NNLO predictions for top-quark pair production with leptonic final states

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The LHC performance is really good.

 \rightarrow large amount of top quark data.

 \rightarrow Observables (XS, differential distributions, mass,...) at % level precision

Top quark pair production is important:

- parameter estimation
- Standard Model precision measurements
- background for many physics searches for SM
- ... and beyond

Necessity of precise theory predictions for production and decay!

$t\overline{t}$ in theory and experiment

Experiment:

- Signature: *b*-jets, leptons, missing energy (depending on the decay channel)
- top-quarks are reconstructed from decay products
- Modeling extremely important
- measurements like $t\bar{t}$ (differential) x-sections rely on extrapolation in fiducial volumes

Theory

- theory of stable on-shell tops well under control: state of the art NNLO $(+EW) \rightarrow$ good modeling of reconstructed top data
- on-shell NWA and off-shell: up to now NLO
 - \rightarrow more realistic final state.
 - \rightarrow omit systematic uncertainties
 - \rightarrow spin information of top accessible!

Theory - resonant top quark pair production

Stable onshell tops, spin summed:

 Total inclusive cross sections @ NNLO+NNLL accuracy

[Czakon, Fiedler, Mitov '13]

Fully differential distributions
 @ NNLO

[Czakon, Fiedler, Heymes, Mitov '16]

+ EW corrections

[Czakon, Heymes, Mitov, Pagani, Tsinikos, Zaro '17]

Unstable tops + spin correlations:

 Approximate NNLO + NNLO decay

[Gao,Papanastasiou '17]



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Towards realistic final states at NNLO

 $\mathsf{Best:} \ \mathsf{full} \ \mathsf{off}\text{-shell} \ \mathsf{NNLO} \leftarrow \mathsf{not} \ \mathsf{feasible} \ \mathsf{yet}$

Here: Narrow-Width-Approximation at NNLO

Necessary ingredients

- Handling of real-radiation contribution:
 - facilitate cancellation of divergences between double-real, real-virtual, double-virtual contributions
 - difficult: double real radiation
 - \rightarrow new implementation of STRIPPER algorithm
- Virtual matrix elements:
 - one-loop \rightarrow no problem here
 - two-loop production and decay matrix elements
 - polarization needed

Motivation: $\Gamma_t \ll m_t$

Narrow-Width-Approximation

- On-shell top-quarks
- Factorization of top-decay
- Separations of QCD corrections
- Keep spin correlations



Polarized matrix elements

$t\bar{t}$ production and decay at NNLO QCD in NWA

Decay Production	LO	NLO	NNLO
LO		Standard NLO	[Bonciani'08] [Asatrian'08] [Beneke'08]
NLO	Standard NLO	Standard NLO	
NNLO	[Long,Czakon,RP '17]		

Polarised $t\bar{t}$ production amplitudes

Gluon channel

 $\mathcal{M} = \epsilon_{1\mu}(p_1)\epsilon_{2\nu}(p_2)M^{\mu\nu}$

 $M^{\mu
u}$ is a rank-2 Lorentz tensor

- Momentum conservation
- Transversality
- Equation of motion
- Parity conservation \rightarrow no γ_5

8 independent structures

(d = 4 dimensions)

$$M^{\mu\nu} = \sum_{j=1}^8 M_j T_j^{\mu\nu}$$

Quark channel



- Two disconnected fermion lines
- Connection by gluons+loops

4 independent structures

$$\mathcal{M} = \sum_{i=1}^{4} M_j T_j$$

with $T_j \sim \bar{v}_2 \Gamma_j u_1 \bar{u}_3 \Gamma'_j v_4$

Two loop polarised $t\bar{t}$ production amplitudes

projection method \rightarrow scalar coefficients with scalar integrals

Master integrals

- reduction of scalar integrals via in-house Laporta implementation
- new partially canonicalised
- numerical treatment of master with help of differential equation
 - ightarrow interpolation grid
- finite remainder functions
- full color and spin information



spin-density coefficients:

Subtraction framework

NNLO subtraction schemes

Handling real radiation contribution in NNLO calculations cancellation of infrared divergences

increasing number of available NNLO calculations with a variety of schemes

- qT-slicing [Catani, Grazzini, '07], [Ferrera, Grazzini, Tramontano, '11], [Catani, Cieri, DeFlorian, Ferrera, Grazzini, '12],
 [Gehrmann, Grazzini, Kallweit, Maierhofer, Manteuffel, Rathley, Torre, '14-15'],
 [Bonciani, Catani, Grazzini, Saresvan, Torre, '14-15'],
- N-jettiness slicing [Gaunt,Stahlhofen,Tackmann,Walsh, '15], [Boughezal,Focke,Giele,Liu,Petriello,'15-'16], [Bougezal,Campell,Ellis,Focke,Giele,Liu,Petriello,'15], [Campell,Ellis,Williams,'16]
- Antenna subtraction [Gehrmann, GehrmannDeRidder, Glover, Heinrich, '05-'08], [Weinzierl, '08, '09],

[Currie,Gehrmann,GehrmannDeRidder,Glover,Pires,'13-'17], [Bernreuther,Bogner,Dekkers,'11,'14], [Abelof,(Dekkers),GehrmannDeRidder,'11-'15], [Abelof,GehrmannDeRidder,Maierhofer,Pozzorini,'14], [Chen,Gehrmann,Glover,Jaquier,'15]

- Colorful subtraction [DelDuca,Somogyi,Troscanyi,'05-'13], [DelDuca,Duhr,Somogyi,Tramontano,Troscanyi,'15]
- Sector-improved residue subtraction (STRIPPER) [Czakon, 10, 11]

[Czakon,Fiedler,Mitov,'13,'15], [Czakon,Heymes,'14] [Czakon,Fiedler,Heymes,Mitov,'16,'17], [Bughezal,Caola,Melnikov,Petriello,Schulze,'13,'14], [Bughezal,Melnikov,Petriello,'11], [Caola,Czernecki,Liang,Melnikov,Szafron,'14], [Bruchseifer,Caola,Melnikov,'13-'14], [Caola, Melnikov, Röntsch,'17]

Outline of the scheme

- decomposition of phase space to disentangle overlapping singularities
- simple extraction of Laurent series in $\boldsymbol{\epsilon}$
- provides a general set of subtraction terms
- numerical treatment of integrated subtraction terms \rightarrow numerical cancellation of ϵ poles
- defined in *d*-dimensions → numerical evaluation not efficient
 - \Rightarrow four-dimensional formulation

Triple collinear factorization



originally: 5 sub-sectors

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Caola, Melnikov, Röntsch [hep-ph:1702.01352v1]

now: 4 sub-sectors

How to improve the STRIPPER subtraction scheme?



New phase space construction: Idea

Goal

Phase space construction with a minimal # of subtraction kinematics

Old construction

- Start with unresolved partons
- · Fill remaining phase space with Born configuration
- \rightarrow Non-minimal # kinematic configurations (e.g. single soft and collinear limits yield different configurations)

New construction

- Start with Born configuration
- Add unresolved partons (u_i)
- Cleverly adjust Born configuration to accommodate the *u_i*

Consequences

Features

- Minimal number of subtraction kinematics
- Only one DU configuration \rightarrow pole cancellation for each Born phase space point
- Expected improved convergence of invariant mass distributions, since $\tilde{q}^2=q^2$

Unintentional features

- Construction in lab frame
- Original construction of 't Hooft Veltman corrections [Czakon,Heymes'14] is spoiled

Implementation

- general (process-independent) STRIPPER implementation
 - new parameterization
 - new four-dimensional construction
- additional input: 1- and 2-loop finite remainder functions
- modifications for NWA:
 - onshell phase spaces
 - additional CS like dipole subtraction for decay part of NLO×NLO contributions (mixed subtractions)

 $pp
ightarrow t \overline{t}
ightarrow b \overline{b} l l' v v'$

differential cross section:

$$\mathrm{d}\sigma = \mathrm{d}\sigma_{t\bar{t}} \times \frac{\mathrm{d}\Gamma_t}{\Gamma_t} \times \frac{\mathrm{d}\Gamma_{\bar{t}}}{\Gamma_t}$$

decays - total width:

$$\Gamma_t = \Gamma(t \to bW^+) \sum_{ff'} \frac{\Gamma(W^+ \to ff')}{\Gamma_W}$$

decays - differential decays:

$$\mathrm{d}\Gamma_t = \mathrm{d}\Gamma(t \to bW^+) \sum_{f \in \{\mathbf{e},\mu\}} \frac{\mathrm{d}\Gamma(W^+ \to f\nu_f)}{\Gamma_W}$$

Consistent treatment of top width

Expansion in α_S :

$$\begin{split} \mathrm{d}\sigma_{t\bar{t}} &= \mathrm{d}\sigma_{t\bar{t}}^{(0)} + \alpha_{s}\mathrm{d}\sigma_{t\bar{t}}^{(1)} + \alpha_{s}^{2}\mathrm{d}\sigma_{t\bar{t}}^{(2)} \\ \mathrm{d}\Gamma_{t(\bar{t})} &= \mathrm{d}\Gamma_{t(\bar{t})}^{(0)} + \alpha_{s}\mathrm{d}\Gamma_{t(\bar{t})}^{(1)} + \alpha_{s}^{2}\mathrm{d}\Gamma_{t(\bar{t})}^{(2)} \\ \Gamma_{t} &= \Gamma_{t}^{(0)} + \alpha_{s}\Gamma_{t}^{(1)} + \alpha_{s}^{2}\Gamma_{t}^{(2)} \end{split}$$

Consistent expansion in α_s :

$$d\sigma^{\text{LO}} \equiv d\sigma^{\text{LO}\times\text{LO}}$$
$$d\sigma^{\text{NLO}} = d\sigma^{\text{NLO}\times\text{LO}} + d\sigma^{\text{LO}\times\text{NLO}} - \frac{2\Gamma_t^{(1)}}{\Gamma_t^{(0)}} d\sigma^{\text{LO}}$$
$$d\sigma^{\text{NNLO}} = d\sigma^{\text{NNLO}\times\text{LO}} + d\sigma^{\text{NLO}\times\text{NLO}} + d\sigma^{\text{LO}\times\text{NLO}}$$
$$- \frac{2\Gamma_t^{(1)}}{\Gamma_t^{(0)}} d\sigma^{\text{NLO}} + \left(\frac{3\Gamma_t^{(1)2}}{\Gamma_t^{(0)2}} - \frac{2\Gamma_t^{(0)}\Gamma_t^{(2)}}{\Gamma_t^{(0)2}}\right) d\sigma^{\text{LO}}$$

Considerations:

• treatment ensures after full incl. integration:

$$\sigma = \sigma_{t\bar{t}} BR(W \to I\nu)$$

 practice: just rescaling lower order contributions

First results

Setup

Setup: CMS

m_t	173.3 GeV
m _W	80.385 GeV
mZ	91.1876 GeV
Γ_W	2.0928 GeV
G _F	$1.16379 \cdot 10^{-5} \text{ Gev}^2$

- comparison to approximate NNLO calculation
- comparison to data provided by CMS [CMS '15]

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prelimiliary results!

Differential distributions



scale variations: $\mu = \mu_R = \mu_F \in [m_t/2, 2m_t]$

Comparison to CMS data - Leptons



scale variations: $\mu = \mu_R = \mu_F \in [m_t/2, 2m_t]$

Comparison to CMS data - b-jets



scale variations: $\mu = \mu_{R} = \mu_{F} \in [m_{t}/2, 2m_{t}]$

Summary

- STRIPPER: fully automated subtraction framework
- with Narrow-Width-Approximation implementation for t and leptonic W decays
- Polarized two-loop matrix elements
- first results for $t\bar{t}$ with leptonic final states

Outlook

- Phenomenological studies:
 - Spin correlations
 - fiducial cross sections
- hadronic W-decays \rightarrow all-jet, lepton + jets channels

Treat resolved particles in 4 dimensions (momenta and polarisations)

- Avoid unnecessary ϵ -orders of the matrix elements
- Avoid growth of dimensionality of phase space integrals

Make resolved phase space 4-dim. using measurement function, e.g.

$$F_n \to F_n \mathcal{N}^{-(n-1)\epsilon} \prod_{i=1}^{n-1} \delta^{(-2\epsilon)}(q_i)$$

This introduces errors of $\mathcal{O}(\epsilon)$ in all contributions! Needed: Separately finite single and double unresolved contributions

- using finiteness of NLO calculation
- shifting terms from single-unresolved to double-unresolved contributions
- corrections calculated in full generality