

Heavy neutrinos at the LHC

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On behalf of the ATLAS and CMS collaborations
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What is a heavy neutrino?

Standard model neutrinos were massless in theory

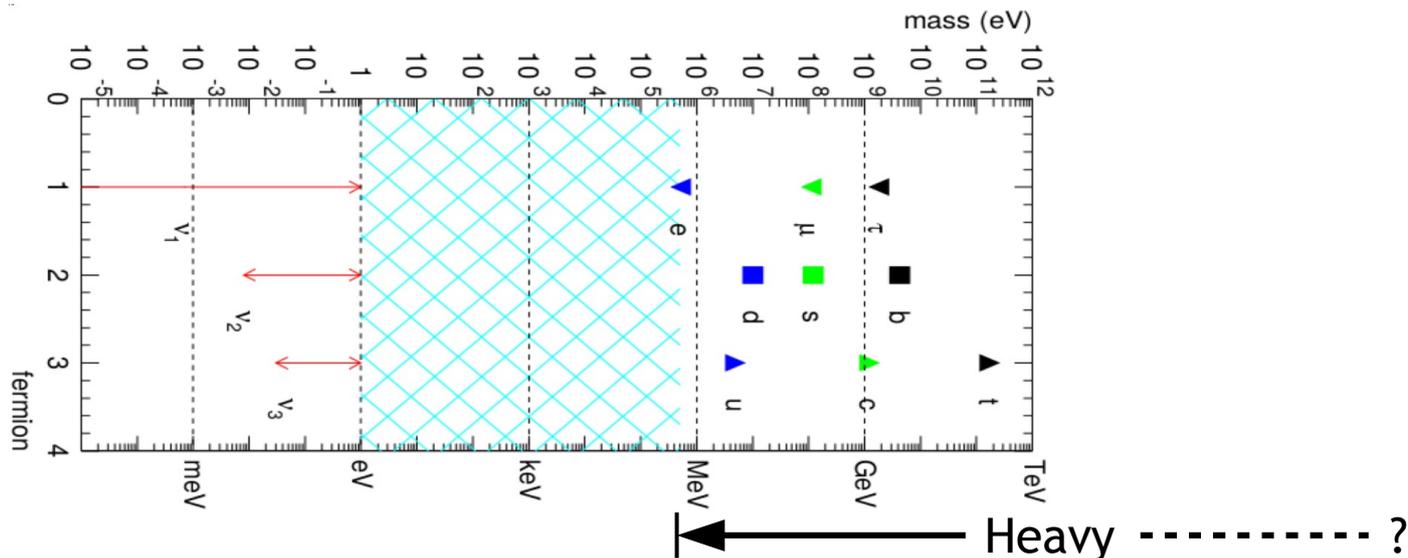
Neutrino oscillation observations show that they have non-zero masses

Nonetheless, neutrinos are “light” with masses $< 1\text{eV}$

- other SM particles start in the $\sim\text{MeV}$ range (or are massless)

“Heavy” implies a mass in the MeV range or (a lot) higher

I’ll mostly be talking about “Heavy Neutral Leptons”



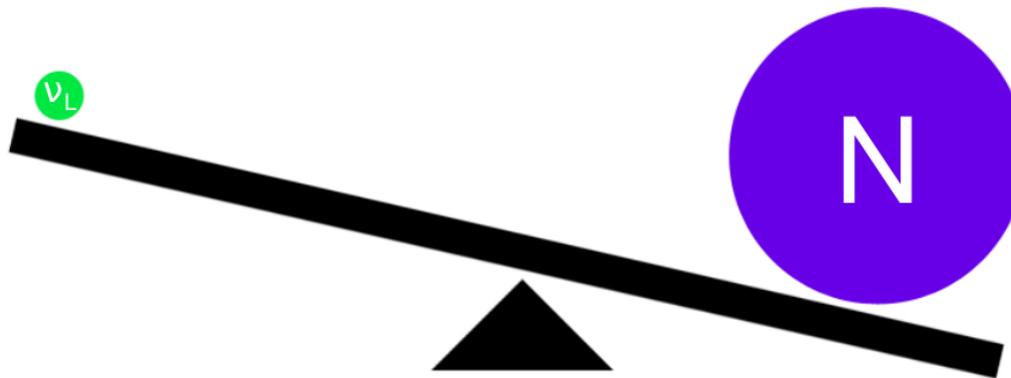
Why heavy neutrinos?

The known, “light” neutrinos are considered to be anomalously light, compared to the other massive fundamental particles

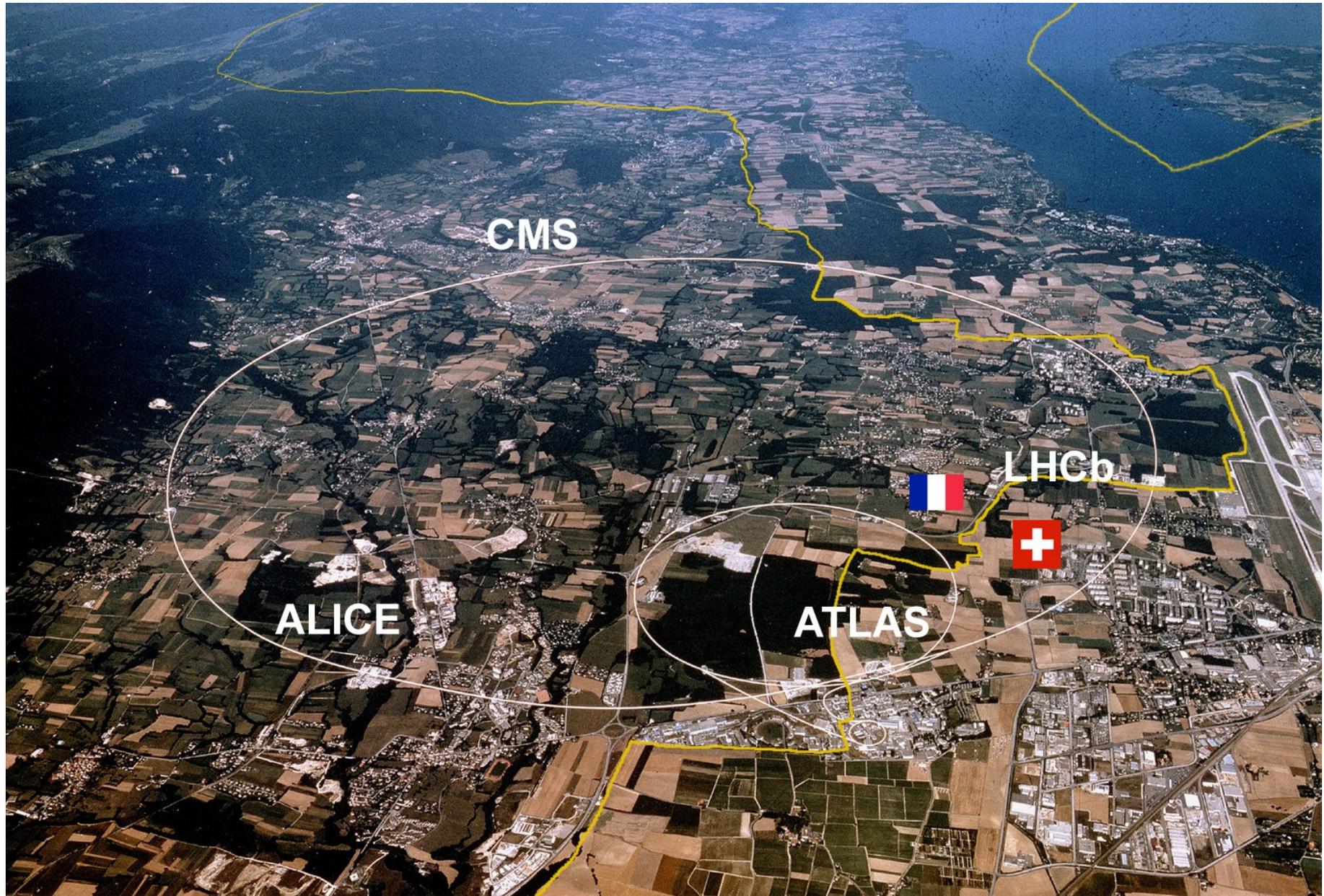
A common explanation for this is the seesaw mechanism, whereby each light neutrino couples to a heavier partner

The “seesaw” part refers to the feature of the theory that a neutrino can be made lighter simply by making a corresponding increase to the mass of their partner

Another nice feature is that the mass mixing matrix could now include some additional CP-violating phases that could explain the matter/antimatter asymmetry of the universe

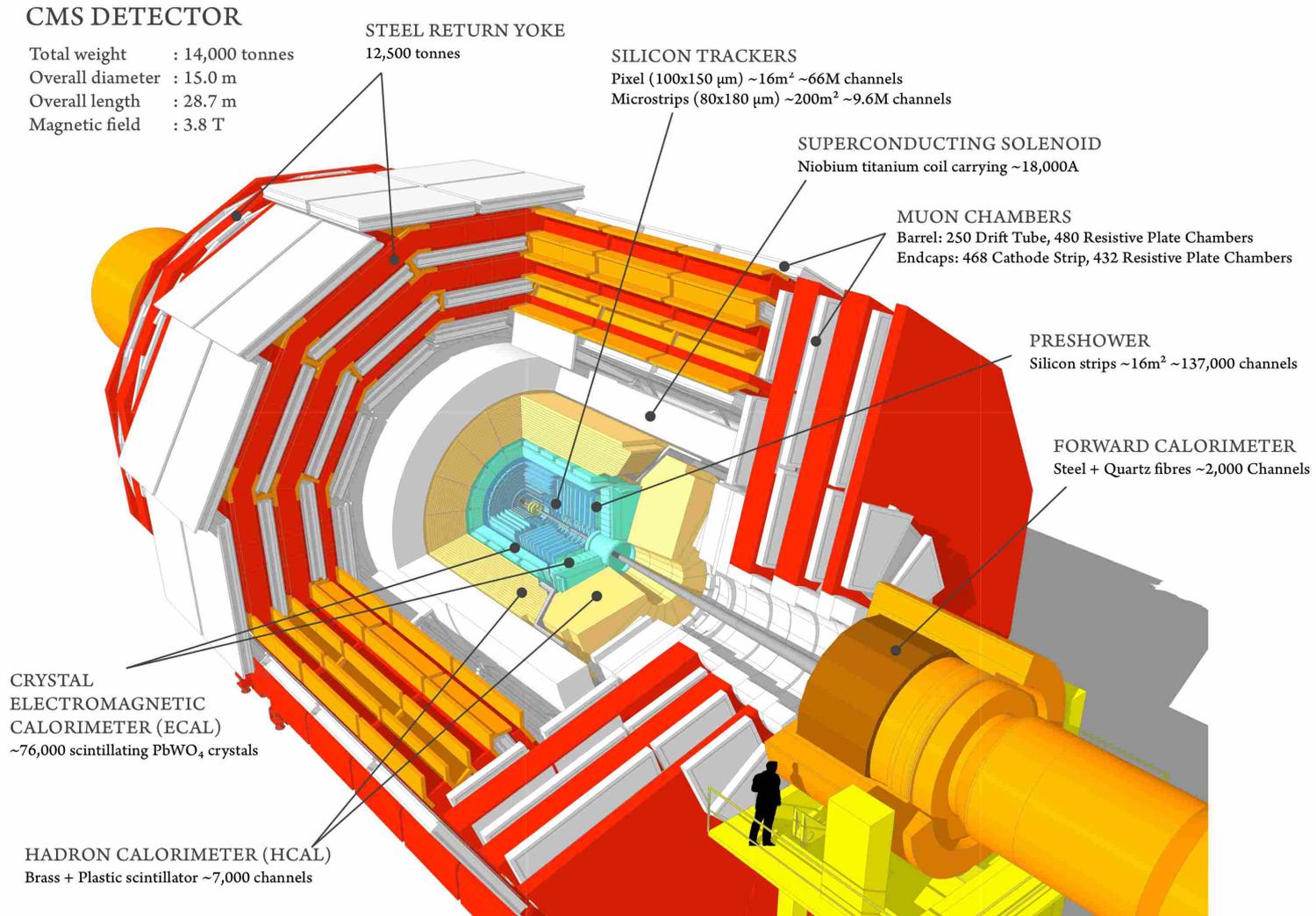


Large Hadron Collider



Prompt decays with CMS

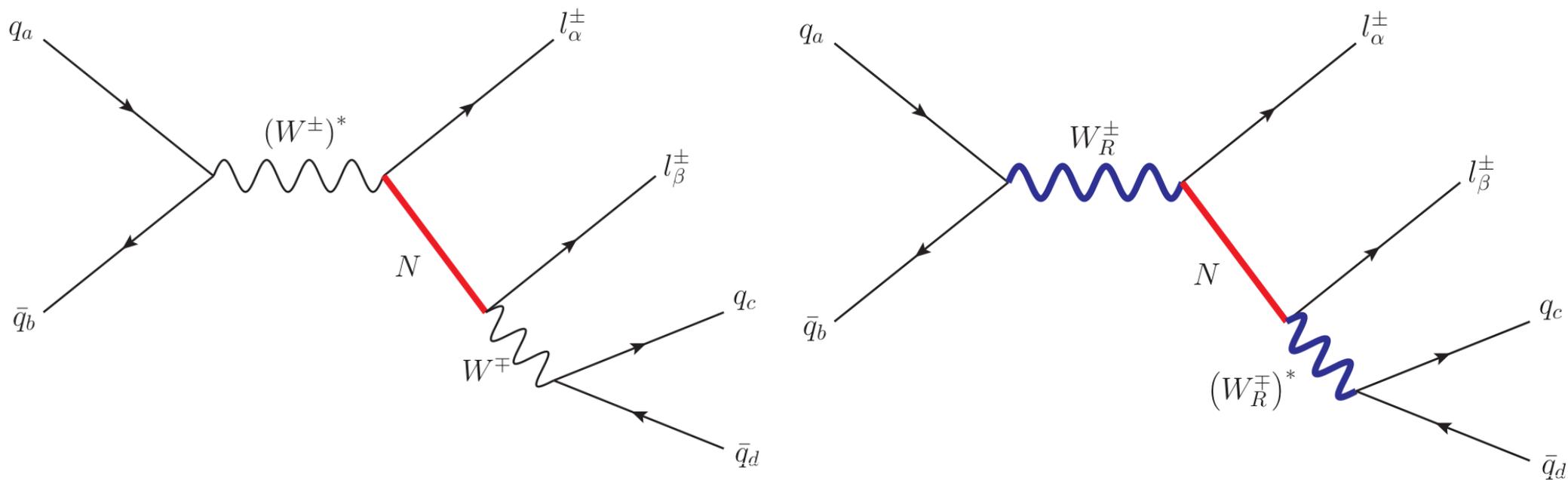
There are several recent results from CMS searching for decays where the heavy neutrino lifetime is short, and so it decays near the interaction point: a “prompt” decay



Prompt decays with CMS

Typical final states are a pair of leptons and a pair of quarks

In Left-Right Symmetric Models (LRSM), the decay chain may include a right-handed W-boson

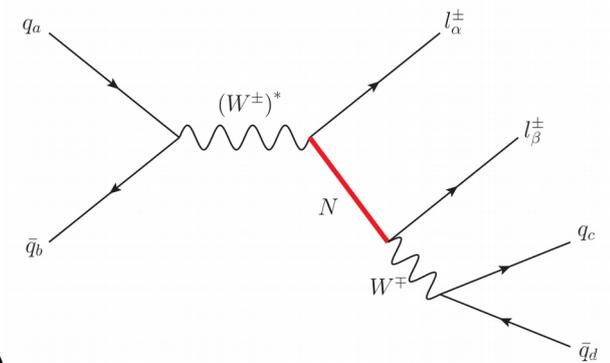


If the heavy neutrino is a Majorana fermion, the final state leptons may have the same charge - selecting for this reduces backgrounds dramatically

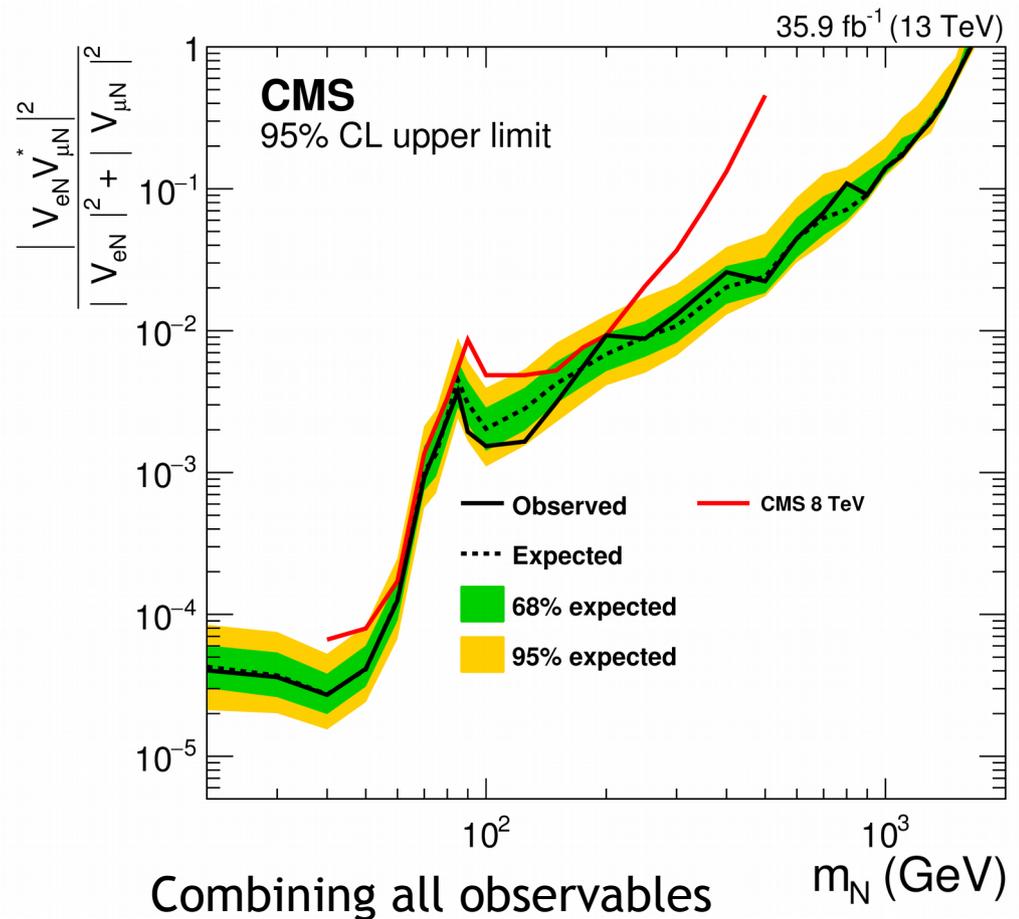
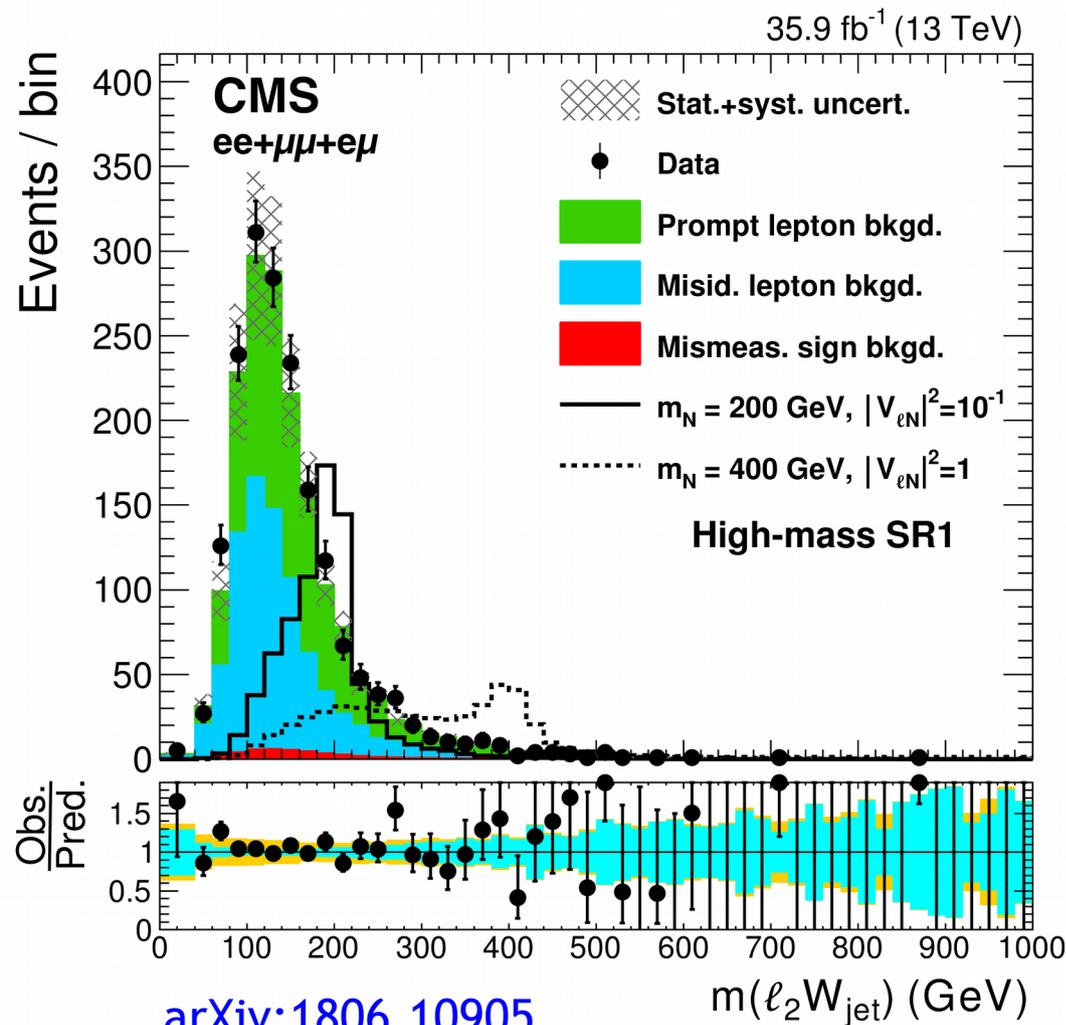
Some models also allow for lepton flavour mixing

Prompt decays with CMS 2018

Searching for a Majorana heavy neutrino, arising from the neutrino minimal standard model (νMSM, Type-1 seesaw)

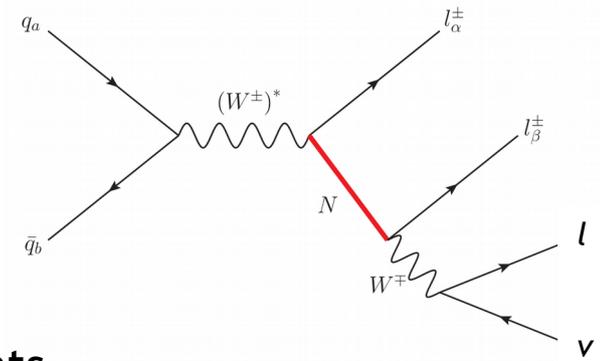


Two leptons with same charge but any flavour (ee, eμ or μμ)
and >1 jet: all flavours combined for different mass final states

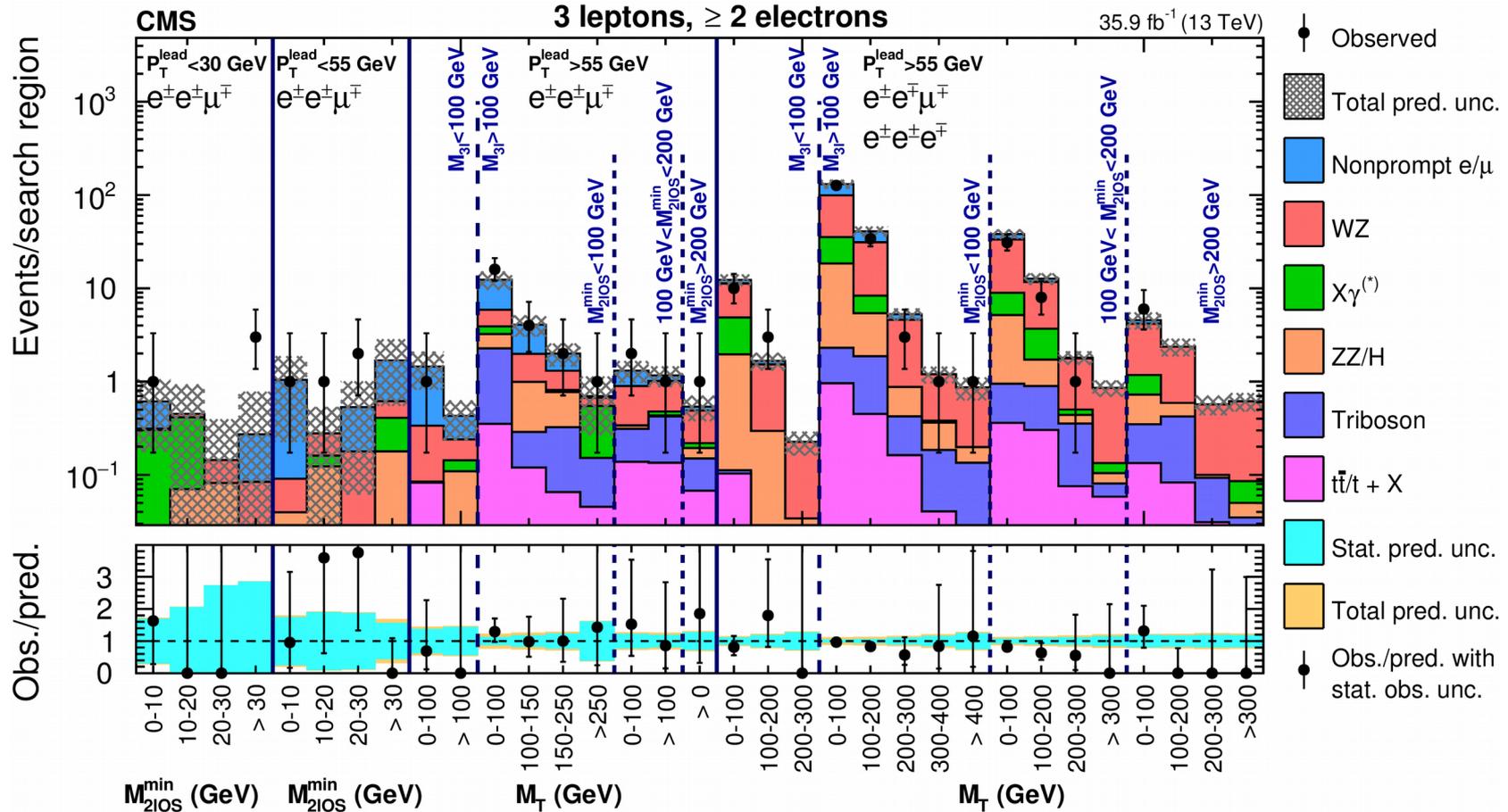


Prompt decays with CMS 2018

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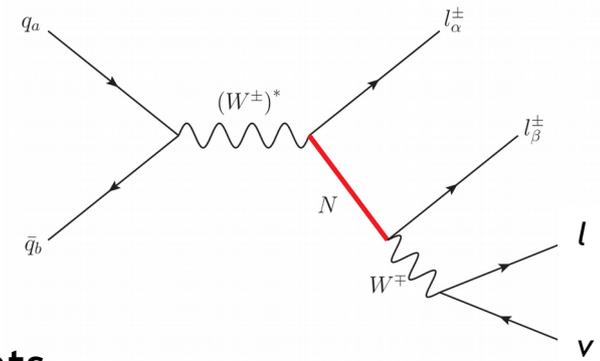


This decay chain ends with $W \rightarrow l\nu$, so the search counts events for several combinations of e and μ in different mass bins ($ee+X$ below)

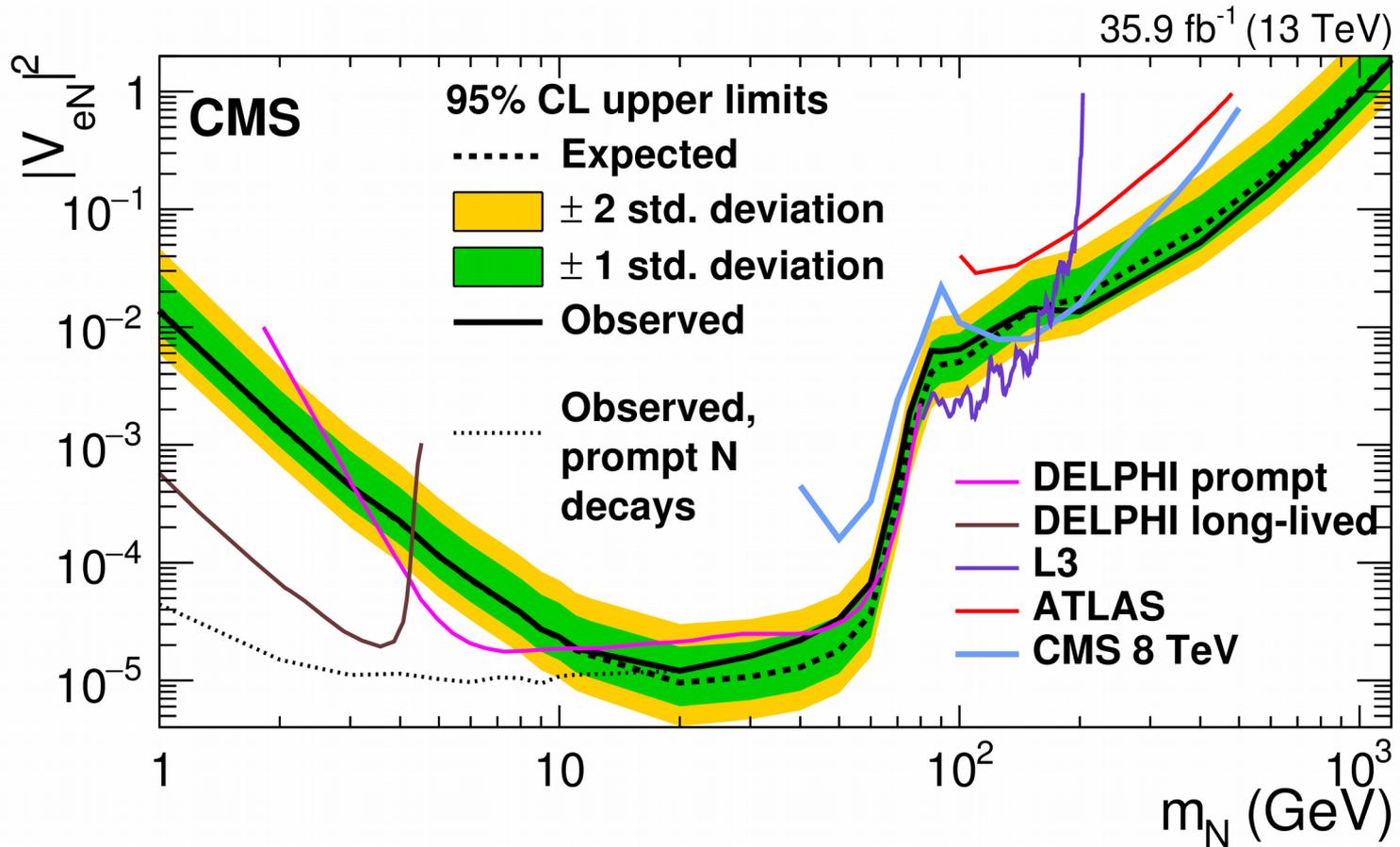


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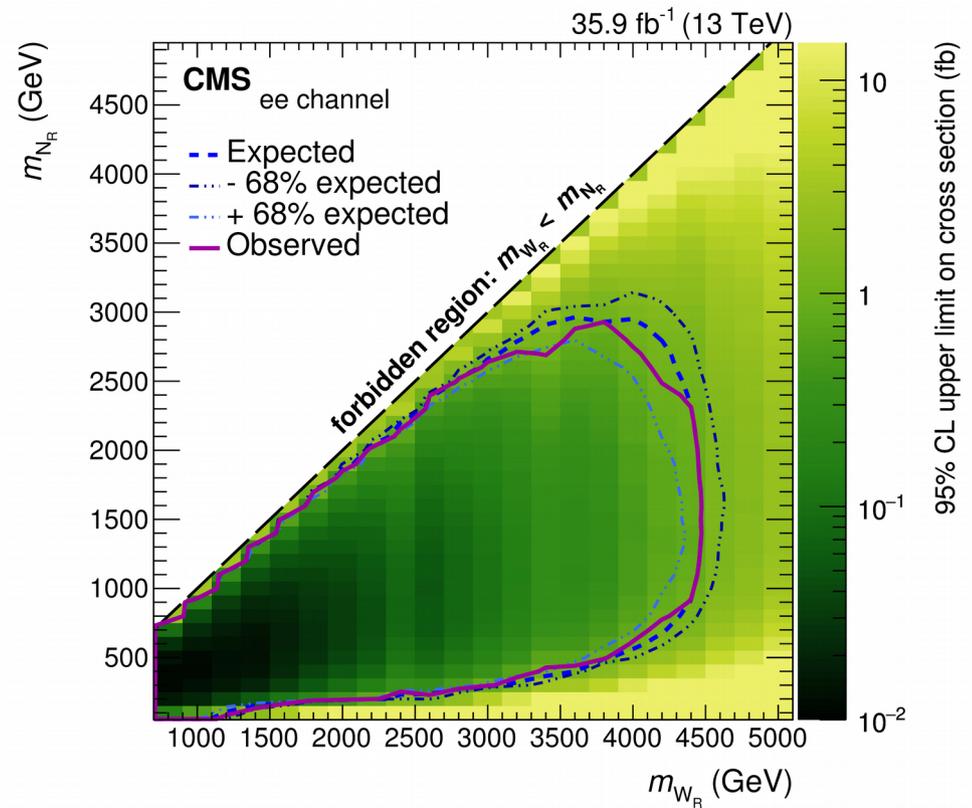
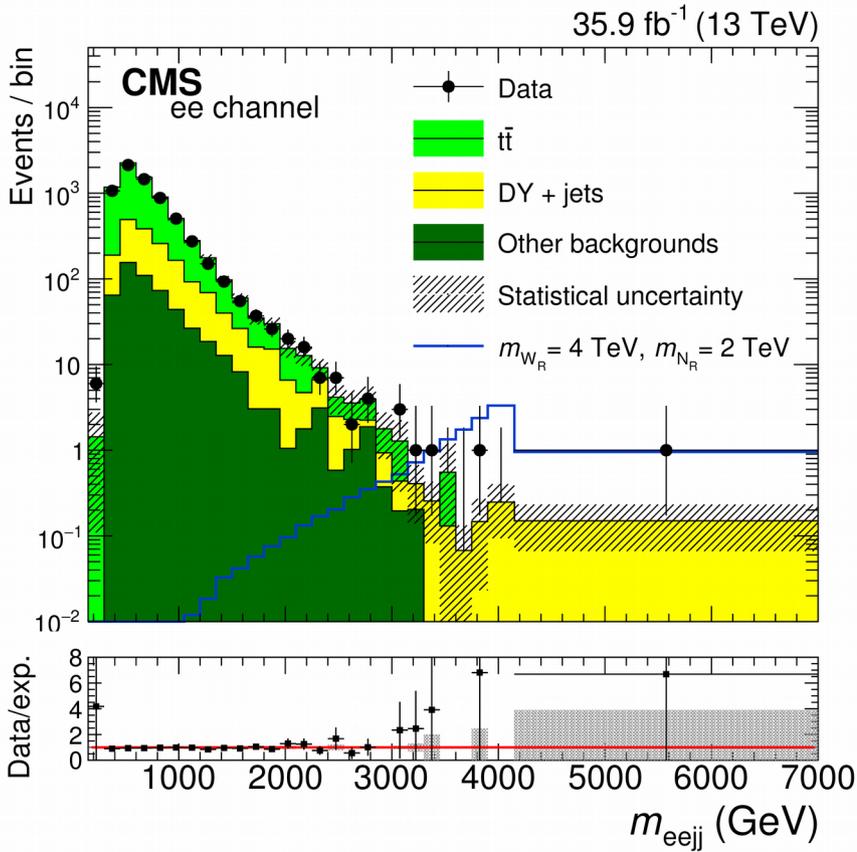
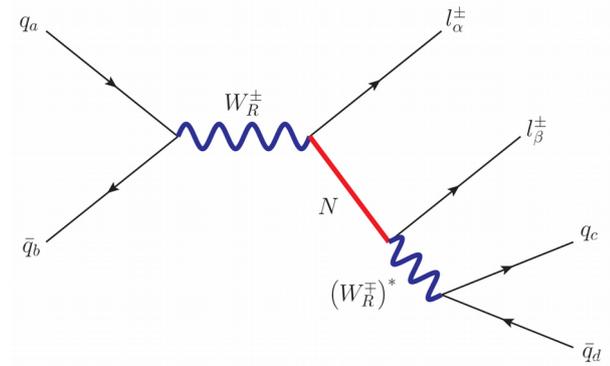


Prompt decays with CMS

2018

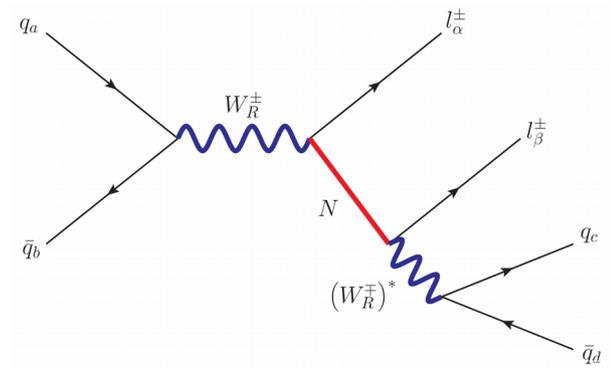
Searching for LRSM decay, final states with 2 leptons of same flavour (ee or $\mu\mu$) and any charge combination

Observable plotted is the mass of the combined final state eejj ($\mu\mu jj$ gives similar results)

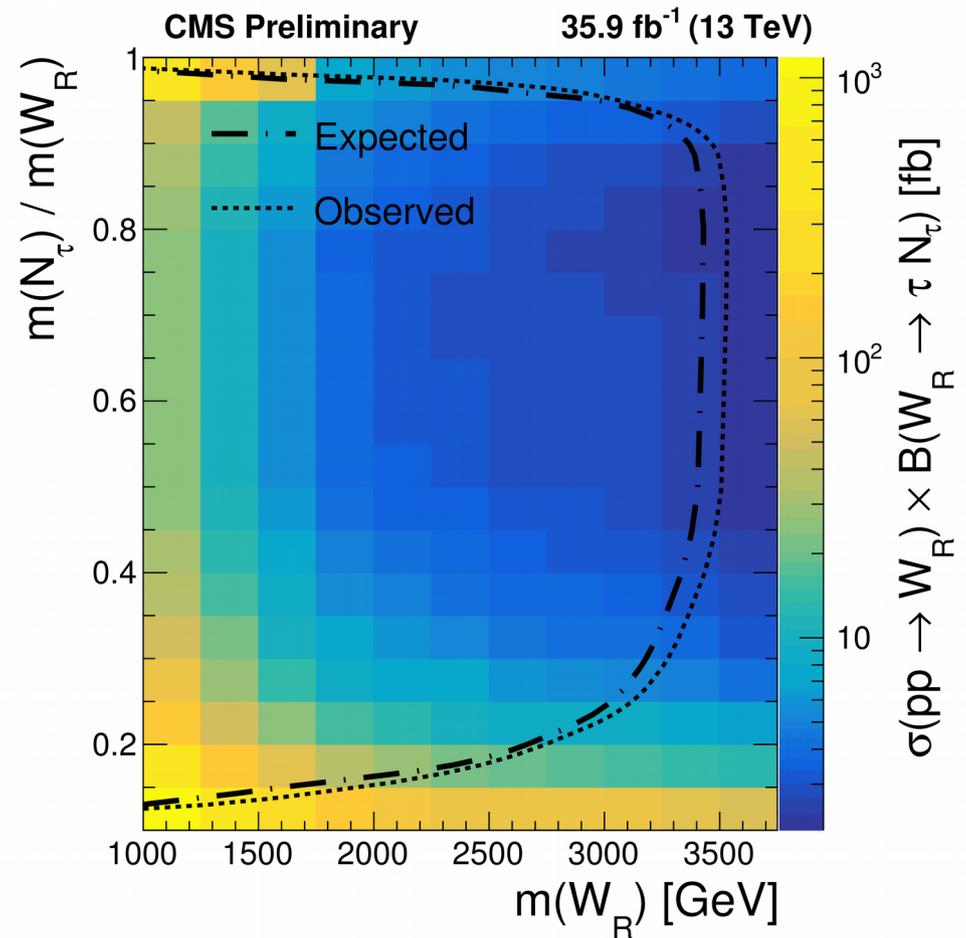
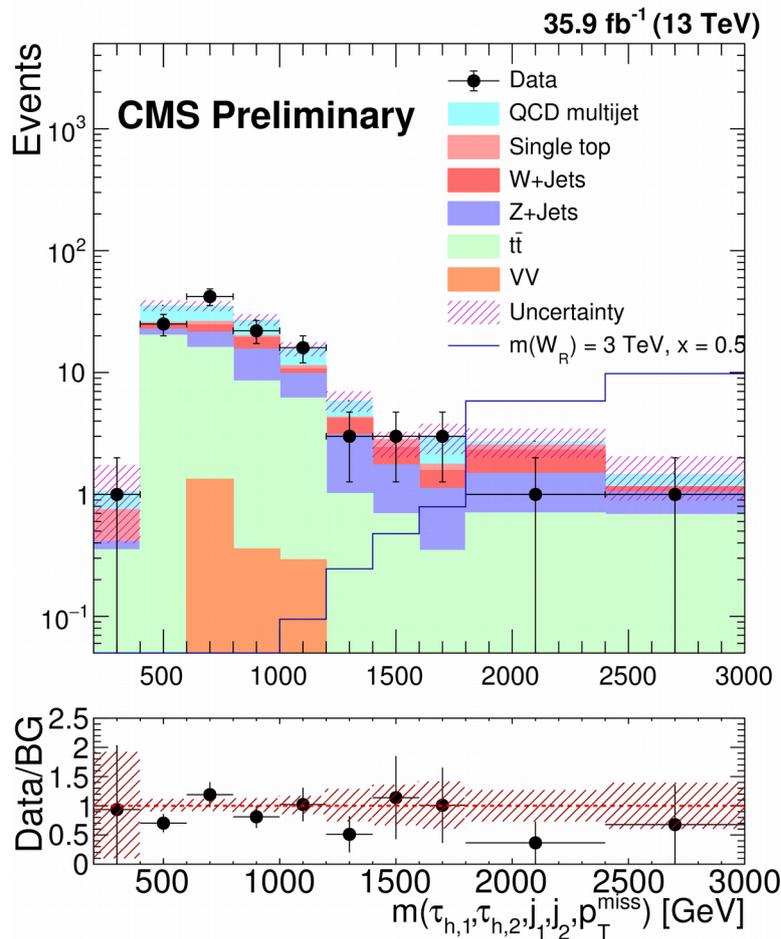


Prompt decays with CMS 2018

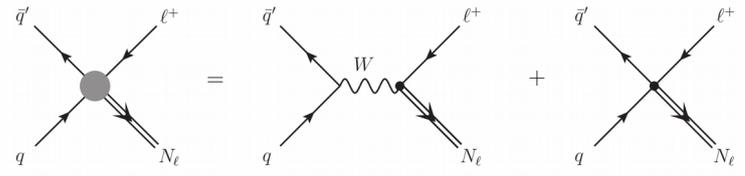
Searching for LRSM decay, final states with 2 τ -leptons decaying hadronically, and 2 jets from the W decay



Observable plotted is the combined mass of the final state particles and missing transverse momentum



Prompt decays with CMS

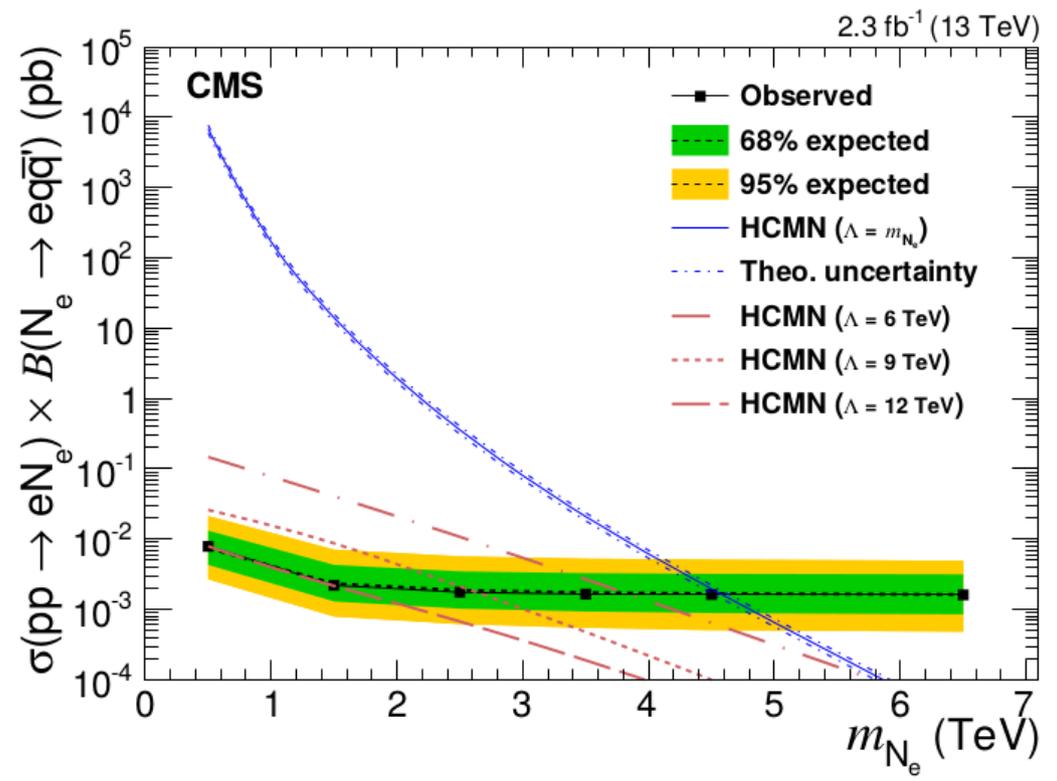
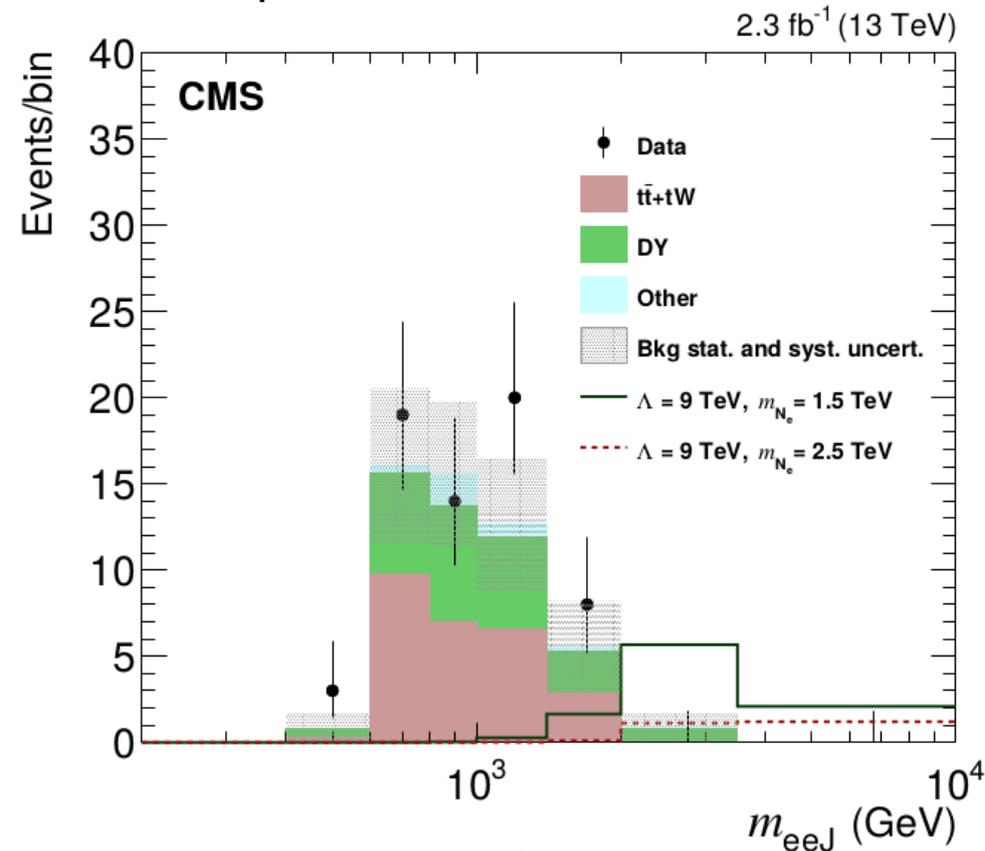


Searching for a heavy composite

Majorana neutrino decay, i.e. the neutrino here is not a fundamental particle

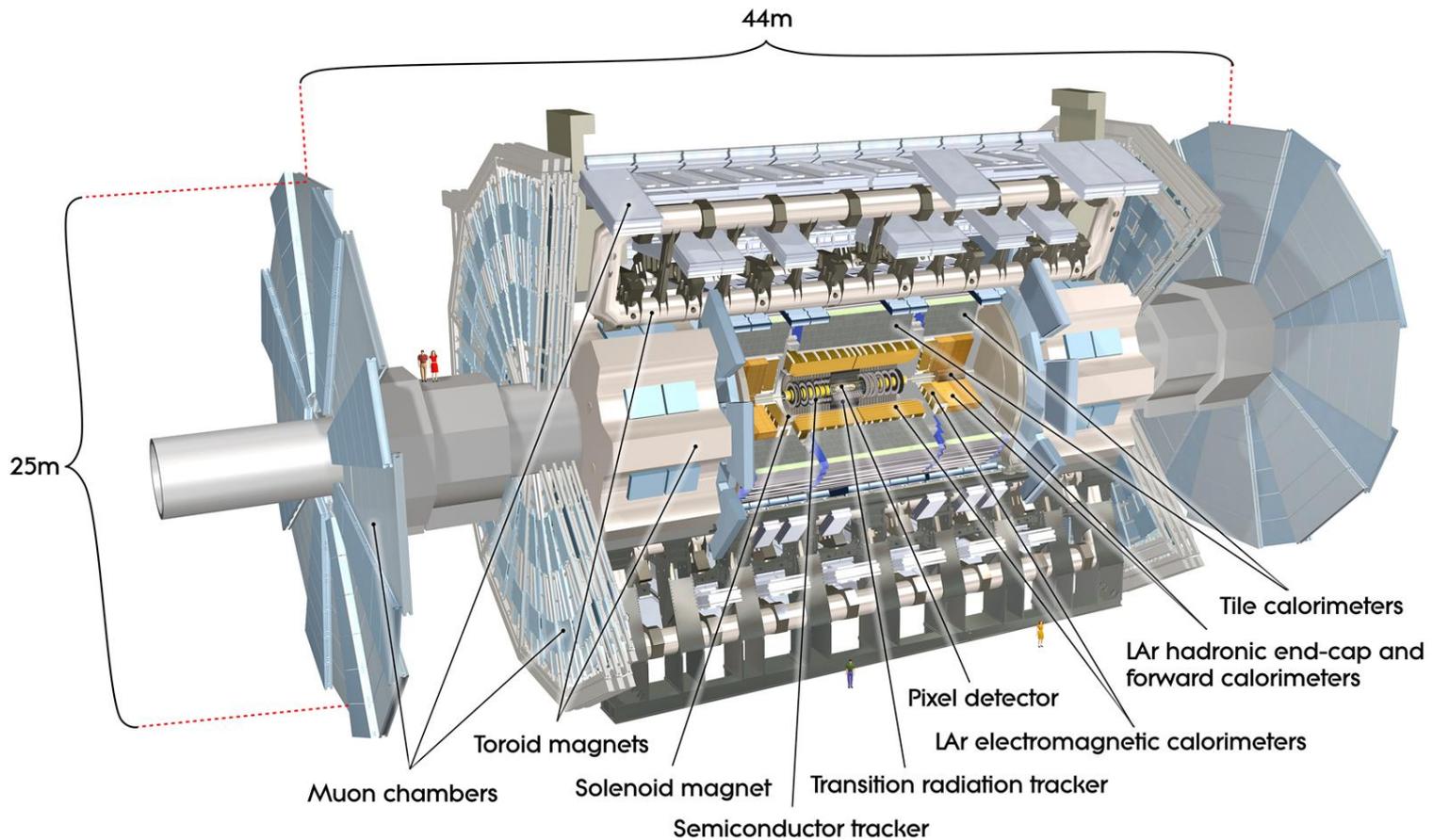
Final states with a pair of either electrons or muons, and a single large jet containing the collimated decay products of the 2 quarks

Limits compared with a variety of different choices of the compositeness scale Λ , and production cross-section dominated by contact-interaction component



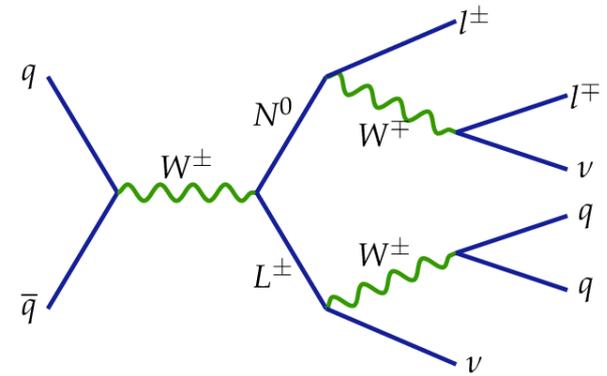
Searches with ATLAS

There's a recent result for Heavy Neutrino prompt decays from ATLAS, as well as a number of Dark Matter searches that are compatible with HN models



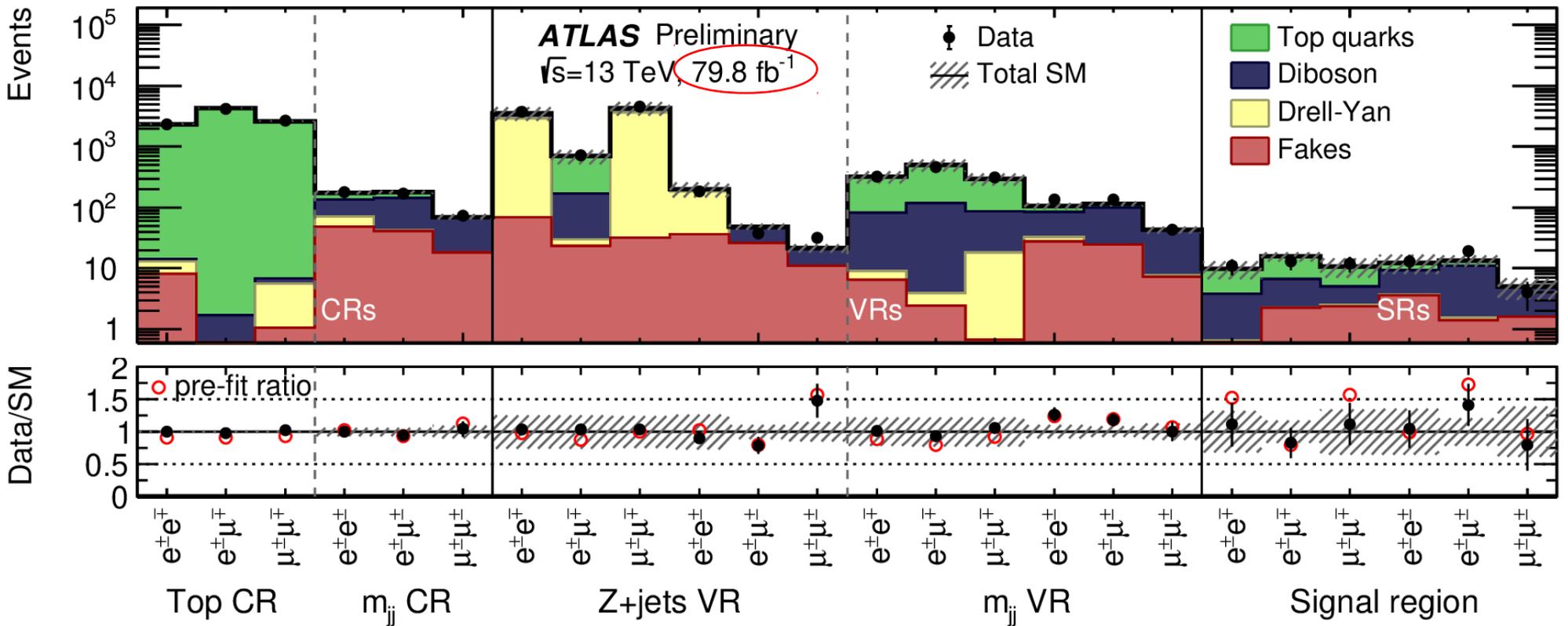
Prompt decays with ATLAS

2018



Searching for Type-3 seesaw mechanism decays (includes a heavy neutral Majorana lepton) in final states with 2 jets, MET, and ee , $e\mu$, or $\mu\mu$

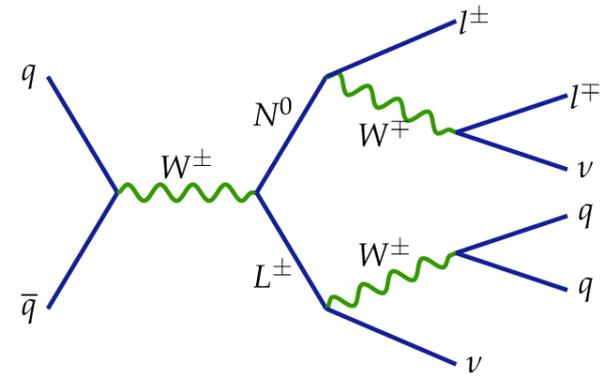
For more information see [Daniela's talk](#) on Tuesday



(CR = control region for background fitting; VR = validation region to test fitting)

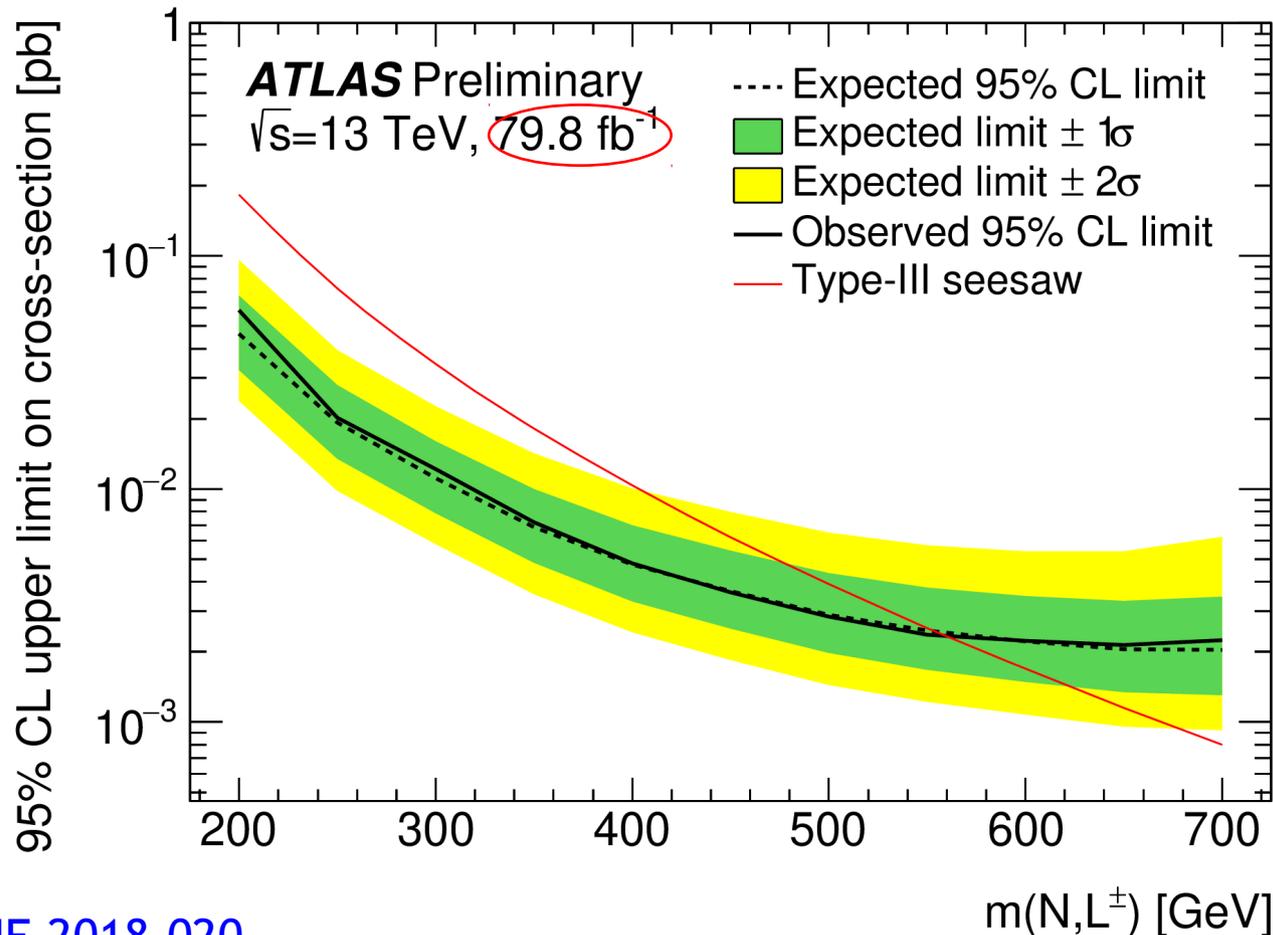
Prompt decays with ATLAS

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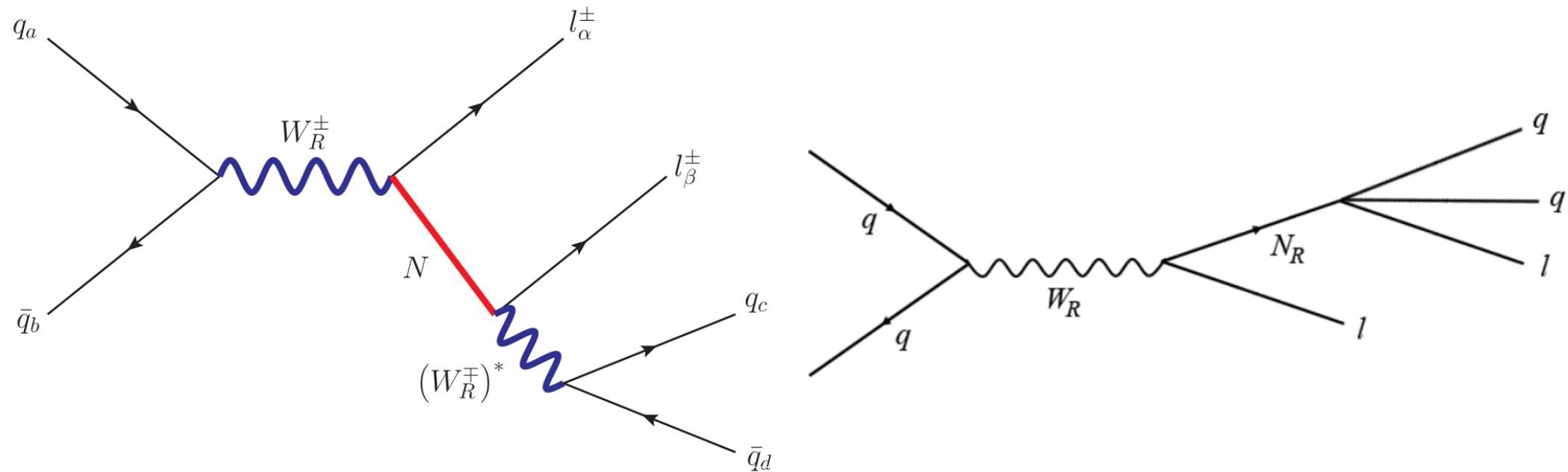
Final state leptons may have arisen from τ decays, exclusion limit assumes the branching fractions to all lepton flavours are equal



A word about modelling

The LRSM decay chain shown on the left has been extensively studied in LHC Run 1, and in the CMS results shown earlier

However, signal simulation in Pythia treated the NR decay as a single vertex (see right) which affects the decay kinematics in many samples



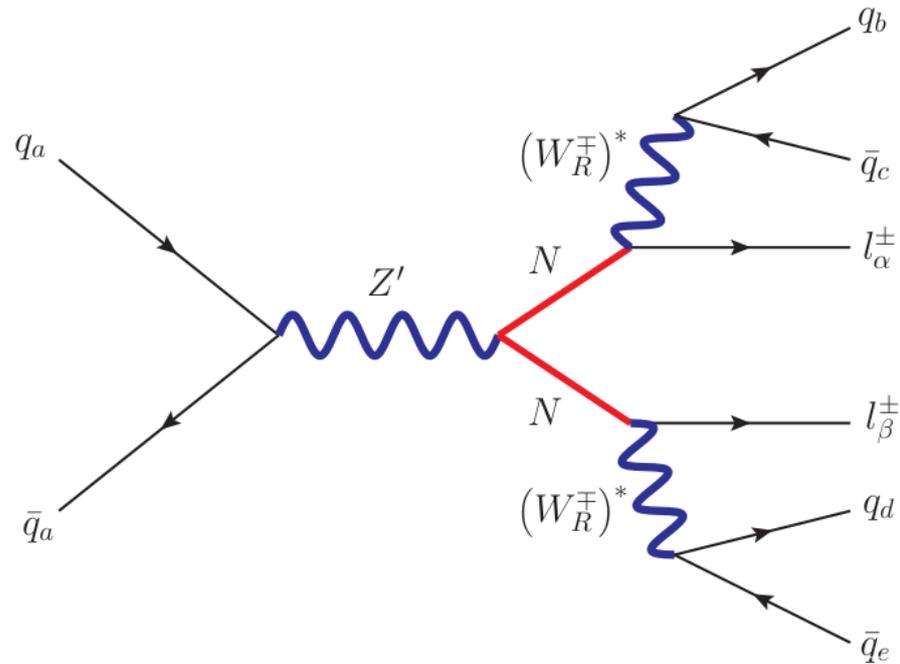
This Pythia model cannot generate signal samples with $m(N_R) > m(W_R)$

Searches are currently underway to set limits with the decay fully modelled in MadGraph, using Run 2 data

Dark Matter?

A stable heavy neutrino would be a potential Dark Matter candidate

Possibility to pair-produce heavy neutrinos from a Z' as shown in the LRSM-motivated diagram below:

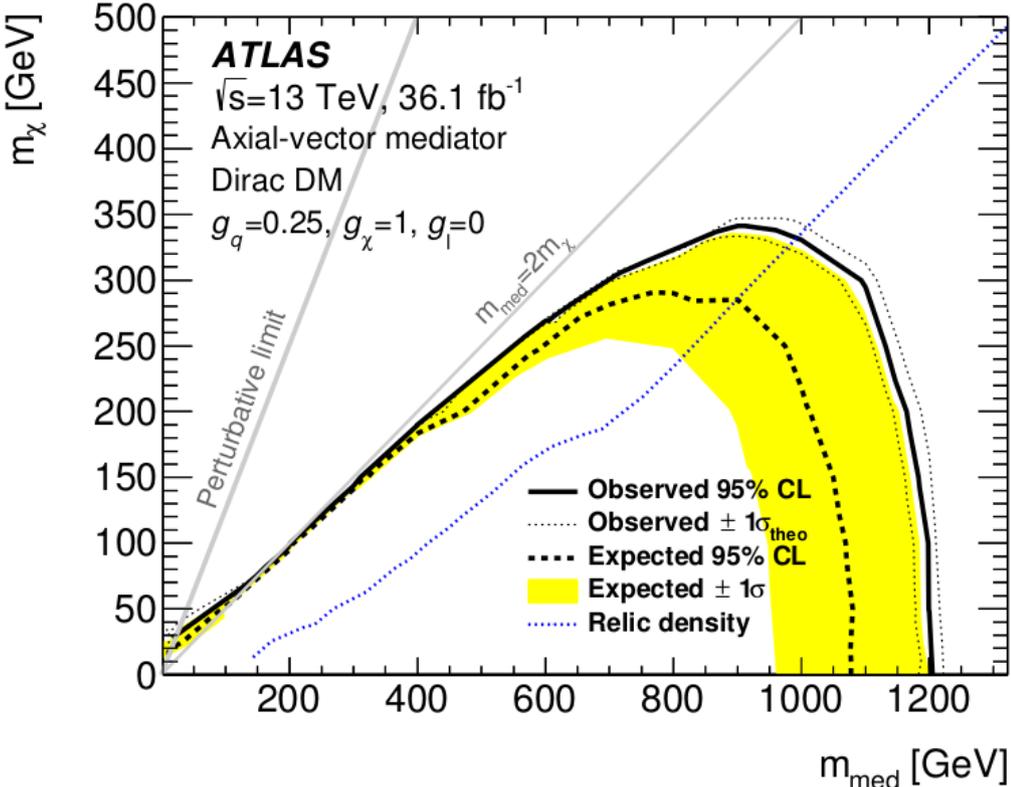
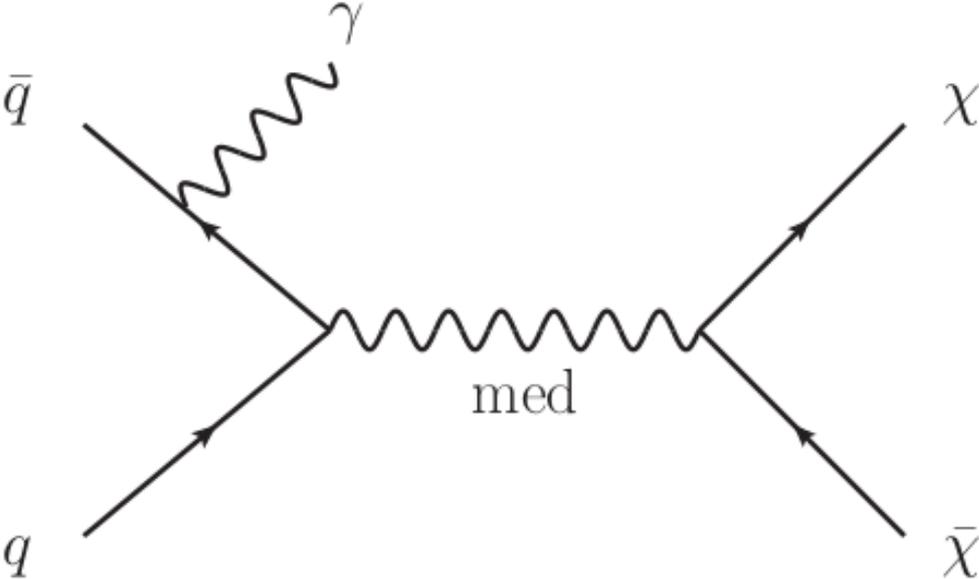


ν MSM introduces a stable heavy neutrino as a Dark Matter candidate (as well as the decays examined on previous slides)

Dark matter searches with ATLAS

A final state with large MET, tagged with a radiated photon

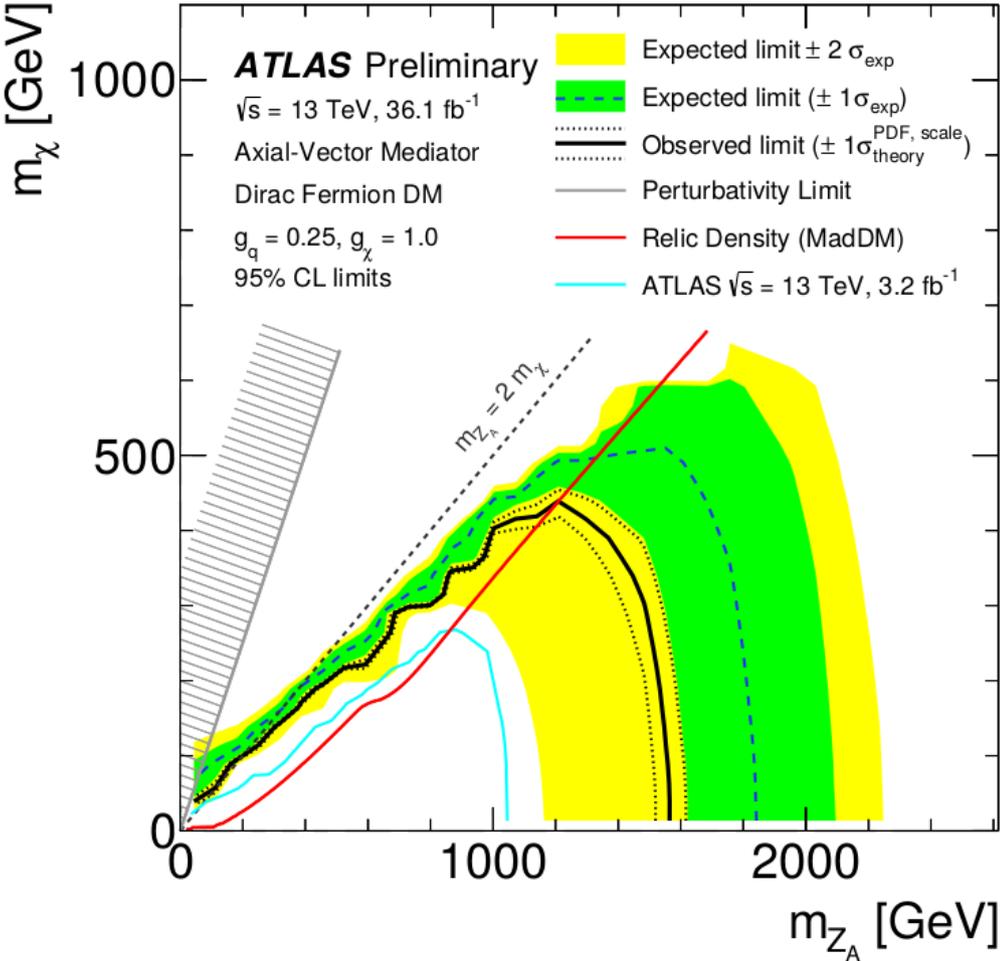
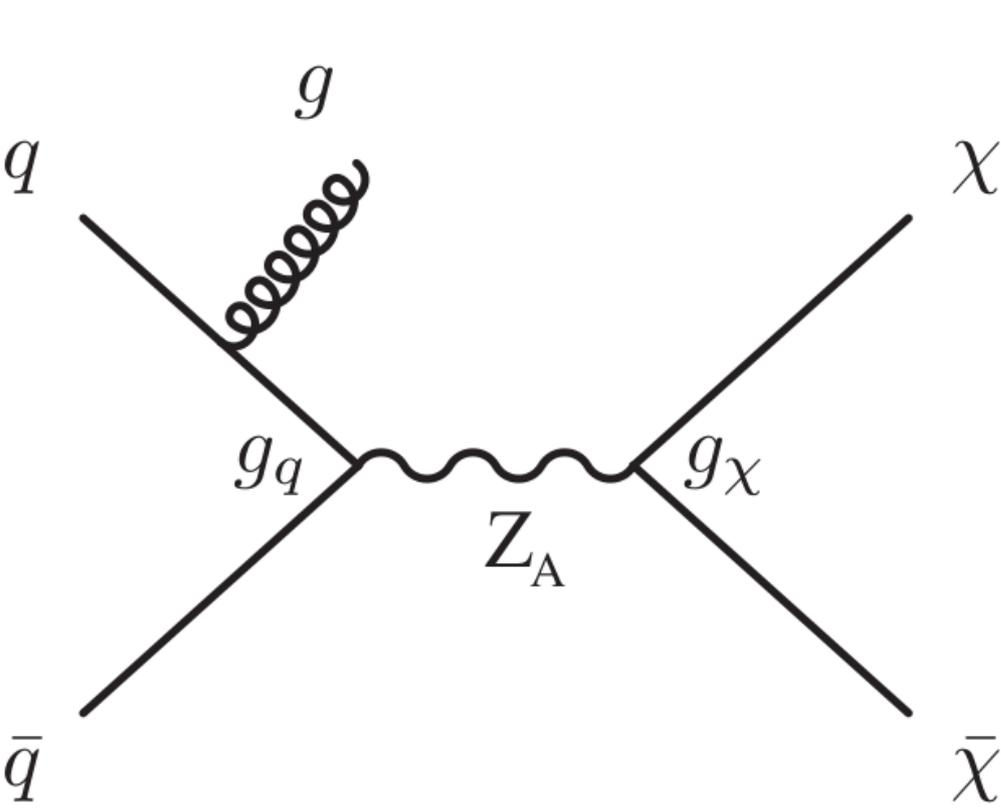
Limit set using a simplified model ([arXiv:1703.05703](https://arxiv.org/abs/1703.05703)) with an axial-vector propagator



Dark matter searches with ATLAS

A final state with large MET, tagged with a radiated gluon jet

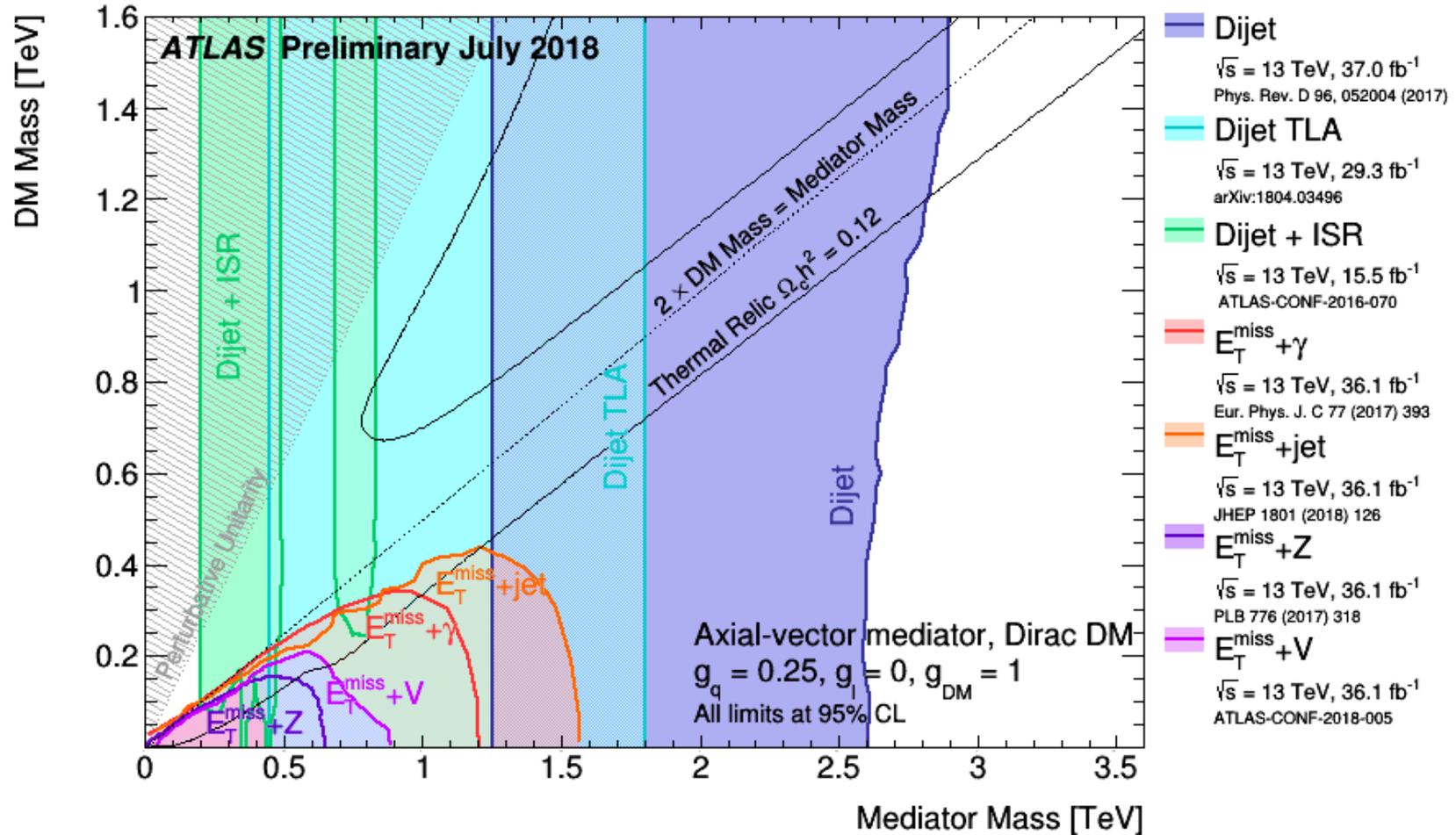
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Dark matter searches with ATLAS

2018

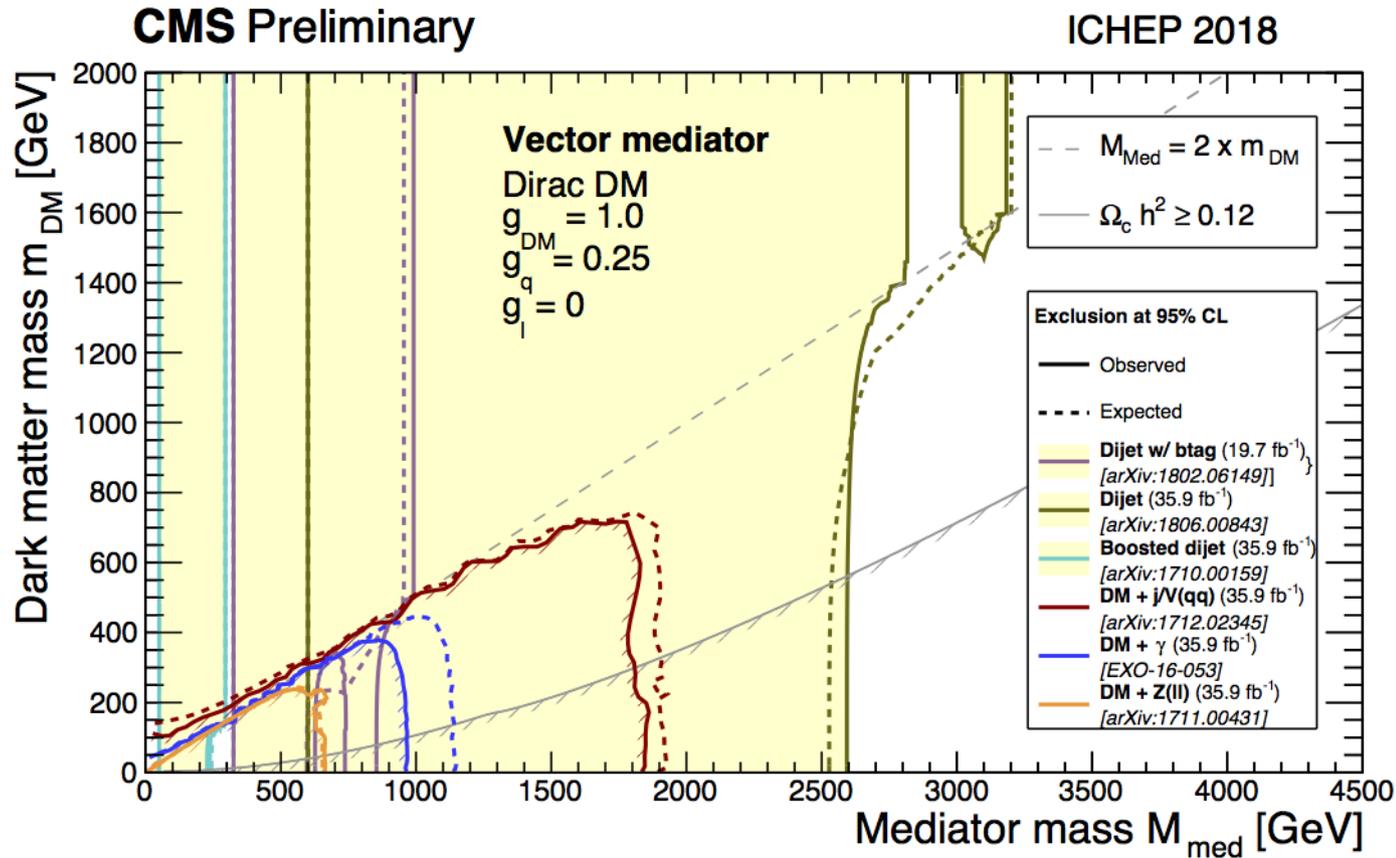
Combined exclusion limits have been produced for these searches



However, note that searches for dijet resonances are much more sensitive, since the propagator in our simplified model must couple to quarks

Dark matter searches with CMS

Combined exclusion limits have been produced for these searches



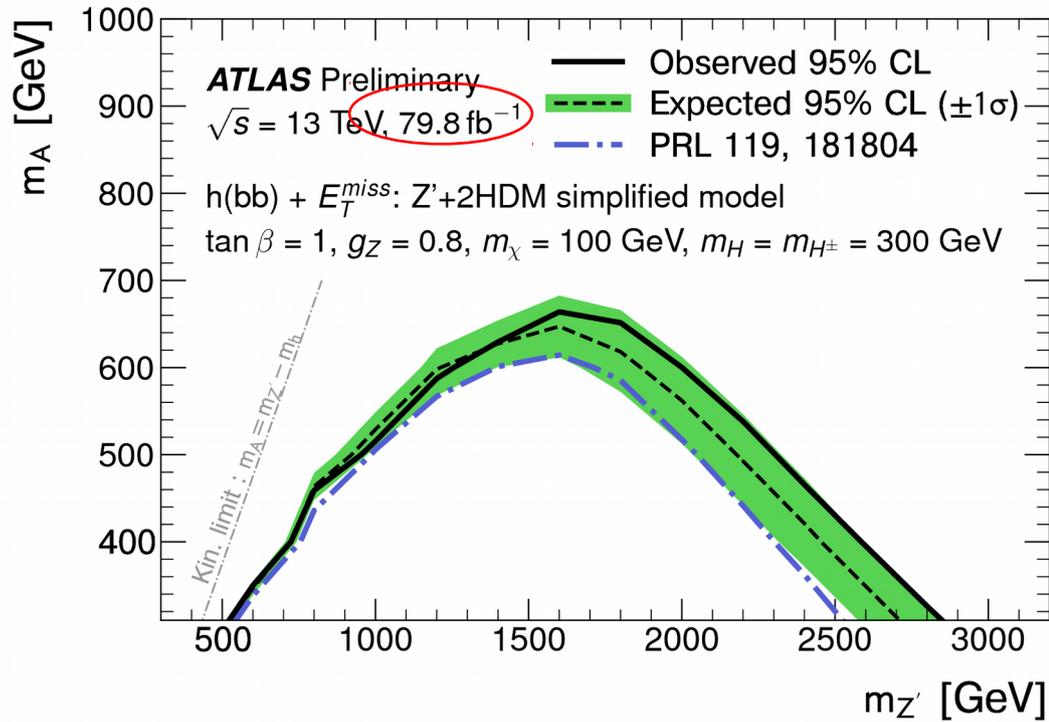
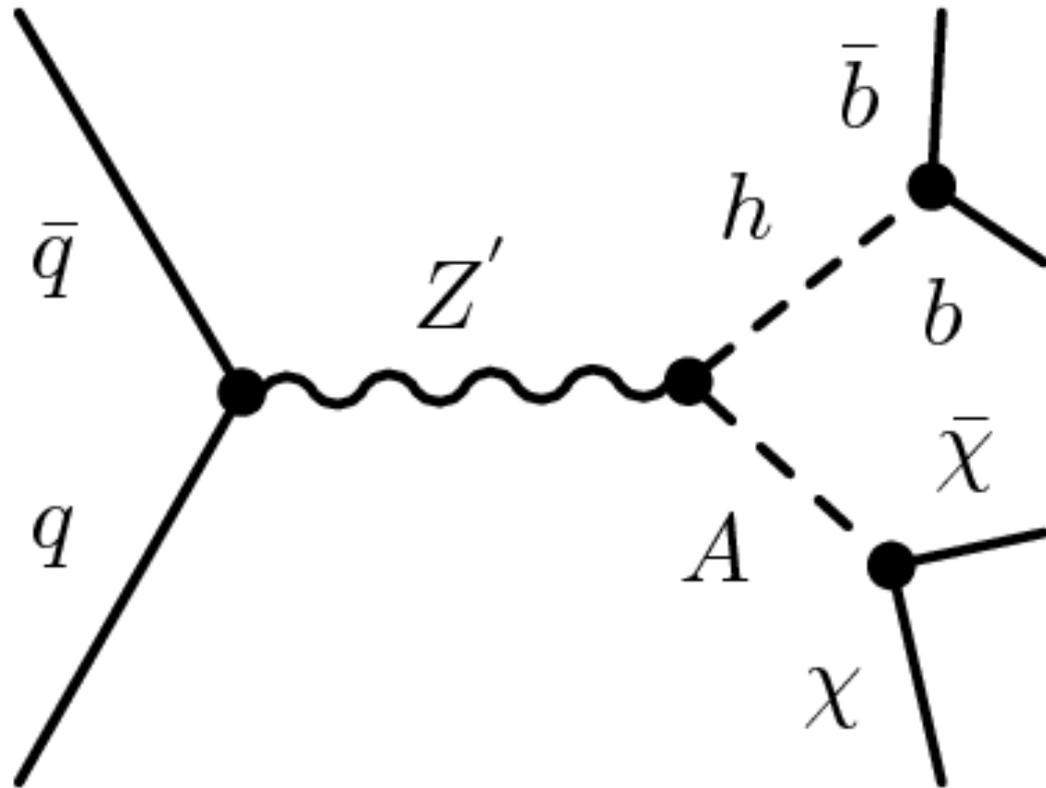
CMS produce a similar combined exclusion plot (here shown for Vector-like mediator)

Dark matter searches with ATLAS

2018

A final state with large MET, and a Higgs boson decaying to a pair of b quarks

Limit set using a 2 Higgs Doublet Model



Supersymmetry at the LHC

Now I'm definitely stretching the definition of "Heavy Neutrino"

There are a vast number of SUSY searches, past and ongoing, which mostly include neutralinos in the final state

ATLAS SUSY Searches* - 95% CL Lower Limits
May 2017

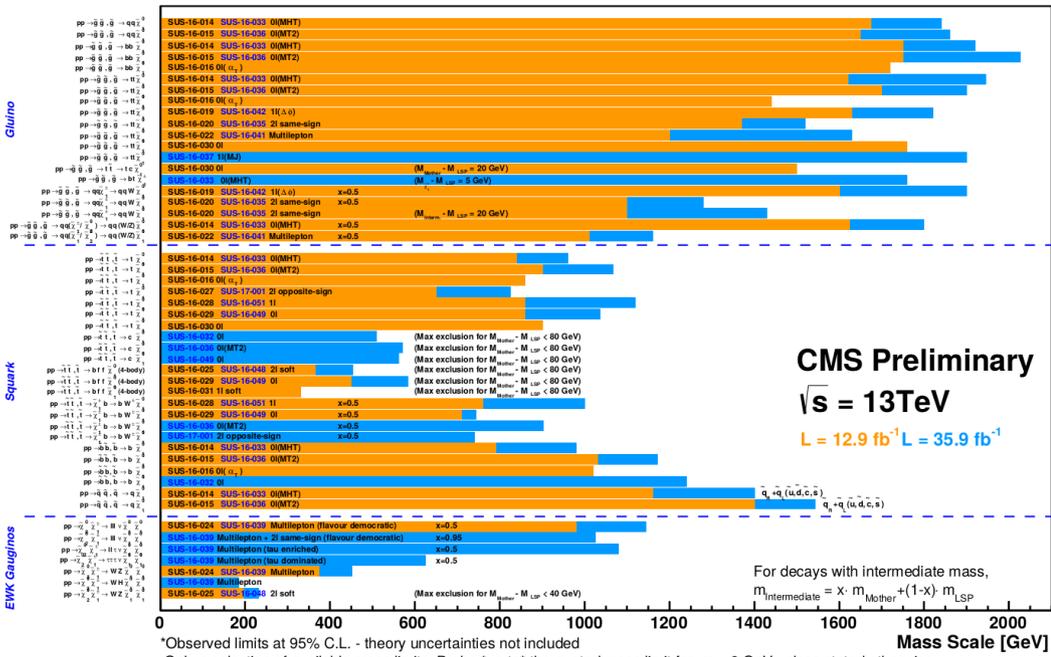
ATLAS Preliminary
 $\sqrt{s} = 7, 8, 13 \text{ TeV}$

Model	$\epsilon, \mu, \tau, \gamma$	Jets	E_{T}^{miss} [$\mathcal{L} d\mathcal{A}(\text{fb}^{-1})$]	Mass limit	$\sqrt{s} = 7, 8 \text{ TeV}$	$\sqrt{s} = 13 \text{ TeV}$	Reference
Inclusive Searches	MSUGRA/CMSSM	$0.3 e, \mu, \tau, \gamma$	2-10 jets+3b	Yes	20.3	1.85 TeV	m $_{\tilde{g}}$ =m $_{\tilde{t}}$
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	0	2-6 jets	Yes	36.1	1.57 TeV	m $_{\tilde{t}}$ >200 GeV, m $_{\tilde{b}}$ (1 st gen.)>200 GeV
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$ (compressed)	mono-jet	1-3 jets	Yes	3.2		m $_{\tilde{g}}$ >m $_{\tilde{t}}$ >5 GeV
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	0	2-6 jets	Yes	36.1	2.02 TeV	m $_{\tilde{t}}$ >200 GeV
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	0	2-6 jets	Yes	36.1	2.01 TeV	m $_{\tilde{t}}$ >200 GeV, m $_{\tilde{b}}$ (1 st gen.)>0.5(m $_{\tilde{t}}$ +m $_{\tilde{b}}$)
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	3 μ	4 jets	Yes	35.1	1.825 TeV	m $_{\tilde{t}}$ >400 GeV
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	0	7-11 jets	Yes	36.1	1.8 TeV	m $_{\tilde{t}}$ >400 GeV
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	1.2 $\tau + 0.1 \ell$	0-2 jets	Yes	3.2	2.0 TeV	m $_{\tilde{t}}$ >150 GeV
	GMSB (\tilde{L} NLSP)	2 γ	1 b	Yes	32	1.85 TeV	m $_{\tilde{g}}$ >950 GeV, c $_{\tilde{t}}$ (NLSP)>0.1 mm, $\mu=0$
	GGM (higgsino-bino NLSP)	2 γ	1 b	Yes	20.3	1.8 TeV	m $_{\tilde{g}}$ >680 GeV, c $_{\tilde{t}}$ (NLSP)>0.1 mm, $\mu=0$
3 rd gen. squarks and production	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	0	3 b	Yes	36.1	1.92 TeV	m $_{\tilde{t}}$ >800 GeV
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	0-1 μ	3 b	Yes	36.1	1.97 TeV	m $_{\tilde{t}}$ >200 GeV
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	1-1 μ	3 b	Yes	20.1	1.37 TeV	m $_{\tilde{t}}$ >300 GeV
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	0	2 b	Yes	36.1	950 GeV	m $_{\tilde{t}}$ >420 GeV
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	2 μ , (SS)	1-2 b	Yes	36.1	275-700 GeV	m $_{\tilde{t}}$ >200 GeV, m $_{\tilde{b}}$ (1 st gen.)>100 GeV
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	0-2 μ , 1 b	1-2 b	Yes	4.7/13.3	117-170 GeV	m $_{\tilde{t}}$ >200 GeV, m $_{\tilde{b}}$ (1 st gen.)>55 GeV
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	0-2 μ , 0-2 jets+2 b	Yes	20/306.1	90-198 GeV	1506/19616, ATLAS CONF-2017-020	m $_{\tilde{t}}$ >1 GeV
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	0	mono-jet	Yes	3.2	90-323 GeV	m $_{\tilde{t}}$ >5 GeV
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	2 μ , (Z)	1 b	Yes	20.3	150-600 GeV	m $_{\tilde{t}}$ >150 GeV
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	3 μ , (Z)	1 b	Yes	36.1	290-790 GeV	m $_{\tilde{t}}$ >60 GeV
EW direct	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	2 μ	0	Yes	36.1	90-440 GeV	m $_{\tilde{t}}$ >0
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	2 μ	0	Yes	36.1	710 GeV	m $_{\tilde{t}}$ >0, m $_{\tilde{b}}$ (1 st gen.)>0.5(m $_{\tilde{t}}$ +m $_{\tilde{b}}$)
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	2 μ	0	Yes	36.1	760 GeV	m $_{\tilde{t}}$ >0, m $_{\tilde{b}}$ (1 st gen.)>0.5(m $_{\tilde{t}}$ +m $_{\tilde{b}}$)
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	2 μ	0	Yes	36.1	1.16 TeV	m $_{\tilde{t}}$ >0, m $_{\tilde{b}}$ (1 st gen.)>0.5(m $_{\tilde{t}}$ +m $_{\tilde{b}}$)
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	2 μ	0	Yes	36.1	590 GeV	m $_{\tilde{t}}$ >0, m $_{\tilde{b}}$ (1 st gen.)>0.5(m $_{\tilde{t}}$ +m $_{\tilde{b}}$)
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	2 μ	0	Yes	36.1	270 GeV	m $_{\tilde{t}}$ >0, m $_{\tilde{b}}$ (1 st gen.)>0.5(m $_{\tilde{t}}$ +m $_{\tilde{b}}$)
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	2 μ	0	Yes	20.3	635 GeV	m $_{\tilde{t}}$ >0, m $_{\tilde{b}}$ (1 st gen.)>0.5(m $_{\tilde{t}}$ +m $_{\tilde{b}}$)
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	2 μ	0	Yes	20.3	440 GeV	m $_{\tilde{t}}$ >0, m $_{\tilde{b}}$ (1 st gen.)>0.5(m $_{\tilde{t}}$ +m $_{\tilde{b}}$)
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	2 μ	0	Yes	20.3	115-370 GeV	c $_{\tilde{t}}$ <1 mm
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	2 μ	0	Yes	20.3	590 GeV	c $_{\tilde{t}}$ <1 mm
Long-lived particles	Direct $\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$ prod. long-lived \tilde{g}	Disapp. trk	1 jet	Yes	36.1	430 GeV	m $_{\tilde{g}}$ >160 MeV, $\tau(\tilde{g})>0.2$ ns
	Direct $\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$ prod. long-lived \tilde{q}	dE/dx trk	Yes	18.4	495 GeV	m $_{\tilde{q}}$ >160 MeV, $\tau(\tilde{q})>15$ ns	
	Stable, stopped \tilde{g} R-hadron	trk	0	Yes	27.9	850 GeV	m $_{\tilde{g}}$ >100 GeV, 10 μ s < c $_{\tilde{g}}$ < 1000 s
	Stable \tilde{g} R-hadron	trk	-	3.2		1.56 TeV	m $_{\tilde{g}}$ >100 GeV, $\tau > 10$ ns
	Metastable \tilde{g} R-hadron	dE/dx trk	-	3.2		1.57 TeV	m $_{\tilde{g}}$ >100 GeV, $\tau > 10$ ns
	GMSB, stable \tilde{g} , long-lived \tilde{g}	trk	-	19.1	537 GeV	10-larg<50	
	GMSB, $\tilde{g} \rightarrow \tilde{g},$ long-lived \tilde{g}	trk	-	20.3	440 GeV	1-c $_{\tilde{g}}$ >3 ns, SPSB model	
	GGM (\tilde{g} NLSP) weak prod. $\tilde{g} \rightarrow \tilde{g} + \gamma$	displ. ex $_{\tilde{g}}$ /mu	-	20.3	1.0 TeV	7-c $_{\tilde{g}}$ >7.740 mm, m $_{\tilde{g}}$ >1.3 TeV	
	GGM (\tilde{g} NLSP) weak prod. $\tilde{g} \rightarrow \tilde{g} + \gamma$	displ. vtx + ex $_{\tilde{g}}$	-	20.3	1.0 TeV	6-c $_{\tilde{g}}$ >7.480 mm, m $_{\tilde{g}}$ >1.1 TeV	
	GGM (\tilde{g} NLSP) weak prod. $\tilde{g} \rightarrow \tilde{g} + \gamma$	displ. vtx + ex $_{\tilde{g}}$	-	20.3	590 GeV	1504/05162	
RPV	LFV $\tilde{g}, \tilde{q} \rightarrow q\tilde{q} + X, \tilde{g}, \tilde{q} \rightarrow q\tilde{q} + \mu\tau$	ex $_{\tilde{g}}$ trk	-	3.2		1.9 TeV	$A_{111} > 0.1, A_{121} > 0.07$
	Bilinear RPV CMSSM	2 μ , (SS)	0-3 b	Yes	20.3	1.45 TeV	m $_{\tilde{g}}$ =m $_{\tilde{t}}$, c $_{\tilde{t}}$ <1 mm
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	4 μ	-	Yes	13.3	1.14 TeV	m $_{\tilde{t}}$ >400 GeV, $A_{113} > 0$ (6, 1, 2)
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	3 μ , $\tau + \gamma$	-	Yes	20.3	450 GeV	m $_{\tilde{t}}$ >200 GeV, $A_{113} > 0$
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	0	4-5 large-R jets	-	14.8	1.08 TeV	BR(\tilde{g})-BR(\tilde{b})>0%
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	0	4-5 large-R jets	-	14.8	1.55 TeV	m $_{\tilde{t}}$ >800 GeV
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	1 μ , 8-10 jets+0-4 b	-	36.1	2.1 TeV	1.85 TeV	m $_{\tilde{t}}$ >1 TeV, $A_{123} > 0$
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	1 μ , 8-10 jets+0-4 b	-	36.1	1.85 TeV	1.85 TeV	m $_{\tilde{t}}$ >1 TeV, $A_{123} > 0$
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	0	2 jets + 2 b	-	15.4	410 GeV	BR(\tilde{g})-BR(\tilde{b})>20%
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{q}$	2 μ , 2 b	-	36.1	450-510 GeV	0.4-1.45 TeV	BR(\tilde{g})-BR(\tilde{b})>20%
Other	Scalar charm, $\tilde{c} \rightarrow c\tilde{c}$	0	2 c	Yes	20.3	510 GeV	m $_{\tilde{c}}$ >200 GeV

*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

Selected CMS SUSY Results* - SMS Interpretation

ICHEP '16 - Moriond '17

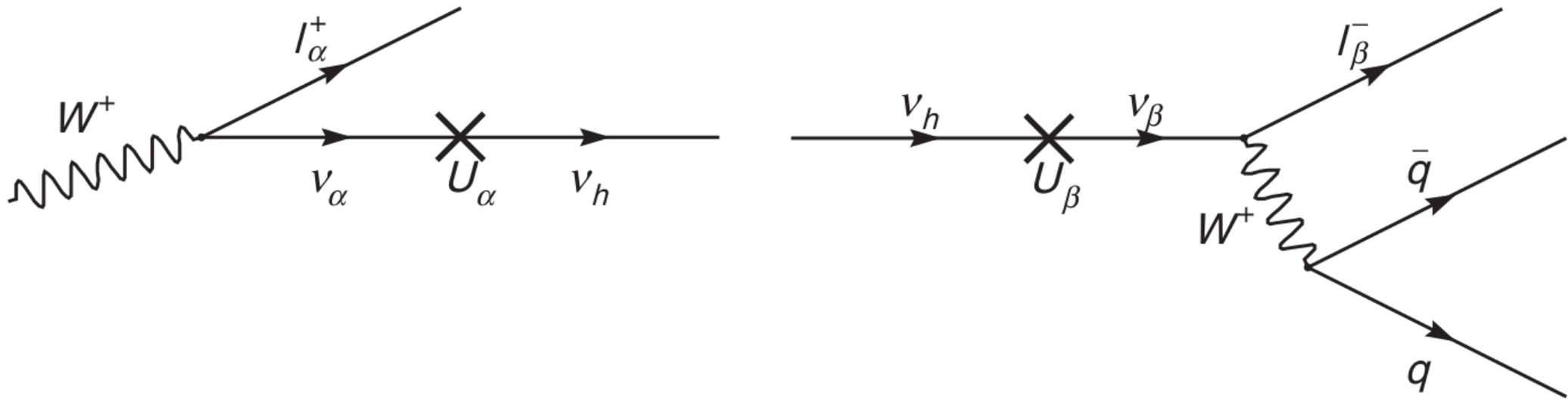


*Observed limits at 95% C.L. - theory uncertainties not included
Only a selection of available mass limits. Probe 'up to' the quoted mass limit for $m_{\tilde{g}} = 0$ GeV unless stated otherwise

Not planning to go into this in any depth, suffice it to say that we're still looking

Displaced vertices

Somewhere in-between prompt decays and long-lived Dark Matter candidates are heavy neutrinos likely to decay after travelling a significant distance



These models don't even rely on the seesaw mechanism - simply adding an extra neutrino mass eigenstate is enough

However, throwing in some sensible values for the coupling strength ($U_\alpha \sim 10^{-4}$), a heavy neutrino with mass 1 GeV decays after $O(10)$ meters

Displaced vertices

There are lots of displaced vertex searches past and ongoing at LHC collider experiments (ATLAS, CMS and LHCb at least)

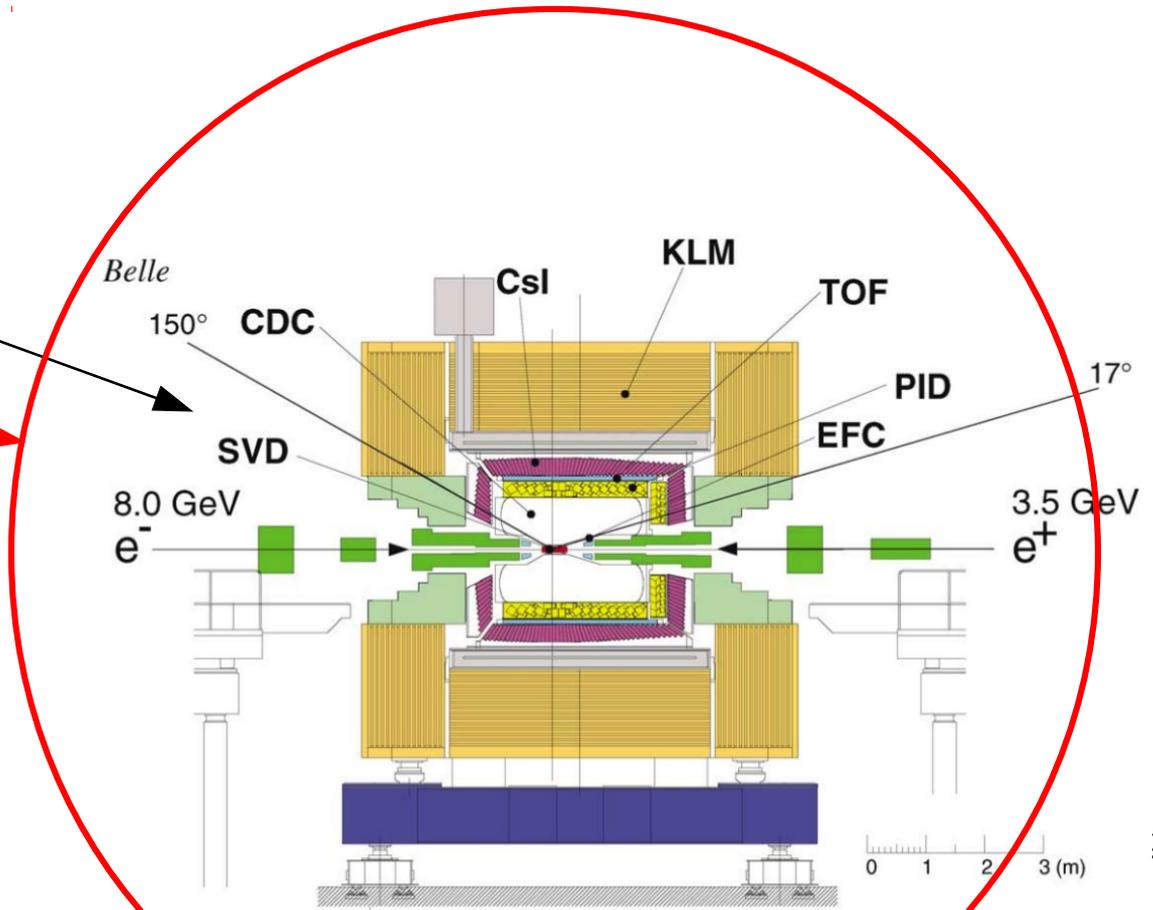
Searches rarely target Heavy Neutrinos, although should be possible to search in the $O(10 \text{ GeV})$ mass region

Credit to Belle (and others) for looking: [arXiv:1301.1105](https://arxiv.org/abs/1301.1105)

Still, the problem remains

Belle detector

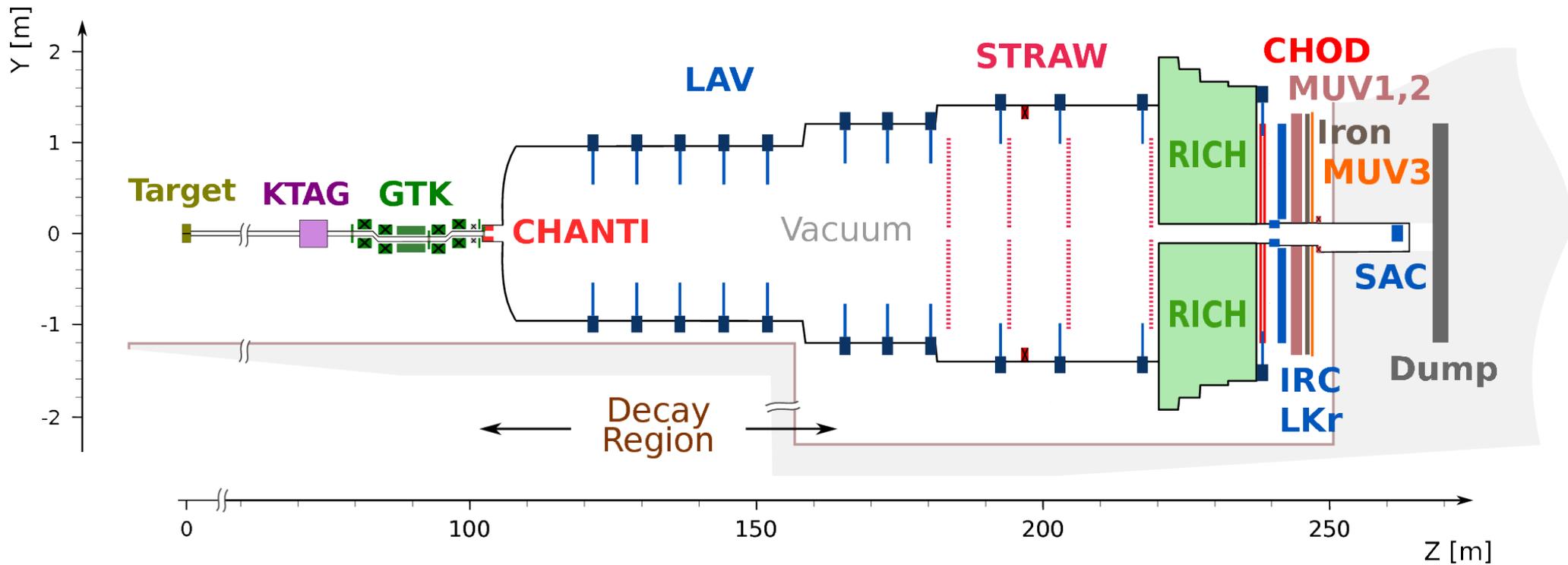
10m radius from
interaction point
(roughly)



NA62 experiment

Rather than trying to instrument an unfeasibly large volume, beam dump experiments need only instrument the region along the beam axis

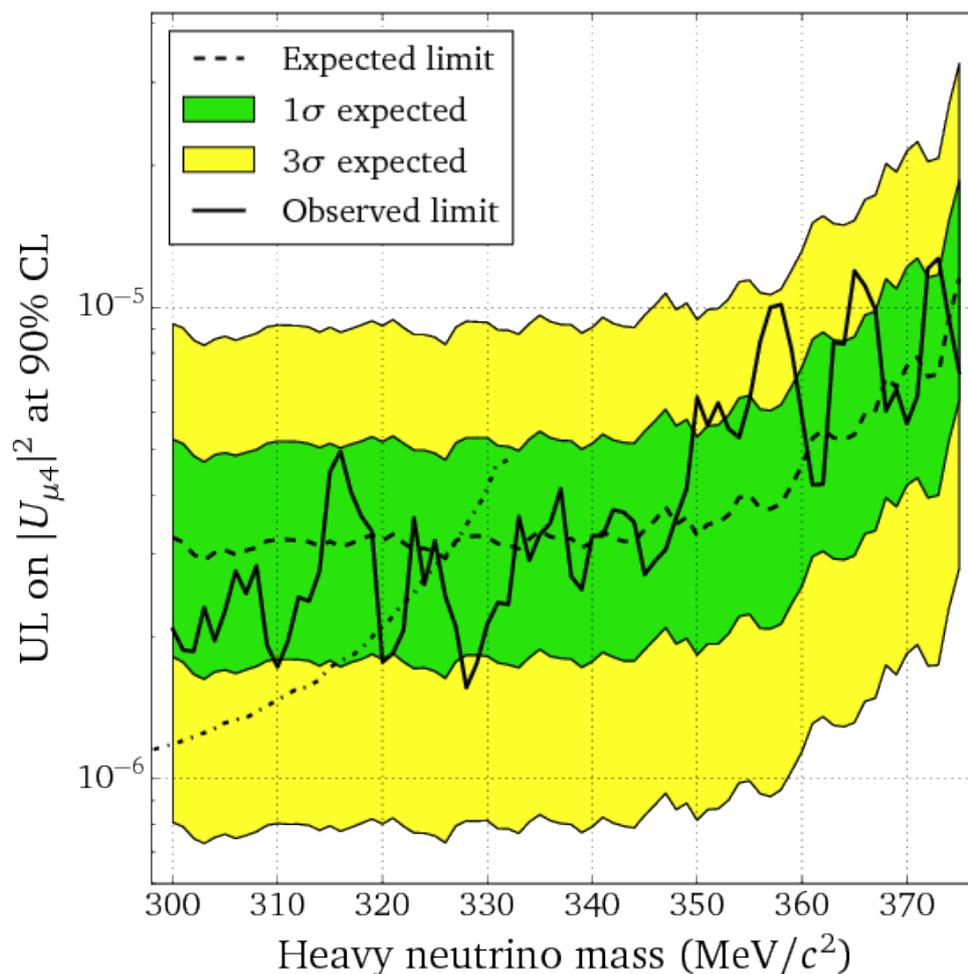
The NA62 experiment at CERN uses 400 GeV protons from the SPS (the same accelerator that feeds the LHC) impinging on a Berillium target



NA62 experiment

NA62 already has a result searching for heavy neutrinos (300-375 MeV) originating from Kaon decays: [arXiv:1705.07510](https://arxiv.org/abs/1705.07510)

Here the heavy neutrino is not expected to decay within 10km of the target, and so the search was simply for single muon and missing mass

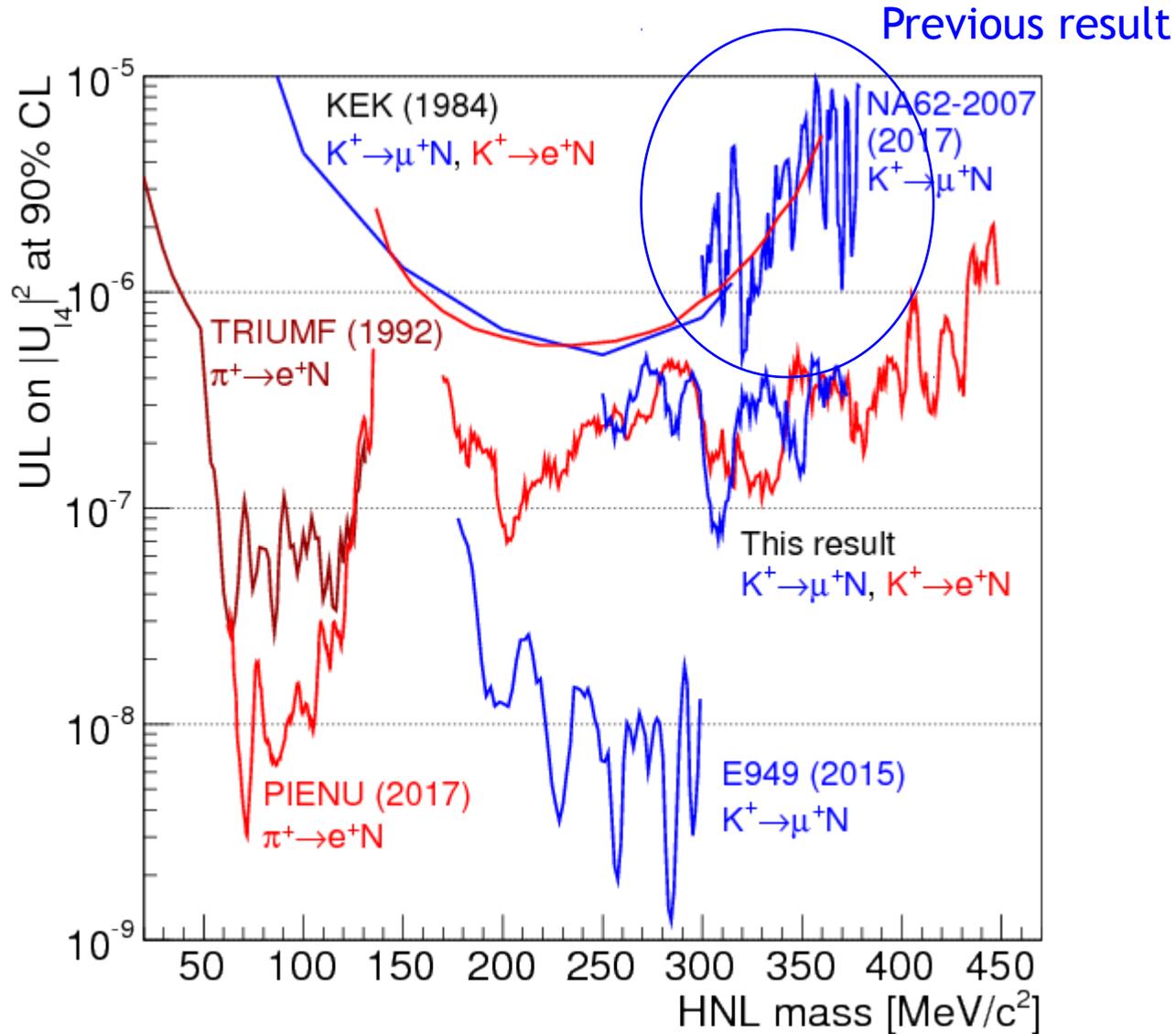


NA62 experiment

2018

Updated result published this year: [arXiv:1712.00297](https://arxiv.org/abs/1712.00297)

Newer dataset, and including electron channel



Summary

Models for prompt decays of Heavy Neutrinos have been (and are still being) studied at the LHC experiments

Heavy Neutrinos - if stable - could be revealed by Dark Matter searches at the LHC

Displaced vertex searches are an evolving field for collider experiments targeting prompt decays, but potential here for further study of Heavy Neutrinos

Beam dump experiments far more sensitive to Heavy Neutrino decays when the couplings are weaker, and so the vertices are more displaced

New experiment planned - SHiP

LHC upgrade schedule on track - higher luminosity for all experiments

Still bigger colliders are possible...

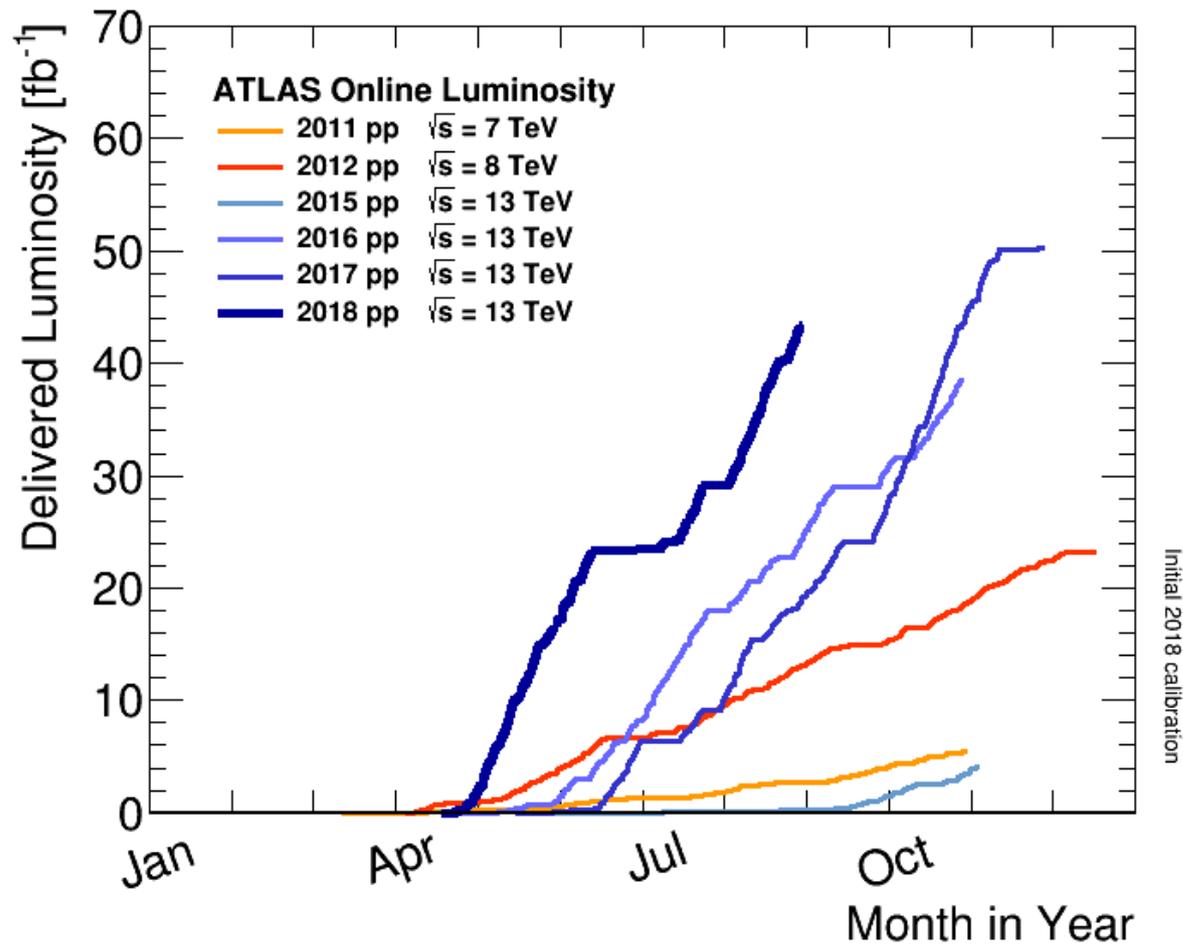
BACKUP

Large Hadron Collider

The LHC is taking pp collision data at 13 TeV

Results with 80 fb^{-1} (2015-2017) are available now

- Full Run 2 data analyses ongoing



Large Hadron Collider

LHC schedule suggests we have 1-2 more months of data-taking this year, taking us to the end of Run 2, and the long shutdown for the LHC upgrade

We are here

	July			Aug				Sep					
Wk	27	28	29	30	31	32	33	34	35	36	37	38	39
Mo	$\beta^* = 90\text{ m}^2$	9	16	23	30	6	13	20	27	3	10	17	24
Tu													
We				MD 2								TS2	
Th										Jeune G.			
Fr											MD 3		
Sa													
Su													

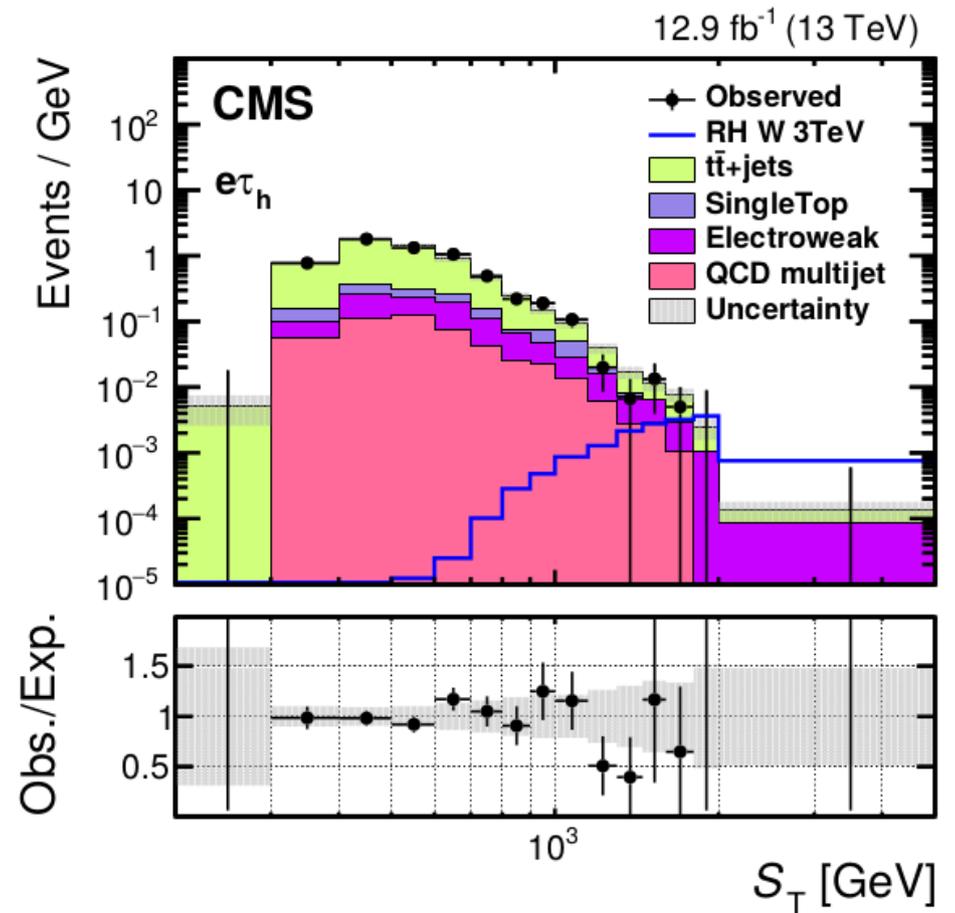
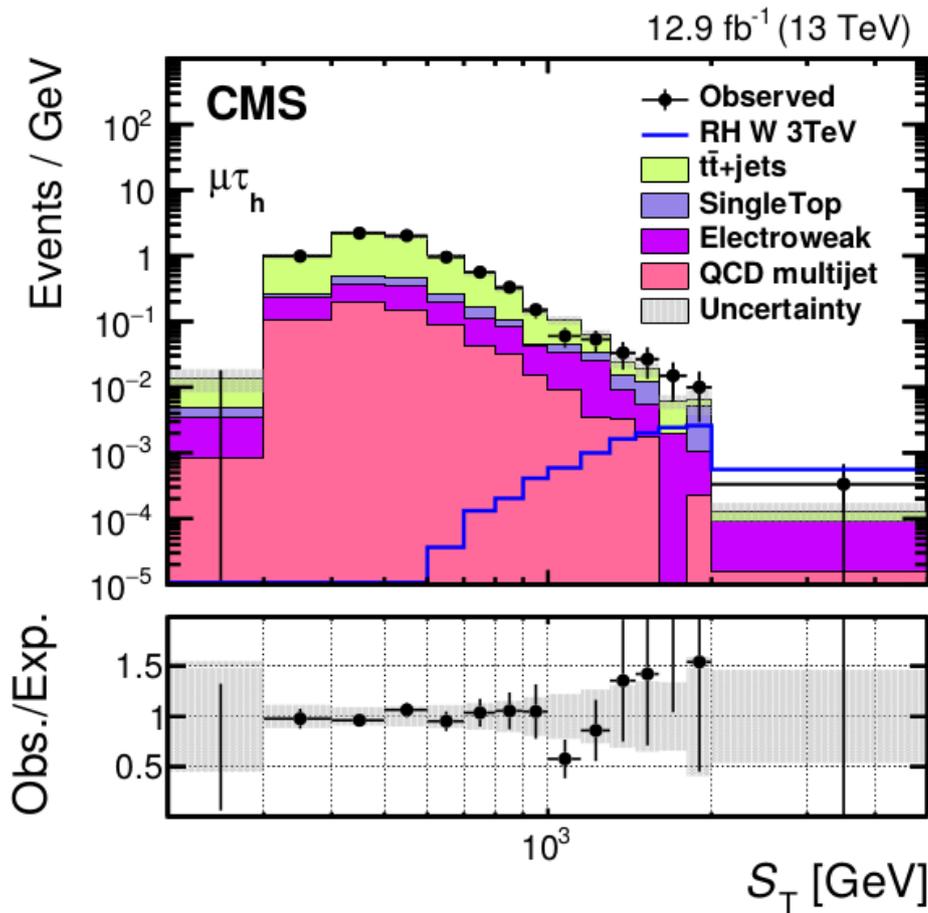
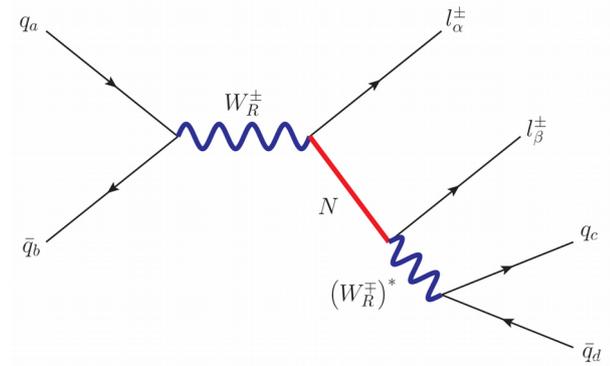
	Oct			Nov				Dec					
Wk	40	41	42	43	44	45	46	47	48	49	50	51	52
Mo	1	8	15	22	29	5	12	19	26	3	10	17	Xmas 24
Tu					MD 4	Ion setting up		MD 5					
We		Special physics run								Powering Tests Magnet Training			
Th					TS3								Long Shutdown 2
Fr													
Sa													
Su				MD 4									

-  Technical Stop
-  Special physics runs (Indicative - schedule to be established)
-  Powering tests
-  Machine development
-  Machine check out
-  Scrubbing (Indicative - dates to be established)
-  Recommissioning with beam
-  Pb - Pb Ion physics run
-  Interleaved commissioning & Intensity ramp up
-  Pb Ion Setting up
-  Proton physics run
-  LINAC 3 Pb oven re-fill

Prompt decays with CMS

Searching for LRSM decay, final states with 2 τ -leptons where one decays hadronically and one leptonically

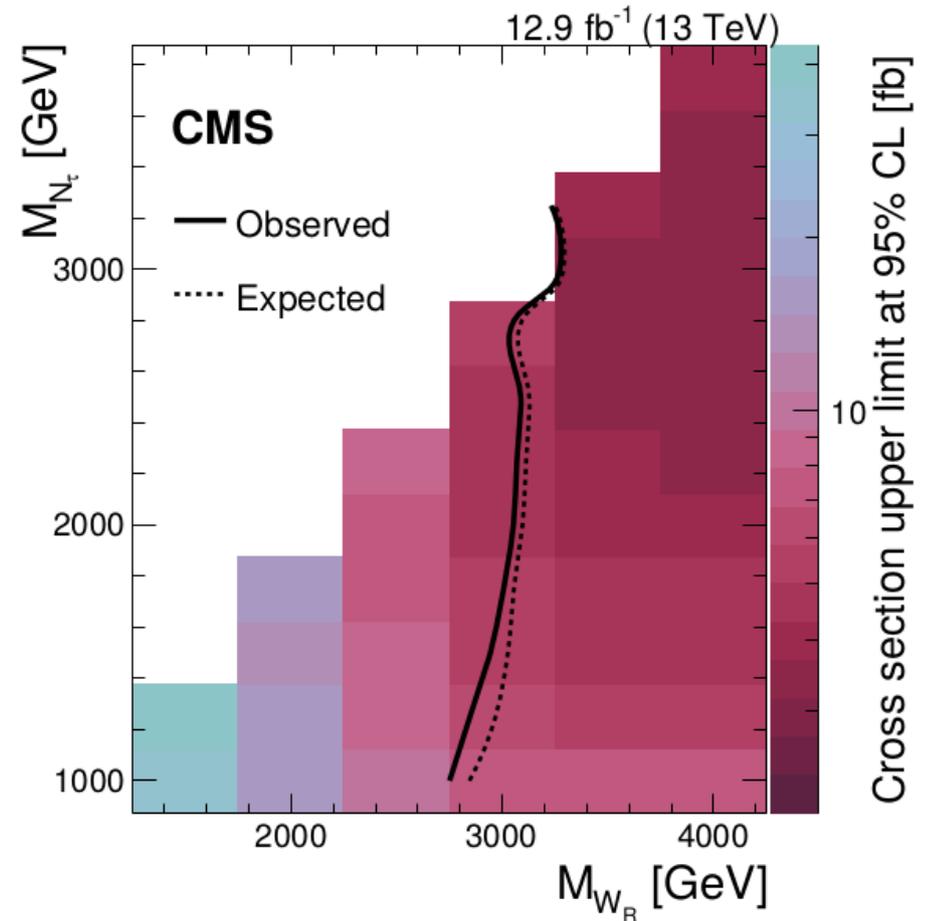
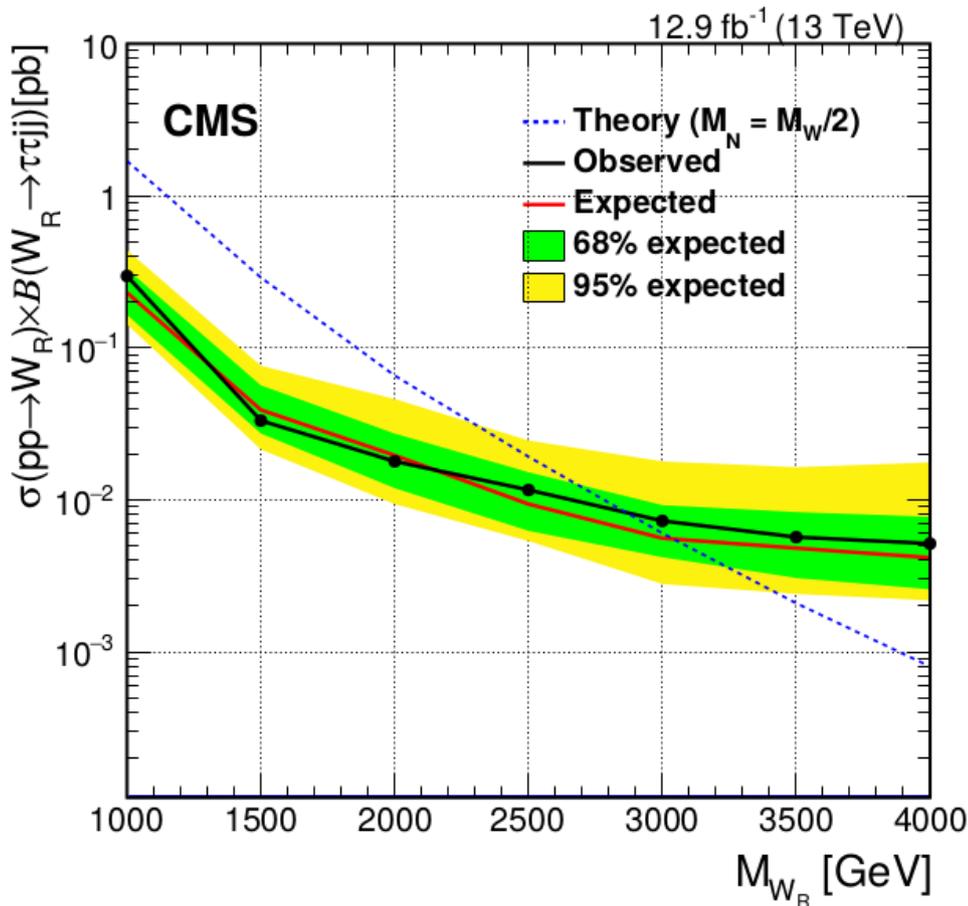
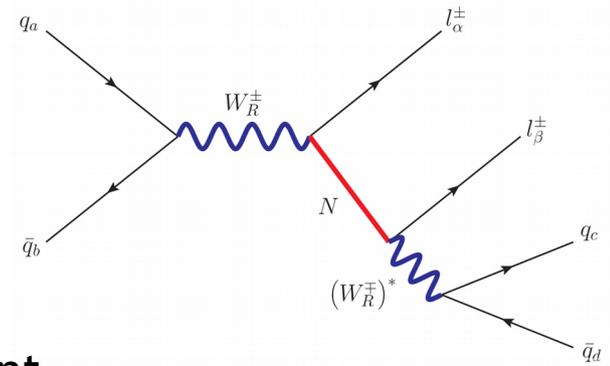
Observable plotted is the scalar sum of final state particle transverse momenta (p_T) and missing transverse energy



Prompt decays with CMS

Searching for LRSM decay, final states with 2 τ -leptons where one decays hadronically and one leptonically

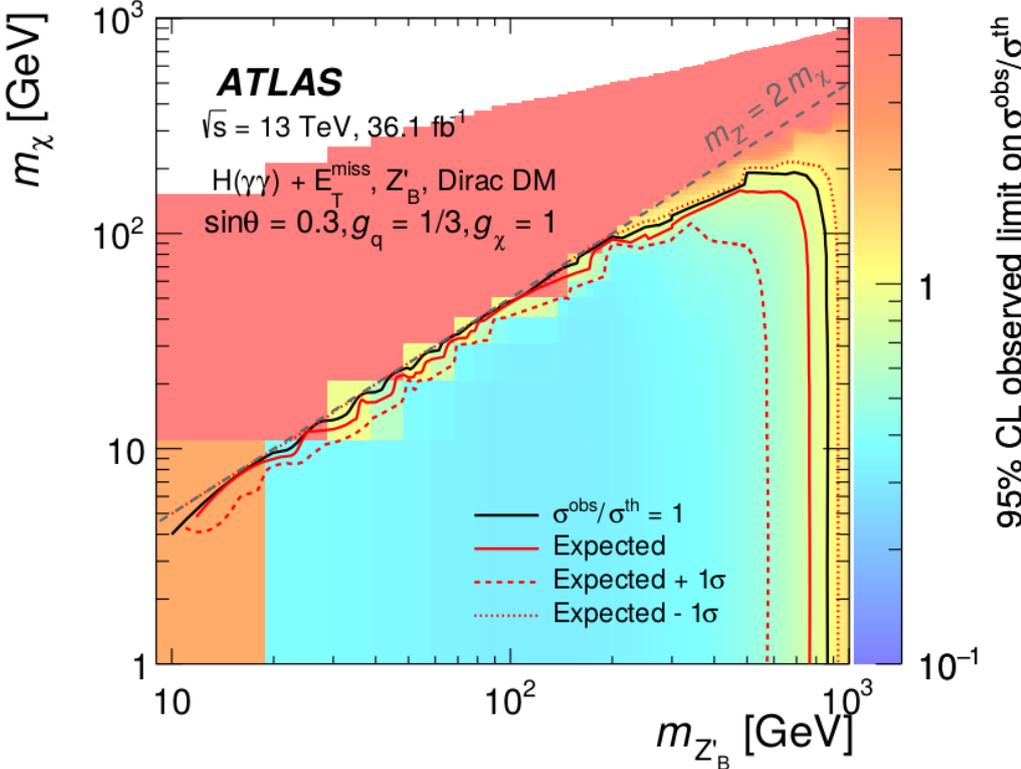
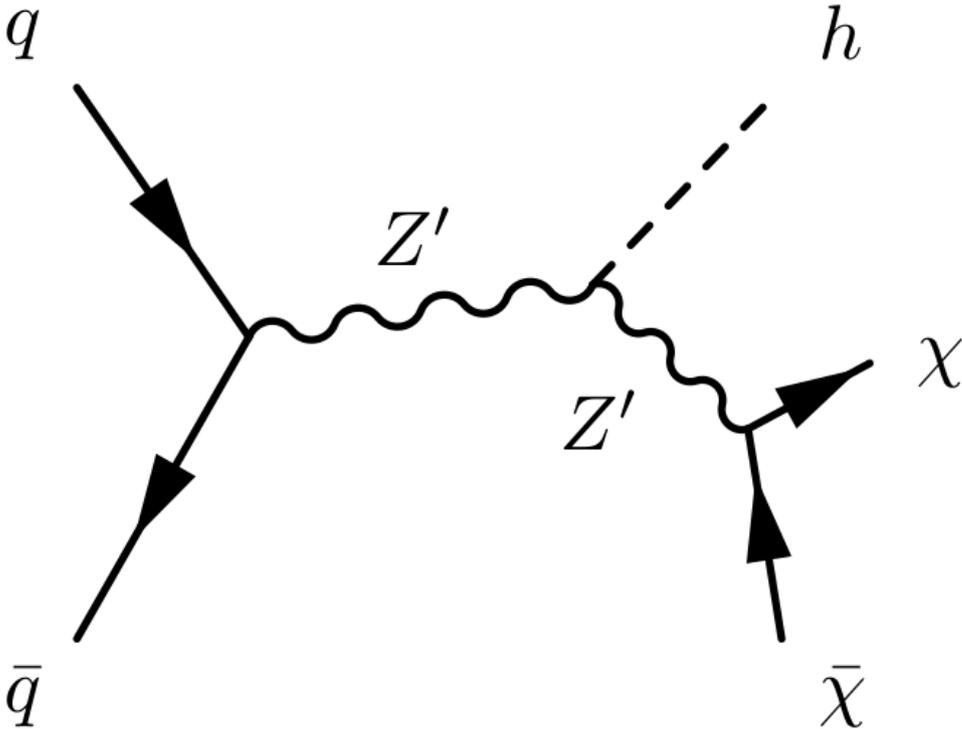
Limit left for fixed $m(N_R):m(W_R)$ ratio, and right for different mass working points



Dark matter searches with ATLAS

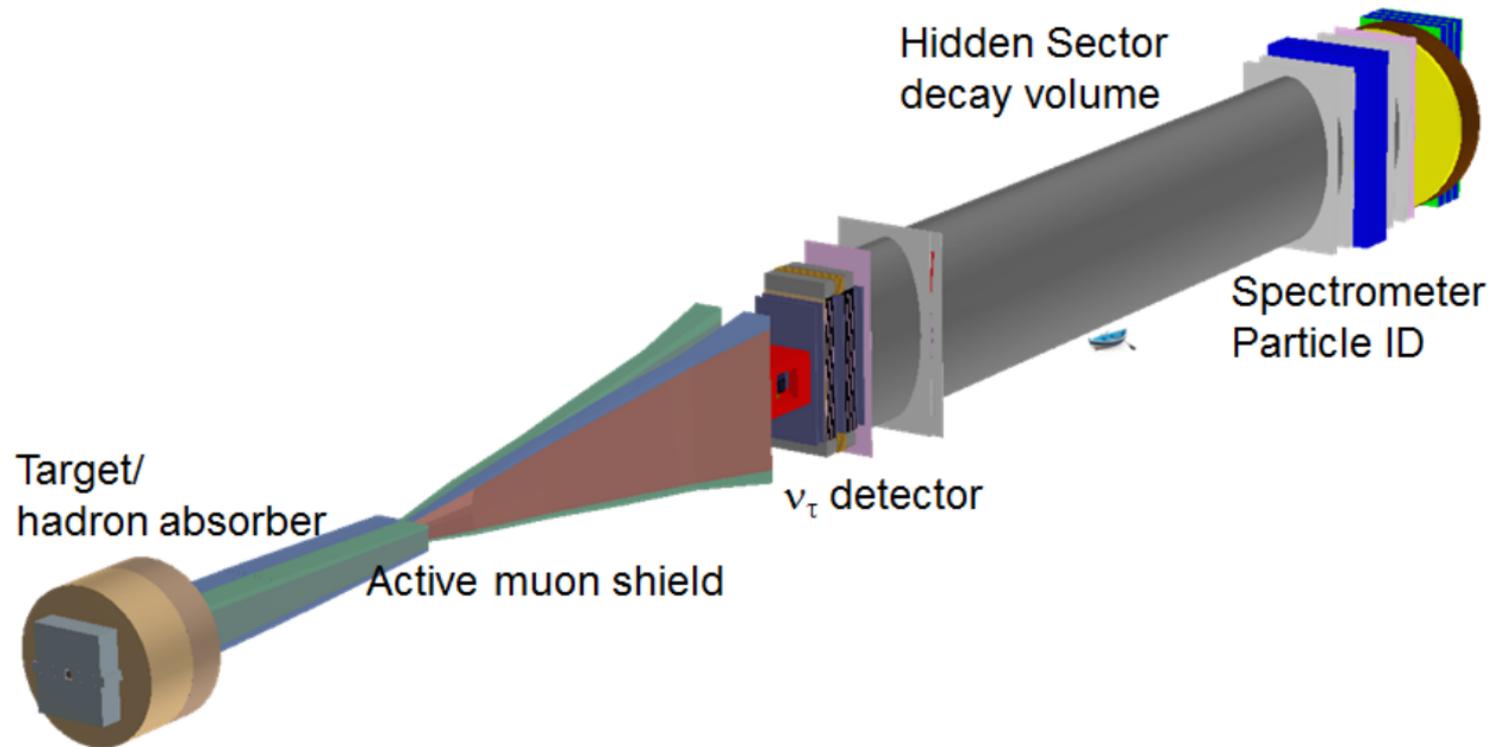
A final state with large MET, and a Higgs boson decaying to a pair of photons

Limit set using a Z' model where Dark Matter only couples to the Z'



Looking into the future...

The proposed **SHiP** experiment has a broadly similar design to NA62, but is specifically intended to search for physics beyond the Standard Model, such as displaced decay vertices from Heavy Neutrinos

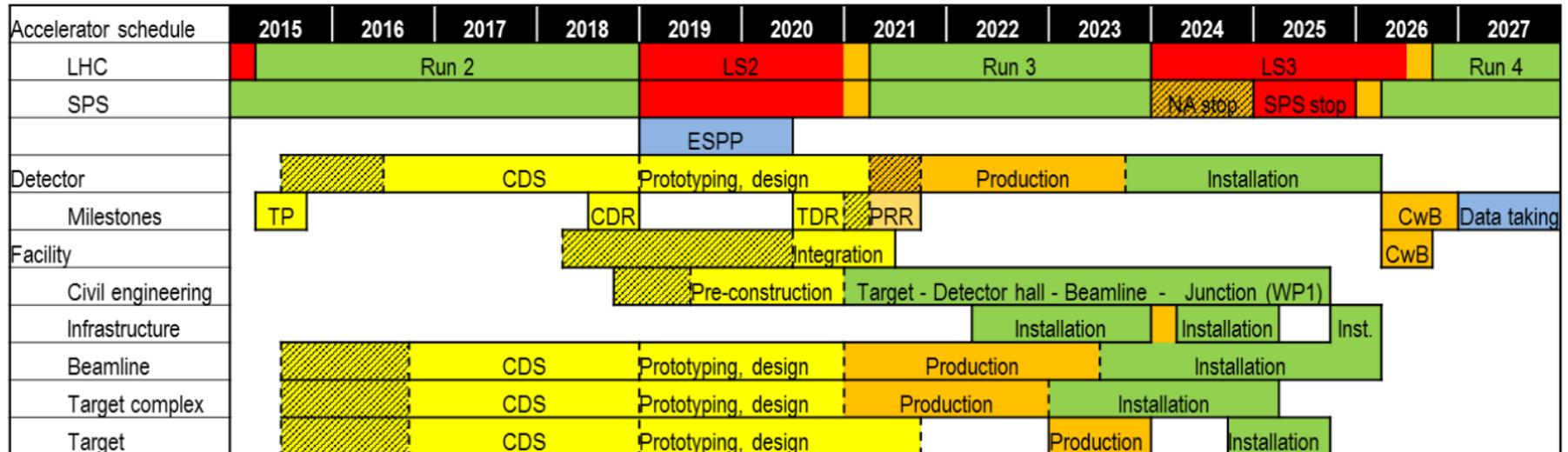


Important features include a magnetic field designed to steer all muons from the target away from the instrumented volume, to reduce background

Looking into the future...

SHiP should eventually be installed using a similar 400 GeV proton beam-line to NA62, and will hopefully be ready for data-taking in 2027

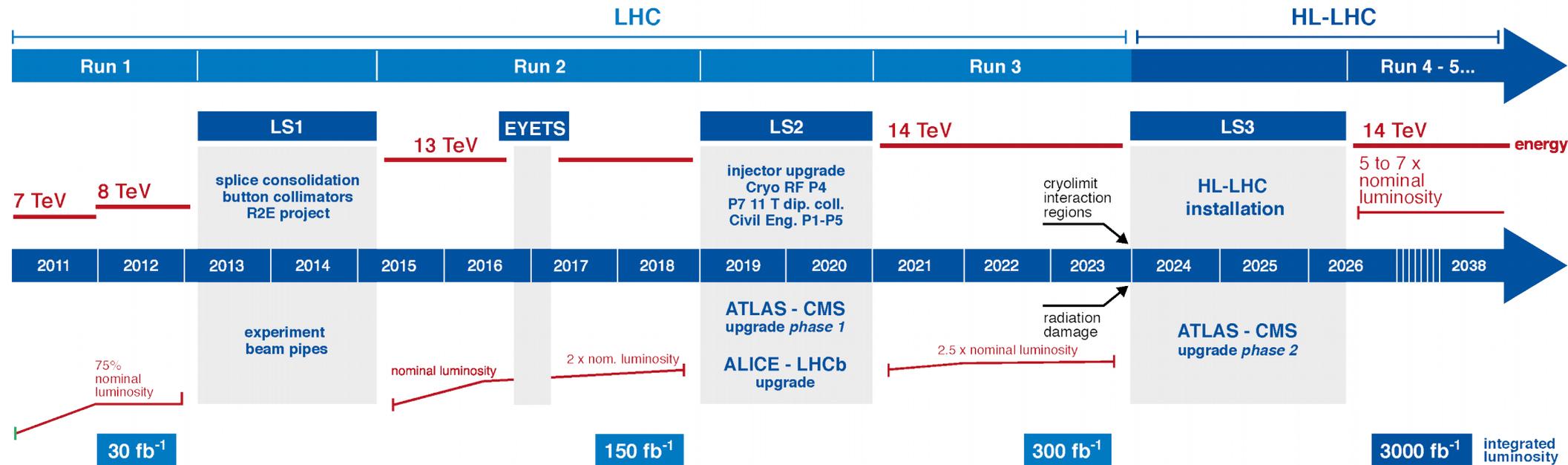
- ▶ Expression of interest 2013
- ▶ Technical proposal (TP) & physics proposal (PP) 2015
- ▶ SPSC and CERN research board recommended we continue to a comprehensive design study (CDS) phase → [Re-optimisation of the entire experiment](#) Now!
- ▶ Part of the CERN Physics beyond colliders (PBC) working group and will be an input to the European strategy meeting (ESPP) in 2019/2020



Looking into the future...

For context, the ShiP timescale is roughly the same as that for the High Luminosity LHC

LHC / HL-LHC Plan



Not that the LHC luminosity isn't already high...

Looking into the future...

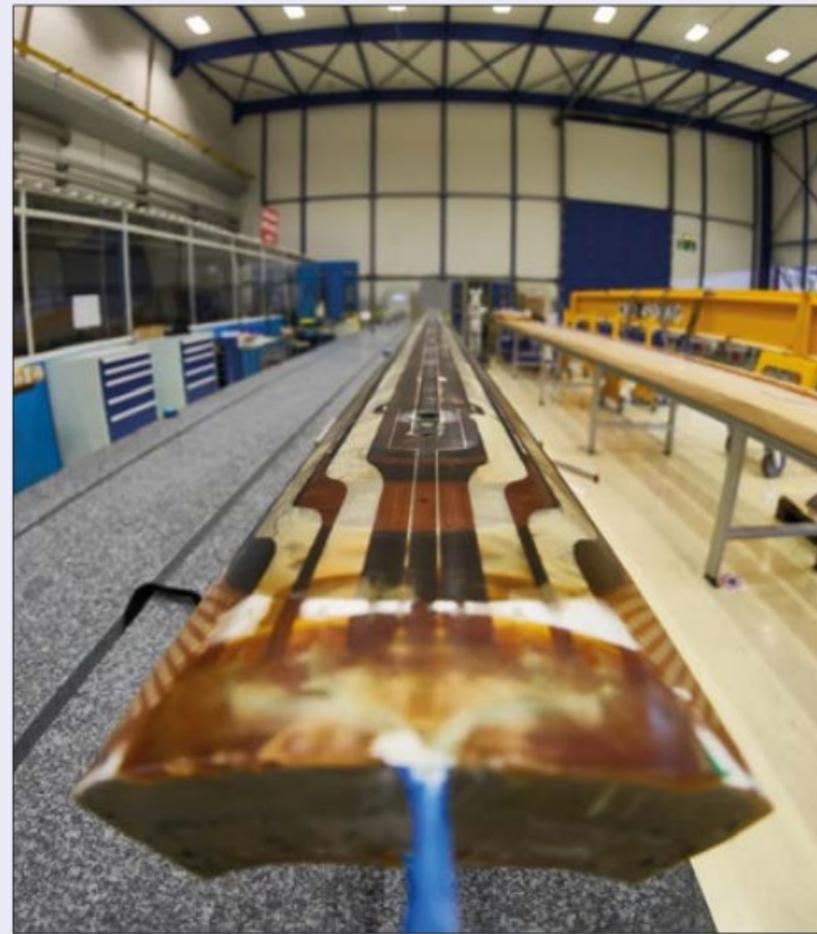
The CERN courier devoted its September issue to superconductivity, and it looks like magnet technology is keeping up with us for now

CERN breaks records with high-field magnets for High-Luminosity LHC

To keep the protons on a circular track at the record-breaking luminosities planned for the LHC upgrade (the HL-LHC) and achieve higher collision energies in future circular colliders, particle physicists need to design and demonstrate the most powerful accelerator magnets ever. The development of the niobium-titanium LHC magnets, currently the highest-field dipole magnets used in a particle accelerator, followed a long road that offered valuable lessons. The HL-LHC is about to change this landscape by relying on niobium tin (Nb_3Sn) to build new high-field magnets for the interaction regions of the ATLAS and CMS experiments. New quadrupoles (called MQFX) and two-in-one dipoles with fields of 11 T will replace the LHC's existing 8 T dipoles in these regions. The main challenge that has prevented the use of Nb_3Sn in accelerator magnets is its brittleness, which can cause permanent degradation under very low intrinsic strain. The tremendous progress of this technology in the past decade led to the successful tests of a full-length 4.5 m-long coil that reached a record nominal field value of 13.4 T at BNL. Meanwhile at CERN, the winding of 7.15 m-long coils has begun.

Several challenges are still to be faced, however, and the next few years will be decisive for declaring production readiness of the MQFX and 11 T magnets. R&D is also ongoing for the development of a Nb_3Sn wire with an improved performance that would allow fields beyond 11 T. It is foreseen that a 14–15 T magnet with real physical aperture will be tested in the US, and this could drive technology for a 16 T magnet for a future circular collider. Based on current experience from the LHC and HL-LHC, we know that the performance requirements for Nb_3Sn for a future circular collider require a large industrial effort to make very large-scale production viable.

● Panagiotis Charitos, CERN.



New long coils for the Nb_3Sn quadrupoles for the HL-LHC.

A little further still...

If you really want to probe the energy and intensity frontiers, you're going to need something big...

