

Insight into thermal modifications of quarkonia from a comparison of continuum-extrapolated lattice results to perturbative QCD

Anna-Lena Kruse

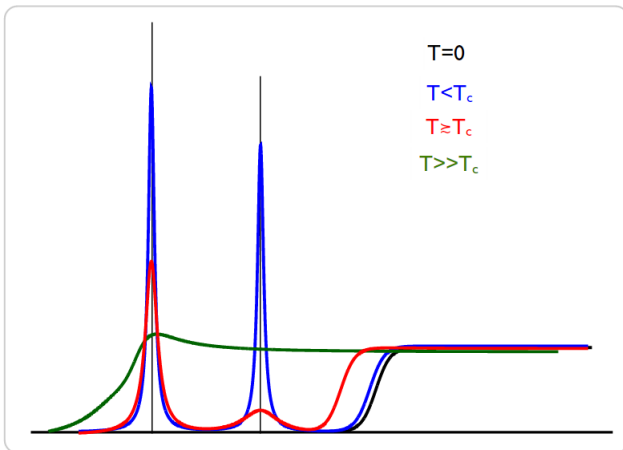
with H.-T. Ding, O. Kaczmarek, M. Laine, H. Ohno & H. Sandmeyer

[arXiv:1710.08858]

Universität Bielefeld



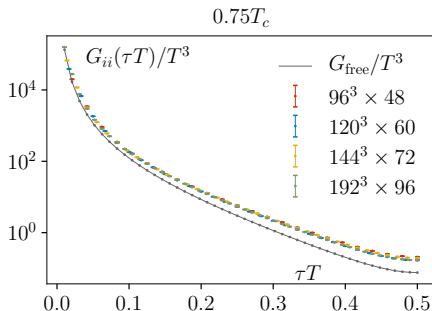
- Information about in-medium properties of quarkonia encoded in spectral function



- Lattice: Correlator
- Pseudoscalar $G_P(\tau) = M_B^2 \int_{\vec{x}} \langle (\bar{\psi} i \gamma_5 \psi)(\tau, \vec{x}) (\bar{\psi} i \gamma_5 \psi)(0, \vec{0}) \rangle_c$
- Vector: $G_{ii}(\tau) = \langle \psi^\dagger(\tau, \vec{x}) \gamma_i \psi(\tau, \vec{x}) \psi^\dagger(0, \vec{0}) \gamma_i \psi(0, \vec{0}) \rangle$
- Related to ρ : $G(\tau) = \int_0^\infty d\omega \rho(\omega) K(\omega, \tau)$

with $K(\omega, \tau) = \frac{\cosh(\omega(\tau - \frac{1}{2T}))}{\sinh(\frac{\omega}{2T})}$

⇒ ill-posed inversion problem



- Pseudoscalar: No transport peak \Rightarrow easier
- Vector:
 - Information on heavy quark diffusion coefficients

$$D = \frac{\pi}{3\chi_q} \lim_{\omega \rightarrow \infty} \sum_{i=3}^3 \frac{\rho_{ii}(\omega, T)}{\omega}$$

- Information on quark number susceptibility

$$G^{00} = \chi_q T$$

Lattice Correlators

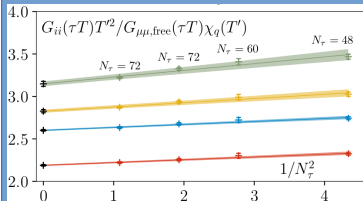
Renormalization

Mass Interpolation



Normalization with G_{free}
and B-Spline Interpolation

Continuum Extrapolation



Perturbatively determined renormalization constants in pseudoscalar channel

Different Options in vector channel

- Perturbative renormalization constants (1 and 2loop available)
- Non-perturbative renormalization constants [Lüscher et al., Phys.Lett. B372(1996)275-282]

$$Z_V^m(g^2, \kappa) = \frac{1 - 0.7663g^2 + 0.0488g^4}{1 - 0.6369g^2} (1 + b_V(g^2)am_q) \text{ with}$$

$$b_V(g^2) = \frac{1 - 0.6518g^2 - 0.1226g^4}{1 - 0.8467g^2}$$

- Renormalization independent ratio with quark number susceptibility

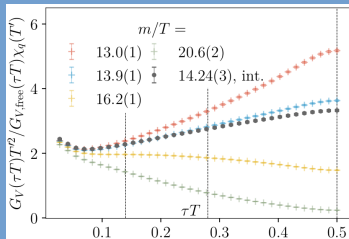
χ_q

$\chi_q T$ given by G^{00} , Choice: χ_q at $T' = 2.25T_c$

Lattice Correlators

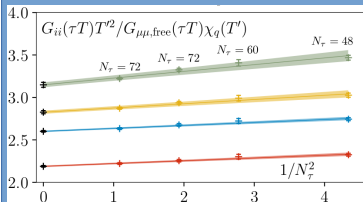
Renormalization

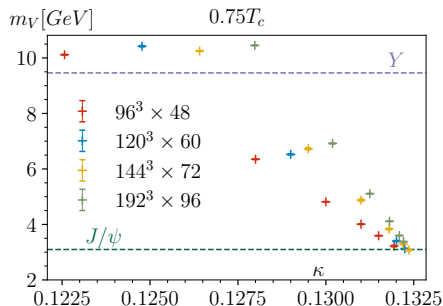
Mass Interpolation



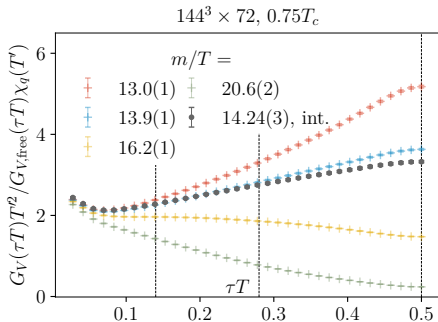
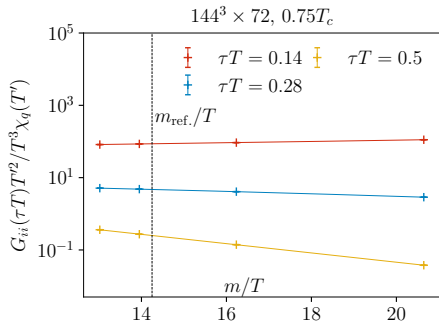
Normalization with G_{free}
and B-Spline Interpolation

Continuum Extrapolation

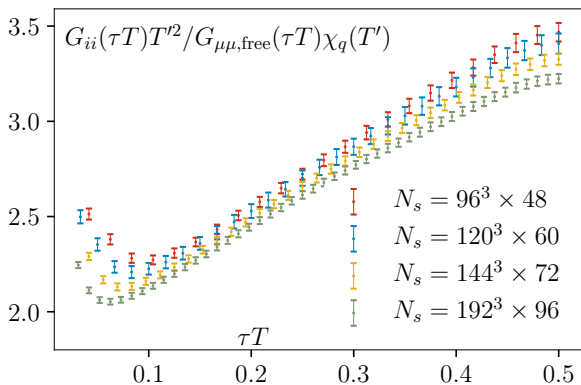




- Problem: Meson mass not equal on different lattices
- Solution: Interpolation of correlators between different masses to a fixed m/T for $T = 0.75T_c$
- Ansatz:
$$\frac{G^{ii}(\tau T, m/T) T'^2}{T^3 \chi_q} = \exp(p(m/T)^2 + qm/T + r)$$



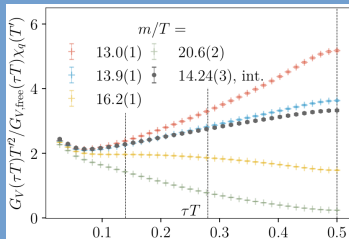
$0.75T_c$



Lattice Correlators

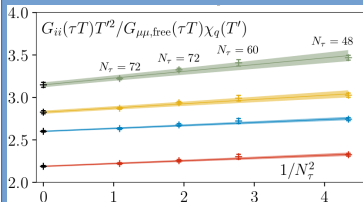
Renormalization

Mass Interpolation

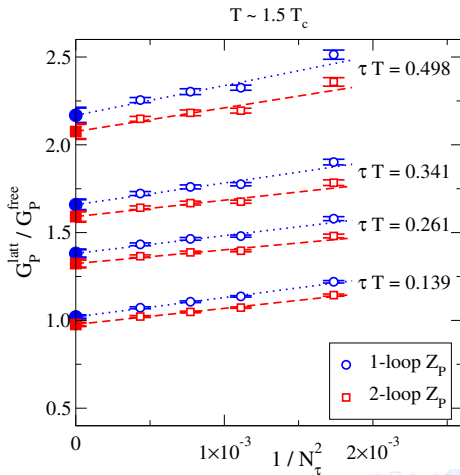


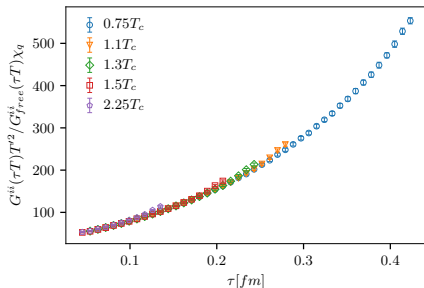
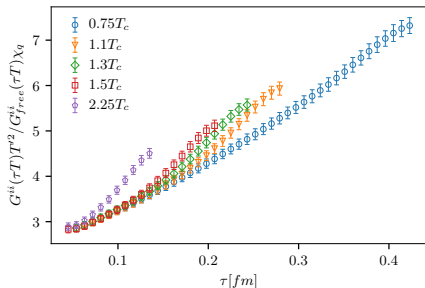
Normalization with G_{free}
and B-Spline Interpolation

Continuum Extrapolation



- Ansatz: $\frac{G^{ii}(\tau T, a) T'^2}{G_{free} T^3 \chi_q} = G^{cont}(\tau T) + \frac{c}{N_\tau^2}$
 or $\frac{G(\tau T, a)}{G_{free}} = G^{cont}(\tau T) + \frac{c}{N_\tau^2}$





- Ultraviolet asymptotics available up to 5-loop level
[Y. Burnier, M. Laine, Eur.Phys.J.C 72 (2012) 1902]
- Threshold region using a finite-temperature real-time static potential
[M.Laine, JHEP 0705:028,2007]
- Matching between regions
[Y. Burnier, H.-T. Ding, O. Kaczmarek, A.-L. Kruse, M. Laine, H. Ohno, H. Sandmeyer, JHEP 1711 (2017) 206]

Comparison of Lattice and Perturbation Theory

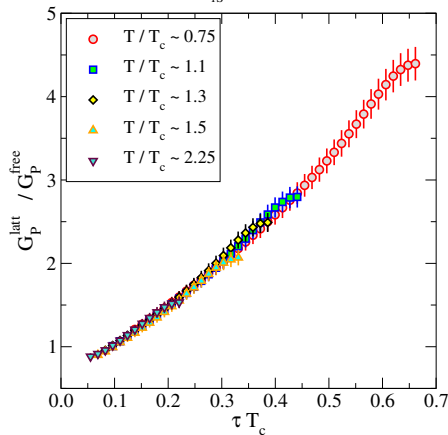
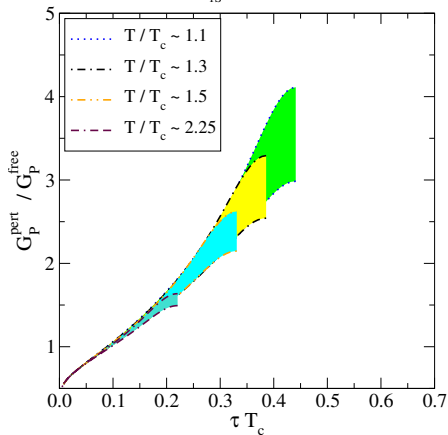
- A first look

$$G(\tau) = \int_0^{\infty} d\omega K(\omega, \tau) \rho(\omega)$$

Charmonium - Pseudoscalar

$M_{1S} \sim 1.5 \text{ GeV}$

$M_{1S} \sim 1.5 \text{ GeV}$



Comparison of Lattice and Perturbation Theory

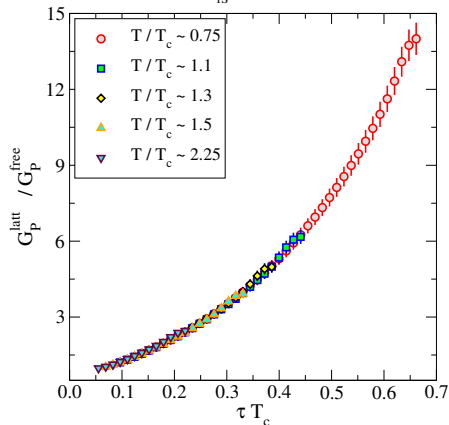
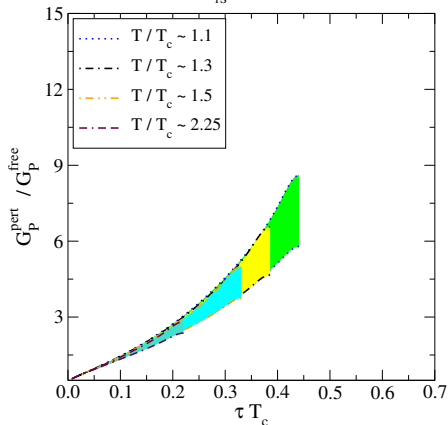
- A first look

$$G(\tau) = \int_0^{\infty} d\omega K(\omega, \tau) \rho(\omega)$$

Bottomonium - Pseudoscalar

$M_{1S} \sim 4.7 \text{ GeV}$

$M_{1S} \sim 4.7 \text{ GeV}$

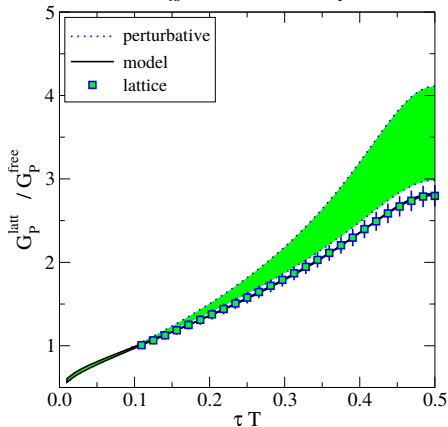


Qualitatively similar results

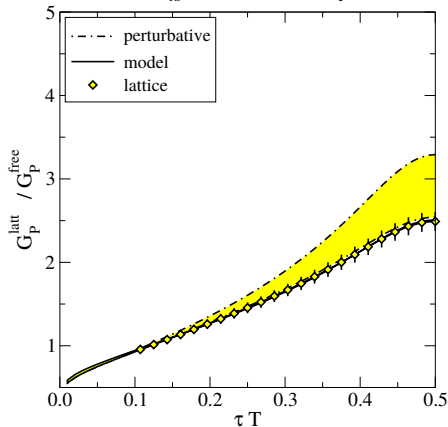
- On Lattice side: Uncertainties in renormalization constants
 - ⇒ Relative normalization might be off
 - ⇒ Introduce overall normalization factor A
- On perturbative side: Relation between M and $m(\bar{\mu}_{ref})$ poorly known
 - ⇒ Threshold location might be off
 - ⇒ Introduce frequency shift B
- Ansatz: $\rho^{model}(\omega) = A\rho^{pert}(\omega - B)$

$$\rho^{model}(\omega) = A\rho^{pert}(\omega - B), \quad G(\tau) = \int_0^\infty d\omega K(\omega, \tau)\rho(\omega)$$

$M_{IS} \sim 1.5 \text{ GeV}, T \sim 1.1 T_c$

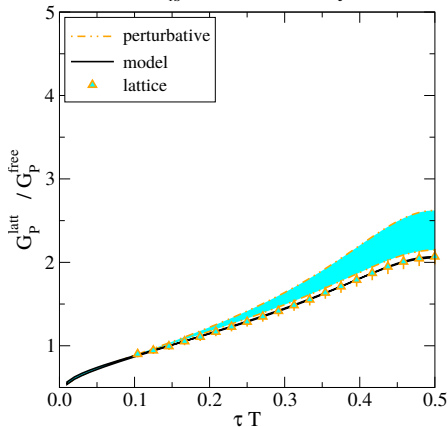


$M_{IS} \sim 1.5 \text{ GeV}, T \sim 1.3 T_c$

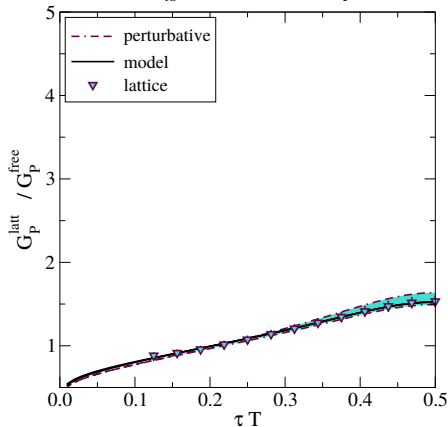


$$\rho^{model}(\omega) = A\rho^{pert}(\omega - B), \quad G(\tau) = \int_0^\infty d\omega K(\omega, \tau)\rho(\omega)$$

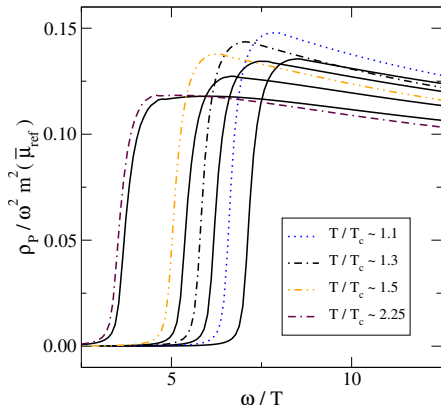
$M_{IS} \sim 1.5 \text{ GeV}, T \sim 1.5 T_c$



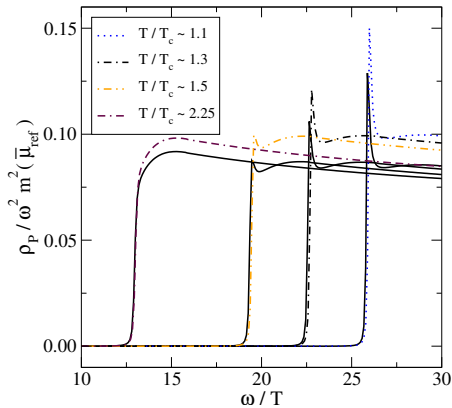
$M_{IS} \sim 1.5 \text{ GeV}, T \sim 2.25 T_c$



$m(\bar{\mu}_{\text{ref}}) = 1 \text{ GeV}$



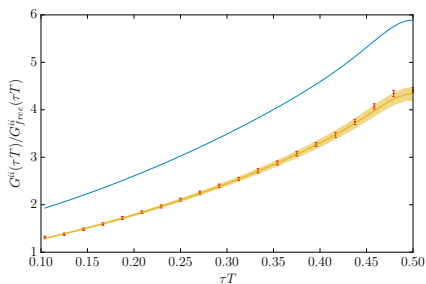
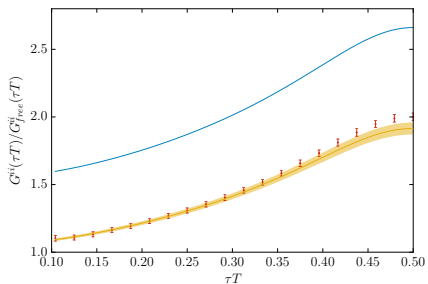
$m(\bar{\mu}_{\text{ref}}) = 5 \text{ GeV}$



dashed: perturbative spectral function
 solid: fitted model spectral function

- Quite good agreement:
 - A close to 1
 - B small
- Now moving on to the vector channel: Shape of the transport peak not known!
But: Integral can be calculated
 \Rightarrow Constant contribution to correlator $\frac{G_{ii}^{const}}{T^3}$ is known.

$1.5T_c$



Fit to small τT compared to lattice data. Difference \Rightarrow Transport peak!

- **Good agreement** in pseudoscalar channel
- **Differences** explained by systematic uncertainties
- **Transport contribution** visible in vector channel

- **Good agreement** in pseudoscalar channel
- **Differences** explained by systematic uncertainties
- **Transport contribution** visible in vector channel

Future:

- Use suitable ansätze to fit the transport peak
- Crosscheck using Bayesian reconstruction methods
[arXiv:1807.06315, arXiv:1712.03341]
- Full QCD study required in future