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Complexity in QFT and Tensor Networks

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Talk based on work:

Complexity Functionals and Complexity Growth Limits in Continuous MERA Circuits. [arXiv:1803.02356]

JMV in collaboration with A. del Campo, UMass, Boston.





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Holographic Complexity. A Tale of two Dualities.

Complexity of quantum states: useful to investigate the duality between entanglement and spacetime geometries in quantum gravity.

Mainly addressed in the context of the AdS/CFT duality.

"Holographic complexity" as the CFT quantity that encodes the continuous evolution of the interior of the black hole [Susskind 16] .

Two proposals for the quantum complexity of boundary states:

Complexity=Volume (CV) conjecture. [Susskind 16] Complexity= Action (CA) conjecture. [Brown et al 16]

Holographic Complexity. A Tale of two Dualities.



The eternal black hole: dual to an entangled state $|\Psi\rangle$ of two CFTs that live on the boundaries.

CV: C_{Ψ} = Volume of the maximal spatial slice.

CA: C_{Ψ} = Action \mathcal{A} of the Wheeler-DeWitt patch. (roughly corresponds to the black hole interior)

CA and Complexification Rates.

The action of the black hole interior ${\cal A}$ increases at a rate such that

$$\frac{d\mathcal{C}_{\Psi}}{dt} = \frac{2M}{\pi\hbar}$$

i.e, saturates a Margolus-Levitin like bound [Brown et al 16].

The ML bound is a universal bound on the speed of evolution of any quantum system (Quantum Speed Limit, QSL).

$$rac{d}{dt} \langle \Psi | \Psi(t)
angle \leq rac{2E_{\Psi}}{\pi\hbar}$$

In Computer Science, ML is a theoretical upper limit on the number of operations that can be performed on unit time.

Black Holes Produce Complexity Fastest.

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CA and Quantum Circuits.



Logic gates (blue) in a quantum circuit (red).

Black hole interiors = quantum circuits that produce $|\Psi\rangle$ at the fastest rate allowed by QM

 ${\cal C}_{|\Psi\rangle}\sim \#\,gates$ required to produce $|\Psi\rangle$ from a simple reference state $|\Omega\rangle.$

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Complexity in QFT: An interesting research program.

• Complexity measures the cost required to prepare $|\Psi\rangle$ from a specific reference state $|\Omega\rangle$ by applying an optimal unitarity U

 $|\Psi
angle pprox U \left|\Omega
ight
angle$

2 U from iterating generators K taken from some elementary set

$$U(u) = \mathcal{P} \exp\left(\int_{u_l}^u K(u') \, du'
ight)$$

 $|\Psi\rangle \approx U(u_F) |\Omega\rangle$

A regularization procedure to deal with ultraviolet divergences.
A measure of complexity.

Entanglement Renormalization. MERA.

[Vidal 2007]



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Entanglement Renormalization. MERA



MERA circuits and Holography.

Proposal [Swingle, 2009] connecting the ideas of:

- AdS/CFT [Maldacena, 1998]
- Holographic Entanglement Entropy [Ryu-Takayanagi, 2006]
- Quantum Renormalization Group [Vidal, 2007]



Continuous MERA circuits





$$|\Psi^{\Lambda}
angle \equiv \mathcal{P}\exp\left(-i\int_{u_{IR}}^{0}d\hat{u}\,\widetilde{K}(\hat{u})
ight)|\Omega
angle$$

where $\widetilde{K}(\hat{u})$ is the entangler operator.

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cMERA Free Boson Theory

For the free boson theory

$$S = \int dt \, dx \, \left((\partial_t \phi)^2 + (\partial_x \phi)^2 - m^2 \phi^2 \right)$$

The entangler operator is given by the Gaussian ansatz

$$\widetilde{K}(u) = \frac{1}{2i} \int dk \left(g(k, u) a_k^{\dagger} a_{-k}^{\dagger} - g^*(k, u) a_k a_{-k} \right)$$

with $g(k, u) = g(u) \Gamma(|k|e^{-u}/\Lambda)$ where $\Gamma(|k|e^{-u}/\Lambda)$ limits the momentum integral to $|k| \leq \Lambda e^u$

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Coherent State Formulation

The state at scale u can be written as a SU(1, 1) gaussian squeezed state

$$|\Psi(u)\rangle = N \exp\left(-\frac{1}{2}\int dk \left[\Phi(k,u) a_{k}^{\dagger}a_{-k}^{\dagger} - \Phi^{*}(k,u) a_{k}a_{-k}\right]\right)|\Omega\rangle$$

with

$$\Phi(k, u) = \int_{u_{IR}}^{u} du' g_k(k, u')$$

• Scale dependent representations of the state are obtained by adding left-right moving modes with $|k| \leq \Lambda e^u$ to $|\Omega\rangle$

cMERA Entanglement.



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cMERA Variational Optimization

The variational parameter $\Phi(k, u)$ of cMERA for the free boson

$$\Phi(k, u) = \left[-\frac{1}{4} \log \frac{k^2 + m^2}{\Lambda^2 + m^2} \right]_{k = \Lambda e^u} = -\frac{1}{4} \log \frac{\Lambda^2 e^{2u} + m^2}{\Lambda^2 + m^2}$$

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cMERA Coherent State Path Integral

$$\mathcal{Z}_{\textit{cMERA}} = \langle \Psi^{\Lambda} | \mathcal{P} \, \exp \left(-i \int_{u_{\textit{IR}}}^{0} ilde{K}(\hat{u}) \, d\hat{u}
ight) | \Omega
angle$$

as a SU(1,1) coherent state path integral.

$$\mathcal{Z}_{cMERA} = \int \mathcal{D}\Phi \mathcal{D}\Phi^* \exp\left(i\mathcal{A}\left[\Phi, \Phi^*\right]\right)$$
$$\mathcal{A}\left[\Phi, \Phi^*\right] = -2\operatorname{Vol}\int_{|k| \le \Lambda} dk \int_{u_{IR}}^0 du \,\Phi^* \partial \Phi$$
$$= -2\operatorname{Vol}\int_{|k| \le \Lambda e^u} dk \int_{u_{IR}}^0 du \,g(u) \,\Phi(k, u)$$

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Our Proposal.

$$\mathcal{C}_{\mathcal{A}}\left[|\Psi^{\Lambda}
angle$$
, $|\Omega
angle
ight]\equiv\mathcal{A}\left[\Phi,\Phi^{*}
ight]_{ ext{on-shell}}$

where *on-shell* indicates that the action is evaluated with the parameter $\Phi(k, u)$ obtained from the cMERA variational procedure.

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Results.

For the massless case

$$\mathcal{C}_{A} = \frac{\text{Vol} \cdot \Lambda}{2}$$

we obtain divergencies \approx CA and CV holographic complexity.

For the massive case

$$\mathcal{C}_{\mathcal{A}} pprox rac{ ext{Vol} \cdot \Lambda}{2} \left[1 - rac{m}{\Lambda} \left(1 - \log \; rac{m}{\Lambda}
ight)
ight]$$

UV divergences are naturally associated with existence of correlations or entanglement down to arbitrarily small length scales.

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Comparison with Circuit Length.

Measures the complexity of a cMERA circuit as a minimal length (in terms of the Fisher information metric) of a path running from $|\Omega\rangle$ to $|\Psi^{\Lambda}\rangle$ [Chapman 18].

$$\mathcal{C}_{\text{cMERA}} = \text{Vol} \, \int_{u_{\text{IR}}}^{0} \mathrm{d}u \, | \, g(u) \, | \, \int_{|k| \le \Lambda e^{u}} dk$$
$$= \text{Vol} \, \int_{|k| \le \Lambda} dk \, | \, \Phi(k, 0) \, |$$

For the massive case

$$\mathcal{C}_{\mathrm{cMERA}} pprox rac{\mathrm{Vol}\cdot\Lambda}{2} \left[1 - rac{m}{2\Lambda}\left(\pi - \arctan rac{m}{\Lambda}
ight)
ight]$$

Complexodynamics. Complexity Growth Limit.

Is \mathcal{C}_A constrained in any sense?

$$\frac{\Delta C_A}{\Delta u} = 2 E_{\Psi}$$

with

$$E_{\Psi} = rac{1}{\Delta u} \int_{u_l}^{u_F} du \left\langle \widetilde{K}(u) \right
angle$$

CGL Margolus-Levitin type bound

Equality as opposed to a lower bound reflects the optimality of \widetilde{K} for the generation of complexity along the cMERA circuit.

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Complexity from Liouville Field Action.

Similar results can be obtained by considering the functional

$$\mathcal{C}_{\mathcal{A}} = \mathcal{A}_{L}[\Phi] = \frac{1}{4} \int dx \, \int_{\epsilon}^{\infty} \, dz \left[4 \left(\partial_{z} \Phi(z) \right)^{2} + \Lambda^{2} \, e^{-4\Phi(z)} \right]$$

 $z = \epsilon e^{-u}$ redefines the cMERA RG coordinate.

In the massless case

$$\Phi(z) = rac{1}{2} \log \Lambda z$$

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Complexity from Liouville Field Action.

From

$$\Phi(z) = rac{1}{2} \log \Lambda - rac{1}{2} arphi_L(z)$$

a Liouville Field Theory Action functional can be derived

$$\mathcal{C}_{\mathcal{A}} = \mathcal{A}_{L}[\varphi_{L}] = \frac{1}{4} \int dx \, \int_{\epsilon}^{\infty} dz \left[\left(\partial_{z} \varphi_{L}(z) \right)^{2} + e^{2\varphi_{L}(z)} \right]$$

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Complexity from Liouville Field Action.

LFT provides a quantum theory of 2D-gravity that is solved by the metric

$$ds^2 = e^{2arphi_L(z)} \left(dz^2 + dx^2
ight)$$

For the massless cMERA $\varphi_L(z) = -\log z$ we obtain the hyperbolic space

$$ds^2 = \frac{1}{z^2} \left(dz^2 + dx^2 \right)$$

cMERA Gates and Vertex Operators.

A cMERA gate \mathcal{G}_{cMERA} is implemented by a Hamiltonian action,

$$\mathcal{G}_{\mathrm{cMERA}}(z) \sim e^{-i\delta \,\widetilde{K}(z)} = e^{-i\delta \,\Phi(z)\,\mathcal{O}} \quad \mathcal{O} \in \mathfrak{su}(1,1)$$



cMERA Gates and Vertex Operators.

We consider the LFT correlation functions

$$\langle \prod_{i} e^{\alpha_{i} \varphi_{L}(z_{i})} \rangle = \int D\varphi_{L} e^{\mathcal{A}_{L}} \prod_{i} e^{\alpha_{i} \varphi_{L}(z_{i})}$$

$$V(\alpha_i, z_i) = e^{\alpha_i \varphi_L(z_i)}$$

The functional integral is dominated by configurations

$$\varphi_L(z) = -\log z \quad |z - z_i| \gg \varepsilon_i$$
$$\varphi_L(z) = -2\alpha_i \log |z - z_i| + \varphi_i + \cdots \quad |z - z_i| < \varepsilon_i$$

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cMERA Gates and Vertex Operators.



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cMERA Gates and Vertex Operators.

After proper regularization [Seiberg 90]

$$\mathcal{A}_L(\alpha_{i,z}) = \mathcal{A}_L(z)_{/\cup \varepsilon_i} + \sum_i \left(\alpha_i \, \varphi_i - 2\alpha_i^2 \log \varepsilon_i^2 \right)$$

the variation of the cMERA circuit complexity is given by

$$\Delta C_{\mathcal{A}} = \sum_{i} \left(\alpha_{i} \, \varphi_{i} - 2\alpha_{i}^{2} \log \varepsilon_{i}^{2} \right)$$

Ground State Entanglement and Liouville Mode.

Entanglement between two neighboring "sites" at scale u_* , is given by $\Phi(u_*)$.

This corresponds to the entanglement of the UV interval $L = \epsilon e^{-u_*}$.



Ground State Entanglement and Liouville Mode.

In the massless case of the free boson theory (c = 1)

$$S(L) = \frac{c}{6} \log \Lambda L$$

$$\frac{6}{c}S(z) = 2\Phi(z) = \log \Lambda - \varphi_L(z)$$

The entanglement structure of the quantum state arises from

$$\partial\left(\frac{6}{c}\,S(z)
ight) = \Lambda e^{-rac{6}{c}\,S(z)}$$

Conclusions.

- Path integral description of "complexity" in cMERA states
- cMERA flow extremize the action functional in the cMERA circuit path integral ≡ Circuit Complexity.
- CA in cMERA saturate Complexity Growth Limits. CGL as new constraints for a cMERA flow. Hints for an holographic description?
- Variational $\Phi \equiv$ Liouville mode in LFT \equiv Ground State Entanglement.

Future.

- Time dependent cMERA circuits.
- FiniteTemperature.



Future.

- Fermions and Bosons. SUSY. Complexity of SUSY states.
- Topological Phases of matter. No adiabatic paths starting from topologically trivial states.



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