Charged pion masses under strong magnetic fields in the NJL model

# CONICET

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*M. Coppola, D. Gomez Dumm and N. N. Scoccola, PLB in press [arXiv:***1802.08041** [hep-ph]]. arXiv.org > hep-ph > arXiv:1802.08041

High Energy Physics - Phenomenology

Charged pion masses under strong magnetic fields in the NJL model

M. Coppola, D. Gomez Dumm, N.N. Scoccola (Submitted on 22 Feb 2018)

### MOTIVATION

> Strong Magnetic Fields: compact stars, heavy ion collisions

> In NJL, calculations for  $m_{\pi^0}$  but not for  $m_{\pi^{\pm}}$  (not <u>exact</u>)

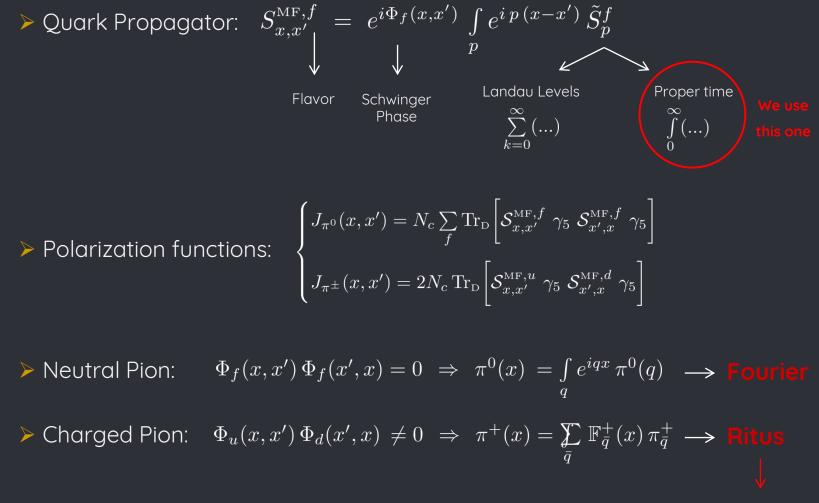
> LQCD calculations for  $m_{\pi^+}$  for small eB<0.4 GeV<sup>2</sup> (2013)

> Recent LQCD results (2017) for both  $m_{\pi}$  for large eB (<4 GeV<sup>2</sup>) with unrealistic  $m_{\pi}(B=0) = 415 \text{ MeV}$ 

> <u>Neutral Pion</u>: Recent NJL calculation (2017) in full RPA with **G(eB)** in Fourier basis that agrees with LQCD but for a set with  $f_{\pi} \sim 80 \text{ MeV}$  in vacuum

> <u>Charged Pion</u>: Schwinger phases do not cancel, no translational invariance, can not use Fourier basis. Recent studies neglect this phases, or use a derivative expansion approximation (inexact for large eB or  $m_{\pi^{\pm}}$ )

### THEORETICAL KEYS (local NJL with B)



Regularization: MFIR scheme with 3D cutoff

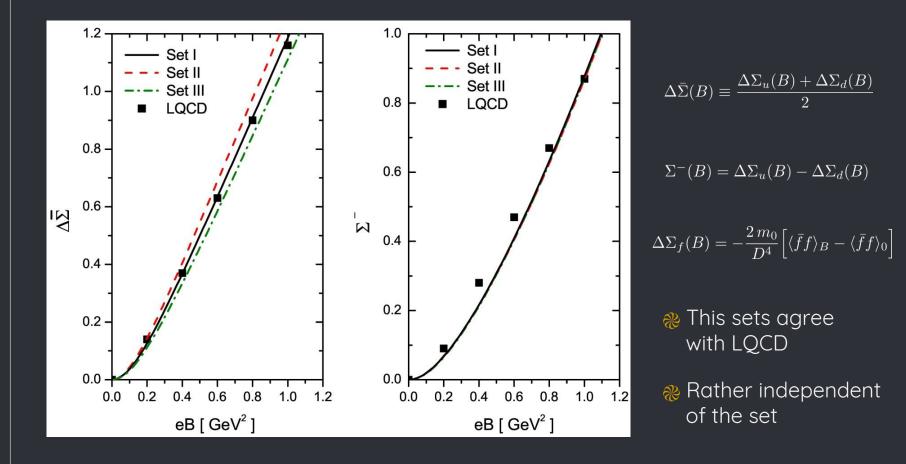
In each basis each  $J_{\pi}(q,q')$ is diagonalized

### RESULTS – QUARK CONDENSATE

JL

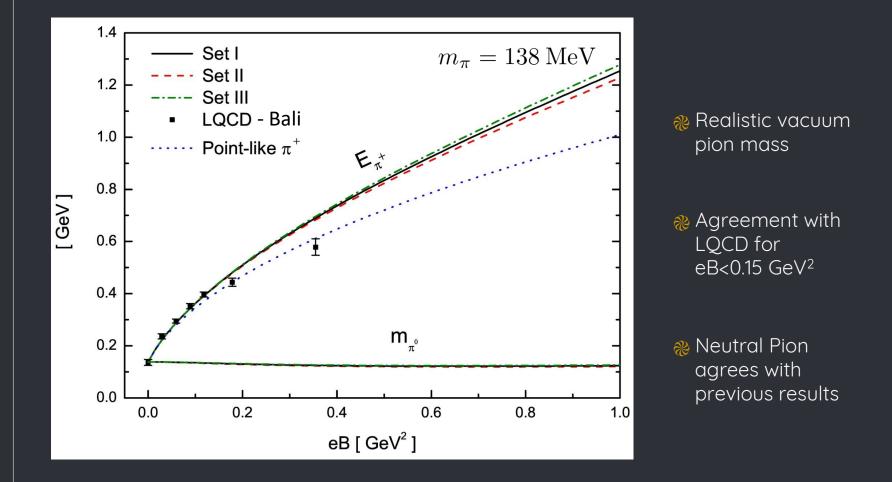
> Three Parameters Sets:  $M_0 = 350, 320, 380, G = cte, \begin{cases} m_\pi = 138 \text{ MeV} \\ f_\pi = 92.4 \text{ MeV} \end{cases}$ 

> In Avancini '17 they use G(eB) to fit LQCD, G=*const.* does not agree



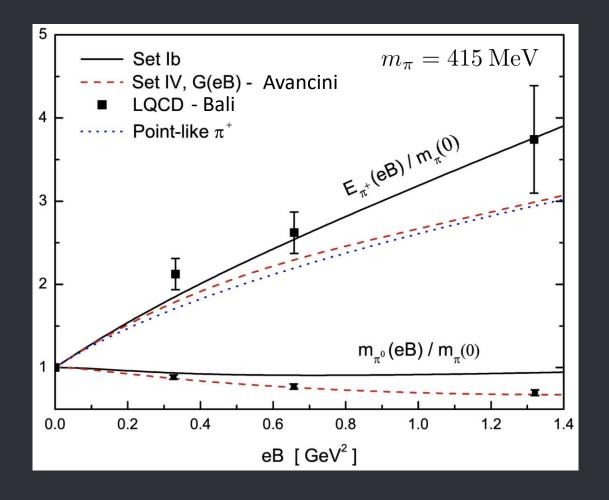
# RESULTS – PION MASSES (138 MeV)

> Instead of 
$$m_{\pi^+}$$
:  $E_{\pi^+}(eB) = \sqrt{m_{\pi^+}^2 + (2k+1)eB + q_3^2} \Big|_{\substack{q_3=0\\k=0}} = \sqrt{m_{\pi^+}^2 + eB}$ 





> LQCD results for unrealistic  $m_{\pi}(B=0) = 415 \text{ MeV}$ > Set Ib = Set I with m<sub>0</sub> increased, as in Set IV of *AVANCINI* '17



<u>Neutral Pion</u> & Low error bars in LQCD Set Ib above LQCD Set IV fitted to LQCD but with  $f_{\pi} \simeq 80 \text{ MeV}$ 

#### Charged Pion

≈ High error bars in LQCD

- Set Ib fits results from spacing extrapolation to the continuum
- Set IV fits results from finite lattice spacing



<u>Quark Condensate</u>: we found Sets with G=*const.* that agree with LQCD

Neutral Pion: Set IV with G(eB) fits LQCD ( $m_{\pi} = 415 \text{ MeV}$ ). Set Ib with G=*const.* lies above LQCD for large eB. We verified that using proper time or Landau levels (as in *Avancini* 17) for  $\tilde{S}_p^f$  lead to the same results.

> <u>Charged Pion</u>: mass results agree with LQCD for eB<0.15 GeV<sup>2</sup>, deviation for larger eB. Normalized mass agrees for Set Ib ( $m_{\pi} = 415 \text{ MeV}$ ), but error bars of LQCD are too large to be conclusive.

> Development of the Ritus method for charged pions. Fully takes into account the translational-breaking effect of the Schwinger phases. Can calculate  $f_{\pi}(eB)$  or  $\Gamma_{\pi}(eB)$ , pion properties at finite T, etc. Thanks!



### Quark propagator

$$\begin{split} \tilde{S}_{p}^{f} &= \int_{0}^{\infty} d\tau \, \exp\left[-\tau \left(M^{2} + p_{\parallel}^{2} + p_{\perp}^{2} \frac{\tanh \tau B_{f}}{\tau B_{f}}\right)\right] \times & \text{Proper} \\ & \left\{\left(M - p_{\parallel} \cdot \gamma_{\parallel}\right) \left[1 + is_{f} \, \gamma_{1} \gamma_{2} \, \tanh \tau B_{f}\right] - \frac{p_{\perp} \cdot \gamma_{\perp}}{\cosh^{2} \tau B_{f}}\right\} & \text{time} \\ \tilde{S}_{p}^{f} &= 2 \, \exp(-p_{\perp}^{2}/|q_{f}B|) \sum_{k=0}^{\infty} \sum_{\lambda=\pm} \frac{1}{d_{k,p_{\parallel}}^{f}} \times & \text{Sum over Landau Levels} \\ & \left[(-1)^{k_{\lambda}} \left(M - p_{\parallel} \cdot \gamma_{\parallel}\right) L_{k_{\lambda}} (2p_{\perp}^{2}/B_{f}) + 2 \, (-1)^{k} \, p_{\perp} \cdot \gamma_{\perp} \, L_{k-1}^{1} (2p_{\perp}^{2}/B_{f})\right] \Delta \end{split}$$

### > Polarization Functions in coordinate space

$$S_{\text{bos}}^{\text{quad}} = \frac{1}{2} \int_{x,x'} \delta\pi(x)^* \left[ \frac{1}{2G} \,\delta^{(4)}(x-x') - J_{\pi}(x,x') \right] \delta\pi(x')$$
$$J_{\pi^0}(x,x') = N_c \sum_f \text{Tr}_{\text{D}} \left[ \mathcal{S}_{x,x'}^{\text{MF},f} \,\gamma_5 \,\mathcal{S}_{x',x}^{\text{MF},f} \,\gamma_5 \right], \quad J_{\pi^{\pm}}(x,x') = 2N_c \,\text{Tr}_{\text{D}} \left[ \mathcal{S}_{x,x'}^{\text{MF},u} \,\gamma_5 \,\mathcal{S}_{x',x}^{\text{MF},d} \,\gamma_5 \right]$$



> Polarization Functions diagonalized in momentum space

$$\begin{split} S_{\pi^{0}}^{\text{quad}} &= \frac{1}{2} \int_{q} \delta \pi^{0}(-q) \, \left[ \frac{1}{2G} - J_{\pi^{0}}(q_{\perp}^{2}, q_{\parallel}^{2}) \right] \delta \pi^{0}(q) & \longrightarrow \quad \frac{1}{2G} - J_{\pi^{0}}^{(reg)}(0, -m_{\pi^{0}}^{2}) = 0 \\ & \downarrow \\ q_{\perp}^{2} &= 0, \ q_{\parallel}^{2} = m_{\pi^{0}}^{2} \end{split}$$

$$S_{\pi^{+}}^{\text{quad}} = \frac{1}{2} \sum_{\bar{q}} (\delta \pi_{\bar{q}}^{+})^{*} \left[ \frac{1}{2G} - J_{\pi^{+}}(k, \Pi^{2}) \right] \delta \pi_{\bar{q}}^{+} \longrightarrow \frac{1}{2G} - J_{\pi^{+}}^{(\text{reg})}(k, -m_{\pi^{+}}^{2}) = 0$$

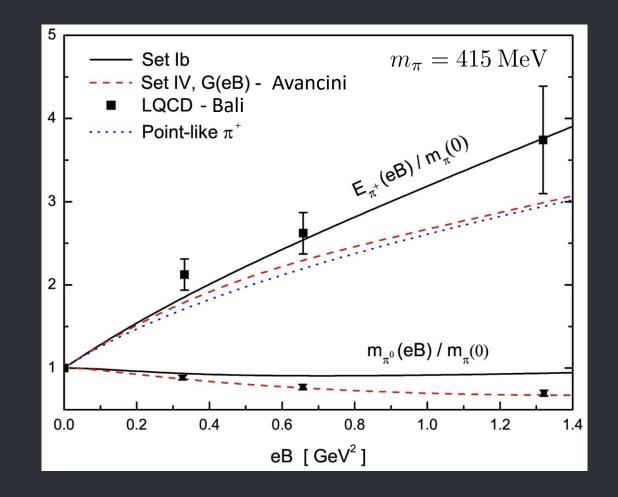
$$\downarrow$$

$$\Pi^{2} = (2k+1) B_{\pi^{+}} + q_{\parallel}^{2} = m_{\pi^{+}}^{2}$$

## NORMALIZED PION MASSES (415 MeV)

JL

& LQCD results from lattice spacing extrapolation to the continuum (the neutral pion fits finite lattice spacing)



### NORMALIZED PION MASSES (415 MeV)

JL

& LQCD results for the meson masses from the lines of constant physics

