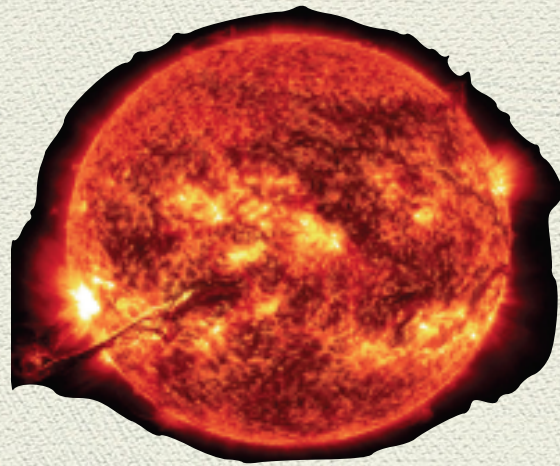


# Neutrino astrophysics

Maria Cristina VOLPE

AstroParticule et Cosmologie (APC), Paris

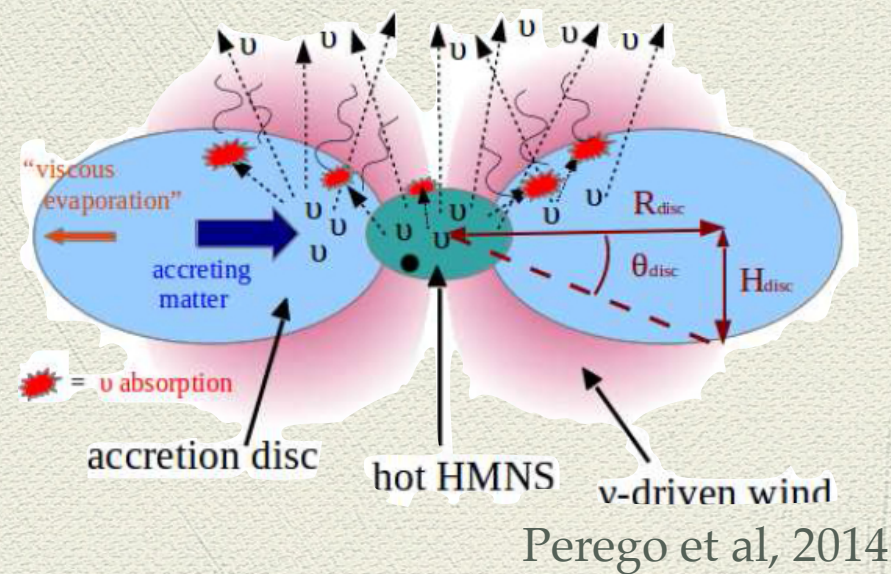
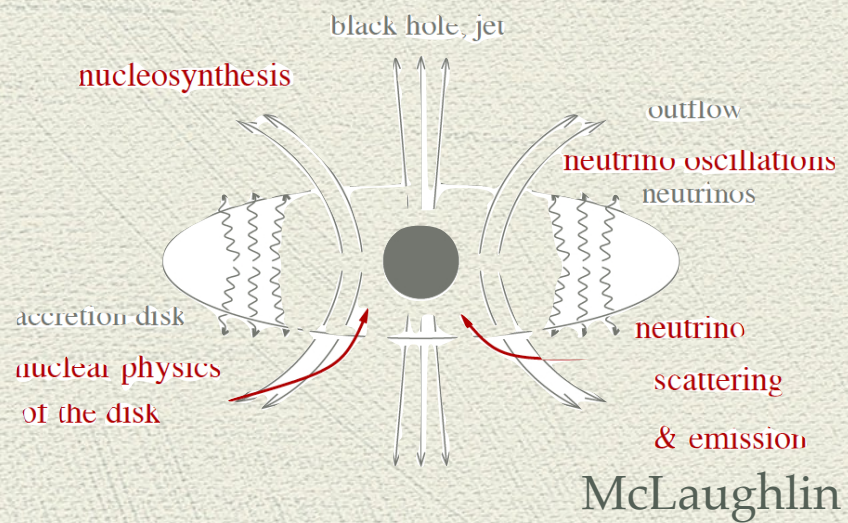




Sun



core-collapse Supernovae



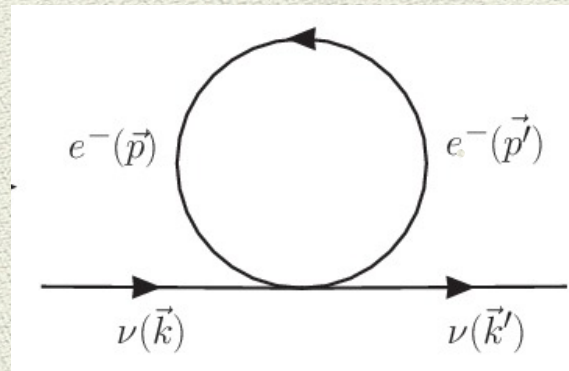
accretion disks around black holes or neutron star mergers remnants



# Solar neutrinos

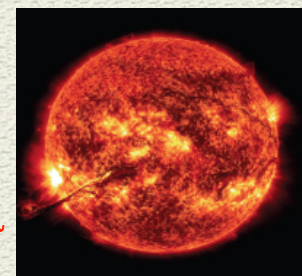
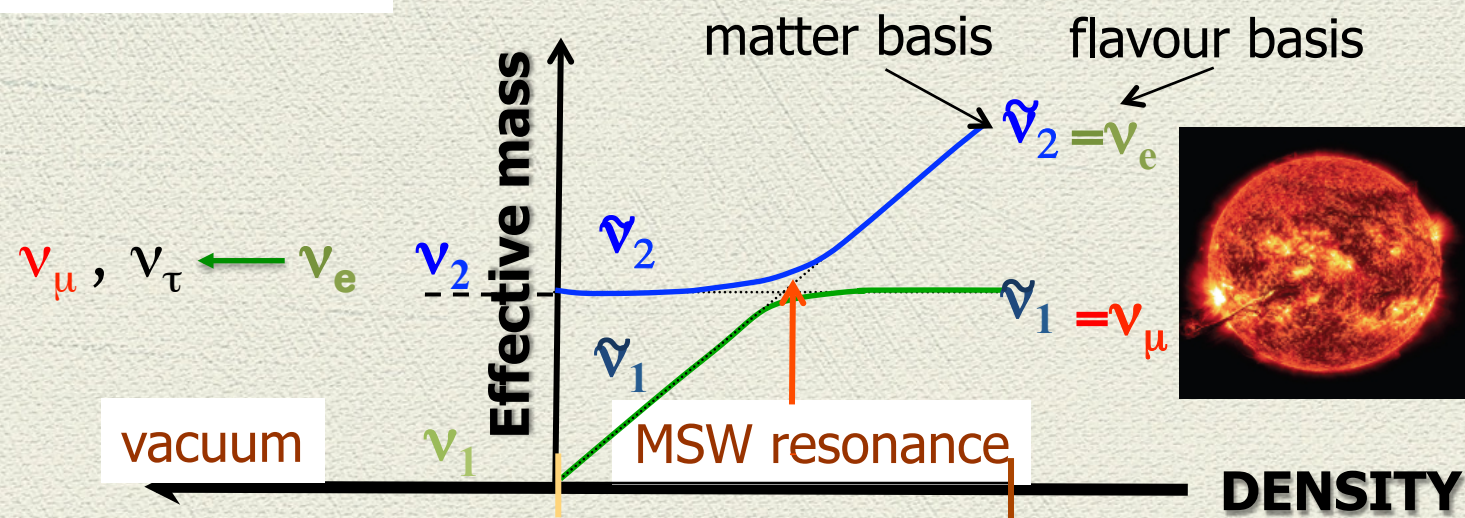
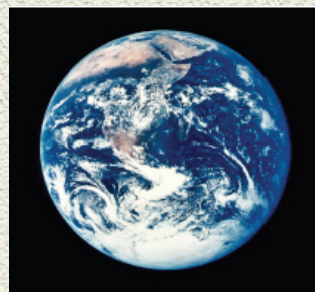
Neutrinos interact with matter and undergo **resonant adiabatic flavor conversion** : the Mikheev-Smirnov-Wolfenstein effect.

Wolfenstein PRD (1978)  
Mikheev, Smirnov(1985)



$$h_{mat} = \sqrt{2}G_F\rho_e$$

mean-field approximation



$$h_\nu = \begin{pmatrix} -\Delta\tilde{m}^2/4E & -i\dot{\theta}_M \\ i\dot{\theta}_M & \Delta\tilde{m}^2/4E \end{pmatrix}$$

neutrino hamiltonian in the matter basis

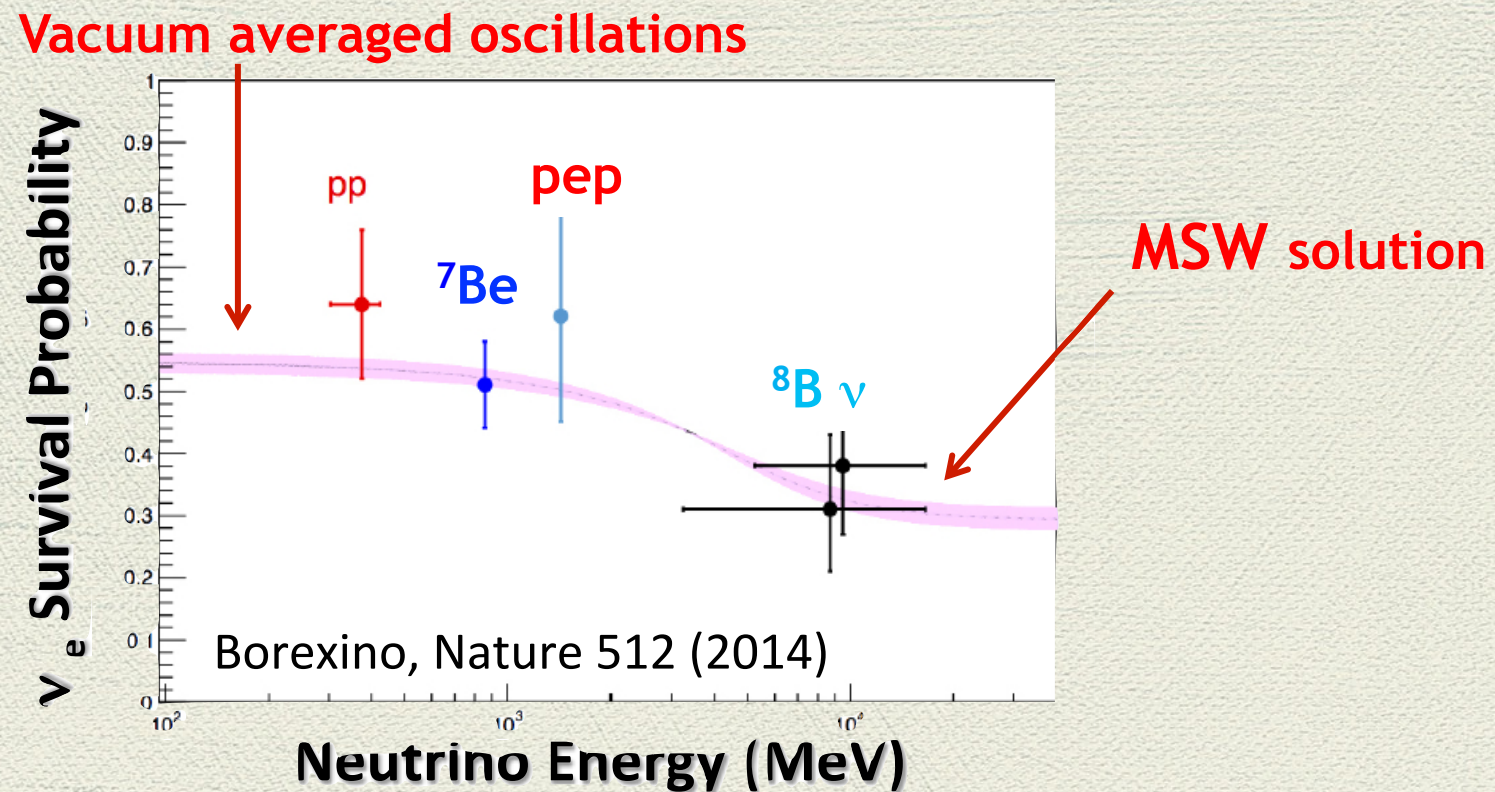
Resonance condition :

$$h_{\nu,11} - h_{\nu,22} \approx 0$$

Adiabaticity :  $\gamma = \frac{|\dot{\theta}_M|}{\Delta\tilde{m}^2/4E} \ll 1$



# Current status of solar neutrino observations



Vacuum-averaged oscillations versus MSW suppression of high energy  ${}^8\text{B}$  neutrinos.

Energy production of low mass main sequence stars confirmed — pp reaction chain.

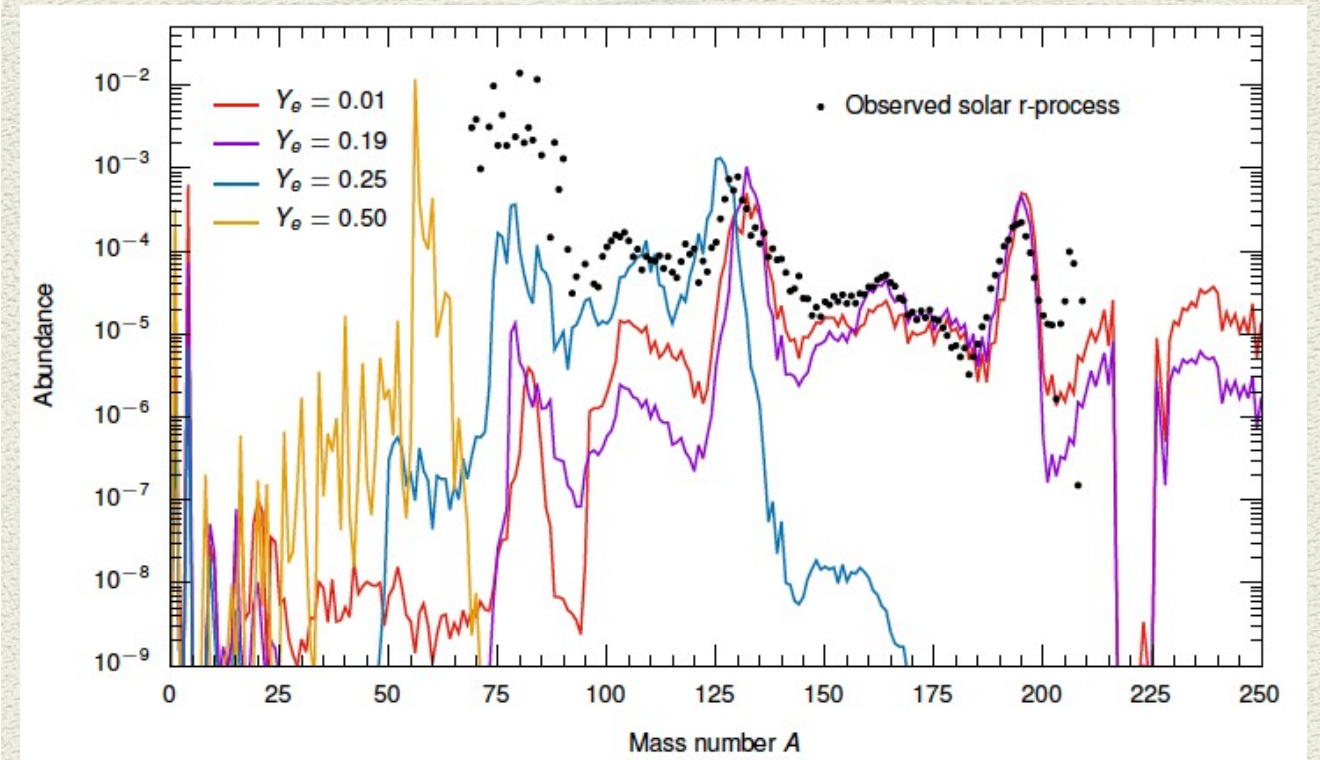
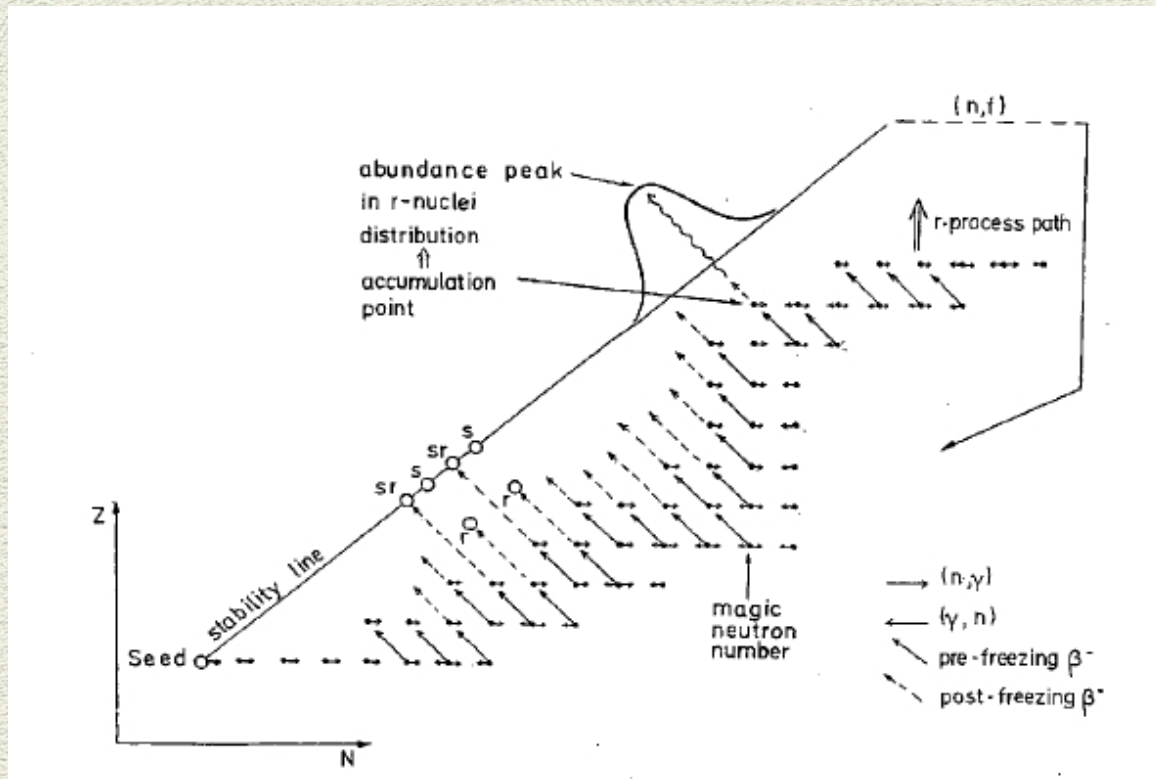
Future measurement of CNO neutrinos - main energy production in massive main sequence stars.



# The site of heavy elements nucleosynthesis is still unknown

- Heavy elements ( $A > 90$ ) are produced in stellar environments during a fast process in presence of a large number of neutrons ( $r$ -process).

$$Y_e = \frac{p}{p + n}$$

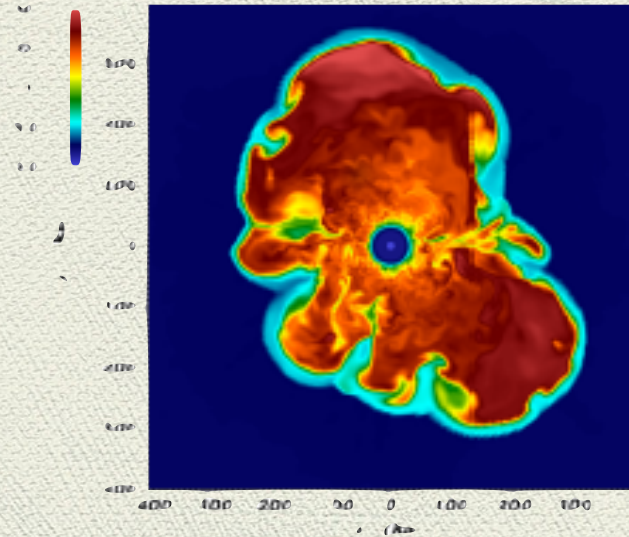


Roberts 2015, arXiv :1508.03133

- Supernovae and accretion disks around compact objects - black holes and binary neutron stars - are candidate sites with lots of neutrinos.



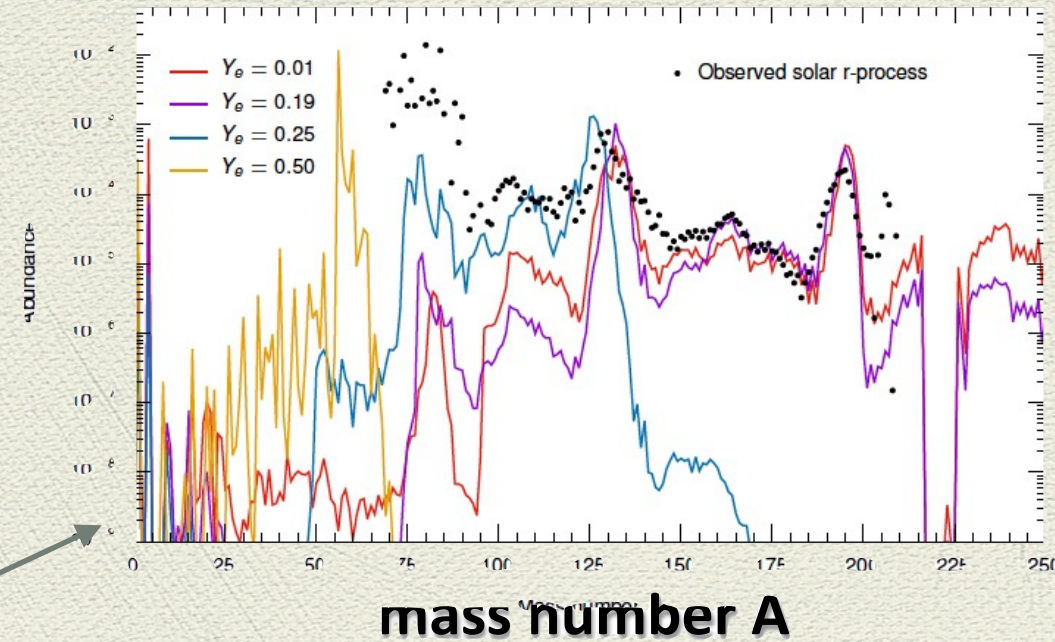
# Neutrinos in dense environments



Mueller et al,  
arXiv: 1705.00620

Iron core-collapse supernova explosions

Understanding  
the role of neutrinos  
and flavor conversion



Heavy elements nucleosynthesis

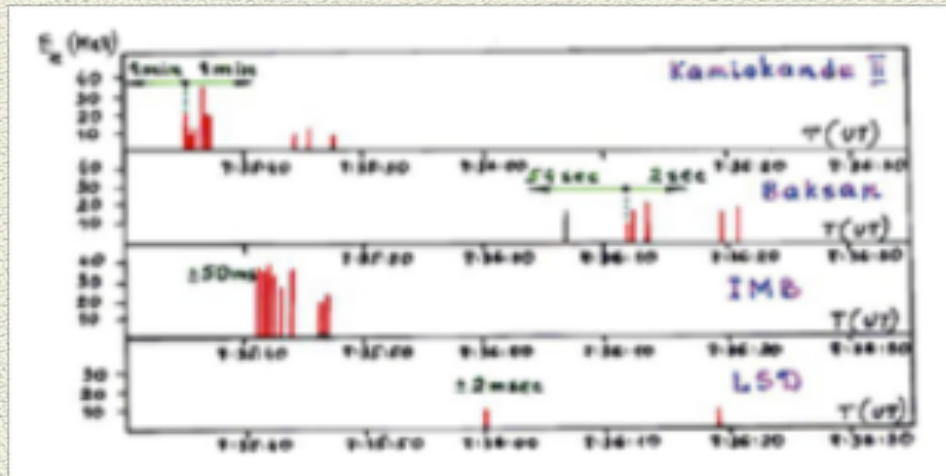
Predicting the neutrino fluxes for future observations :

- an (extra)galactic supernova - we expect  $10^4$ - $10^6$  events at 10 kpc
- the diffuse supernova neutrino background - EGADS project (Super-K + Gd)



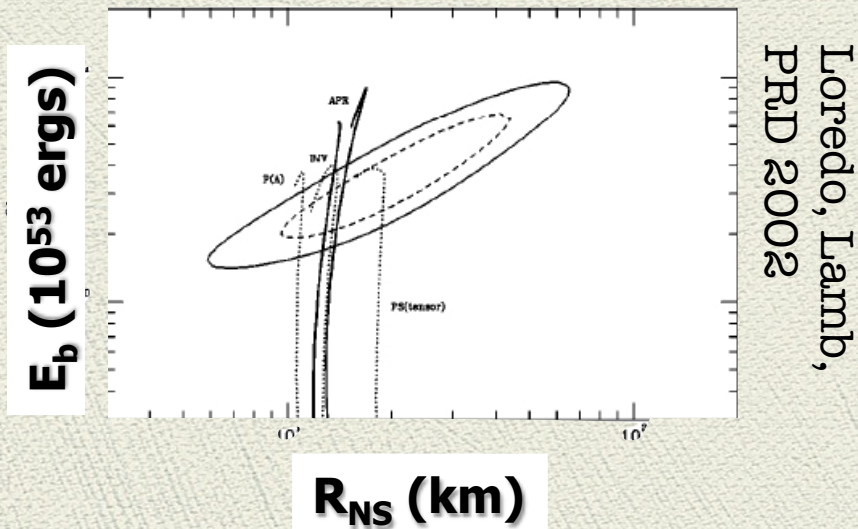
# Neutrinos from core-collapse Supernovae

Suzuki, A. J. Conf. Phys. 120 (2008)

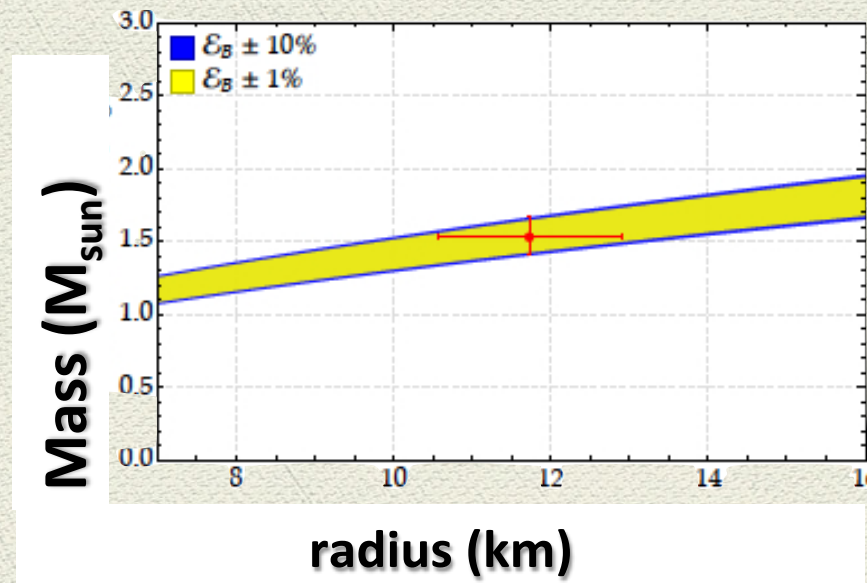


Sanduleak 69<sup>o</sup>202, a blue super-giant in Large Magellanic Cloud, at 50 kpc, no remnant found so far.

**SN1987A**: Delayed explosion mechanism favored over the prompt one.



Loredo, Lamb, PRD 2002

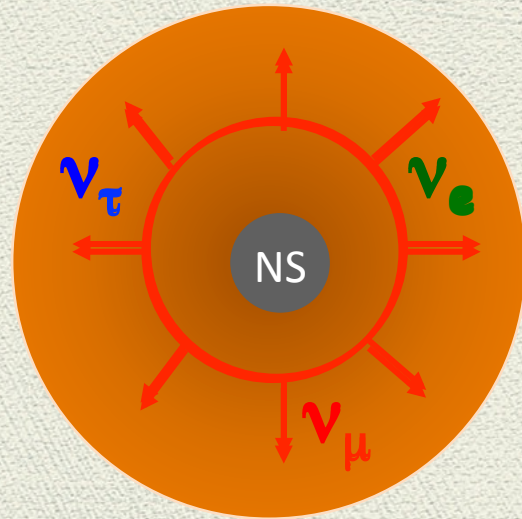


Gallo Rosso et al, JCAP1711 (2017)

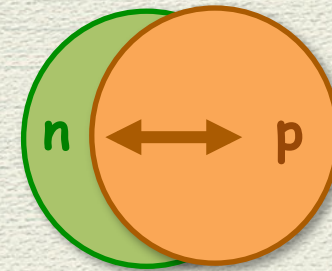
Predictions for Super-K and Hyper-K in case of a supernova at 10 kpc



# Neutrino flavor evolution in dense environments : a many-body problem



$\nu$  in stars or accretion disks



atomic nucleus

$10^{57}$

weak

unbound

$$\rho_{ji} = \langle a_i^\dagger a_j \rangle \text{ neutrinos}$$

$$\bar{\rho}_{ji} = \langle b_i^\dagger b_j \rangle \text{ anti-neutrinos}$$

$N$

interaction

system

density

200

strong

bound

$$\rho = \langle a^\dagger a \rangle \text{ neutrons}$$

$$\text{protons}$$



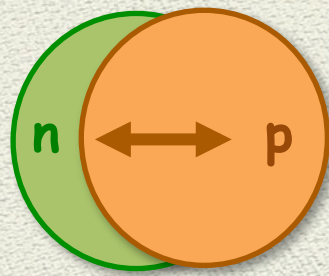
## To determine the dynamics

$$\rho_1 = \langle a^\dagger a \rangle \quad \rho_{12} = \langle a^\dagger a^\dagger a a \rangle \quad \rho_{123} = \langle a^\dagger a^\dagger a^\dagger a a a \rangle \quad \dots$$

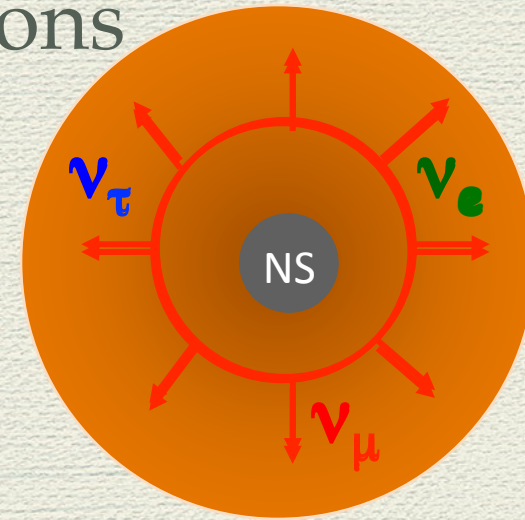
one-body density      two-body      three-body      N-body



# The neutrino evolution equations



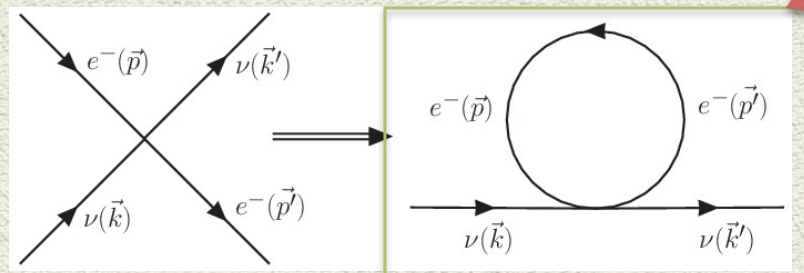
$$\rho = \langle a^\dagger a \rangle \quad \text{one-body density matrix}$$



$$\rho = \begin{pmatrix} \langle a_{\nu_e}^\dagger a_{\nu_e} \rangle & \langle a_{\nu_\mu}^\dagger a_{\nu_e} \rangle \\ \langle a_{\nu_e}^\dagger a_{\nu_\mu} \rangle & \langle a_{\nu_\mu}^\dagger a_{\nu_\mu} \rangle \end{pmatrix}$$

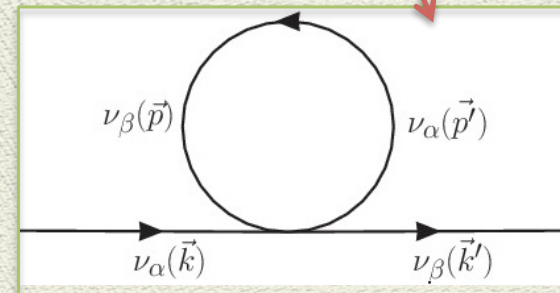
**BBGKY hierarchy : mean-field and beyond**

$$i\dot{\rho} = [h(\rho), \rho] \quad h = h_{vac} + h_{mat} + h_{\nu\nu}(\rho)$$



neutrino-matter

$$h_{mat} = \sqrt{2}G_F \rho_e$$



neutrino self-interactions

**non-linear term**

$$h_{\nu\nu} = \sqrt{2}G_F \sum_{\alpha} \left[ \int (1 - \hat{q} \cdot \hat{p}) \times [dn_{\nu_{\alpha}} \rho_{\nu_{\alpha}}(\vec{p}) - dn_{\bar{\nu}_{\alpha}} \bar{\rho}_{\bar{\nu}_{\alpha}}(\vec{p})] \right],$$

Volpe, Väänänen, Espinoza. PRD 87 (2013)

Volpe, «Neutrino quantum kinetic equations », Int. J. Mod. Phys.E24(2015)

MEAN-FIELD approximation

**Novel conversion phenomena** in dense media due to the neutrino self-interaction



# Nucleosynthesis and neutrinos

- Neutrinos influence the neutron richness of the material through :



that sets the **neutron-to-proton ratio** or **electron fraction** -  $Y_e = \frac{p}{p+n}$

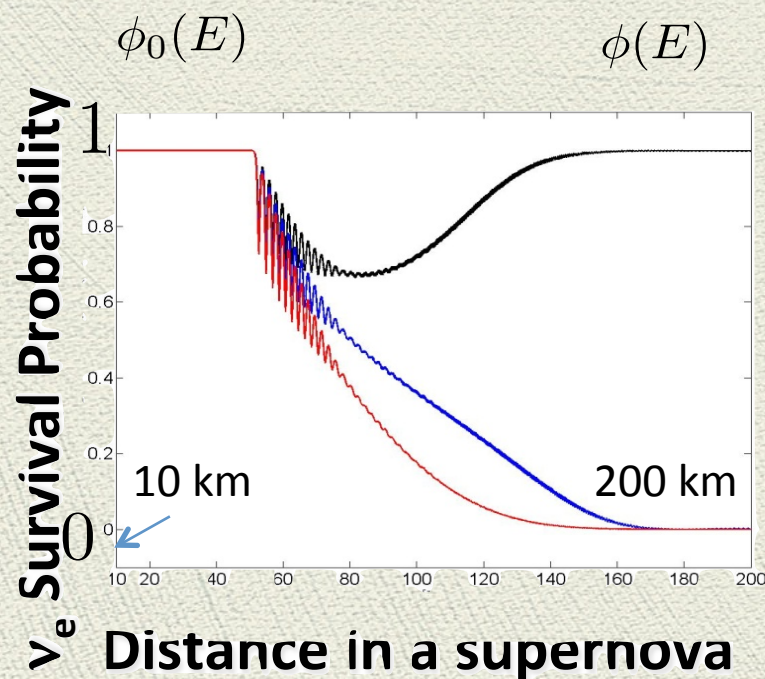
- Neutrino flavor evolution influences  $Y_e$  because flavor modification produces spectral swapping(s).

## AN EXAMPLE

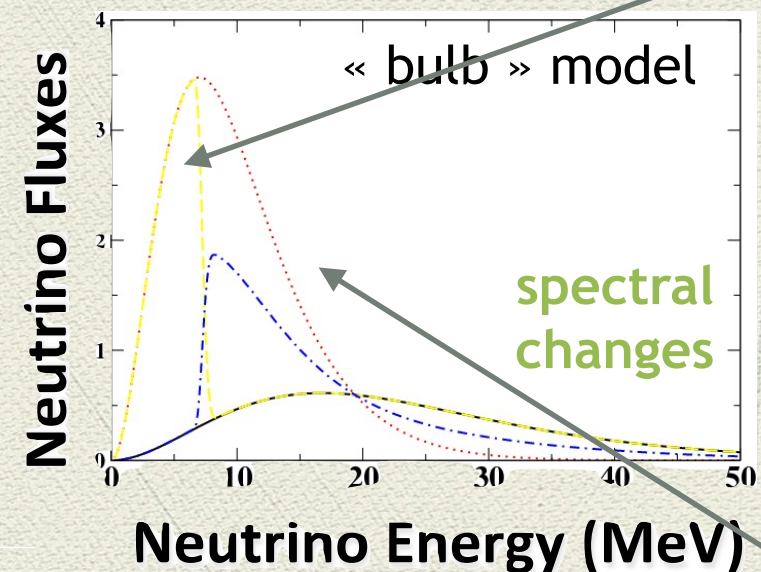
$$\phi_{\nu_e}(E) = P(\nu_e \rightarrow \nu_e)\phi_{\nu_e}^0(E) + [1 - P(\nu_e \rightarrow \nu_e)]\phi_{\nu_x}^0(E)$$

neutrino fluxes at  
the neutrino sphere

neutrino fluxes at  
200 km from it



If  $P(\nu_e \rightarrow \nu_e) \approx 1$   $\phi_{\nu_e}(E) = \phi_{\nu_e}^0(E)$

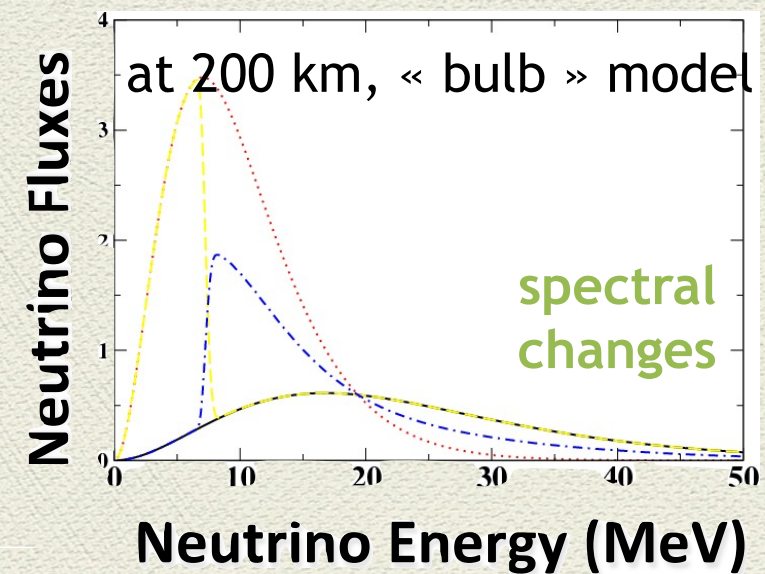
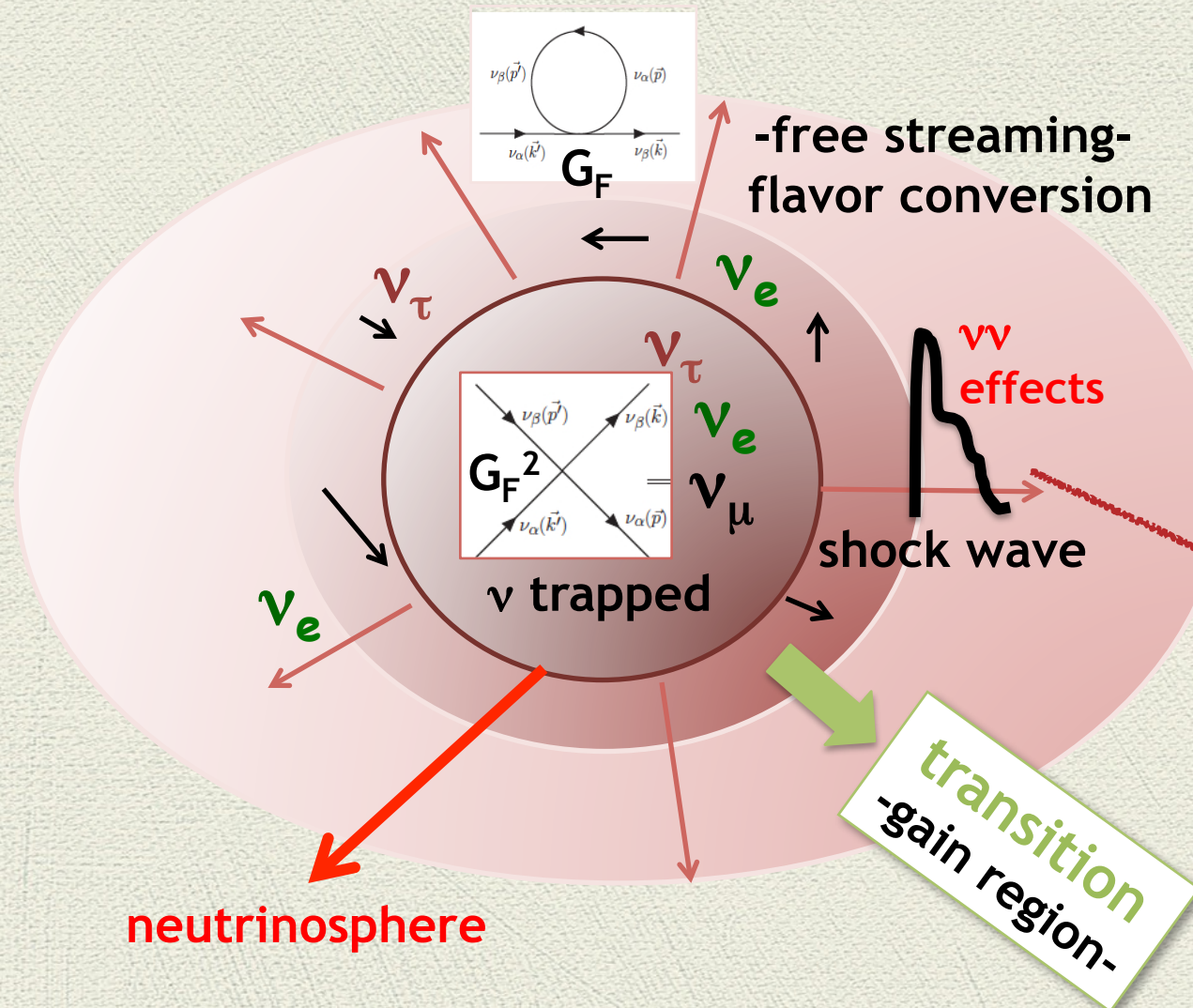
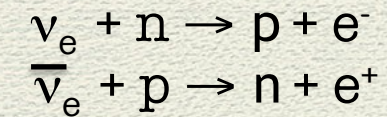


If  $P(\nu_e \rightarrow \nu_e) \approx 0$   $\phi_{\nu_e}(E) = \phi_{\nu_x}^0(E)$



# Supernovae explosions and flavor evolution

The heating rate, behind the shock, could be enhanced by spectral changes of the neutrino fluxes.

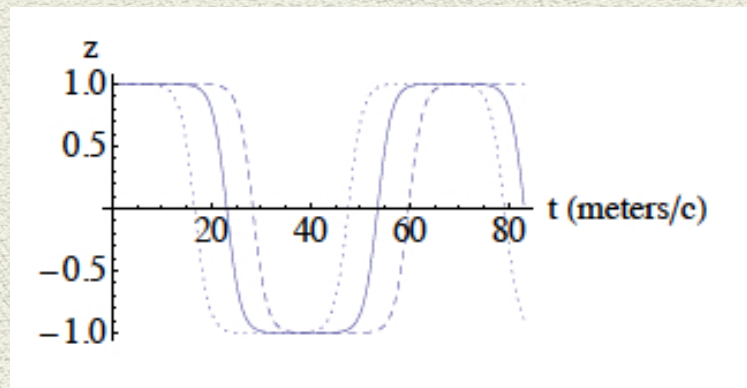


Sharp transition from the dense (Boltzmann) to the dilute (mean-field) region.

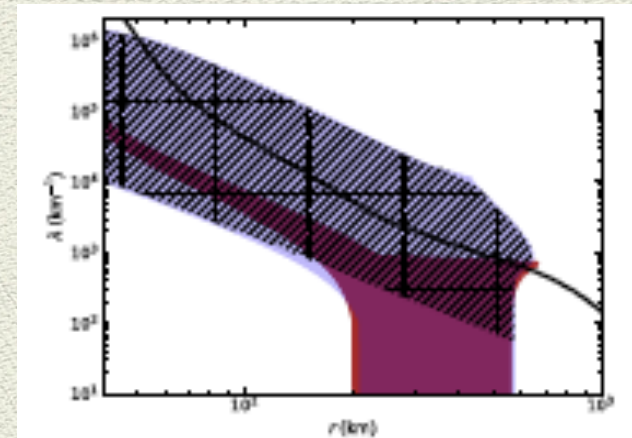


# Improved description of the « transition » region

- Appearance of « fast » conversion modes on short distance scales, if emission at the neutrino-sphere is anisotropic.



Sawyer, PRL108 (2016)



Abbar, Duan, arXiv:1712.07013

- Corrections to the evolution equations from correlators with helicity change, due to neutrino mass.

$$\xi = \langle a_+^+ a_- \rangle$$

$$\mathcal{R} = \begin{pmatrix} \rho & \xi \\ \xi^* & \bar{\rho} \end{pmatrix} \quad \mathcal{H} = \begin{pmatrix} h & \Phi \\ \Phi^* & \bar{h} \end{pmatrix}$$

$\mathcal{R}$  and  $\mathcal{H}$  have helicity and flavor structure ( $2\mathcal{N}_f \times 2\mathcal{N}_f$ ).

$\Phi$  couples  $\nu$  with  $\bar{\nu}$   
helicity (or spin) coherence

$$\Phi \sim (h_{\text{mat}}^{\text{perp}} + h_{\nu\nu}^{\text{perp}}) \times m/2E$$

Vlasenko, Fuller, Cirigliano, PRD89 (2014)

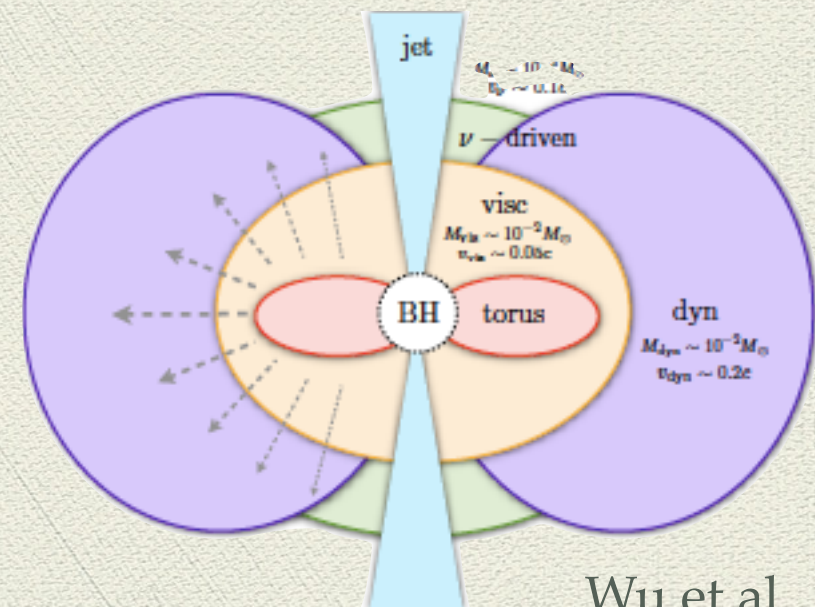
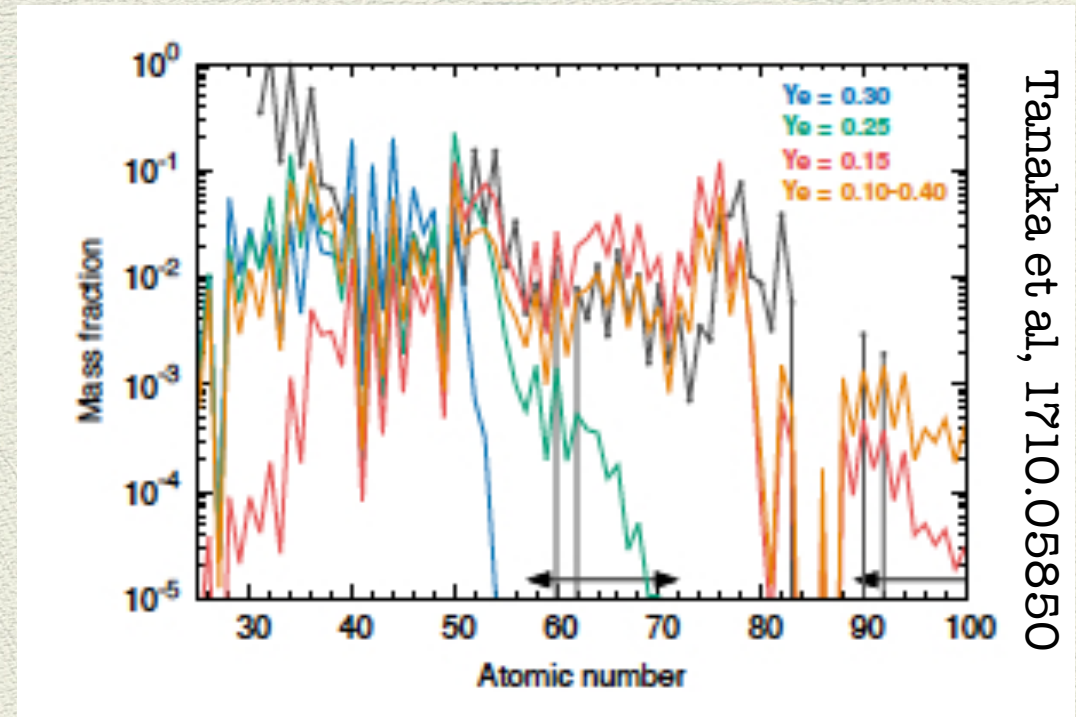
Serreau, Volpe, PRD90 (2014)

A first calculation has shown helicity coherence could modify flavor evolution significantly.



# GW170817 and the kilonova observations

- The recent observation of gravitational waves from binary neutron star mergers, in coincidence with a short gamma-ray-burst and a kilonova.
- The electromagnetic signal covering has a red and a blue component that indicate the presence of r-process elements and in particular lanthanides in the ejecta.
- Lanthanides elements are extremely sensitive to the electron fraction  $Y_e$ . If  $Y_e > 0.25$  lanthanides are not synthesized. Neutrinos drive  $Y_e$  to large values of  $Y_e$ .
- Observations and comparisons with binary neutron star mergers models tell us that there are dynamical ejecta from the early merging phase, with  $Y_e < 0.25$ , and ejecta from neutrino-driven winds in the late time post-merger phase with  $Y_e > 0.25$ .



Wu et al., 2017

Flavor evolution in  $\nu$ -driven winds?



# Helicity coherence in binary neutron star merger remnants

Chatelain, Volpe, PRD 95 (2017), 1611.01862

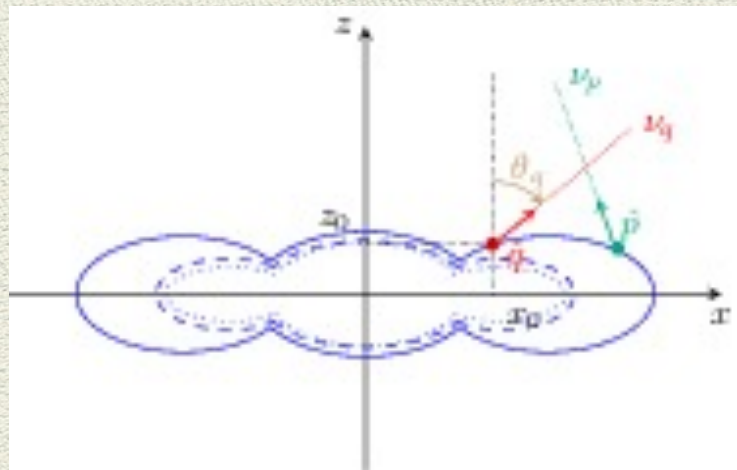
Investigated the role of correlators with helicity change in an extended mean-field description.  
Solved the evolution equations with

$$\mathcal{H} = \begin{pmatrix} h & \Phi \\ \Phi^* & \bar{h} \end{pmatrix}$$

In two flavors  
 $\mathcal{H}$  is a 4 x 4 matrix.

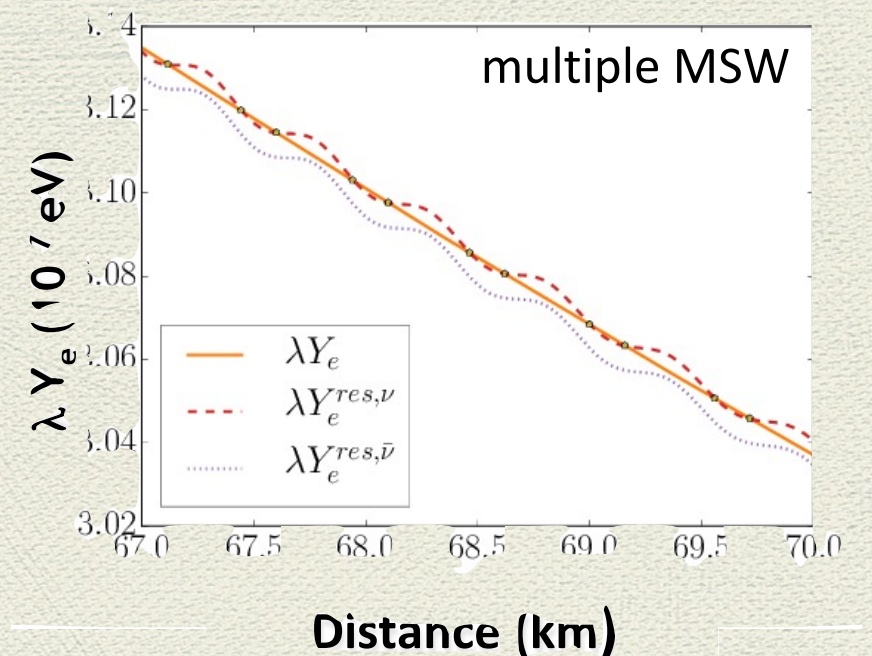
$\Phi$  couples  $\nu$  with  $\bar{\nu}$   
helicity coherence

$$\Phi = (h_{\text{mat}}^{\text{perp}} + h_{\nu\nu}^{\text{perp}}) \times m/2E$$



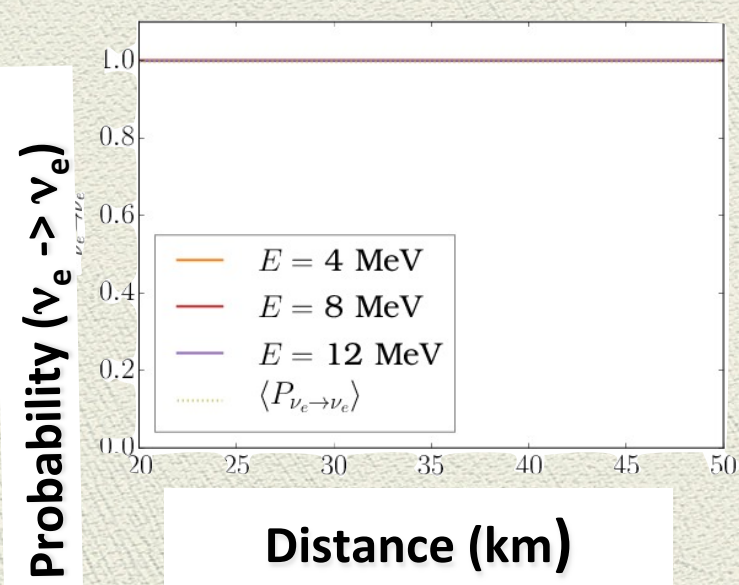
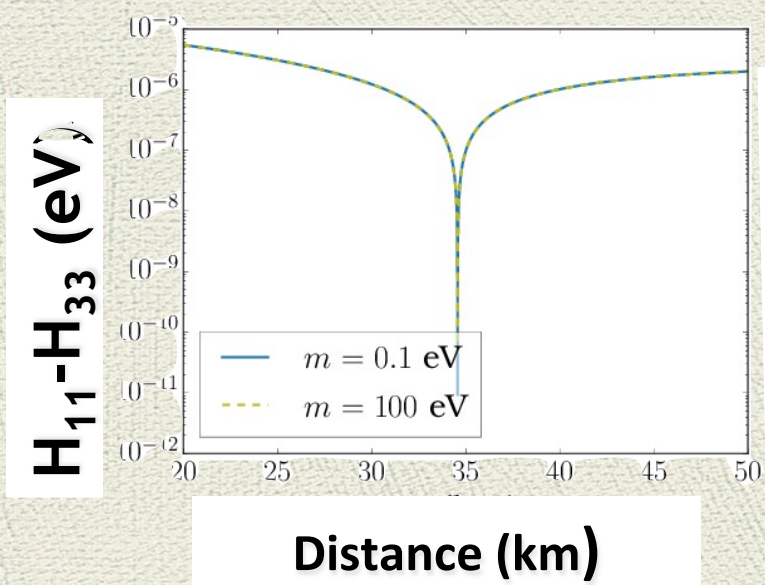
	$\langle E_\nu \rangle$	$L_\nu$	$R_\nu$ (km)
$\nu_e$	10.6	15	84
$\bar{\nu}_e$	15.3	30	60
$\nu_x$	17.3	8	58

MeV  $10^{51}$  erg/s



Resonance condition fulfilled  
but adiabaticity not enough  
to modify the flavor content.

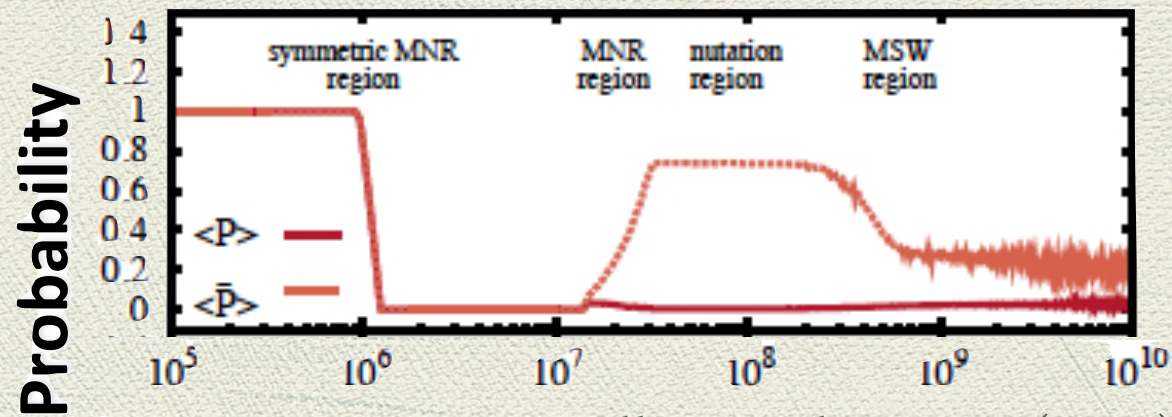
contrary to the findings in Vlasenko,  
Fuller, Cirigliano, 1406.6724



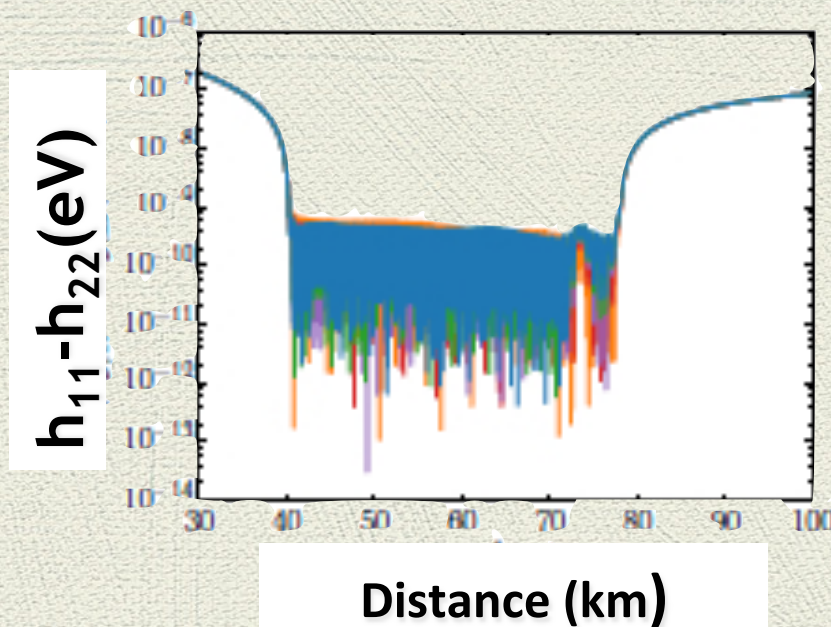
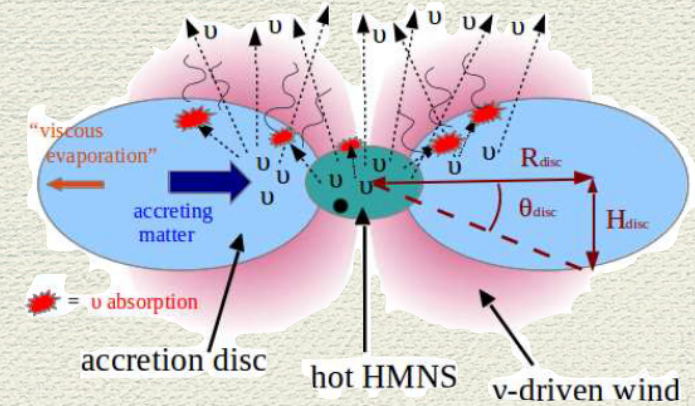


# Matter-neutrino resonance in binary neutron star mergers

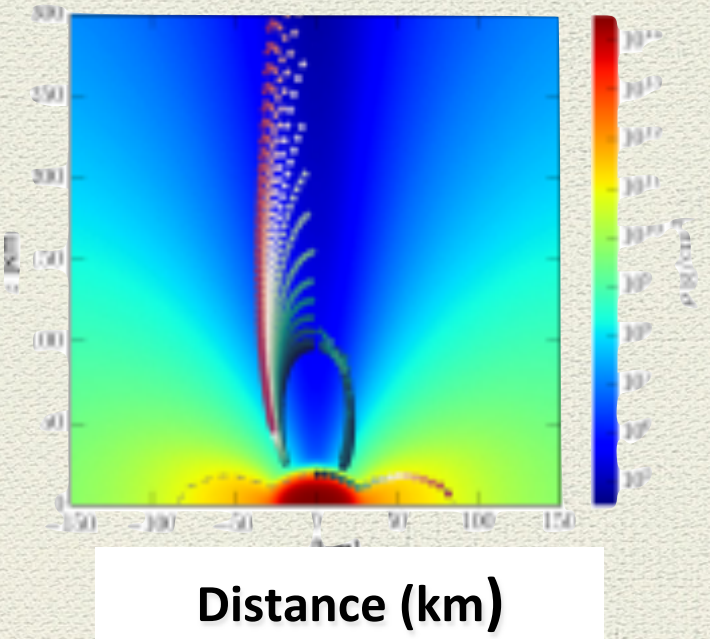
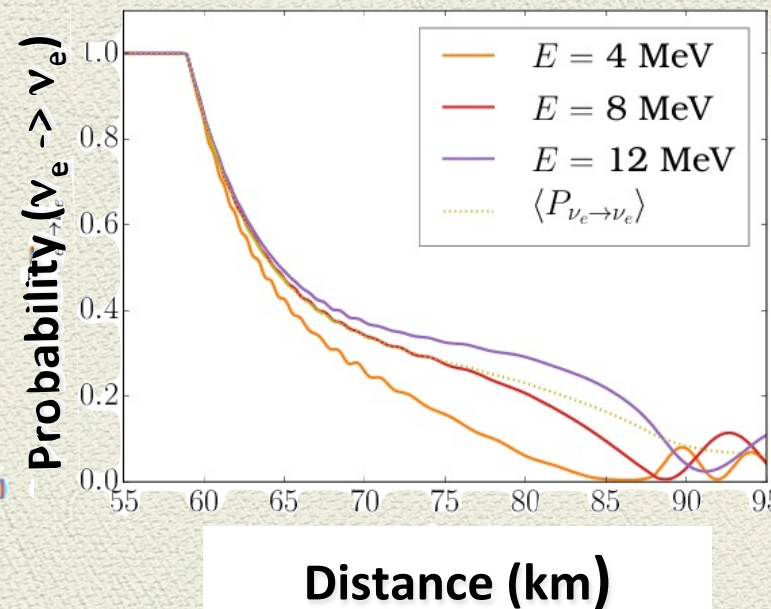
In binary neutron star mergers, the electron antineutrino excess can produce a cancellation of the matter and self-interaction potentials. This is known as the **Matter-Neutrino Resonance**.



Malkus et al, PRD86 (2002)



Chatelain, Volpe, PRD 95 (2017), 1611.01862



Frensel et al., PRD95 (2017)

Resonance condition fulfilled, adiabatic evolution, flavor modified for electron neutrinos and anti-neutrinos. MNR resonances occur in the polar region.



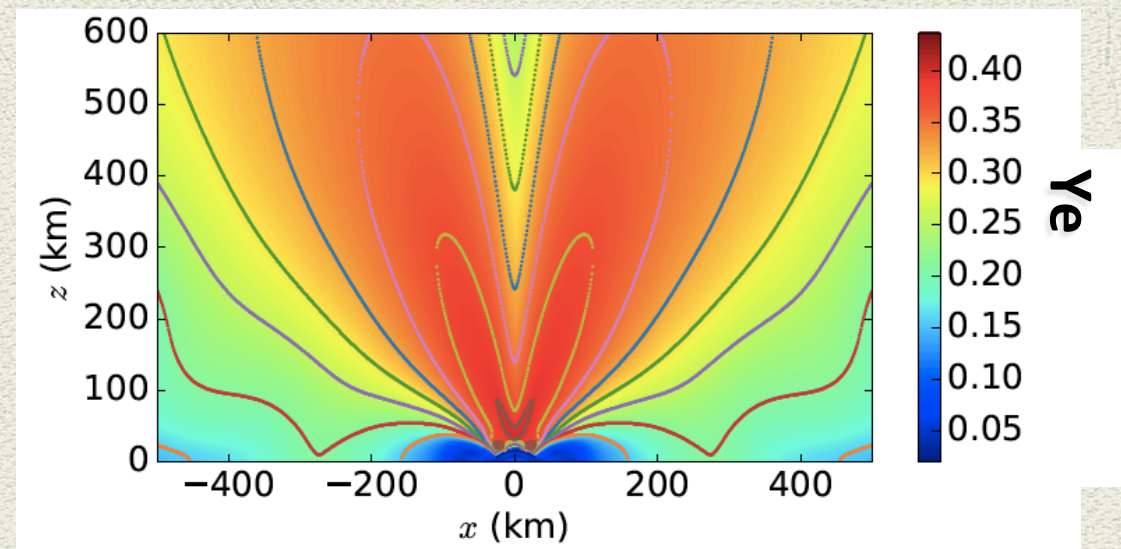
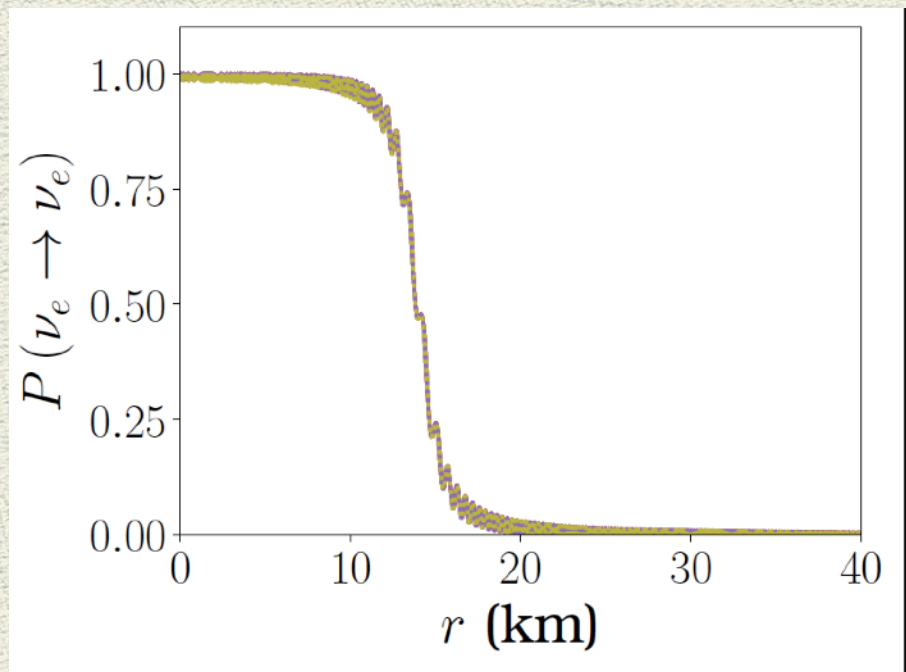
# Flavor evolution and non-standard interactions

$$\left( \begin{array}{l|l} |\epsilon_{ee}| < 2.5 & |\epsilon_{e\tau}| < 1.7 \\ \hline & |\epsilon_{\tau\tau}| < 9.0 \end{array} \right).$$

$$h_{\text{NSI}} = \lambda \left( \begin{array}{l|l} \left(\frac{Y_{\odot} - Y_e}{Y_{\odot}}\right) \delta\epsilon^n & (3 + Y_e)\epsilon_0 \\ \hline (3 + Y_e)\epsilon_0^* & 0 \end{array} \right).$$

The **I-resonance** is due to a cancellation between the standard and non-standard matter terms.

It can be seen also as a **synchronized MSW** resonance, where all effective spins in flavor space undergo the resonance coherently.

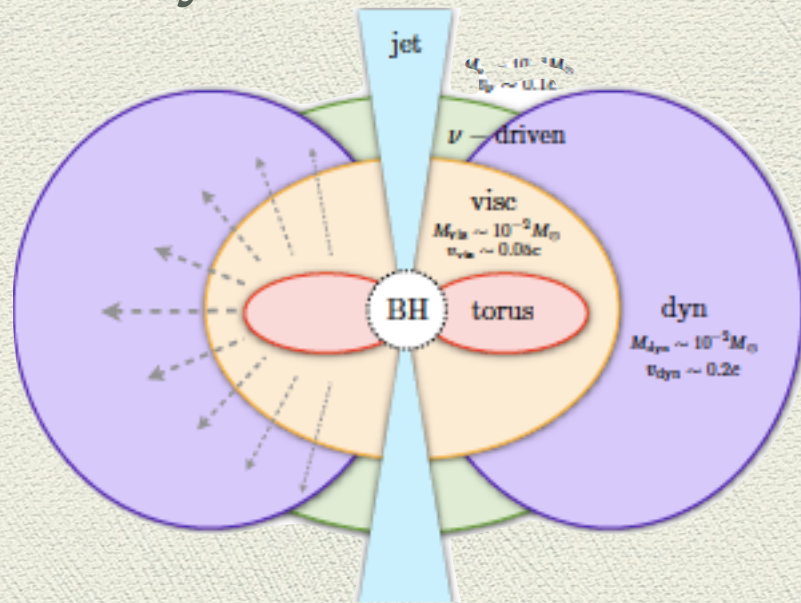


Chatelain, Volpe, PRD98 (2018)

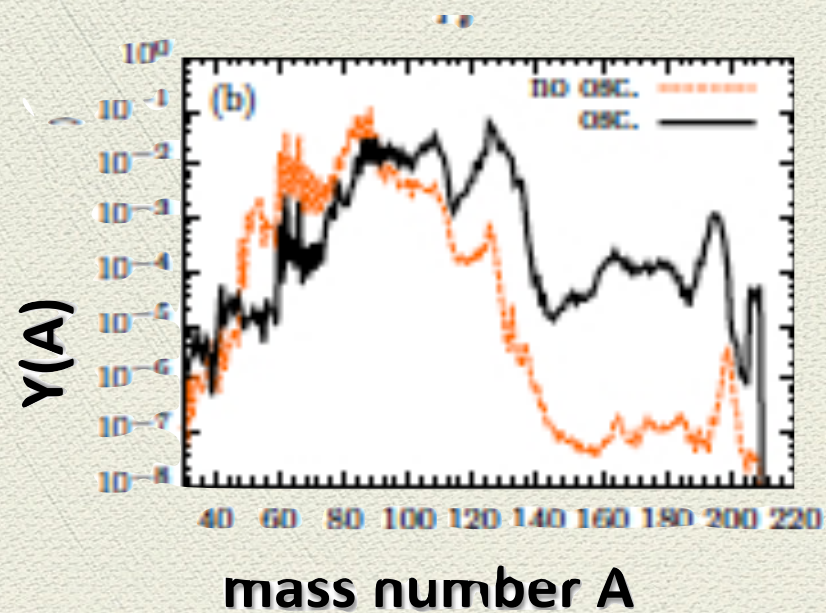
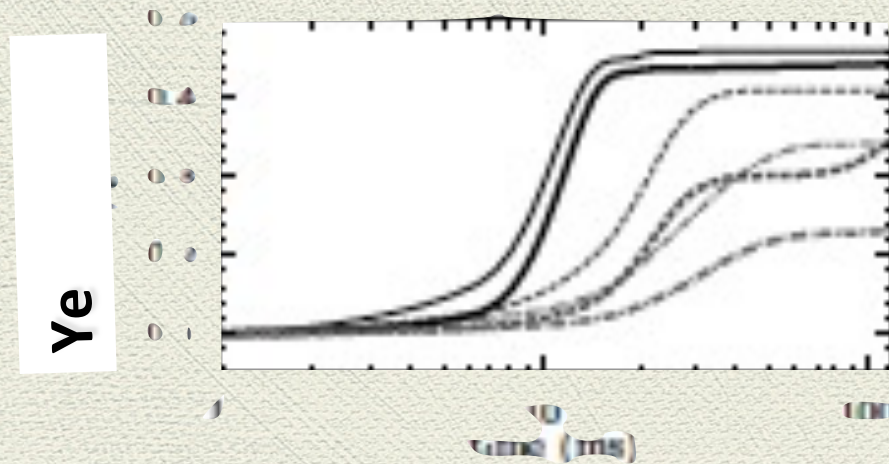
The I-resonance produces flavor modification nearby the neutrino sphere.



# Nucleosynthesis in neutrino-driven winds and kilonovae



« Fast modes » might bring an equilibration of the neutrino fluxes. Here an example of the impact on nucleosynthesis, in a schematic calculation of the flavor evolution.




Wu et al., PRD96 (2017)


Flavor evolution here tends to decrease  $Y_e$  and favor a strong r-process.




## Conclusions

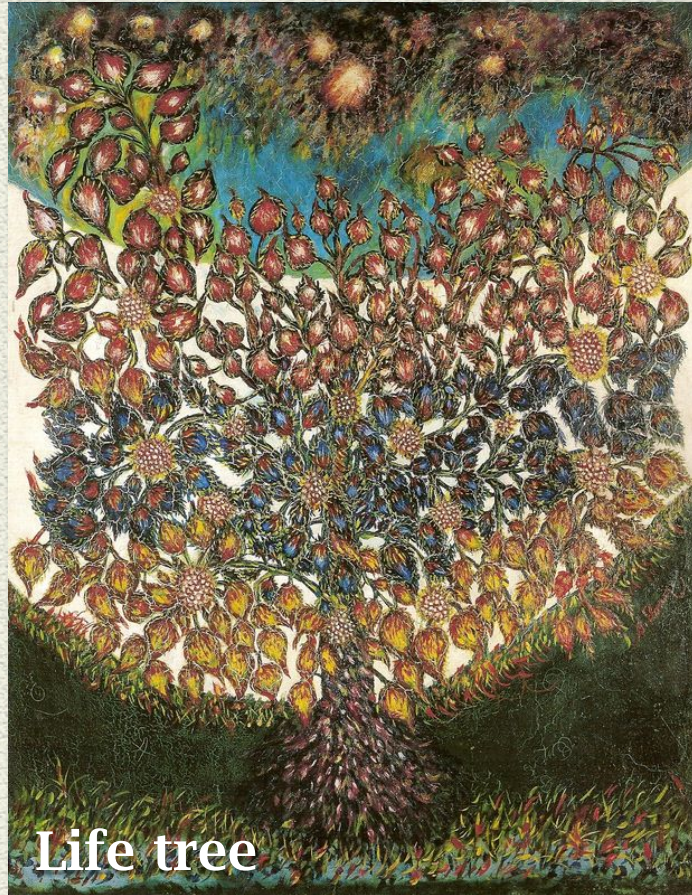
 Neutrino propagating in dense media modify their flavor in intriguing ways, because of neutrino self-interactions (but also shock waves, turbulence, ...)

 Intense activity to unravel the conditions for the occurrence of these modes and their nature. Fast modes on short time scales as well as other collective modes can modify the shock dynamics and determine the neutron-to-proton ratio.

 Neutrino flavor conversion in binary neutron star mergers can influence r-process nucleosynthesis in neutrino-driven winds. To definitely assess its influence is important to understand their contribution in the ejecta of kilonovae and the associated electromagnetic signal.

 Many open issues remain, including the role of decoherence, the influence of gravity nearby compact objects, the role of symmetry breaking, the impact of « fast » modes and of collisions, ...





Life tree

*Seraphine de Senlis*