

Quarkonia production and dissociation in a Langevin approach

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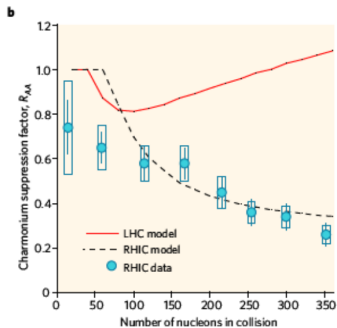
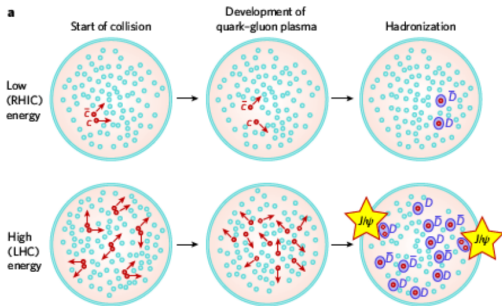
Hot Quarks 18, Texel September 07-14, 2018

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Motivation

- develop a classical model for J/ψ dissociation and recombination
- quarkonia are an important tool for the investigation of the QGP
- signal for deconfinement [T. Matsui and H. Satz 1986 *Phys. Lett. B* **178** 416]
- dissociation due to collisions and color screening
- recombination at higher beam energies

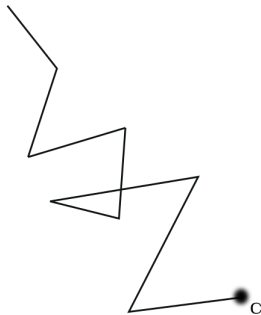


[from J. Stachel, Talk at EMMI workshop Feb/13/18]

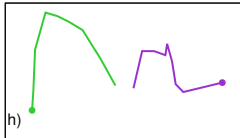
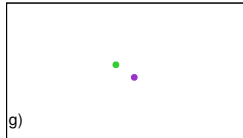
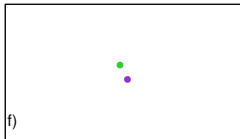
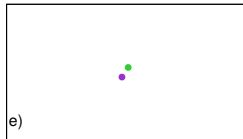
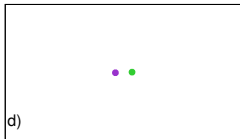
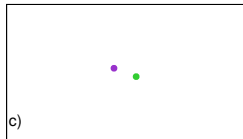
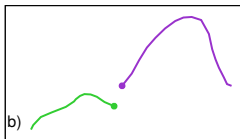
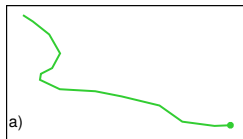
Langevin equation

Langevin equation

- Langevin equation describes Brownian motion (heavy particles in a light medium)
- particle performs a 'random walk'
- $\frac{d\mathbf{p}}{dt} = -\gamma\mathbf{p} + F(\mathbf{r}) + \zeta$
- $\frac{d\mathbf{x}}{dt} = \frac{\mathbf{p}}{m}$



Langevin Equation



(a) trajectory of single c quark

(b) trajectory of a $c\bar{c}$ pair

(c-g) time sequence of a bound $c\bar{c}$ pair

(h) trajectory of a dissociating bound state

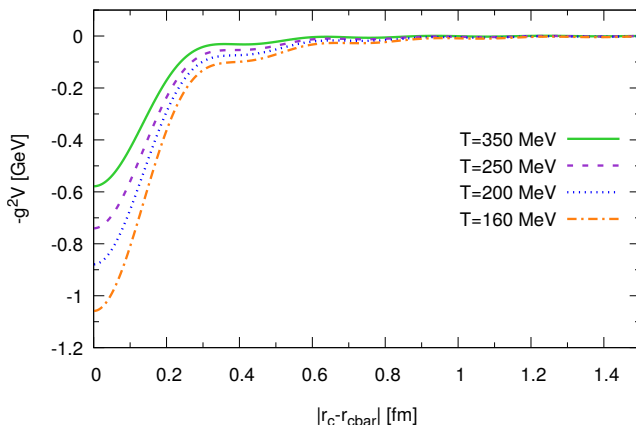
Formalism

$c\bar{c}$ -potential with cut-off $\Lambda = 4$ GeV

- drag coefficient γ and potential for $F(\mathbf{r}) = -\nabla V(r)$

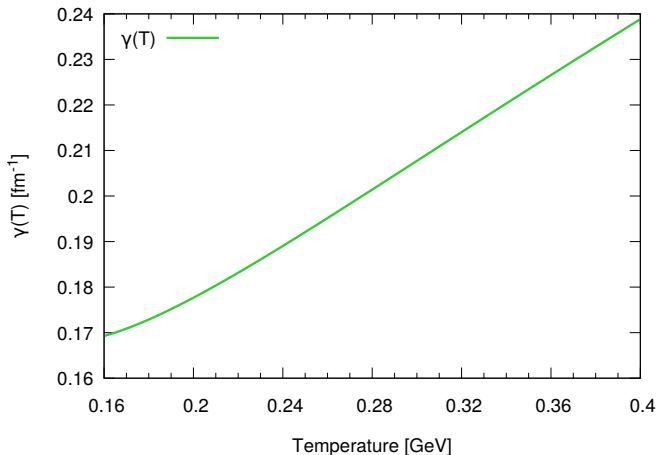
from [J. P. Blaizot, D. De Boni, P. Faccioli, and G. Garberoglio 2016 *Nucl. Phys. A* 946 49]

- $g^2 = 4\pi\alpha_s = \frac{4\pi\alpha_s(T_C)}{1 + \text{Cln}\left(\frac{T}{T_C}\right)}$ $C = 0.76$, $T_C = 160\text{MeV}$, $\alpha_s(T_C) = 0.5$



Friction coefficient

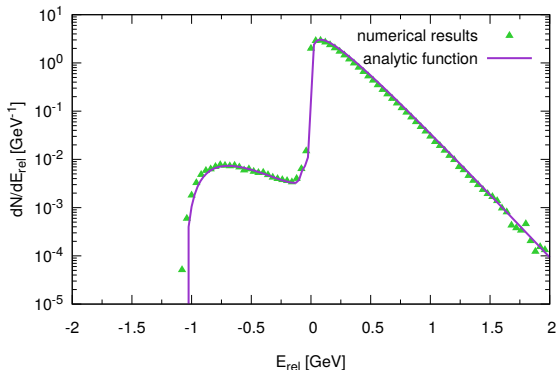
$$\bullet \gamma = \frac{m_D^2 g^2}{24\pi m_Q} \left(\ln \left(1 + \frac{\Lambda^2}{m_D^2} \right) - \frac{\frac{\Lambda^2}{m_D^2}}{\frac{\Lambda^2}{m_D^2} + 1} \right)^{-1}$$



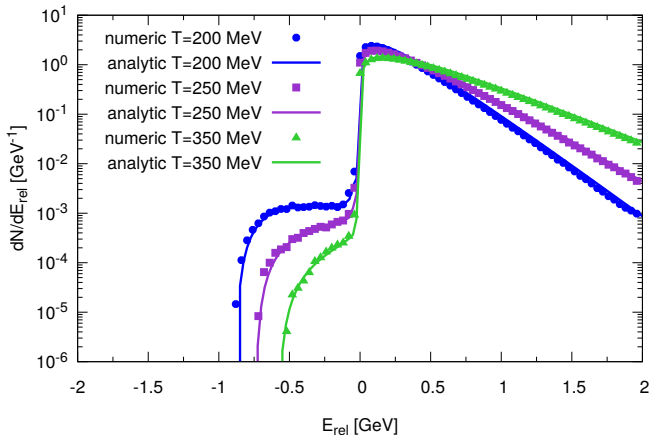
Numerical calculations

Energy distribution $T=160$ MeV

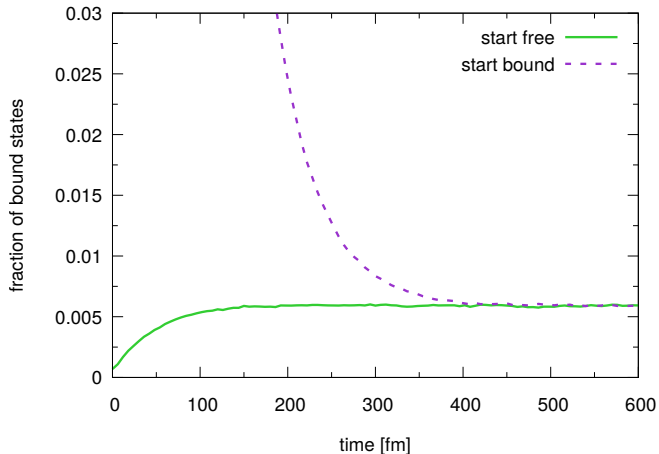
$$\begin{aligned}\frac{dN}{dE_{\text{rel}}} &= C \int_{\mathbb{R}^3} d^3\mathbf{r} \int_{\mathbb{R}^3} d^3\mathbf{p}_{\text{rel}} \delta(E_{\text{rel}} - H_{\text{rel}}) \exp\left(-\frac{H_{\text{rel}}}{T}\right) \\ &= (4\pi)^2 (2\mu)^{3/2} C \int_0^R dr r^2 \sqrt{E_{\text{rel}} - V(r)} \exp\left(-\frac{E_{\text{rel}}}{T}\right)\end{aligned}$$



Energy distribution at different temperatures

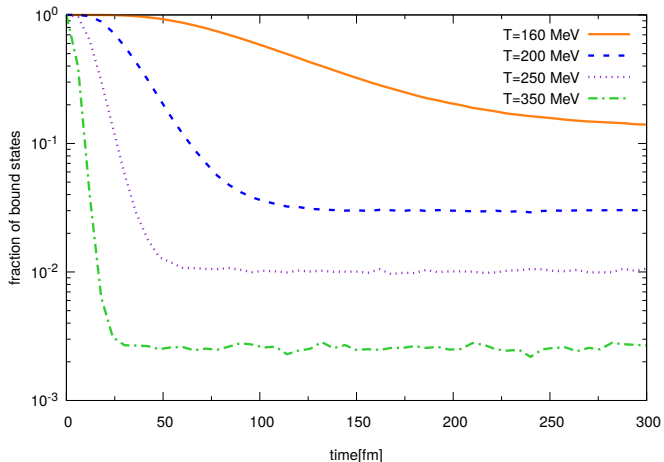


Comparison of initial conditions



- $T = 160$ MeV, single pair, $m_c = 1.8$ GeV, rigid box

Fraction of bound states at different temperatures



- $T = 160$ MeV, five $c\bar{c}$ pairs, $m_c = 1.8$ GeV, rigid box

Comparison with an equilibrated thermodynamic system

- number of J/ψ in the system can be calculated with the grand canonical partition function
- $T = 160$ MeV, single pair, $m_c = 1.8$ GeV, rigid box

Volume	grand canonical fraction of J/ψ	numerical fraction of J/ψ
8^3 fm^3	0.0066	0.0059
10^3 fm^3	0.0035	0.0029
12^3 fm^3	0.002	0.0017

- Summary

- model for J/ψ dissociation and recombination
- motion of heavy quarks realized via Langevin equation
- model passes all equilibrium box tests
- great scope for further research

- Outlook

- expanding fireball
- include distance-dependence in γ
- use different potential models (confinement)
- include quantum effects for example using Wigner function
[C., Young and E. Shuryak 2009 *Phys. Rev. C* **79** 034907]
- long-time goal: full in-medium quantum Langevin treatment