

High Precision Electro-Weak Quantum Loops: Standard Model predictions for the top-quark and Brout-Englert-Higgs scalar masses, later confirmed by their discovery, contributed directly to the awards of Two Nobel Prizes in Physics

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The Standard Model (SM), augmented by massive neutrinos and classical general relativity, has never failed a legitimate experimental or observational test. The stakes for LEP, SLC, Tevatron and LHC were immense! Either the SM's reliable high-accuracy predictions for the top-quark and Brout-Englert-Higgs scalar (BEH) masses were true, and experimentally/observationally verified, or we had to change this most basic SM paradigm. Instead, with the experimental discovery of the top-quark in 1995 and the BEH scalar in 2012, the High Precision Electro-Weak Quantum Loop (HP-EW-QL) structure of the SM was completely vindicated.

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I. THE CORE THEORY IS VINDICATED

The Standard Model [162–164], augmented with right-handed neutrinos and neutrino masses (the Minimal νSM [165–167]) and classical general relativity, is the most powerful, accurate, predictive, successful and experimentally/observationally verified theory known to science. It has never failed a legitimate experimental or observational test. Frank Wilczek [355] calls it the “Core Theory”.

With discovery of the Brout-Englert-Higgs (BEH) scalar (aka “Higgs”), the SM has now been experimentally completely vindicated. Most persuasive is its renormalizable/unitary quantum loop structure, tested with High-Precision Electro-Weak Quantum Loops (HP-EW-QL), notably at CERN's LEP and SLAC's SLC e^+e^- colliders with center-of-mass energy on/near the Z -pole. The most powerful HP-EW-QL tests are:

- Successful 1984 Z -pole physics 1-loop HP-EW-QL prediction for the top-quark mass [670, 682] led to the award of the 1999 Nobel Prize in Physics for the renormalizability/unitarity of spontaneously broken Yang Mills gauge theories [367];
- Successful 1984 Z -pole physics (1-loop [670, 682], augmented by certain 2-loop [858]) HP-EW-QL prediction for m_{BEH} . The eventual ultra-high-accuracy July 2011 prediction, in Figure 4, led to its July 4, 2012 LHC discovery at $m_{BEH} = 125$ GeV, and contributed to the award of the 2013 Nobel Prize in Physics [368].

- 2-loop prediction for the W^\pm mass [175, 386, 625, 670, 682, 858], also contributed to the successful predictions of the top-quark and BEH masses, and the associated 1999 and 2013 Nobel Prizes [367, 368].
- The 1-loop predictions for the W^\pm and Z^0 masses, using neutrino-hadron data, contributed to their 1983 discovery, and the 1984 Nobel Prize in Physics [176]
- The classification [178, 354, 369, 676, 740, 746, 779?] of “oblique” HP-EW-QL has also experimentally ruled out vast swaths of “Beyond the Standard Model (BSM)” theoretical high energy physics (HEP).

Renormalizability, unitarity, infra-red finiteness and practical calculability for the Standard Model was conceived/invented and brought to mathematical and computational maturity (by 1979) by the seminal work of: G. 't Hooft, M.J.G. Veltman, L.D.. Fadeev, V.N. Popov, J. Goldstone, Y. Nambu, A. Salam, S. Weinberg, J.C. Ward, J.C. Taylor, A.A. Slavnov, C. Becchi, A. Rouet, R. Stora, I.V. Tyutin, D. Yennie, S. Frautschi, A. Surra, F. Behrends, T. Kinoshita, A. Sirlin, G. Passarino and M. Consoli.

The theoretical part of the field is now widely regarded as having been completed [155]: i.e. as a scientific and mathematical field of theoretical discovery.

But, since the 1-loop and certain 2-loop quantum structure was directly and reliably derived from the SM, the experimental stakes in 1994 were immense: either the Standard Model's reliable predictions for m_{top}, m_{BEH} were true (and eventually verified), or we had to change this most basic paradigm. Instead, with the discovery of the top-quark in 1995, and the SM BEH scalar in 2012, the quantum-loop structure of the SM was completely vindicated.

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HP-EW-QL is now so conclusively experimentally validated that it is widely regarded as an engineering and technology skill, used, for example, to calculate LHC signals and backgrounds.

But there are those of us who believe that the Standard Model still has a few tricks [381] up its sleeve: e.g. new mathematical symmetry properties [382–385, 660–663, 769, 770, 948–952] and game-changing scientific discoveries [771–776].

II. OBLIQUE HP-EW-QL

The 2-point A (photon) and Z self-energies, $Z-A$ mixing and W^\pm self-energy 1-particle-irreducible connected truncated Green's functions are first defined, by separating out gauge coupling constants, in term of the $SU(2)_L$ isospin current \vec{J}^μ and electromagnetic current J_Q^μ :

$$\begin{aligned}\Pi_{AA}(q^2) &= q^2 e_*^2 \Pi'_{QQ} \\ \Pi_{ZA}(q^2) &= \frac{e_*^2}{s_* c_*} [\Pi_{3Q} - q^2 s_*^2 \Pi'_{3Q}] \\ \Pi_{3Q}(q^2) &= q^2 \Pi'_{3Q} + \Pi_{3Q}^L \\ \Pi_{ZZ}(q^2) &= \frac{e_*^2}{s_*^2 c_*^2} [\Pi_{33} - 2s_*^2 \Pi_{3Q} + q^2 s_*^4 \Pi'_{3Q}] \\ \Pi_{WW}(q^2) &= \frac{e_*^2}{s_*^2} \Pi_{+-}\end{aligned}\tag{1}$$

We suppress Lorentz indices, for the moment, by writing the 2-point functions in Landau gauge. The running $*$ couplings are defined below.

A. Gauge-independent running couplings

When the more general scalar sector considered in [740] is specialized to a single BEH doublet

$$\begin{aligned}\frac{1}{e_*^2(q^2)} &= \frac{1}{e_*^2(\mu^2)} - \text{Re}\left(\Pi'_{QQ} + 2\Gamma'(\xi)\right)_{q^2} \\ &\quad + \text{Re}\left(\Pi'_{QQ} + 2\Gamma'(\xi)\right)_{q^2=\mu^2} \\ \frac{1}{4\sqrt{2}G_\mu^*(q^2)} &= \frac{1}{4\sqrt{2}G_\mu^*(\mu^2)} \\ &\quad + \text{Re}\left(-\Pi_{+-} + q^2 \Pi'_{3Q} + 2\Pi_{3Q}^L(\xi)\right)_{q^2} \\ &\quad - \text{Re}\left(-\Pi_{+-} + q^2 \Pi'_{3Q} + 2\Pi_{3Q}^L(\xi)\right)_{q^2=\mu^2} \\ \frac{1}{\rho^*(q^2)} &= 1 - 4\sqrt{2}G_\mu^*(q^2) \text{Re}\left(\Pi_{33} - \Pi_{+-}\right)_{q^2}\end{aligned}$$

$$\begin{aligned}\frac{1}{g_*'^2(q^2)} &= \frac{1}{g_*'^2(\mu^2)} - \text{Re}\left(\Pi'_{QQ} - \Pi'_{3Q}\right)_{q^2} \\ &\quad + \text{Re}\left(\Pi'_{QQ} - \Pi'_{3Q}\right)_{q^2=\mu^2}\end{aligned}\tag{2}$$

where μ^2 indicate an arbitrary set of 3 renormalization points. Written this way, the 1-loop running couplings are manifestly independent of renormalization scheme. With the further definitions

$$\begin{aligned}e_*^2 &\equiv \frac{g_*^2 g_*'^2}{g_*^2 + g_*'^2} \equiv g_*^2 s_*^2 \equiv g_*'^2 c_*^2 \\ s_*^2 + c_*^2 &= 1\end{aligned}\tag{3}$$

$s_*^2(q^2)$ is the \sin^2 of the running mixing angle. The leading UV logs in (2,3) match those of the renormalization group beta functions, but also include all of the analytic and infra-red terms from the vector self-energies.

Gauge-independence of running couplings:

Although the transverse vector 2-point functions Π'_{QQ}, Π_{3Q} are gauge-independent, the longitudinal parts of massive W^\pm, Z and Z -photon mixing 2-point functions, $\Pi_{+-}, \Pi_{33}, \Pi_{3Q}^L$, are gauge-dependent. The gauge-dependent vertex-function $\Gamma'(\xi; q^2)$ and longitudinal $\Pi_{3Q}^L(\xi; q^2)$ are related (with oblique 1-loop improvements)

$$\Pi_{3Q}^L = -\frac{e_*^2}{4\sqrt{2}G_\mu^* s_*^2} \Gamma'\tag{4}$$

and are **universal**: i.e. they do not depend on the external particles of the particular S-Matrix element being considered. Their job is to cancel the longitudinal parts of the Z, W^\pm and $Z-A$ mixing 2-point functions, thus ensuring manifest gauge-independence [740] of the effective running couplings in G. 't Hooft's R_ξ gauges. G. Degrandi and A. Sirlin [391], and independently M. Kuroda, G. Moulhaka and D. Schildknecht [5], (much) later confirmed the gauge-independence of the $*$ -scheme.

B. Improved Born Approximation: Standard Model high precision electroweak oblique quantum loops, for 4-fermion processes which are 1-gauge-boson reducible by cutting a massive W^\pm a massive Z^0 , or a massless photon

Let I_3 and Q be the initial weak isospin and electric charge, I_3' and Q' the final. Let I_+ and I_- be the charge-raising and charge-lowering operators, respectively. [740] shows that (apart from certain non-trivial imaginary parts which we suppress for pedagogical reasons [740]) the neutral-current and charged-current matrix elements are:

$$\mathcal{M}_{NC}^*(q^2) = \frac{e_*^2 Q Q'}{q^2 [1 - i \text{Im} \Pi_{AA}^*]}\tag{5}$$

$$\begin{aligned} & + \frac{e_*^2}{s_*^2 c_*^2} \frac{[I_3 - (s_*^2)Q][I'_3 - (s_*^2)Q']}{q^2 + \frac{e_*^2}{s_*^2 c_*^2} \frac{1}{4\sqrt{2}G_\mu^* \rho^*} - is \left(\frac{\Gamma_Z^*}{\sqrt{s}} \right)} \\ \mathcal{M}_{CC}^*(q^2) & = \frac{e_*^2}{2s_*^2} \frac{I_+ I'_- + I_- I'_+}{q^2 + \frac{e_*^2}{s_*^2 c_*^2} \frac{1}{4\sqrt{2}G_\mu^*} - is \left(\frac{\Gamma_W^*}{\sqrt{s}} \right)} \end{aligned}$$

where the s-dependent Z total width ¹, and W^\pm total width (Hollik105)

$$\begin{aligned} s \left(\frac{\Gamma_Z^*}{\sqrt{s}} \right) & = \frac{e_*^2}{s_*^2 c_*^2} \text{Im} [\Pi_{33} - 2s_*^2 \Pi_{3Q} + q^2 s_*^4 \Pi'_{QQ}] \\ s \left(\frac{\Gamma_W^*}{\sqrt{s}} \right) & = \frac{e_*^2}{s_*^2} \text{Im} \Pi_{+-} \end{aligned} \quad (6)$$

are easily broken up into the various partial widths, and match the partial widths generated by the numerators in (5).

The effective neutral and charged current 1-loop **“improved Born approximation”** [740] 4-fermion matrix elements (2,3,5,8), (i.e. supplemented by certain non-trivial residual resummed 1-loop imaginary contributions [740]) obey the optical theorem for all 4-fermion t-channel processes, as well as for s-channel $|q^2|$ processes below the $t\bar{t}$, $t\bar{b}$, $b\bar{t}$, W^+W^- , ZZ , $W^\pm H$ and ZH mass thresholds.

$$\begin{aligned} 2\text{Im} [\mathcal{M}_{NC}^*] & = \mathcal{M}_{NC}^{*\dagger} \mathcal{M}_{NC}^* \\ 2\text{Im} [\mathcal{M}_{CC}^*] & = \mathcal{M}_{CC}^{*\dagger} \mathcal{M}_{CC}^* \end{aligned} \quad (7)$$

That is, of course, as it should be, and first established [740] the correctness of the approximate Breit-Wigner Z, W^\pm line-shapes, with s-dependent Z, W^\pm widths. But (5)’s obedience of the optical theorem also ensures accurate neutrino counting on the Z peak, as well as the correct partial-width behavior there.

$$\begin{aligned} \mathcal{M}_{NC}^*(-M_Z^2) & = \frac{i[I_3 - s_*^2 Q][I'_3 - s_*^2 Q']}{\text{Im} [\Pi_{33} - 2s_*^2 \Pi_{3Q} + q^2 s_*^4 \Pi'_{QQ}]} \\ \sigma_{NC}^*(-M_Z^2) & \sim \frac{1}{M_Z^2} \frac{\Gamma_{InitialFermion}^*}{\Gamma_{Total}^*} \frac{\Gamma_{FinalFermion}^*}{\Gamma_{Total}^*} \\ \mathcal{M}_{CC}^*(-M_W^2) & = \frac{i}{2} \frac{I_+ I'_- + I_- I'_+}{\text{Im} \Pi_{+-}} \end{aligned} \quad (8)$$

¹ The total Z width Z , QCD correction for the light quarks, with 69 calculated up to third order in s , except for the mb-dependent singlet terms, which are known to $O(2s)70,71$. For a review of the QCD corrections to the Z width, see 72. QCD corrections were first derived for the leading term of $O(s \text{ G m}2t)$ 73 and were subsequently completed by the $O(s)$ correction to the $\log mt/MW$ term 74 and the residual terms of $O(s)$ 75. Partial width electroweak 1-loop couplings 76. These non-factorizable corrections sufficiently inclusive final states, i.e. for loose cuts to the invariant mass of the secondary fermions 77.

The simple form of the improved Born approximation [740] (2,3,5,8) allows the Standard Model and its dominant oblique loops, together with certain Beyond the Standard Model (BSM) physics [740], to be fit to experimental data with ease, and also with clear and accurate interpretation and dissemination of HP-EW-QL information to the wider physics community.

C. α, G_μ, M_Z renormalization scheme

We need to fix 3 experimentally measured ² inputs

$$\begin{aligned} \frac{1}{e_*^2(q^2)} & = \frac{1}{e_*^2(0)} - \Delta_Q(q^2) \\ \frac{1}{4\sqrt{2}G_\mu^*(q^2)} & = \frac{1}{4\sqrt{2}G_\mu^*(0)} - \Delta_+(q^2) \\ M_Z^2 & = M_{Z Pole}^2 = \frac{e_*^2(-M_Z^2)}{s_*^2(-M_Z^2)c_*^2(-M_Z^2)} \\ & \times \frac{1}{4\sqrt{2}G_\mu^*(-M_Z^2)\rho^*(-M_Z^2)} \end{aligned} \quad (9)$$

and **calculate** the $SU(2)_{L-R}$ custodial isospin symmetry breaking function

$$\begin{aligned} \frac{1}{4\sqrt{2}G_\mu^*(q^2)\rho^*(q^2)} & = \frac{1}{4\sqrt{2}G_\mu^*(0)\rho^*(0)} - \Delta_3(q^2) \\ \frac{1}{\rho^*(0)} & = 1 - 4\sqrt{2}G_\mu^*(0) [\Pi_{33}(0) - \Pi_{+-}(0)] \end{aligned} \quad (10)$$

while making use of certain auxiliary functions

$$\begin{aligned} \Delta_Q(q^2) & = \text{Re} \left[\Pi'_{QQ} + 2\Gamma'(\xi) \right]_{q^2} \\ & - \text{Re} \left[\Pi'_{QQ} + 2\Gamma'(\xi) \right]_{q^2=0} \\ \Delta_3 & = \text{Re} \left[\Pi_{33} - q^2 \Pi'_{3Q} - 2\Pi_{3Q}^L(\xi) \right]_{q^2} \end{aligned}$$

² Electro-weak renormalization schemes:

- On-shell scheme $\sin^2(\theta_W) = s_W^2 = 1 - M_W^2/M_Z^2$, $\alpha_{QED}(0), G_\mu(0)$ [307, 386, 387, 422, 558, 612, 613, 792, 793, 804, 887, 963–965, 968?, 969]
- $\alpha_{QED}(0), G_\mu(0), M_Z$ scheme [4, 676, 740, 972]
- MSBar scheme 1922
- $\sin^2(\theta_W) = 1 - M_W^2/M_Z^2$ with mixing angle deduced from neutrino-electron scattering 23
- Effective running couplings 24,25, [740].
- Scheme-dependence allows estimate of missing higher-order contributions (see e.g. 26 for a comprehensive study)

$$\begin{aligned}
& -Re \left[\Pi_{33} - q^2 \Pi'_{3Q} - 2\Pi_{3Q}^L(\xi) \right]_{q^2=0} \\
\Delta_+ = & Re \left[\Pi_{+-} - q^2 \Pi'_{3Q} - 2\Pi_{3Q}^L(\xi) \right]_{q^2} \\
& -Re \left[\Pi_{+-} - q^2 \Pi'_{3Q} - 2\Pi_{3Q}^L(\xi) \right]_{q^2=0}
\end{aligned} \tag{11}$$

The α, G_μ, M_Z renormalization scheme specifies experimental input data from atomic Thomson scattering, the muon lifetime³ and the precise Z mass:

$$\begin{aligned}
\frac{e_*^2(0)}{4\pi} &= \alpha_{QED}(0) = (137.03602)^{-1} \\
G_\mu^*(0) &= G_\mu + \text{CertainGaugeIndependent} \\
&\quad \text{Vertices, Boxes, Radiation} \\
G_\mu &= 1.16637 \times 10^{-5} GeV^{-2} \\
M_Z &= 91.1876 \pm 0.0021 GeV
\end{aligned} \tag{12}$$

D. High precision OBLIQUE electro-weak quantum loop experimental predictions

HP-EW-QL LEP/SLC asymmetry experiments are predicted by

$$s_*^2(-M_Z^2)c_*^2(-M_Z^2) = \frac{e_*^2(-M_Z^2)}{4\sqrt{2}G_\mu^*(-M_Z^2)\rho^*(-M_Z^2)M_Z^2} \tag{13}$$

Running $e_*^2(q^2), s_*^2(q^2), G_\mu^*(q^2)$ and $\rho^*(q^2)$ down, from $q^2 = -M_Z^2$, to $q^2 = -M_W^2$ predicts the W^\pm mass.

$$M_W^2 = \frac{e_*^2(-M_W^2)}{s_*^2(-M_W^2)c_*^2(-M_W^2)4\sqrt{2}G_\mu^*(-M_W^2)} \tag{14}$$

Running $s_*^2(q^2)$ down further to low space-like values predicts (including quark and lepton thresholds) ν -hadron scattering and parity violation in atoms.

$$\begin{aligned}
s_*^2(0)c_*^2(0) &= \left[\frac{s_*^2(0)c_*^2(0)}{s_*^2(-M_Z^2)c_*^2(-M_Z^2)} \right] \\
&\quad \times s_*^2(-M_Z^2)c_*^2(-M_Z^2)
\end{aligned} \tag{15}$$

Note that the **running** of $e_*^2(q^2), s_*^2(q^2)$ does not depend on m_{BEH} , and only mildly on m_{top} . Running $s_*^2(q^2)$ for time-like and spacelike regions was first shown, for $M_Z, m_{top} = 93, 40$ GeV, in Figure 12 of [740].

³ Fermi coupling constant G_μ . Originally, the μ -lifetime τ_μ has been calculated within the framework of the effective 4-point Fermi interaction. Beyond the well-known 1-loop QED corrections 51, the 2-loop QED corrections in the Fermi model have been calculated quite recently 52

E. Hadronic uncertainty for 4-lepton processes

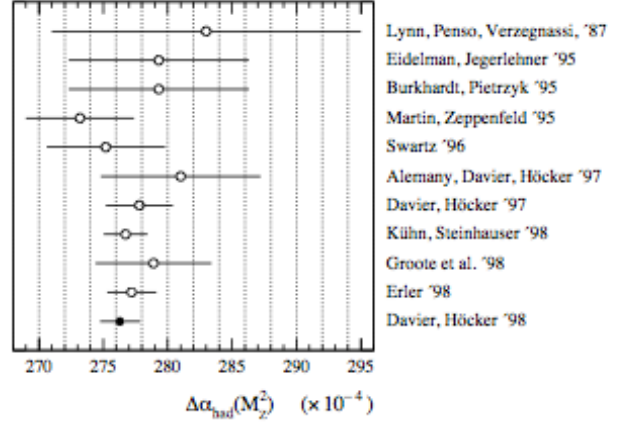


FIG. 1

Hadronic vacuum polarization dispersion relation value 29,30 agrees with another independent analysis 31. Improve uncertainty with τ -decays 32. Recently other attempts have been made to increase the precision of $\alpha_*(M_Z^2)$ 33,35,36,37 by theory-driven analyses of the dispersion integral based on perturbative QCD. Quark-mass-dependent $\mathcal{O}(\alpha_{QCD}^2)$ QCD corrections 38 near b and c thresholds (39 in the massless approximation). Minimize error from less reliable regions 36 and 33: e.g. by rescaling open charm regions 35. Separate $\alpha_*(M_Z^2)$ in un-subtracted $\bar{M}\bar{S}$ in 37 gives comparable error. Ref 34 gives history of dispersion relation starts by quoting Lynn, A. Penso, Verzegnassi 1987. Higher-dimensional operators in the operator product expansion 40 probe non-perturbative contributions, showing they are negligibly small 33,34. Recent preliminary measurements of R at BES at 2.6 and 3.3 GeV show values slightly lower than the previous data 41,2. Value $\Delta\alpha_{hadronic} = 0.0280 \pm 0.0007$ independent of theoretical assumptions on QCD at lower energies and thus less sensitive to potential systematic effects not under consideration now 42.

III. STANDARD MODEL HIGH-PRECISION PREDICTIONS

B.W.Lynn and R.G.Stuart [670, 682] successfully predicted both the precise top-quark and Brout-Englert-Higgs (BEH) masses, m_{top} and m_{BEH} respectively, from High-Precision Electro-Weak Quantum-Loops (HP-EW-QL) on Z-pole in 1984, with their prediction to be compared to future

LEP/SLD Z-pole data ⁴.

Inversely, when m_{top} and m_{BEH} are input into HP-EW-QL theory, LEP, SLD, CDF, D0 and NuTeV experiments vindicate the re-summed 1-loop and 2-loop quantum structure of the Standard Model [155].⁵

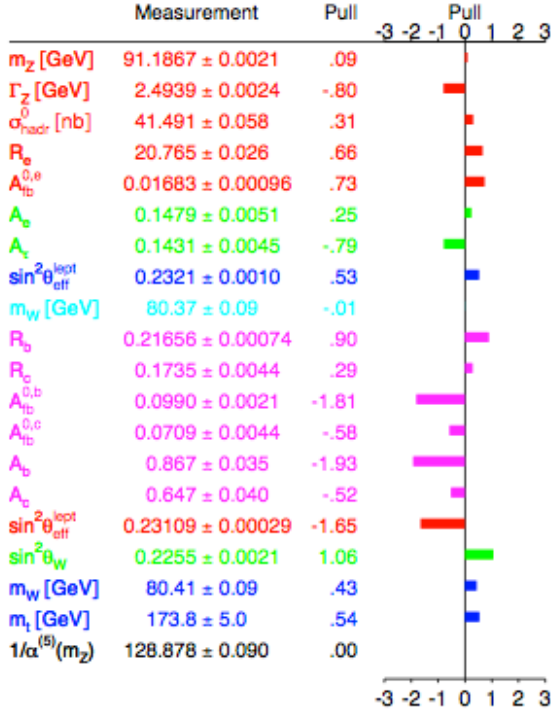


FIG. 2

⁴ Certain of [682]’s (α_{QED}, G_μ, M_Z renormalization scheme) Z-pole results were mutually directly confirmed with M. Veltman (for G. Passarino), W. Hollik and M. Consoli [558, 625, 968]: unfortunately, because they used the on-shell renormalization scheme for the weak mixing angle $s_W^2 = 1 - \frac{M_W^2}{M_Z^2}$, those authors were unable to (and do not) suggest or display the extreme HP-EW-QL sensitivity to m_{top} and m_{BEH} on Z-pole, which later contributed most directly to those particles’ experimental discovery.

⁵ Note that the experimental value for l points at the presence of genuine electroweak corrections by 3.5 standard deviations, an observation that has been persisting already for several years 83. The inclusion of the two-loop electroweak corrections m_{2t} from 61 yields a sizeable positive contribution to s_{2l} , see Figure 4. The inclusion of this term hence strengthens the upper bound on MH. The inclusion of the two-loop electroweak corrections m_{2t} from 61 yields a sizeable positive contribution to s_{2l} , see Figure 4. The inclusion of this term hence strengthens the upper bound on MH.

A. Theory fundamentals

1. Invention of symbolic manipulation programming (in 1965 on punch cards!), expressly for HP-EW-QL, by M.J.G. Veltman [568, 569]
2. Renormalizability and unitarity in the Standard Model [179, 548–553, 618].
3. Theory fundamentals for practical calculations in the Standard Model [554–557, 564]
4. Translation of subtle theory to experimental possibilities [558?], [625, 670, 682]
5. Constraints on BEH mass from vacuum stability, absence of Landau pole, and lattice gauge theory [90–104, 404]
6. Avoidance of unphysical negative quartic couplings from the negative top quark contribution to its beta function [95–99]

B. Vertices, boxes, higher orders

There is therefore no real need here to “pinch”, or worry about, the 1-loop gauge-dependence of vertex and box diagrams: ‘*-scheme’ Γ and Π_{3Q}^L counterterms will cancel that separately.

1. 1998 report of the Particle Data Group 53, MW, MZ of the vector bosons, expressed in terms of α and G; in 1-loop order it is given by: r higher than 1-loop (see also the contribution by Kühn 54 to these proceedings leading log resummation 55 of
2. resummation of the leading m_{2t} contribution 56 in terms of the complete $O(s)$ corrections to the self energies are available 57,58. Non-leading QCD corrections to r are also available 59. non-leading higher-order terms containing mass singularities of the type $2 \log(MZ/m_f)$ from light fermions are incorporated 60.
3. subleading $G_{2m_{2t}}$ MZ2 contribution of the electroweak 2-loop order 61
4. Meanwhile exact results have been derived for the Higgs-dependence of the fermionic 2-loop corrections in r 62, and comparisons were performed with those obtained via the top mass expansion 63. Pure fermion-loop contributions (n fermion loops at n -loop order) have also been investigated 63,64.

5. Effective Z boson couplings: 1-loop diagrams without virtual photons, the non-QED or weak corrections. written in terms of fermion-dependent overall normalizations f and effective mixing angles s_2^f in the NC vertices (see e.g. [65] G2m2t MZ2 for the leptonic mixing angle s_2^l have also been obtained in the meantime, as well as for [66]).
6. difference between the d and b couplings calculated perturbatively, including the complete 1-loop order term [67] leading electroweak 2-loop contribution of $O(G^2 m_4^2)$ [46,68]
8. Another recent analysis [90] (for earlier studies see [91,92]) based on the data set of summer 1998 yields a Higgs mass $M_H = 107^{+67}_{-45}$ GeV indirect determination of the Higgs mass range has shown that the Higgs is light, with its mass well below the non-perturbative regime.
9. For LEP 2, an W [93] error of about 40 MeV in M_W can be reached .
10. The present lower limit [95]

C. Electromagnetism and bremsstrahlung

1. Electro-magnetic $\alpha_{leptonic}$: The 2-loop correction [27], and also the 3-loop contribution is now available [28].
2. The radiator function $H(k)$ with soft-photon resummation and the exact $O(2)$ result for initial-state QED corrections is given in ref [78] . It has been improved recently by the $O(3)$ term [79].
3. Bhabha scattering in the forward direction=luminosity, higher-order QED corrections [80]. BHLUMI are an important step in pinning down the theoretical error from 0.11

D. Fitting theory to experiment

1. Computational integration of theoretical knowledge [672]
2. Electroweak precision data on Z and W bosons [1, 2] and $p\bar{p}$ collider Tevatron [2, 936? ? -940]
3. Discovery of the top quark at Tevatron [941?]. 1998 mass measurement [3]
4. The leptonic mixing angle determined via ALR by the SLD experiment [82] and the s_2^l average from LEP, still differ by 2.8 standard deviations.
5. W mass prediction [80.372], experiment [80.39 0.06]
6. Standard model global fits: The FORTRAN codes ZFITTER87 and TOPAZ088 (both used in Figure ?? , good agreement between the predictions from the two independent programs [89].
7. The 1998 upper limit to the Higgs mass at the [95]

E. Other processes

1. History of theoretical physics: this paper and [354, 358]
2. s_2^W the ratio M_W / M_Z can indirectly be measured in deep-inelastic neutrino-nucleon scattering. The average from the experiments CCFR, CDHS and CHARM [85] with the recent NUTEV result [86]
3. Parity violations in atoms by B.W. Lynn, and independently W.J. Marciano and A. Sirlin, in 1982 [418, 726, 727, 747, 765], later confirmed experimentally [?] .
4. Production and decay of W bosons [72-88, 403, 825-827, 896, 897]
5. BEH decays $H \rightarrow W^+W^-$, $H \rightarrow ZZ$, $H \rightarrow f\bar{f}$ [105-109, 980-986]
6. The decay $B \rightarrow X_s\gamma$ [110-122, 723]
7. Muon anomalous magnetic moment [12, 13, 25, 48, 123, 130-136, 399, 573-577, 724, 725, 877]. Theoretical uncertainties from virtual hadrons: the dominant error is from vacuum polarization; that involving light-by-light scattering has now been established with acceptable uncertainty.

IV. SUCCESSFUL PREDICTION OF THE TOP QUARK MASS FROM LEP/SLD DATA

B.W. Lynn and R.G. Stuart's successful 1984 Z-pole physics 1-loop HP-EW-QL prediction for the top-quark mass [670, 682] fit to EXPOSTAR [671-674] led to its discovery, and contributed directly to the award of the 1999 Nobel Prize in Physics for the renormalizability and unitarity of spontaneously broken Yang Mills gauge theories [367];

A. EXPOSTAR HP-EW-QL prediction for the top quark mass, based on ALEPH rapidity, confirmed by experimental discovery at FNAL

The 1991/1992 published literature’s HP-EW-QL predictions for the top-quark mass were:

- To quote from P. Langacker and Mingxing Luo, “Implications of precision electroweak experiments for $m_t, \rho_0, \sin^2(\theta_W)$ and grand unification” [1008]: “The implications of precision Z-pole, W-mass, and weak-neutral-current data for $SU(2) \times U(1)$ models are described. Within the minimal model one finds ... the top-quark mass is predicted to be $m_{top}^{1991WorldAverage1} = 124_{-34-15}^{+28+20} GeV$, where the second uncertainty is from m_H , with $m_t < 174(182) GeV$ at 90(95)% C.L.”
- They were supported by [363] which fit Z-pole, FNAL top-quark non-discovery, ν -hadron collisions, M_W, M_Z and parity violations in atoms data with $m_{top}^{WorldAverage2} = 122_{-20}^{+25} GeV$ and $m_{BEH}^{WorldAverage2} = 65_{-4}^{+245} GeV$
- The LEP 1991 data combined fit, in G. Alexander et al., the LEP Collaborations ALEPH, DELPHI, L3 and OPAL, “Electroweak Parameters of the Z^0 Resonance and the Standard Model” [174]., combined the 4 LEP experiments’ separate fits to $\approx 650,000 Z^0$ decays to hadrons and charged leptons in 1989/1990 LEP1 data; Results: $m_{top;\alpha_{QCD}Unconstrained}^{LEP1991} = 94_{-24}^{+53+23} GeV$, and $m_{top;\alpha_{QCD}Constrained}^{LEP1991} = 124_{-56-21}^{+40+21} GeV$ when α_{QCD} was constrained to $\alpha_{QCD} = 0.118 \pm 0.008$ [362].

Meanwhile, in a completely separate HP-EW-QL effort, EXPOSTAR [671–674] was written to contain: all LEP1 and SLC 4-fermion processes (especially on/near Z pole), including Bhabha scattering; complete 1-loop vertex and box corrections; dispersion relation inclusion of light-hadron contributions to vector self-energies; all-orders resummed 1-loop oblique [740]; certain 2-loop electroweak; photon radiation including hard and resummed/exponentiated soft/co-linear photon emission; leading 2-loop $\mathcal{O}(\alpha_{QCD} m_{top}^2)$; minimal use of Monte-Carlos for hard wide-angle hard photons and box diagrams; appropriate perturbative α_{QCD} corrections for final-state heavy quarks and total hadrons.

Using EXPOSTAR, “On estimating the top quark mass from global analyses of e^+e^- collision data” [672] first identified two previously unknown

ALEPH, DELPHI, L3, OPAL, and SLD Monte-Carlo-based systematic analysis errors: background rejection (Figure 3 in [672]); and efficiency vs. inverse-purity (Figure 2 in [672]). These we repaired by cutting particle-production data in rapidity (i.e. rather than acolinearity) and working in the e^+e^- collision frame.

As a bonus, the resulting almost-Monte-Carlo-free speed-up factor of $\approx 50,000$ -250,000 allowed [673, 674] to fit M_Z, m_{top}, m_{BEH} directly to ALEPH data, without reference to baroque definitions of non-physical running $s_*^2(q^2), \bar{s}^2(q^2), s_{MS}^2(q^2)$ or other (e.g. $e_*^2(q^2), G_\mu^*(q^2), \rho^*(q^2), s_W^2 = 1 - \frac{M_W^2}{M_Z^2}$, etc.) intermediate quantities.

After [673] fit 1990 ALEPH-LEP1 data with EXPOSTAR [671], D. Levinthal and F. Bird (with B.W. Lynn and R.G. Stuart) [674], “Fits of 1990/1991 (ALEPH) Data with EXPOSTAR, including non-Standard Model Extensions” followed with significantly more ALEPH data.

In stark contrast to the world-physics-community consensus $m_{top} \approx 124 GeV$ [174, 363, 1008], our rapidity-cut EXPOSTAR fits of the 1990/1991 ALEPH data (directly to $M_Z, m_{top}, m_{BEH}, \alpha_{QCD}$) raised the top-quark mass dramatically to $m_{top;\alpha_{QCD}Unconstrained}^{RapidityALEPH} = 206_{-37}^{+31} GeV$. When α_{QCD} was constrained to the then-ALEPH value $\alpha_{QCD}^{ALEPH1991} = 0.125 \pm 0.005$, our $m_{top;\alpha_{QCD}Constrained}^{RapidityALEPH} = 181_{-34}^{+29} GeV$ was ~ 60 GeV over the then-ALEPH acollinearity-fit value [675] for the same data set: i.e. much closer to the top-quark’s 1995 experimental discovery-mass and since-improved mass measurements.

B. Successful prediction for the Top Quark Mass, based on LEP rapidity

The ALEPH (M. Martinez, J. Harton, A. Blondel et. al. [364]) and ZFITTER collaborations compared EXPOSTAR and ZFITTER to check/confirm 1-loop correctness in the early days of ZFITTER [357, 817]. The other 4 LEP/SLC experiments and ZFITTER [817] later also adopted our rapidity/collision frame analysis and speed-up factor.

In 1993, ALEPH [365, 671] predicted (the 2nd errors in [365, 366] are from m_{BEH}) $m_{top}^{ALEPH1993} = 156_{-25-22}^{+22+17}$ from $\approx 520,000$ ALEPH 1989-1992 Z-decays, ν -hadron collisions, and $\frac{M_W}{M_Z}$ in $p\bar{p}$ collisions.

In 1994, M. Martinez [366, 671, 817] predicted $m_{top}^{LEP1993} = 173_{-13-20}^{+12+18}$ from $\approx 2,000,000$ Z-decays in 1989-1993 LEP data,

The top-quark was discovered, close to D. Levinthal, F.Bird, B.W. Lynn and R.G. Stuart’s

1992 [671–674], ALEPH’s 1993 [365, 671], and M. Martinez’s 1994 [366, 671, 672, 817] predicted masses, in 1995 at FNAL (CDF, D0) with (eventually) $m_{top}^{Experiment} = 172.44 \pm 0.13 \pm 0.47 GeV$.

V. BROUT-ENGLERT-HIGGS MASS

Once the top-quark mass was known, i.e. after its discovery in 1995, the Standard Model HP-EW-QL could be used [670, 682, 858] to predict the SM Brout-Englert-Higgs [168–173] mass.

A. 2011 Prediction of the BEH mass

The LEP, SLC and CLEO HP-EW-QL fits, and FNAL top-quark discovery and BEH scalar search data, as of July 2011, are displayed in Figure 3 [357].

- The blue contour shows the 2-loop HP-EW-QL Z-pole and Nutev physics m_{BEH} prediction: (1-loop [670, 682] augmented by certain 2-loop [155, 858]) fit to data from LEP (ALEPH, DELPHI, L3, OPAL), SLC(SLD) and CLEO(Nutev), with SM theory from ZFITTER [356–358, 793, 817–819] and TOPAZ0 [?],
- The green belt is the top-quark mass from FNAL (CDF, D0)
- The yellow regions are excluded by LEP2 (ALEPH, DELPHI, L3, OPAL) and FNAL(CDF, D0) searches.

Putting this all together yields the 2-loop HP-EW-QL Z-pole-physics m_{BEH} prediction [670, 682, 858] fit to data from LEP, FNAL and NuTeV, with ZFITTER and TOPAZ0 SM theory [356–358, 793, 817–819] in Figure 4.

The successful July 2011 prediction in Figure 4 led directly to the July 4, 2012 discovery of the BEH scalar at the LHC, at (eventually) $m_{BEH} = 125.09 \pm 0.24 GeV$: it was confirmed as the Standard Model BEH on 14 March 2013. The HP-EW-QL prediction in Figure 4 contributed to the award of the 2013 Nobel Prize in Physics [368] to F. Englert and P. Higgs.

D. Bardin, M. Bilenky, P. Khristova, M. Jack, L. Kalinovskaya, A. Olshevsky, S. Riemann and T. Riemann, were later honored as First Scientific Award of JINR, Dubna, 19 January 2001, refereed by Lev Borisovich Okun [819].

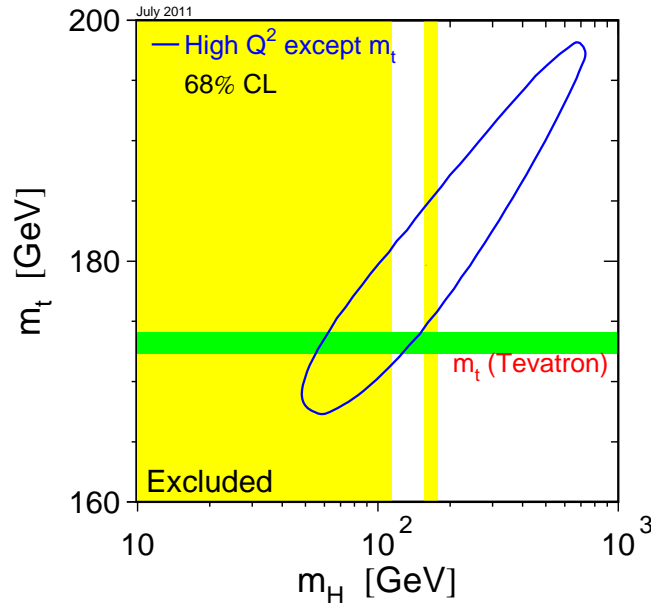


FIG. 3

B. ATLAS and CMS BEH scalar discovery papers’ reference bibliographies

J.D. Wells, “The theoretical physics ecosystem behind the discovery of the Higgs boson” [358], gives a nice review of a part of the physics used by (and explicitly referenced by) ATLAS and CMS, to analyse their vast data. Define **LHC’s background-data region** as the area to the right-hand-side of Figure 4, extended from $m_{BEH} \sim 200 GeV$ to $m_{BEH} \sim 1 TeV$: i.e. mostly focussed on “Beyond the Standard Model” theorist-imagined heavy BEH-like scalars’ production and decay channels, of which LHC has (so far) demonstrated experimentally that Nature makes no use. Its complement, **LHC’s discovery-data region**, between $m_{BEH} \sim 114 GeV$ and $m_{BEH} \sim 200 GeV$ includes the white regions. But the Standard Model HP-EW-QL m_{BEH} prediction [357, 670, 682, 817, 858] (parabolic lines) was shown (to 99% C.L.) to lie in the discovery-data region soon after the 1995 experimental discovery of the top-quark, i.e. at least a decade before the 2010 turn-on of the LHC.

It is therefore astonishing to read Wells’ incorrect assertion: “What made the discussion regarding the (SM) Higgs boson particularly difficult is that its mass was not known a priori ... One had to be prepared for all possibilities of the (Brout-Englert) Higgs boson mass all the way from

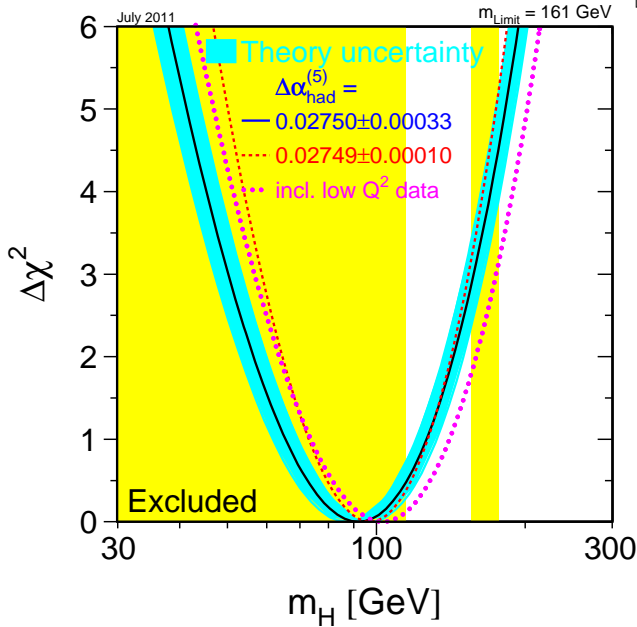


FIG. 4

$M_{BEH} \sim m_{electron}$ to $M_{BEH} \approx 1TeV$ ". In exact disagreement, Figures 3 and 4 show that $114GeV \leq M_{BEH} \leq 161GeV$ to high C.L.!

In-explicably, Wells' "ecosystem", and the ATLAS and CMS theory references he studies, do not contain a single reference to the huge exhaustively-documented $\sim 15,000$ man-year (pre-LHC-design) effort (ALEPH, DELPHI, L3, OPAL, SLD, CDF, D0, NuTeV experiments; and an army of HP-EW-QL theorists in the bibliography below) made to (successfully!) predict the Standard Model BEH mass from HP-EW-QL.

C. Implications for CDF, D0 and the Tevatron

Wolfgang Hollik [155] presents a comprehensive summary and detailed explanation of state-of-the-art theoretical Standard Model HP-EW-QL as of December 1998. Hollik and his references (included within the exhaustive bibliography below) make the case (already airtight and fully globally socialized before The 1998 International Conference on High Energy Physics (ICHEP1998) [155]) for the accuracy and reliability of the theoretical Standard Model HP-EW-QL prediction of the BEH mass; i.e. that precise SM prediction later confirmed by the 2012 experimental discovery.

All of this shows that the BEH scalar could have

been discovered at the Tevatron by Summer 2006. [944].

- LEP2 BEH scalar Higgs exclusion searches (Figure ??, 1998) show $M_{BEH} > 114GeV$
- CDF and D0 BEH scalar Higgs exclusion searches (Figure ??, 2006) $M_{BEH} \neq 149$ to 180 GeV [942];
- To high C.L., the HP-EW-QL prediction (i.e. by naively chopping off the LEP and FNAL BEH-exclusion regions) is $114GeV < M_{BEH} < 149GeV$.
- But, to quote J. Wells [358], "... the richest and most complicated mass range of $120GeV \leq M_{BEH} \leq 160GeV$... enabled competing branching fractions of the decays of the Higgs bosons into many states, such as b quarks [361], leptons, WW^* and ZZ^* ... (which) means that no one final state is necessarily the dominant signal of the Higgs boson ... the initial Higgs boson discovery (at LHC) was ultimately made by two channels, $H \rightarrow \gamma\gamma$ and $H \rightarrow ZZ^* \rightarrow (e^+e^-e^+e^-, e^+e^-\mu^+\mu^-, \mu^+\mu^-\mu^+\mu^-)$ ".

Figures ?? and ?? also show that the BEH could have been studied and, in time, confirmed as the Standard Model BEH, at the Tevatron.

- Multiple final states in the $114GeV \leq M_{BEH} \leq 161GeV$ region predicted by HP-EW-QL "... also provides in time a variety of Higgs study opportunities, since the nature of the Higgs boson can be discerned better by having many accessible decay channels rather than just one or two [358]."

One cannot help but wonder whether [that is, in an alternate landscape universe where proposed CDF and D0 detector upgrades (e.g. vertex detectors, etc. [944]) as well as any necessary Tevatron upgrades, were approved at FNAL after the discovery of the top-quark] the Standard Model BEH scalar would have been discovered, studied and confirmed a decade before the LHC did so.

VI. HP-EW-QL W^\pm MASS

A. Sirlin's precise 1980 1-loop prediction for the W^\pm mass [175, 386, 625, 670, 682], augmented by certain 2-loop contributions [858], contributed directly to the successful HP-EW-QL predictions of the top-quark and BEH masses, and the associated 1999 and 2013 Nobel Prizes [367, 368].

VII. NON-DECOUPLING OF VERY HEAVY “BEYOND THE STANDARD MODEL” (BSM) MATTER

The classification [178, 354, 676, 740, 746] of hypothesized BSM contributions to oblique HP-EW-QL has experimentally ruled out vast swaths of imagined BSM physics. Recent experimental exclusion of such BSM physics at CERN’s Large Hadron Collider has confirmed these theoretical predictions.

B.W. Lynn, M.E. Peskin and R.G. Stuart [178, 676, 746] first identified and named “oblique” loop corrections; first showed the extreme LEP1/SLC sensitivity to 1-loop non-decoupling effects of very heavy BSM particles (i.e. $SU(8) \times SU(8)$ and $O(16)$ Technicolor, heavy quarks and leptons, N=1 SUSY squarks and sleptons, N=1 SUGRA gauginos and Higgsinos, etc.); first proposed that such high precision indirect sensitivity be part of the practical and official discovery strategies of LEP1 and SLC. Table 1 is reproduced from [676]: its last 2 rows show achieved uncertainties.

There are 3 dimensionless oblique gauge-independent HP-EW-QL objects from which very heavy particles do not necessarily de-couple:

$$\begin{aligned} \Delta'_3(q^2) &= \frac{\Delta_3}{q^2}; & \Delta'_+(q^2) &= \frac{\Delta_+}{q^2}; \\ \rho^*(0) &= \rho_{Veltman}; \end{aligned} \quad (16)$$

The 1st and 2nd functions were introduced in July 1987 [740]. The 3rd, the parameter $\rho^*(0)$, is equal to M. Einhorn, D.R.T. Jones and M.J.G. Veltman, and separately M.J.G. Veltman and D. Ross’s, famous custodial symmetry breaking $\rho_{Veltman}$ parameter.⁶

Three “non-de-coupling” parameters were introduced [740] to capture the effects of very heavy $M^2 \gg M_Z^2$ new BSM fermions, split scalars and Technicolor theories:

$$\begin{aligned} \Delta'_3(-M_Z^2) &= \left[\frac{\Delta_3}{q^2} \right]_{-M_Z^2}; & \Delta'_+(-M_Z^2) &= \left[\frac{\Delta_+}{q^2} \right]_{-M_Z^2} \\ \rho^*(0) &= \rho_{Veltman}; \end{aligned} \quad (17)$$

⁶ ρ -parameter, originally defined as the ratio of the neutral to the charged current strength in neutrino scattering 43, is unity in the standard model at the tree level, but gets a deviation $\Delta\rho$ from 1 by radiative corrections. 43; $\frac{1}{\rho^*(q^2)} = 1 - \Delta\rho(q^2)$ [740] in oblique-loop re-summation. Main contribution is from (t, b) doublet 44; electroweak 2-loop part 45,46. $\delta\rho_{QCD}$ is the QCD correction to the leading $G_\mu m_{top}^2$ term 47,48; 3-loop coefficient 48 For large Higgs masses 1-loop $\sim \log M_H$ 49; the 2-loop contribution 50 shows a dependence 2-loop $\sim M_H^2$

and were common knowledge at 1987, 1988, 1989 and 1990 LEP workshops, seminars and international conferences [?].

For example, [740] [Eq. 6.7] first showed the constant 1-oblique-loop contribution to modern S , which would be due to a very heavy non-decoupling BSM fermion multiplet (e.g. a 4th generation of quarks and leptons, gauginos, Higgsinos) of isospin weight i :

$$\Delta'_3(-M_Z^2) = \frac{i(i+1)(2i+1)N_{colors}}{144\pi^2} \quad (18)$$

LEP/SLC therefore rules out⁷ a 6th quark/lepton generation with degenerate masses.

Technicolor theories with N_{TC} Techni-Color and N_{TF} Techni-Flavor fermions embed the 3 Nambu-Goldstone bosons (NGB) (i.e. resulting from $SU(2)_L \times U(1)_Y \rightarrow U(1)_{QED}$ symmetry breaking) within a Techni-Multiplet of pseudo-NGB pseudo-scalars with low-momentum non-linear sigma model interactions. Beginning in 1990, in order to prove the non-trivial gauge-independence (and non-decoupling) of such theories, using a dispersion relation appropriate to the resulting tower of Techni-Hadrons, later versions of [740]’s oblique-loop classification scheme were introduced [369–372, 785, 901, 923]. Pierre Ramond ([354] pg. 206 and Appendix 5) shows their 1-to-1 correspondence with [740]’s oblique-correction classification scheme of heavy non-decoupling BSM physics.

$$\begin{aligned} \Delta'_3(0) &\equiv -\frac{1}{16\pi}S \\ \Delta'_+(0) - \Delta'_3(0) &\equiv \frac{1}{16\pi}U \\ \frac{1}{\rho_{Veltman}} &\equiv 1 - \alpha_{QED}(0)T \end{aligned} \quad (19)$$

Using their gauge-independent Techni-Hadron dispersion relation, [369] gets Technicolor S

$$S \approx 0.3 \frac{N_{TF} N_{TC}}{2 \cdot 3} \quad (20)$$

and Extended-Technicolor (which attempted to predict quark and lepton masses) models’ T with $N_{Colors} = 3$

$$T \approx \frac{N_{Colors} m_{top}^2}{16\pi^2 s_*^2(0) c_*^2(0) M_Z^2} \frac{4N_{TC}}{9} \quad (21)$$

Then, using then-available HP-EW-QL data, [369] ruled out QCD-like Technicolor theories with a large Technisector.

⁷ Experimental non-existence of such BSM physics was 1st shown in a 1991 fit to ALEPH Z-pole data [673, 674].

Table 1: Oblique 1-Loop Responses of Various LEP/SLC Asymmetries and the W^\pm Mass

	$\delta A_{LR}^{e^+e^- \rightarrow \mu^+\mu^-} = \delta A_{\tau Polarization}^{e^+e^- \rightarrow \tau^+\tau^-}$ $= \delta A_{LR}^{e^+e^- \rightarrow Hadrons}$	$\delta A_{FB}^{e^+e^- \rightarrow \mu^+\mu^-}$ $= \delta A_{FB}^{e^+e^- \rightarrow \tau^+\tau^-}$	$\delta M_W (MeV)$
STANDARD MODEL			
<i>Heavy Quark Pair</i>			
Large \vec{I} Splitting	0.02	0.01	300
Degenerate	-0.004	-0.002	-42
<i>Heavy Lepton Pair</i>			
Large \vec{I} Splitting	0.012	0.006	300
Degenerate	-0.0013	-0.0006	-14
N=1 SUSY SM			
<i>Heavy Squark Pair</i>			
Large \vec{I} Splitting	0.02	0.01	300
Degenerate	0	0	0
<i>Heavy Slepton Pair</i>			
Large \vec{I} Splitting	0.012	0.006	300
Degenerate	0	0	0
N=1 SUGRA SM			
<i>Winos</i>			
$m_{3/2} \ll 100 GeV$	0.005	0.0025	100
$m_{3/2} \gg 100 GeV$	< 0.001	< 0.001	< 10
TECHNICOLOR			
<i>(No Dispersion Relation)</i>			
$SU(8) \times SU(8)$	-0.04	-0.018	-500
$O(16)$	-0.07	-0.032	-500
UNCERTAINTY			
Hadronic Theoretical	± 0.0033	± 0.0016	± 25
Experimental	± 0.0033	± 0.0016	± 40

Later 1992 analysis of HP-EW-QL 1990/1991 LEP data ruled out all then-existing Technicolor and Extended-Technicolor models.

VIII. HP-EW-QL STANDARD MODEL VINDICATED: RUNNING $s_*^2(q^2)$

Appendix A: Guide to the Bibliography

The High Precision Electro-Weak Quantum Loop literature is immense! For example, a SLAC IN-

SPIRE query of the HP-EW-QL literature, for Date ≤ 1988 , reveals a list of 338 theoretical physicists [177] then working in the field. The bibliography below is meant to be exhaustive for Date ≤ 1998 , as per SLAC INSPIRE, Cornell arXiv and W. Hollik [155]. We apologize if we have in-advertently left anyone out.

In addition, we have included all references in [155], i.e. included in W. Hollik's wonderful detailed and explanatory Plenary Session review and summary (as of December 1998) of the state of HP-EW-QL at the 1998 International Conference on High Energy (ICHEP98). After eliminating double-

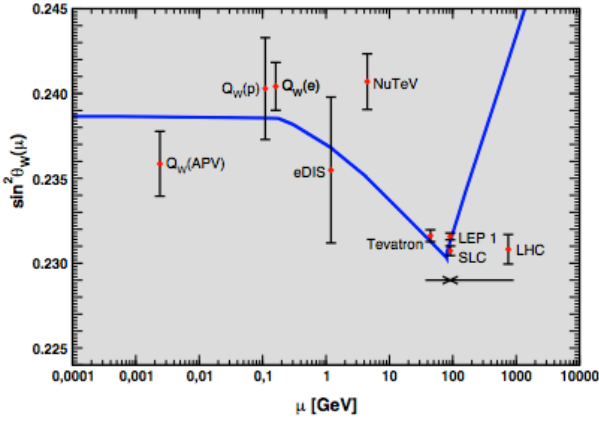


FIG. 5

Standard Model ideas and basics, along with (most of) their practical implementation and execution were already completed by the end of 1988.

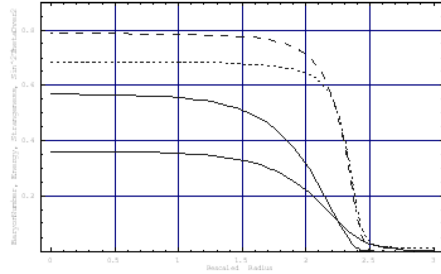


FIG. 7: Standard Model HP-EW-QL 1-loop [670, 682] and 2-loop [858] oblique diagrams containing a virtual top-quark

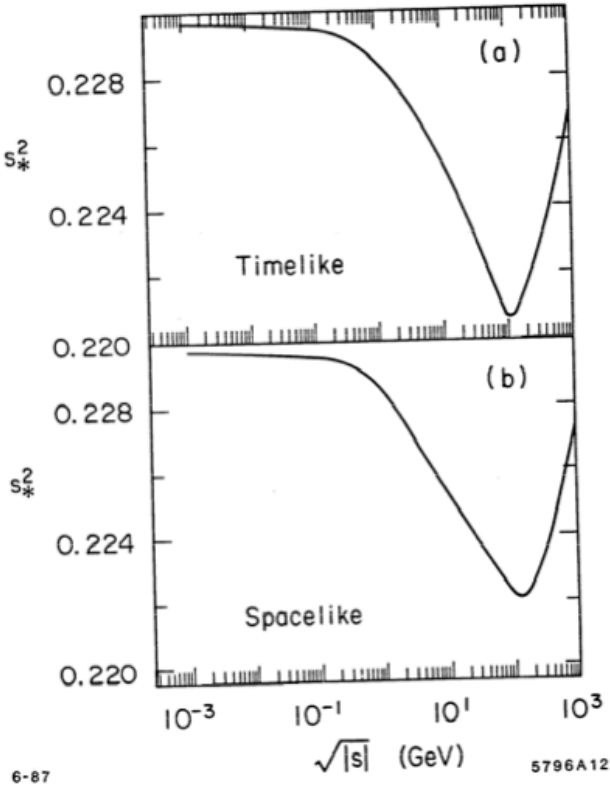


FIG. 6

counting, we generate 989 references by the end of 1998.

We have also included selected papers, books, experimental results, experimental analysis and Nobel citations for $1999 \leq \text{Date} \leq \text{present}$.

For $\text{Date} \leq 1988$, HP-EW-QL was a thriving scientific and mathematical mainstream and target of discovery, but its important theoretical Stan-

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