

# Relativistic Corrections within the Integral Transform Approach

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# How to extend the reliability of n.r. *ab initio* results for e.w. cross sections to high energy/momentum

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- ★ General considerations on the lepton-nucleus hadron tensor and the inclusive response functions  $R^{\mu\nu}(\mathbf{q},\omega)$

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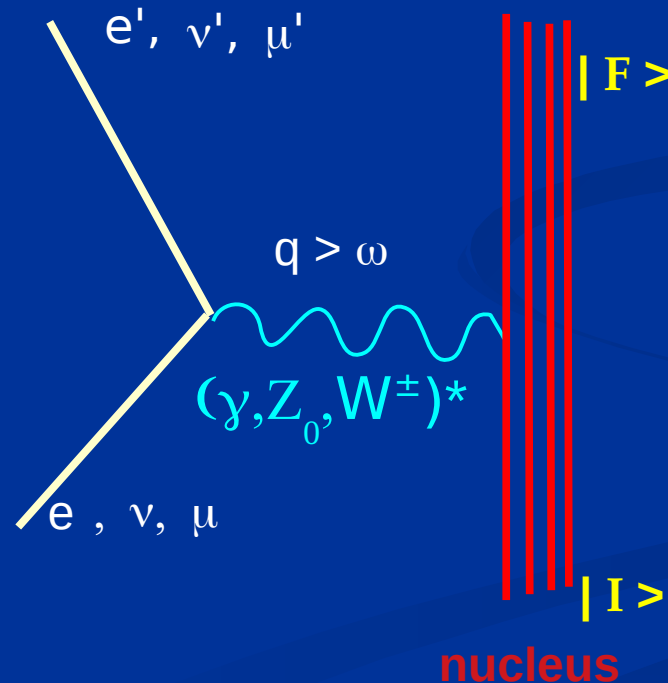
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- ★ From frame dependence to frame independence
- ★ Test on the (e,e’) scattering

# Physics of e.w. Interactions (with nuclei)





at **1st order P.T.** the crucial quantity in the cross section is the **Hadron Tensor**

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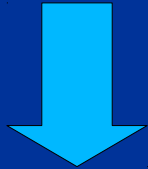
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- $|F\rangle$  can be a bound or a continuum (scattering) state
- $\delta^4$  expresses the energy-momentum conservation

If  $|I\rangle$  is the g.s.  $|0\rangle$  and  $|F\rangle$  is “inclusive”

r

$$W^{\mu\nu} = \langle I | J^\mu | F \rangle \langle F | J^\nu | I \rangle \times \delta^4$$



$$R^{\mu\nu}(\vec{q}, \omega) = \sum_n \langle 0 | J^\mu(\vec{q}) | n \rangle \langle n | J^\nu(\vec{q}) | 0 \rangle^2 \times \delta(\omega - E_n + E_0)$$

# Notice!

- ★ The 3-momentum transfer  $\vec{q}$  originates from the 3-momentum delta-function  $\delta^3$  which now involves the **c.m. of the nucleus**

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- ★ Therefore the non relativistic problem

$$\mathbf{H} |\mathbf{n}\rangle = \mathbf{E}_n |\mathbf{n}\rangle$$

has to be solved, referred to the “internal” **(i.e. translation/galileian invariant)** dynamics

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# The consistent n.r. “ab initio” approach in nuclear physics

- Take as input an Hamiltonian with protons and neutrons as d.o.f. interacting with *realistic*  $V_{\text{NN}}$  (i.e. reproducing *NN* cross sections with  $\chi/\text{datum} \sim 1$ )

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taking into account the **full many-body dynamics**, respecting **translation/Galileian invariance**. No approximation, but **controlling the numerical accuracy**

# The big problems:

- How to solve the Hamiltonian for  $|F\rangle$ , namely the **many-body scattering state**, when the nucleus breaks into pieces, (*known as “final state interaction” FSI!*)

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*[Notice: due to “good” asymptotic boundary conditions the ground state  $|0\rangle$  can be calculated with controlled accuracy, at least up to medium heavy systems ]*

# The big problems:

- How to solve the Hamiltonian for  $|F\rangle$ , namely the **many-body scattering state**, when the nucleus breaks into pieces, (*known as “final state interaction” FSI!*)
- Up to which energy/momentum can one push the *ab initio* **non relativistic** treatment of the **dynamics** in  $R^{\mu\nu}(q, \omega)$ ??



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## The solution:

- The integral transform approach

# The big problem:

- Up to which energy/momentum can one push the *ab initio* non relativistic treatment of the dynamics in  $R^{\mu\nu}(q, \omega)$ ??

# The solution:

- Analyze the frame dependence, choose the “*right frame*” and the “*proper rel. input kinematics*”

**... and test on electromagnetic  
(e,e') cross section**

# The integral transform approach

# Integral transform approaches

There are many examples in physics where one uses  
“integral transform approaches”

$$\Phi = K R$$

accessible object

object of interest

The diagram illustrates the equation  $\Phi = K R$ . A cyan arrow points from the text 'accessible object' to the symbol  $\Phi$ . A red arrow points from the symbol  $R$  to the text 'object of interest'. The symbol  $K$  is white, and the symbol  $R$  is red.

There are many classes of problems that are difficult to solve in their original representations. An integral transform "maps" an equation from its **original "domain"** into **another domain**. Manipulating and solving the equation in the **target domain** is sometimes much easier than manipulation and solution in the **original domain**. The solution is then **mapped back** to the original domain with the inverse of the integral transform.

↓ KERNEL

$$\Phi(\tau) = \int d\omega K(\omega, \tau) R(\omega)$$

One is able to calculate  $\Phi(\tau)$  but wants  $R(\omega)$ ,  
which is the quantity of direct physical meaning.



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Warning:

The “inversion” of  $\Phi(\tau)$  may be problematic (“**ill posed problem**”)

a “good” Kernel has to satisfy two requirements

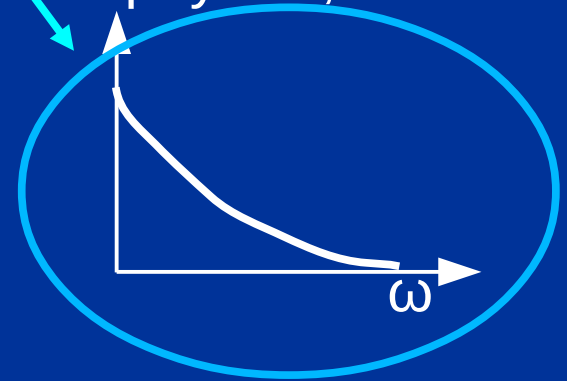
- 1) one must be able to calculate the integral transform
- 2) one must be able to invert the transform controlling the instabilities (*“ill-posedness”*)

Two examples in the literature:



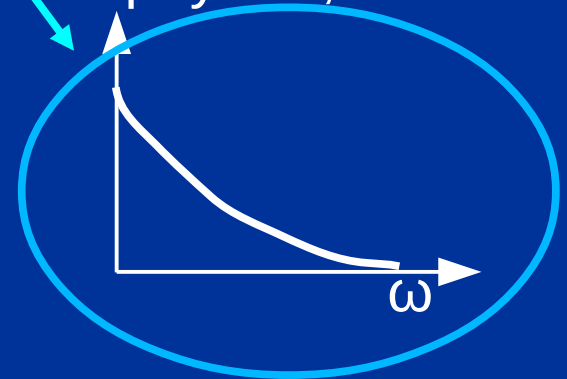
**Exponential Kernel:**  $K(\omega, \tau) = e^{-\omega \tau}$   $\tau$  real

- used in condensed matter physics, nuclear physics, lattice QCD, ...
- Degree of *ill-posedness* : **high**
- $\Phi(\tau)$  calculated by **GFMC**



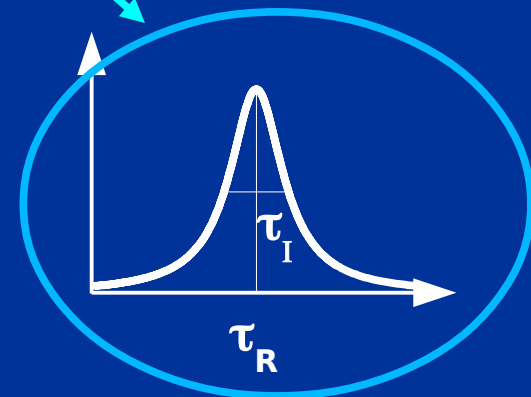
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**Lorentzian Kernel:**  $K(\omega, \tau) = [(\omega - \tau) (\omega - \tau)^*]^{-1}$   
complex =  $\tau_R + \tau_I$

- used in nuclear physics
- Degree of *ill-posedness* : **low**
- $\Phi(\tau)$  calculated via **matrix diagonalization on bound** basis functions



# One can calculate

$$\Phi^{\mu\nu}(q, \tau) = \int d\omega K(\omega, \tau) R^{\mu\nu}(q, \omega)$$

...and then invert  $\Phi$



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the many-body scattering problem of calculating  $|F\rangle$  is avoided!!!

How important are **relativistic effects**  
as **q** increases?

# The analysis of frame dependence

**One criteria to judge the importance of relativistic effects is the frame dependence of the results**



# The electron scattering (e,e') response functions in various frames

$\rho(q)_{fr}$

$$R_L^{fr}(q, \omega^{fr}) = \sum_n \langle 0 | J^0(q)_{fr} | n \rangle \langle n | J^0(q)_{fr} | 0 \rangle^2 \times \delta(\omega_{fr} - E_n^{fr} + E_0^{fr})$$

$$R_T^{fr}(q, \omega^{fr}) = \sum_n \langle 0 | J_{\perp}(q)_{fr} | n \rangle \langle n | J_{\perp}(q)_{fr} | 0 \rangle^2 \times \delta(\omega_{fr} - E_n^{fr} + E_0^{fr})$$

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**LAB:** initially nucleons have momenta  $\mathbf{p}_i \approx 0$

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**ANB:** initially nucleons have momenta  $p_i \approx -q/2$

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They are connected to the response functions  
in the **LAB frame**  
( where they are measured ! )

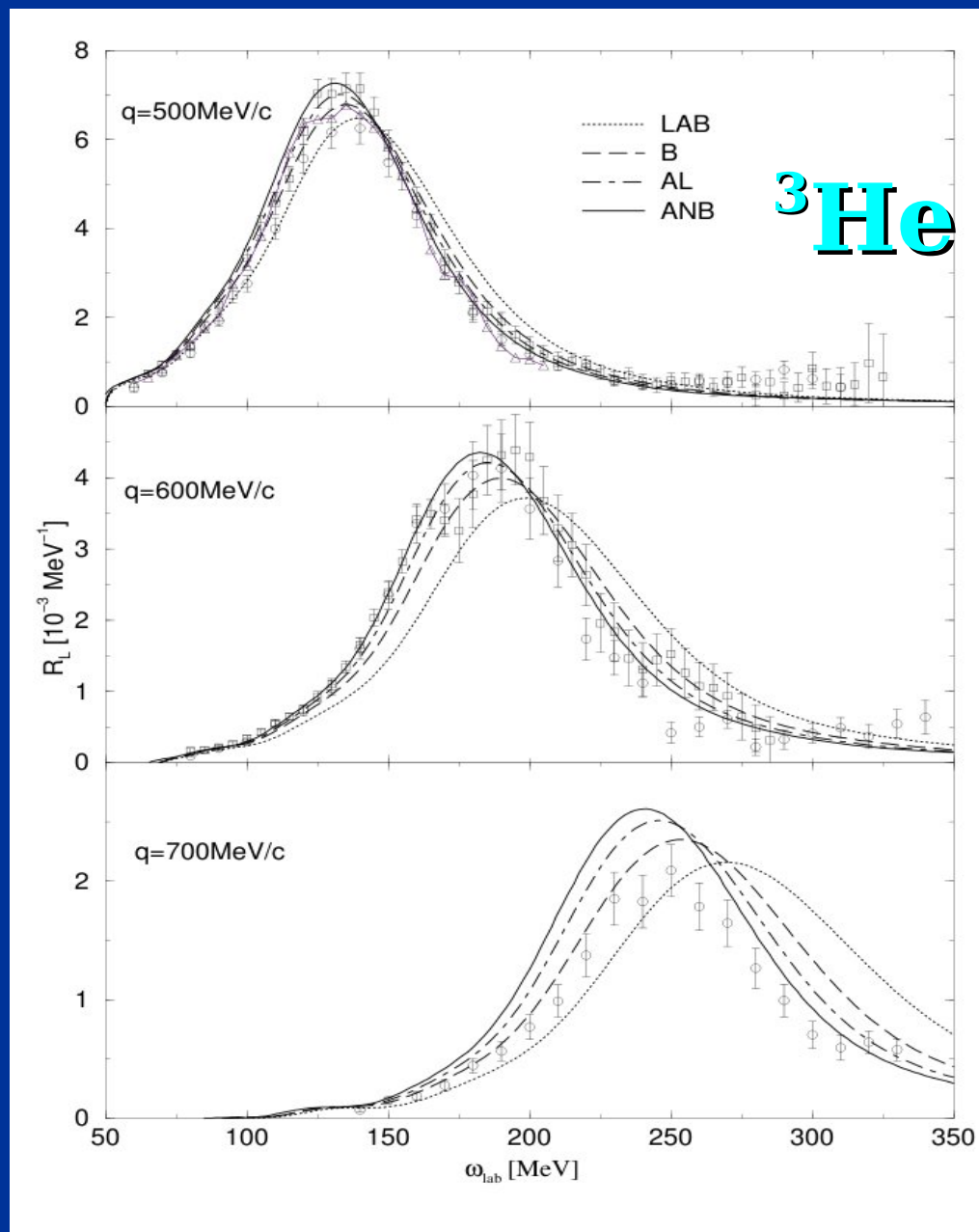
$$R_L^{\text{LAB}}(q, \omega) = \frac{q^2}{q_{fr}^2} \frac{E_i^{fr}}{M_T} R_L^{fr}(q_{fr}, \omega_{fr})$$

$$R_T^{\text{LAB}}(q, \omega) = \frac{E_i^{fr}}{M_T} R_T^{fr}(q_{fr}, \omega_{fr})$$

# Longitudinal response of $^3\text{He}$

$$R_L(q, \omega)$$

Large frame dependence!!!



V.Efros, W.Leidemann, G.O., E.L.Tomusiak PRC 72 (2005) 011002

**Is there an easy way to cure it?**



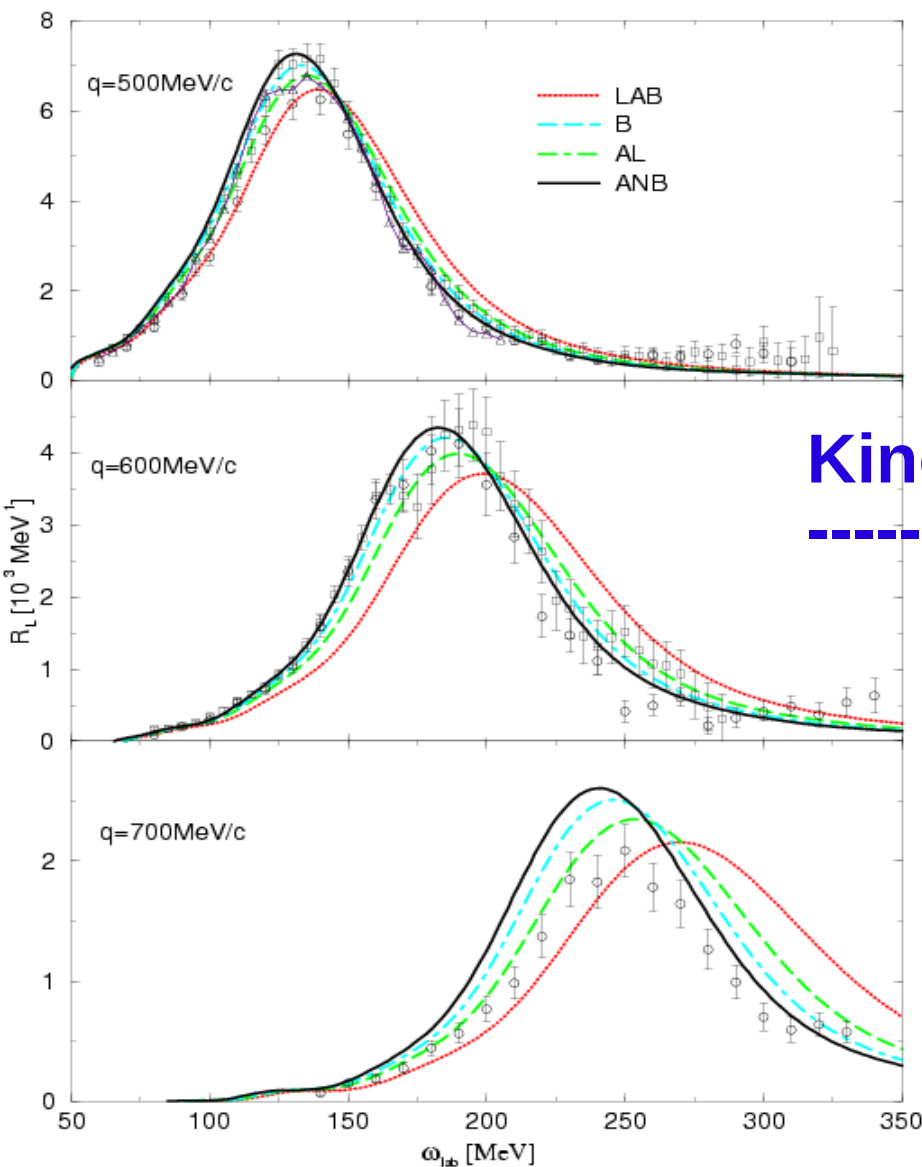
Is there an easy way to cure it?

*use in each frame the **kinematical inputs**  
corresponding to the  
**quasi elastic 2-body** assumption i.e.  
 **$p + (A-1)$ -system***

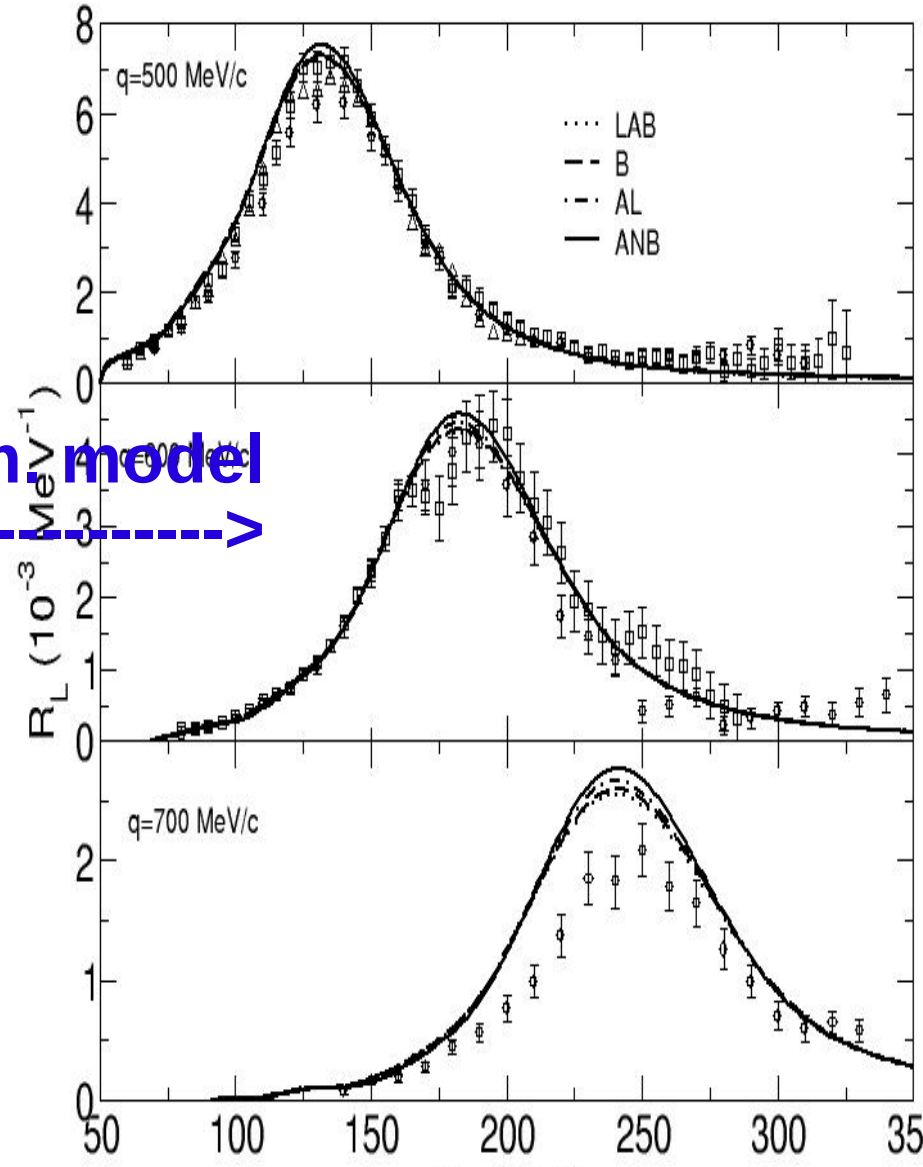
The **relative momentum**  $p_{\text{rel}}$  of the 2 bodies ( $p + (A-1)$ )  
can be calculated in each frame in a  
**relativistically correct** way.

The **energy** of the final state (the input of a **non relativistic** dynamical calculation) is then taken in its  
**non relativistic form**  $p_{\text{rel}}^2 / 2 \mu$

# Longitudinal response of $^3\text{He}$ $R_L(q, \omega)$



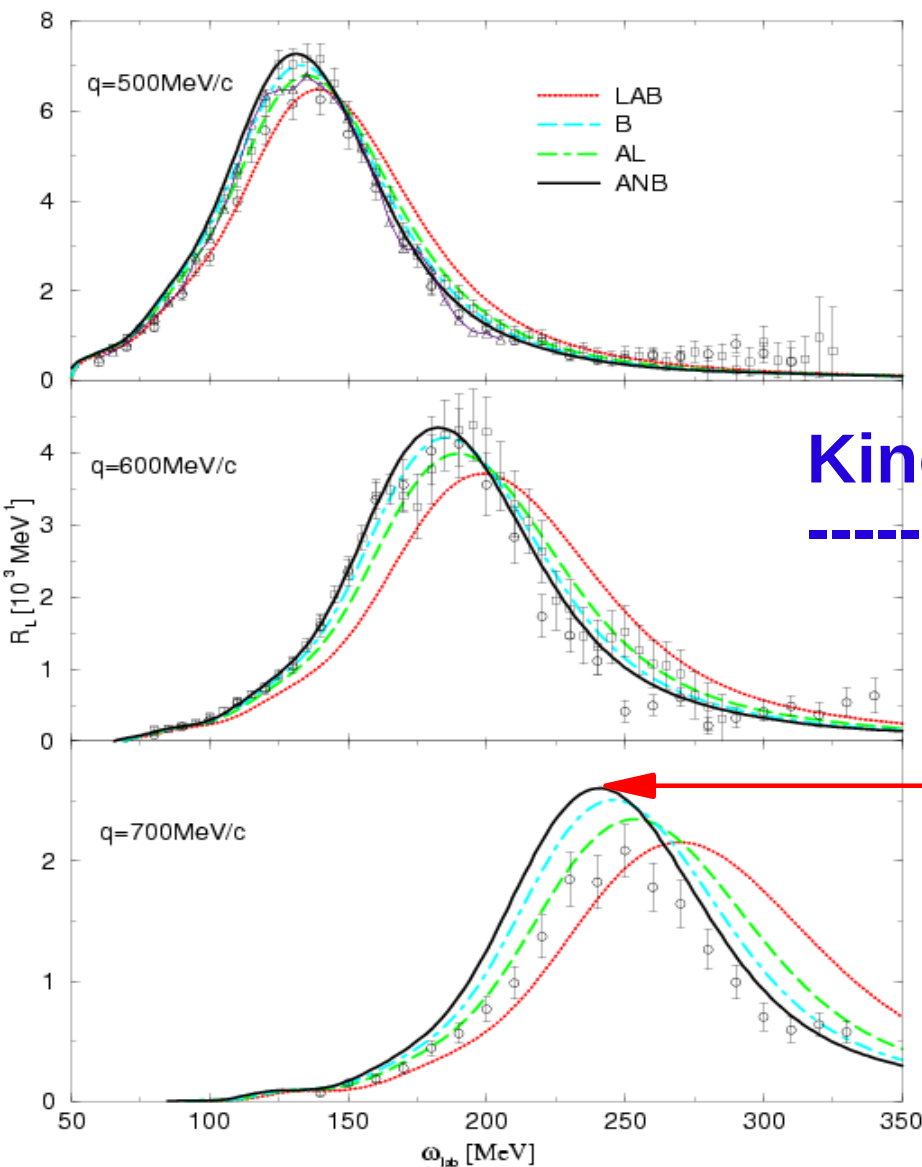
Kinem model



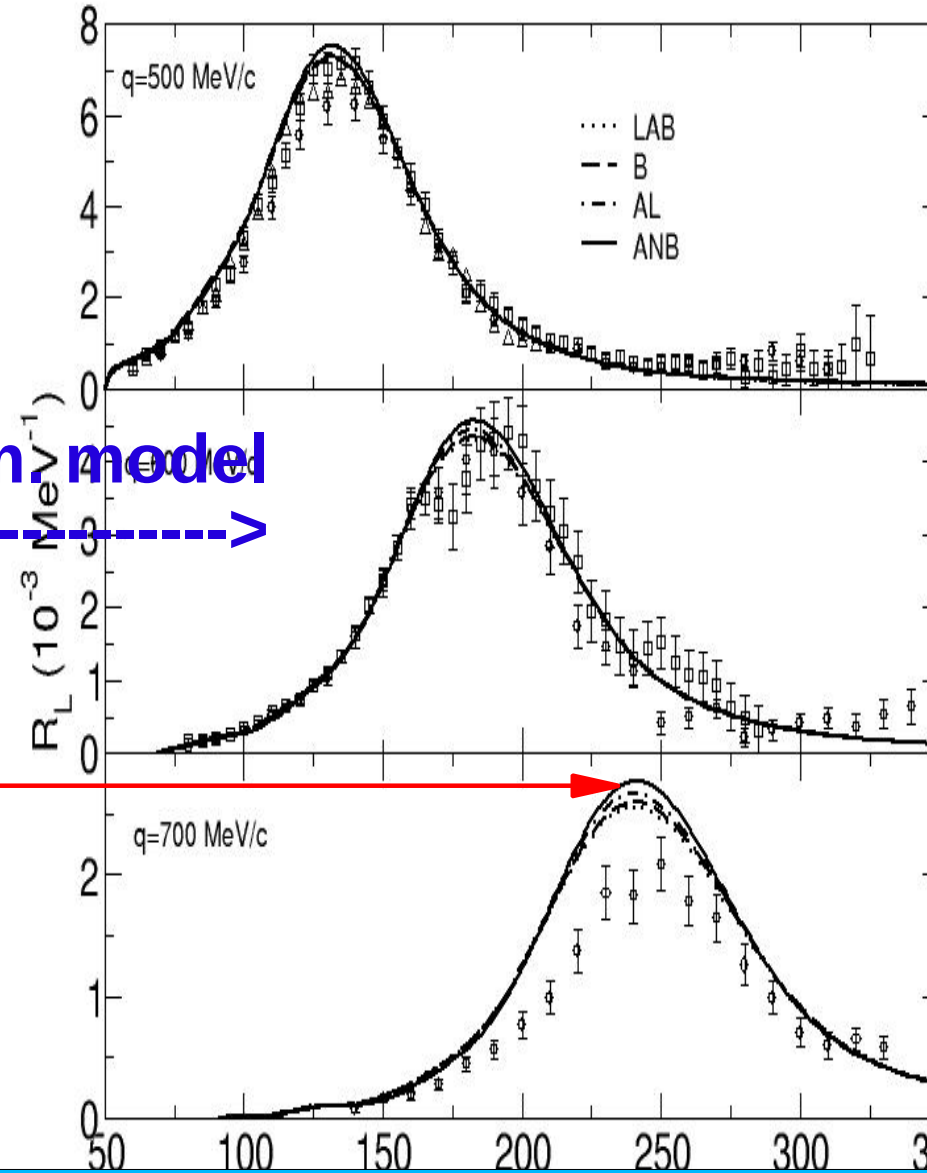
**remark:**

Of the 4 frames the **ANB** result is the **less affected** by the **relativistically correct** kinematical model.

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$$T \approx \frac{p^2}{2m} - \frac{p^4}{8m^3} + \dots \quad \frac{\Delta T}{T} \approx \frac{p^2}{4m^2}$$



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$$\text{LAB} : \frac{\Delta T}{T} \approx \frac{q^2}{4m^2} \quad \text{ANB} : \frac{\Delta T}{T} \approx \frac{q^2}{16m^2} \quad !!!$$

Moreover: the **peak position** in the **ANB** frame is always **relativistically correct**, in fact in general:

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**rel. different from n.r. !!!**

**ANB** :  $\omega_{\text{peak}} \cong T(q/2) - T(q/2) = 0$

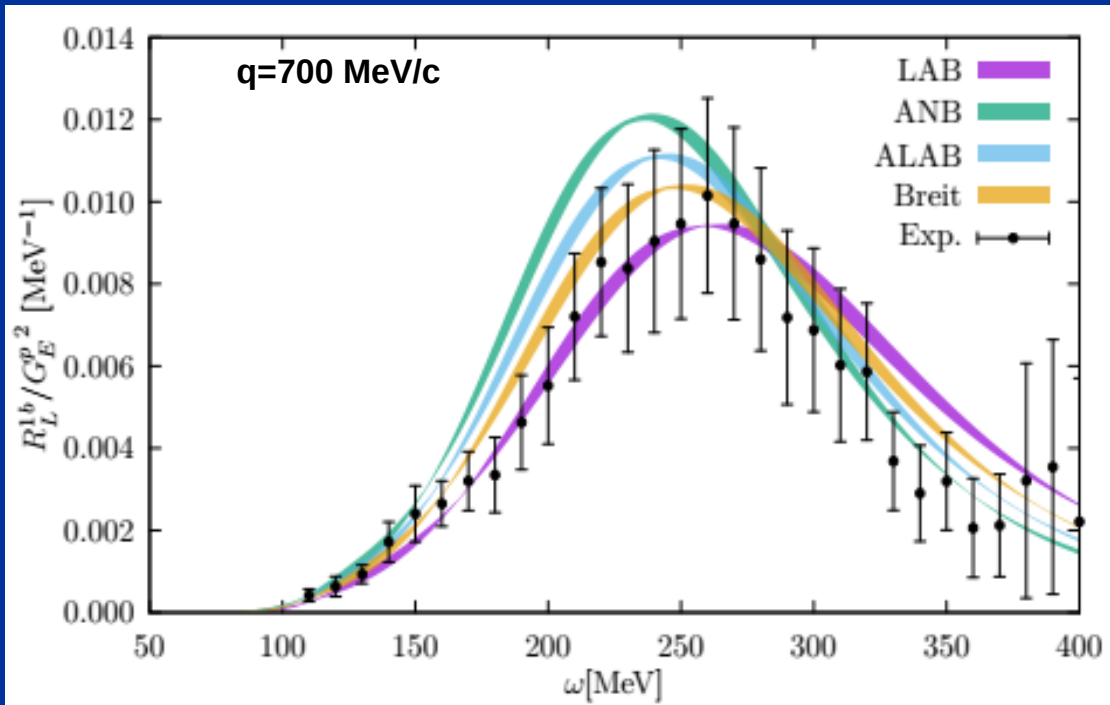
**rel. equal to n.r.  
always correct !!!**

# The test on the ${}^4\text{He}(e,e')$ cross section

# $R_L(q, \omega)$ of ${}^4\text{He}$

Integral transform calculation with

$$K(\omega, \tau) = e^{-\omega \tau} \quad \tau \text{ real, GFMC, full FSI}$$



Large frame  
dependence  
also in  ${}^4\text{He}$ !

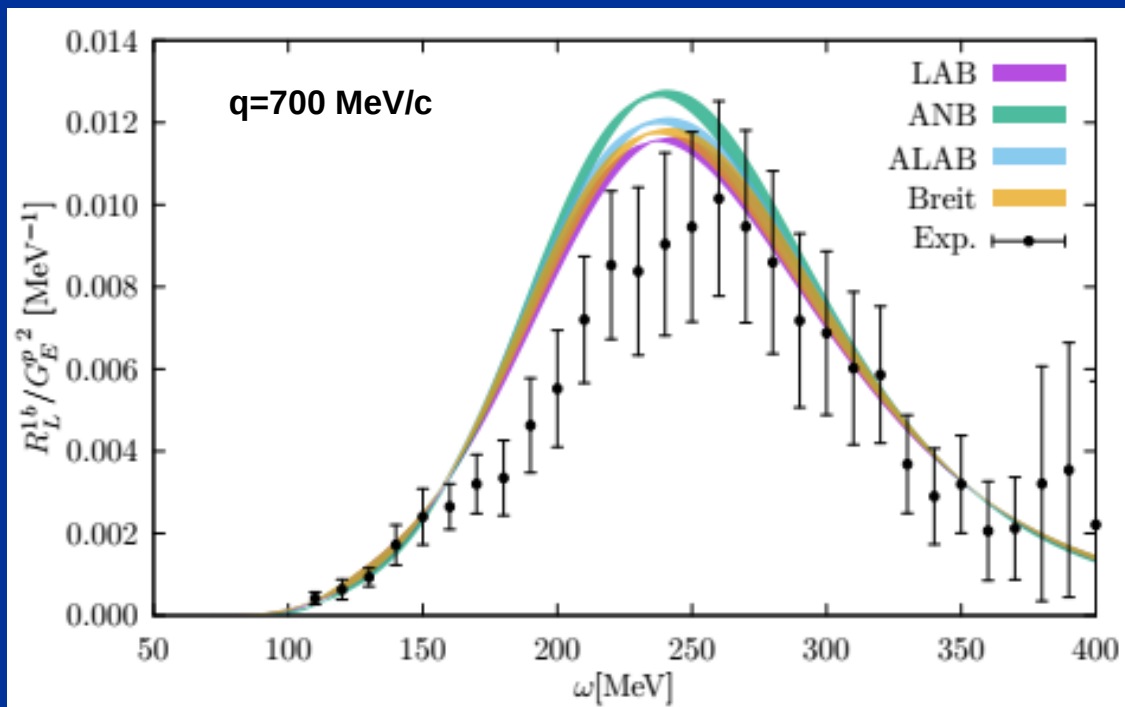
N.Rocco, W.Leidemann, A. Lovato, G.O.  
Phys. Rev. C 97, 055501 (2018)

Assuming q.e. kinematics *[2-body break-up 1-(A-1)]*  
one can treat the relativistic kinematical inputs  
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frame  
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much  
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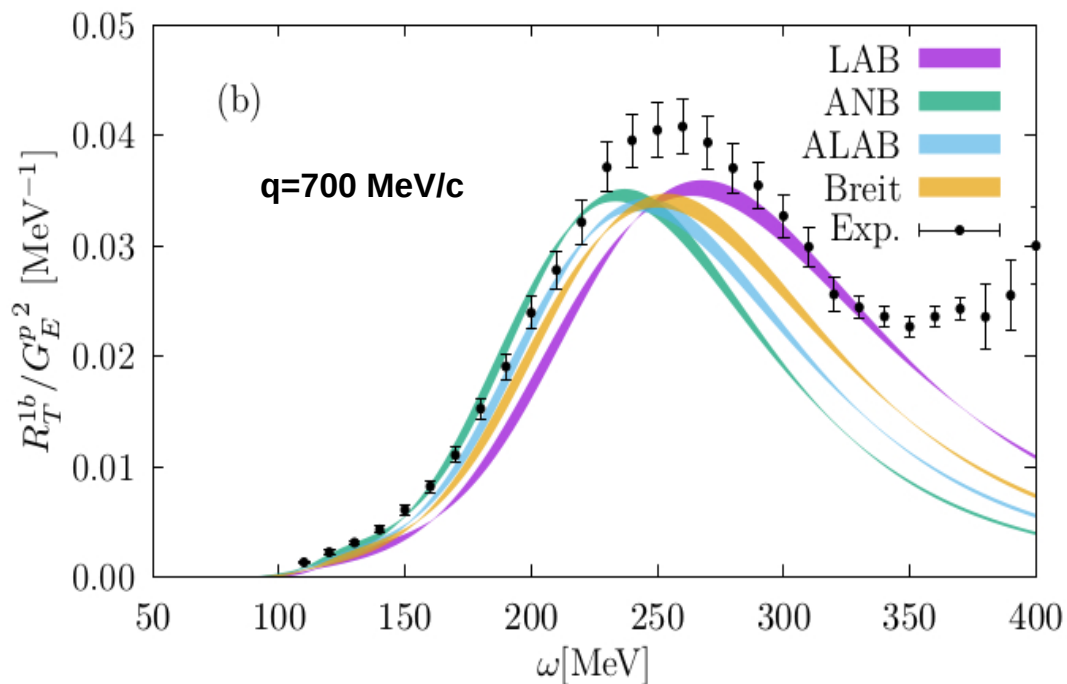
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# $R_T(q, \omega)$ of ${}^4\text{He}$

Integral transform calculation with

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Large frame  
dependence  
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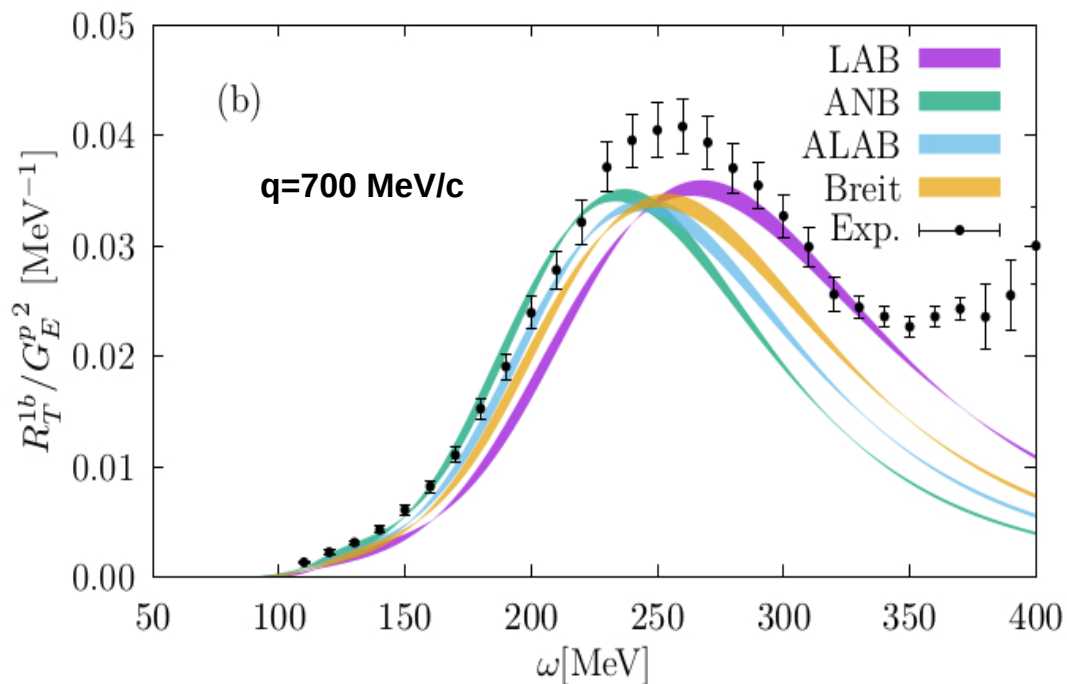
# $R_T(q, \omega)$ of ${}^4\text{He}$

Integral transform calculation with

$K(\omega, \tau) = e^{-\omega \tau}$   $\tau$  real, GFMC, **full FSI**

1-body + 2-body currents

No pion production



Large frame  
dependence  
also in  ${}^4\text{He}$ !

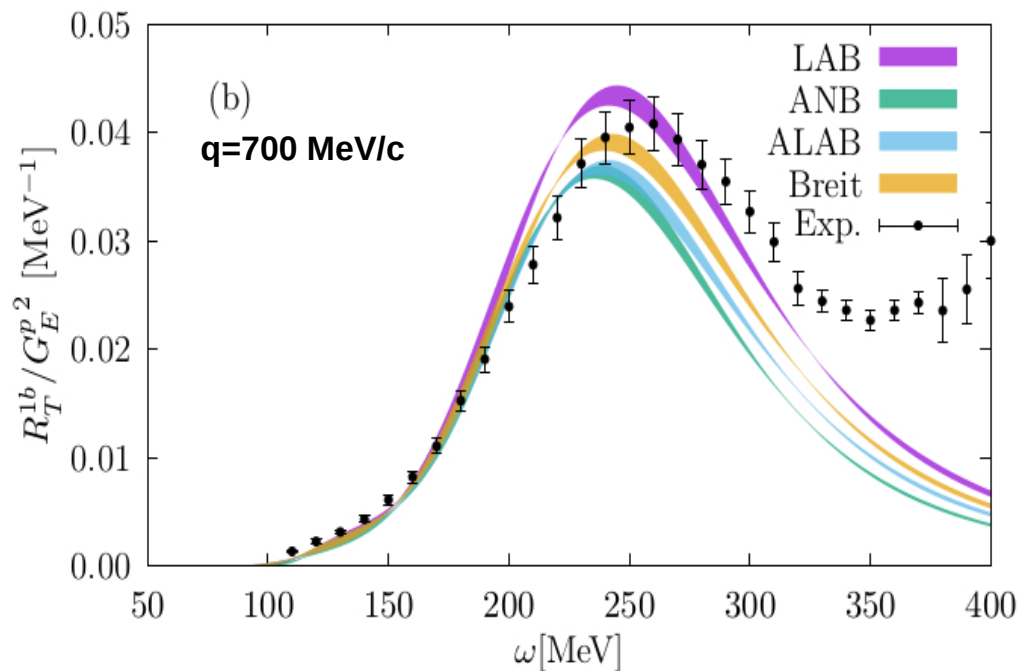
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Integral transform calculation with  
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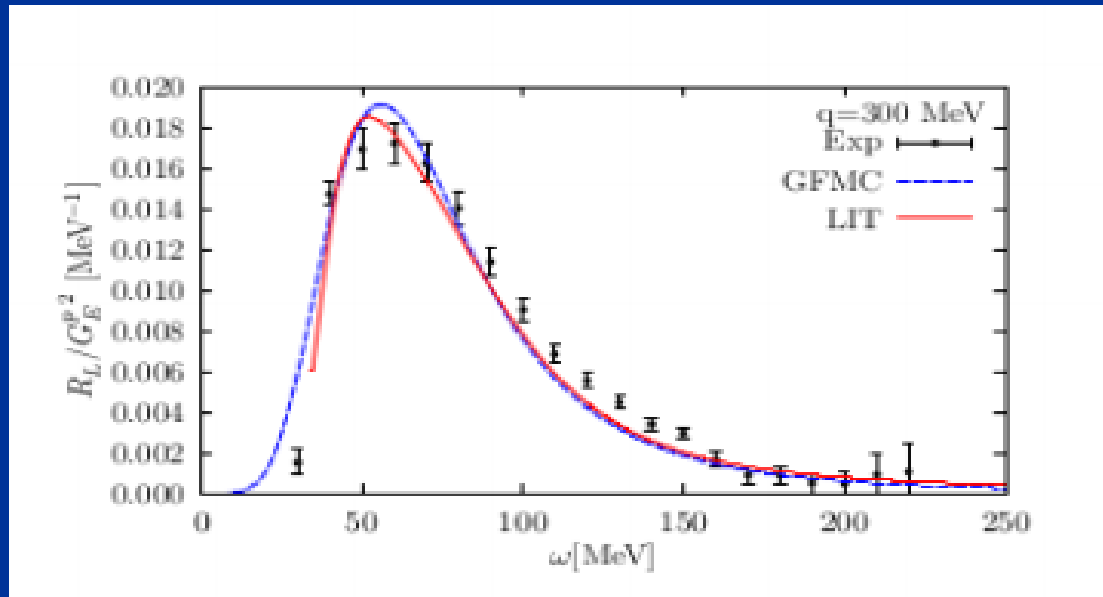
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# estimation of accuracy



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**Blue line: Exponential Kernel**

**Red line: Lorentzian Kernel:**

# results on total cross section

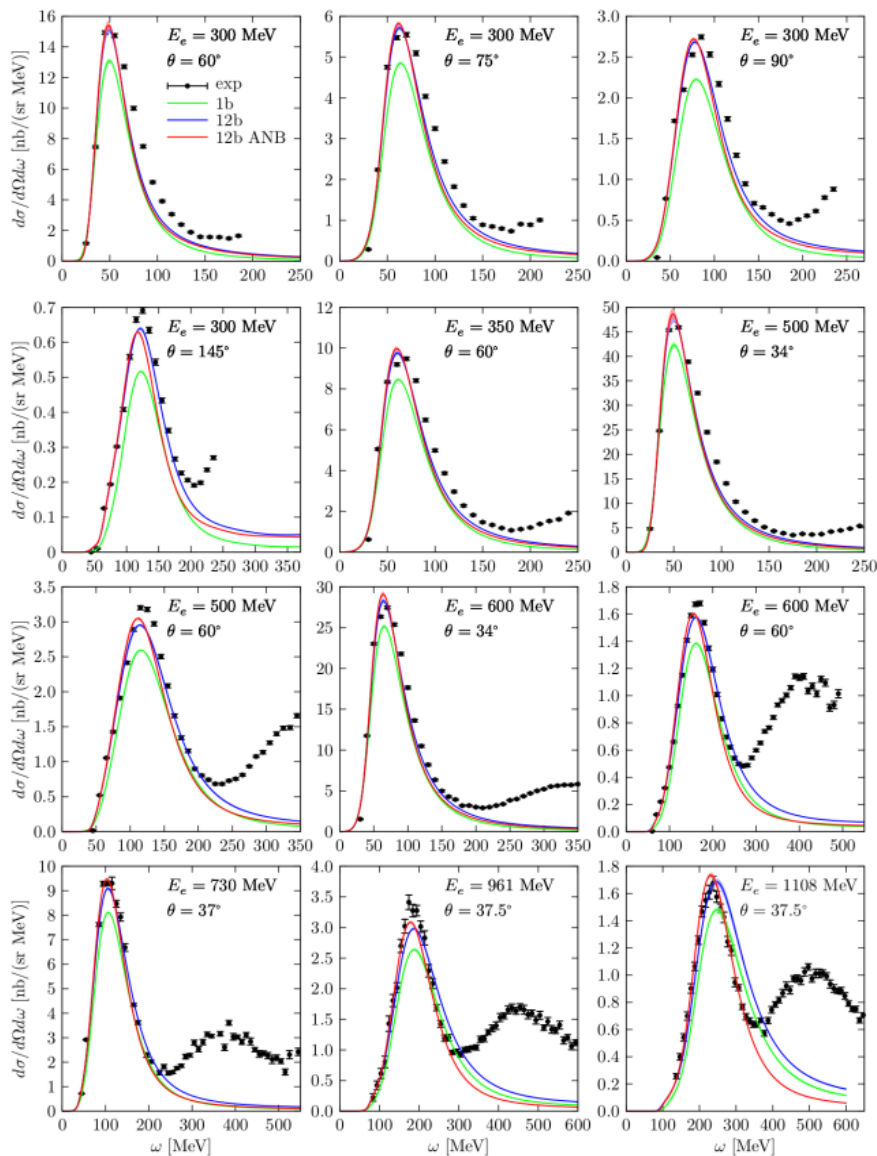


FIG. 7. Double-differential electron-<sup>4</sup>He cross sections for different values of incident electron energy and scattering angle. The green and blue lines correspond to GFMC calculation where only one- body and one- plus two-body contributions in the electromagnetic currents are accounted for. The red line indicates one plus two-body current results obtained in the ANB frame, employing the two-body fragment model to account for relativistic kinematics. The experimental data are taken from Ref. [? ].

- Many contributions at different  $q$  !
- Very high computational effort demanded
- Smart interpolation via scaling variable performed
- see

N.Rocco, W.Leidemann, A. Lovato, G.O. Phys. Rev. C 97, 055501 (2018)

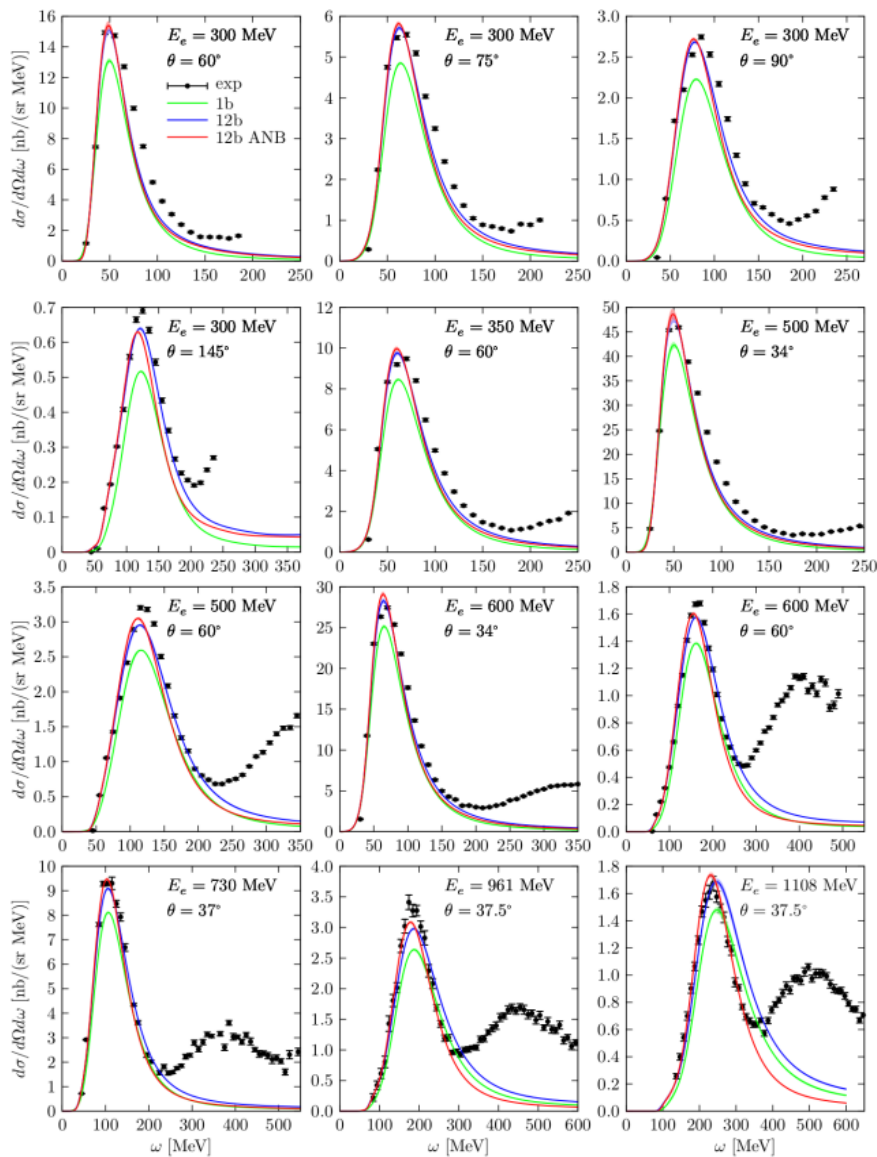


FIG. 7. Double-differential electron- $^4\text{He}$  cross sections for different values of incident electron energy and scattering angle. The green and blue lines correspond to GFMC calculation were only one- body and one- plus two-body contributions in the electromagnetic currents are accounted for. The red line indicates one plus two-body current results obtained in the ANB frame, employing the two-body fragment model to account for relativistic kinematics. The experimental data are taken from Ref. [? ].

Test on electron  
scattering data:  
Very good!

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Phys. Rev. C 97, 055501 (2018)



**One can extend to high  $q$  the applicability  
of an *ab initio* n.r. calculation  
by choosing the right frame**

# Summarizing:

- Ab initio **non relativistic** calculations of the  $(e,e')$  nuclear cross section can be performed considering the **full realistic potential dynamics** both in the initial and in the **final** state

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- **One and two-body currents** (with relativistic corrections up to  $(q/m)^2$  are included
- **Frame dependence** is much reduced, using quasi elastic kinematics
- In q.e. regime relativistic effects are minimized in the **ANB frame**

# Conclusion and outlook

- The **test** of the described approach on **(e,e')** measured cross section turns out to be **very good**
- Then one can use the same approach for **neutrino** scattering
- Heavier targets than  $^4\text{He}$  can also be treated
- Work on  $^{12}\text{C}$  is in progress

# Results obtained with

- **Noemi Rocco** (Argonne Nat. Lab.)
- Alessandro Lovato (INFN Trento)
- Winfried Leidemann (Univ. Trento)
- Victor Efros (Kurchatov Centre Moscow)
- Ed Tomusiak (Univ. Victoria Canada)