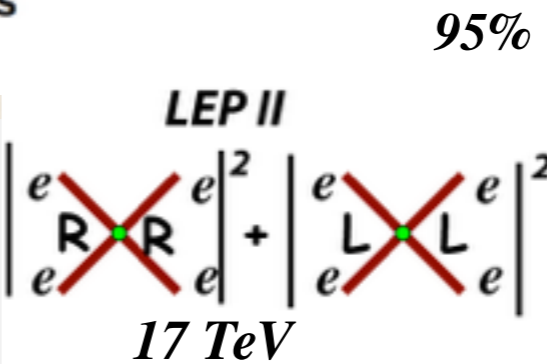
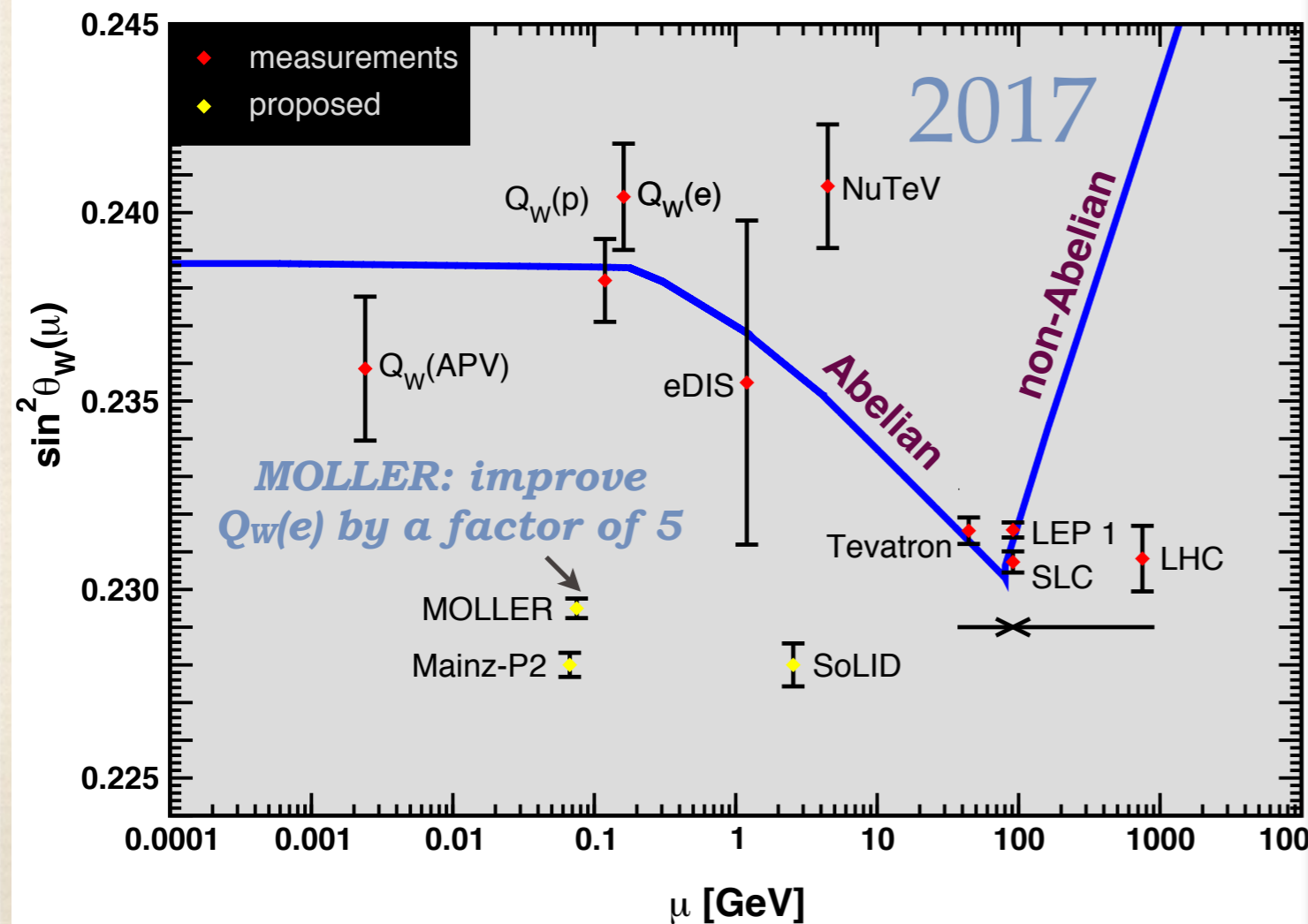
PARTICLE PHYSICS

Electrons are not ambidextrous

Andrzej Czarnecki and William J. Marciano

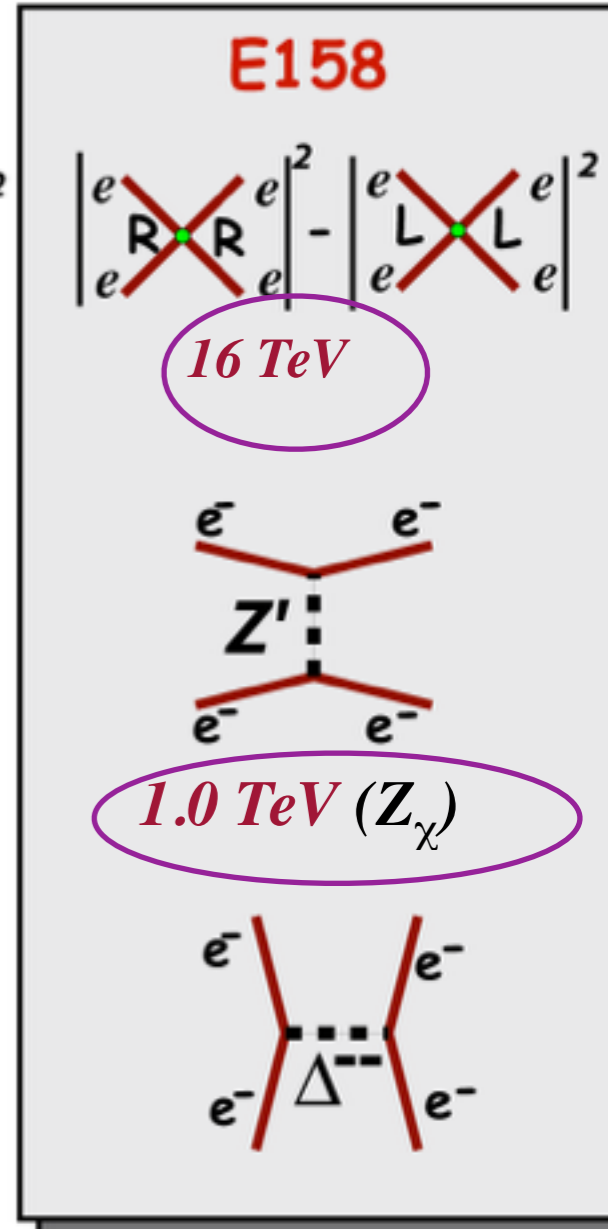
The best low-energy measurement yet obtained of the electroweak mixing angle — a central parameter of the standard model of particle physics — is the last hurrah for Stanford's powerful two-mile linear accelerator.

Limits on "New" Physics



doubly charged scalar exchange

$$0.01 \cdot G_F$$



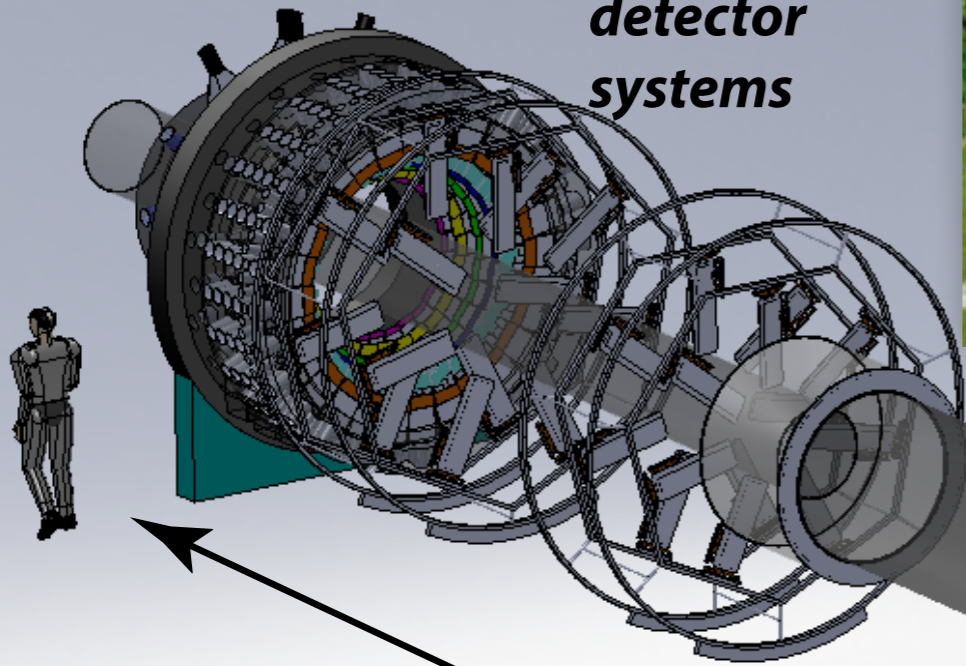
Parity-Violating Fixed Target 11 GeV electron-electron (Møller) scattering

MOLLER at JLab

Unique opportunity leveraging the 12 GeV Upgrade investment



detector systems



hybrid toroid

Evolutionary progression to extraordinary luminosity and electron beam stability with high longitudinal beam polarization

$60 \mu\text{A}$ 90% polarized electrons

Special purpose installation in Hall A

28 m

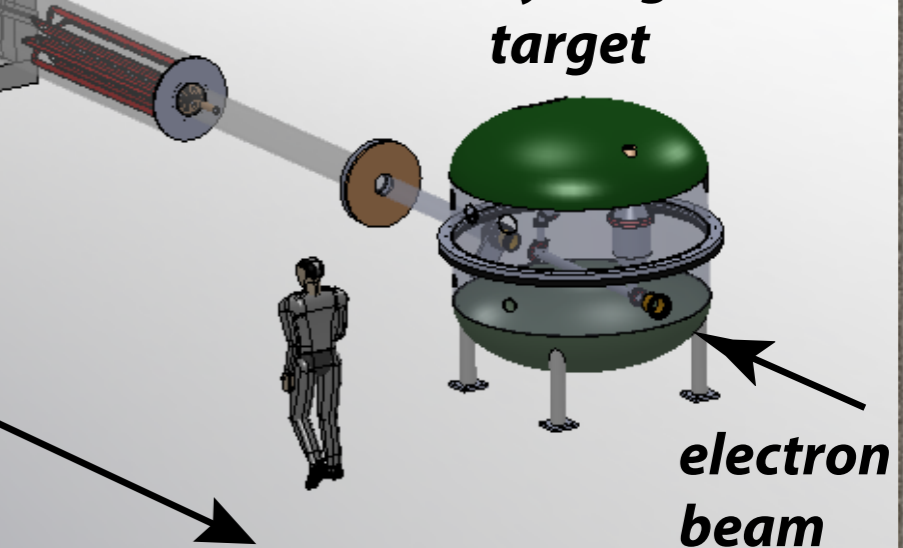
upstream toroid

liquid hydrogen target

$A_{PV} = 35 \text{ ppb}$
 $\delta(A_{PV}) = 0.73 \text{ parts per billion}$

$\delta(Q^e_W) = \pm 2.1 \% \text{ (stat.)} \pm 1.1 \% \text{ (syst.)}$

$\delta(\sin^2\theta_W) = \pm 0.00028$



95% C. L. Reach

Comparison with e^+e^- Collisions

Best reach on purely leptonic contact interaction amplitudes: LEP200

$$\mathcal{L}_{e_1 e_2} = \sum_{i,j=L,R} \frac{g_{ij}^2}{2\Lambda^2} \bar{e}_i \gamma_\mu e_i \bar{e}_j \gamma^\mu e_j$$

$$g_{ij} = 4\pi\eta_{ij}$$

Model	η_{LL}^f	η_{RR}^f	η_{LR}^f	η_{RL}^f
LL^\pm	± 1	0	0	0
RR^\pm	0	± 1	0	0
VV^\pm	± 1	± 1	± 1	± 1

LEP200 Reach

$$\Lambda_{LL}^{ee} \sim 8.3 \text{ TeV}$$

E158 Reach

$$\Lambda_{LL}^{ee} \sim 12 \text{ TeV}$$

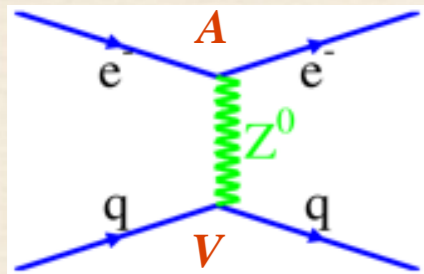
MOLLER Reach

$$\Lambda_{LL}^{ee} \sim 27 \text{ TeV}$$

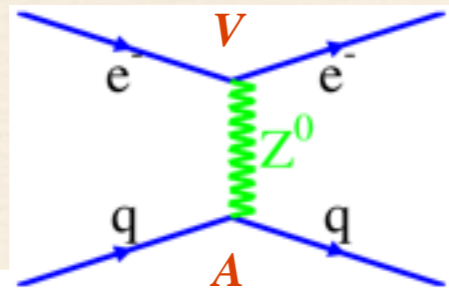
MOLLER is accessing discovery space that cannot be reached until the advent of a new lepton collider or neutrino factory

Awaiting green light from DOE to start construction ~ 2020-23

Semi-Leptonic Weak Neutral Current Interactions



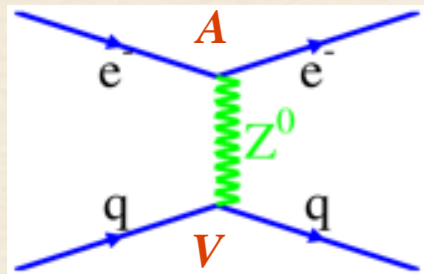
$$C_{1i} \equiv 2g_A^e g_V^i$$



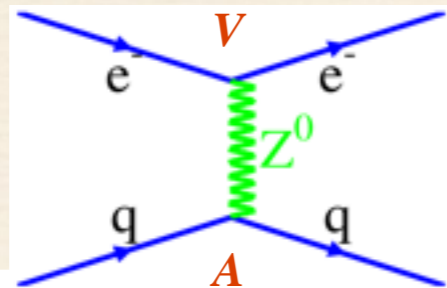
$$C_{2i} \equiv 2g_V^e g_A^i$$

$$\begin{aligned} \mathcal{L}^{PV} = & \frac{G_F}{\sqrt{2}} [\bar{e}\gamma^\mu\gamma_5 e (C_{1u}\bar{u}\gamma_\mu u + C_{1d}\bar{d}\gamma_\mu d) \\ & + \bar{e}\gamma^\mu e (C_{2u}\bar{u}\gamma_\mu\gamma_5 u + C_{2d}\bar{d}\gamma_\mu\gamma_5 d)] \\ & + C_{ee} (e\gamma^\mu\gamma_5 e \bar{e}\gamma_\mu e) \end{aligned}$$

Semi-Leptonic Weak Neutral Current Interactions



$$C_{1i} \equiv 2g_A^e g_V^i$$



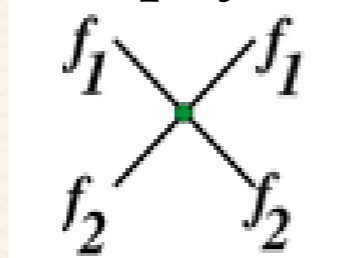
$$C_{2i} \equiv 2g_V^e g_A^i$$

$$\mathcal{L}^{PV} = \frac{G_F}{\sqrt{2}} [\bar{e}\gamma^\mu\gamma_5 e (C_{1u}\bar{u}\gamma_\mu u + C_{1d}\bar{d}\gamma_\mu d) + \bar{e}\gamma^\mu e (C_{2u}\bar{u}\gamma_\mu\gamma_5 u + C_{2d}\bar{d}\gamma_\mu\gamma_5 d)] + C_{ee} (e\gamma^\mu\gamma_5 e \bar{e}\gamma_\mu e)$$

C_{1u}	$=$	$-\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W$	\approx	-0.19
C_{1d}	$=$	$\frac{1}{2} - \frac{2}{3} \sin^2 \theta_W$	\approx	0.35
C_{2u}	$=$	$-\frac{1}{2} + 2 \sin^2 \theta_W$	\approx	-0.04
C_{2d}	$=$	$\frac{1}{2} - 2 \sin^2 \theta_W$	\approx	0.04

new physics

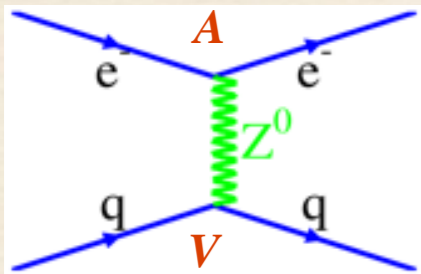
+



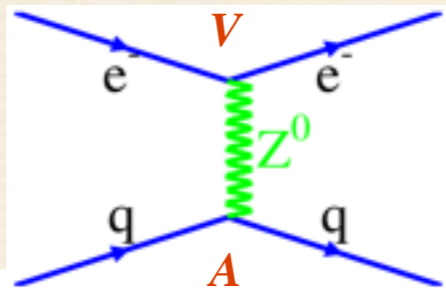
$\mathcal{L}_{f_1 f_2} =$

$$\sum_{i,j=L,R} \frac{(g_{ij}^{12})^2}{\Lambda_{ij}^2} \bar{f}_{1i}\gamma_\mu f_{1i} \bar{f}_{2j}\gamma_\mu f_{2j}$$

Semi-Leptonic Weak Neutral Current Interactions



$$C_{1i} \equiv 2g_A^e g_V^i$$

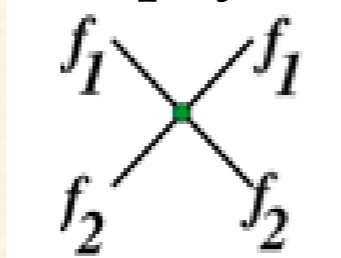


$$C_{2i} \equiv 2g_V^e g_A^i$$

$$\mathcal{L}^{PV} = \frac{G_F}{\sqrt{2}} [\bar{e}\gamma^\mu\gamma_5 e (C_{1u}\bar{u}\gamma_\mu u + C_{1d}\bar{d}\gamma_\mu d) + \bar{e}\gamma^\mu e (C_{2u}\bar{u}\gamma_\mu\gamma_5 u + C_{2d}\bar{d}\gamma_\mu\gamma_5 d)] + C_{ee} (e\gamma^\mu\gamma_5 e \bar{e}\gamma_\mu e)$$

C_{1u}	$=$	$-\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W$	\approx	-0.19
C_{1d}	$=$	$\frac{1}{2} - \frac{2}{3} \sin^2 \theta_W$	\approx	0.35
C_{2u}	$=$	$-\frac{1}{2} + 2 \sin^2 \theta_W$	\approx	-0.04
C_{2d}	$=$	$\frac{1}{2} - 2 \sin^2 \theta_W$	\approx	0.04

new physics

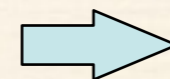


+

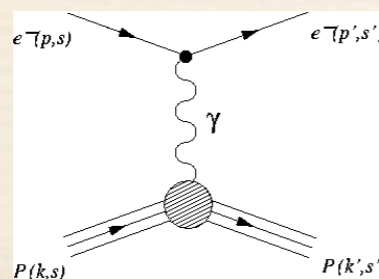
$$\mathcal{L}_{f_1 f_2} =$$

$$\sum_{i,j=L,R} \frac{(g_{ij}^{12})^2}{\Lambda_{ij}^2} \bar{f}_{1i}\gamma_\mu f_{1i} \bar{f}_{2j}\gamma_\mu f_{2j}$$

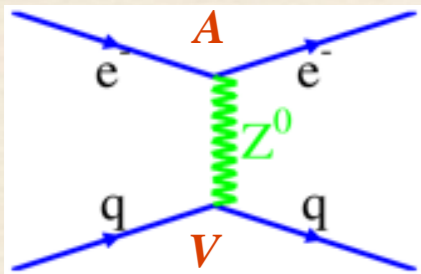
$$C_{1q} \propto (g_{RR}^{eq})^2 + (g_{RL}^{eq})^2 - (g_{LR}^{eq})^2 - (g_{LL}^{eq})^2$$



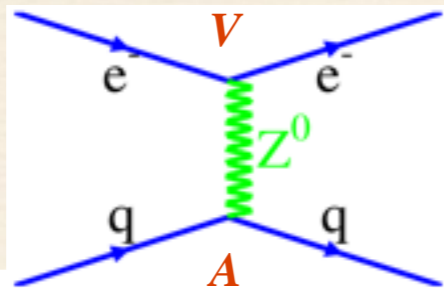
PV elastic e-N scattering, Atomic parity violation



Semi-Leptonic Weak Neutral Current Interactions



$$C_{1i} \equiv 2g_A^e g_V^i$$

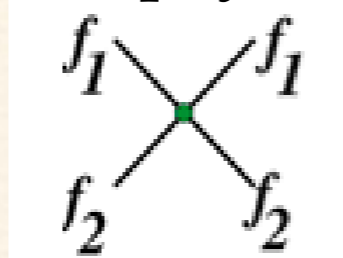


$$C_{2i} \equiv 2g_V^e g_A^i$$

$$\mathcal{L}^{PV} = \frac{G_F}{\sqrt{2}} [\bar{e}\gamma^\mu\gamma_5 e (C_{1u}\bar{u}\gamma_\mu u + C_{1d}\bar{d}\gamma_\mu d) + \bar{e}\gamma^\mu e (C_{2u}\bar{u}\gamma_\mu\gamma_5 u + C_{2d}\bar{d}\gamma_\mu\gamma_5 d)] + C_{ee}(e\gamma^\mu\gamma_5 e\bar{e}\gamma_\mu e)$$

C_{1u}	$=$	$-\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W$	\approx	-0.19
C_{1d}	$=$	$\frac{1}{2} - \frac{2}{3} \sin^2 \theta_W$	\approx	0.35
C_{2u}	$=$	$-\frac{1}{2} + 2 \sin^2 \theta_W$	\approx	-0.04
C_{2d}	$=$	$\frac{1}{2} - 2 \sin^2 \theta_W$	\approx	0.04

new physics



+

$$\mathcal{L}_{f_1 f_2} =$$

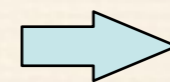
$$\sum_{i,j=L,R} \frac{(g_{ij}^{12})^2}{\Lambda_{ij}^2} \bar{f}_{1i}\gamma_\mu f_{1i} \bar{f}_{2j}\gamma_\mu f_{2j}$$

$$C_{1q} \propto (g_{RR}^{eq})^2 + (g_{RL}^{eq})^2 - (g_{LR}^{eq})^2 - (g_{LL}^{eq})^2$$

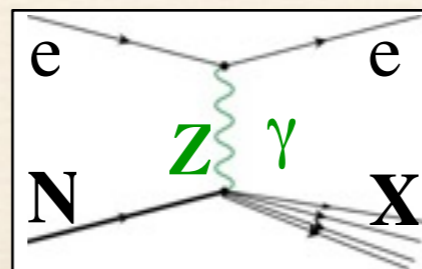
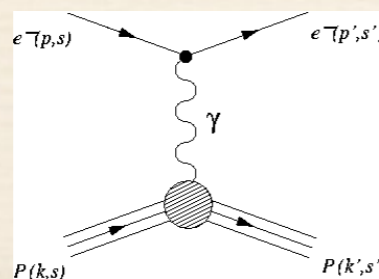


PV elastic e-N scattering, Atomic parity violation

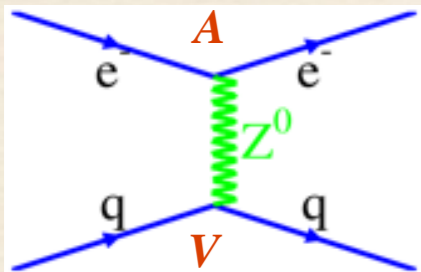
$$C_{2q} \propto (g_{RR}^{eq})^2 - (g_{RL}^{eq})^2 + (g_{LR}^{eq})^2 - (g_{LL}^{eq})^2$$



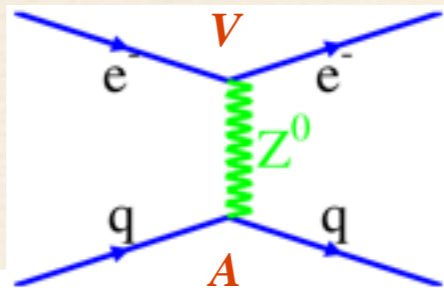
PV deep inelastic scattering



Semi-Leptonic Weak Neutral Current Interactions



$$C_{1i} \equiv 2g_A^e g_V^i$$

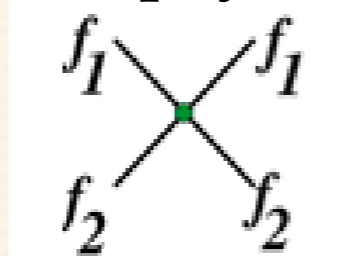


$$C_{2i} \equiv 2g_V^e g_A^i$$

$$\mathcal{L}^{PV} = \frac{G_F}{\sqrt{2}} [\bar{e}\gamma^\mu\gamma_5 e (C_{1u}\bar{u}\gamma_\mu u + C_{1d}\bar{d}\gamma_\mu d) + \bar{e}\gamma^\mu e (C_{2u}\bar{u}\gamma_\mu\gamma_5 u + C_{2d}\bar{d}\gamma_\mu\gamma_5 d)] + C_{ee} (e\gamma^\mu\gamma_5 e \bar{e}\gamma_\mu e)$$

C_{1u}	$=$	$-\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W$	\approx	-0.19
C_{1d}	$=$	$\frac{1}{2} - \frac{2}{3} \sin^2 \theta_W$	\approx	0.35
C_{2u}	$=$	$-\frac{1}{2} + 2 \sin^2 \theta_W$	\approx	-0.04
C_{2d}	$=$	$\frac{1}{2} - 2 \sin^2 \theta_W$	\approx	0.04

new physics



+

$$\mathcal{L}_{f_1 f_2} =$$

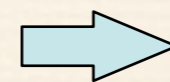
$$\sum_{i,j=L,R} \frac{(g_{ij}^{12})^2}{\Lambda_{ij}^2} \bar{f}_{1i}\gamma_\mu f_{1i} \bar{f}_{2j}\gamma_\mu f_{2j}$$

$$C_{1q} \propto (g_{RR}^{eq})^2 + (g_{RL}^{eq})^2 - (g_{LR}^{eq})^2 - (g_{LL}^{eq})^2$$



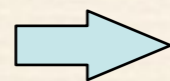
PV elastic e-N scattering, Atomic parity violation

$$C_{2q} \propto (g_{RR}^{eq})^2 - (g_{RL}^{eq})^2 + (g_{LR}^{eq})^2 - (g_{LL}^{eq})^2$$

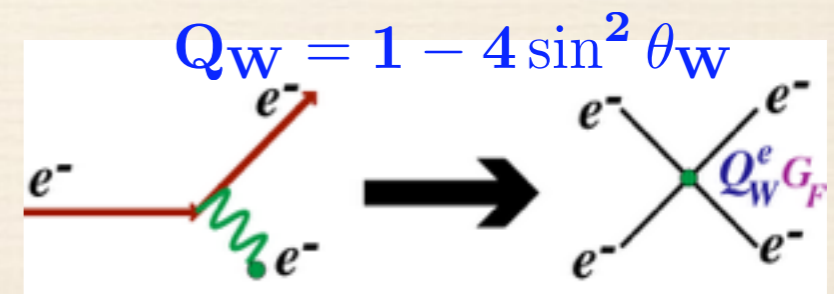
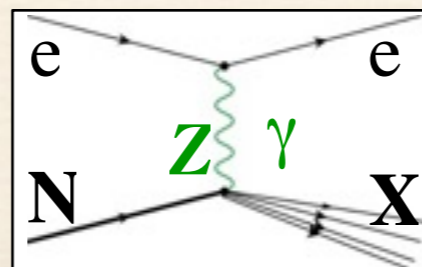
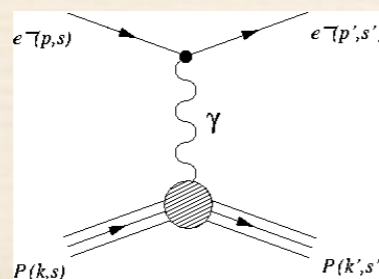


PV deep inelastic scattering

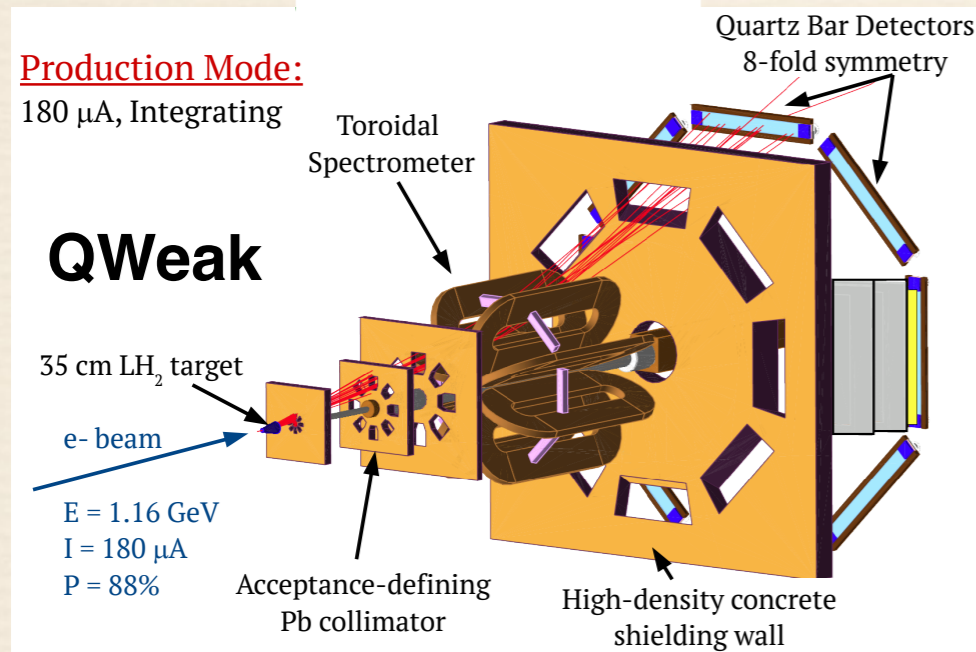
$$C_{ee} \propto (g_{RR}^{ee})^2 - (g_{LL}^{ee})^2$$



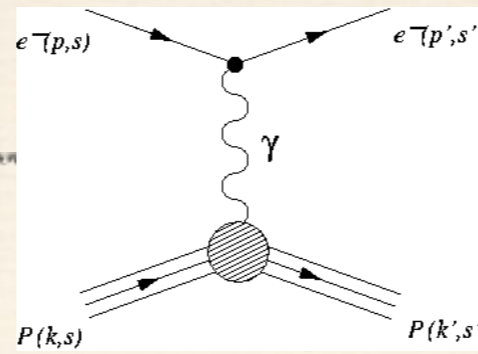
PV Møller scattering



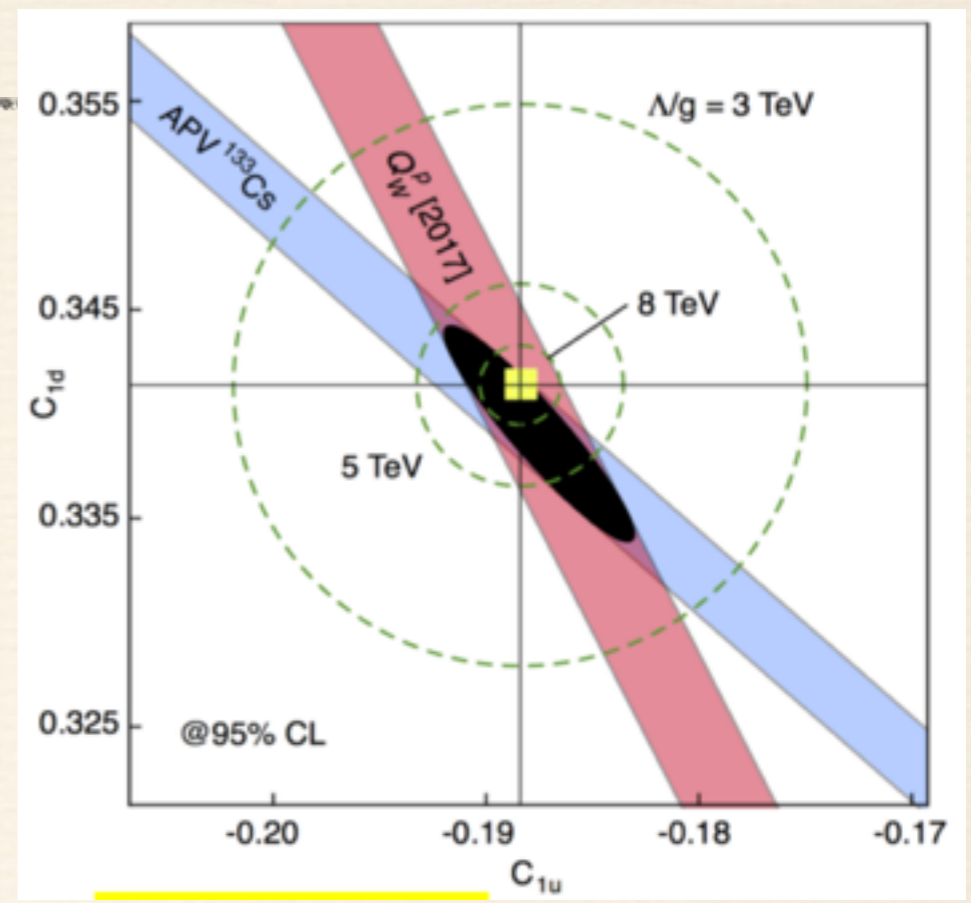
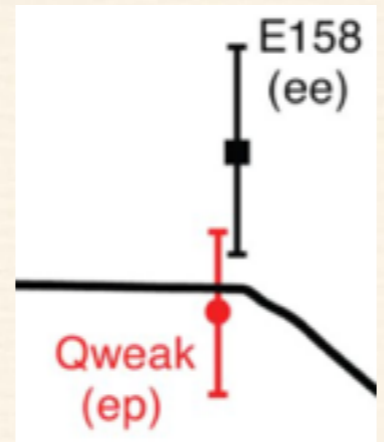
Recent Past and Future



$$A_{ep} = -226.5 \pm 7.3(\text{stat}) \pm 5.8(\text{syst}) \text{ ppb}$$



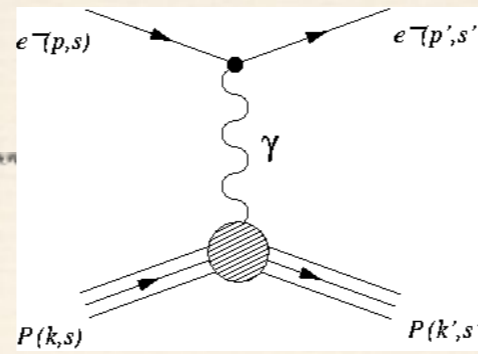
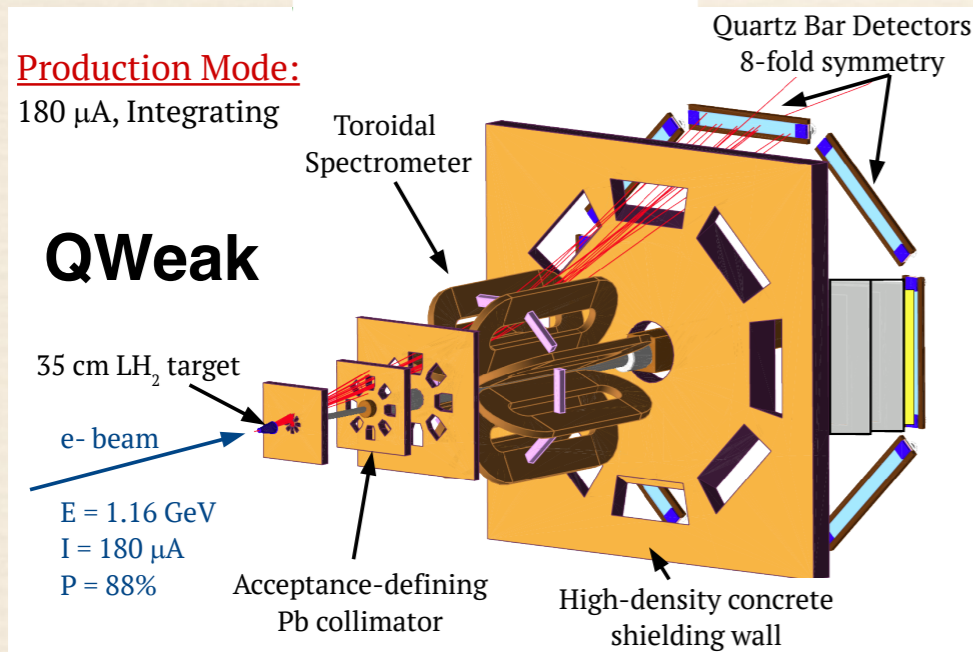
Nature 557 (2018)
no.7704, 207-211



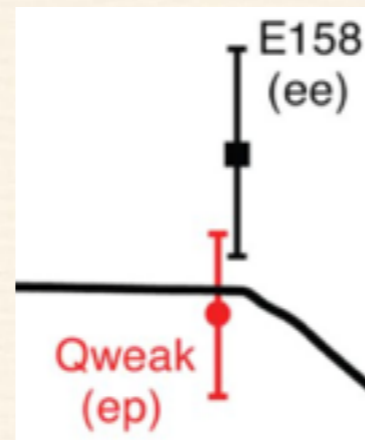
~2022

mid-2020s

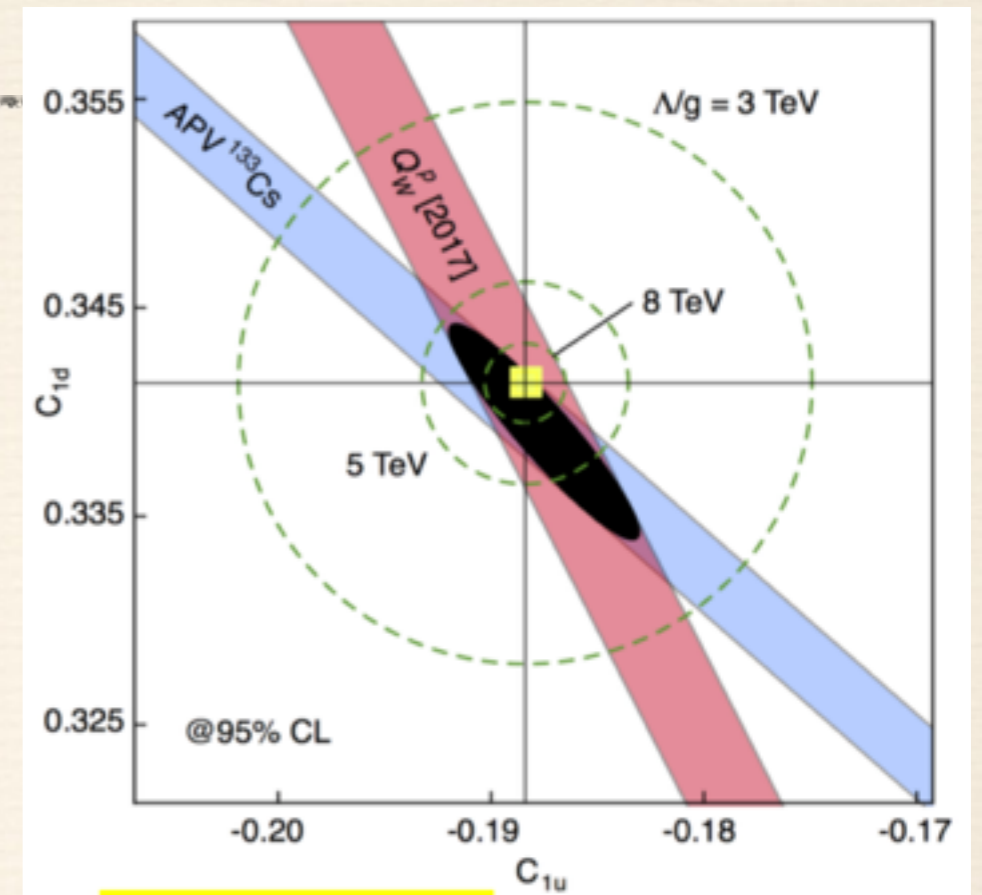
Recent Past and Future



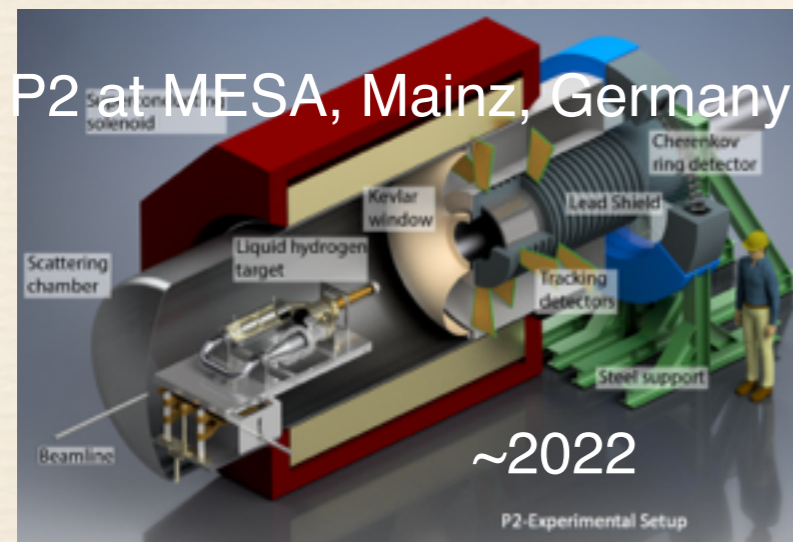
Nature 557 (2018)
no.7704, 207-211



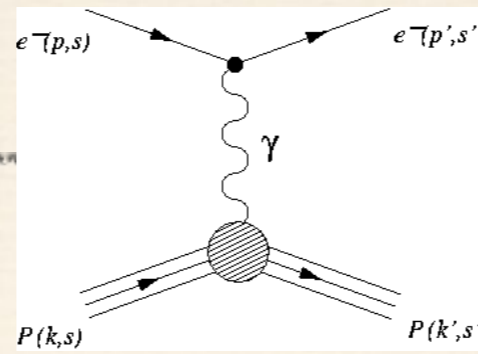
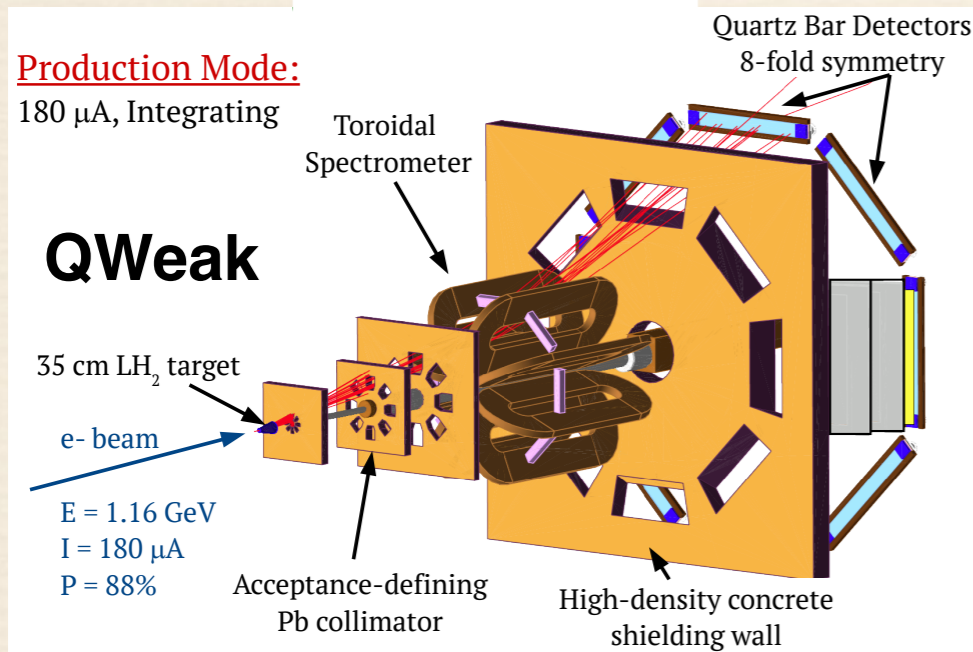
$$A_{ep} = -226.5 \pm 7.3(\text{stat}) \pm 5.8(\text{syst}) \text{ ppb}$$



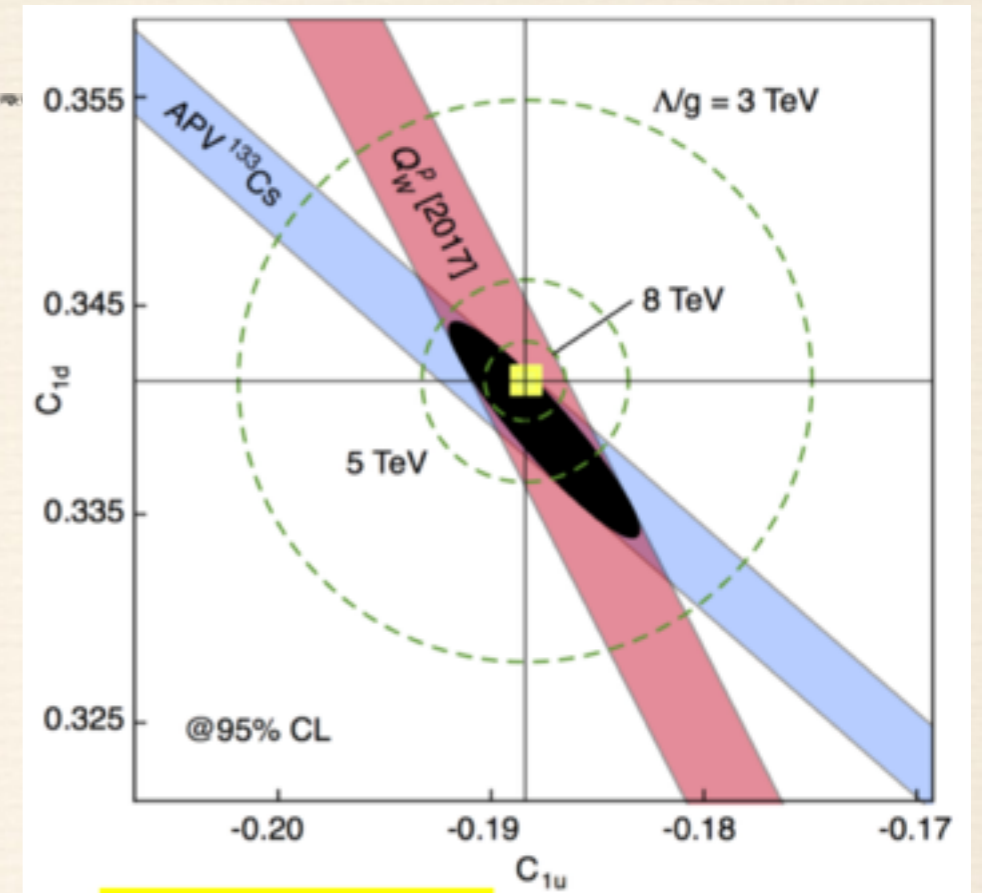
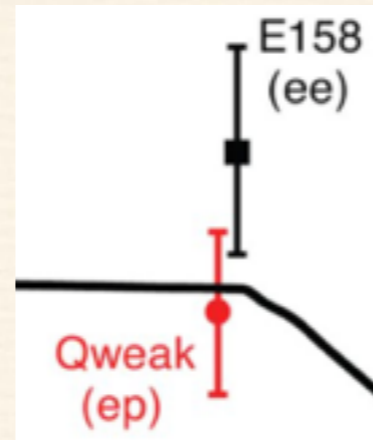
mid-2020s



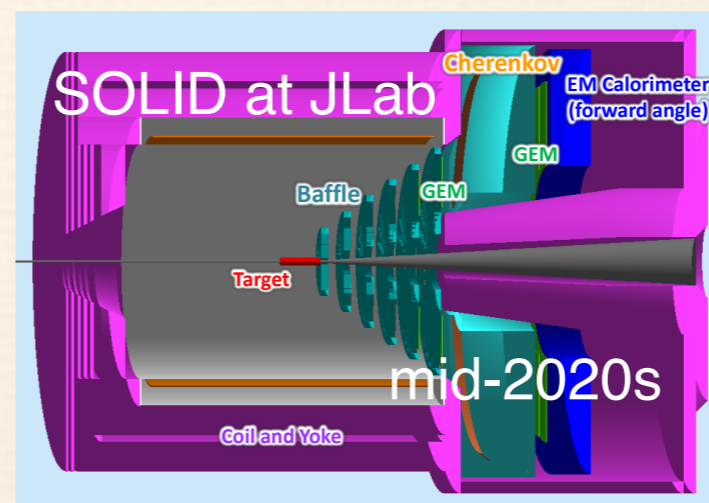
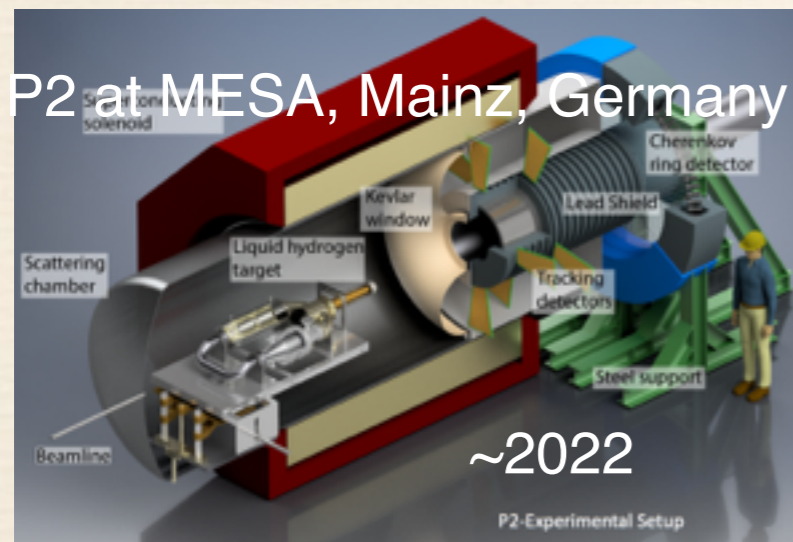
Recent Past and Future



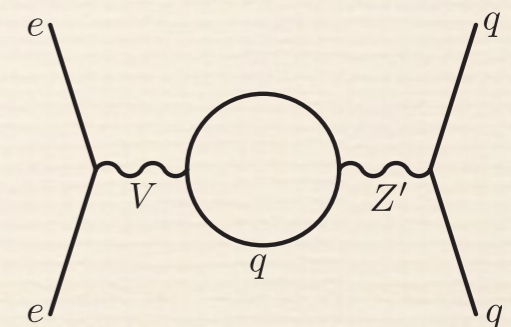
Nature 557 (2018)
no.7704, 207-211



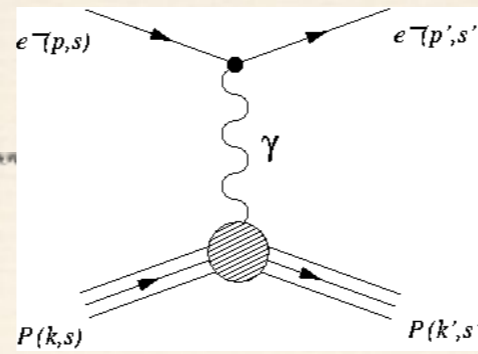
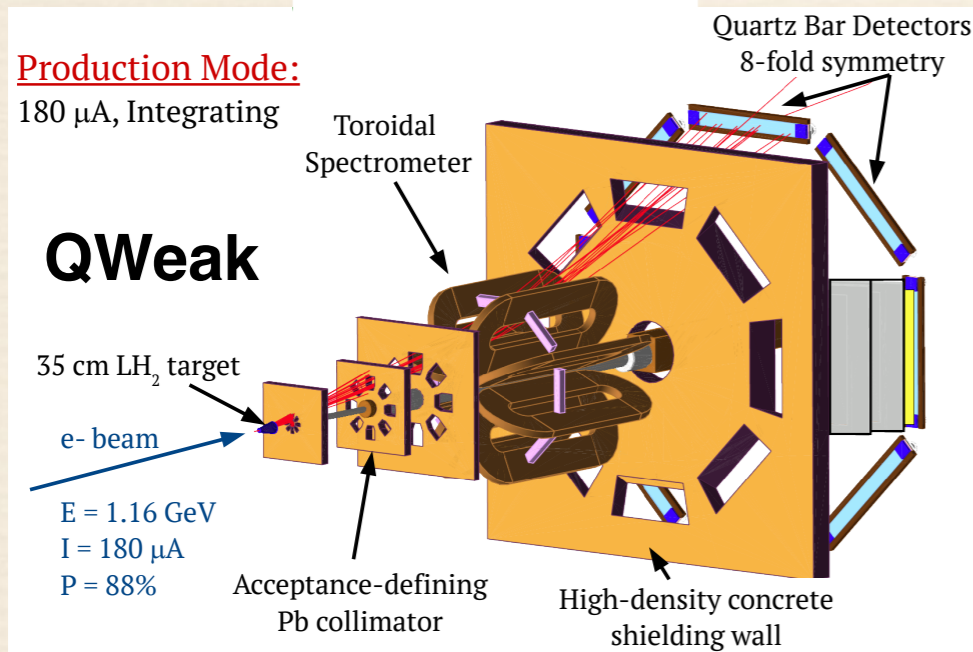
$$A_{ep} = -226.5 \pm 7.3(\text{stat}) \pm 5.8(\text{syst}) \text{ ppb}$$



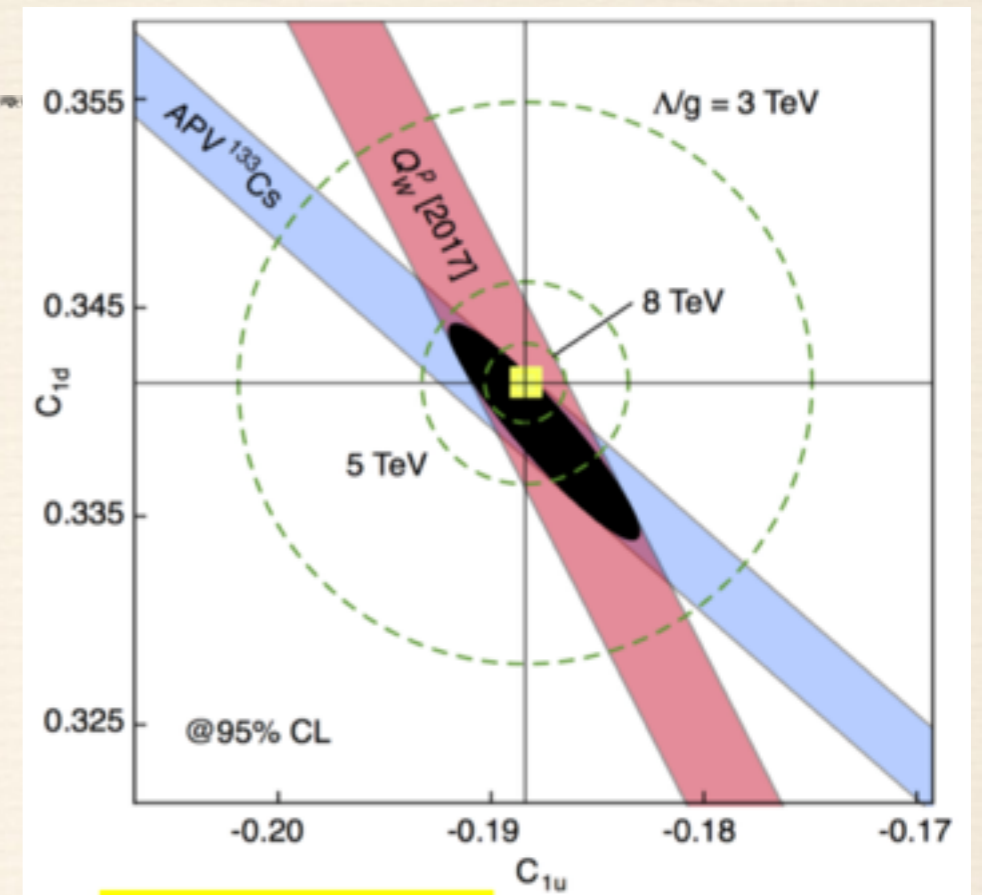
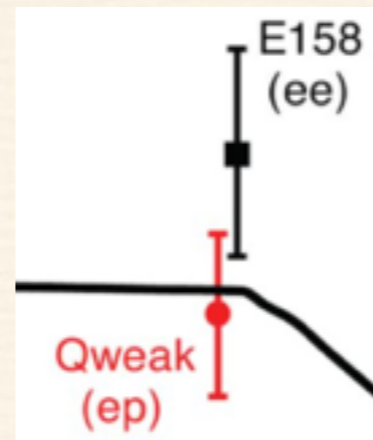
Leptophobic Z'



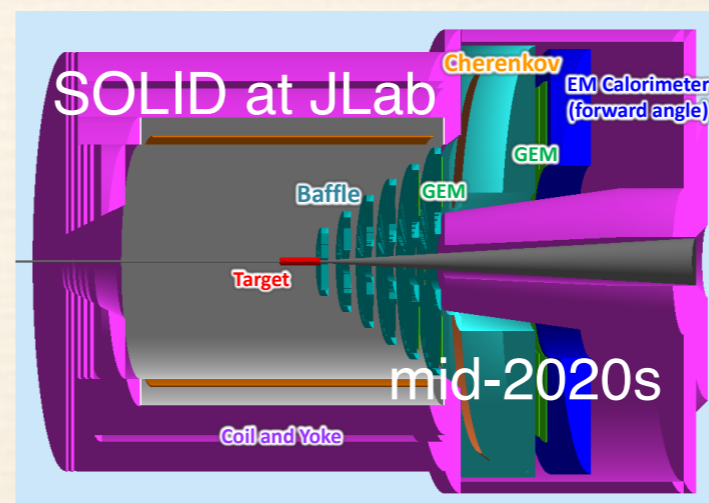
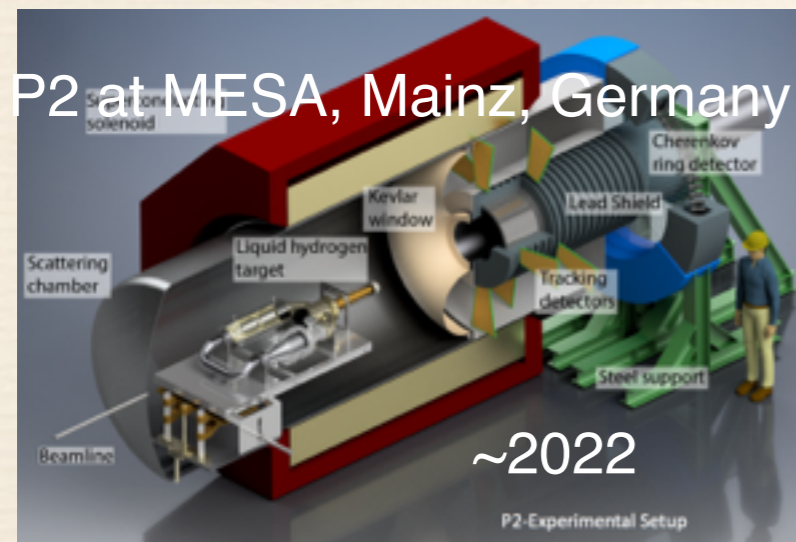
Recent Past and Future



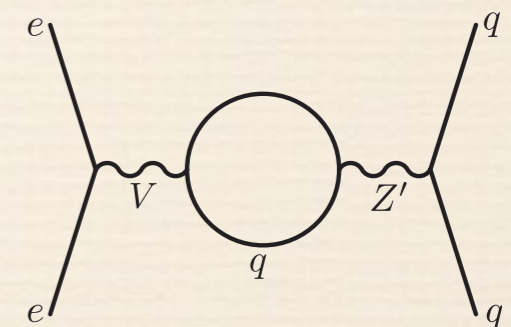
Nature 557 (2018)
no.7704, 207-211



$$A_{ep} = -226.5 \pm 7.3(\text{stat}) \pm 5.8(\text{syst}) \text{ ppb}$$

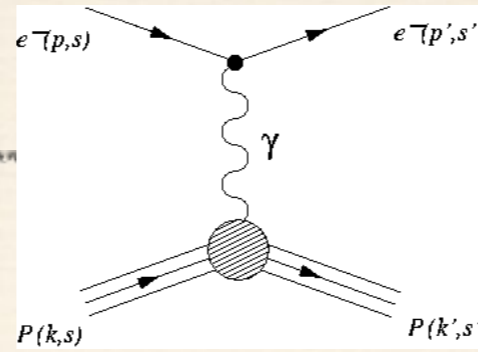
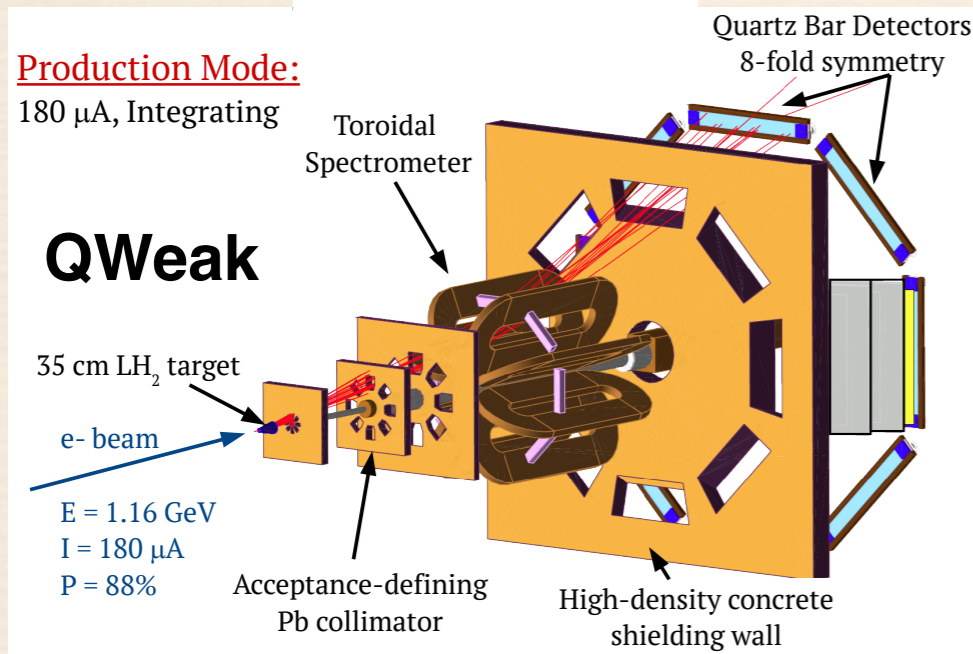


Leptophobic Z'

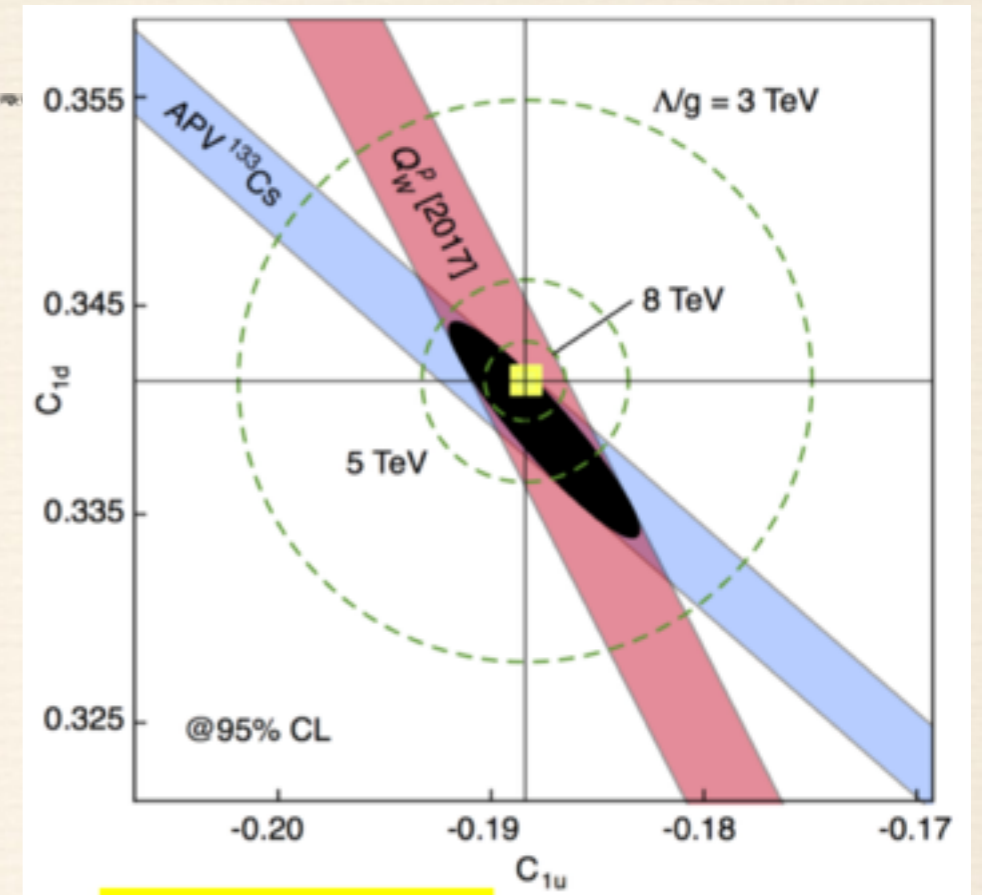
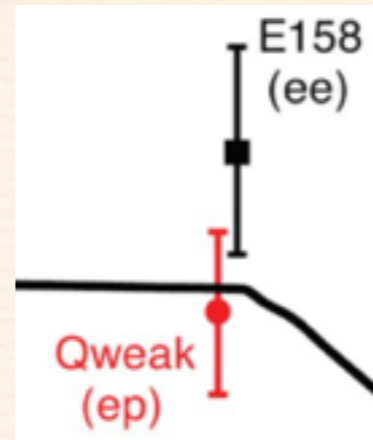


Possible new neutrino scattering measurements at the DUNE (FNAL) Near Detector

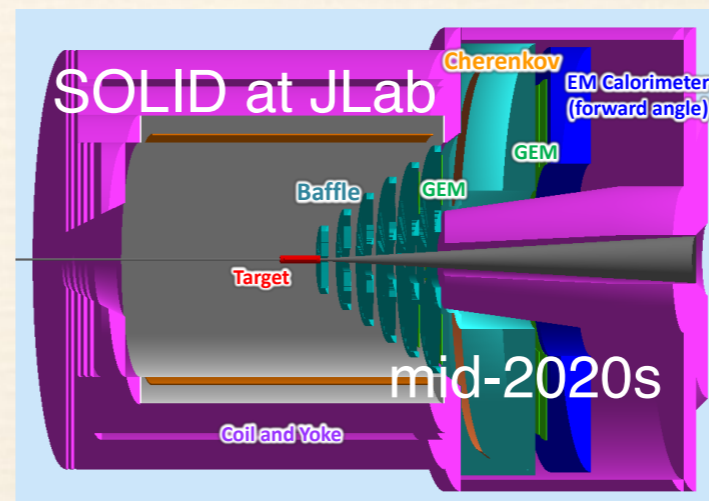
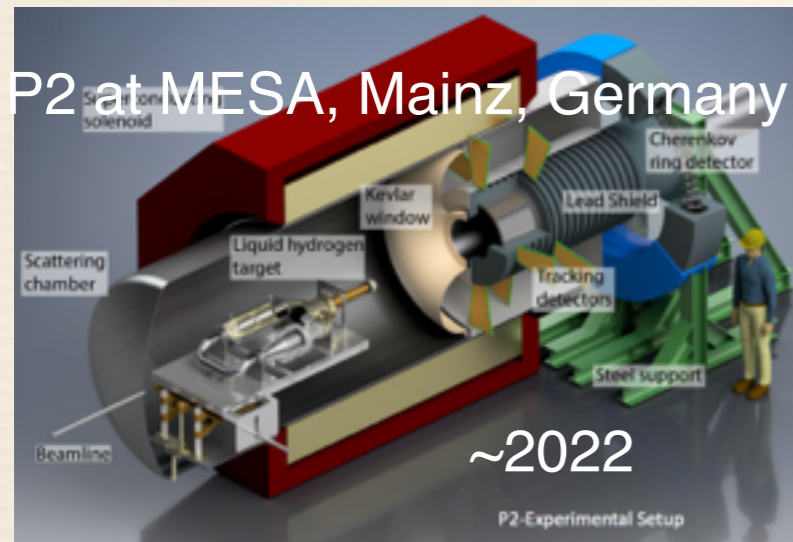
Recent Past and Future



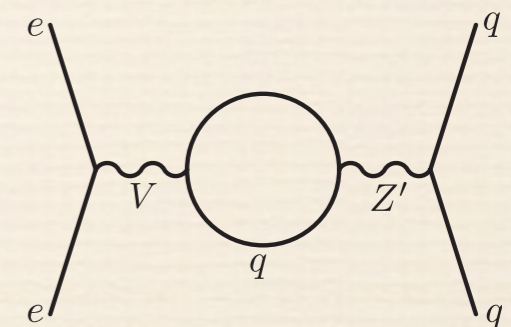
Nature 557 (2018)
no.7704, 207-211



$$A_{ep} = -226.5 \pm 7.3(\text{stat}) \pm 5.8(\text{syst}) \text{ ppb}$$



Leptophobic Z'



Possible new neutrino scattering measurements at the DUNE (FNAL) Near Detector

New Initiatives in atomic parity violation: Fr (TRIUMF), Yb, Dy (Berkeley/Mainz), Ra-ion (KVI)

Neutrinoless Double Beta Decay: Physics Motivation

A Model of Leptons: SU(2)_L X U(1)_Y

$$e^-, \mu^-, \tau^- \Rightarrow Q = -e;$$

$$\nu_e, \nu_\mu, \nu_\tau \Rightarrow Q = 0$$

$$\begin{pmatrix} \nu \\ l^- \end{pmatrix}_L \quad l^-_R \quad \nu_R$$

$\pm \frac{1}{2} \quad 0 \quad 0 \Rightarrow T_3$

$$Q = T_3 + Y/2$$

$$Y_{\nu_L} = -1 \quad Y_{\nu_R} = 0$$

**Right-handed neutrino has
no gauge interactions**

A Model of Leptons: SU(2)_L X U(1)_Y

$$e^-, \mu^-, \tau^- \Rightarrow Q = -e;$$

$$\nu_e, \nu_\mu, \nu_\tau \Rightarrow Q = 0$$

$$\begin{pmatrix} \nu \\ l^- \end{pmatrix}_L \quad l^-_R \quad \nu_R$$

$\pm \frac{1}{2} \quad 0 \quad 0 \Rightarrow T_3$

Helicity $\equiv \vec{p} \cdot \vec{\Sigma} \equiv h = \pm 1$

$$Q = T_3 + Y/2$$

$$Y_{\nu_L} = -1 \quad Y_{\nu_R} = 0$$

**Right-handed neutrino has
no gauge interactions**

Chirality $\equiv \frac{1 \pm \gamma^5}{2} \equiv P_{L,R}$

A Model of Leptons: SU(2)_L X U(1)_Y

$$e^-, \mu^-, \tau^- \Rightarrow Q = -e;$$

$$\nu_e, \nu_\mu, \nu_\tau \Rightarrow Q = 0$$

$$\begin{pmatrix} \nu \\ l^- \end{pmatrix}_L \quad l^-_R \quad \nu_R$$

$\pm \frac{1}{2} \quad 0 \quad 0 \Rightarrow T_3$

$$Q = T_3 + Y/2$$

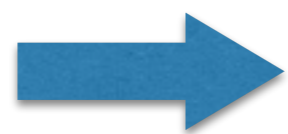
$$Y_{\nu_L} = -1 \quad Y_{\nu_R} = 0$$

Right-handed neutrino has no gauge interactions

Helicity $\equiv \vec{p} \cdot \vec{\Sigma} \equiv h = \pm 1$

Chirality $\equiv \frac{1 \pm \gamma^5}{2} \equiv P_{L,R}$

For a massless particle
(or ultra-relativistic limit)



helicity = chirality

A Model of Leptons: SU(2)_L X U(1)_Y

$$e^-, \mu^-, \tau^- \Rightarrow Q = -e;$$

$$Q = T_3 + Y/2$$

$$\nu_e, \nu_\mu, \nu_\tau \Rightarrow Q = 0$$

$$Y_{\nu_L} = -1 \quad Y_{\nu_R} = 0$$

$$\begin{pmatrix} \nu \\ l^- \end{pmatrix}_L \quad l^-_R \quad \nu_R$$

$\pm \frac{1}{2} \quad 0 \quad 0 \Rightarrow T_3$

**Right-handed neutrino has
no gauge interactions**

Helicity $\equiv \vec{p} \cdot \vec{\Sigma} \equiv h = \pm 1$

Chirality $\equiv \frac{1 \pm \gamma^5}{2} \equiv P_{L,R}$

For a massless particle
(or ultra-relativistic limit)

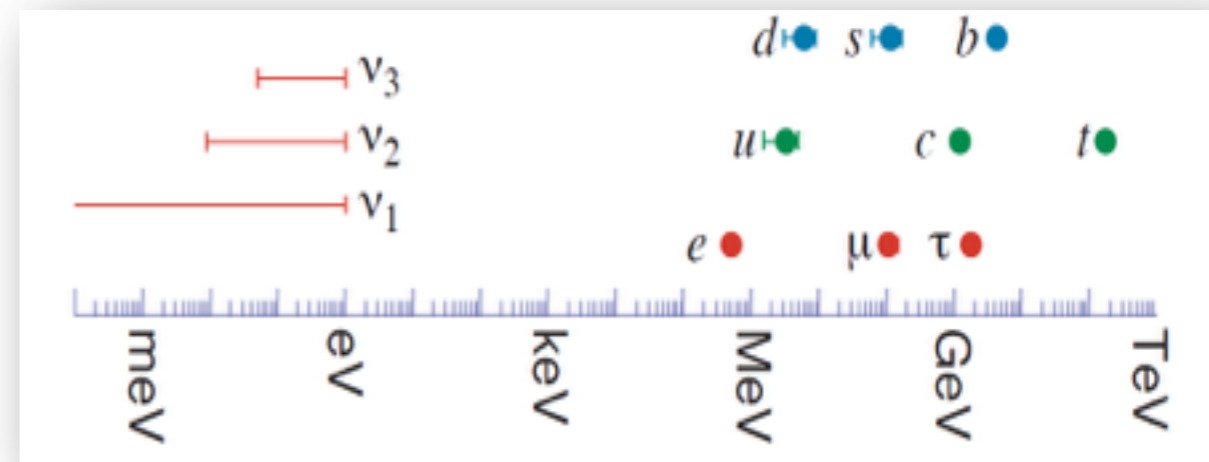
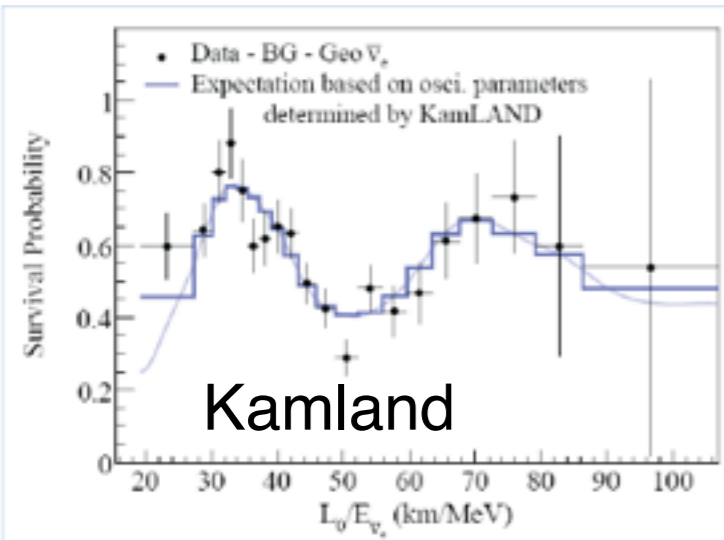


helicity = chirality

**Original formulation of the Standard Model:
 ν massless and no right-handed state**

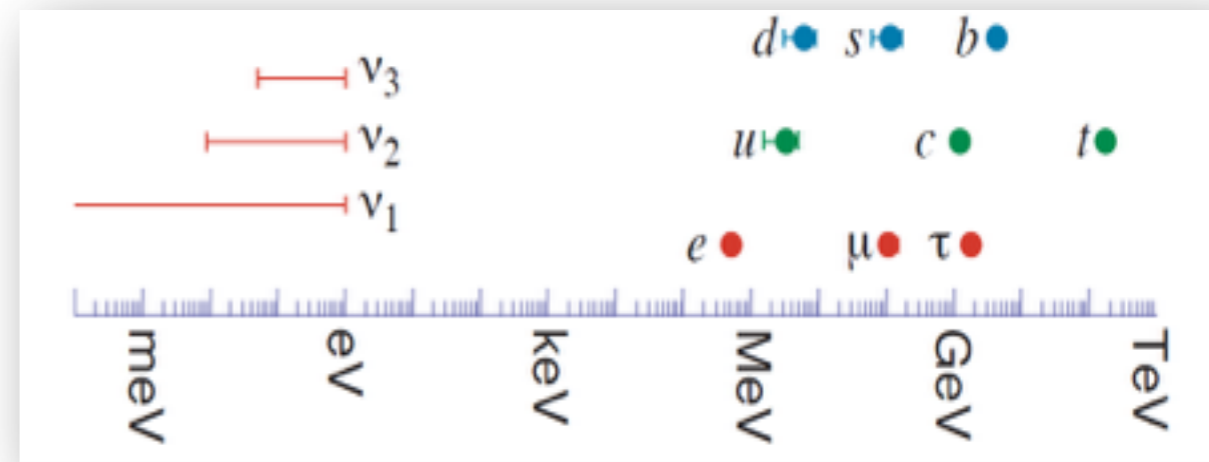
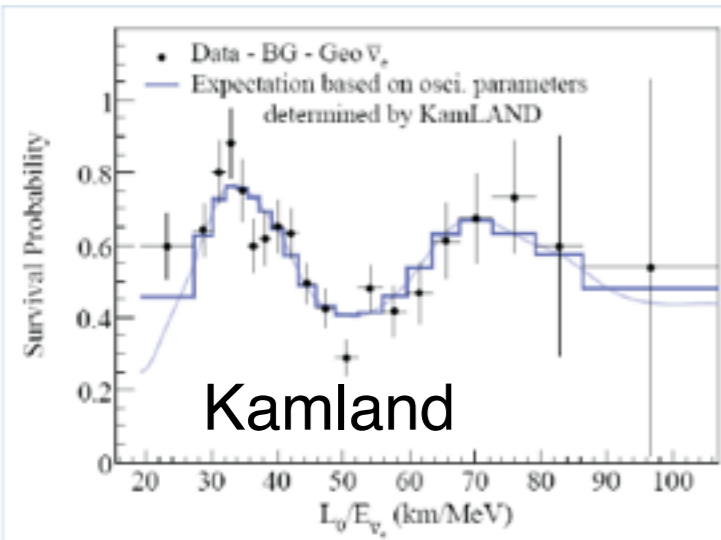
Postulate the Massive Right-Handed Neutrino

- Why is neutrino mass so small?
- How small is it?
- What is the mass generating mechanism?
- And...



Postulate the Massive Right-Handed Neutrino

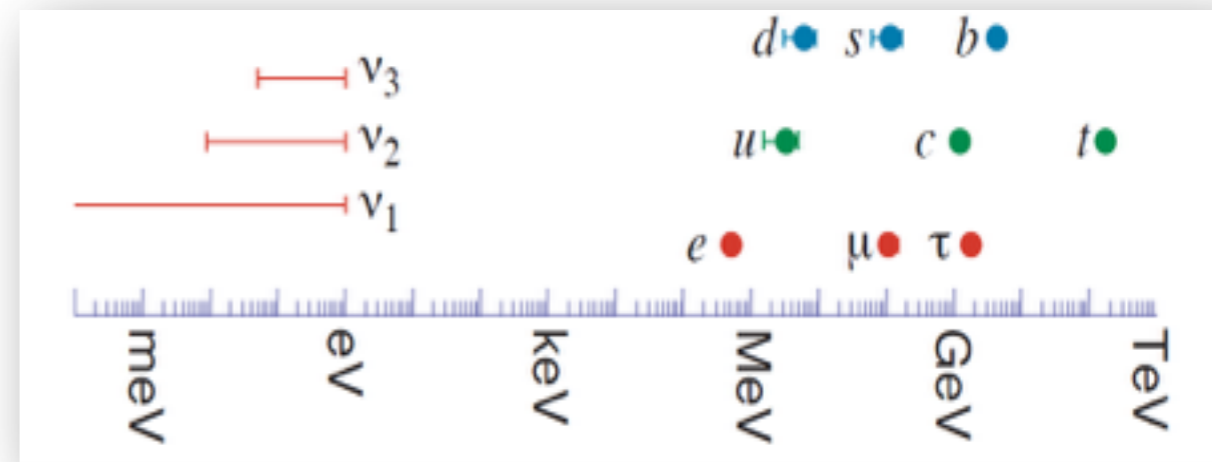
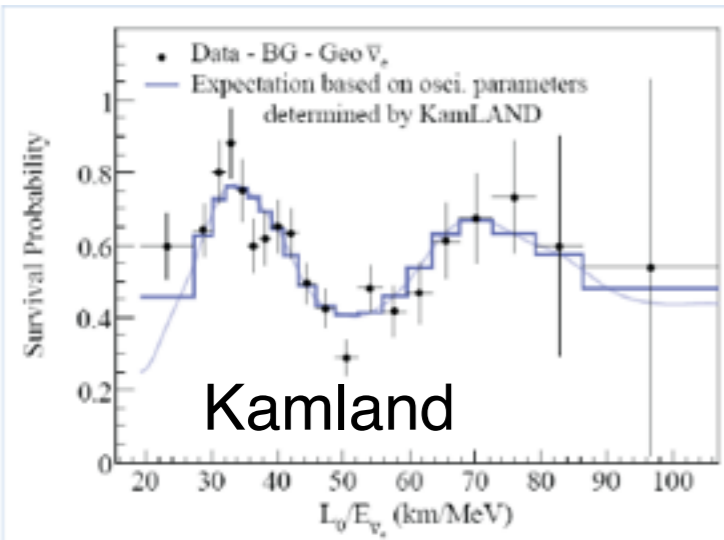
- Why is neutrino mass so small?
- How small is it?
- What is the mass generating mechanism?
- And...



CPT transformation: *left-handed* particle to *right-handed anti-particle*

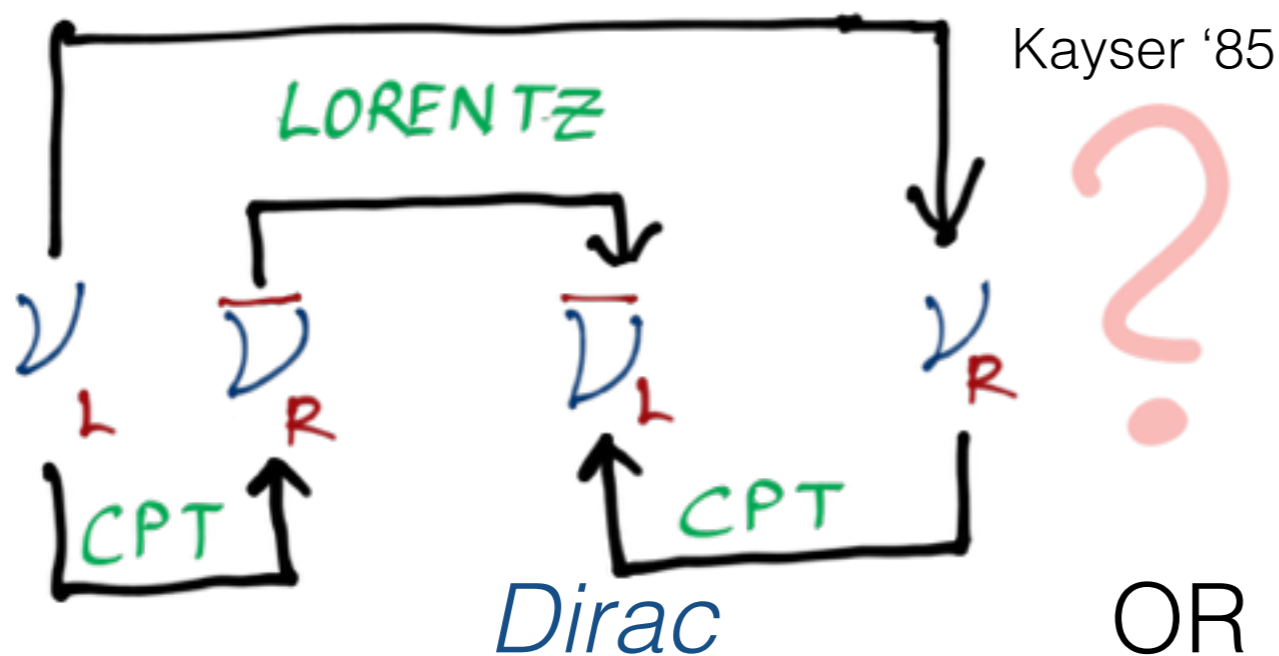
Postulate the Massive Right-Handed Neutrino

- Why is neutrino mass so small?
- How small is it?
- What is the mass generating mechanism?
- And...

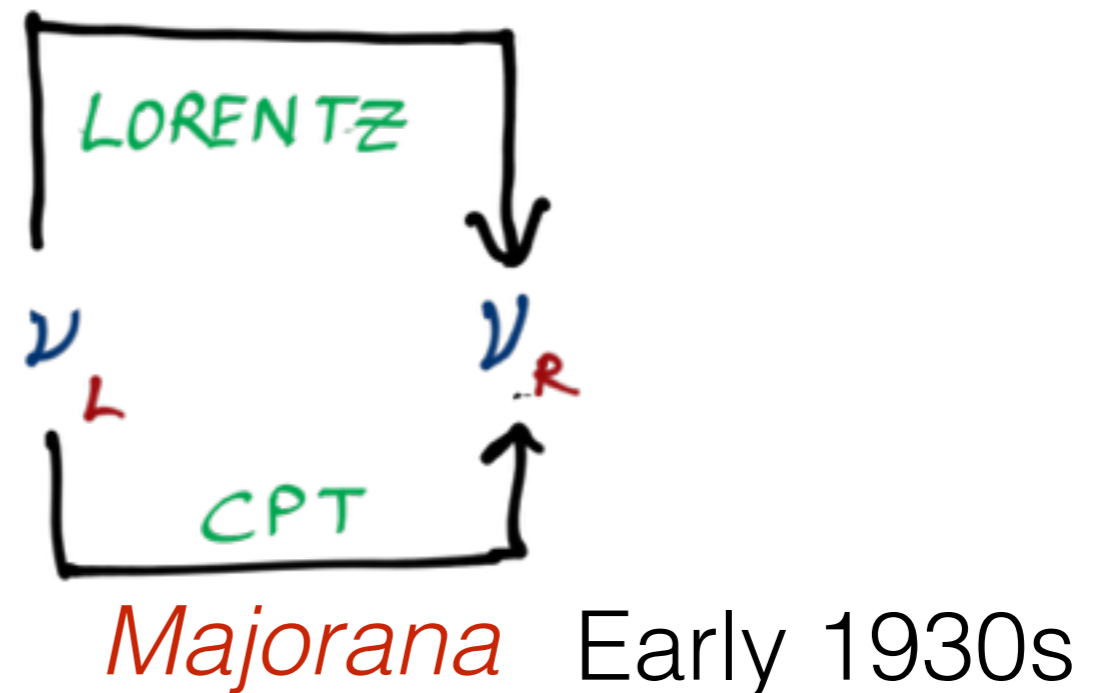


CPT transformation: *left-handed* particle to *right-handed anti-particle*

A profound question:



OR



What is the Discovery Experiment?

Neutral Current interactions have subtle differences

Kayser '82

But

Dirac-Majorana Confusion Theorem: the difference between ν_D and ν_M interactions vanishes in the ultra-relativistic limit

Exotic possibilities beyond Standard Model V-A

Nevertheless

What is the Discovery Experiment?

Neutral Current interactions have subtle differences

Kayser '82

But

Dirac-Majorana Confusion Theorem: the difference between ν_D and ν_M interactions vanishes in the ultra-relativistic limit

Exotic possibilities beyond Standard Model V-A

Nevertheless

The most pragmatic approach to discover the Majorana nature of neutrinos is to search for **Lepton Number Violation (LNV)**

Practically: discover **Neutrinoless Double-Beta Decay ($0\nu\beta\beta$)**

What is the Discovery Experiment?

Neutral Current interactions have subtle differences

Kayser '82

But

Dirac-Majorana Confusion Theorem: the difference between ν_D and ν_M interactions vanishes in the ultra-relativistic limit

Exotic possibilities beyond Standard Model V-A

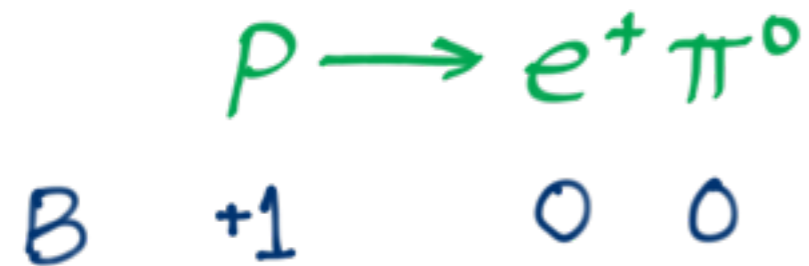
Nevertheless

The most pragmatic approach to discover the Majorana nature of neutrinos is to search for **Lepton Number Violation (LNV)**

Practically: discover **Neutrinoless Double-Beta Decay ($0\nu\beta\beta$)**

Why do we care so much?

Baryon & Lepton Number



Proton Decay

Forbidden if B is conserved

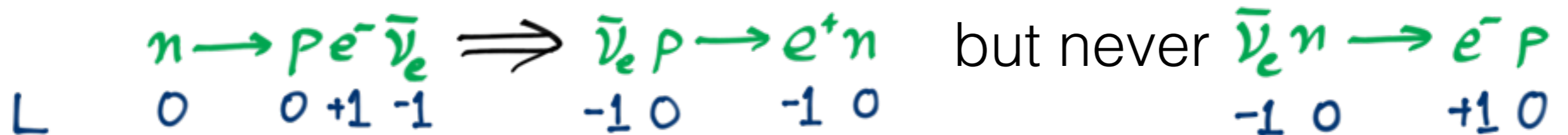
Baryon & Lepton Number



Proton Decay

$$B \quad +1 \quad 0 \quad 0$$

Forbidden if B is conserved



$$0 \quad 0 \quad +1 \quad -1 \quad -1 \quad 0 \quad -1 \quad 0$$

$$-1 \quad 0 \quad +1 \quad 0$$

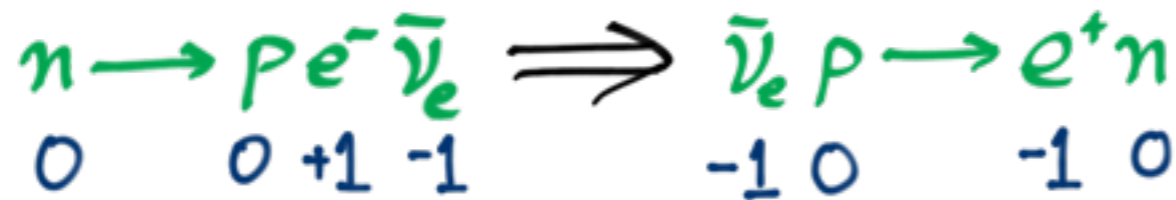
Baryon & Lepton Number



Proton Decay

$$B \quad +1 \quad 0 \quad 0$$

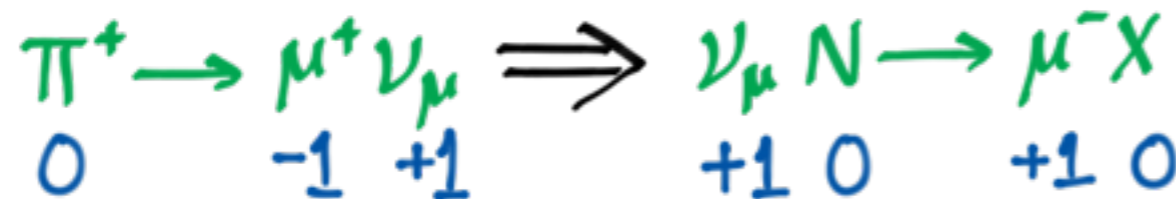
Forbidden if B is conserved



but never $\bar{\nu}_e n \rightarrow e^- p$

$$L \quad 0 \quad 0 \quad +1 \quad -1 \quad -1 \quad 0 \quad -1 \quad 0$$

$$-1 \quad 0 \quad +1 \quad 0$$



but never $\mu^+ X$

$$L \quad 0 \quad -1 \quad +1 \quad +1 \quad 0 \quad +1 \quad 0$$

$$-1 \quad 0$$

Baryon & Lepton Number



Proton Decay

$$B \quad +1 \quad 0 \quad 0$$

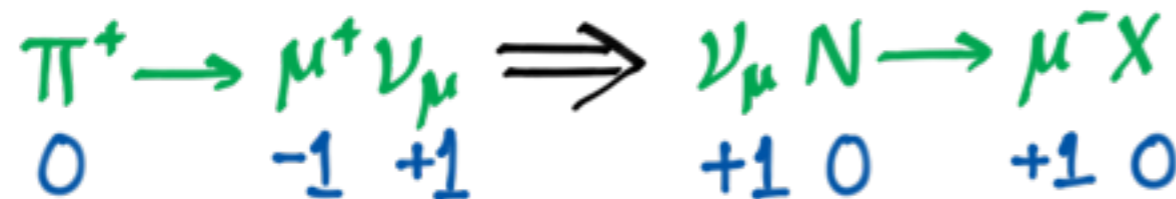
Forbidden if B is conserved



but never $\bar{\nu}_e n \rightarrow e^- p$

$$L \quad 0 \quad 0 \quad +1 \quad -1 \quad -1 \quad 0 \quad -1 \quad 0$$

$$-1 \quad 0 \quad +1 \quad 0$$



but never $\mu^+ X$

$$L \quad 0 \quad -1 \quad +1 \quad +1 \quad 0 \quad +1 \quad 0$$

$$-1 \quad 0$$

Introduce Lepton Number:

$$L_{e^-} = L_{\nu_e} = -L_{e^+} = -L_{\bar{\nu}_e} = +1$$

This is encoded into the Standard Model Feynman Rules

Is L Conserved?

If L ν_R \longleftrightarrow $\bar{\nu}_R$ Majorana Neutrino **L is violated**
 $+1$ -1

- *Only B-L strictly conserved in the Standard Model*
- *B+L is violated due to anomalies*
- *No fundamental reason to expect B and L to be conserved (assuming only 4 forces in Nature)*

Is L Conserved?

If L ν_R \longleftrightarrow $\bar{\nu}_R$ Majorana Neutrino **L is violated**
 $+1$ -1

- *Only B-L strictly conserved in the Standard Model*
- *B+L is violated due to anomalies*
- *No fundamental reason to expect B and L to be conserved (assuming only 4 forces in Nature)*

Neutrinos only interact via the parity-violating weak interaction:

Chirality can explain all observed weak interaction phenomena

Is L Conserved?

If $\nu_R \leftrightarrow \bar{\nu}_R$ Majorana Neutrino **L is violated**
 $L \quad \begin{matrix} \nu_R \\ +1 \end{matrix} \quad \begin{matrix} \bar{\nu}_R \\ -1 \end{matrix}$

- *Only B-L strictly conserved in the Standard Model*
- *B+L is violated due to anomalies*
- *No fundamental reason to expect B and L to be conserved (assuming only 4 forces in Nature)*

Neutrinos only interact via the parity-violating weak interaction:

Chirality can explain all observed weak interaction phenomena

Majorana neutrinos: possible explanation of light neutrino masses

Matter-antimatter asymmetry.... Mass generation beyond the Higgs...

Is L Conserved?

If $\nu_R \leftrightarrow \bar{\nu}_R$ Majorana Neutrino **L is violated**

L $\begin{matrix} \nu_R \\ +1 \end{matrix} \leftrightarrow \begin{matrix} \bar{\nu}_R \\ -1 \end{matrix}$

- *Only B-L strictly conserved in the Standard Model*
- *B+L is violated due to anomalies*
- *No fundamental reason to expect B and L to be conserved (assuming only 4 forces in Nature)*

Neutrinos only interact via the parity-violating weak interaction:

Chirality can explain all observed weak interaction phenomena

Majorana neutrinos: possible explanation of light neutrino masses

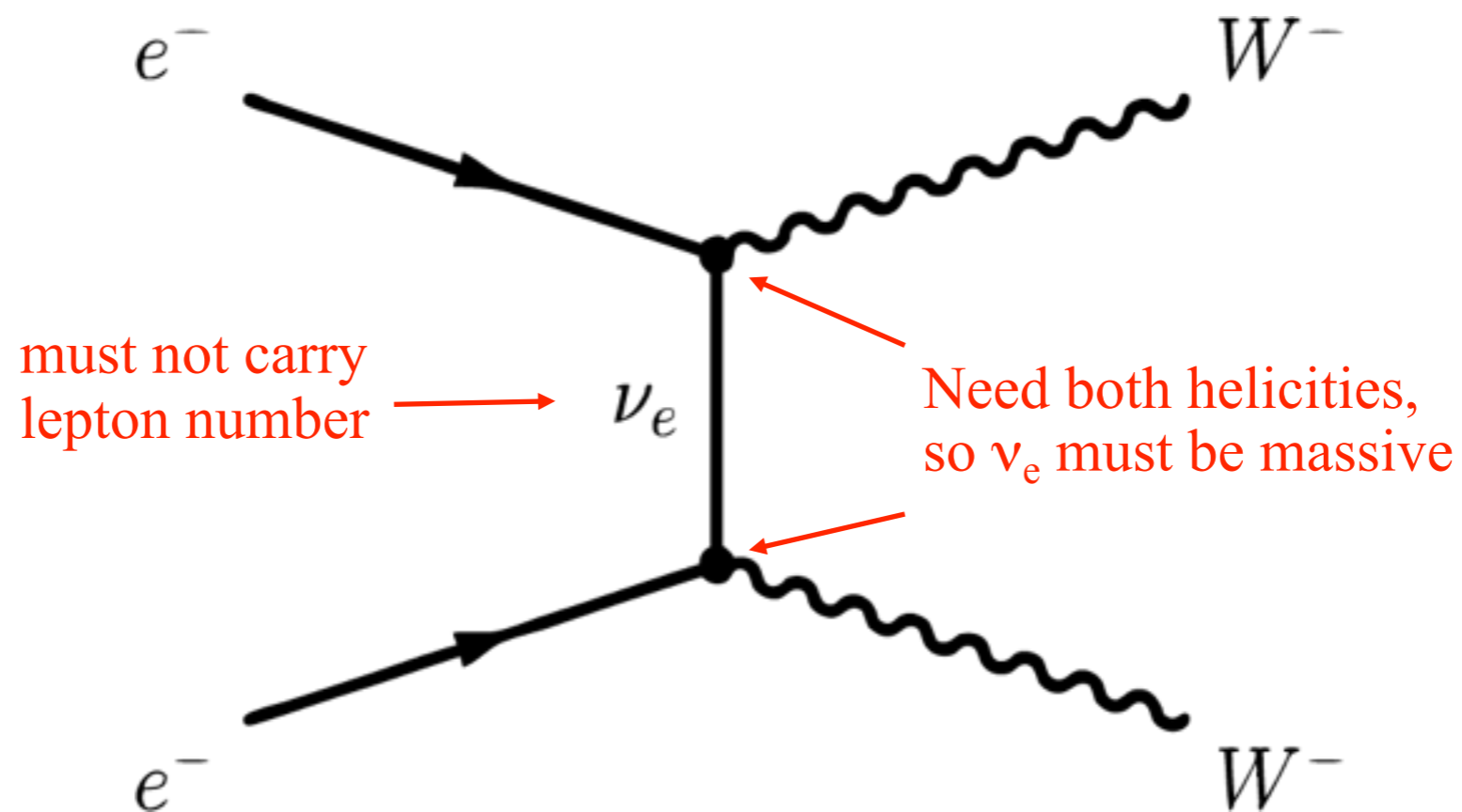
Matter-antimatter asymmetry.... Mass generation beyond the Higgs...

In any case: a new heavy scale for physics beyond the SM

A Gedanken Experiment

L

+1 +1 0 0



must not carry
lepton number

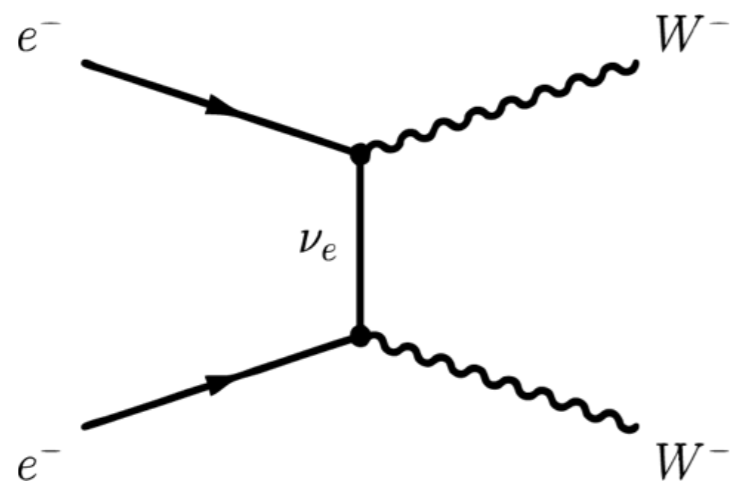
Need both helicities,
so ν_e must be massive

Lepton number changes by two units: $\Delta L=2$

For light neutrinos, this cross-section is unobservably small

Virtual W's Instead

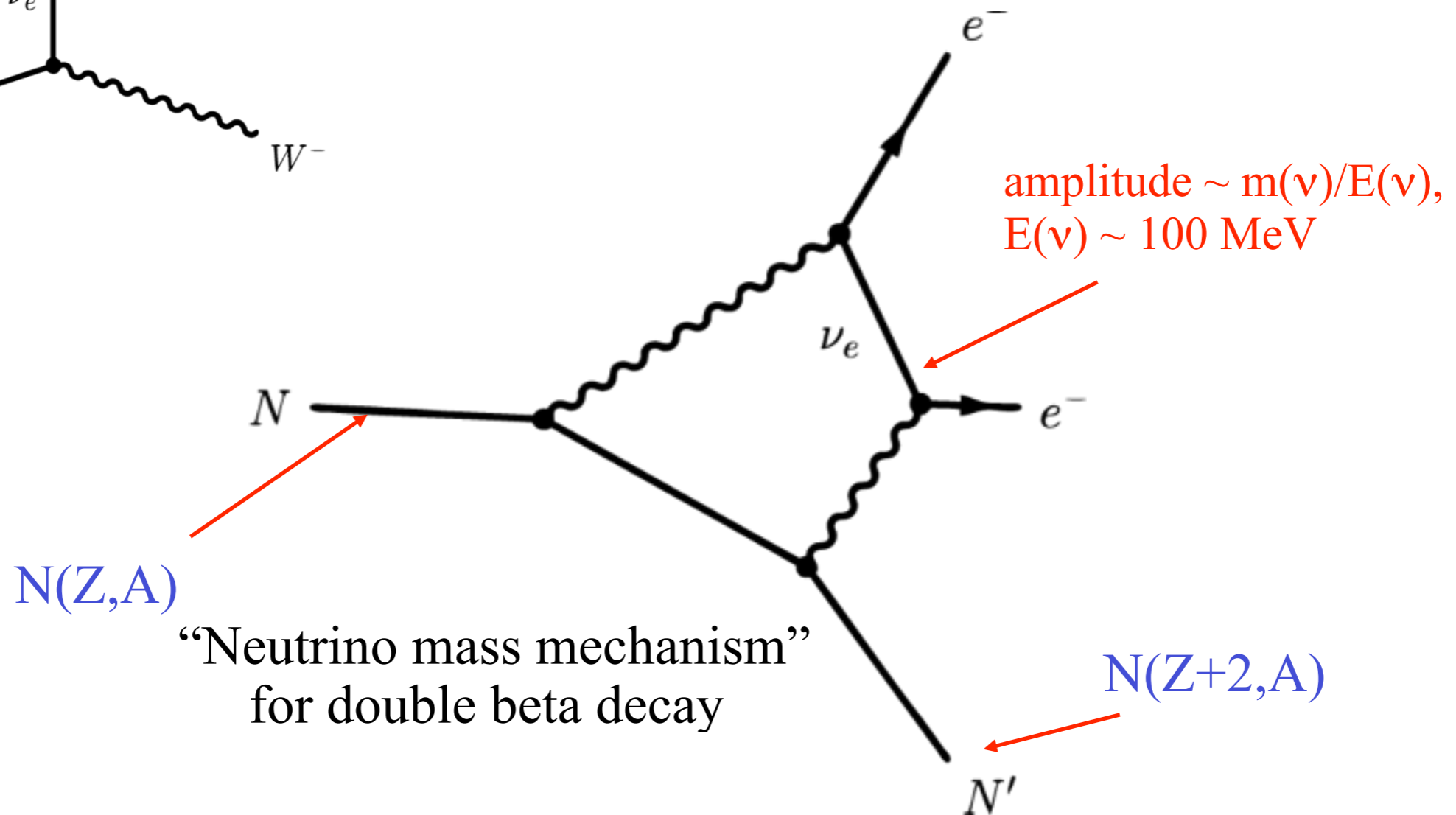
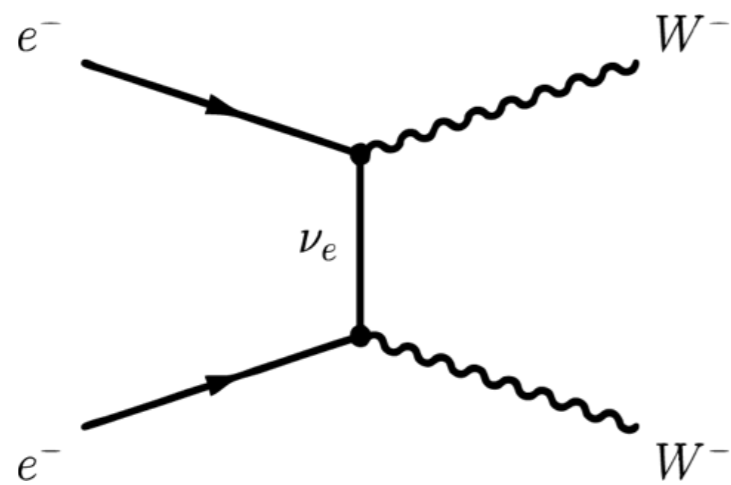
Lepton number changes by two units: $\Delta L=2$



Virtual W's Instead

Lepton number changes by two units: $\Delta L=2$

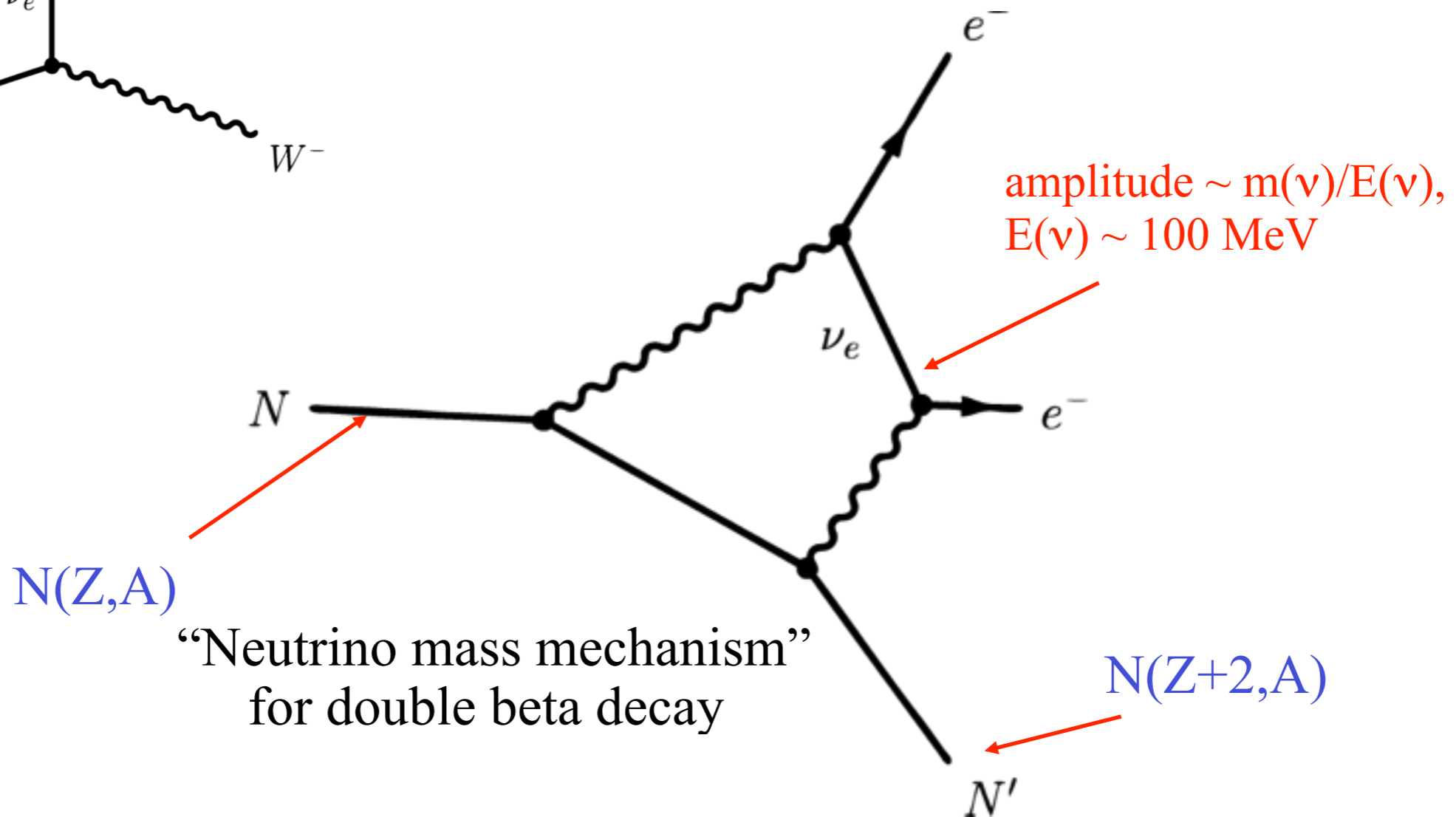
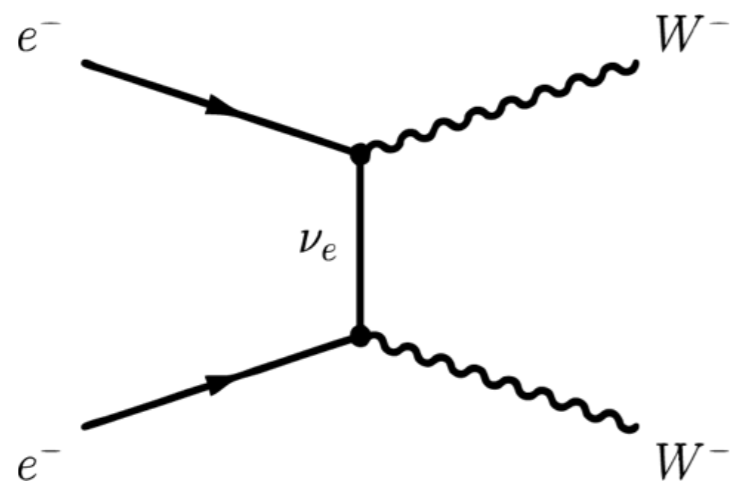
$$(N, Z) \rightarrow (N - 2, Z + 2) + e^- + e^-$$



Virtual W's Instead

Lepton number changes by two units: $\Delta L=2$

$$(N, Z) \rightarrow (N - 2, Z + 2) + e^- + e^-$$



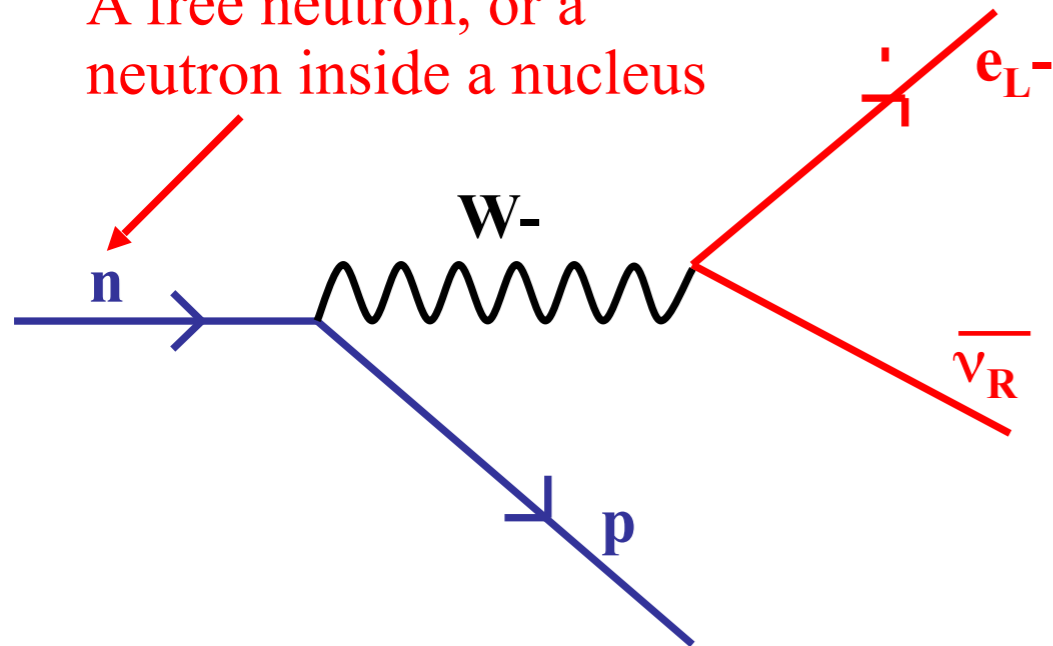
“Neutrino mass mechanism”
for double beta decay

Racah and Furry suggested this was possible for Majorana particles in 1937 soon after Majorana published his theory!

Lepton Number Conserving Standard Model Process

2ν Double Beta Decay

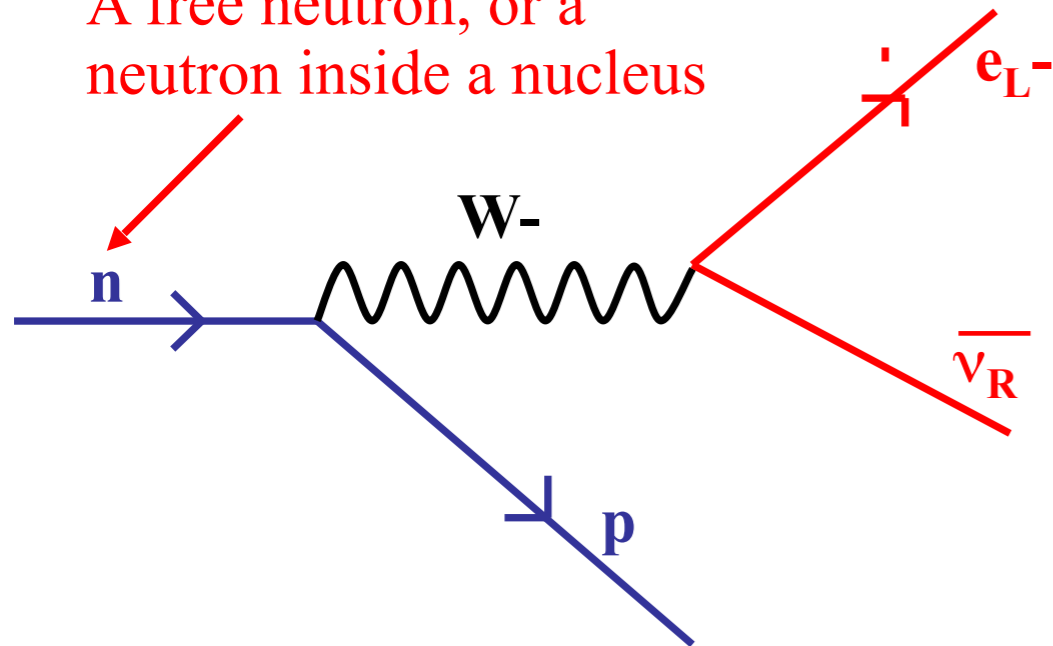
A free neutron, or a
neutron inside a nucleus



Nuclear Beta Decay

2ν Double Beta Decay

A free neutron, or a
neutron inside a nucleus



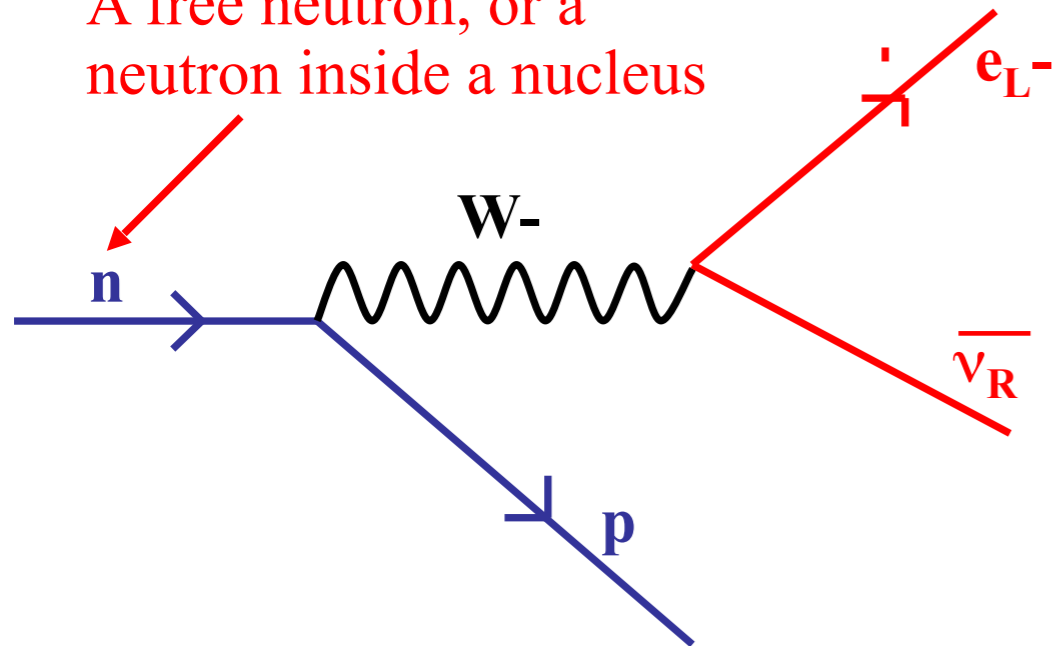
Nuclear Beta Decay

Nuclear Double-Beta
Decay with the emission
of two neutrinos

Lepton Number Conserving Standard Model Process

2ν Double Beta Decay

A free neutron, or a neutron inside a nucleus



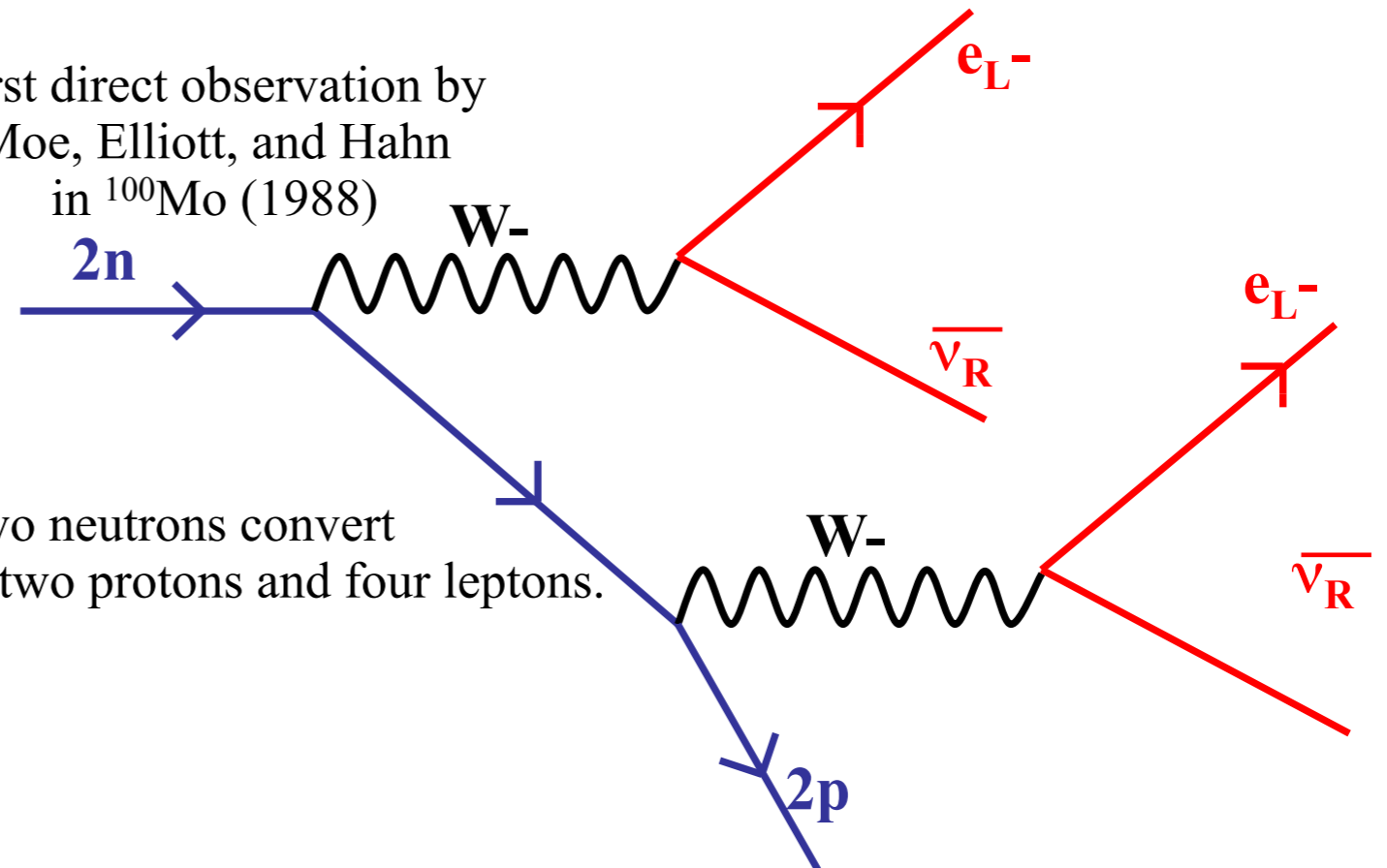
Nuclear Beta Decay

Nuclear Double-Beta Decay with the emission of two neutrinos

Suggested by Maria Goeppert-Mayer in 1935!

First direct observation by Moe, Elliott, and Hahn in ^{100}Mo (1988)

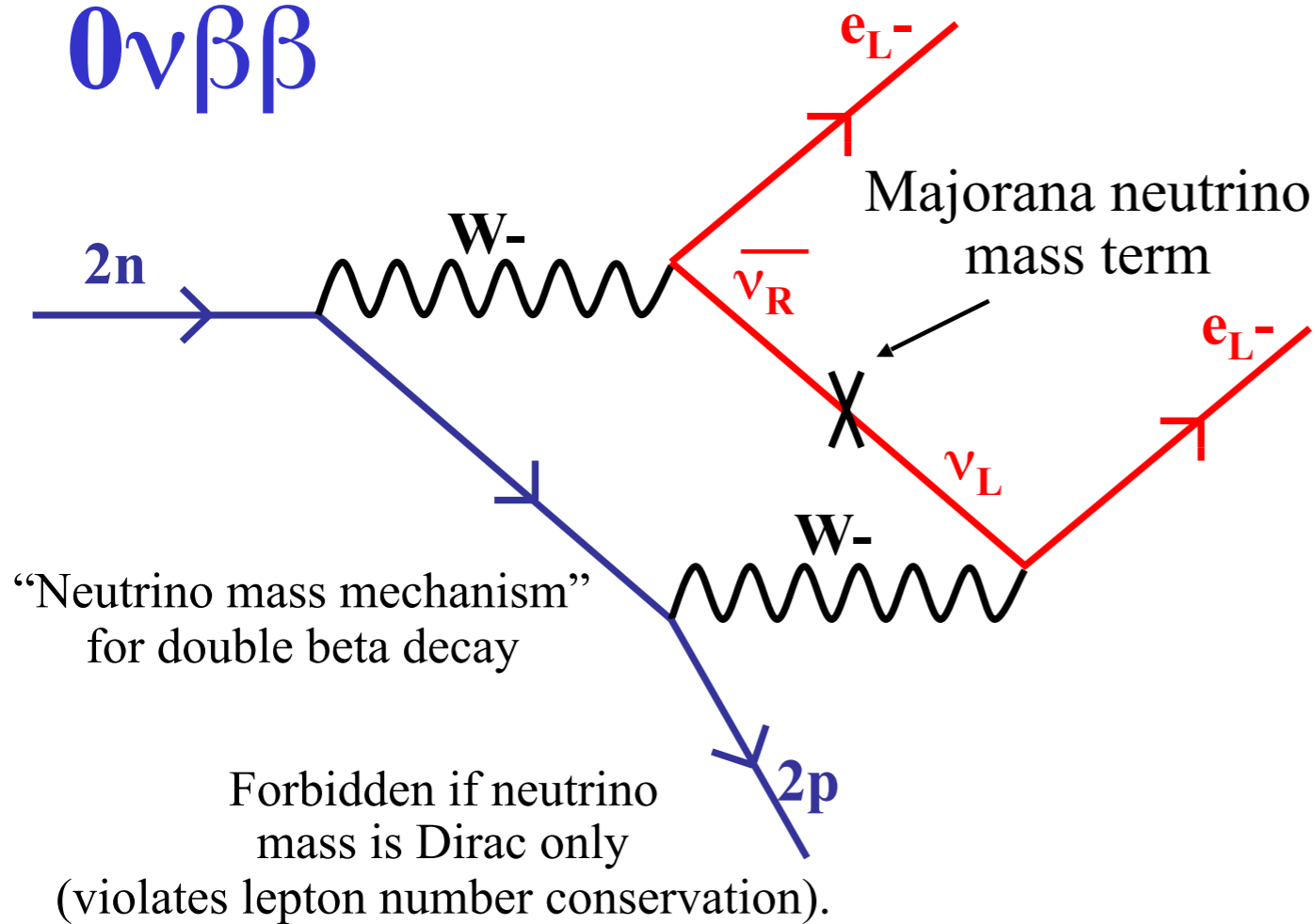
Two neutrons convert to two protons and four leptons.



0ν Double Beta Decay

$$(N, Z) \rightarrow (N - 2, Z + 2) + e^- + e^-$$

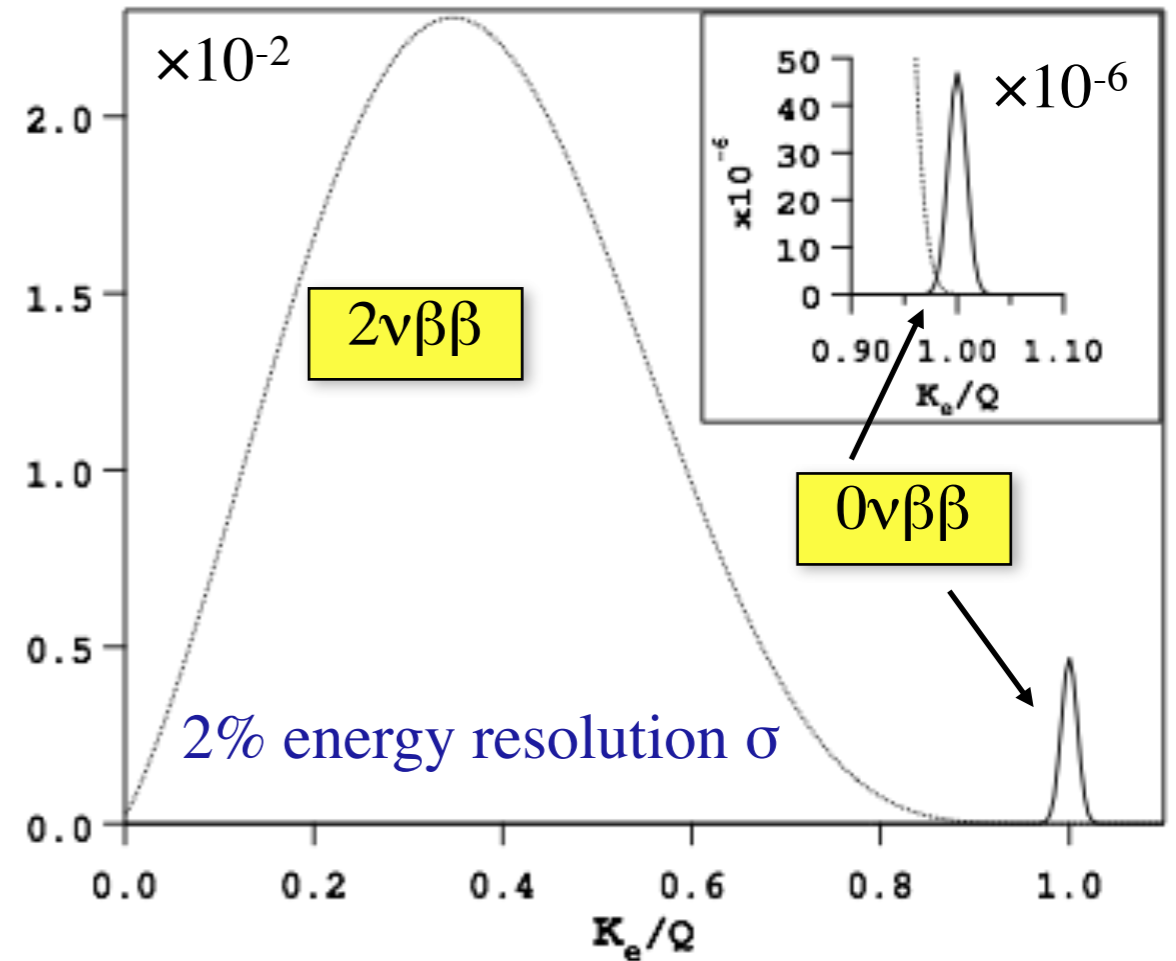
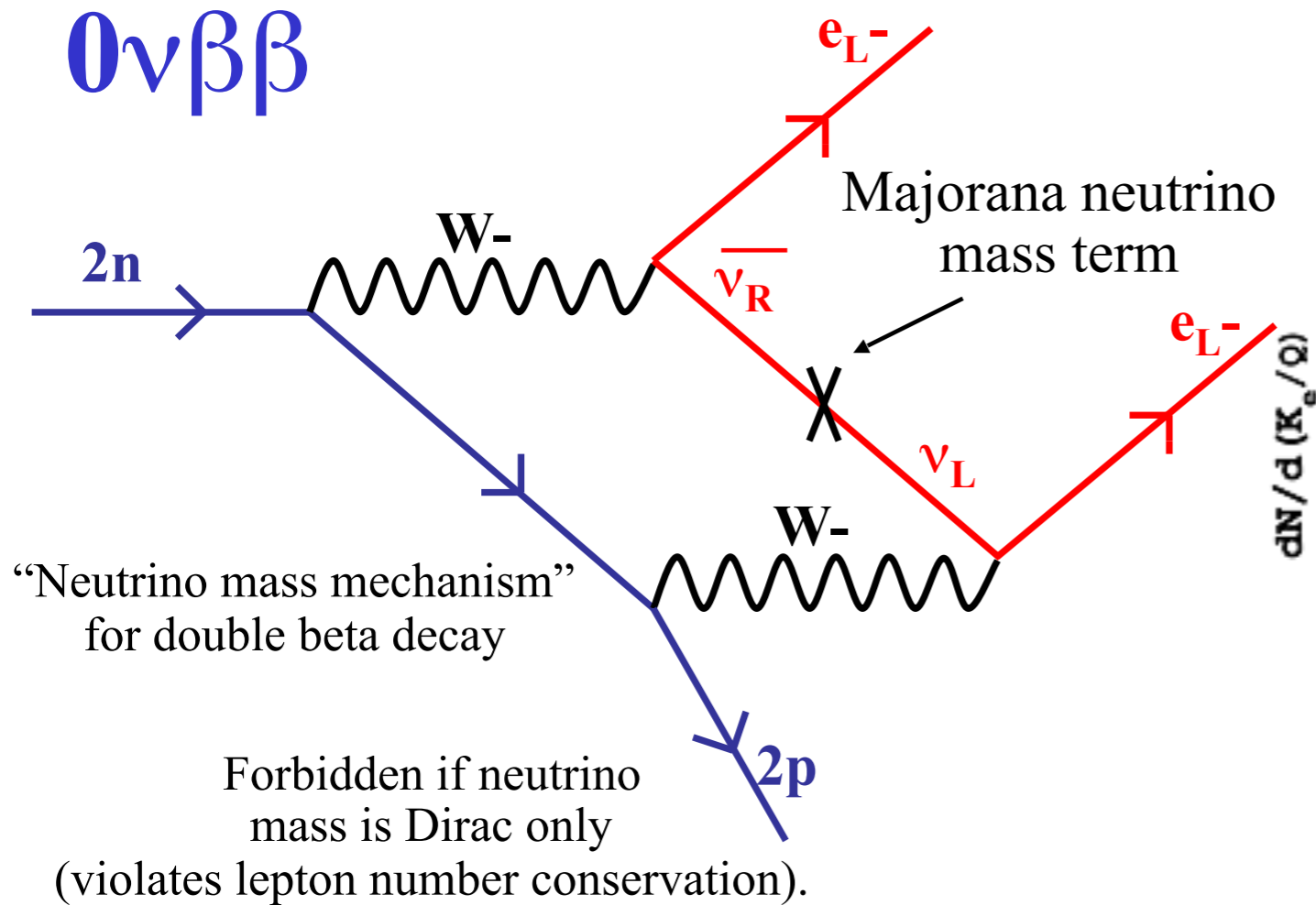
$0\nu\beta\beta$



0ν Double Beta Decay

Experimental Signature

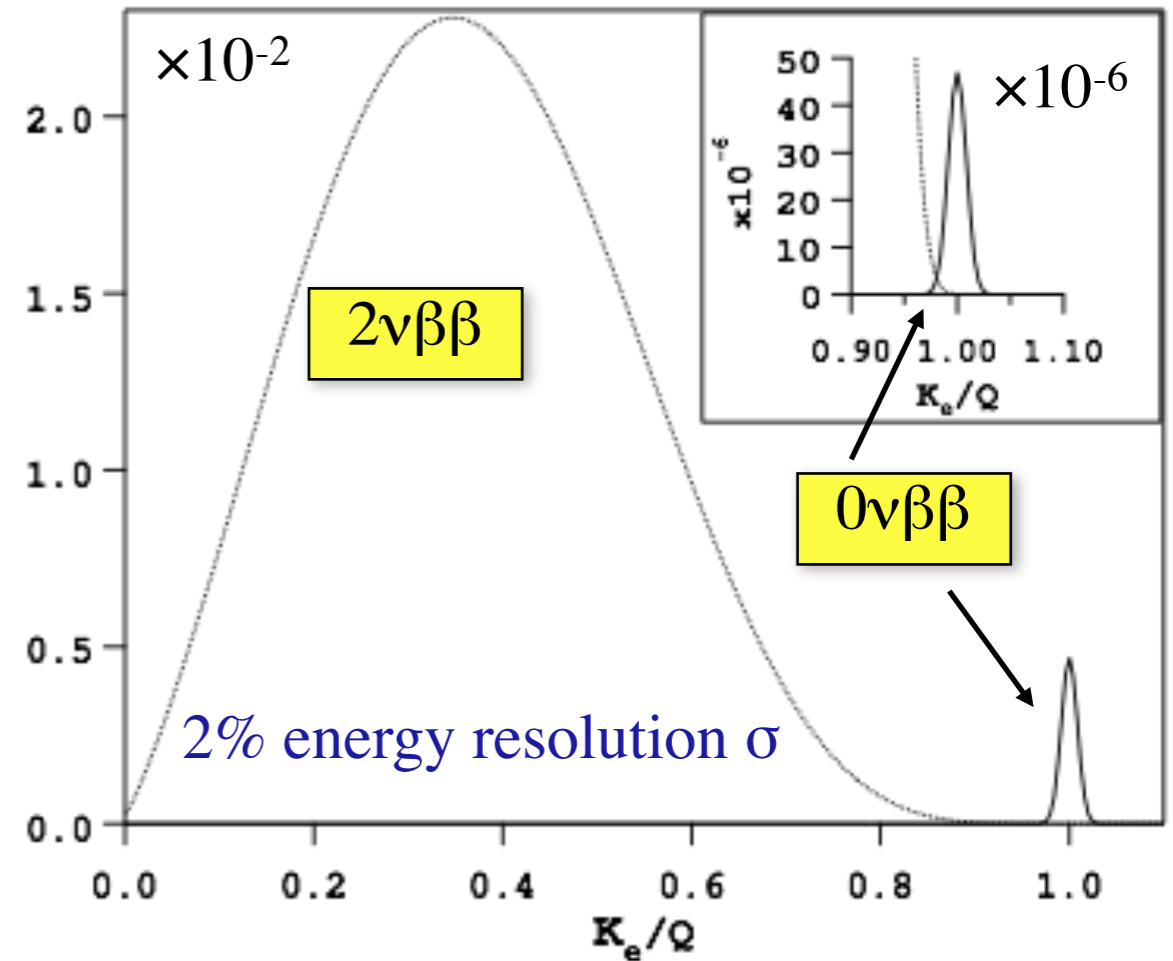
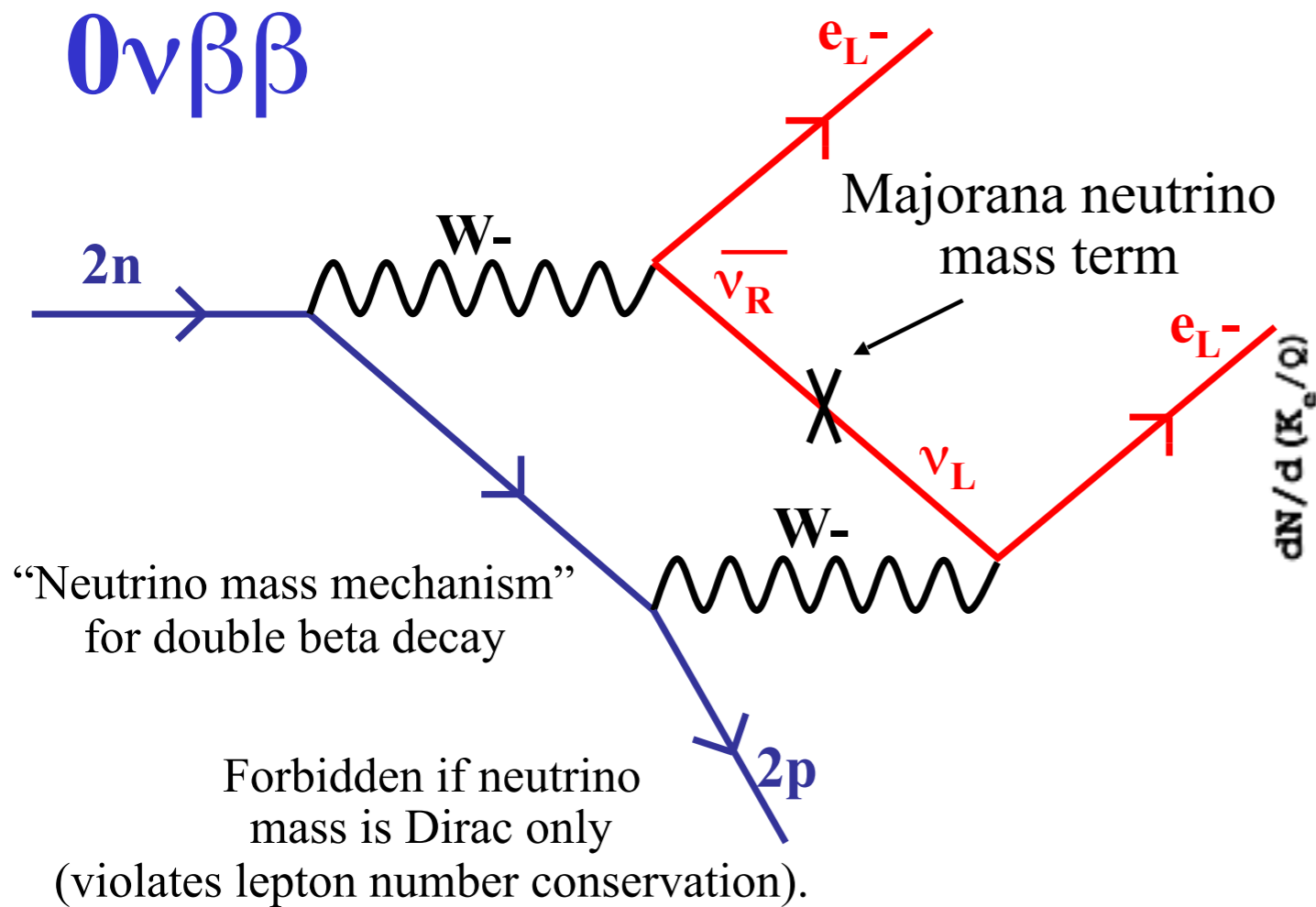
$$(N, Z) \rightarrow (N - 2, Z + 2) + e^- + e^-$$



0ν Double Beta Decay

Experimental Signature

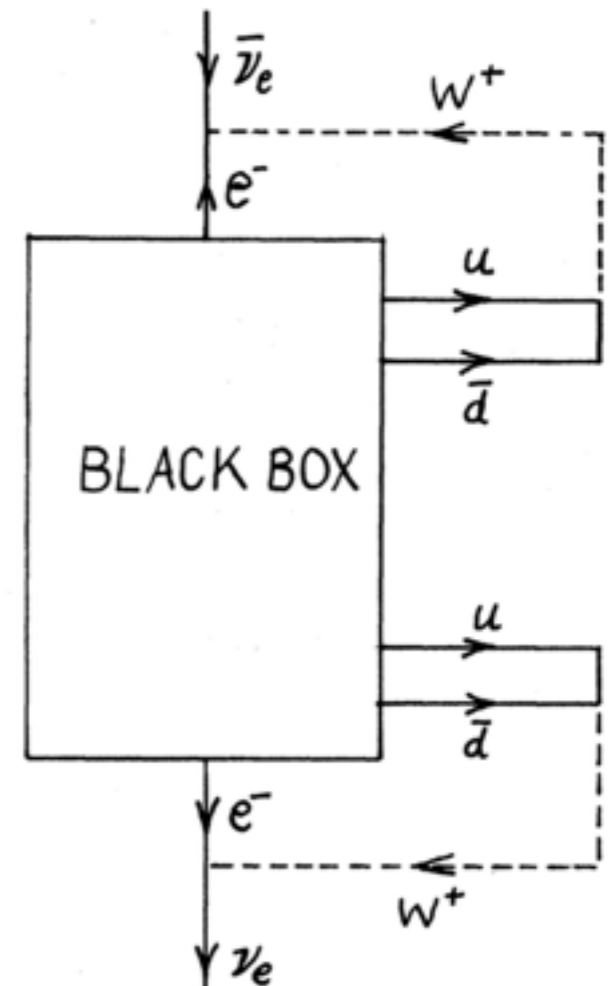
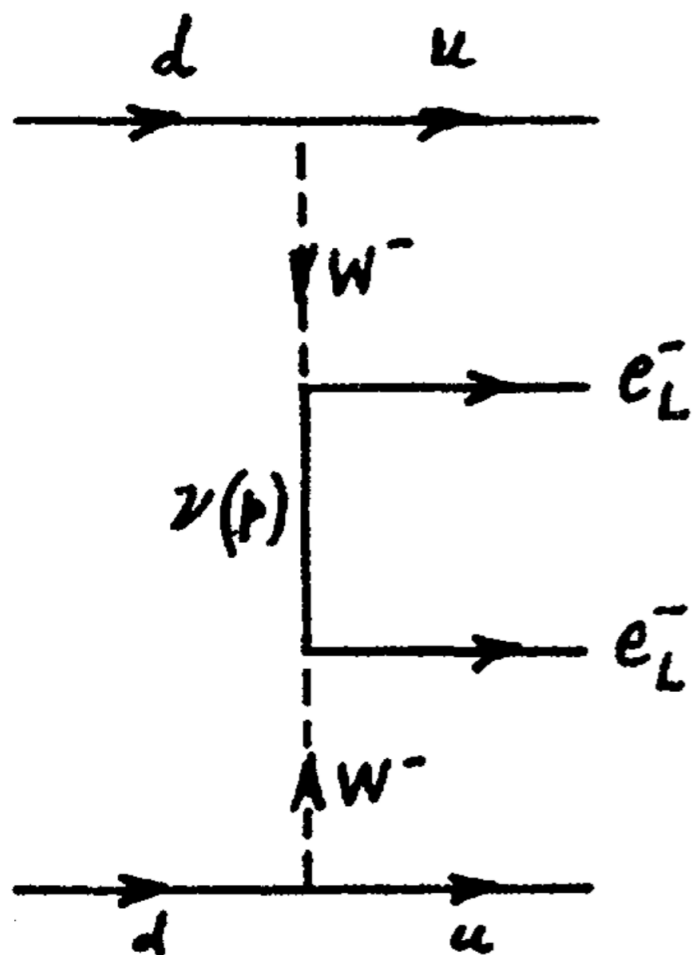
$$(N, Z) \rightarrow (N - 2, Z + 2) + e^- + e^-$$



If observed, it would unambiguously signal that Lepton Number is NOT a conserved quantity, and that neutrinos are Majorana particles i.e. their own anti-particles

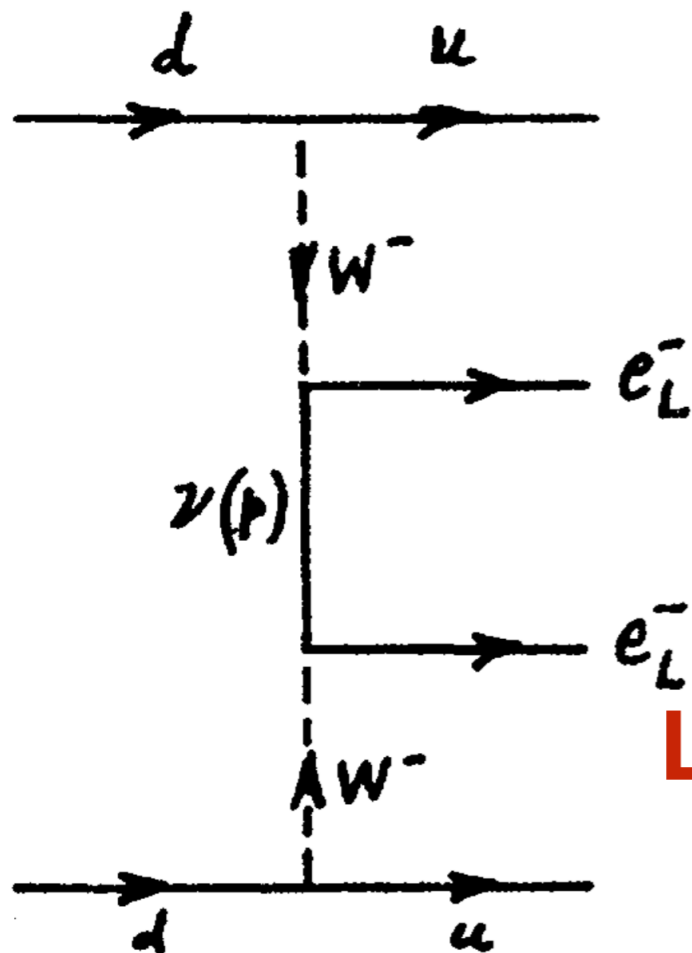
A Theorem

*If neutrinoless double-beta decay occurs, there exists a way to convert an anti-neutrino to a neutrino, a **Majorana mass amplitude***



A Theorem

*If neutrinoless double-beta decay occurs, there exists a way to convert an anti-neutrino to a neutrino, a **Majorana mass amplitude***

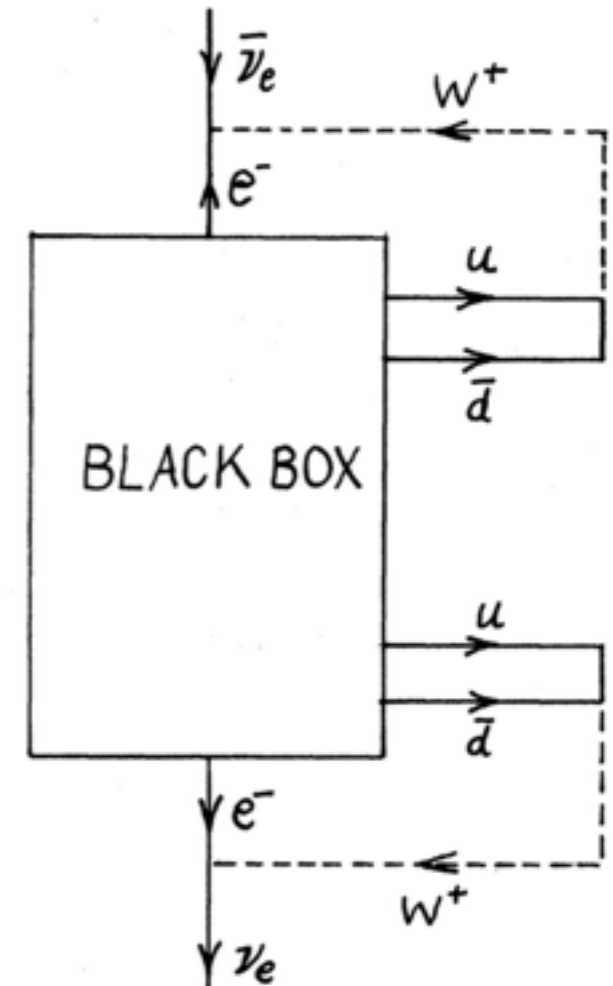


No caveats:

$0\nu\beta\beta$



**Lepton Number Violation
and
Majorana Neutrinos**



Choosing a Nuclide

**Typical $2\nu\beta\beta$ half-life is very long:
second-order weak process**

$$\frac{1}{T_{\frac{1}{2}}^{0\nu}} = G^{2\nu}(Q, Z) |M^{2\nu}|^2$$

$$\frac{1}{G^{2\nu}} \simeq 10^{20} \text{ years}$$

Choosing a Nuclide

Typical $2\nu\beta\beta$ half-life is very long:
second-order weak process

$$\frac{1}{T_{\frac{1}{2}}^{0\nu}} = G^{2\nu}(Q, Z) |M^{2\nu}|^2$$

$$\frac{1}{G^{2\nu}} \simeq 10^{20} \text{ years}$$

Choose nuclei where single beta decay forbidden

**but double-beta
decay is possible**

Choosing a Nuclide

**Typical $2\nu\beta\beta$ half-life is very long:
second-order weak process**

$$\frac{1}{T_{\frac{1}{2}}^{0\nu}} = G^{2\nu}(Q, Z) |M^{2\nu}|^2$$

*Atomic mass affected by nuclear pairing term:
even A nuclei occupy 2 parabolas,
even-even below odd-odd*

$$\frac{1}{G^{2\nu}} \simeq 10^{20} \text{ years}$$

Choose nuclei where single beta decay forbidden

**but double-beta
decay is possible**

Choosing a Nuclide

**Typical $2\nu\beta\beta$ half-life is very long:
second-order weak process**

$$\frac{1}{T_{\frac{1}{2}}^{0\nu}} = G^{2\nu}(Q, Z) |M^{2\nu}|^2$$

*Atomic mass affected by nuclear pairing term:
even A nuclei occupy 2 parabolas,
even-even below odd-odd*

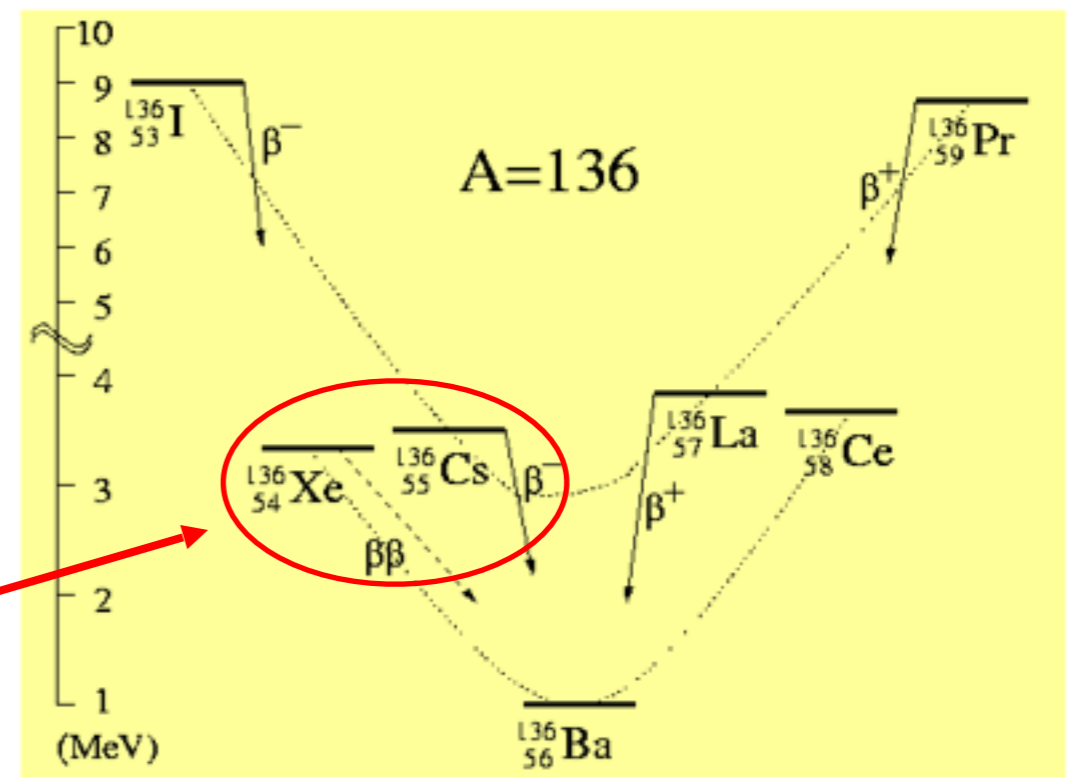
$$\frac{1}{G^{2\nu}} \simeq 10^{20} \text{ years}$$

Choose nuclei where single beta decay forbidden

**but double-beta
decay is possible**

Double-beta decay:

*a second-order process
only detectable if first
order beta decay is
energetically forbidden*



Choosing a Nuclide

Typical $2\nu\beta\beta$ half-life is very long:
second-order weak process

$$\frac{1}{T_{\frac{1}{2}}^{0\nu}} = G^{2\nu}(Q, Z) |M^{2\nu}|^2$$

Atomic mass affected by nuclear pairing term:
even A nuclei occupy 2 parabolas,
even-even below odd-odd

$$\frac{1}{G^{2\nu}} \simeq 10^{20} \text{ years}$$

Candidate Q Abund.
(MeV) (%)

$^{48}\text{Ca} \rightarrow ^{48}\text{Ti}$	4.271	0.187
$^{76}\text{Ge} \rightarrow ^{76}\text{Se}$	2.040	7.8
$^{82}\text{Se} \rightarrow ^{82}\text{Kr}$	2.995	9.2
$^{96}\text{Zr} \rightarrow ^{96}\text{Mo}$	3.350	2.8
$^{100}\text{Mo} \rightarrow ^{100}\text{Ru}$	3.034	9.6
$^{110}\text{Pd} \rightarrow ^{110}\text{Cd}$	2.013	11.8
$^{116}\text{Cd} \rightarrow ^{116}\text{Sn}$	2.802	7.5
$^{124}\text{Sn} \rightarrow ^{124}\text{Te}$	2.228	5.64
$^{130}\text{Te} \rightarrow ^{130}\text{Xe}$	2.533	34.5
$^{136}\text{Xe} \rightarrow ^{136}\text{Ba}$	2.479	8.9
$^{150}\text{Nd} \rightarrow ^{150}\text{Sm}$	3.367	5.6

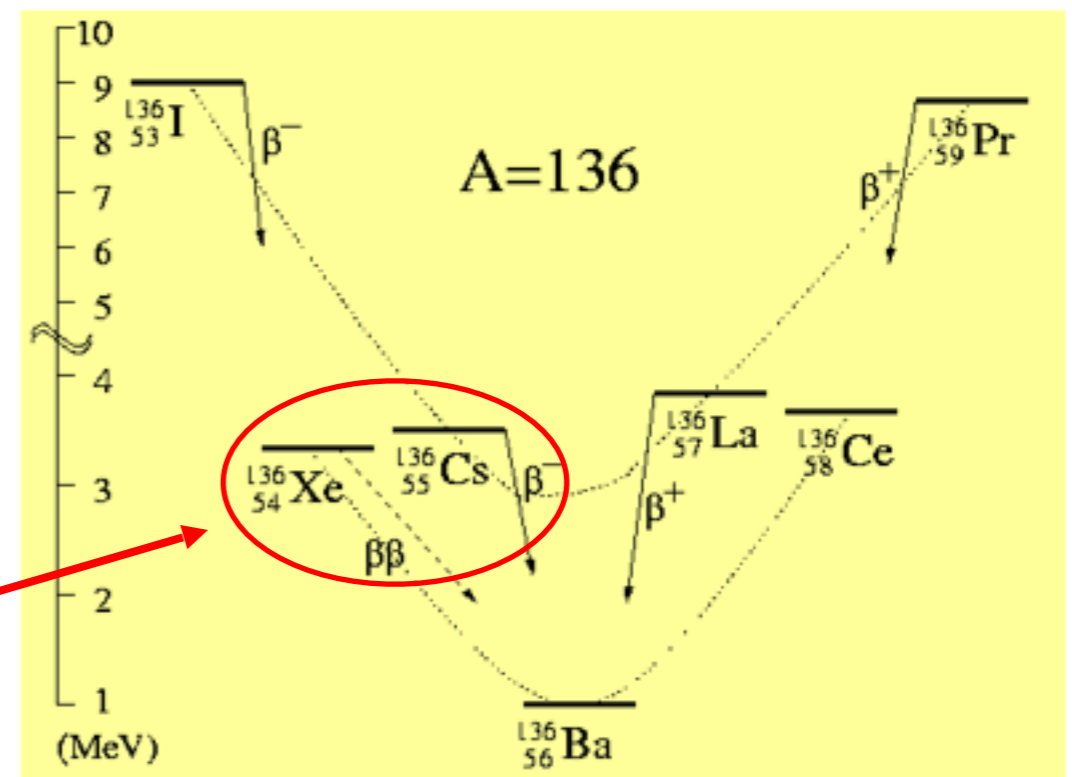
Choose nuclei where single beta decay forbidden

but double-beta decay is possible

Candidate nuclei with $Q > 2$ MeV

Double-beta decay:

a second-order process only detectable if first order beta decay is energetically forbidden



Decay Rate for $0\nu\beta\beta$

$$\Gamma^{0\nu} = G(Q, Z) |M(A, Z)\eta|^2$$

Transition
Probability

$$\propto \frac{m}{Q^2} \quad (Q \sim m_e)$$

Phase Space
Factor $G \sim G_F^4 g_A^4 m_e^5$

Nuclear Matrix
Element

$$M(A, Z)$$

Particle Physics
of the Black Box η

Decay Rate for $0\nu\beta\beta$

$$\Gamma^{0\nu} = G(Q, Z) |M(A, Z) \eta|^2$$

Transition
Probability

$$\propto \frac{m}{Q^2} \quad (Q \sim m_e)$$

Phase Space
Factor

$$G \sim G_F^4 g_A^4 m_e^5$$

Nuclear Matrix
Element

$$M(A, Z)$$

Particle Physics
of the Black Box

$$\eta$$

For light neutrino exchange

All 3 neutrinos will contribute: $\eta \sim m \rightarrow \langle m_{\beta\beta} \rangle = \sum_i U_{ie}^2 m_i$

PMNS Matrix



Decay Rate for $0\nu\beta\beta$

$$\Gamma^{0\nu} = G(Q, Z) |M(A, Z) \eta|^2$$

Transition
Probability

$$\propto \frac{m}{Q^2} \quad (Q \sim m_e)$$

Phase Space
Factor

$$G \sim G_F^4 g_A^4 m_e^5$$

Nuclear Matrix
Element

$$M(A, Z)$$

Particle Physics
of the Black Box

$$\eta$$

For light neutrino exchange

All 3 neutrinos will contribute: $\eta \sim m \rightarrow \langle m_{\beta\beta} \rangle = \sum_i U_{ie}^2 m_i$

$$m_{\beta\beta} \sim 1 \text{ eV} \implies T_{1/2} \sim 10^{24} \text{ years}$$

$$m_{\beta\beta} \sim 0.1 \text{ eV} \implies T_{1/2} \sim 10^{26} \text{ years}$$

$$m_{\beta\beta} \sim 0.01 \text{ eV} \implies T_{1/2} \sim 10^{28} \text{ years}$$

PMNS Matrix



$$\sum_i U_{ie}^2 m_i$$

Decay Rate for $0\nu\beta\beta$

$$\Gamma^{0\nu} = G(Q, Z) |M(A, Z) \eta|^2$$

Transition Probability

$$\propto \frac{m}{Q^2} \quad (Q \sim m_e)$$

Phase Space Factor $G \sim G_F^4 g_A^4 m_e^5$

Nuclear Matrix Element

$$M(A, Z)$$

Particle Physics of the Black Box η

For light neutrino exchange

All 3 neutrinos will contribute: $\eta \sim m \rightarrow \langle m_{\beta\beta} \rangle = \sum_i U_{ie}^2 m_i$

PMNS Matrix



~ 10 kg $m_{\beta\beta} \sim 1$ eV $\implies T_{1/2} \sim 10^{24}$ years

Ruled out

~ 100 kg $m_{\beta\beta} \sim 0.1$ eV $\implies T_{1/2} \sim 10^{26}$ years

Current sensitivity

~ 1000 kg $m_{\beta\beta} \sim 0.01$ eV $\implies T_{1/2} \sim 10^{28}$ years

Next Generation

Decay Rate for $0\nu\beta\beta$

$$\Gamma^{0\nu} = G(Q, Z) |M(A, Z) \eta|^2$$

Transition Probability

$$\propto \frac{m}{Q^2} \quad (Q \sim m_e)$$

Phase Space Factor $G \sim G_F^4 g_A^4 m_e^5$

Nuclear Matrix Element

$$M(A, Z)$$

Particle Physics of the Black Box η

For light neutrino exchange

← BUT.....

PMNS Matrix

All 3 neutrinos will contribute: $\eta \sim m \rightarrow \langle m_{\beta\beta} \rangle = \sum_i U_{ie}^2 m_i$

$\sim 10 \text{ kg} \quad m_{\beta\beta} \sim 1 \text{ eV} \implies T_{1/2} \sim 10^{24} \text{ years}$

Ruled out

$\sim 100 \text{ kg} \quad m_{\beta\beta} \sim 0.1 \text{ eV} \implies T_{1/2} \sim 10^{26} \text{ years}$

Current sensitivity

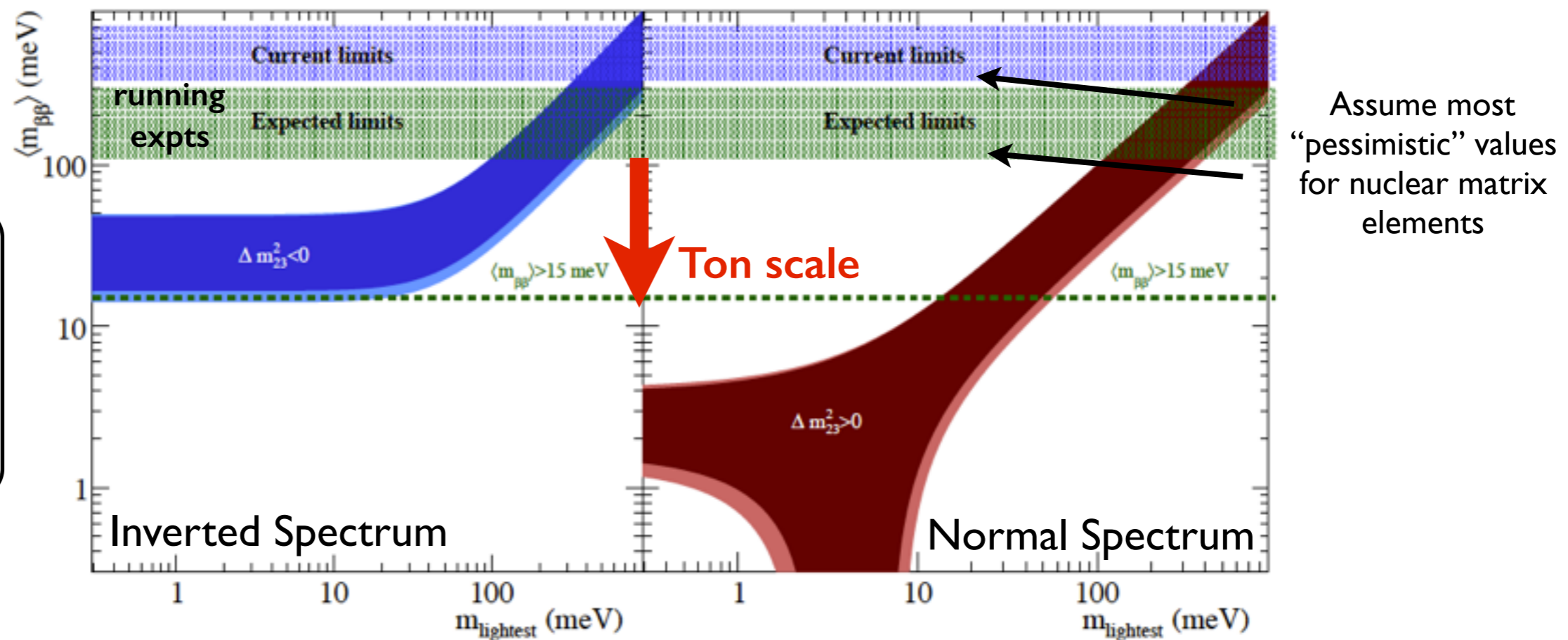
$\sim 1000 \text{ kg} \quad m_{\beta\beta} \sim 0.01 \text{ eV} \implies T_{1/2} \sim 10^{28} \text{ years}$

Next Generation

Various Possibilities for the Black Box

V. Cirigliano

- ton-scale $0\nu\beta\beta$ probes LNV from variety mechanisms, involving different scales (M) and coupling strengths (g)

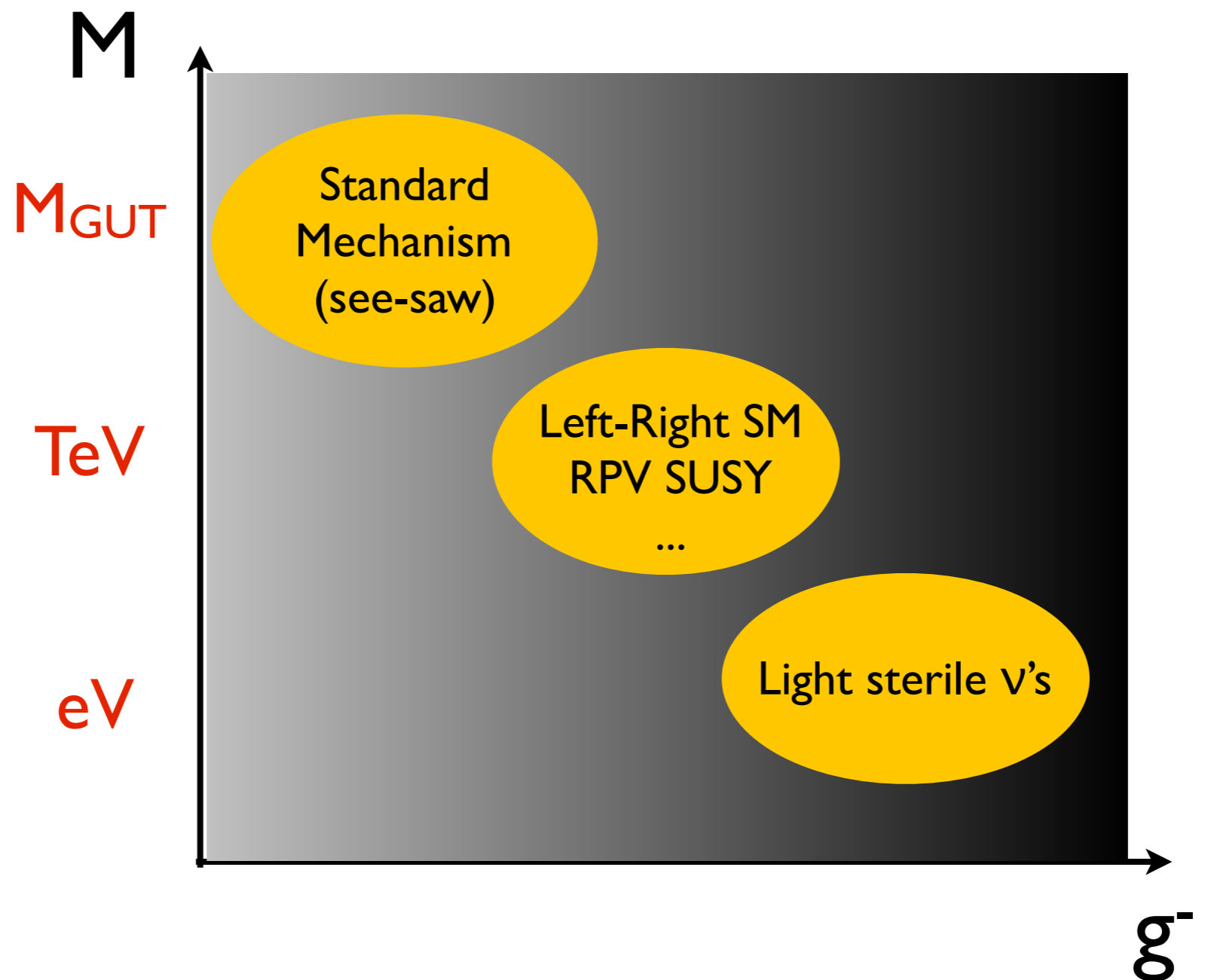


- Discovery possible for **inverted spectrum** OR **$m_{\text{lightest}} > 50$ meV**

Various Possibilities for the Black Box

V. Cirigliano

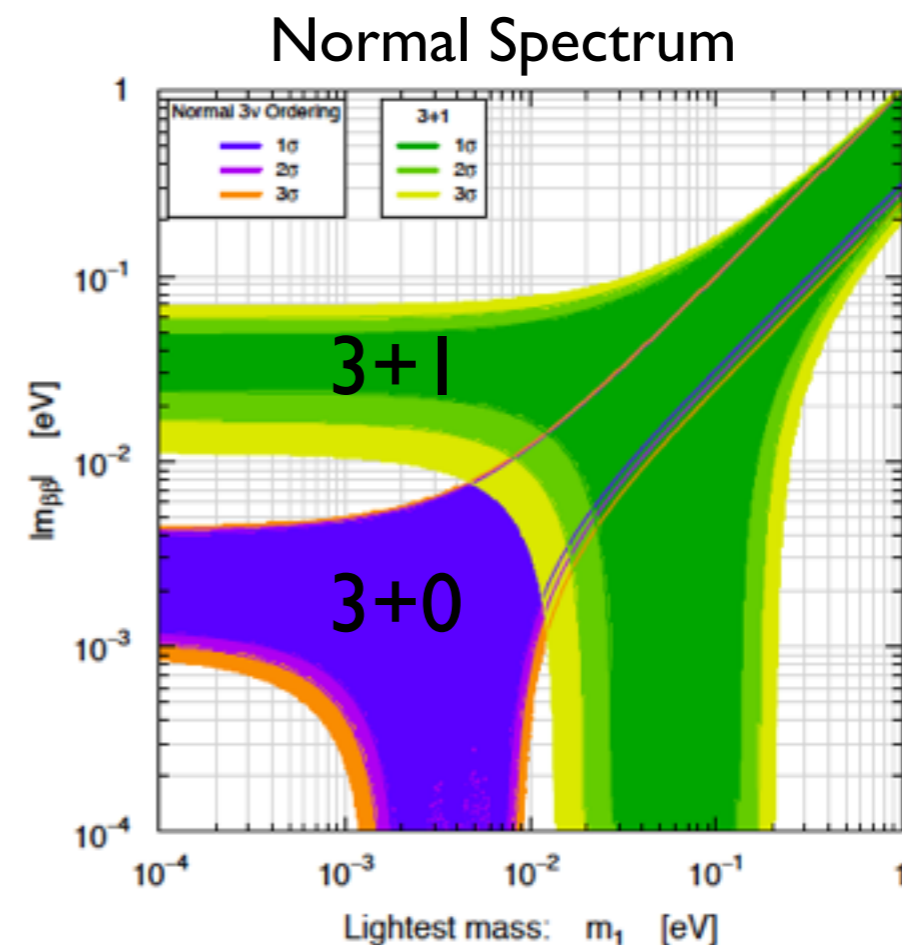
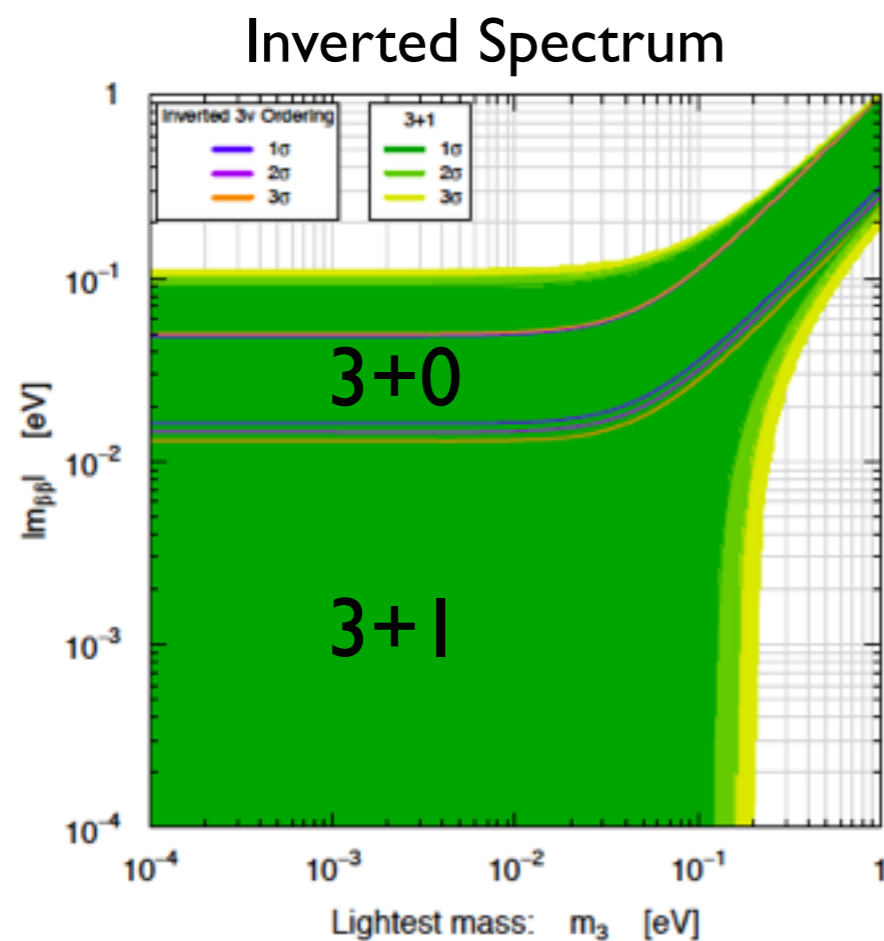
- ton-scale $0\nu\beta\beta$ probes LNV from variety mechanisms, involving different scales (M) and coupling strengths (g)



Various Possibilities for the Black Box

V. Cirigliano

- **Low scale seesaw**: intriguing example with one light sterile ν_R with mass ($\sim eV$) and mixing (~ 0.1) to fit short baseline anomalies
- Extra contribution to effective mass



Giunti-Zavanin
2015

Usual phenomenology turned around!!

Neutrinoless Double Beta Decay: Experimental Status

Signal and Background

An experimental challenge of rare events

Most measured half-lives of $2\nu\beta\beta$ are $O(10^{21})$ years

- Compare to lifetime of Universe: 10^{10} years
- Compare to Avogadro's number 6×10^{23}
- Mole of isotope will produce ~ 1 decay/day

If it exists, half-lives of $0\nu\beta\beta$ would be longer
(^{136}Xe limits is $> 10^{25}$ years)

Half life (years)	Signal (cts/tonne-year)
10^{25}	500
5×10^{26}	10
5×10^{27}	1
5×10^{28}	0.1

Natural radioactivity: a nanogram produces more than 1 decay/day!

Cosmogenically induced radioactivity exacerbates technical challenge

$$\left[T_{1/2}^{0\nu} \right] \propto \epsilon_{ff} \cdot I_{abundance} \cdot \text{Source Mass} \cdot \text{Time}$$

background free

$$\left[T_{1/2}^{0\nu} \right] \propto \epsilon_{ff} \cdot I_{abundance} \cdot \sqrt{\frac{\text{Source Mass} \cdot \text{Time}}{\text{Bkg} \cdot \Delta E}}$$

background limited

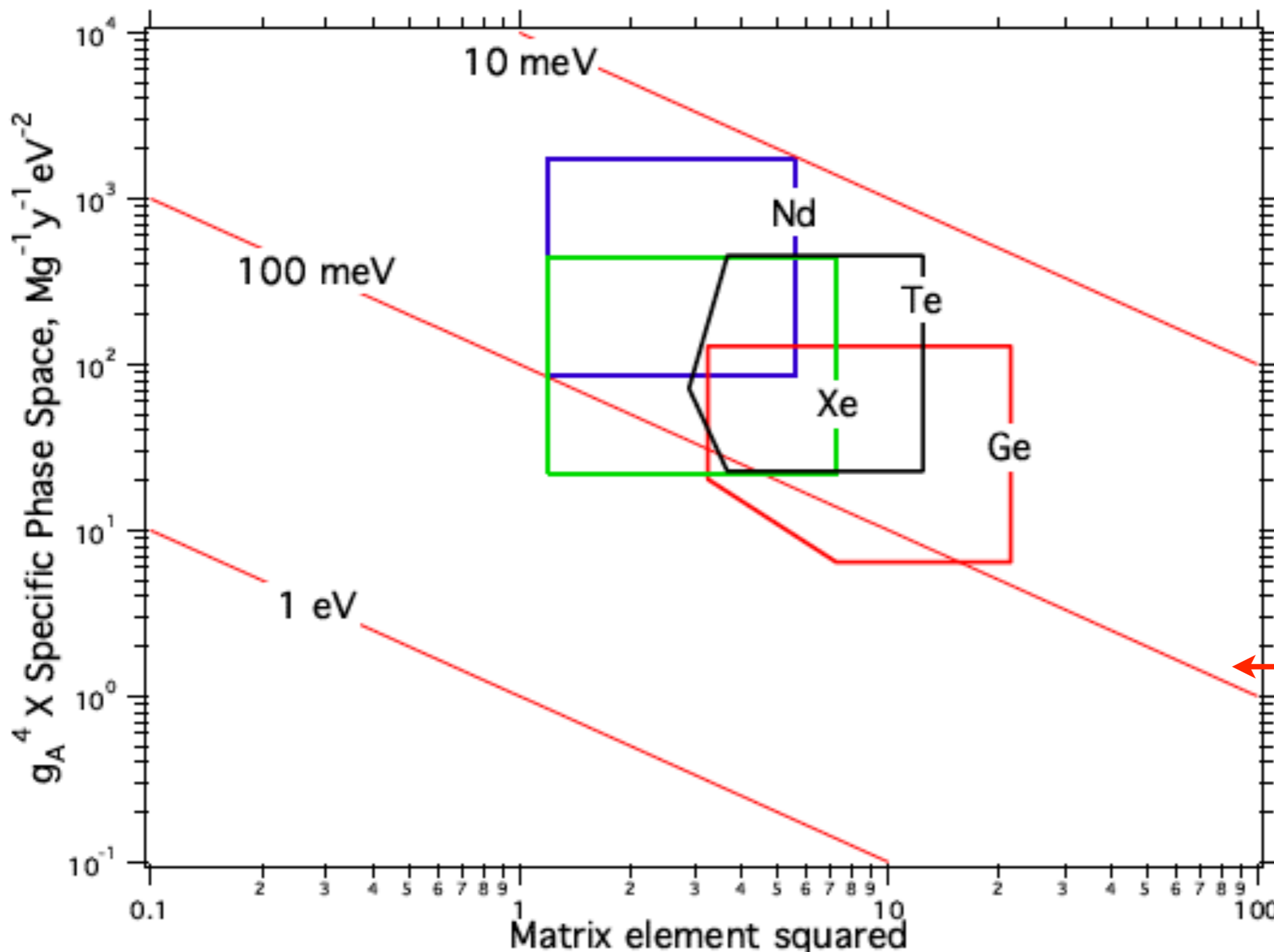
backgrounds do not always scale with detector mass

Favorite Isotope?

For Ge, Te, Xe, Nd

← uncertainty on NME^2 →

R.G.H. Robertson, MPL A
28 (2013) 1350021
 (arXiv 1301.1323)



↑
 uncertainty on
 value of g_A^4
 ↓

← Signal of
 1 cnt/t-y for
 corresponding
 values of NME
 and g_A

The Experimental Challenge

$0\nu\beta\beta$ source with
high isotopic abundance

Detector with
high detection efficiency
good energy resolution
low-background

Experiment
long exposure time
large total mass of isotope

To reach IH region requires sensitivities of

$0\nu\beta\beta$ $T_{1/2} \sim 10^{27} - 10^{28}$ years

$(2\nu\beta\beta$ $T_{1/2} \sim 10^{19} - 10^{21}$ years)

$$T_{1/2}^{0\nu} \text{ sensitivity} \propto a \cdot \epsilon \sqrt{\frac{M \cdot t}{b \cdot \delta E}}$$

a = source isotopic abundance

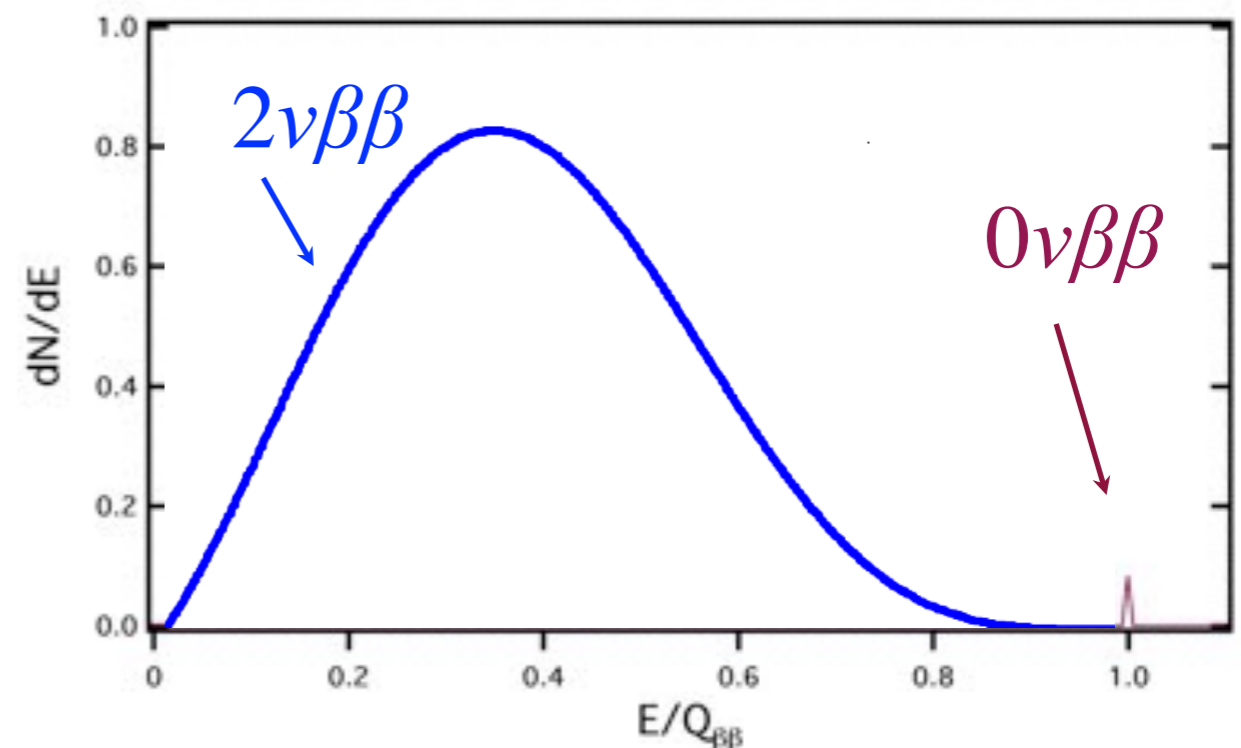
ϵ = detection efficiency

M = total mass

t = exposure time

b = background rate at $0\nu\beta\beta$ energy

δE = energy resolution



Background Strategies

Potential Backgrounds

- Primordial, **natural radioactivity** in detector components: U, Th, K
- Backgrounds from **cosmogenic activation** while material is above ground ($\beta\beta$ -isotope or shield specific, ^{60}Co , ^3H ...)
- Backgrounds from the **surrounding environment**:
 - external γ , (α, n) , (n, α) , Rn plate-out, etc.
- **μ -induced backgrounds** generated at depth:
 - Cu, Pb($n, n' \gamma$), $\beta\beta$ -decay specific(n, n), (n, γ), direct μ
- **2 neutrino double beta decay** (irreducible, E resolution dependent)
- neutrino backgrounds (negligible)

Reduce Backgrounds

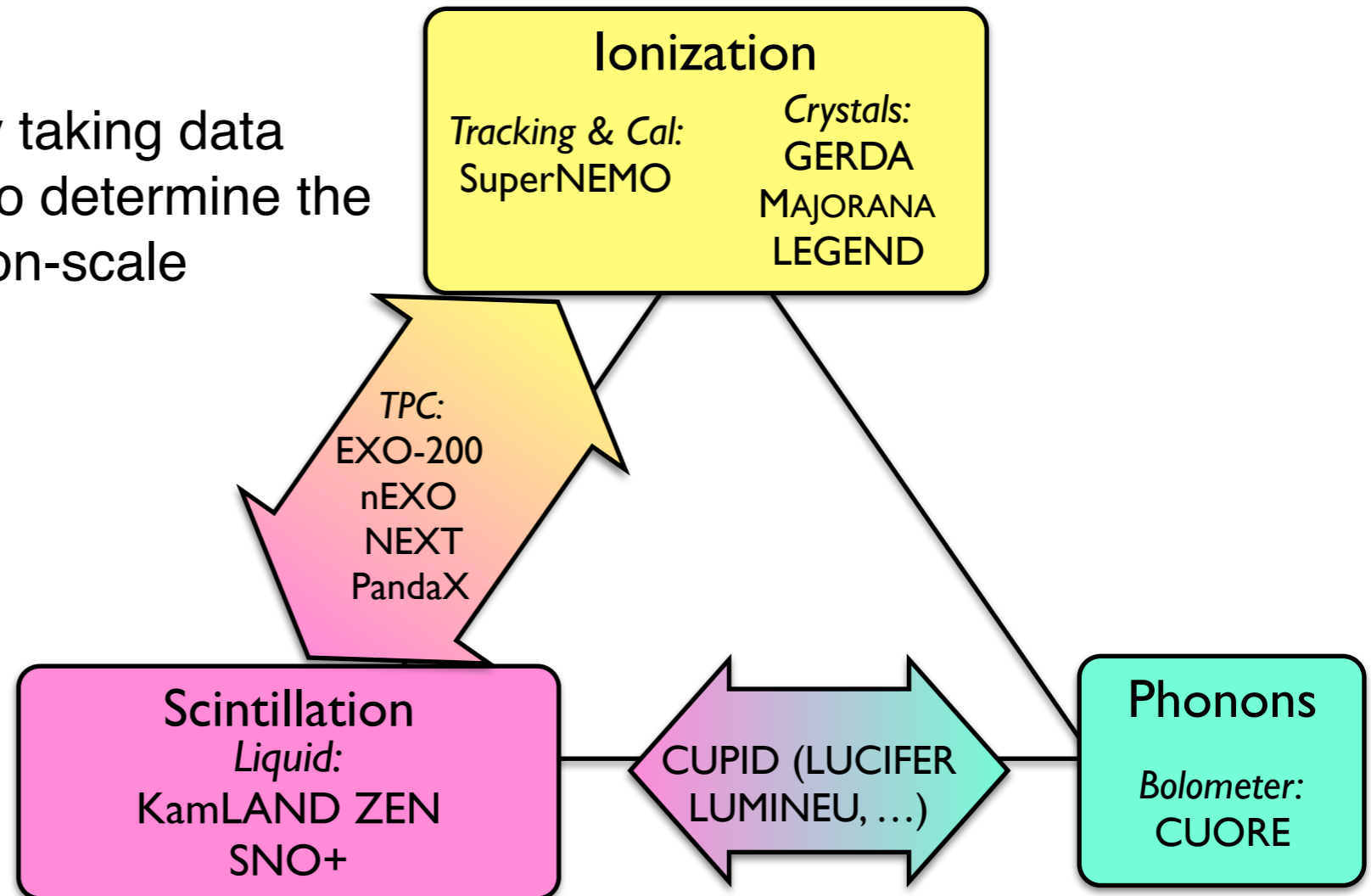
- ultra-pure materials
- shielding
- deep underground
- ...

Discriminate Backgrounds

- energy resolution
- tracking (even topology)
- fiducial fits
- pulse shape discrimination (PSD)
- particle ID
- ...

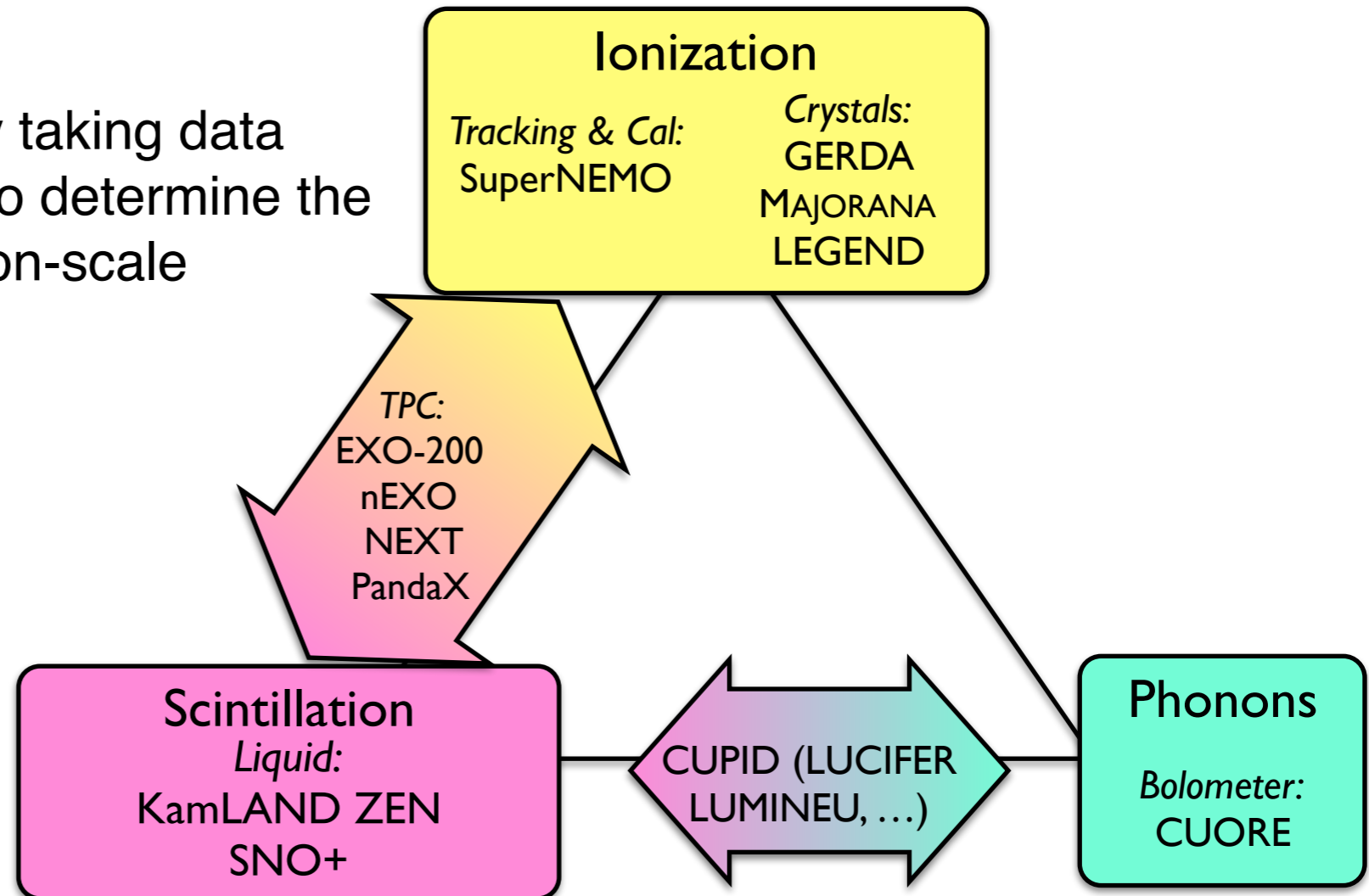
Multi-Prong Detection Strategy

- 100 kg class experiments currently taking data
- In parallel, major R&D under way to determine the optimum path to discovery at the ton-scale



Multi-Prong Detection Strategy

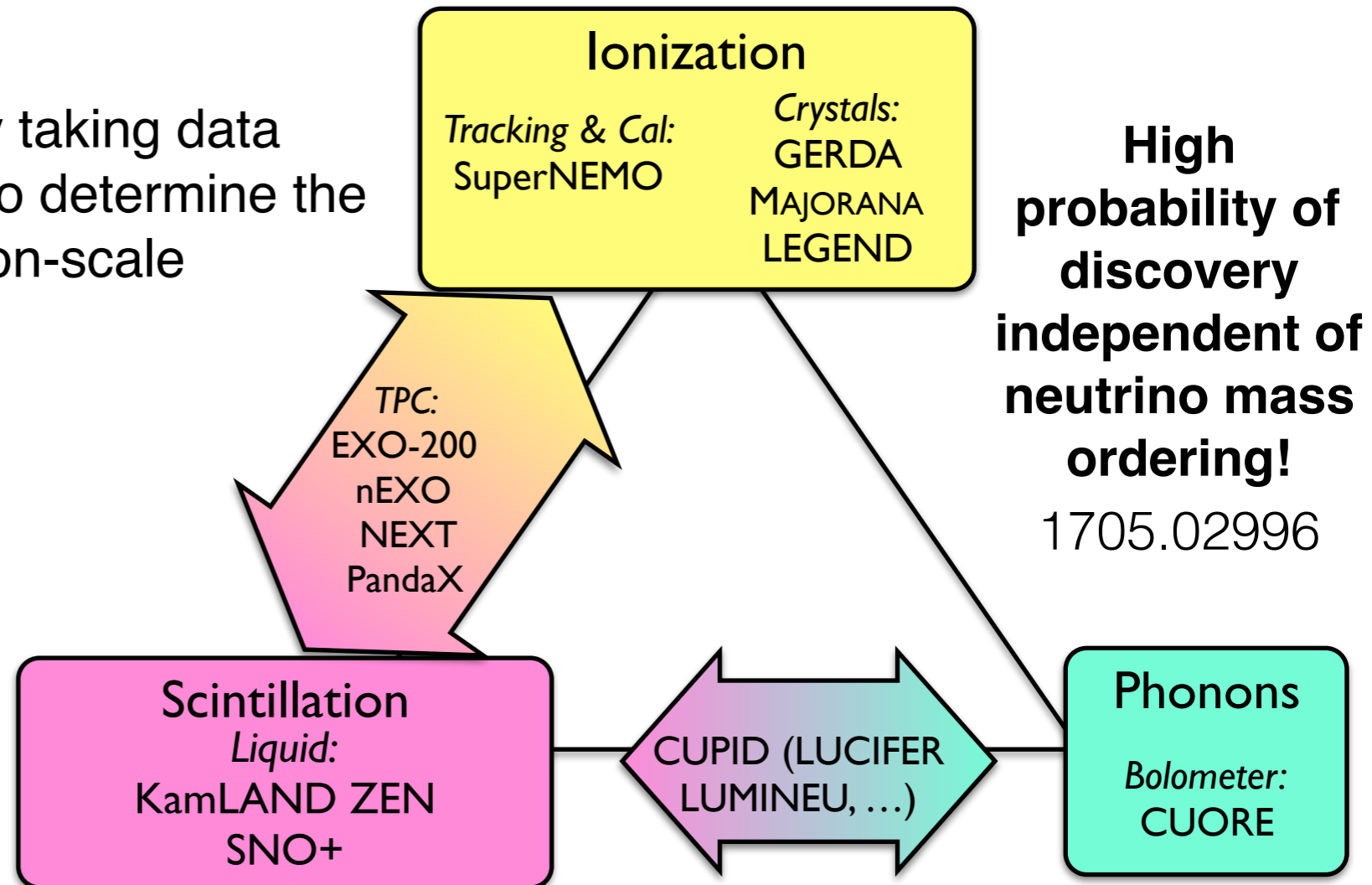
- 100 kg class experiments currently taking data
- In parallel, major R&D under way to determine the optimum path to discovery at the ton-scale



- Ton-scale $0\nu\beta\beta$ searches ($T_{1/2} > 10^{27-28}$ yr) probe at unprecedented levels LNV from a variety of mechanisms
- If light Majorana neutrinos are responsible for $0\nu\beta\beta$, then absolute neutrino mass scale determination within reach of ton-scale experiments

Multi-Prong Detection Strategy

- 100 kg class experiments currently taking data
- In parallel, major R&D under way to determine the optimum path to discovery at the ton-scale



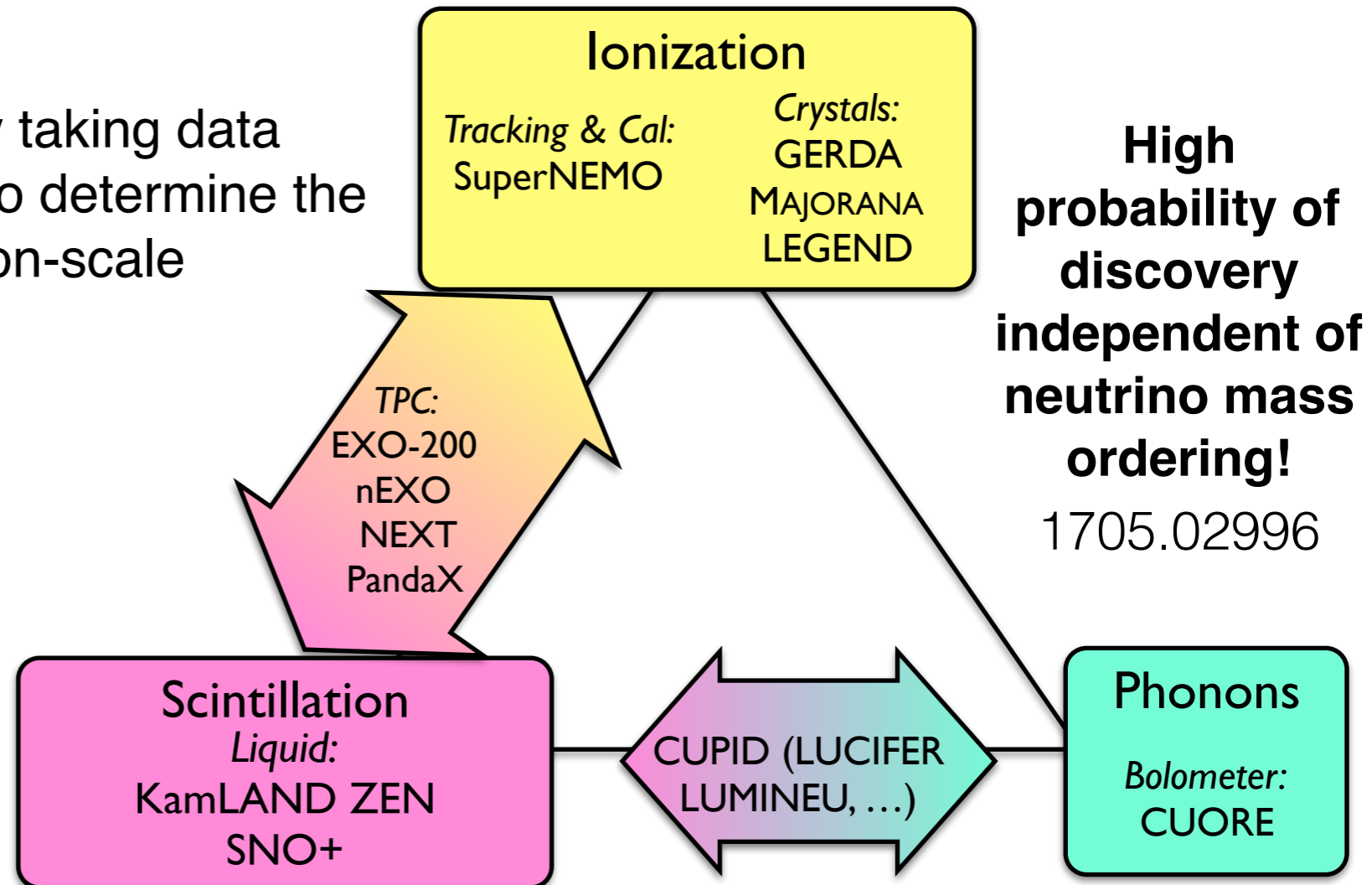
- Ton-scale $0\nu\beta\beta$ searches ($T_{1/2} > 10^{27-28}$ yr) probe at unprecedented levels LNV from a variety of mechanisms
- If light Majorana neutrinos are responsible for $0\nu\beta\beta$, then absolute neutrino mass scale determination within reach of ton-scale experiments

Multi-Prong Detection Strategy

- 100 kg class experiments currently taking data
- In parallel, major R&D under way to determine the optimum path to discovery at the ton-scale

^{76}Ge , ^{130}Te , ^{136}Xe

Leading isotopes for successful ton-scale designs: most promising technologies have source = detector



- Ton-scale $0\nu\beta\beta$ searches ($T_{1/2} > 10^{27-28}$ yr) probe at unprecedented levels LNV from a variety of mechanisms
- If light Majorana neutrinos are responsible for $0\nu\beta\beta$, then absolute neutrino mass scale determination within reach of ton-scale experiments

International Program

Previous Expts.
 $T_{1/2} \sim 10^{24}$ y
 (~ 1 eV)
 ~kg scale



Quasi-degenerate
 $T_{1/2} \sim 10^{25} - 10^{26}$ y
 (~100 meV)
 30 - 200 kg
 ~8 expts

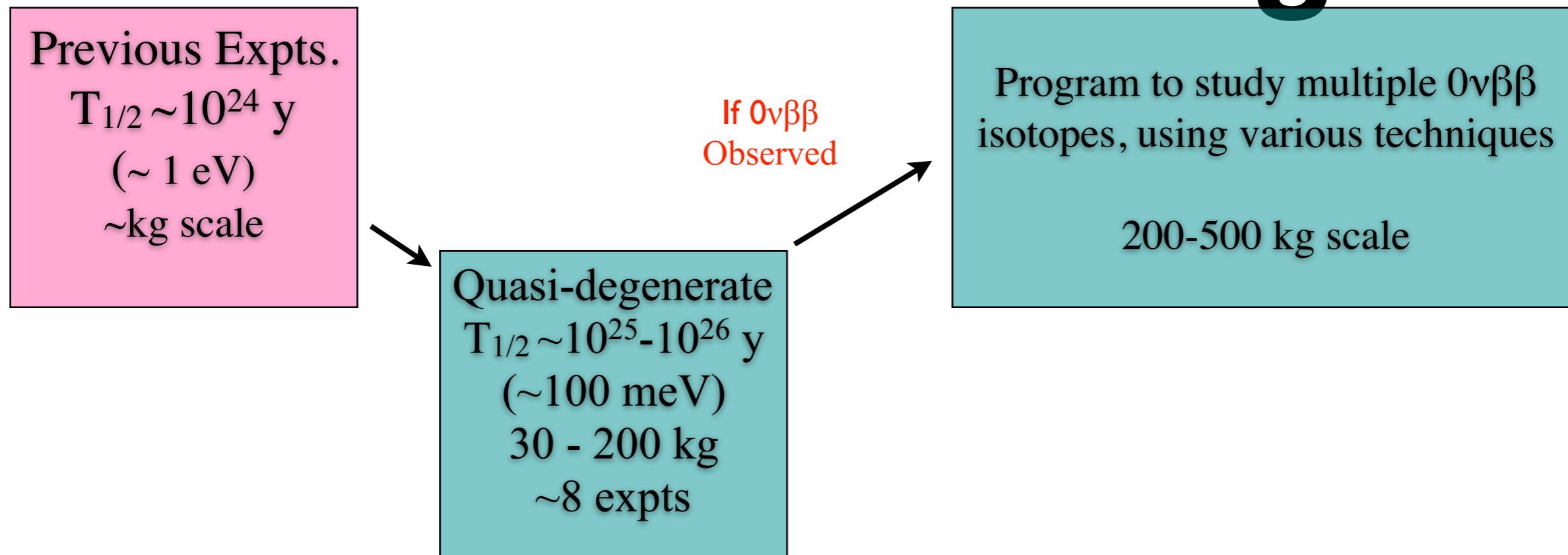
1980 - 2007

2007 - 2019

2019 - mid-2020's



International Program



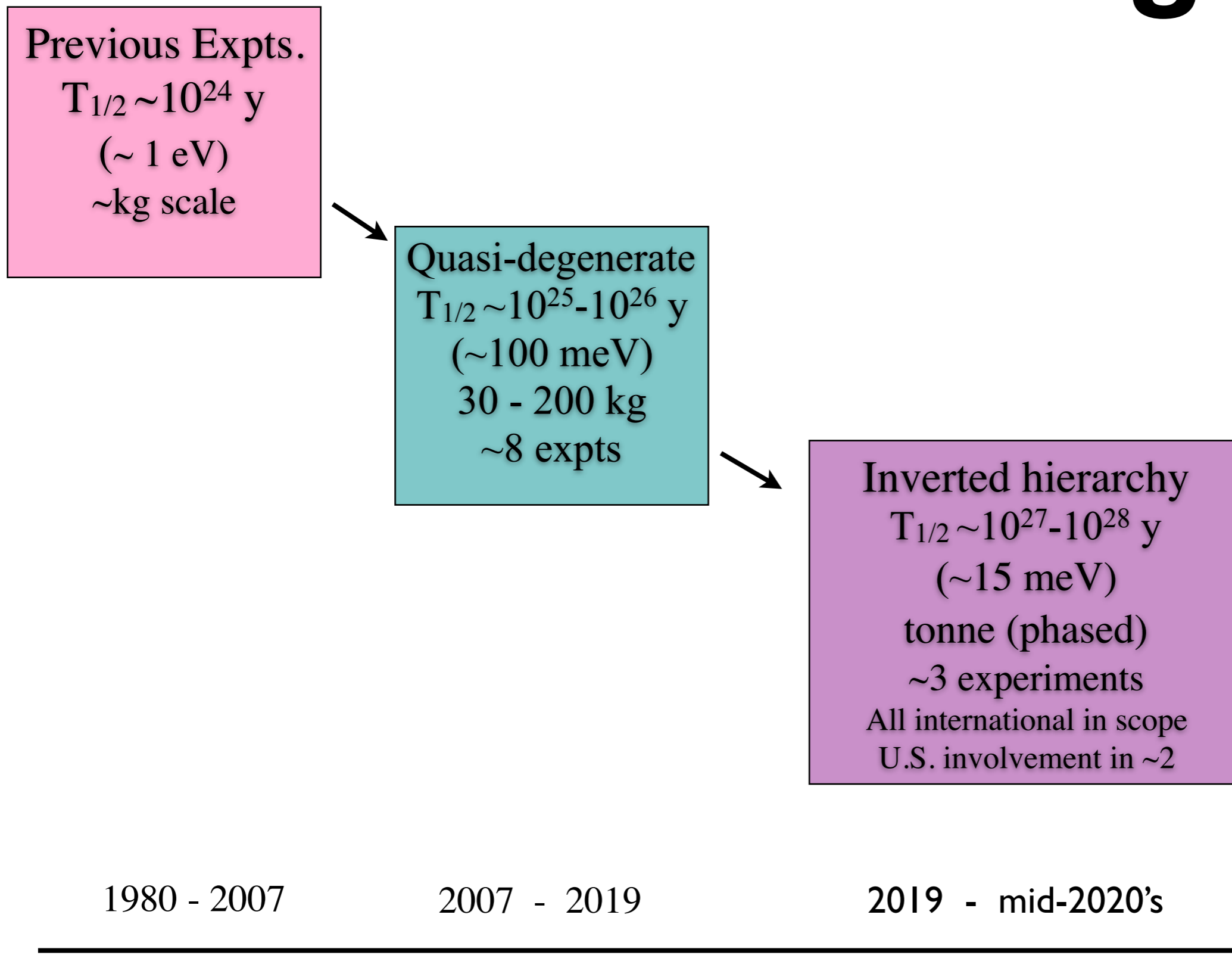
1980 - 2007

2007 - 2019

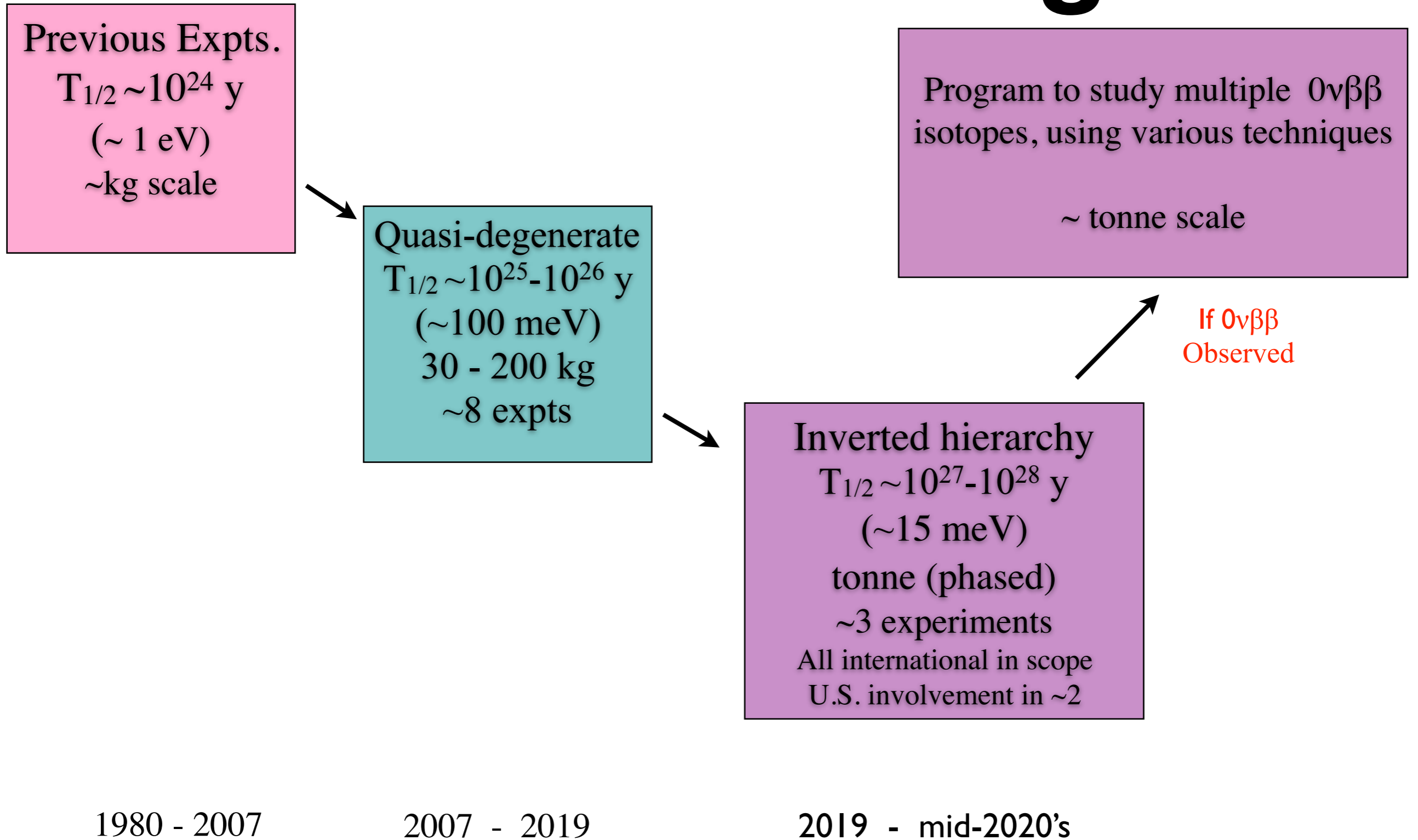
2019 - mid-2020's



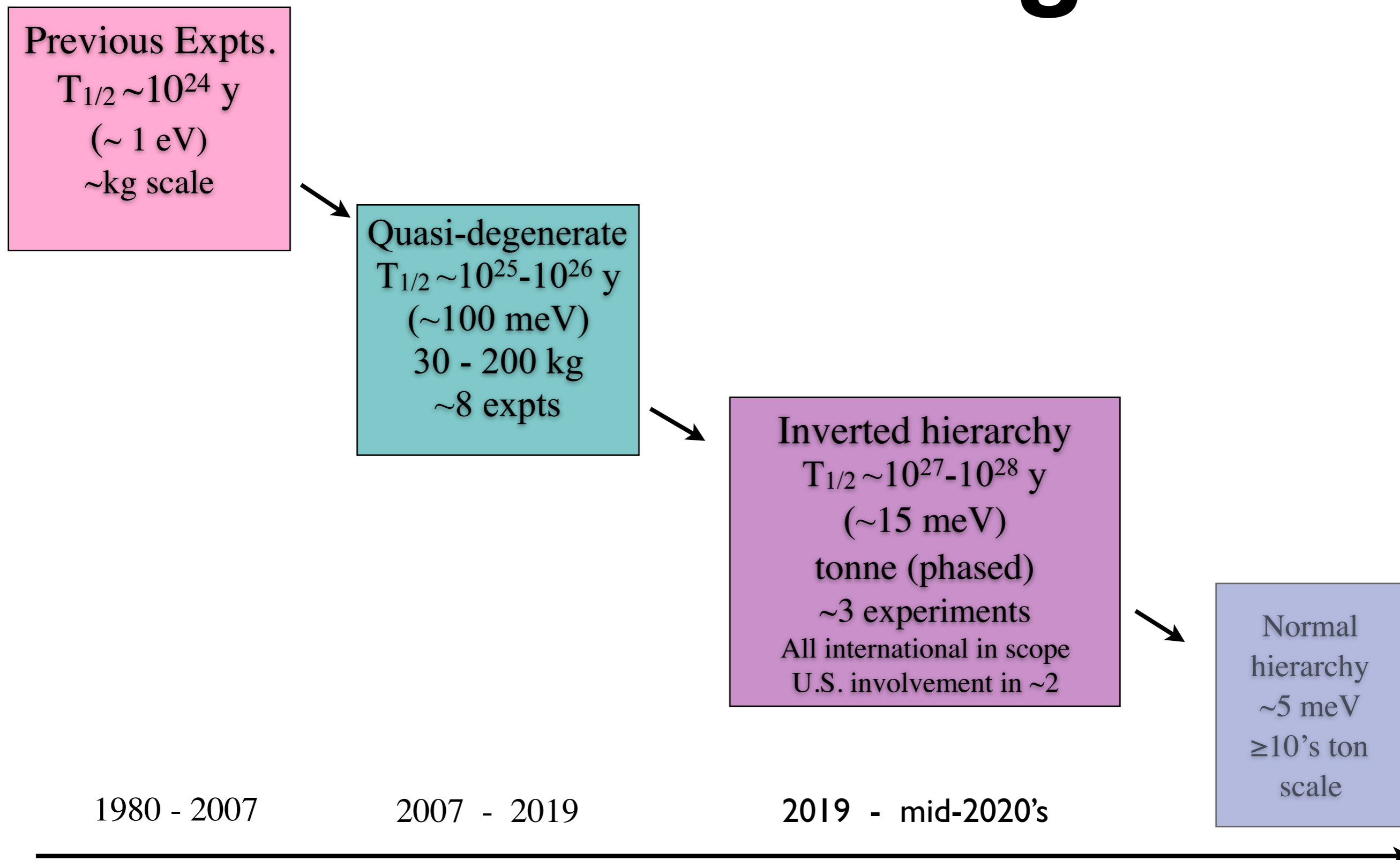
International Program



International Program

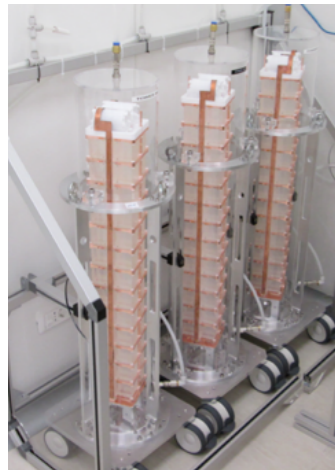


International Program

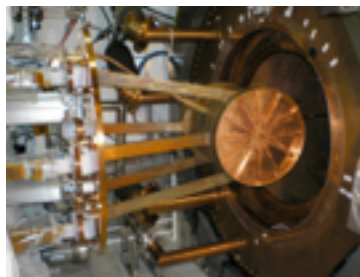


World Program

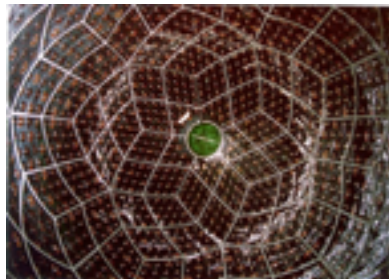
CUORE



EXO200



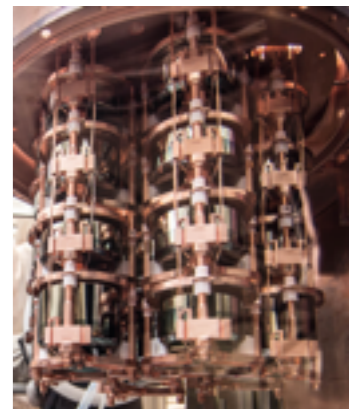
KamLAND Zen



GERDA



MAJORANA



SNO+



Collaboration	Isotope	Technique	mass ($0\nu\beta\beta$ isotope)	Status
CANDLES	Ca-48	305 kg CaF ₂ crystals - liq. scint	0.3 kg	Construction
CARVEL	Ca-48	⁴⁸ CaWO ₄ crystal scint.	~ ton	R&D
GERDA I	Ge-76	Ge diodes in LAr	15 kg	Complete
GERDA II	Ge-76	Point contact Ge in LAr	31	Operating
MAJORANA DEMONSTRATOR	Ge-76	Point contact Ge	25 kg	Operating
LEGEND	Ge-76	Point contact	~ ton	R&D
NEMO3	Mo-100 Se-82	Foils with tracking	6.9 kg 0.9 kg	Complete
SuperNEMO Demonstrator	Se-82	Foils with tracking	7 kg	Construction
SuperNEMO	Se-82	Foils with tracking	100 kg	R&D
LUCIFER (CUPID)	Se-82	ZnSe scint. bolometer	18 kg	R&D
AMoRE	Mo-100	CaMoO ₄ scint. bolometer	1.5 - 200 kg	R&D
LUMINEU (CUPID)	Mo-100	ZnMoO ₄ / Li ₂ MoO ₄ scint. bolometer	1.5 - 5 kg	R&D
COBRA	Cd-114,116	CdZnTe detectors	10 kg	R&D
CUORICINO, CUORE-0	Te-130	TeO ₂ Bolometer	10 kg, 11 kg	Complete
CUORE	Te-130	TeO ₂ Bolometer	206 kg	Operating
CUPID	Te-130	TeO ₂ Bolometer & scint.	~ ton	R&D
SNO+	Te-130	0.3% ^{nat} Te suspended in Scint	160 kg	Construction
EXO200	Xe-136	Xe liquid TPC	79 kg	Operating
nEXO	Xe-136	Xe liquid TPC	~ ton	R&D
KamLAND-Zen (I, II)	Xe-136	2.7% in liquid scint.	380 kg	Complete
KamLAND2-Zen	Xe-136	2.7% in liquid scint.	750 kg	Upgrade
NEXT-NEW	Xe-136	High pressure Xe TPC	5 kg	Operating
NEXT	Xe-136	High pressure Xe TPC	100 kg - ton	R&D
PandaX - 1k	Xe-136	High pressure Xe TPC	~ ton	R&D
DCBA	Nd-150	Nd foils & tracking chambers	20 kg	R&D

Ton Scale Experiments

- Active international collaborations building on current efforts.
 - ^{76}Ge : **LEGEND**, HPGE crystals, ~ton (builds on GERDA & MAJORANA)
 - ^{82}Se : SuperNEMO : Se foils, tracking and calorimeter, 100 kg scale
 - ^{100}Mo : AMoRE : CaMoO_4 scint. bolometer, 200 kg scale
 - ^{136}Xe : **nEXO** — Liquid TPC, 5 tons
 - NEXT — High pressure gas TPC, ton scale
 - PandaX - III — High pressure gas TPC, ton scale
 - KamLAND-Zen — ^{136}Xe in scintillator, 800 kg scale
 - LZ — $^{\text{nat}}\text{Xe}$ liquid TPC, 7 tons, operating 2019
 - ^{130}Te : **CUPID (CUORE with Particle ID)** — Bolometer - Scintillation
 - SNO+ Phase I & II — ^{130}Te in scintillator
- Experiments can be done in a staged (phased) approach. Most are considering stepwise increments.
- Isotope enrichment (^{76}Ge , ^{82}Se , ^{136}Xe) requires time and \$s.
- Potential underground lab sites
 - SNOLAB, JingPing, Gran Sasso, SURF, CanFranc, Frejus, Kamioka, ANDES, Y2L

**Specific Examples:
EXO-200 and future nEXO,
LEGEND and CUPID**

Advantages of ^{136}Xe

Isotopic enrichment easier & known: *Xe is a gas and ^{136}Xe is the heaviest isotope.*

Xenon is “reusable”: *can be re-purified (noble gas: relatively easy) during measurement and easily recycled into a different detector (no crystal growth)*

... replace ^{136}Xe with $^{\text{nat}}\text{Xe}$ if signal observed

Monolithic detector: *LXe is self shielding, surface contamination minimized.*

Minimal cosmogenic activation: *no long lived radioactive isotopes of Xe.*

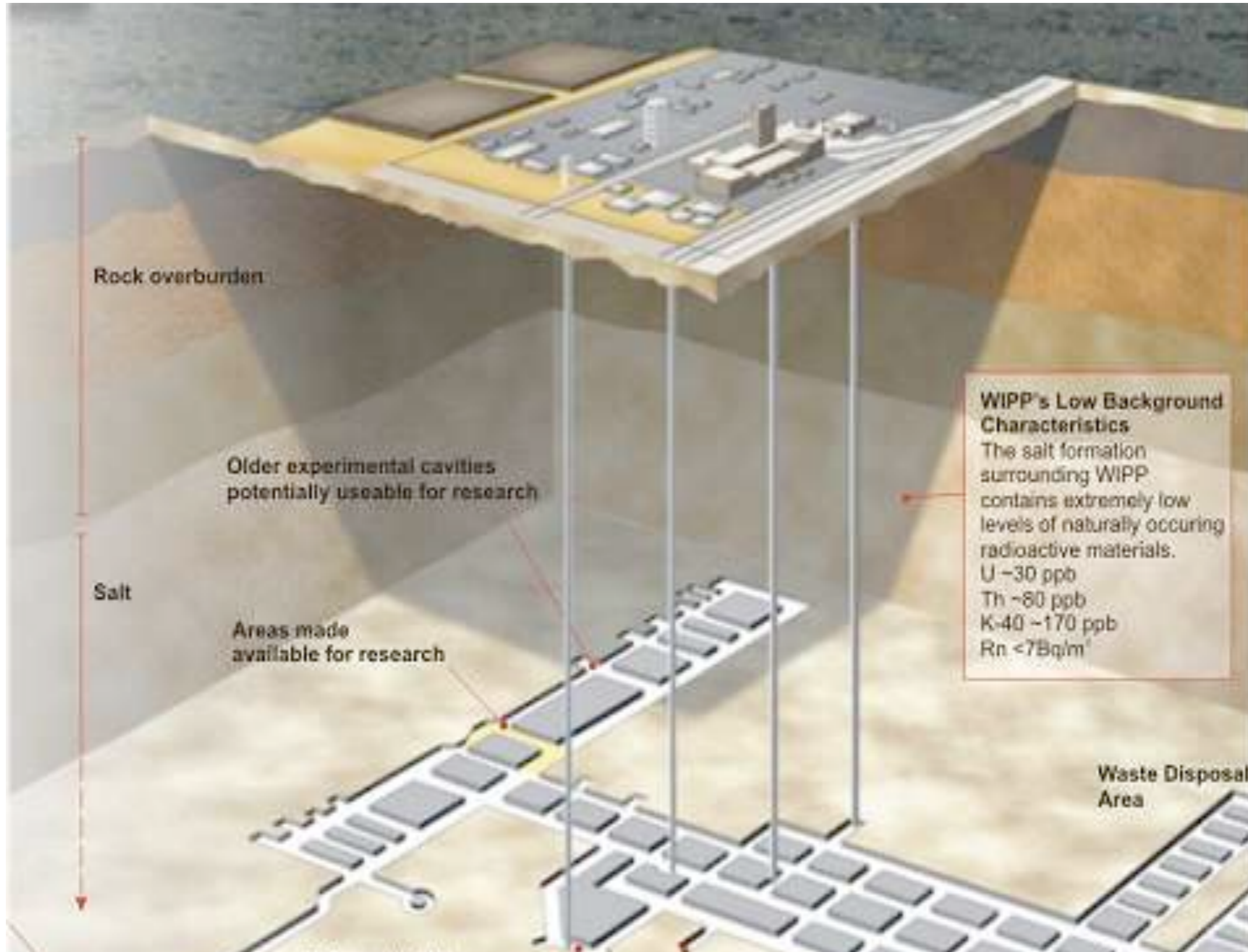
Energy resolution in LXe improved: *scintillation light + ionization anti-correlation.*

Standard $2\nu\beta\beta$ is slow! (see later): *get away with modest energy resolution*

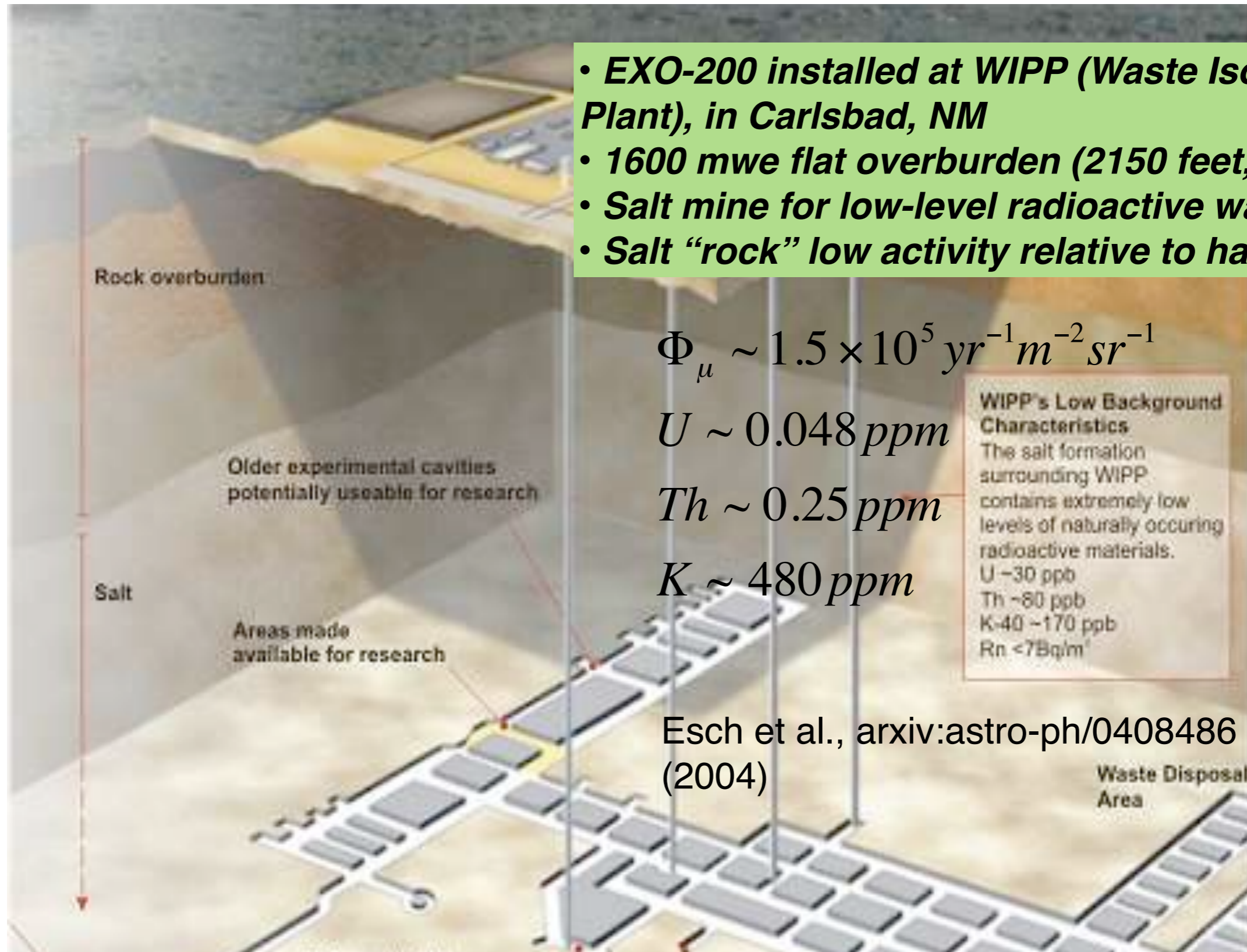
... admits a novel coincidence technique: *background reduction by Ba tagging*

... potentially access normal hierarchy

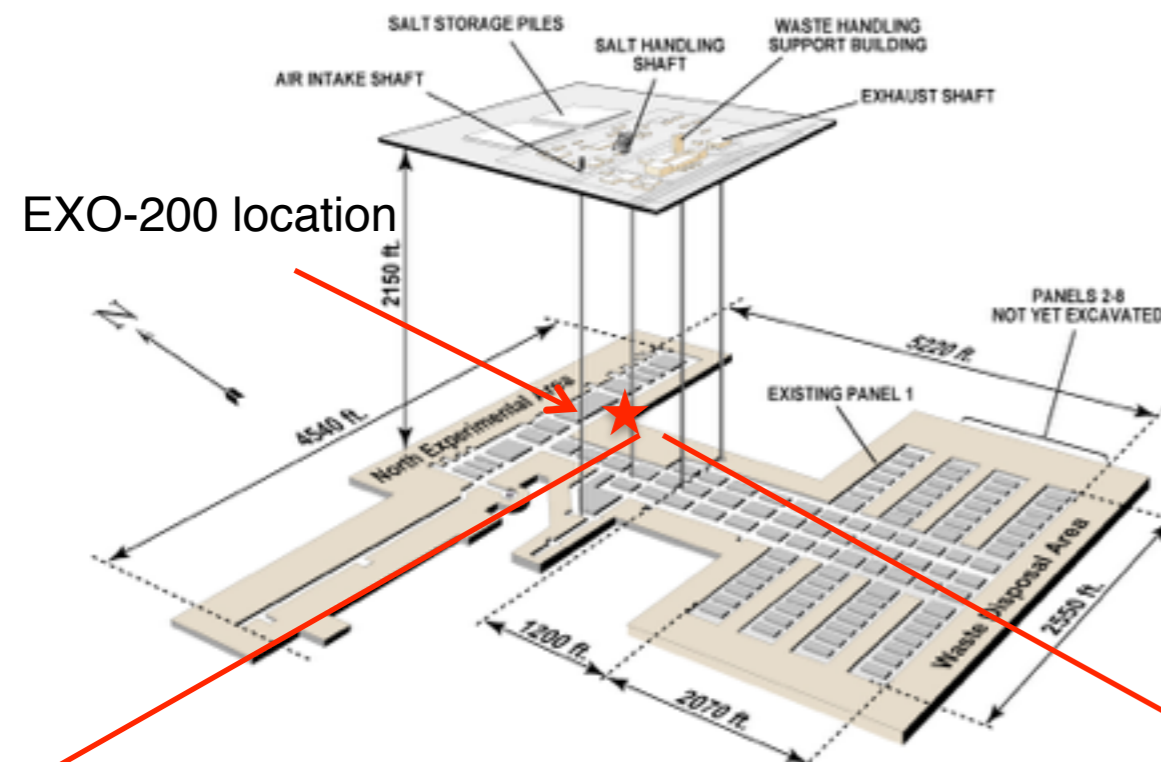
EXO-200 at WIPP



EXO-200 at WIPP



EXO-200 at WIPP



- *EXO-200 installed at WIPP (Waste Isolation Pilot Plant), in Carlsbad, NM*
- *1600 mwe flat overburden (2150 feet, 650 m)*
- *Salt mine for low-level radioactive waste storage*
- *Salt “rock” low activity relative to hard-rock mine*

$$\Phi_{\mu} \sim 1.5 \times 10^5 \text{ yr}^{-1} \text{ m}^{-2} \text{ sr}^{-1}$$

$$U \sim 0.048 \text{ ppm}$$

$$Th \sim 0.25 \text{ ppm}$$

$$K \sim 480 \text{ ppm}$$

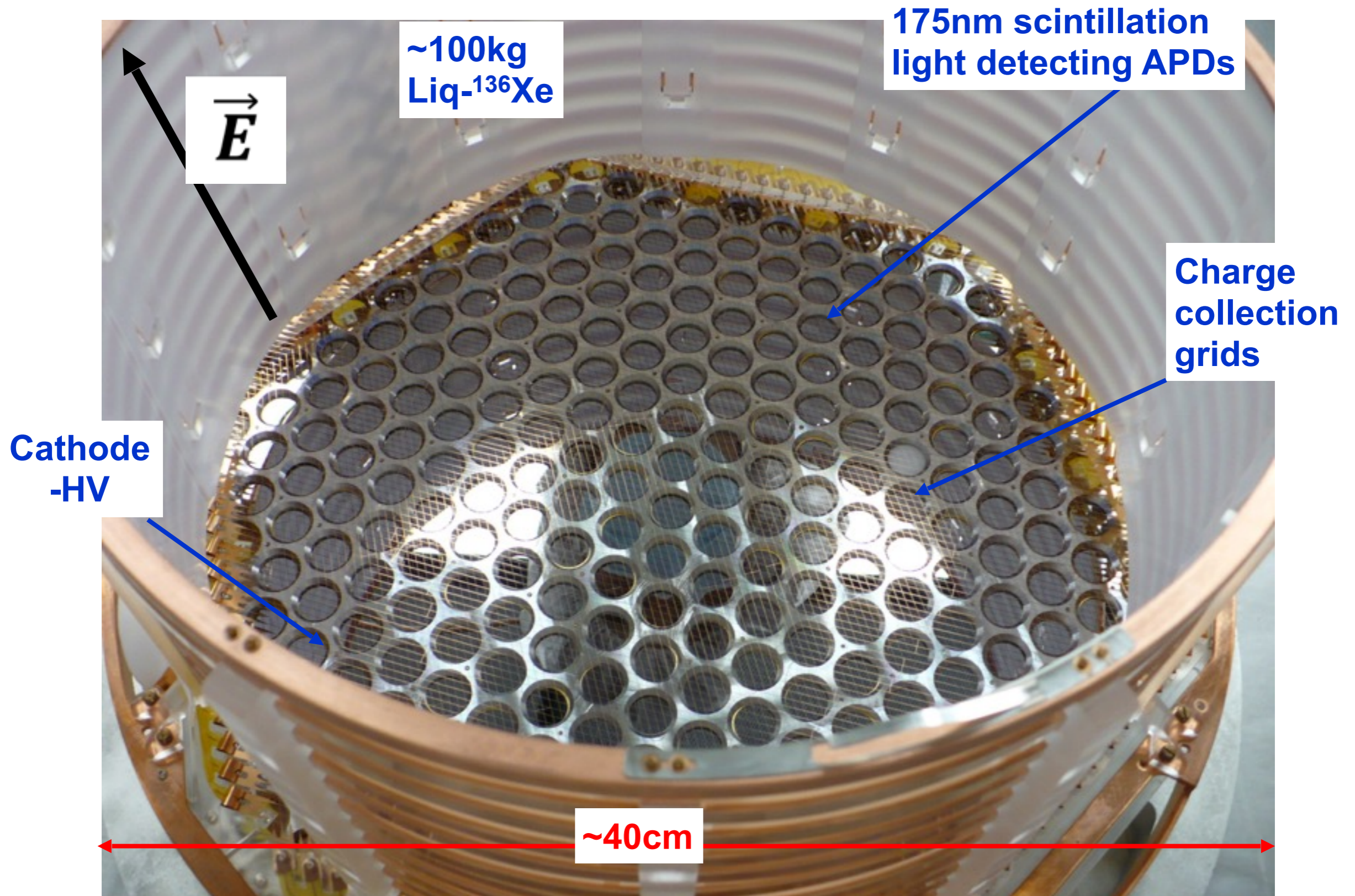
WIPP's Low Background Characteristics
 The salt formation surrounding WIPP contains extremely low levels of naturally occurring radioactive materials.
 U ~30 ppb
 Th ~80 ppb
 K-40 ~170 ppb
 Rn <7Bq/m³

Esch et al., arxiv:astro-ph/0408486 (2004)

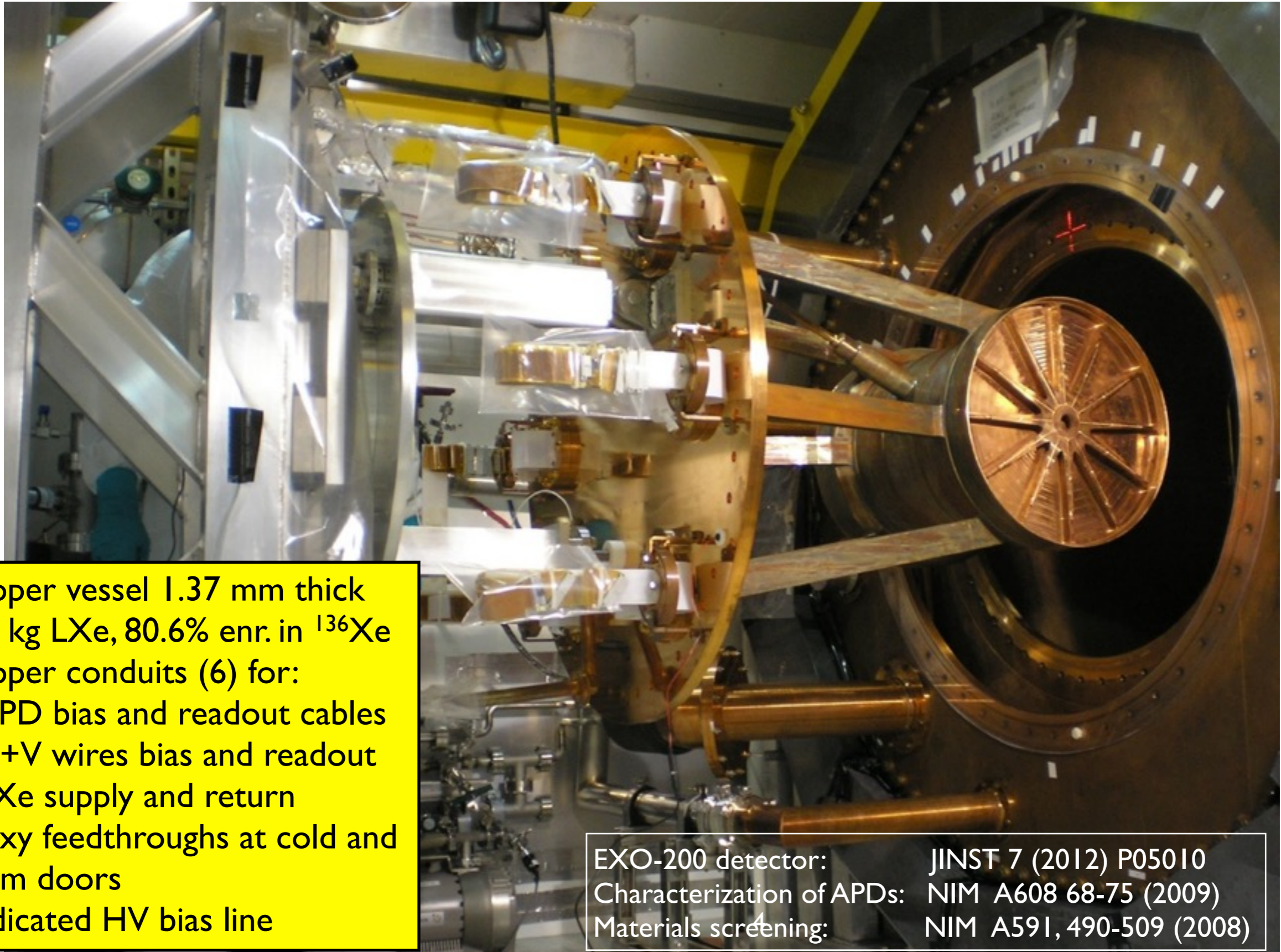
Waste Disposal Area



The EXO-200 TPC



TPC Entering the Cryostat



Copper vessel 1.37 mm thick
175 kg LXe, 80.6% enr. in ^{136}Xe
Copper conduits (6) for:

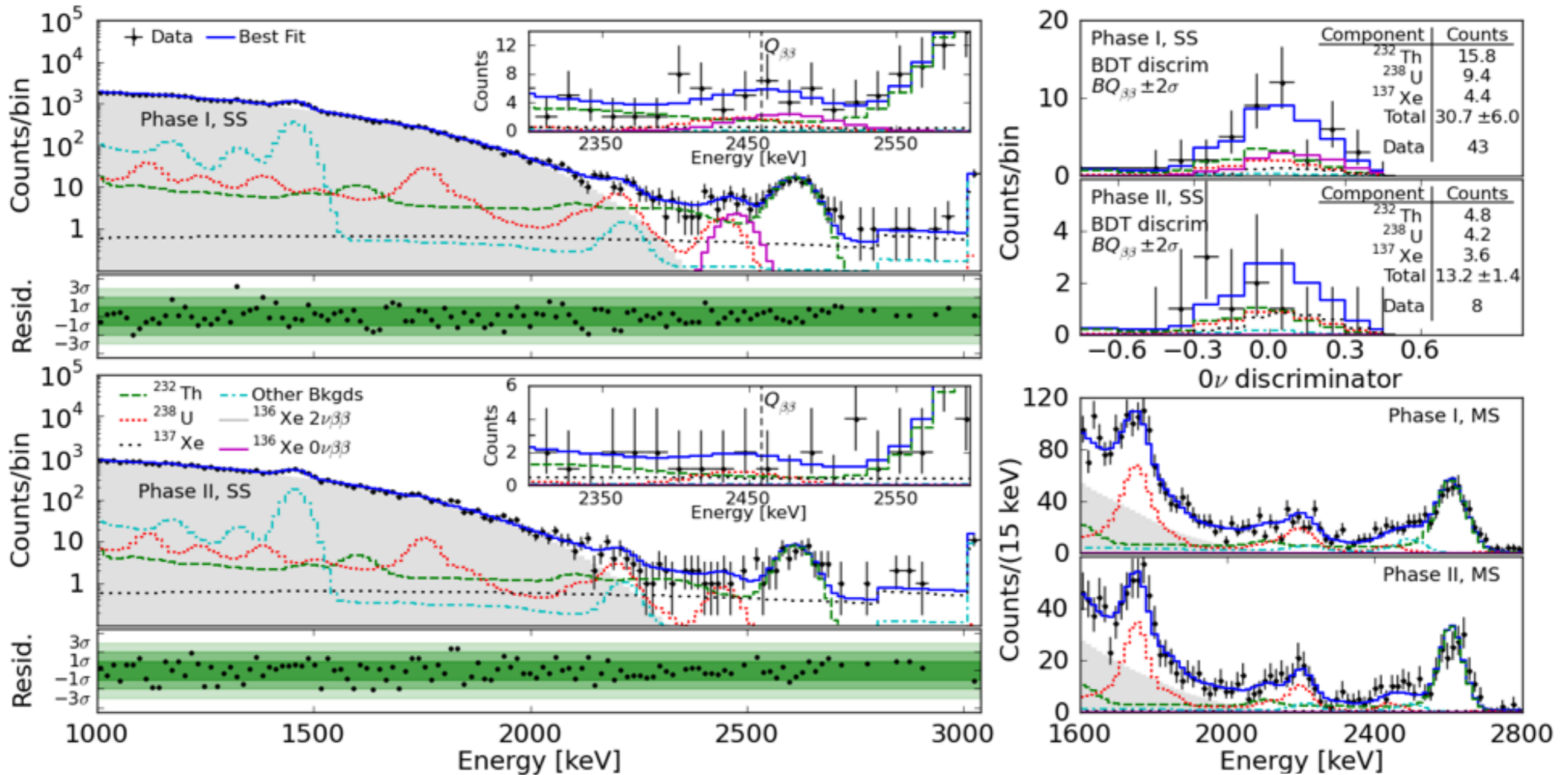
- APD bias and readout cables
- U+V wires bias and readout
- LXe supply and return

Epoxy feedthroughs at cold and warm doors
Dedicated HV bias line

EXO-200 detector: JINST 7 (2012) P05010
Characterization of APDs: NIM A608 68-75 (2009)
Materials screening: NIM A591, 490-509 (2008)

EXO-200: Recent Results

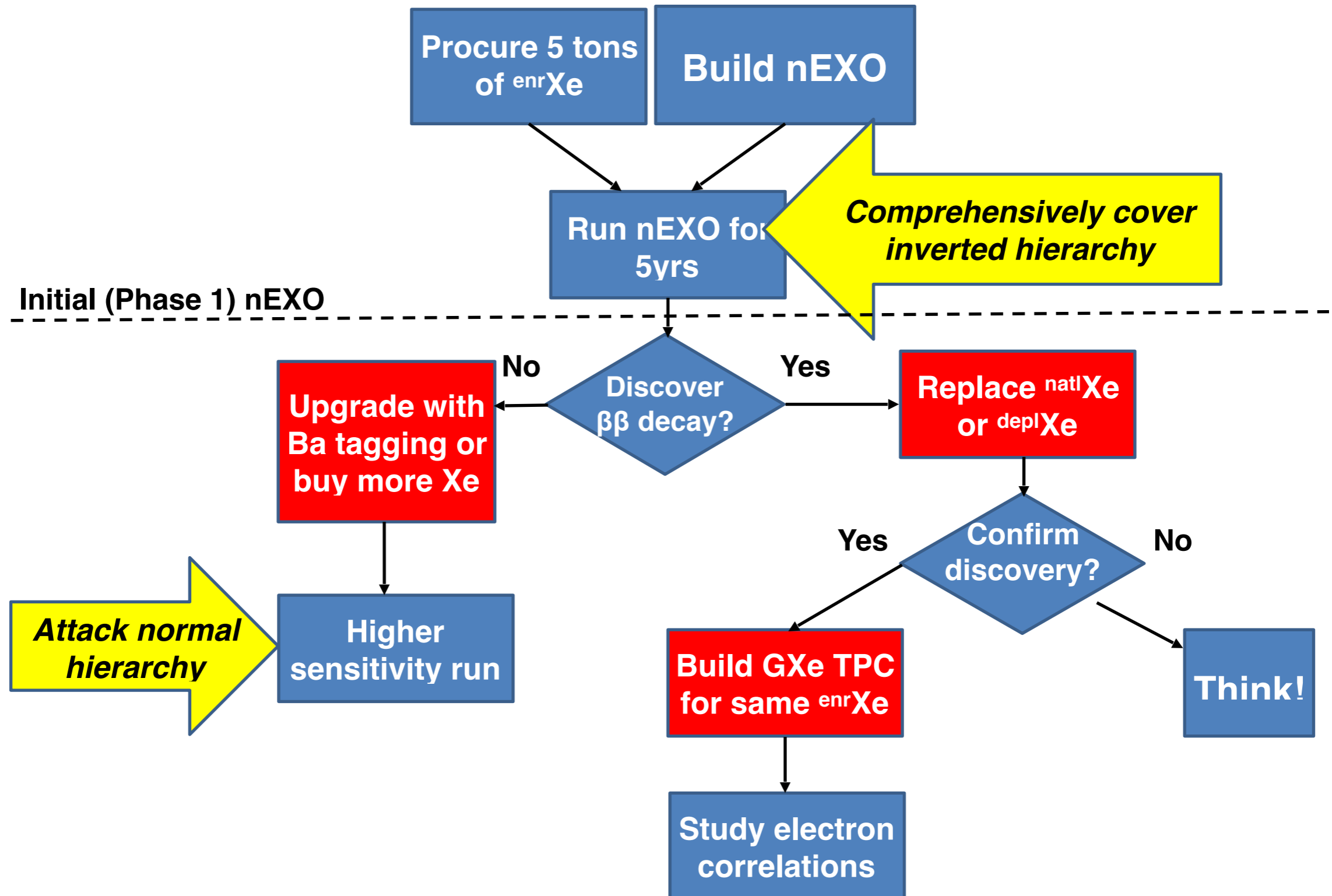
- Background model + data \rightarrow maximum likelihood fit
- Combine Phase I + Phase II profiles



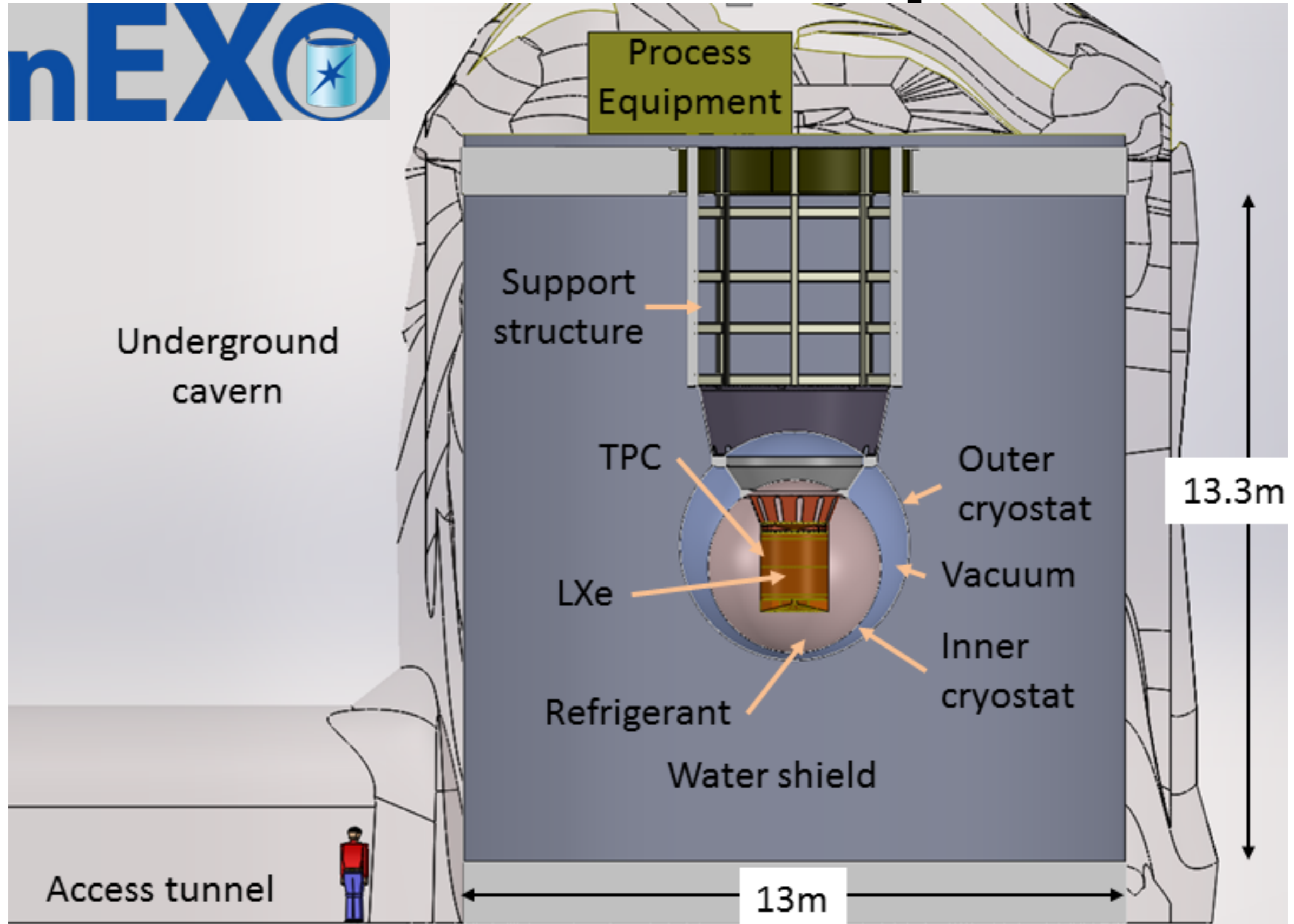
	Sensitivity (yr)	90% CL Limit (yr)	$\langle m_{\beta\beta} \rangle$ (meV)
PRL 120 072701 (2018)	3.8×10^{25}	1.8×10^{25}	147-398

nEXO Strategy

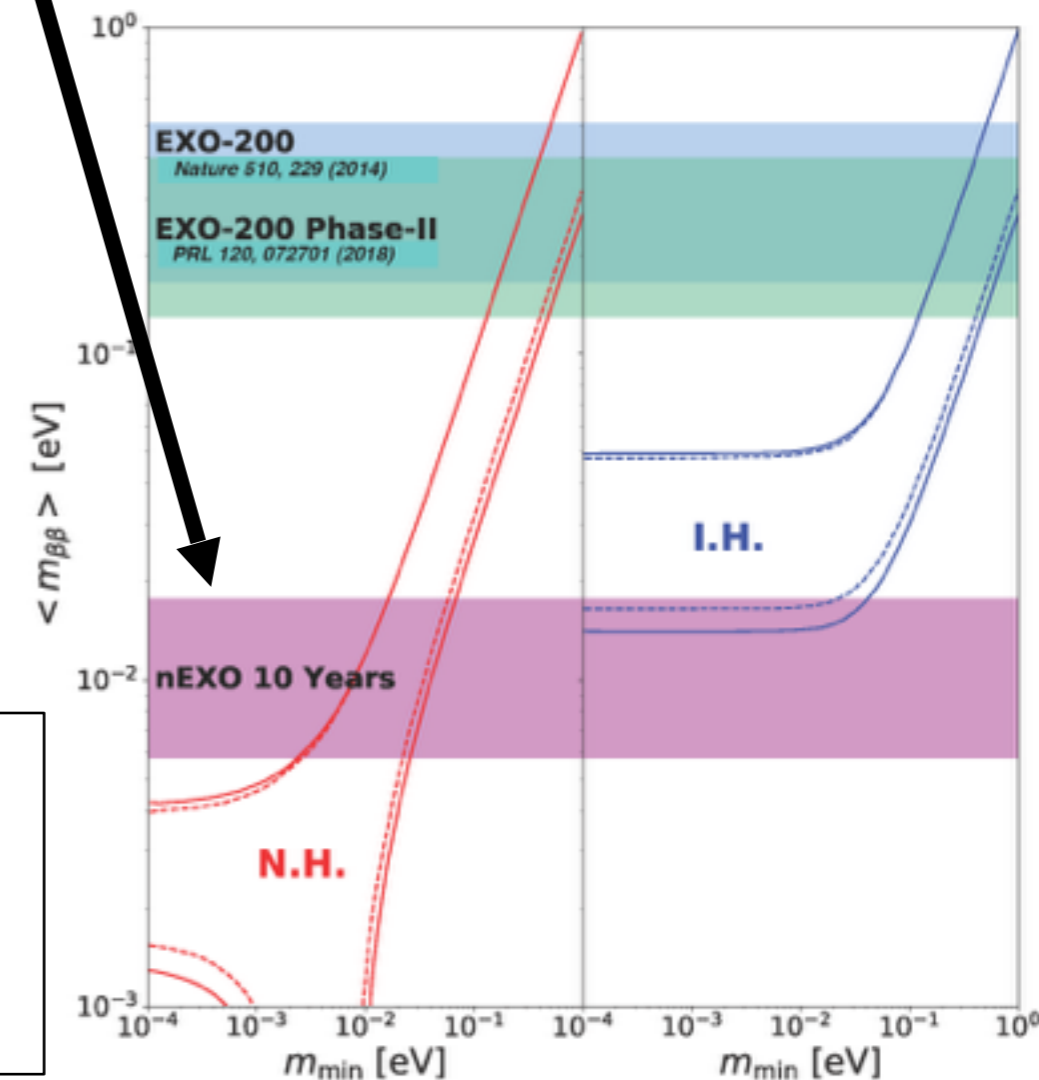
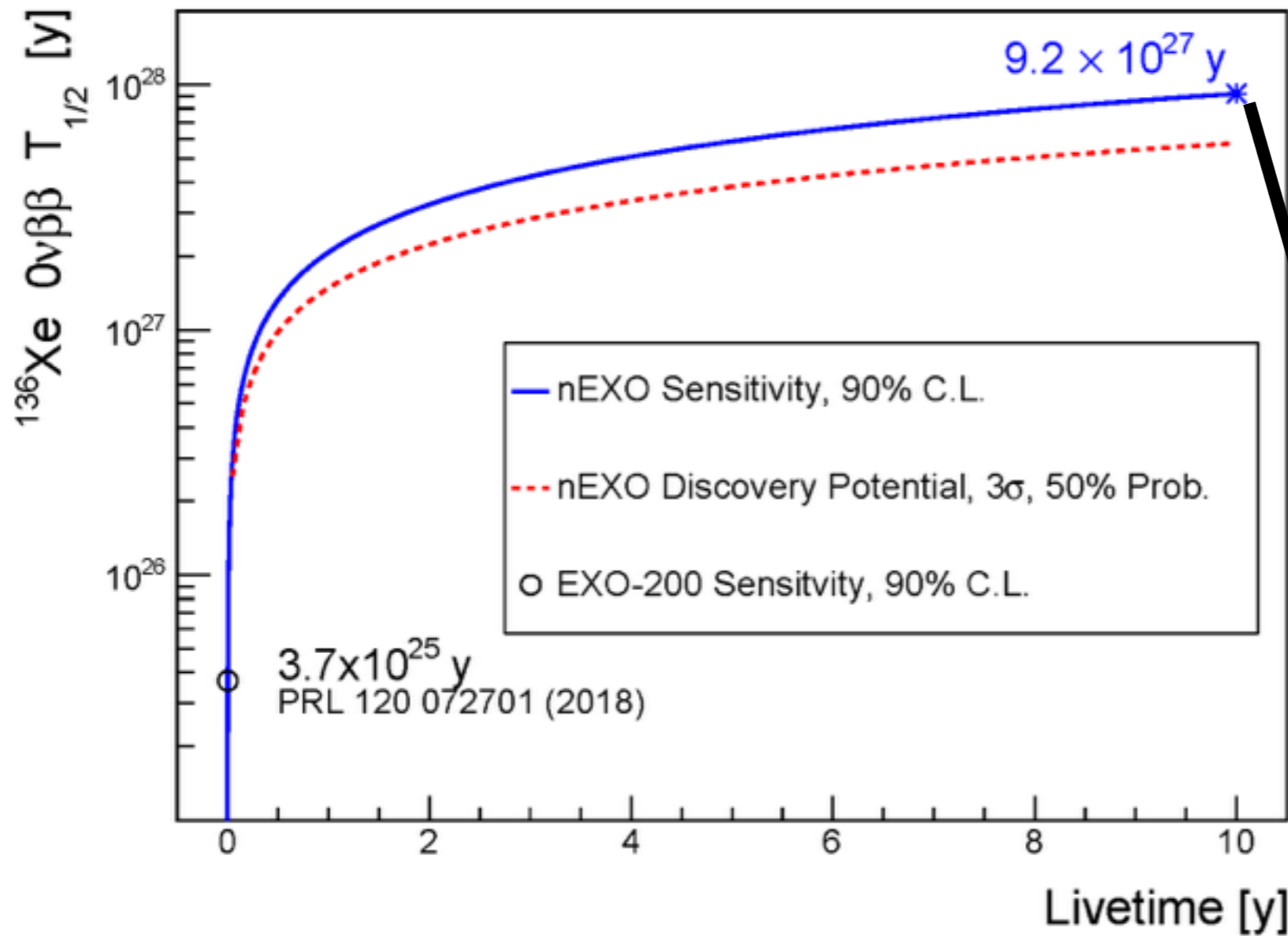
Flexible program based on the initial nEXO investment



nEXO Concept



nEXO: Discovery Reach



- $g_A = g_A^{\text{free}} = -1.2723$

- Band is the envelope of NME:

EDF: T.R. Rodríguez and G. Martínez-Pinedo, PRL 105, 252503 (2010)

ISM: J. Menendez et al., Nucl Phys A 818, 139 (2009)

IBM-2: J. Barea, J. Kotila, and F. Iachello, PRC 91, 034304 (2015)

QRPA: F. Šimkovic et al., PRC 87 045501 (2013)

SkyrmeQRPA: M.T. Mustonen and J. Engel PRC 87 064302 (2013)

^{76}Ge : LEGEND

Mission: The collaboration aims to develop a phased, ^{76}Ge -based double-beta decay experimental program with discovery potential at a half-life significantly longer than 10^{27} years, using existing resources as appropriate to expedite physics results.

Select best technologies, based on what has been learned from GERDA and the MAJORANA DEMONSTRATOR, as well as contributions from other groups and experiments.

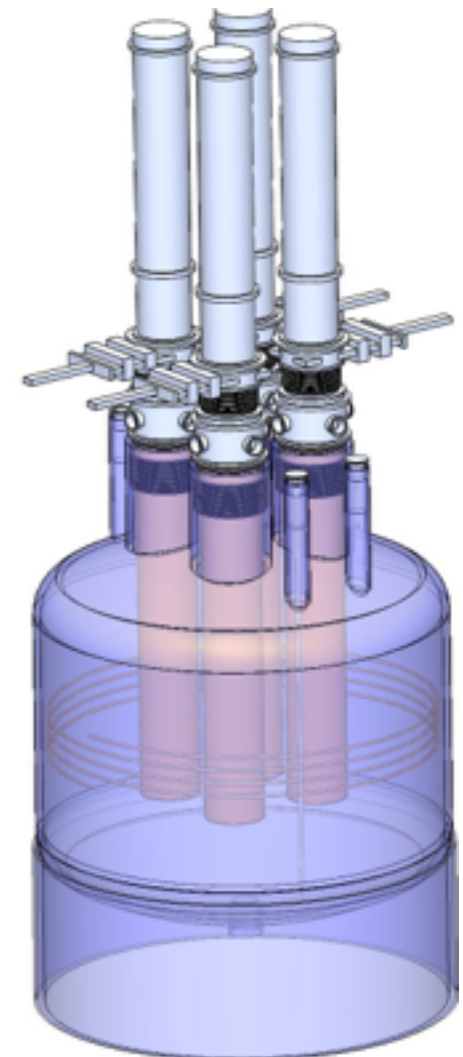
First Phase:

- (up to) 200 kg
- modification of existing GERDA infrastructure at LNGS
- BG goal (x5 lower) 0.6 c / (FWHM t y)
- start by 2021



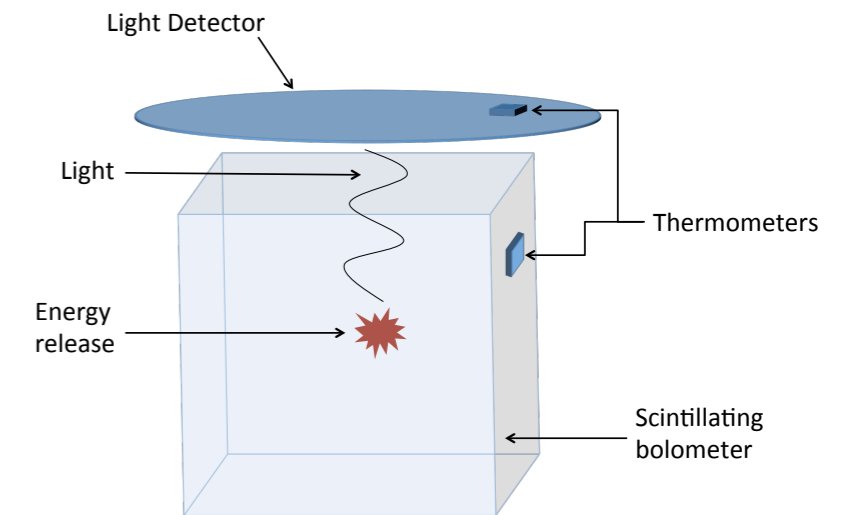
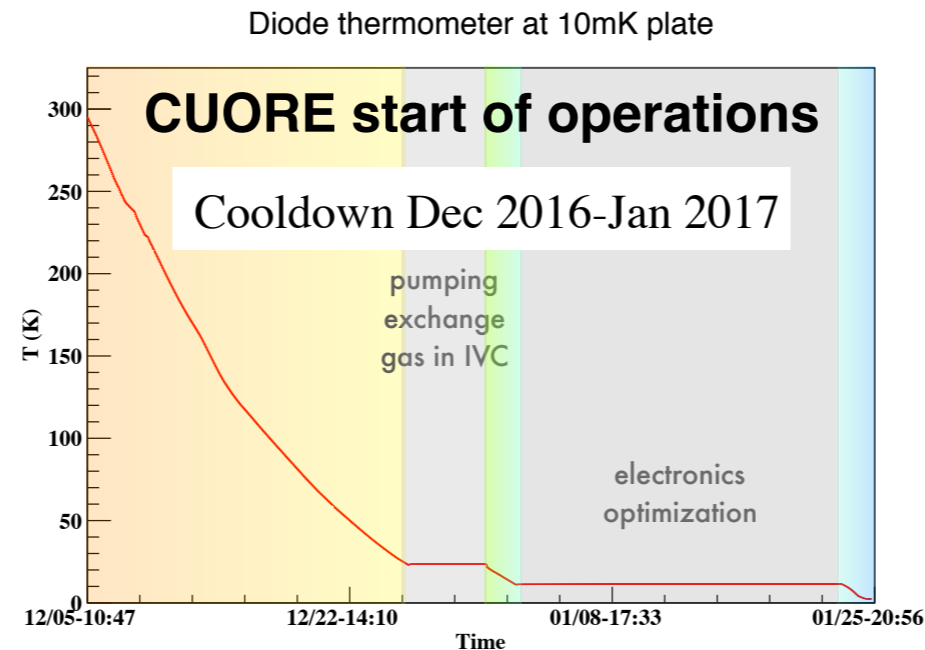
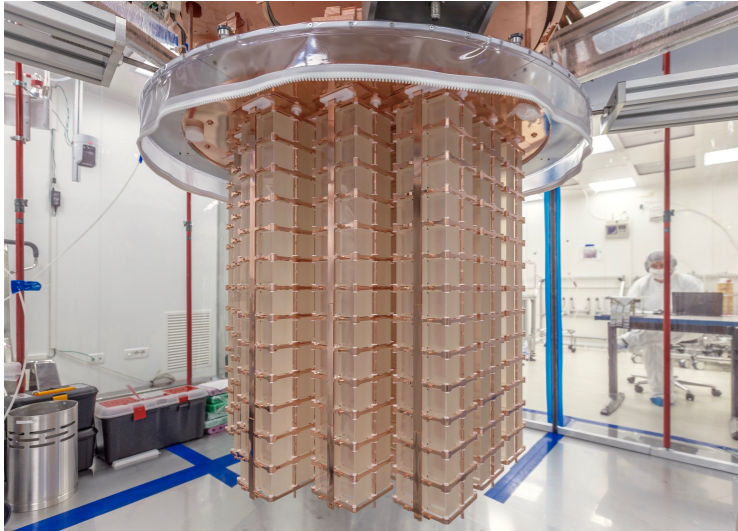
Subsequent Stages:

- 1000 kg (staged)
- timeline connected to U.S. DOE down selection process
- BG: goal (x30 lower) 0.1 c / (FWHM t y)
- Location: TBD
- Required depth ($^{77\text{m}}\text{Ge}$) under investigation



^{130}Te : CUORE/CUPID

CUORE detectors installed



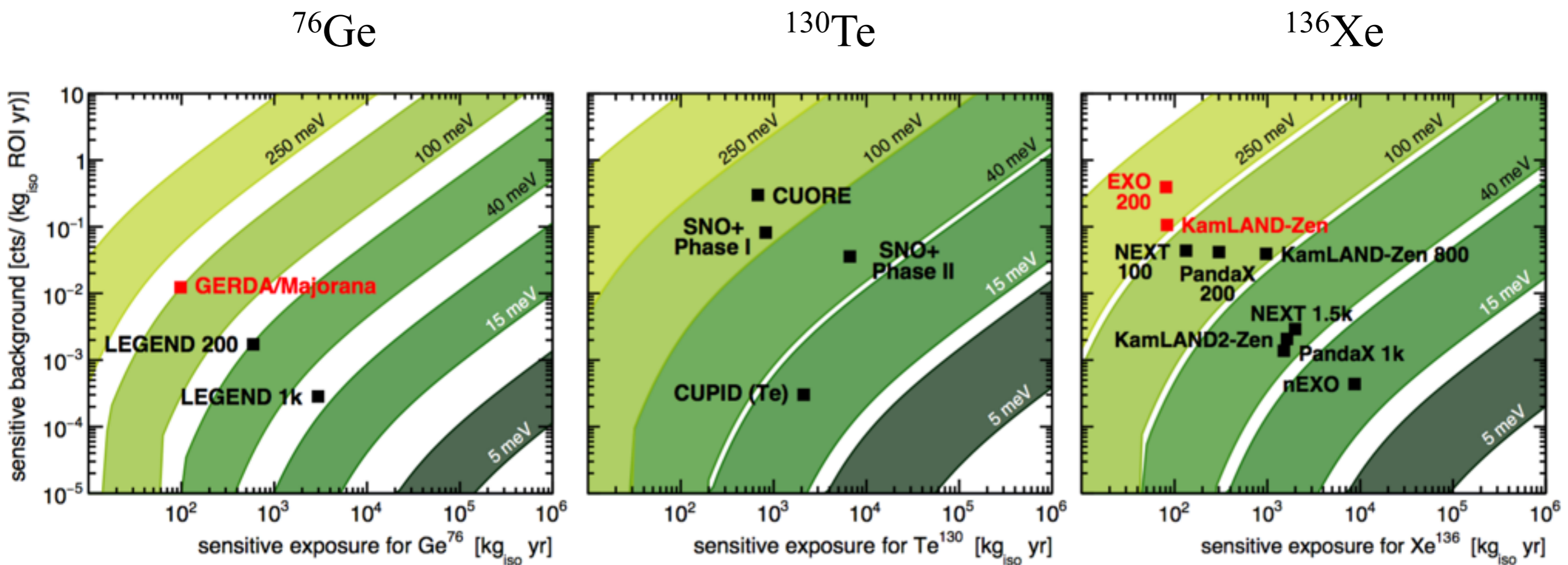
Next-generation bolometric tonne-scale experiment based on the CUORE design, proven CUORE cryogenics

- CUORE Milestones:
 - Tower installation: Jul-Aug 2016
 - Cryostat closeout: Nov 2016
 - Cooldown: Dec-Jan 2016
 - Commissioning and initial performance optimization: Jan-May 2017
 - First science run: May 2017
- Cryostat performs very well: base $T < 7$ mK
- $>95\%$ of detectors operational
- First data reported in Summer 2017

- Intense CUPID R&D effort in the next 2-3 years
 - ☞ US focus: $^{130}\text{TeO}_2$ enrichment and purification, high-resolution sensors for Cherenkov light
 - ☞ Complementary European efforts
 - ☞ Background goal is 0.1 cts/ROI-t-yr; achieve sensitivity to the full Inverted Hierarchy
 - ☞ Other important R&D: detailed background analysis, cosmogenic backgrounds @ LNGS — to be addressed before downselect
 - ☞ Worldwide efforts: 8 countries, 32 institutions
 - ☞ Data from CUORE and pilot detectors will drive technology and isotope choice

Discovery Sensitivity Comparison

Discovery probability of next-generation neutrinoless double-beta decay experiments
 Matteo Agostini, Giovanni Benato, and Jason Detwiler arXiv:1705.02996v3



Red : Achieved Backgrounds; Black : Projected Backgrounds

Width of bands based on range of NME values

Summary and Outlook

- ◆ **Low Energy Weak Neutral Current Interactions**
 - ★ Central to our understanding of the Standard Model
 - ★ Remains relevant for BSM searches, especially flavor-diagonal
- ◆ **Parity-Violating Electron Scattering in the next decade**
 - ★ Technical progress has enabled unprecedented precision
 - ★ Flagship experiments at electron accelerators
 - ★ Fundamental Nuclear/Nucleon as well as EW/BSM physics
- ◆ **Atomic Parity Violation and Neutrino Scattering**
 - ★ The low Q region might become important in the future!
 - ★ Neutrino scattering should be investigated in any case...
- ◆ **Neutrinoless Double Beta Decay Searches**
 - ★ The Majorana nature of neutrinos: pressing BSM question
 - ★ Discovery of lepton number violation would have wide implications
 - ★ The next generation experiments will attack the inverted hierarchy
 - ★ R&D to go further to the normal hierarchy has already begun