Near-future prospects for nuclear PDFs

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Latest available nPDF parametrizations in 2018

	00	10		45	10
	EPS09	DSSZ12	KA15	NCTEQ15	EPPS16
Order in α_s	NLO	NLO	NNLO	NLO	NLO
DIS in ℓ [−] +A	✓	✓	$\overline{}$	✓	\checkmark
Drell-Yan in p+A	✓	✓	✓	✓	✓
RHIC pions d+Au	✓	✓		✓	\checkmark
Neutrino-nucleus DIS		✓			✓
Drell-Yan in π + A					✓
LHC p+Pb dijets					✓
LHC p+Pb W, Z					√
Q cut in DIS	1.3 GeV	1 GeV	1 GeV	$2\mathrm{GeV}$	$1.3\mathrm{GeV}$
datapoints	929	1579	1479	708	1811
free parameters	15	25	16	16	20
error analysis	Hessian	Hessian	Hessian	Hessian	Hessian
error tolerance $\Delta\chi^2$	50	30	N.N	35	52
proton baseline PDFs	CTEQ6.1	MSTW2008	JR09	стео6м-like	CT14NLO
Heavy-quark effects		✓		✓	√
Flavour separation				partial	full
Reference	JHEP 0904 065	PR D85 074028	PR D93, 014026	PR D93 085037	EPJ C77 163

Expect "SOON": Andrés-Zurita NNLO nPDFs [https://indico.cern.ch/event/639067/contributions/2642447/]

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This talk restricted to:

- \bullet CMS $\sqrt{s}=5\,\mathrm{TeV}$ p-p and p-Pb dijets
- Towards including LHCb D-meson measurements

 \bullet Normalized dijet cross sections at $\sqrt{s}=5\,\mathrm{TeV}$ in p-p [CMS PAS HIN-16-003]

$$\frac{d\sigma^{\mathrm{pp}}(\eta_{\mathrm{dijet}}, p_{\mathrm{T}}^{\mathrm{average}})}{\int d\sigma^{\mathrm{pp}}(\eta_{\mathrm{dijet}}, p_{\mathrm{T}}^{\mathrm{average}}) d\eta_{\mathrm{dijet}}} \qquad p_{\mathrm{T}}^{\mathrm{average}} = (p_{\mathrm{T}}^{\mathrm{leading}} + p_{\mathrm{T}}^{\mathrm{subleading}})/2$$

$$\eta = (\eta^{\mathrm{leading}} + \eta^{\mathrm{subleading}})/2$$

$$\frac{(\eta_{\mathrm{sw}}) = 5.02 \, \mathrm{TeV}}{\rho_{\mathrm{reliminary}}} \qquad \frac{\rho_{\mathrm{reliminary}}}{\rho_{\mathrm{reliminary}}} \qquad \frac{\rho_{\mathrm{preliminary}}}{\rho_{\mathrm{reliminary}}} \qquad \frac{\rho_{\mathrm{preliminary}}}{\rho_{\mathrm{reliminary}}} \qquad \frac{\rho_{\mathrm{preliminary}}}{\rho_{\mathrm{reliminary}}} \qquad \frac{\rho_{\mathrm{preliminary}}}{\rho_{\mathrm{preliminary}}} \qquad \frac{\rho_{\mathrm{preliminary}}}{\rho_{\mathrm{preliminary}$$

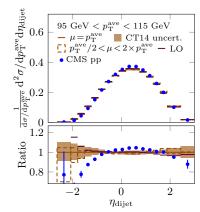
• The preliminary data not well reproduced by the current PDFs

 \bullet Normalized dijet cross sections at $\sqrt{s}=5\,\mathrm{TeV}$ in p-p [CMS PAS HIN-16-003]

$$\frac{d\sigma^{\rm pp}(\eta_{\rm dijet}, p_{\rm T}^{\rm average})}{\int d\sigma^{\rm pp}(\eta_{\rm dijet}, p_{\rm T}^{\rm average})d\eta_{\rm dijet}}$$

$$\begin{split} p_{\mathrm{T}}^{\mathrm{average}} &= (p_{\mathrm{T}}^{\mathrm{leading}} + p_{\mathrm{T}}^{\mathrm{subleading}})/2 \\ \eta &= (\eta^{\mathrm{leading}} + \eta^{\mathrm{subleading}})/2 \end{split}$$

- Normalization apparently suppresses the experimental systematic uncertainties
- NLO calculations are wider in $\eta_{\rm dijet}$ than the (preliminary) data
- The scale uncertainty is very small for $\eta_{
 m dijet} \in [-1,2]$ would not expect large NNLO effects. Already NLO to LO difference is small near $\eta_{
 m dijet} \sim 0$

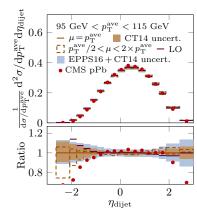


ullet Normalized dijet cross sections at $\sqrt{s}=5\,\mathrm{TeV}$ in p-Pb [CMS PAS HIN-16-003]

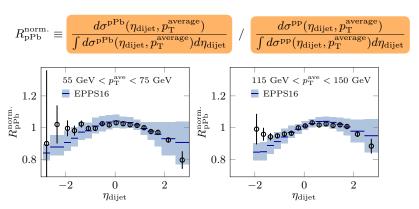
$$\frac{d\sigma^{\rm PPb}(\eta_{\rm dijet},p_{\rm T}^{\rm average})}{\int d\sigma^{\rm PPb}(\eta_{\rm dijet},p_{\rm T}^{\rm average})d\eta_{\rm dijet}}$$

$$\begin{split} p_{\mathrm{T}}^{\mathrm{average}} &= (p_{\mathrm{T}}^{\mathrm{leading}} + p_{\mathrm{T}}^{\mathrm{subleading}})/2 \\ \eta &= (\eta^{\mathrm{leading}} + \eta^{\mathrm{subleading}})/2 \end{split}$$

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- Similar situation in p-Pb measurements

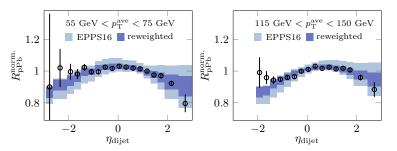


Nuclear modification — a ratio of ratios



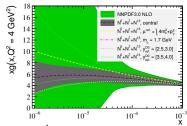
- ullet The preliminary $R_{
 m pPb}^{
 m norm.}$ data decently described by EPPS16
- The data uncertainties beat EPPS16 by far
- ullet Some deviations in the backward direction (probing **very** large- x_{Pb} region)

 The effect of these CMS preliminary data estimated by an improved PDF-reweighting/profiling method (check out P. Paakkinen in DIS'18)

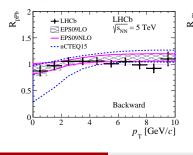


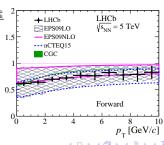
- The preliminary data promise a major effect in EPPS16 even more dramatic for nCTEQ15 which has larger gluon uncertainties at large x.
- Have to settle down the issues (discussed in the previous slides) with the normalized spectra before can include these data in global nPDF fits.

- The potential of D (and B) meson production has been demonstrated in p-p [PRL 118 072001, EPJ C75 396]
- Good gluon resolution based on including data down to $p_{\rm T}^{\rm D}=0$



• Recent $R_{\rm pPb}$ data from LHCb [JHEP 1710 (2017) 090] show compelling evidence of small-x shadowing



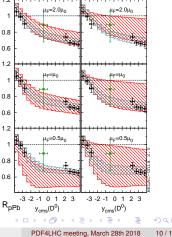


Idea introduced in [EPJ C77 (2017) 1] and then applied in [ARXIV:1712.07024]:

$$d\sigma(\mathrm{D}^0) = f_g(x_1, Q_f^2) \otimes \frac{d\sigma_{gg}^{\mathrm{D}^0}(Q_f^2, Q_r^2)}{\otimes f_g(x_2, Q_f^2)}$$

Fit the coefficient functions to p-p data

- Neglects all but the gluon-gluon channel
 - Close to fixed-flavour number scheme (FFNS)
 - EPPS16/nCTEQ15 are not FFNS PDFs...
- Based on $2 \rightarrow 2$ kinematics
 - May bias the x_2 distributions to overly low x
- I would say a more appropriate treatment requires a general-mass variable flavour number (GM-VFNS) approach



In FFNS, the heavy quarks are produced in three partonic processes

$$g + g \rightarrow Q$$
, $q + \overline{q} \rightarrow Q$, $q + g \rightarrow Q$

Phenomenological fragmentation functions (FFs) for $Q \to D$ transition

• FFNS cross sections diverge as $\sim \log(p_{\rm T}^2/m^2)$. In GM-VFNS these logs are resummed into heavy-quark PDFs and scale-dependent FFs

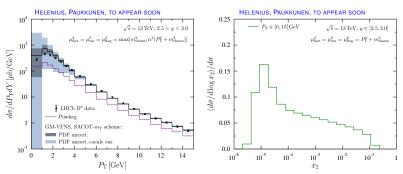
$$\frac{d\sigma(h_1+h_2\to D^0+X)}{dP_{\mathrm{T}}dY} = \sum_{ijk} \int_{z_{\mathrm{min}}}^1 \frac{dz}{z} \int_{x_{\mathrm{1}}^{\mathrm{min}}}^1 dx_1 \int_{x_{\mathrm{2}}^{\mathrm{min}}}^1 dx_2$$

$$f_i^{h_1}(x_1,\mu_{\mathrm{fact}}^2) \quad \frac{d\hat{\sigma}^{ij\to k}(x_1,x_2,m,\mu_{\mathrm{ren}}^2,\mu_{\mathrm{fact}}^2,\mu_{\mathrm{frag}}^2)}{dp_{\mathrm{T}}dy} \quad f_j^{h_2}(x_2,\mu_{\mathrm{fact}}^2) \quad D_{k\to D^0}(z,\mu_{\mathrm{frag}}^2)$$

Coefficient functions behave as FFNS at low $p_{\rm T}$, as zero-mass matrix elements at high $p_{\rm T}$

Scale-dependent, universal FFs

 LHCb p-p cross sections fairly well reproduced by GM-VFNS approach in SACOT-m_T scheme — a generalization of SACOT-χ to hadroproduction



- ullet x_2 distributions in GM-VFNS peaked at low x a long tail towards large x
- ullet Sizable theory uncertainties at low p_{T} : GM-VFNS scheme dependence, ambiguous fragmentation variable z, scale uncertainties, etc...
 - \Rightarrow Will need to set a cut $p_{\rm T} \gtrsim 5\,{\rm GeV}$ or so

Summary

• The (preliminary) CMS $\sqrt{s}=5\,{
m TeV}$ dijet $R_{
m pPb}$ data promises a huge potential to constrain nuclear gluons at large-x

However, the $\eta, p_{\rm T}$ dependence of the (preliminary) p-p and p-Pb spectra are not well reproduced — only the ratio seems OK

- ⇒ Have to sort out these discrepancies before one can confidently include the data in global fits
- Potential of LHCb D-meson measurements in p-Pb are significant at small-x

GM-VFNS approach can describe the p-p data down to $p_{\rm T}=0$ but the theory uncertainties are huge — requires a cut $p_{\rm T}\gtrsim 5~{\rm GeV}$ or so

 \Rightarrow Retain sensitivity to small-x but at higher interaction scale