



UNIVERSITÄT
HEIDELBERG
Zukunft. Seit 1386.



INCLUSIVE PROMPT PHOTON PRODUCTION FROM THE CGC

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In collaboration with
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1

MOTIVATION

2

FRAMEWORK: CGC

3

POWER COUNTING?

4

SOME RESULTS

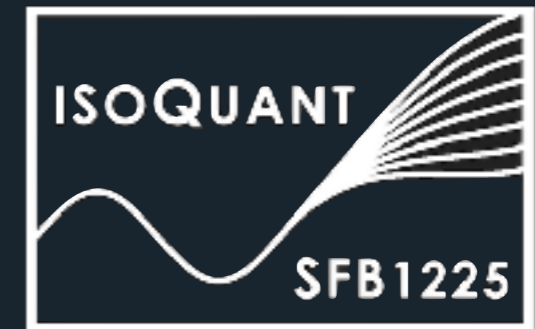
5

**SUMMARY AND
OUTLOOK**



GLOBAL GOAL

Understand (nuclear) matter under extreme conditions

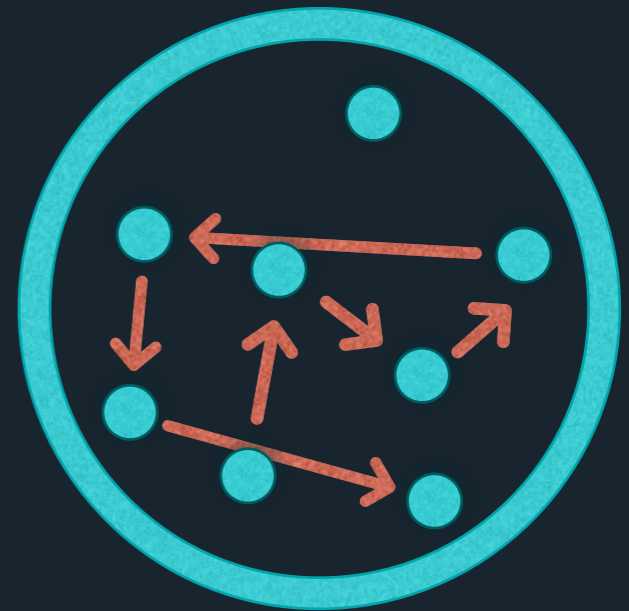


LOCAL GOAL

Use photons as probes for saturation
in dilute-dense collisions

LOCAL GOAL

Use photons as probes for saturation
in dilute-dense collisions



WHY PHOTONS?

No strong interaction

LOCAL GOAL

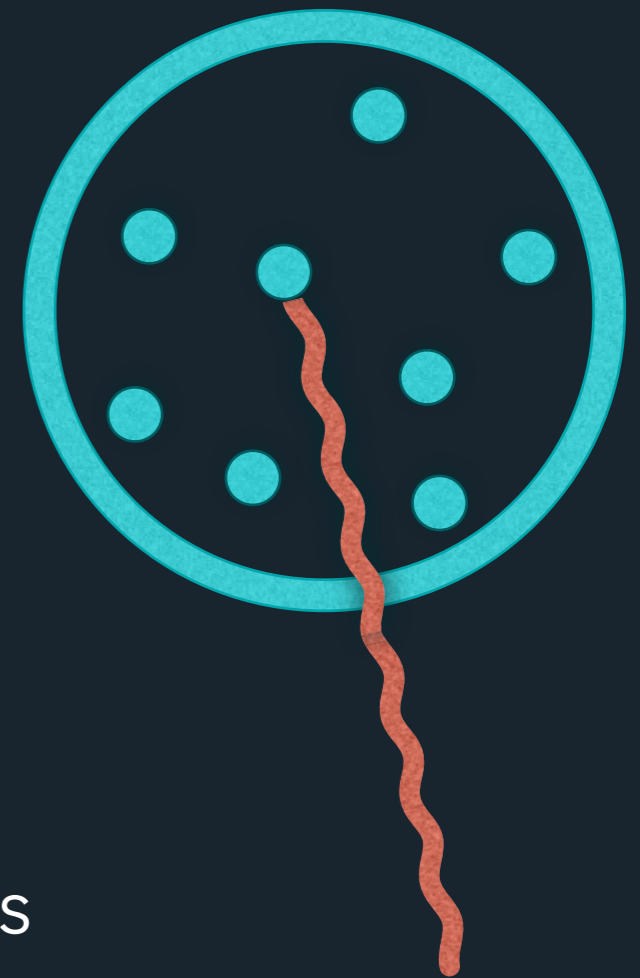
Use photons as probes for saturation
in dilute-dense collisions

WHY PHOTONS?

No strong interaction

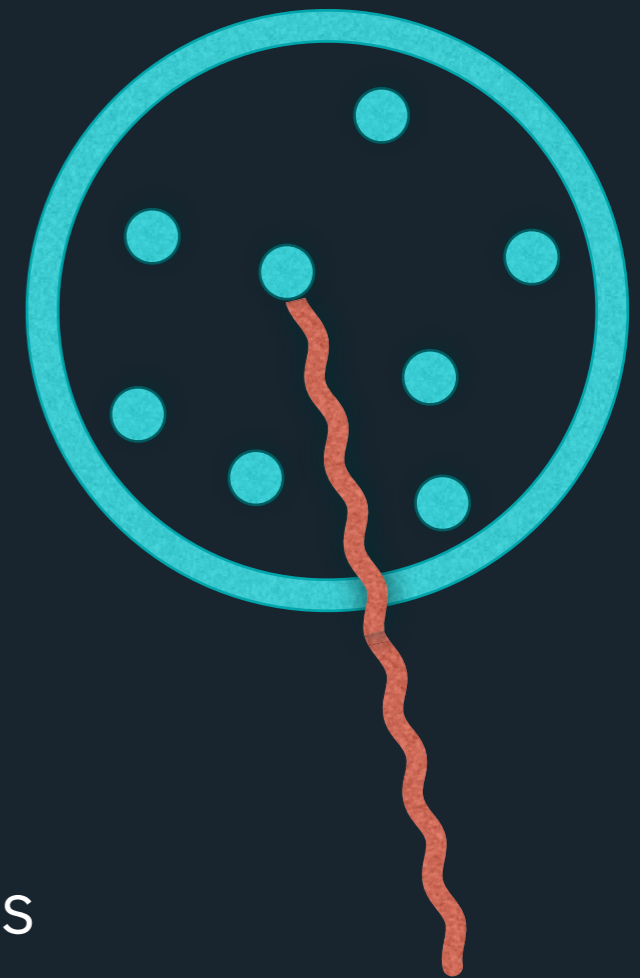


Clean Probes



LOCAL GOAL

Use photons as probes for saturation
in dilute-dense collisions



WHY PHOTONS?

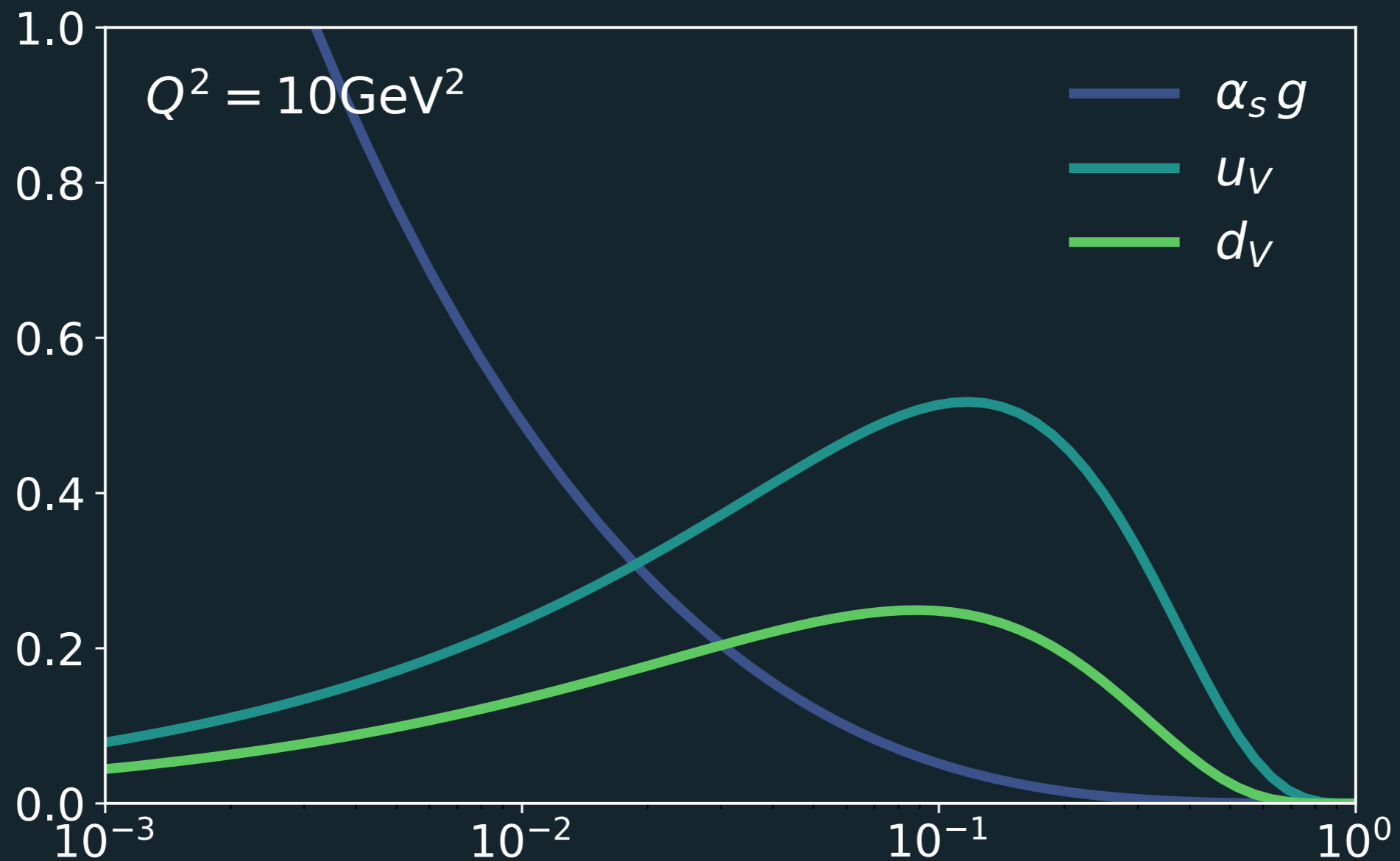
No strong interaction → Clean Probes

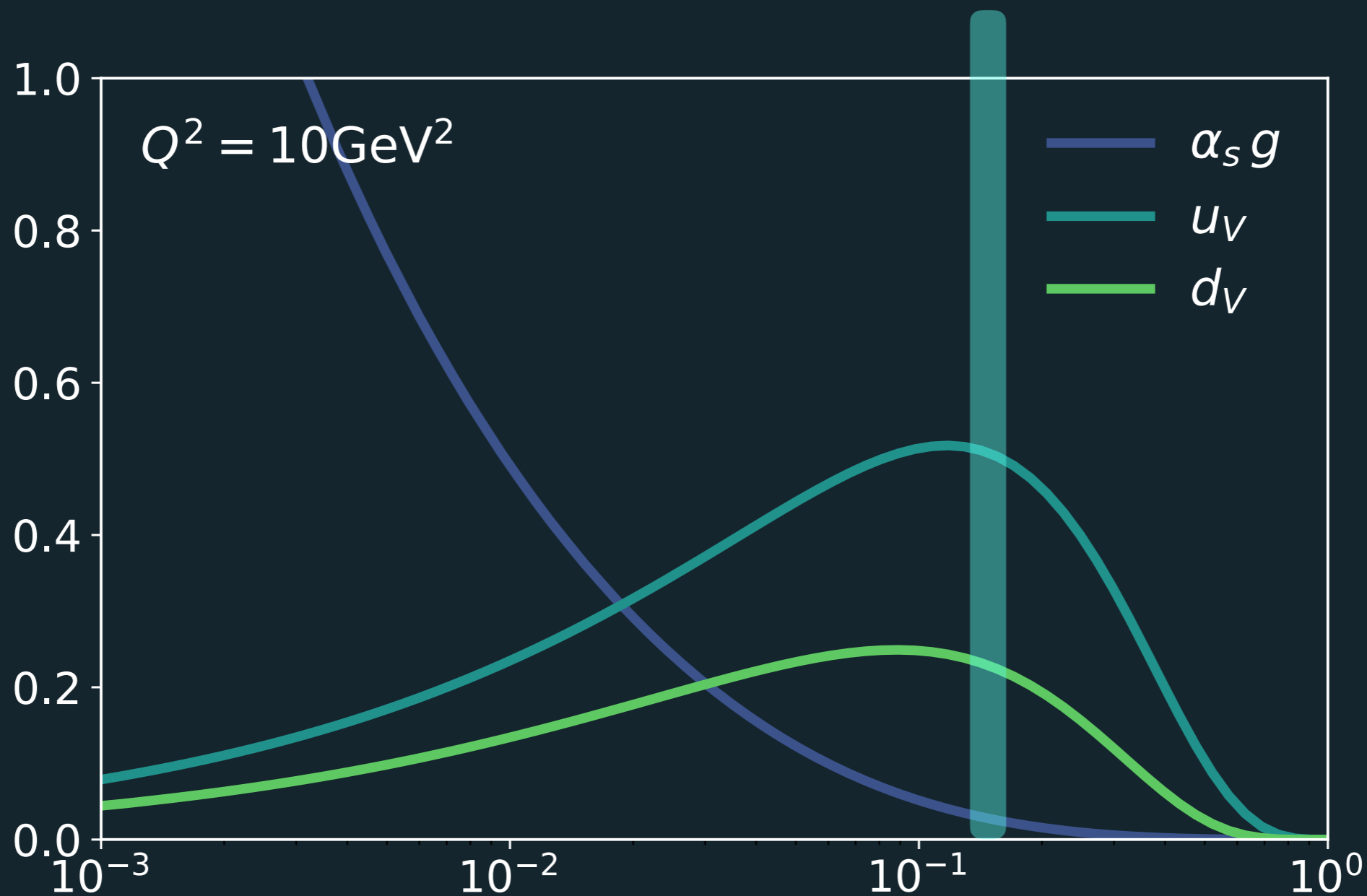
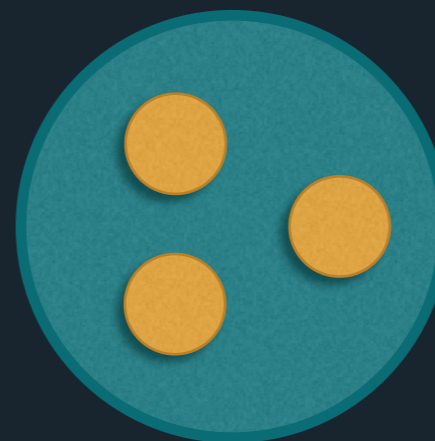
WHY DILUTE-DENSE?

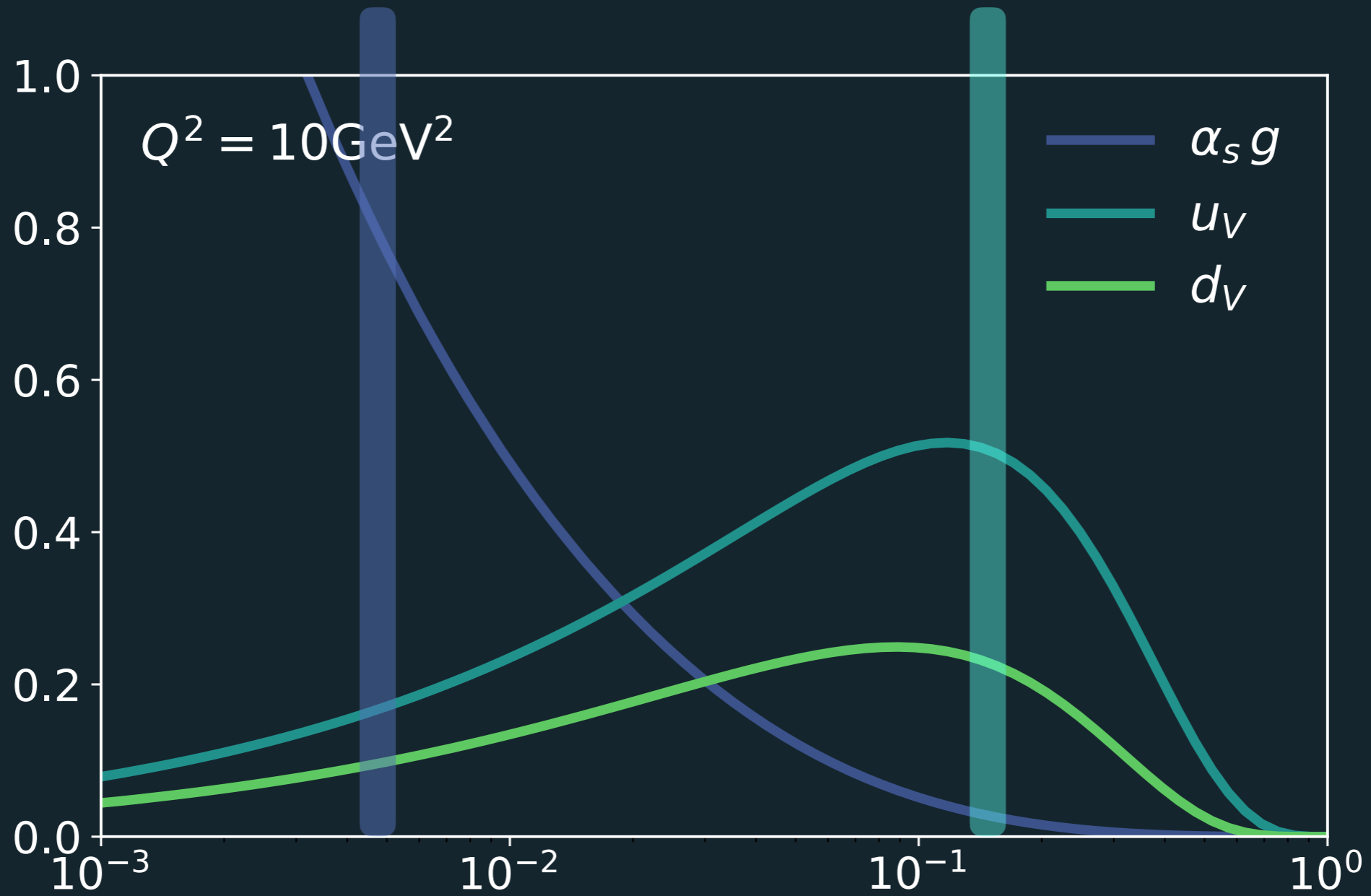
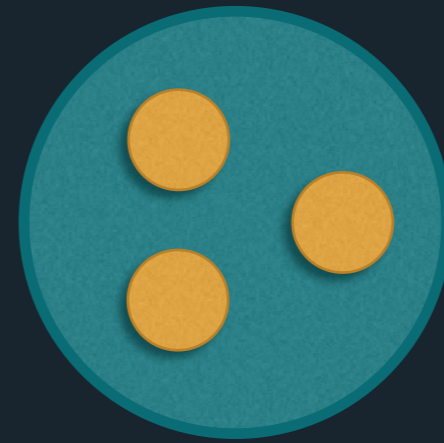
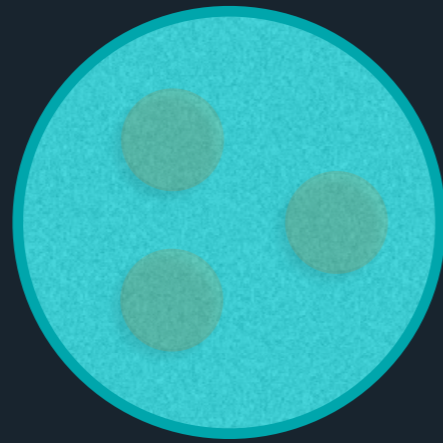
Use the 'known' to probe the 'unknown'

**WHAT IS
DILUTE-DENSE?**

PARTON DISTRIBUTION FUNCTIONS







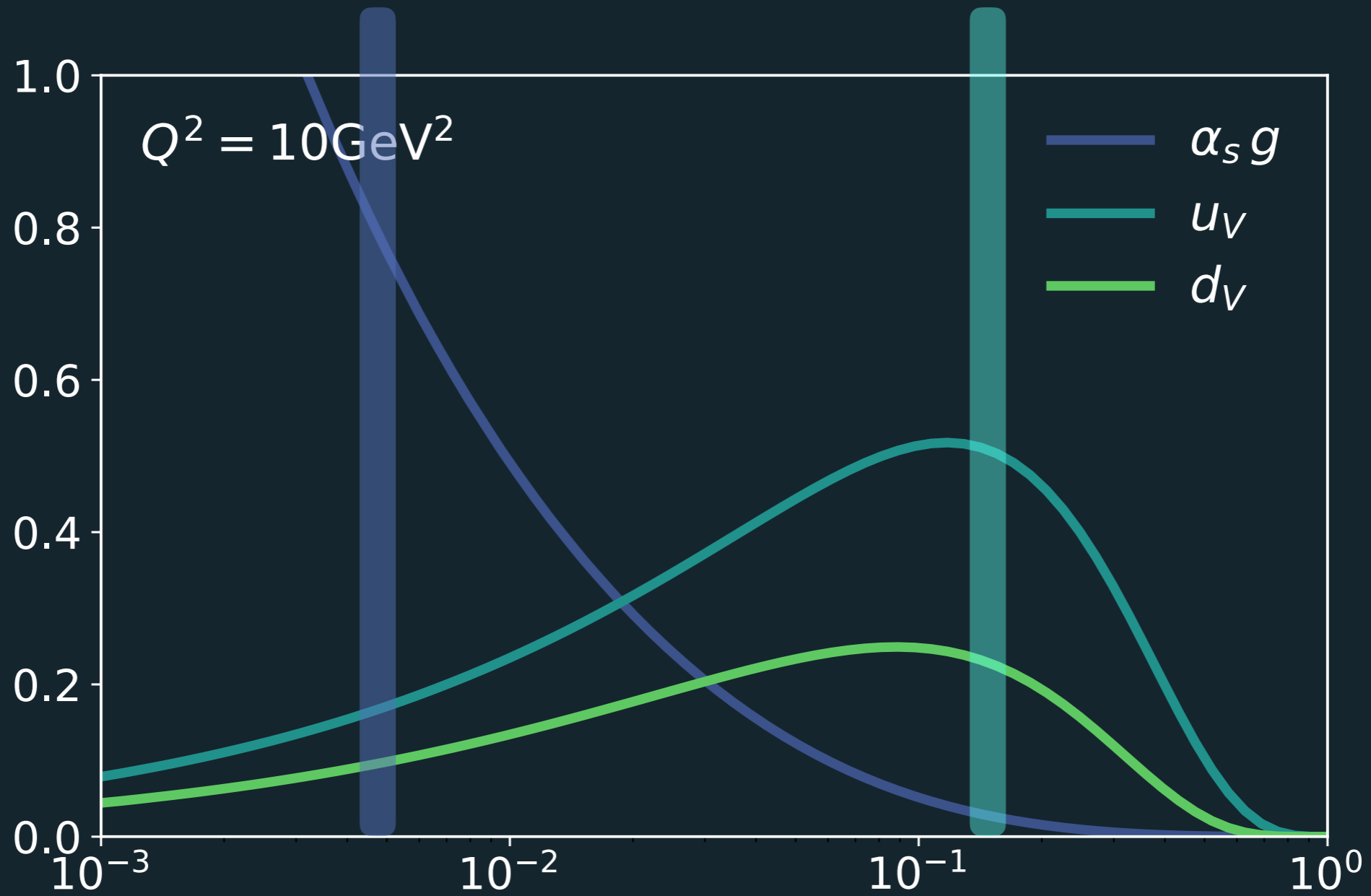
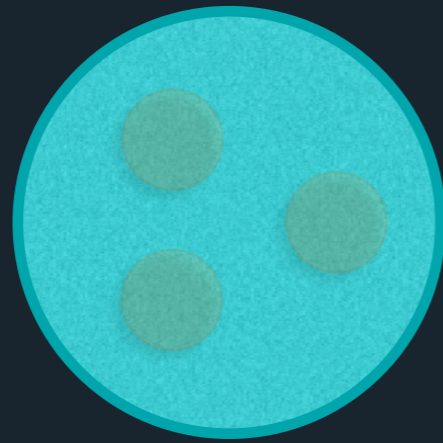
GLUON RISE

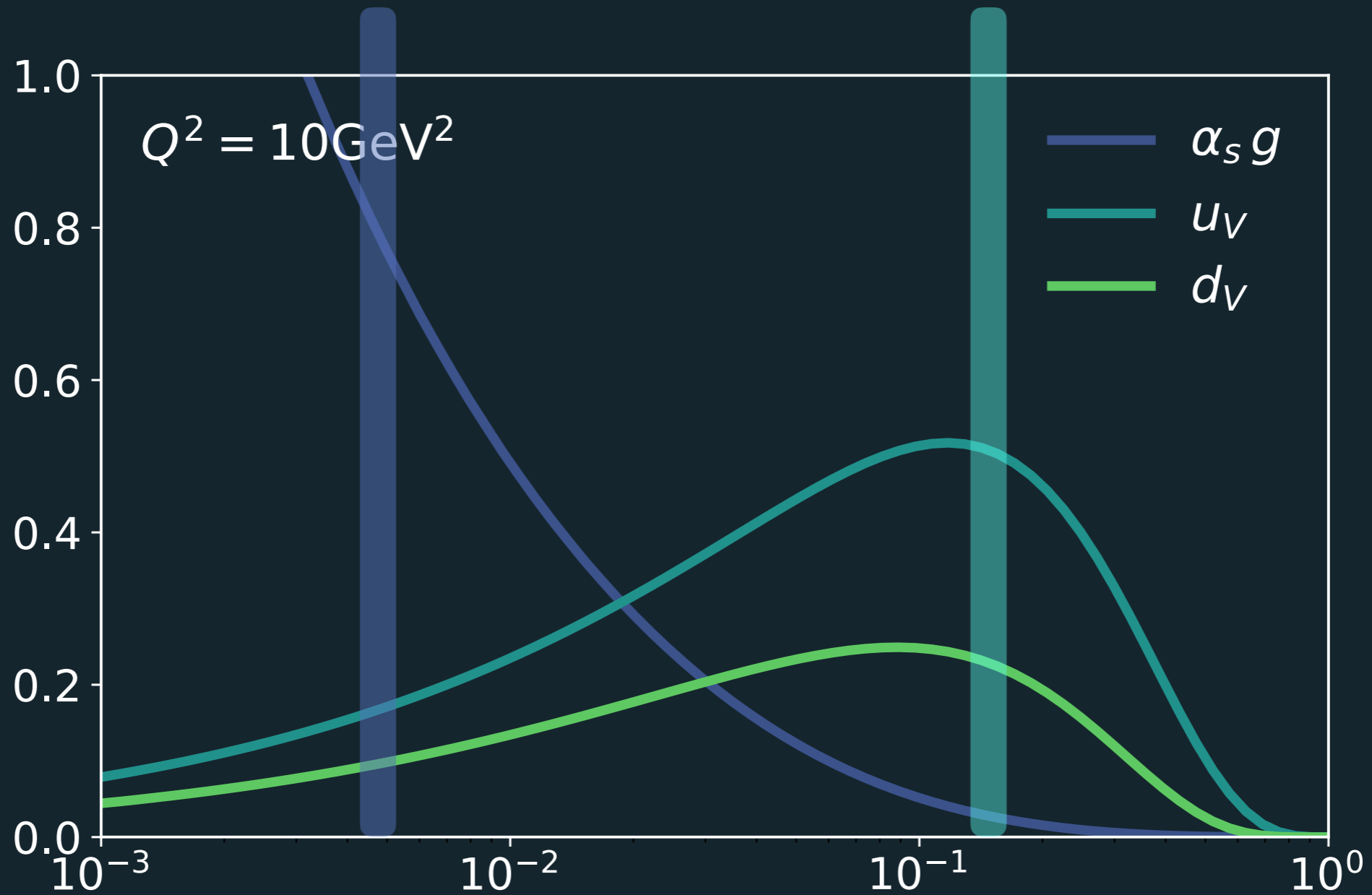
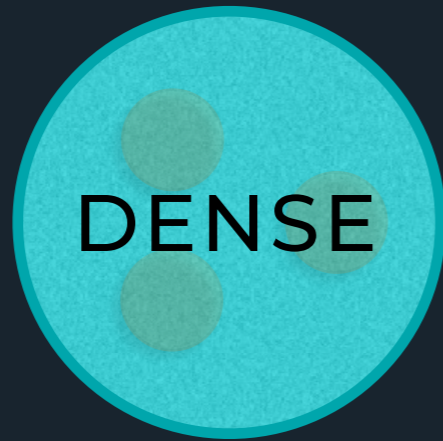


GLUON RECOMBINATION



SATURATION







HOW
TO
CATCH
A
GLUON
SOUP?

1

MOTIVATION

2

FRAMEWORK: CGC

3

POWER COUNTING?

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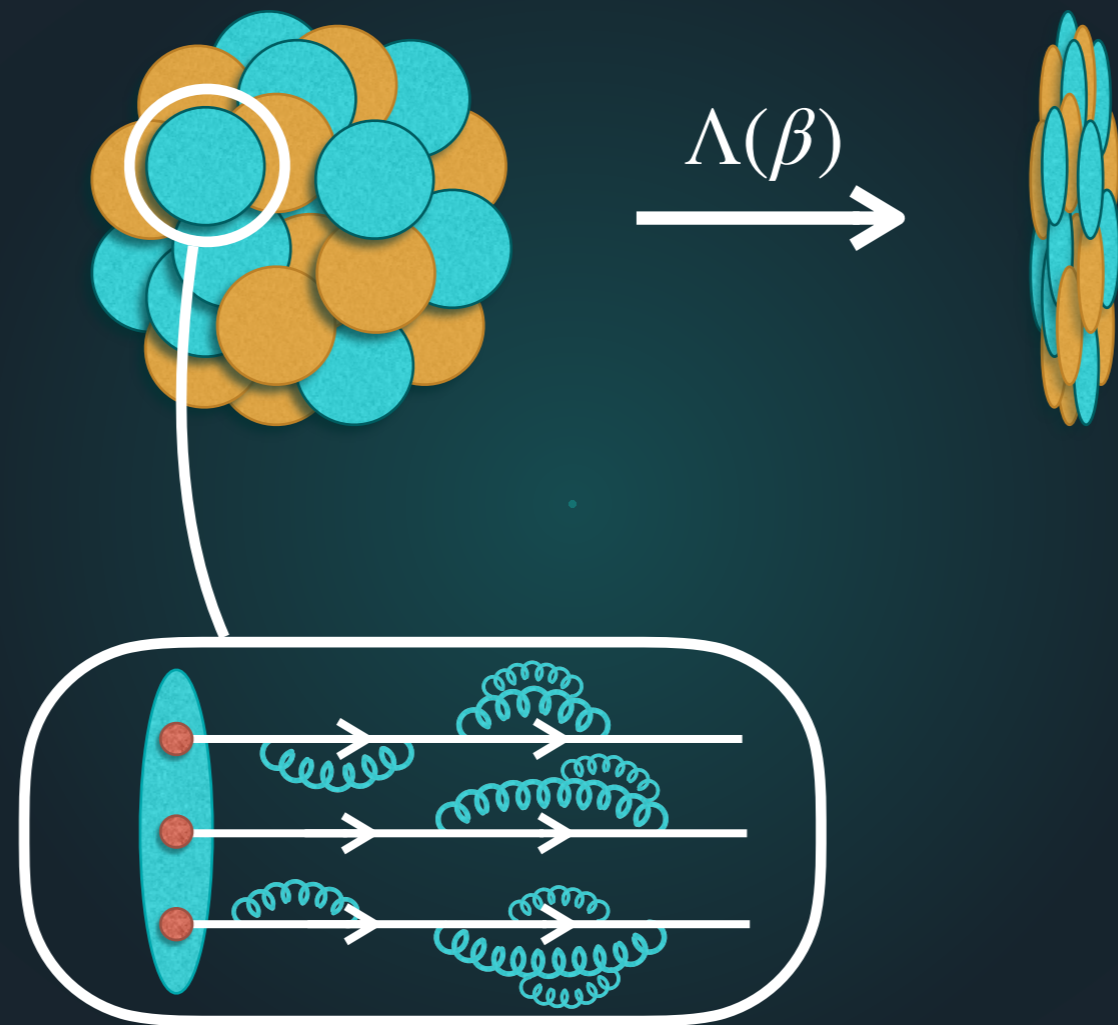
SOME RESULTS

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**SUMMARY AND
OUTLOOK**



COLOR GLASS CONDENSATE



Phys.Rev. D49 (1994) 2233-2241
Phys.Rev. D49 (1994) 3352-3355
Phys.Rev. D50 (1994) 2225-2233

COLOR GLASS CONDENSATE

GLUE

COLOR GLASS CONDENSATE

GLUE



Soft Partons

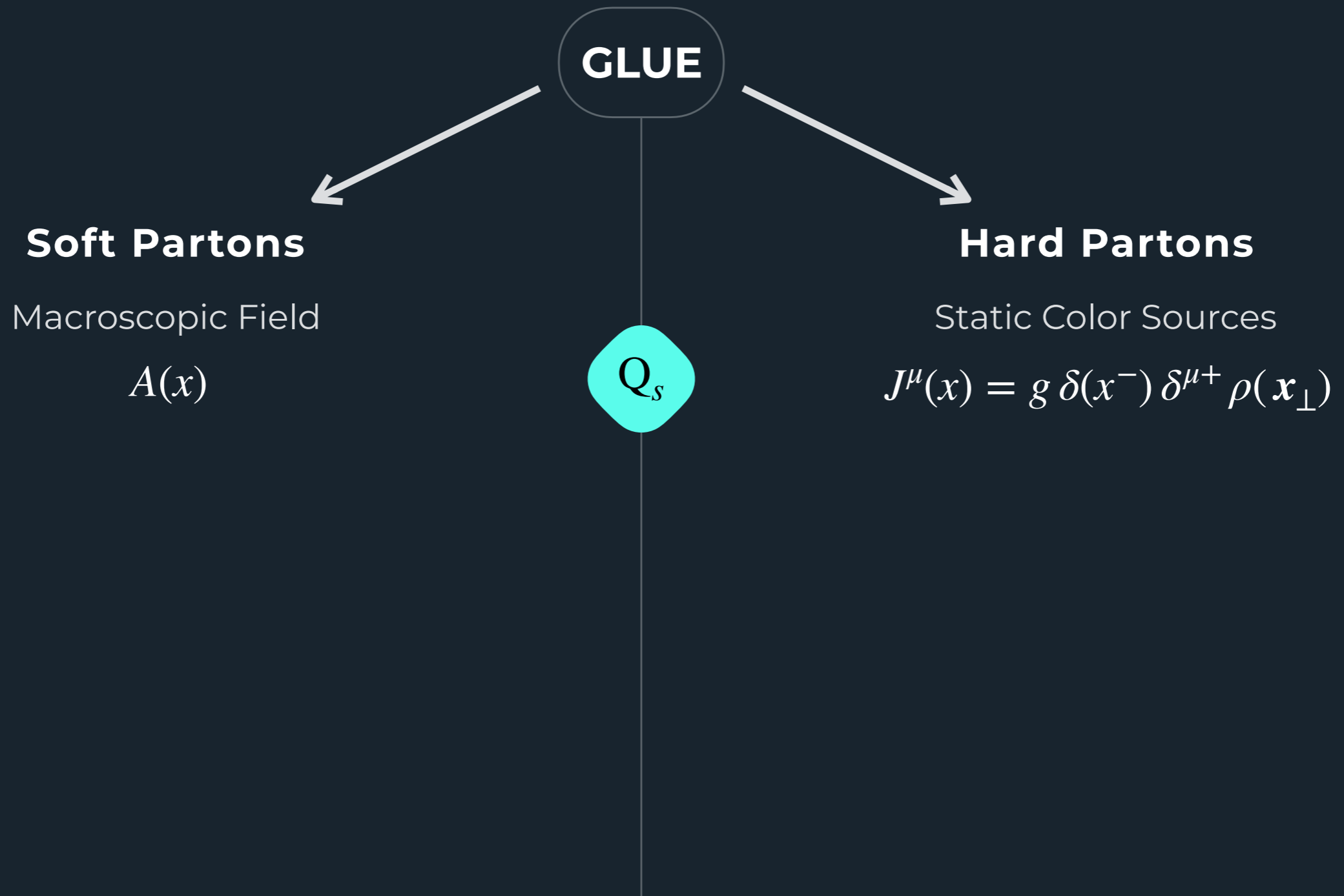
Macroscopic Field

$A(x)$

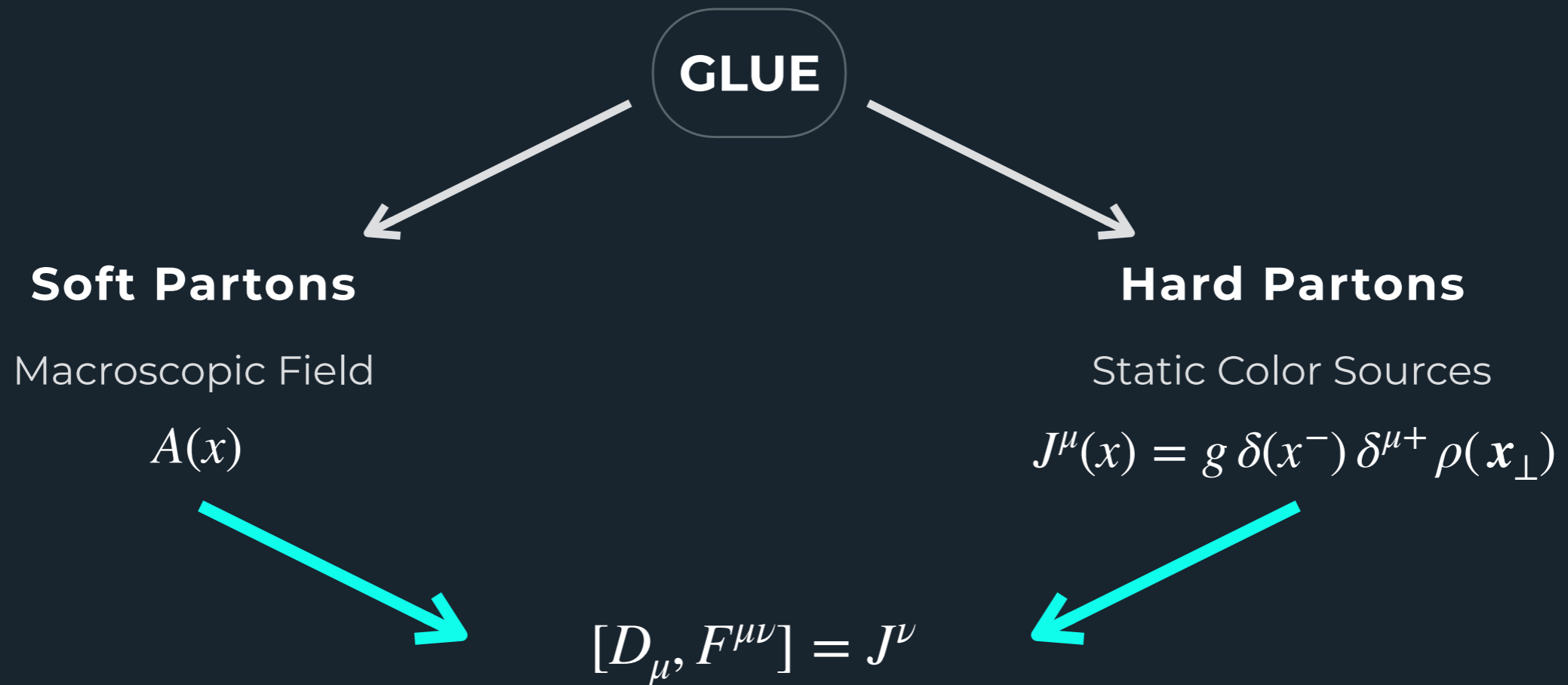
COLOR GLASS CONDENSATE



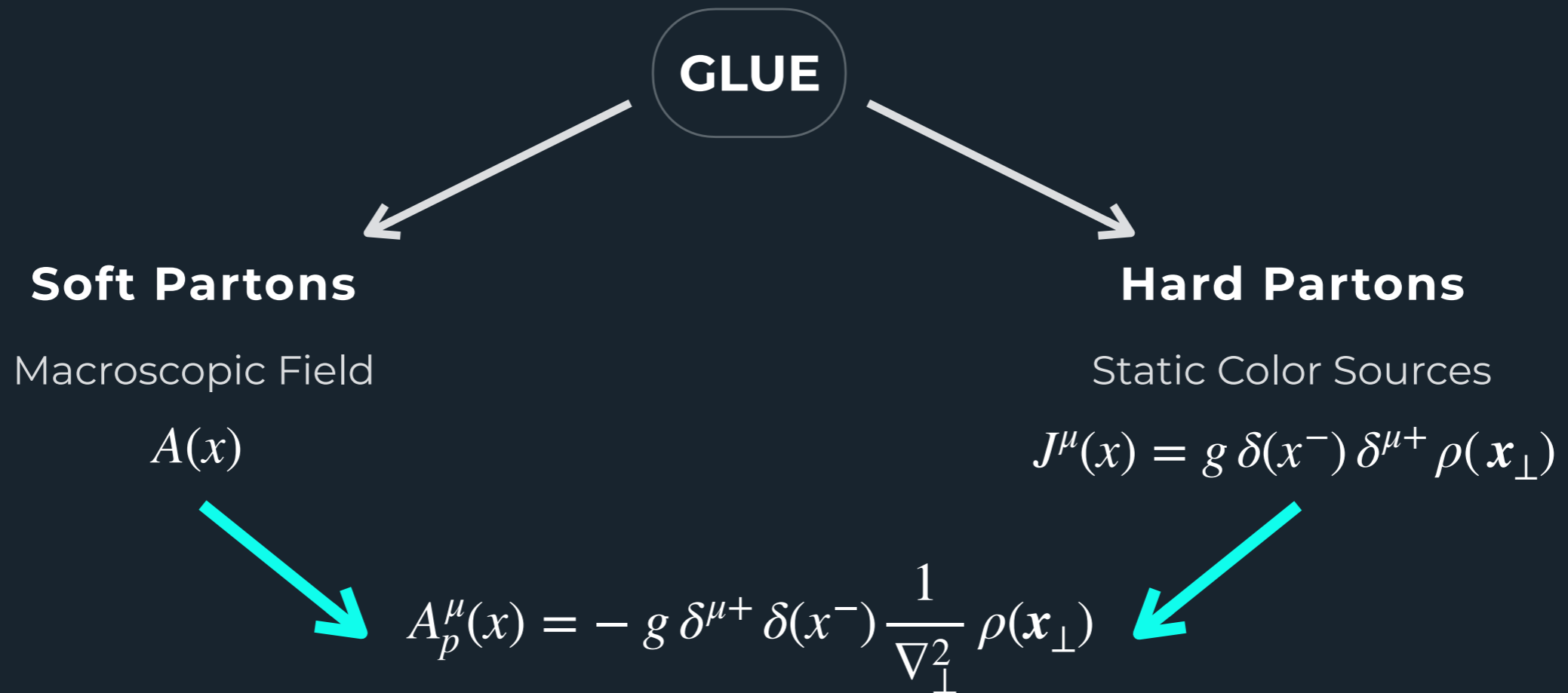
COLOR GLASS CONDENSATE



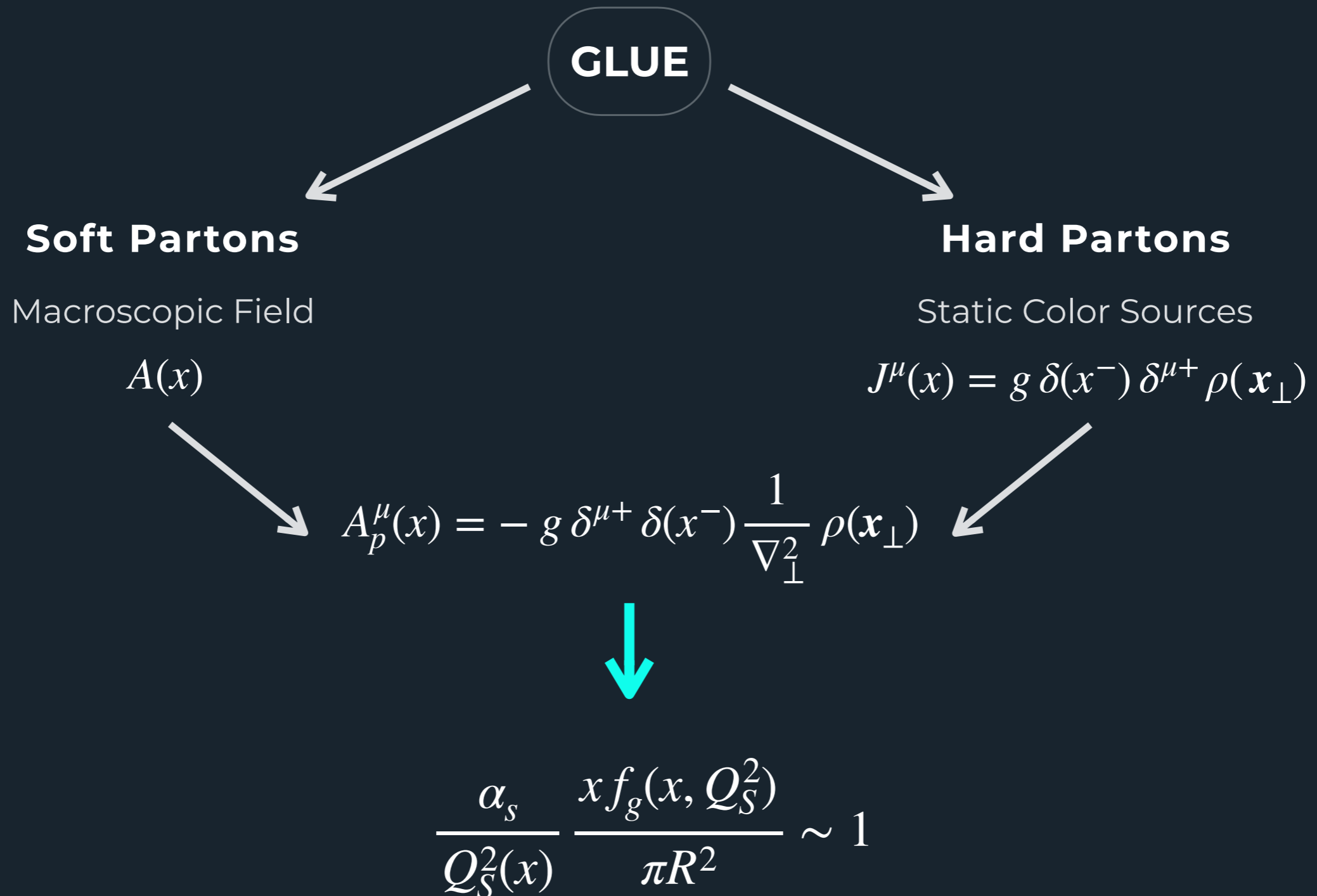
COLOR GLASS CONDENSATE



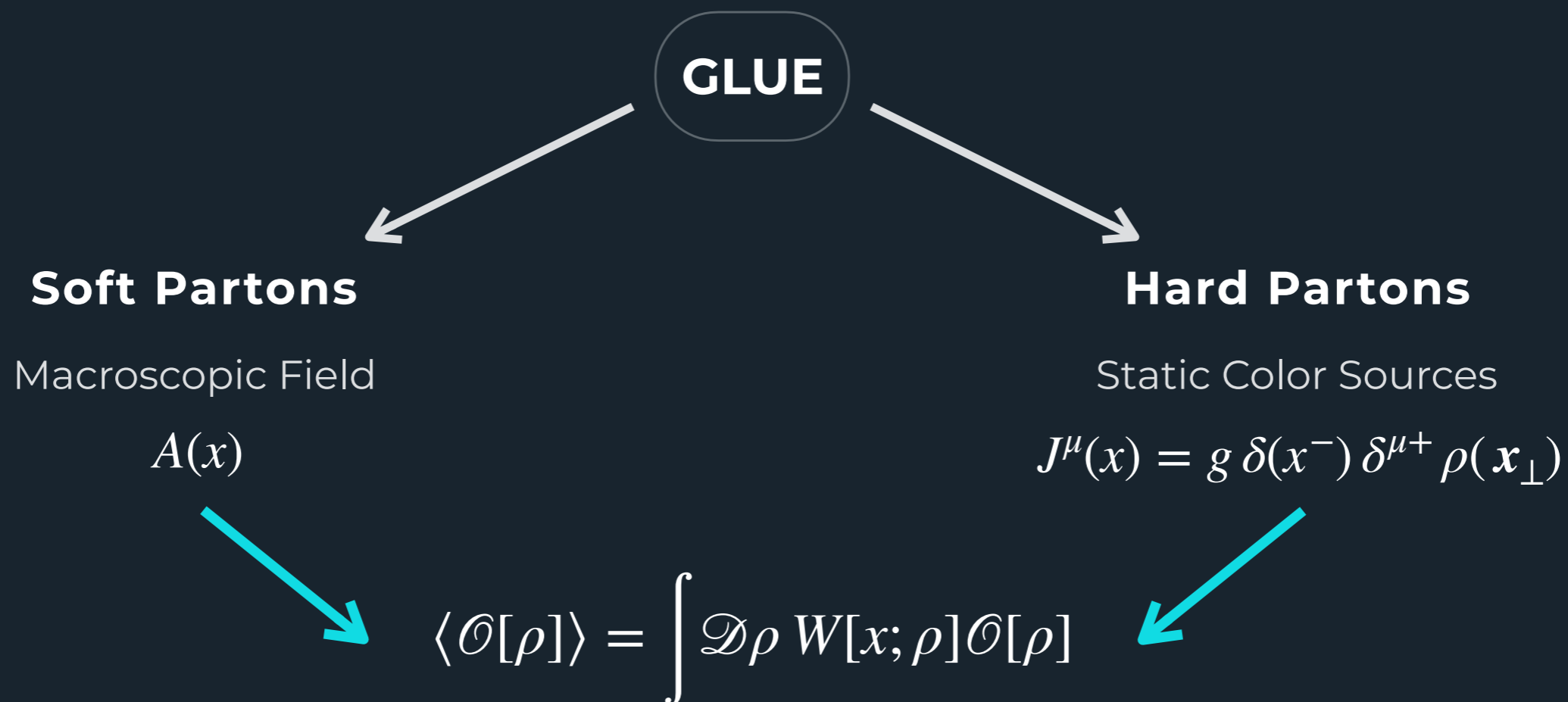
COLOR GLASS CONDENSATE



COLOR GLASS CONDENSATE

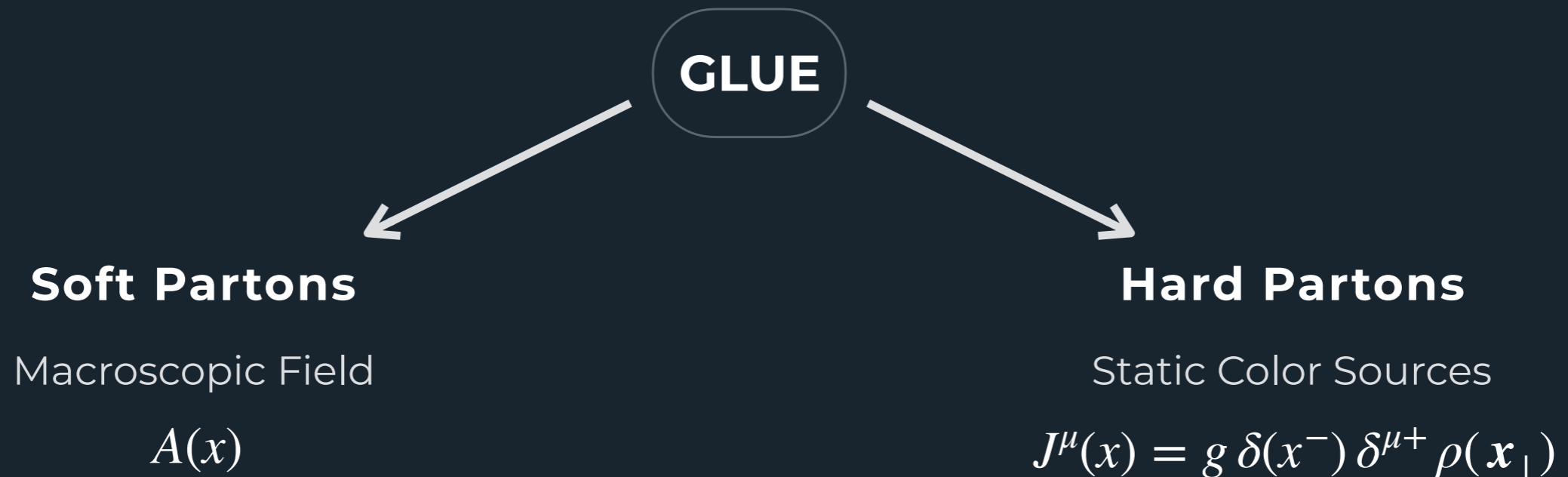


COLOR GLASS CONDENSATE



$W[x; \rho]$: gauge invariant probability distribution

COLOR GLASS CONDENSATE

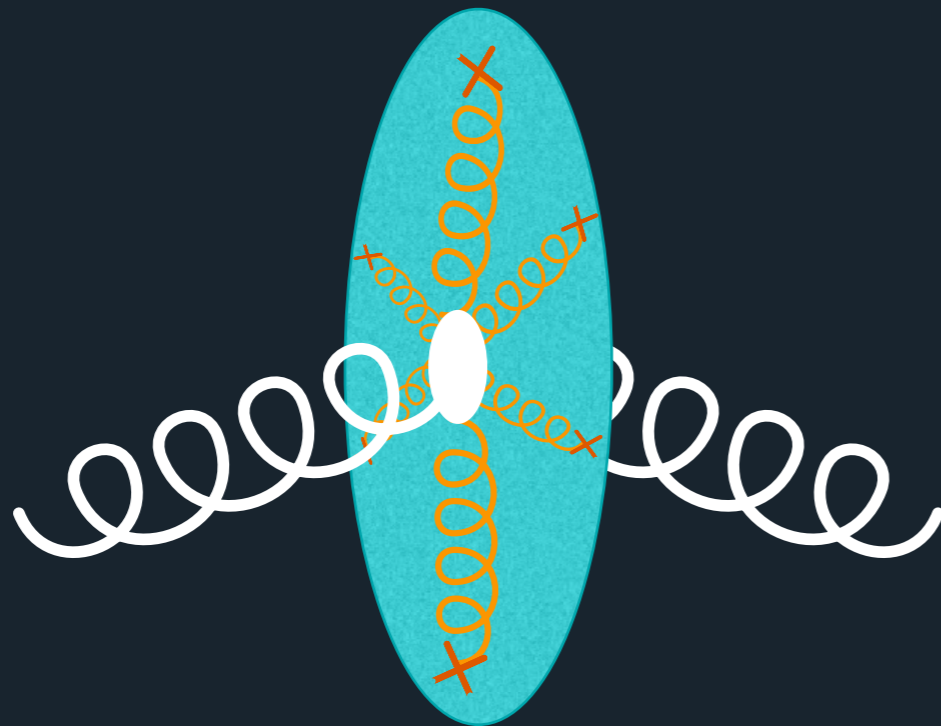


SPECIAL CASE

McLerran-Venugopalan Model

$$\langle \rho^a(\mathbf{x}_\perp) \rho^b(\mathbf{y}_\perp) \rangle = g^2 \delta^{ab} \mu^2 \delta^{(2)}(\mathbf{x}_\perp - \mathbf{y}_\perp)$$

PROPAGATION: GLUE



Is the gluon field (projectile) modified by the nuclear CGC?



Multiple Scatterings

$$A \times \text{oooooooo} = \rho_p \times \text{oooooooo} + A_p \times \text{oooooooo} \otimes \text{oooooooo}$$

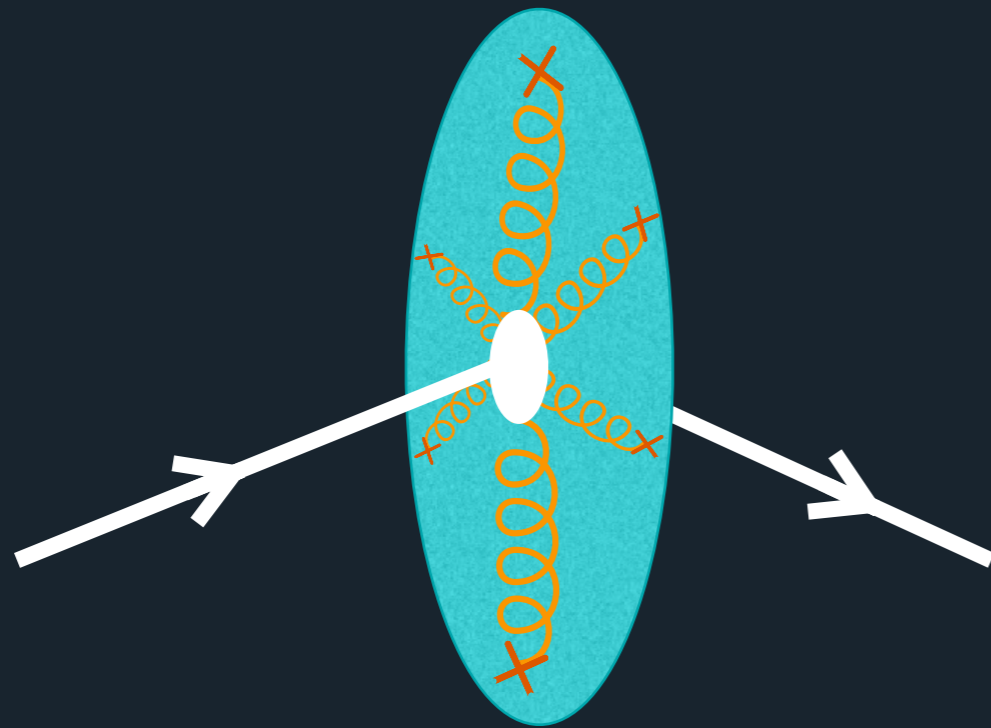


$$A^\mu(q) = A^\mu(q) + \frac{ig}{q^2 + iq^+\epsilon} \int_{k_\perp} \int_{x_\perp} e^{i(q_\perp - k_\perp) \cdot x} C^\mu(q, k_\perp) U(x_\perp) \frac{\rho_p(k_\perp)}{k_\perp^2}$$

Nucl.Phys. A743 (2004) 57-91

Nucl.Phys. A743 (2004) 13-56

PROPAGATION: QUARKS



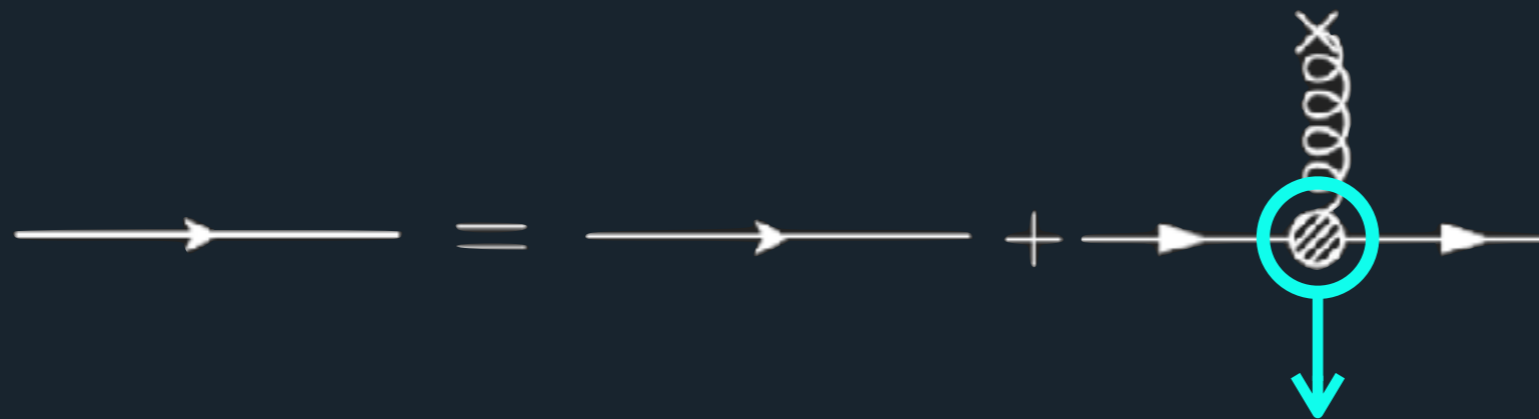
How do quarks behave through the CGC?



Multiple Scatterings



Dressed propagator



$$\mathcal{T}(k, p) = 2\pi \operatorname{sgn}(p^+) \gamma^+ \int_{x_\perp} e^{i\mathbf{x}_\perp \cdot \mathbf{k}_\perp} [\tilde{U}(\mathbf{x}_\perp) - 1]$$

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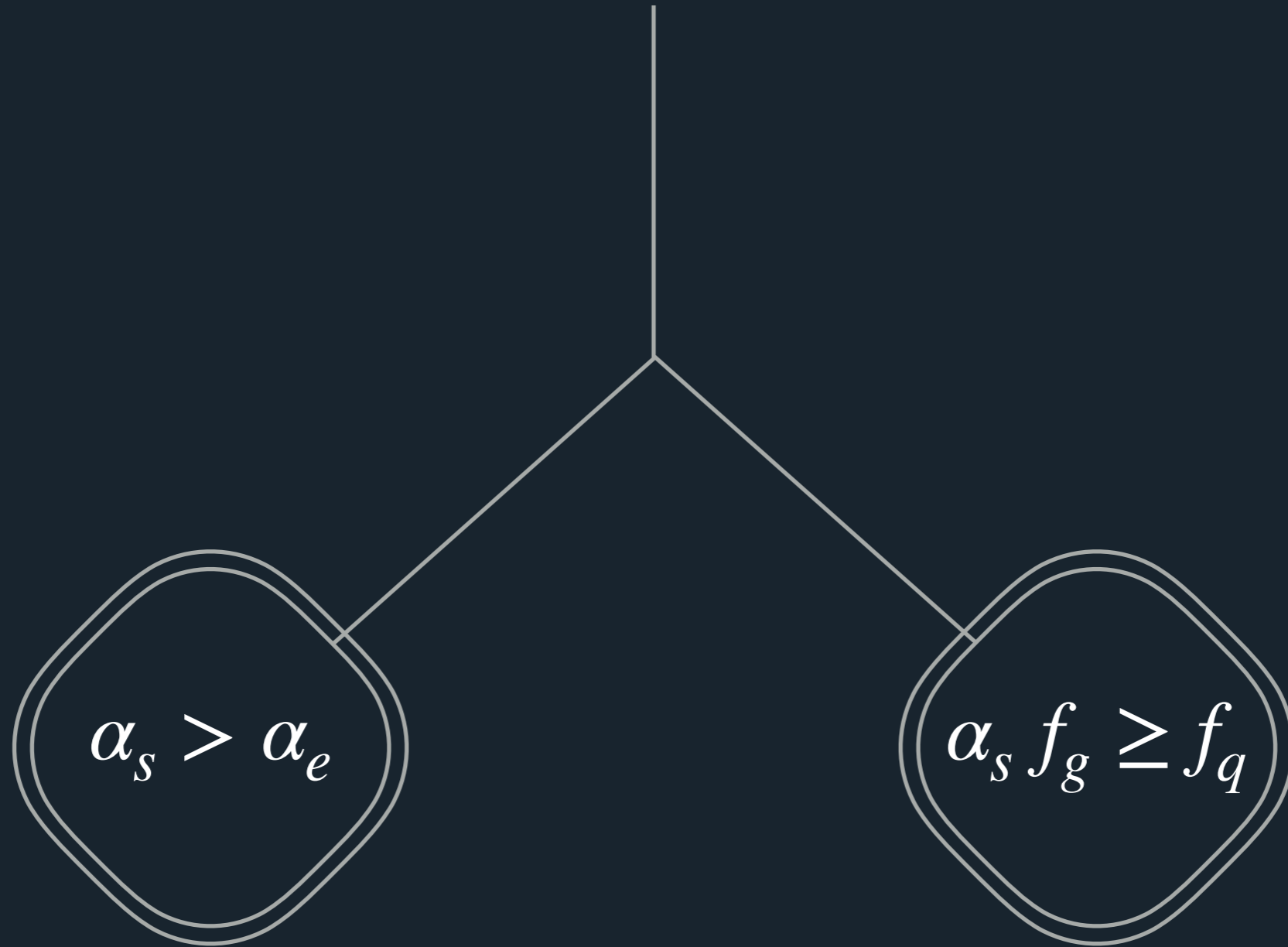
**SUMMARY AND
OUTLOOK**



POWER COUNTING

$\alpha_s > \alpha_e$

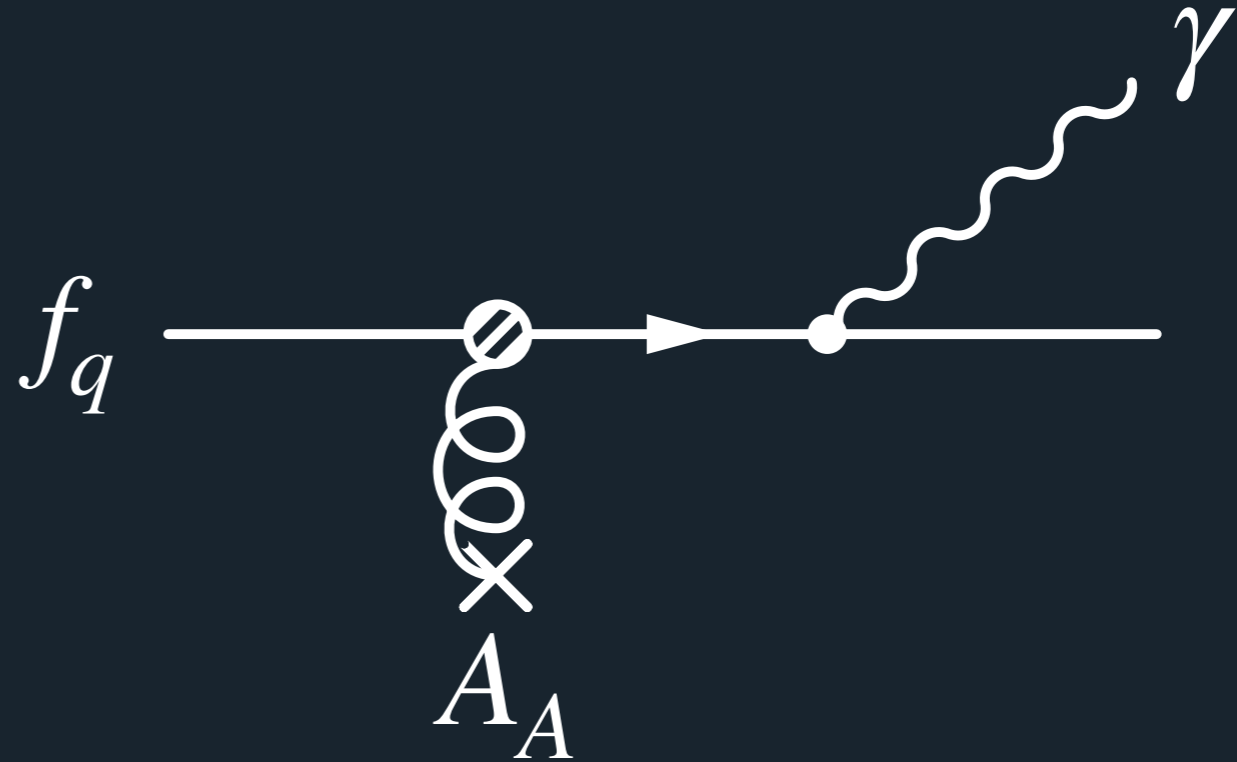
POWER COUNTING



LO [$\mathcal{O}(\alpha_e)$]

$$x \sim 10^{-2}$$

$$\alpha_s f_g \geq f_q$$



$$\frac{d\sigma}{d^2k_{\gamma\perp} d\eta_\gamma} = \frac{\alpha\alpha_s^2 q_f^2}{(2\pi)^8 C_F} \int_{\eta_q} \int_{\mathbf{q}_\perp} \int_{x_p} (2\pi) \delta(l^+ - q^+ - k_\gamma^+)$$

$$\times f_{q,p}(x_p, Q^2) \mathcal{N}_A(x_A, \mathbf{q}_\perp + \mathbf{k}_{\gamma\perp}) \theta_{LO}(q, k_\gamma)$$

Phys.Rev. D97 (2018) 054023

Phys. Rev. C 59 (1999) 1609

Phys. Rev. D 66 (2002) 014021

Nucl. Phys. A 741 (2004) 358

with
$$N(x_0, k) = \frac{1}{N_C} \langle U(k) U^\dagger(0) \rangle$$

LO [$\mathcal{O}(\alpha_e)$]



$$x \sim 10^{-2}$$

$$\alpha_s f_g \geq f_q$$

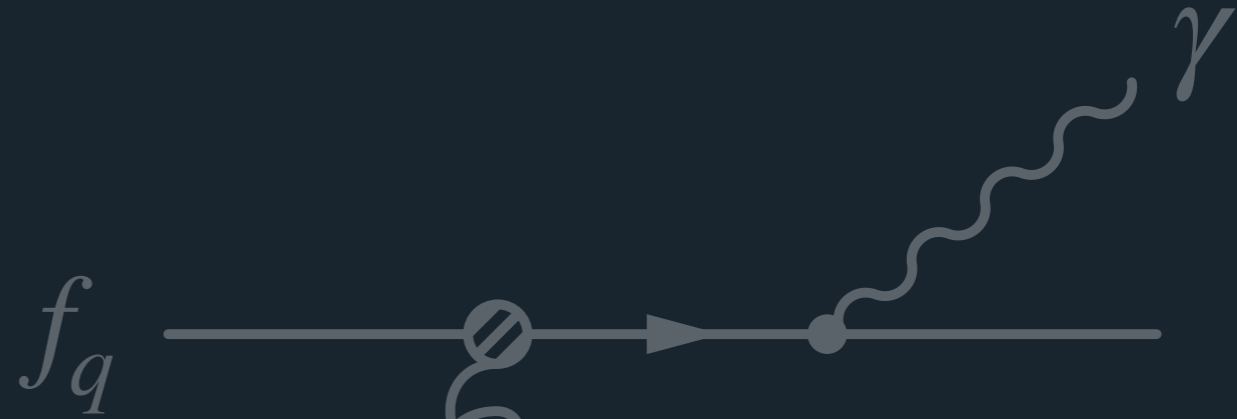
SUPPRESSED

$$\frac{d\sigma}{d^2k_{\gamma\perp} d\eta_\gamma} = \frac{\alpha\alpha_s^2 q_f^2}{(2\pi)^8 C_F} \int_{\eta_q} \int_{q_\perp} \int_{x_p} (2\pi) \delta(l^+ - q^+ - k_\gamma^+)$$

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LO [$\mathcal{O}(\alpha_e)$]



$$x \sim 10^{-2}$$

$$\alpha_s f_g \geq f_q$$

SUPPRESSED
(DEPENDENT ON X)

$$\frac{\alpha_s}{d^2k_{\gamma\perp} d\eta_\gamma} = \frac{\alpha_s^2}{(2\pi)^8 C_F} \int_{\eta_q} \int_{q_\perp} \int_{x_p} (2\pi) \delta(l^+ - q^+ - k_\gamma^+)$$

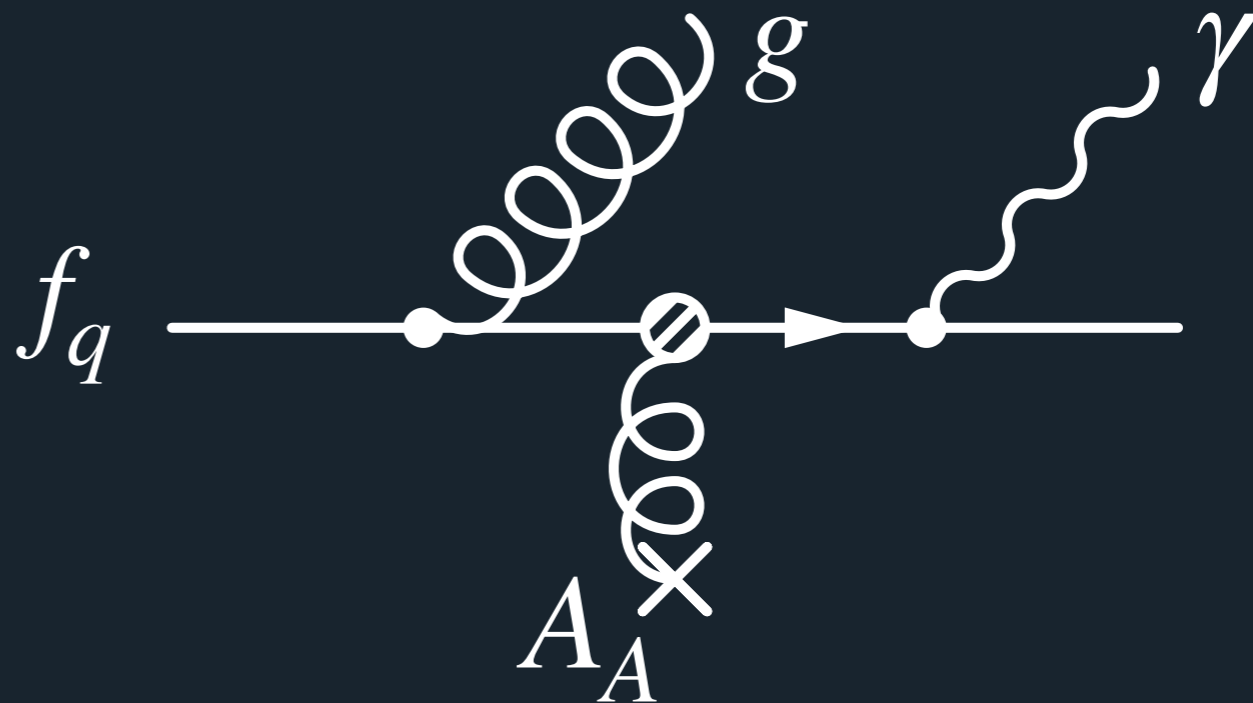
$$\times f_{q,p}(x_p, Q^2) \mathcal{N}_A(x_A, q_\perp + k_{\gamma\perp}) \theta_{LO}(q, k_\gamma)$$

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 Nucl. Phys. A 741 (2004) 358

NLO I [$\mathcal{O}(\alpha_s\alpha_e)$]

$$x \sim 10^{-2}$$

$$\alpha_s f_g \geq f_q$$



*** For inclusive photon, it can be included as evolution of the quark distributions

NLO II [$\mathcal{O}(\alpha_s\alpha_e)$]

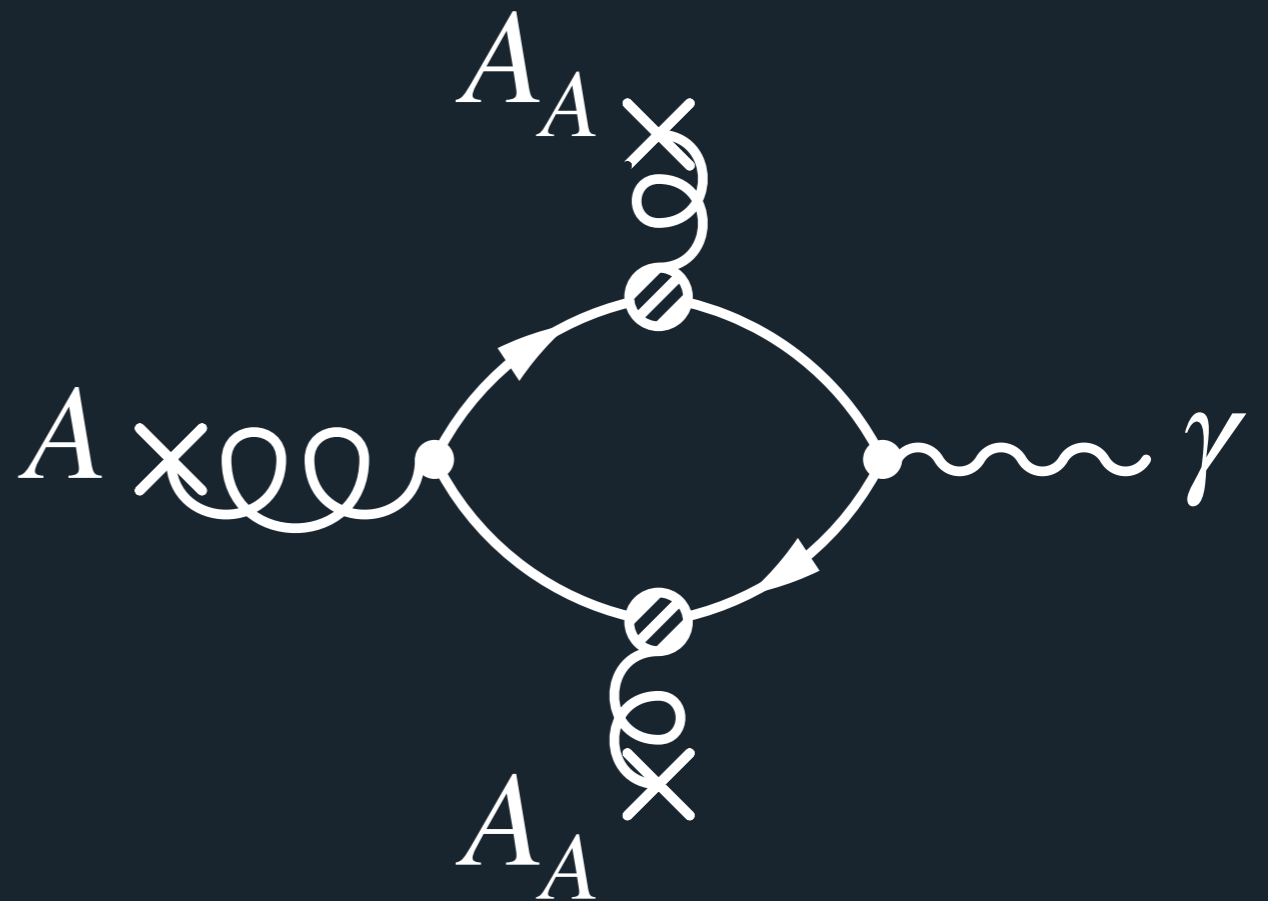
$$x \sim 10^{-2}$$

Enhanced by

$$\alpha_s f_g \geq f_q$$

Kinematically constrained

$$\text{Dominated by } k_{\perp}^2 = Q_{S,A}^2$$



NLO II [$\mathcal{O}(\alpha_s \alpha_e)$]

$$x \sim 10^{-2}$$

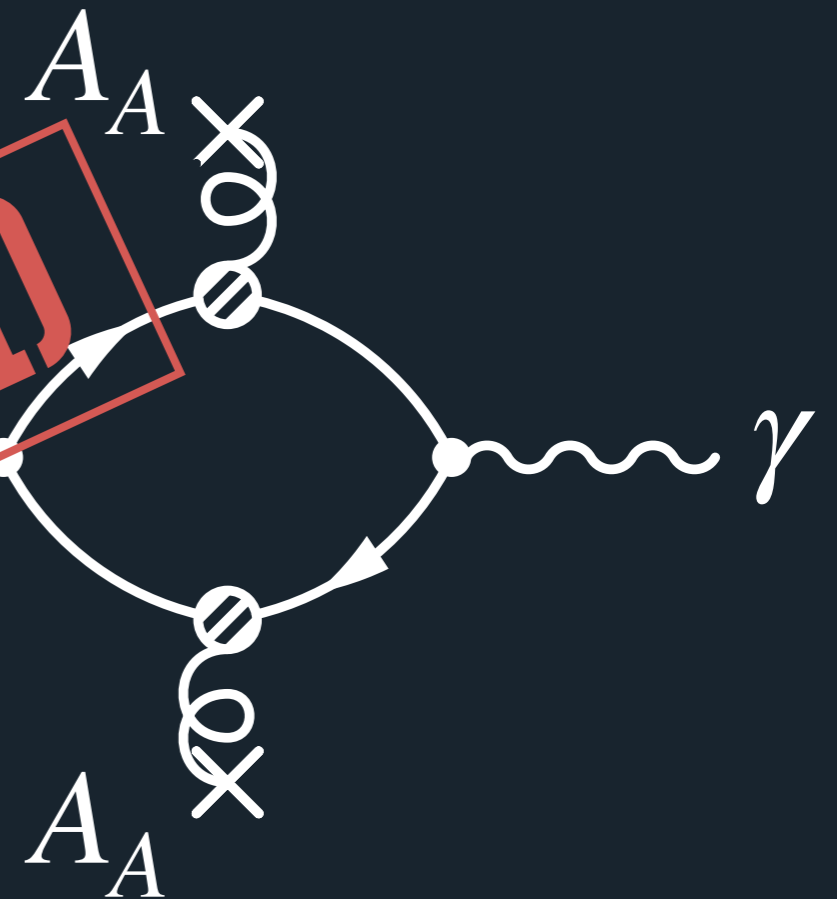
Enhanced by

$$\alpha_s f_g \geq f_q$$

Kinematically constrained

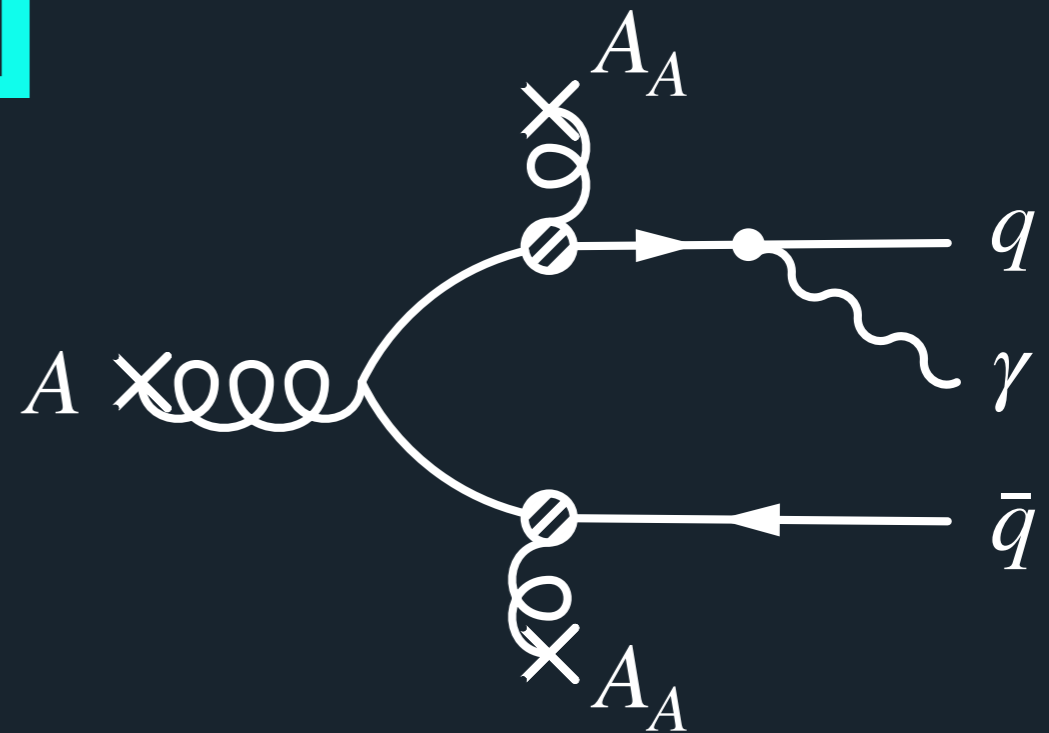
Dominated by $k_{\perp}^2 = Q_s^2$

SUPPRESSED



NLO III [$\mathcal{O}(\alpha_s \alpha_e)$]

$x \sim 10^{-2}$
 Enhanced by
 $\alpha_s f_g \geq f_q$



$$\frac{d\sigma}{d^2\mathbf{k}_{\gamma\perp} d\eta_\gamma} = \frac{\alpha\alpha_s^2 q_f^2}{(2\pi)^8 C_F} \int_{\eta_q, \eta_p} \int_{\mathbf{q}_\perp, \mathbf{p}_\perp, \mathbf{k}_\perp, \mathbf{k}_{1\perp}, \mathbf{k}_{2\perp}} \delta^{(2)}(\mathbf{P}_\perp - \mathbf{k}_{1\perp} - \mathbf{k}_{2\perp}) \frac{\varphi_p(\mathbf{k}_{1\perp})}{k_{1\perp}^2}$$

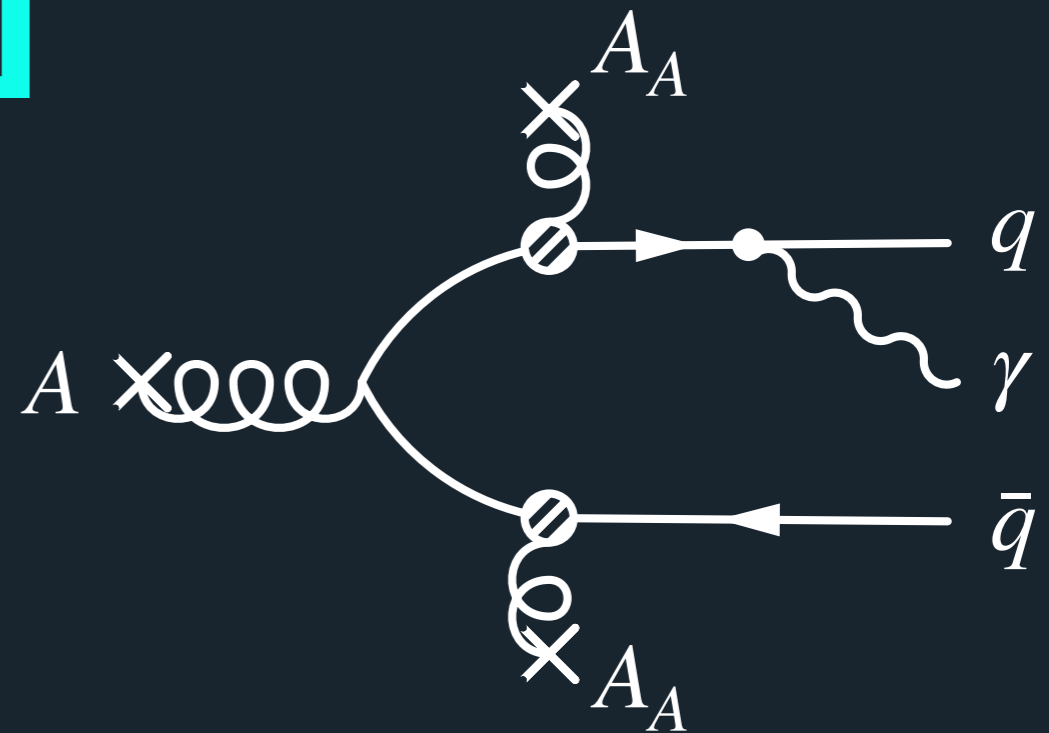
$$\times \left\{ \tau_{g,g}(\mathbf{k}_{1\perp}) \frac{\phi_A^{g,g}(\mathbf{k}_{2\perp})}{k_{2\perp}} \right.$$

$$+ 2\tau_{g,qq}(\mathbf{k}_{1\perp}, \mathbf{k}_\perp) \frac{\phi_A^{qq,g}(\mathbf{k}_\perp, \mathbf{k}_{2\perp})}{k_{2\perp}}$$

$$\left. + \tau_{qq,qq}(\mathbf{k}_{1\perp}, \mathbf{k}_\perp, \mathbf{k}_\perp) \frac{\phi_A^{qq,qq}(\mathbf{k}_\perp, \mathbf{k}_{2\perp})}{k_{2\perp}} \right\}$$

NLO III [$\mathcal{O}(\alpha_s \alpha_e)$]

$x \sim 10^{-2}$
Enhanced by
 $\alpha_s f_g \geq f_q$

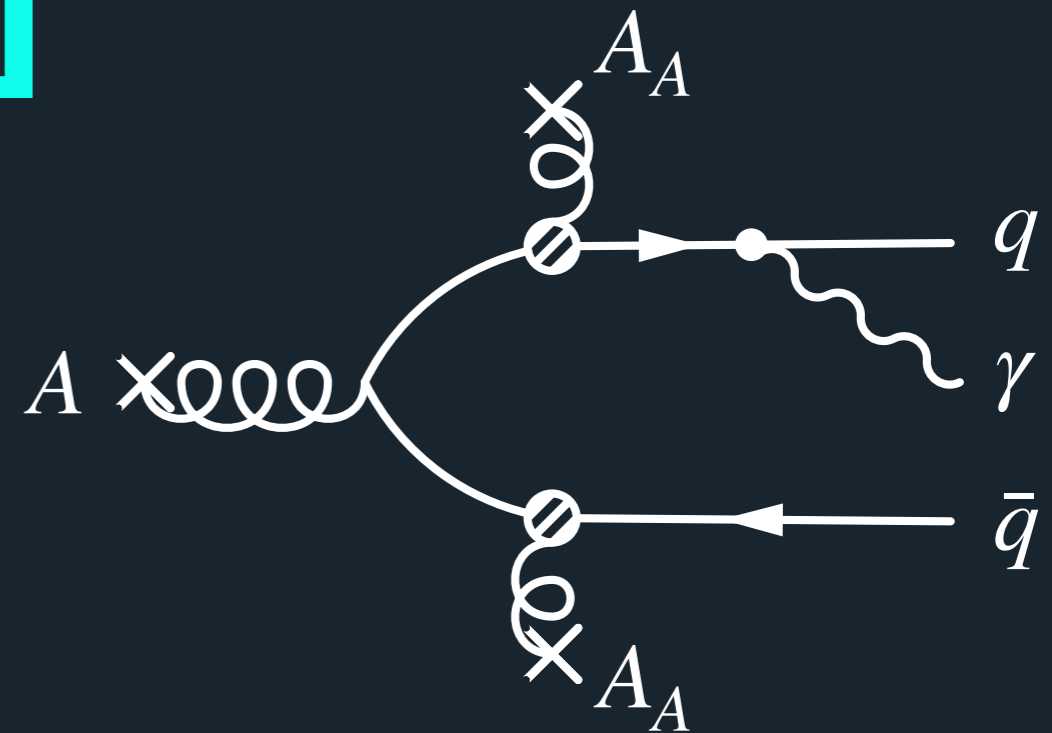


$$\frac{d\sigma}{d^2\mathbf{k}_{\gamma\perp} d\eta_\gamma} = \frac{\alpha\alpha_s^2 q_f^2}{(2\pi)^8 C_F} \int_{\eta_q, \eta_p} \int_{\mathbf{q}_\perp, \mathbf{p}_\perp, \mathbf{k}_\perp, \mathbf{k}_{1\perp}, \mathbf{k}_{2\perp}} \delta^{(2)}(\mathbf{P}_\perp - \mathbf{k}_{1\perp} - \mathbf{k}_{2\perp}) \frac{\varphi_p(\mathbf{k}_{1\perp})}{k_{1\perp}^2} \times \theta_{NLO}(\mathbf{k}_{1\perp}, \mathbf{k}_\perp, \mathbf{k}'_\perp) \mathcal{N}(x_0, \mathbf{k}_{2\perp}) \mathcal{N}(x_0, \mathbf{k}_\perp - \mathbf{k}_{2\perp})$$

AT LARGE N_C

NLO III [$\mathcal{O}(\alpha_s \alpha_e)$]

$x \sim 10^{-2}$
 Enhanced by
 $\alpha_s f_g \geq f_q$



$$\frac{d\sigma}{d^2k_{\gamma\perp} d\eta_\gamma} = \frac{\alpha\alpha_s^2 q_f^2}{(2\pi)^8 C_F} \int_{\eta_q, \eta_p} \int_{\mathbf{q}_\perp, \mathbf{p}_\perp, \mathbf{k}_{1\perp}, \mathbf{k}_{2\perp}} \delta^{(2)}(\mathbf{P}_\perp - \mathbf{k}_{1\perp} - \mathbf{k}_{2\perp}) \frac{\varphi_p(\mathbf{k}_{1\perp})}{k_{1\perp}^2}$$

$$\times \theta_{LT}(\mathbf{k}_{1\perp}, \mathbf{k}'_{\perp}) \tilde{\mathcal{N}}(x_0, \mathbf{k}_{2\perp})$$

AT LT

EVOLUTION

$$\frac{\partial \mathcal{N}(\mathbf{r}, x)}{\partial \log(x_0/x)} = \int d^2 \mathbf{r}_1 K^{run}(\mathbf{r}, \mathbf{r}_1, \mathbf{r}_2) [\mathcal{N}(\mathbf{r}_1, x) + \mathcal{N}(\mathbf{r}_1, x) - \mathcal{N}(\mathbf{r}, x) - \mathcal{N}(\mathbf{r}_1, x) \mathcal{N}(\mathbf{r}_2, x)]$$

EVOLUTION

$$\frac{\partial \mathcal{N}(\mathbf{r}, x)}{\partial \log(x_0/x)} = \int d^2 \mathbf{r}_1 K^{run}(\mathbf{r}, \mathbf{r}_1, \mathbf{r}_2) [\mathcal{N}(\mathbf{r}_1, x) + \mathcal{N}(\mathbf{r}_1, x) - \mathcal{N}(\mathbf{r}, x) - \mathcal{N}(\mathbf{r}_1, x) \mathcal{N}(\mathbf{r}_2, x)]$$



KERNEL

EVOLUTION

$$\frac{\partial \mathcal{N}(\mathbf{r}, x)}{\partial \log(x_0/x)} = \int d^2 \mathbf{r}_1 K^{run}(\mathbf{r}, \mathbf{r}_1, \mathbf{r}_2) [\mathcal{N}(\mathbf{r}_1, x) + \mathcal{N}(\mathbf{r}_1, x) - \mathcal{N}(\mathbf{r}, x) - \mathcal{N}(\mathbf{r}_1, x) \mathcal{N}(\mathbf{r}_2, x)]$$



KERNEL



NON LINEAR

EVOLUTION

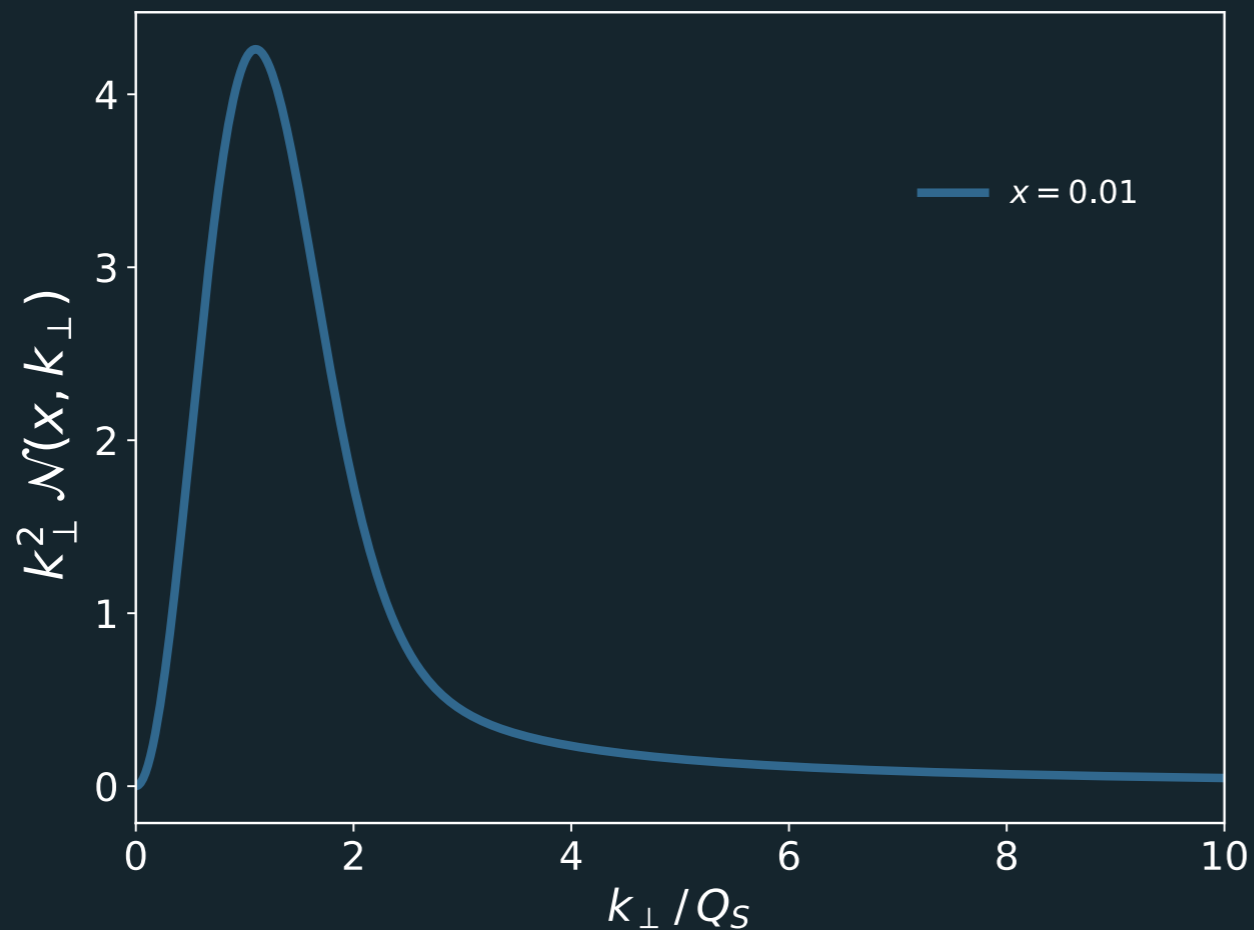
$$\frac{\partial \mathcal{N}(\mathbf{r}, x)}{\partial \log(x_0/x)} = \int d^2\mathbf{r}_1 K^{run}(\mathbf{r}, \mathbf{r}_1, \mathbf{r}_2) [\mathcal{N}(\mathbf{r}_1, x) + \mathcal{N}(\mathbf{r}_2, x) - \mathcal{N}(\mathbf{r}, x) - \mathcal{N}(\mathbf{r}_1, x) \mathcal{N}(\mathbf{r}_2, x)]$$



KERNEL



NON LINEAR



Initial condition:
MV MODEL



EVOLUTION

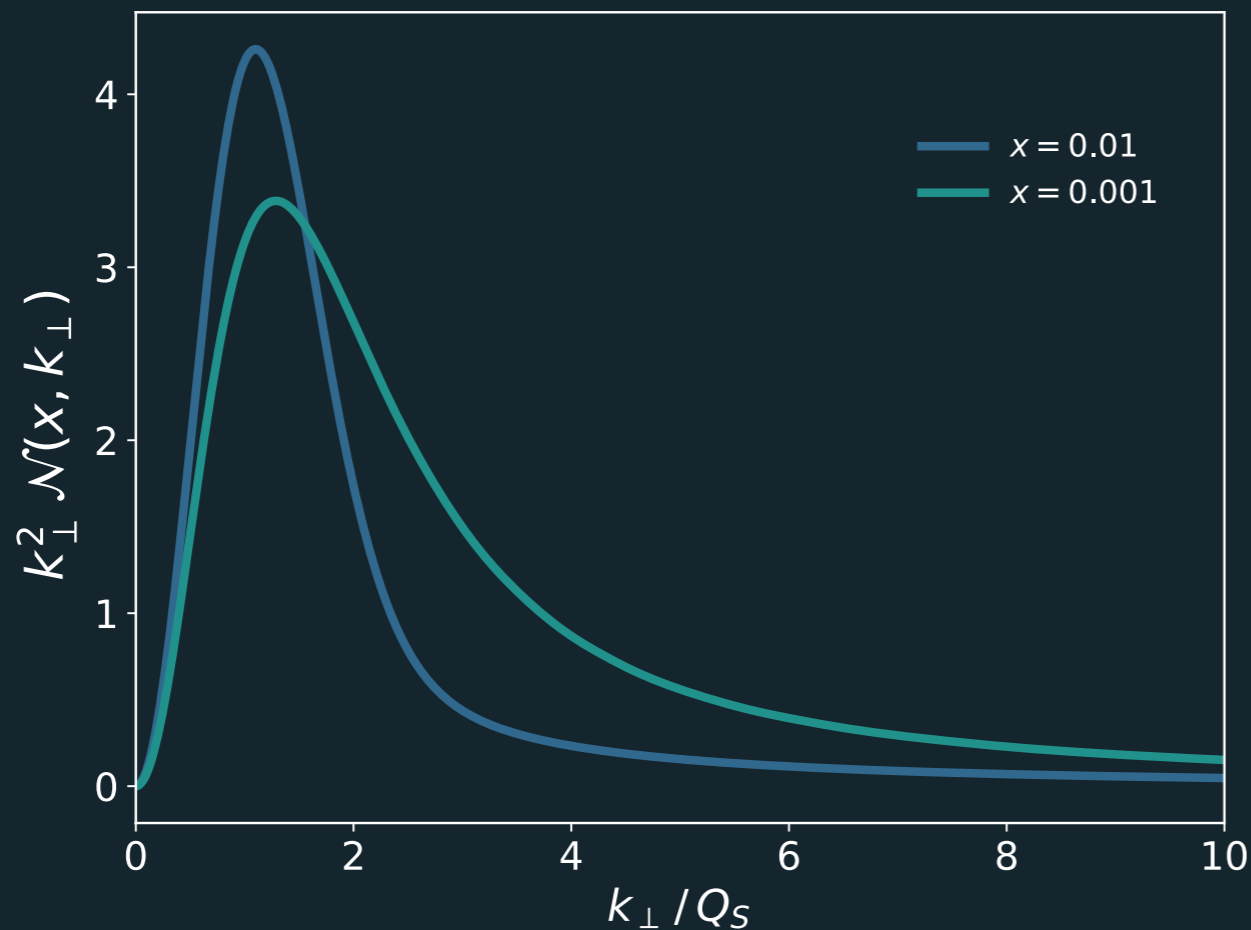
$$\frac{\partial \mathcal{N}(\mathbf{r}, x)}{\partial \log(x_0/x)} = \int d^2 \mathbf{r}_1 K^{run}(\mathbf{r}, \mathbf{r}_1, \mathbf{r}_2) [\mathcal{N}(\mathbf{r}_1, x) + \mathcal{N}(\mathbf{r}_2, x) - \mathcal{N}(\mathbf{r}, x) - \mathcal{N}(\mathbf{r}_1, x) \mathcal{N}(\mathbf{r}_2, x)]$$



KERNEL



NON LINEAR



Initial condition:
MV MODEL



EVOLUTION

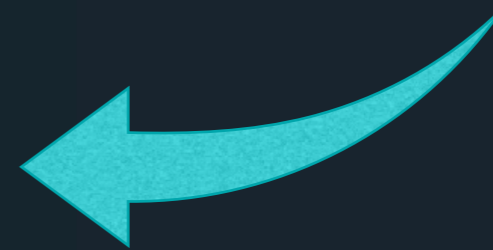
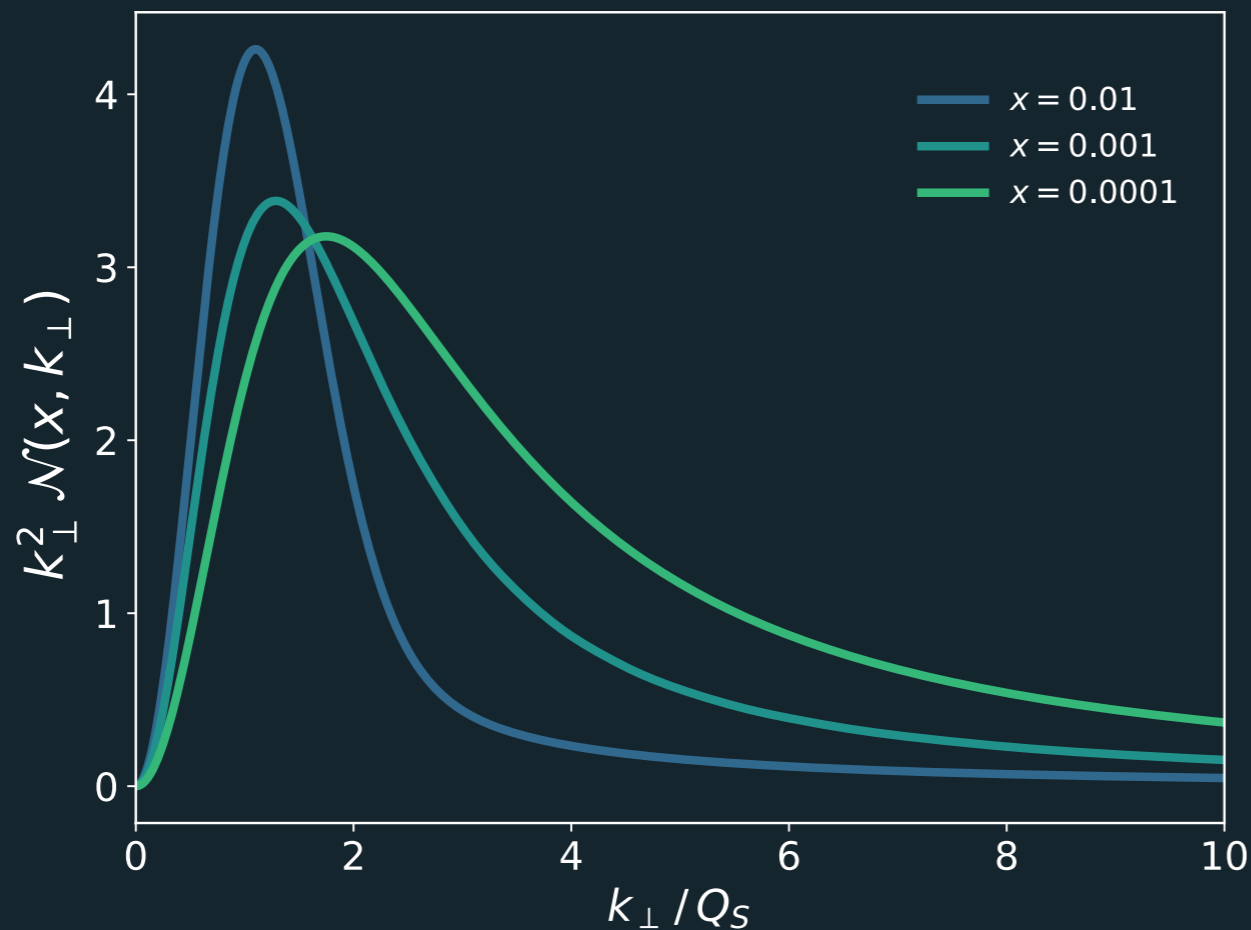
$$\frac{\partial \mathcal{N}(\mathbf{r}, x)}{\partial \log(x_0/x)} = \int d^2 \mathbf{r}_1 K^{run}(\mathbf{r}, \mathbf{r}_1, \mathbf{r}_2) [\mathcal{N}(\mathbf{r}_1, x) + \mathcal{N}(\mathbf{r}_2, x) - \mathcal{N}(\mathbf{r}, x) - \mathcal{N}(\mathbf{r}_1, x) \mathcal{N}(\mathbf{r}_2, x)]$$



KERNEL



NON LINEAR



EVOLUTION

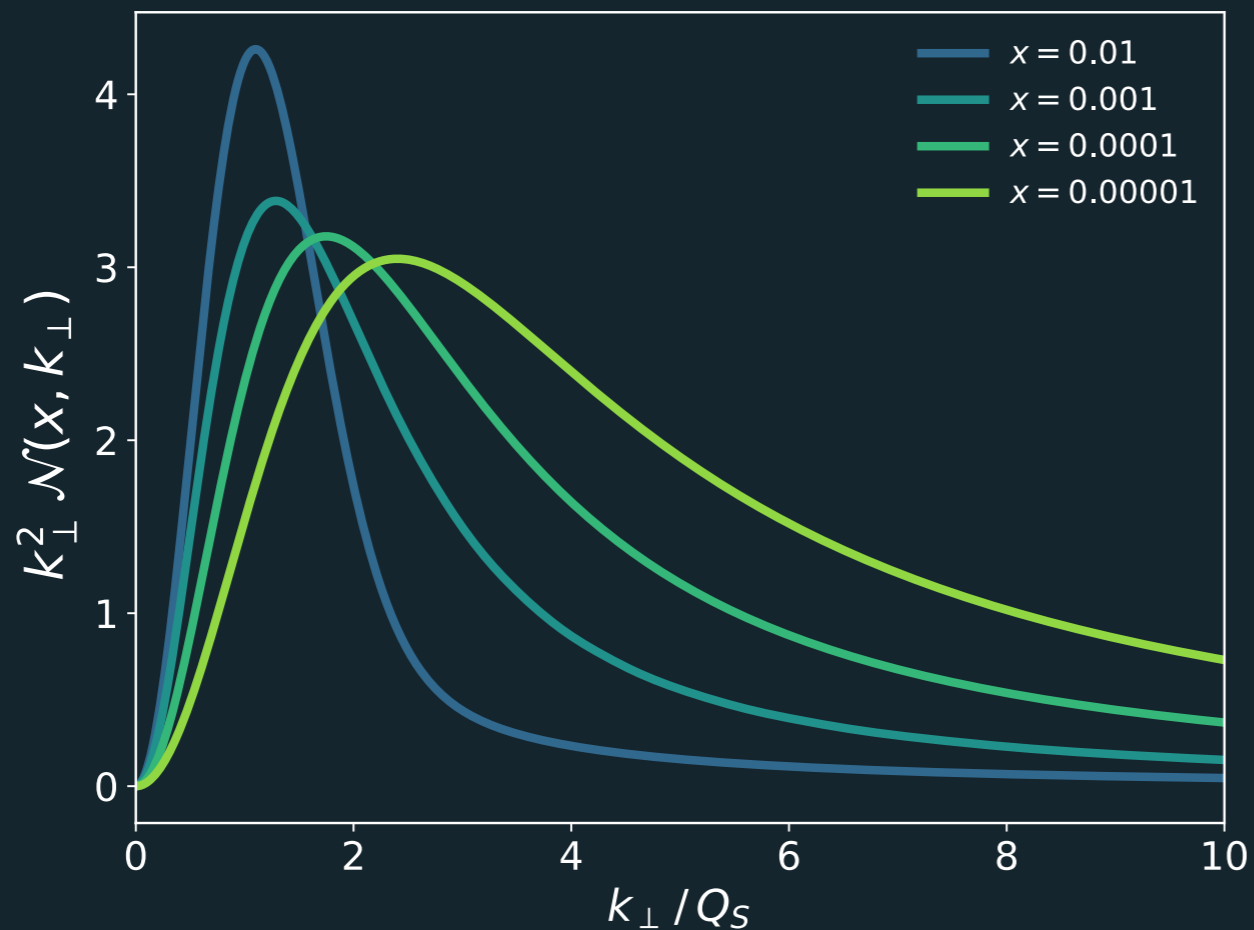
$$\frac{\partial \mathcal{N}(\mathbf{r}, x)}{\partial \log(x_0/x)} = \int d^2 \mathbf{r}_1 K^{run}(\mathbf{r}, \mathbf{r}_1, \mathbf{r}_2) [\mathcal{N}(\mathbf{r}_1, x) + \mathcal{N}(\mathbf{r}_2, x) - \mathcal{N}(\mathbf{r}, x) - \mathcal{N}(\mathbf{r}_1, x) \mathcal{N}(\mathbf{r}_2, x)]$$



KERNEL



NON LINEAR



Initial condition:
MV MODEL



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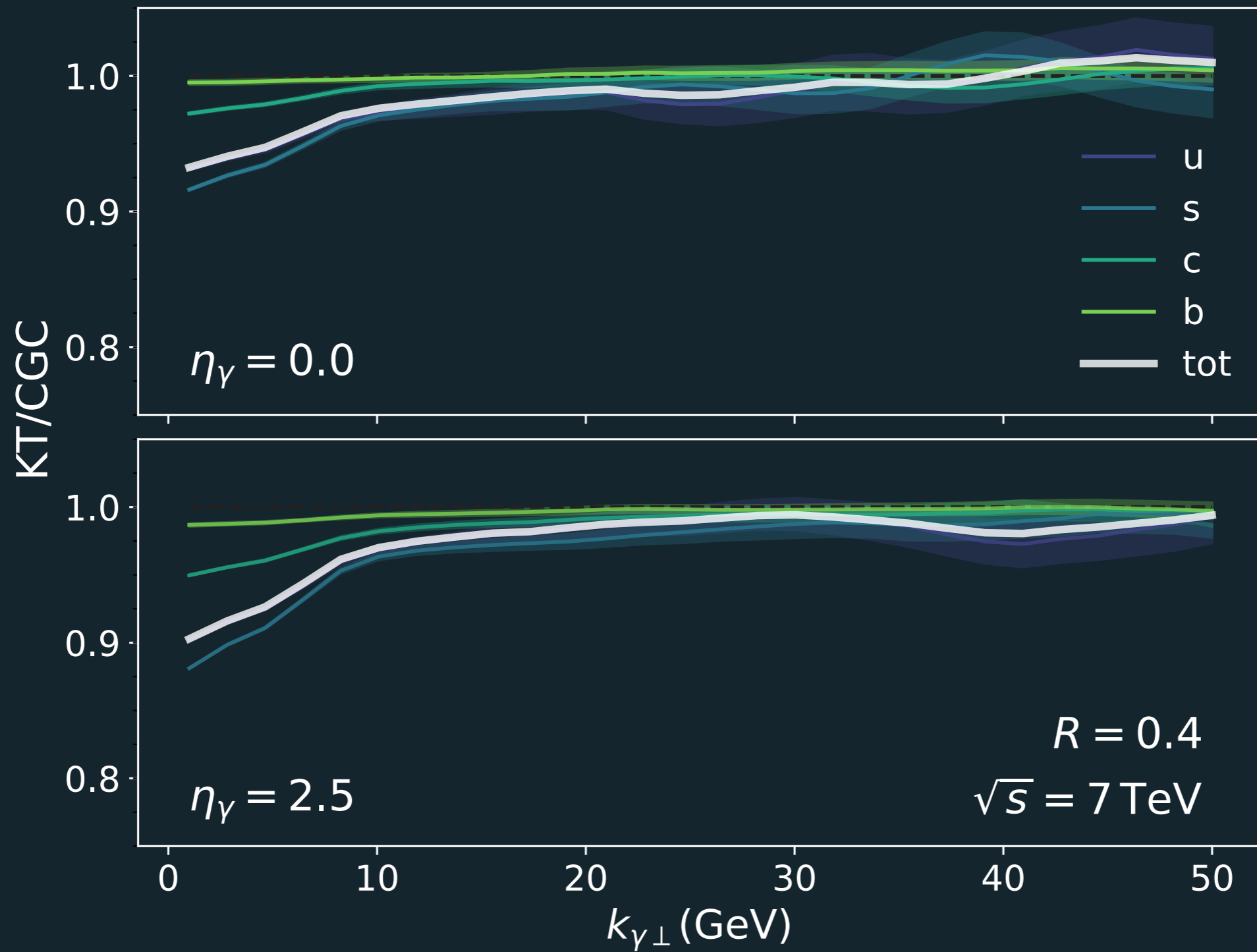
SOME RESULTS

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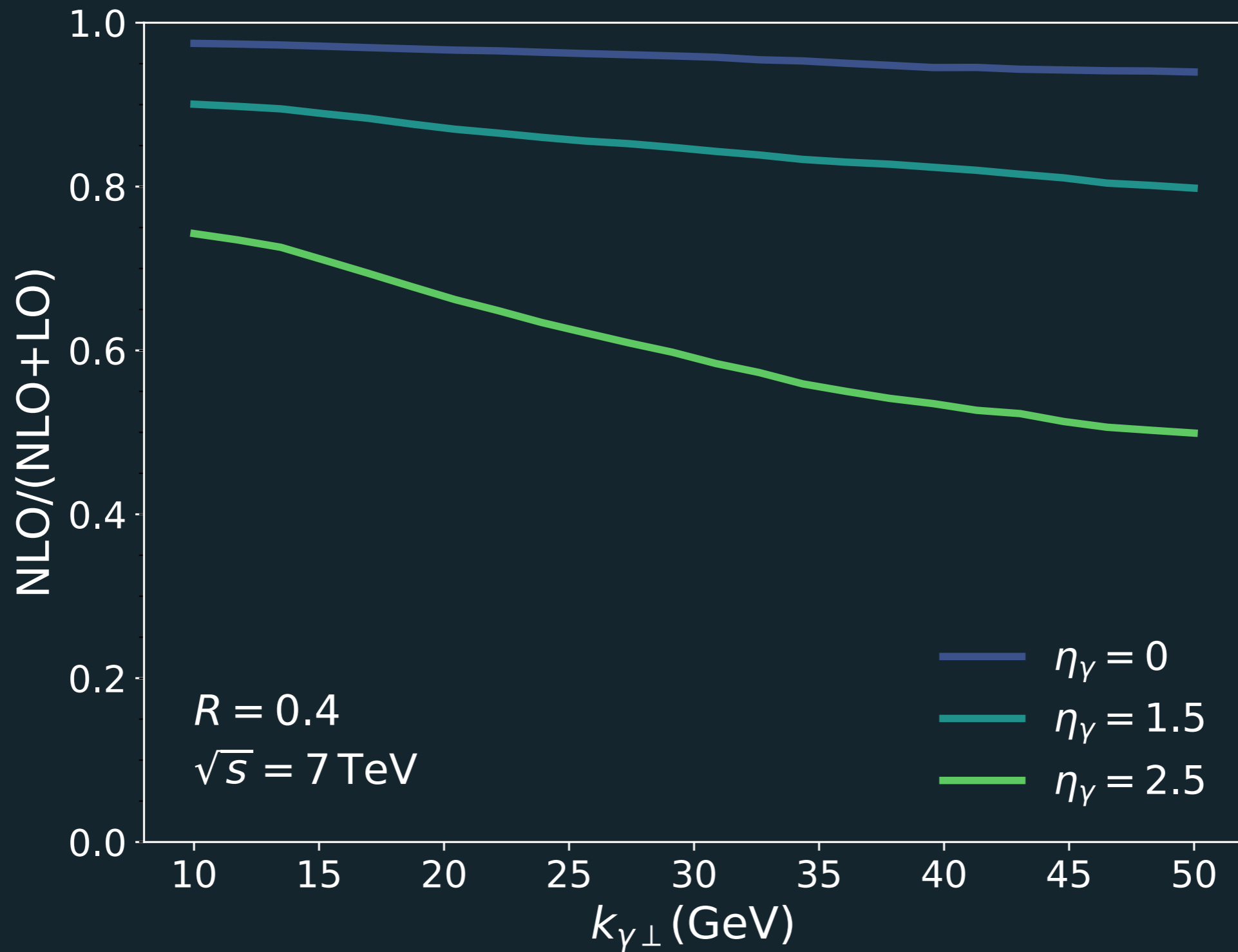
**SUMMARY AND
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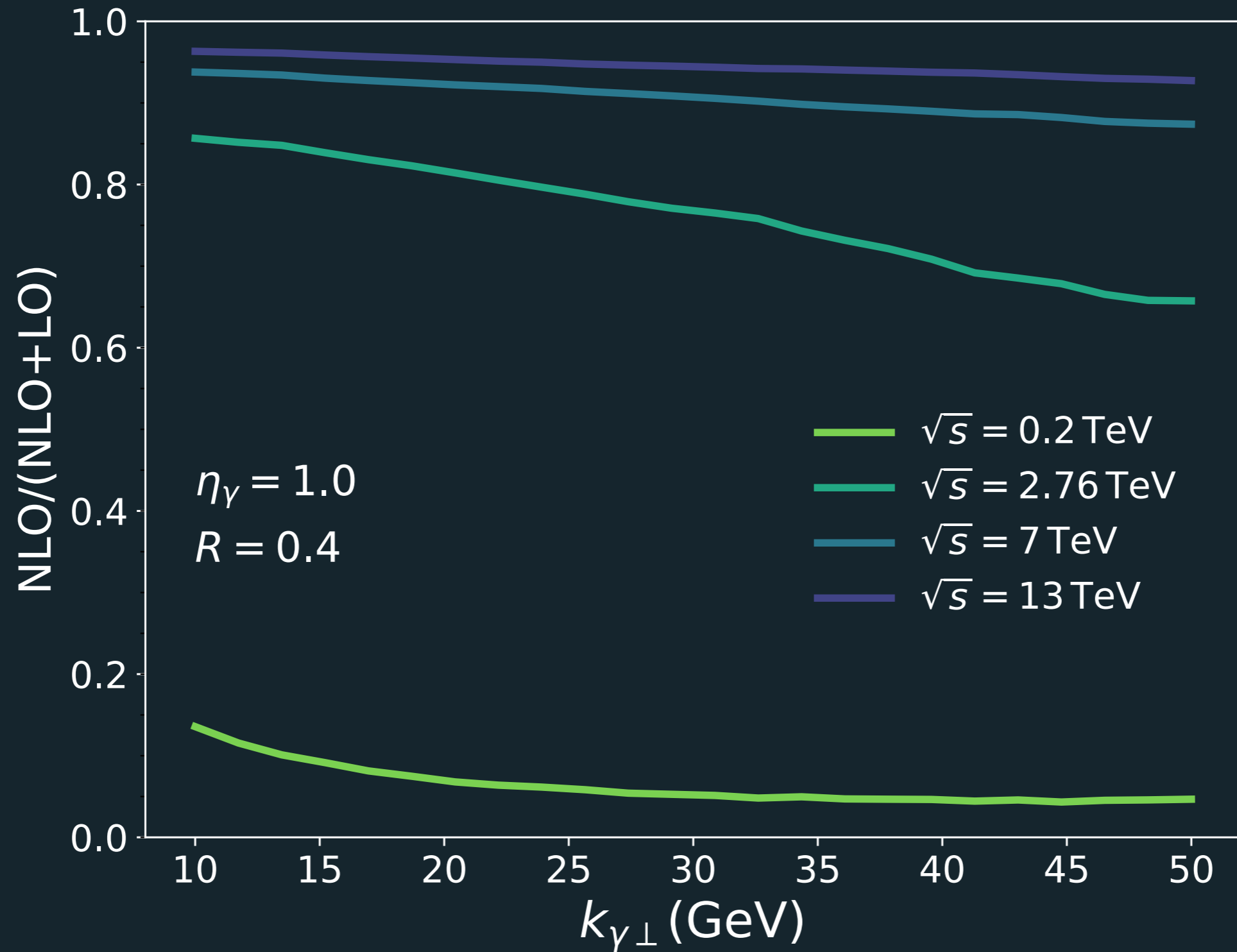
APPROXIMATION



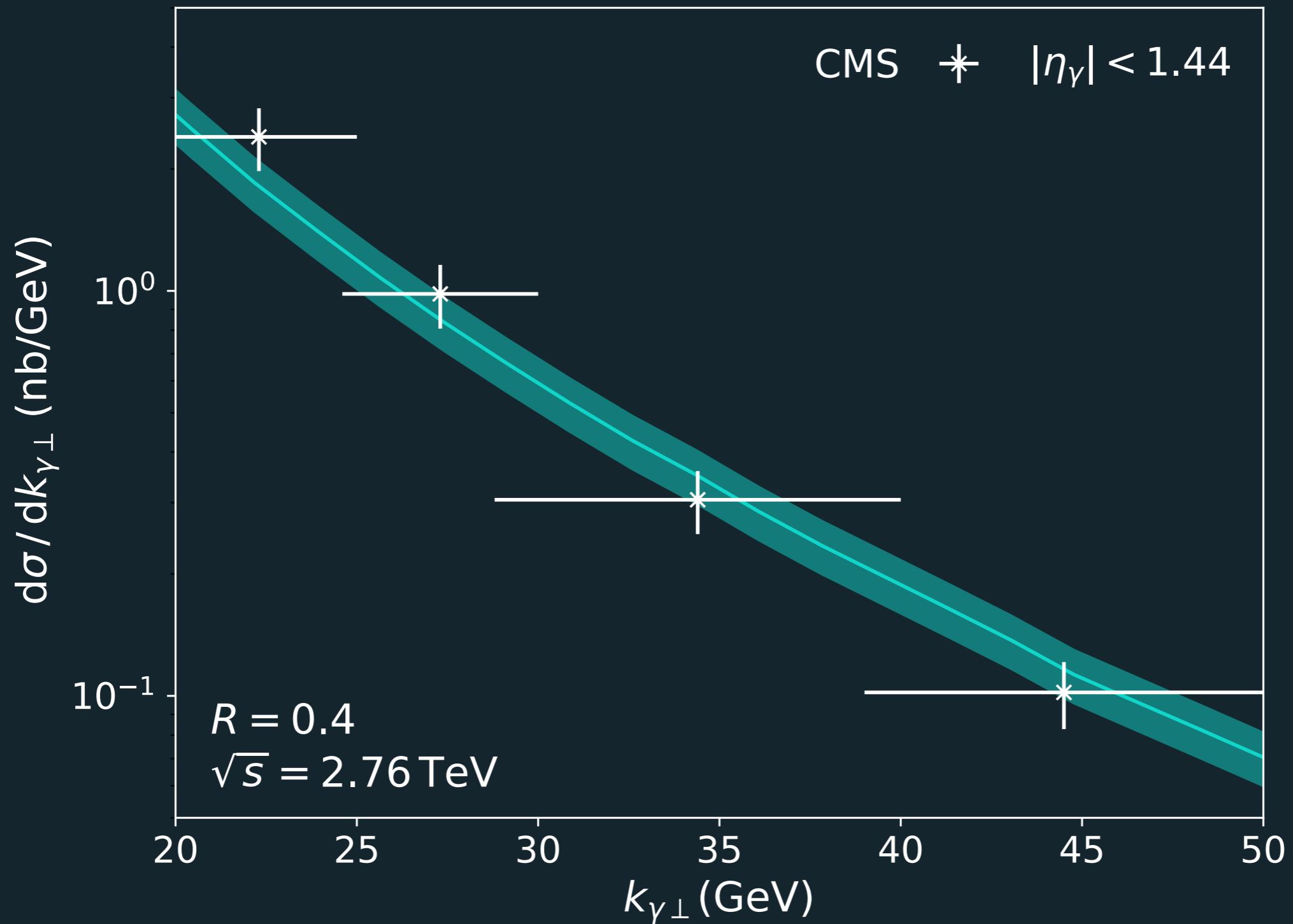
LO VS NLO: RAPIDITY



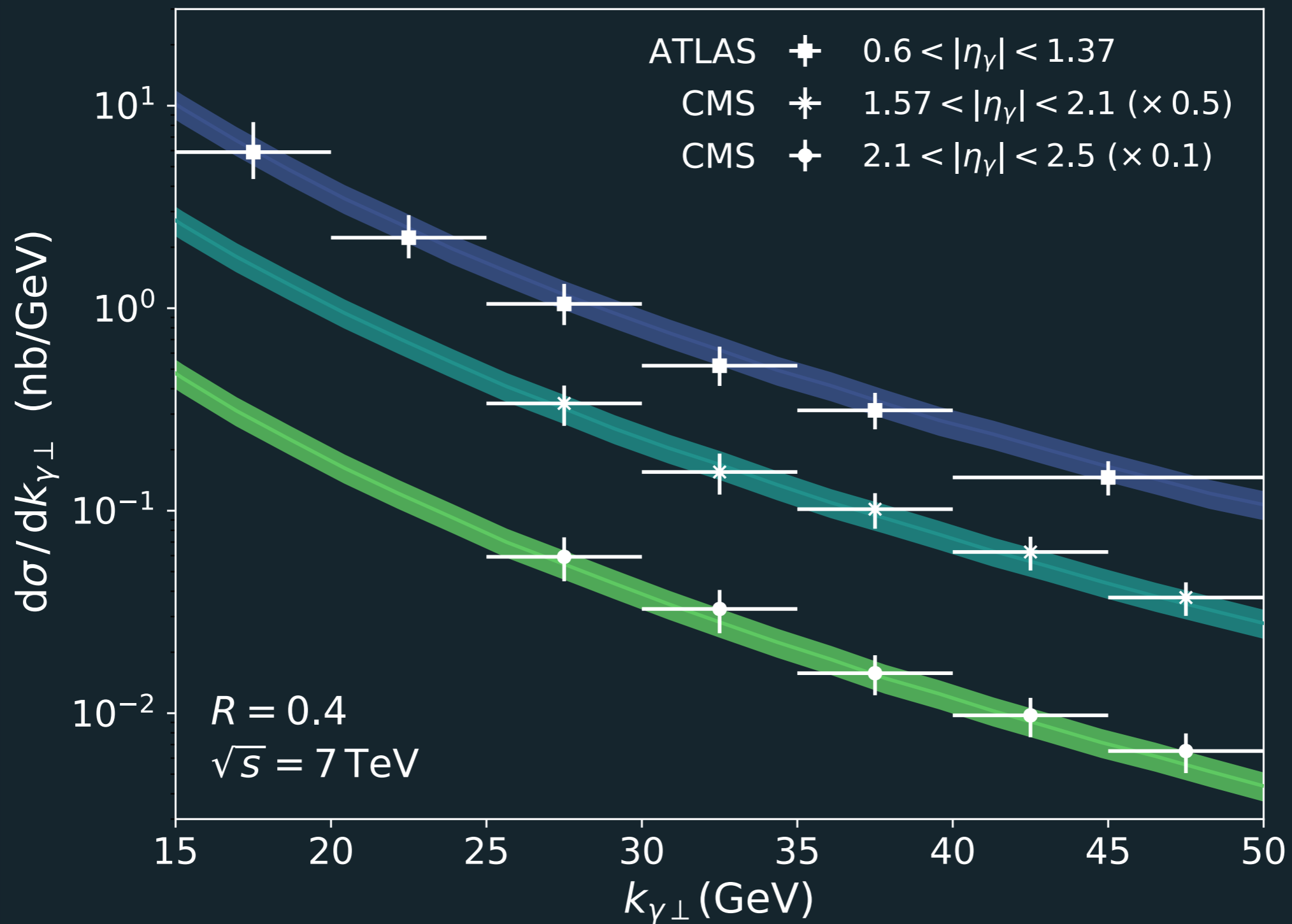
LO VS NLO: ENERGY



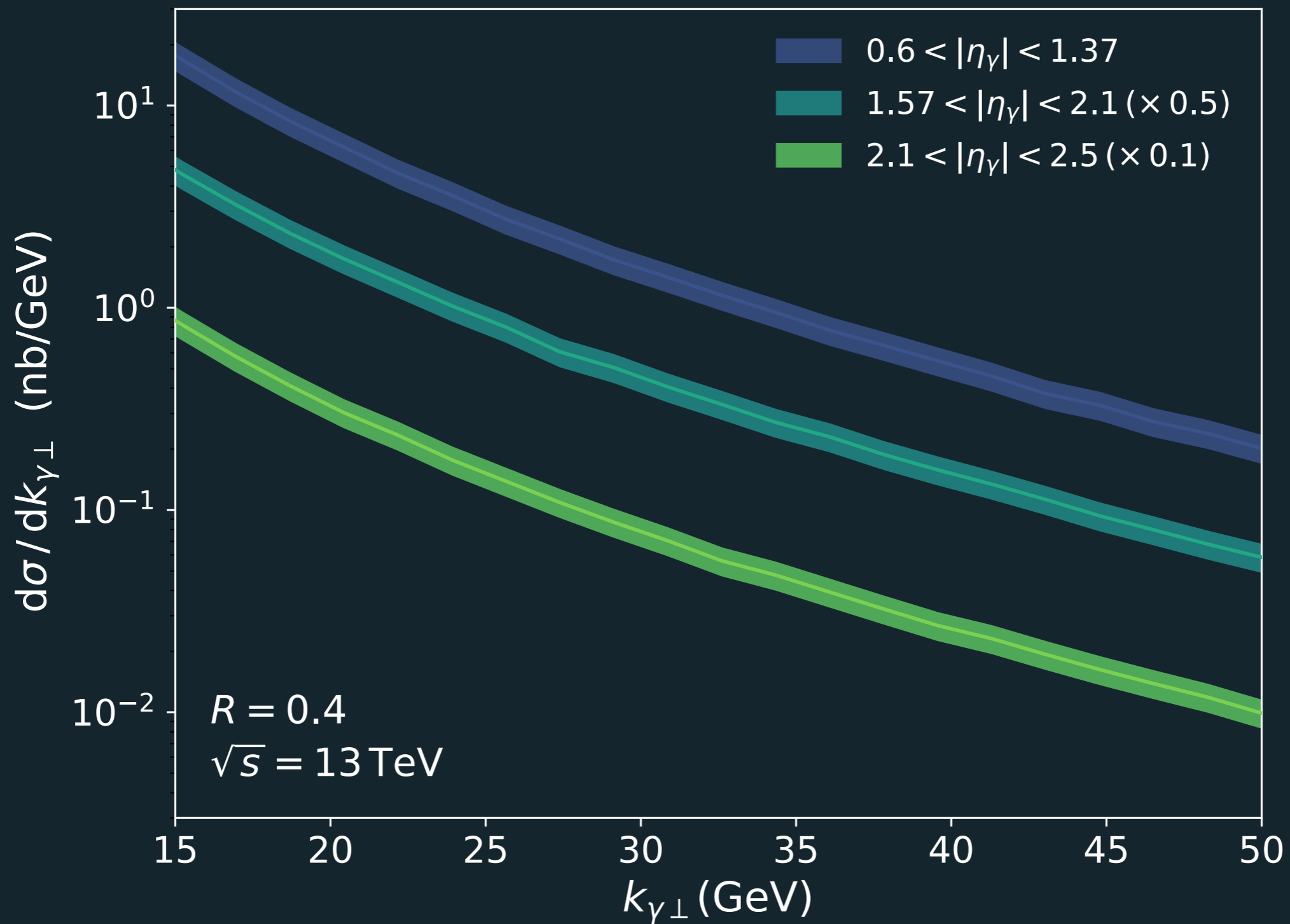
THEORY VS EXPERIMENT



THEORY VS EXPERIMENT



PREDICTION



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SUMMARY

- A** Saturation modifies the emission process thanks to multi-particle scatterings
- B** Complete analytical result at NLO, that is $\mathcal{O}(\alpha_s\alpha_e)$
- C** CGC formalism yields correct limits to the pQCD results
- D** CGC k_{\perp} -factorized yields good agreement with the experimental data

OUTLOOK

A

Comparison to current and future p+A experimental data

B

Photon hadron correlations are sensitive to saturation

c

More exciting studies

- ◆ NNLO?
- ◆ Higher particle correlations?
- ◆ Saturation and Anomalies?



**THANK
YOU**

✉ GARCIA@THPHYS.UNI-HEIDELBERG.DE