Experience at SPring-8 with Beam Position Monitors for Measuring Second-Order Moments of Charged Particle Beams

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Introduction

- SPring-8 linac BPM was designed so that it could be simply understood with respect to the signal detection. Its transverse cross section keeps cylindrical shape as much as possible.
- This features enables us to consider electrostatic interaction only and to calculate moments easily, i.e.,
 - * singal detection is simplified as the detection of wall current,
 - * a field calculation of the BPM with circular cross-section can be done **analytically**, especially for the higher order moment calculatoin.
- Consider cylindrical coordinate system
- -> Analysis can be separated longitudinally and transversely.



Longitudinal wall current analysis





Multi-particle system (M particles)

Because of the superposition principle, the electric field distribution of multi-paricle system $Er(R, \theta)$ is sum of all single particle distributions $Esr(R, \theta)s$, i.e.;

 $E_{r}(R,\theta) = \sum_{i=1}^{M} E_{sr}(R,\theta)$ $=\sum_{n=1}^{M} \frac{\lambda}{2\pi R\epsilon_{n}} \left(1 + 2\sum_{n=1}^{\infty} \frac{p_{Nn} \cos n\theta + q_{Nn} \sin n\theta}{R^{n}} \right)$ $=\frac{\lambda}{2\pi R\varepsilon_0}\left(M+2\sum_{n=1}^{\infty}\frac{\sum_{N=1}^{M}p_{Nn}\cos n\theta+\sum_{N=1}^{M}q_{Nn}\sin n\theta}{R^n}\right)$ $=\frac{M\lambda}{2\pi R\varepsilon_0}\left(1+2\sum_{n=1}^{\infty}\frac{\frac{1}{M}\sum_{N=1}^{m}p_{Nn}\cos n\theta+\frac{1}{M}\sum_{N=1}^{M}q_{Nn}\sin n\theta}{R^n}\right)$ $=\frac{\Lambda}{2\pi R\varepsilon} \left(1+2\sum_{n=1}^{\infty} \frac{P_n \cos n\theta + Q_n \sin n\theta}{R^n}\right)$

 $\Lambda = M\lambda$ $(P_n) = \frac{1}{M} \sum_{N=1}^{M} p_{Nn} = \frac{1}{M} \sum_{N=1}^{M} b_N^n \cos n\beta_N [m^n]$ $(Q_n) = \frac{1}{M} \sum_{N=1}^{M} q_{Nn} = \frac{1}{M} \sum_{N=1}^{M} b_N^n \sin n\beta_N [m^n]$ nth-order absolute

cosine (sine) moment

NOTE : Absolute moments are measurable physical quantities using BPM.

Separation of centroid and relative moments from absolute moment

We only treat the cosine component here, because sine componet is derived same way.

Multi-particle system (M particles) must have the centroid which is located at (bG, βG).

Position vector of Nth particle (b_N , β_N) can be decomposed into Charged a position vector of the centroid (bg, β g) and a remaining Particle 3 vector (bNg, β Ng) so as to satisfy, $b_3 \beta_3$ $[b_N \cos \beta_N = b_G \cos \beta_G + b_{Na} \cos \beta_{Na}]$ Centroid Charged ³⁹ Particle 2 $b_N \sin \beta_N = b_G \sin \beta_G + b_{Nq} \sin \beta_{Nq}$ Charged Therefore nth-order absolute Nth single Particle $b_2\beta_2$ particle cosine moment p_{Nn} is expressed as, Χ $p_{N1} = b_N \cos \beta_N = b_G \cos \beta_G + b_{Ng} \cos \beta_{Ng} = p_{G1} + p_{Ng1}$ $(p_{N2}) = (b_N^2 \cos 2\beta_N) = (b_G^2 \cos 2\beta_G) + (2b_G b_{Ng}) \cos (\beta_G + \beta_{Ng}) + (b_{Ng}^2 \cos 2\beta_{Ng})$ **D**2g β_{2g} $= (p_{G2}) + 2b_G b_{Nq} \cos(\beta_G + \beta_{Nq}) + (p_{Nq2})$ Typical three particle system (M=3) **p**_{Gn} : nth-order centroid $p_{N3} = b_N^3 \cos 3\beta_N$ cosine moment $= b_G^3 \cos 3\beta_G + 3b_G^2 b_{Nq} \cos (2\beta_G + \beta_{Nq}) + 3b_G b_{Nq}^2 \cos (\beta_G + 2\beta_{Nq}) + b_{Nq}^3 \cos 3\beta_{Nq}$ **p**_{Ngn} : Nth particle nth-order relative $= (p_{G3} + 3b_G^2 b_{Ng} \cos(2\beta_G + \beta_{Ng}) + 3b_G b_{Ng}^2 \cos(\beta_G + 2\beta_{Ng}) + (p_{Ng3})$ cosine moment

Absolute, centroid and relative moments of a multi-particle system

nth-order
absolute
moments
$$P_{1} = \frac{1}{M} \sum_{N=1}^{M} p_{N1} = \frac{1}{M} \sum_{N=1}^{M} p_{G1} + \frac{1}{M} \sum_{N=1}^{M} p_{N01} = p_{G1}$$

$$P_{2} = \frac{1}{M} \sum_{N=1}^{M} p_{N2} = \frac{1}{M} \sum_{N=1}^{M} p_{G2} + \frac{1}{M} \sum_{N=1}^{M} 2b_{G}b_{Ng} \cos(\beta_{G} + \beta_{Ng}) + \frac{1}{M} \sum_{N=1}^{M} p_{Ng2} = p_{G2} + P_{g2}$$

$$P_{3} = \frac{1}{M} \sum_{N=1}^{M} p_{N3} = \frac{1}{M} \sum_{N=1}^{M} p_{G3} + \frac{1}{M} \sum_{N=1}^{M} 3b_{G}^{2}b_{Ng} \cos(2\beta_{G} + \beta_{Ng})$$

$$+ \frac{1}{M} \sum_{N=1}^{M} 3b_{G}b_{Ng}^{2} \cos(\beta_{G} + 2\beta_{Ng}) + \frac{1}{M} \sum_{N=1}^{M} p_{Ng3}$$

$$= p_{G3} + \frac{1}{M} \sum_{N=1}^{M} 3b_{G}b_{Ng}^{2} \cos(\beta_{G} + 2\beta_{Ng}) + \frac{1}{M} \sum_{N=1}^{M} p_{Ng3}$$

$$= p_{G3} + \frac{1}{M} \sum_{N=1}^{M} 3b_{G}b_{Ng}^{2} \cos(\beta_{G} + 2\beta_{Ng}) + P_{g3}$$
Note : $\frac{1}{M} \sum_{N=1}^{M} p_{Gn} = p_{Gn}$, $\frac{1}{M} \sum_{N=1}^{M} q_{Gn} = q_{Gn}$ nth-order centroid moments
 $\frac{1}{M} \sum_{N=1}^{M} p_{Ng1} = \frac{1}{M} \sum_{N=1}^{M} b_{Ng} \cos\beta_{Ng} = 0$, $\frac{1}{M} \sum_{N=1}^{M} q_{Ng1} = \frac{1}{M} \sum_{N=1}^{M} b_{Ng} \sin\beta_{Ng} = 0$
 $P_{gn} = \frac{1}{M} \sum_{N=1}^{M} p_{Ngn}$, $Q_{gn} = \frac{1}{M} \sum_{N=1}^{M} q_{Ngn}$ nth-order relative moments

Definition of size agn and argument αgn of nth-order relative moment

$$\begin{aligned} \text{Definition } a_{gn} &= \sqrt[2n]{P_{gn}^2 + Q_{gn}^2} \quad [m], \quad \alpha_{gn} = \frac{1}{n} \cos^{-1} \frac{P_{gn}}{a_{gn}^n} \quad \&\& \quad \alpha_{gn} = \frac{1}{n} \sin^{-1} \frac{Q_{gn}}{a_{gn}^n} \quad (0 \le \alpha_{gn} < \frac{2\pi}{n}) \\ P_1 &= p_{G1}, \quad Q_1 = q_{G1}, \\ P_2 &= p_{G2} + P_{g2}, \quad Q_2 = q_{G2} + Q_{g2}, \\ P_3 &= p_{G3} + 3b_G a_{g2}^2 \cos(\beta_G + 2\alpha_{g2}) + P_{g3}, \quad Q_3 = q_{G3} + 3b_G a_{g2}^2 \sin(\beta_G + 2\alpha_{g2}) + Q_{g3}, \\ P_4 &= p_{G4} + 6b_G^2 a_{g2}^2 \cos(2\beta_G + 2\alpha_{g2}) + 4b_G a_{g3}^3 \cos(\beta_G + 3\alpha_{g3}) + P_{g4}, \\ Q_4 &= q_{G4} + 6b_G^2 a_{g2}^2 \sin(2\beta_G + 2\alpha_{g2}) + 4b_G a_{g3}^3 \sin(\beta_G + 3\alpha_{g3}) + Q_{g4}, \end{aligned}$$

When a beam is shifted by steering magnet, if we suppose a beam position is changed, but a beam shape (distribution) is not changed, following corollary is obtained.

- * Absolute moments, Pn, Qn, vary.
- * Centroid moments, pGn, qGn, vary.
- * Relative moments, Pgn, Qgn, do not vary.
- * If relative moments vary, something (formulation, calculation, coefficient, measurement, value of moment, correction or etc.) must be invalid.

We carefully take difference over sum of Vds.



Above R....s are named effective aperture radii ($R \leq R$). Smaller apature radius corresponds larger contribution. The smallest order term, \bigcirc and \bigcirc , is dominant. Difference over sum, \bigcirc , is a measurand.

Values of effective aperture radii [mm]

R_{C1P1u}	18.69	R_{S2Q2u}	17.59
R_{C1P5u}	17.50	R_{S2Q4u}	17.39
$R_{_{C1P2d}}$	23.16	$R_{\scriptscriptstyle S2P2d}$	23.16
R_{C1P4d}	19.95	$R_{\scriptscriptstyle S2P4d}$	19.95

Consequently the smallest order moment is expressed as,

$$\begin{aligned} P_{1} &= C_{1} \frac{R_{C1P1u}}{2} + C_{1} \left(\frac{R_{C1P1u}}{R_{C1P2d}}^{2} P_{2} - \frac{R_{C1P1u}}{R_{C1P4d}}^{4} P_{4} + \cdots \right) + \frac{R_{C1P1u}}{R_{C1P5u}}^{5} P_{5} + \cdots \quad [m], \\ C_{1} &= \frac{V_{1} - V_{3} - V_{4} + V_{6}}{V_{1} + V_{3} + V_{4} + V_{6}} \\ Q_{1} &= S_{1} \frac{R_{S1O1u}}{2} + S_{1} \left(\frac{R_{S101u}}{R_{S1P2d}}^{2} P_{2} - \frac{R_{S101u}}{R_{S1P4d}}^{4} P_{4} + \cdots \right) - \frac{R_{S101u}}{R_{S103u}}^{3} Q_{3} - \frac{R_{S101u}}{R_{S103u}}^{5} Q_{5} + \cdots \quad [m], \\ P_{2} &= C_{2} \frac{R_{C2P2u}^{2}}{2} + C_{2} \left(-\frac{R_{C2P2u}^{2}}{R_{C2P2d}}^{2} P_{2} + \frac{R_{C2P2u}^{2}}{R_{C2P4d}}^{4} P_{4} + \cdots \right) + \frac{R_{C2P2u}^{2}}{R_{C2P4u}^{4}} P_{4} + \cdots \quad [m^{2}], \\ Q_{2} &= S_{2} \frac{R_{S202u}^{2}}{2} + S_{2} \left(\frac{R_{S202u}^{2}}{R_{S2P2d}^{2}} P_{2} - \frac{R_{S202u}^{2}}{R_{S2P4d}^{4}} P_{4} + \cdots \right) - \frac{R_{S202u}^{2}}{R_{S204u}^{4}} Q_{4} + \cdots \quad [m^{2}], \\ Q_{3} &= S_{3} \frac{R_{3303u}^{3}}{2} + \cdots \quad [m^{3}], \\ \vdots \text{ Dominant term} \\ &: \text{ Correction terms} \\ \vdots \text{ Correction terms} \\ \end{aligned}$$

: Measurand

Q₃ (absolute moments) are simple. However we actually need relative moments, these formuale are modified so as to be expressed using a combination of P₁, Q₁, P_{g2}, Q_{g2}, Q_{g3} (relative moments). From now we confine calculations up to 5th-order.

Expressions of absolute moments with correction terms which are described by relative moments

$$P_{1} = C_{1} \frac{R_{C1P1u}}{2} + C_{1} \frac{R_{C1P1u}}{R_{C1P2d}^{2}} \left\{ \left(P_{1}^{2} - Q_{1}^{2}\right) + P_{g2} \right\} + C_{1} \frac{R_{C1P1u}}{R_{C1P4d}^{4}} \left\{ \left(P_{1}^{4} - 6P_{1}^{2}Q_{1}^{2} + Q_{1}^{4}\right) + 6\left(P_{1}^{2} - Q_{1}^{2}\right)P_{g2} - 12P_{1}Q_{1}Q_{g2} + \frac{4P_{1}P_{g3}}{4P_{g3}} - 4Q_{1}Q_{g3} + \frac{P_{g4}}{P_{g4}} \right\} + \frac{R_{C1P1u}}{R_{C1P5u}^{5}} \left\{ \left(P_{1}^{5} - 10P_{1}^{3}Q_{1}^{2} + 5P_{1}Q_{1}^{4}\right) + 10\left(P_{1}^{3} - 3P_{1}Q_{1}^{2}\right)P_{g2} - 10\left(3P_{1}^{2}Q_{1} - Q_{1}^{3}\right)Q_{g2} + \frac{10\left(P_{1}^{2} - Q_{1}^{2}\right)P_{g3}}{4P_{g3}} - 20P_{1}Q_{1}Q_{g3} + 5P_{1}P_{g4} - 5Q_{1}Q_{g4} + P_{g5} \right\},$$

$$Q_{1} = S_{1} \frac{R_{S1Q1u}}{2} + S_{1} \frac{R_{S1Q1u}}{R_{S1P2d}^{2}} \left\{ \left(P_{1}^{2} - Q_{1}^{2}\right) + P_{g2} \right\} - S_{1} \frac{R_{S1Q1u}}{R_{S1P4d}^{4}} \left\{ \left(P_{1}^{4} - 6P_{1}^{2}Q_{1}^{2} + Q_{1}^{4}\right) + 6\left(P_{1}^{2} - Q_{1}^{2}\right)P_{g2} - 12P_{1}Q_{1}Q_{g2} + \frac{4P_{1}P_{g3}}{4P_{g3}} - 4Q_{1}Q_{g3} + \frac{P_{g4}}{P_{g4}} \right\} - \frac{R_{S1Q1u}}{R_{S1P4d}^{3}} \left\{ \left(3P_{1}^{2} - Q_{1}^{3}\right) + S_{1} \frac{R_{S1Q1u}}{R_{S1Q2u}^{2}} + S_{1} \frac{R_{S1Q1u}}{R_{S1Q2u}^{3}} \right\} - \frac{R_{S1Q1u}}{R_{S1Q3u}^{3}} \left\{ \left(3P_{1}^{2}Q_{1} - Q_{1}^{3}\right) + 3Q_{1}P_{g2} + 3P_{1}Q_{g2} + Q_{g3} \right\} - \frac{R_{S1Q1u}}{R_{S1Q3u}^{3}} \left\{ \left(5P_{1}^{4}Q_{1} - 10P_{1}^{2}Q_{1}^{3} + Q_{1}^{5}\right) + 10\left(3P_{1}^{2}Q_{1} - Q_{1}^{3}\right)P_{g2} + 10\left(P_{1}^{3} - 3P_{1}Q_{1}^{2}\right)Q_{g2} + \frac{20P_{1}Q_{1}P_{g3}}{4P_{g3}} + 10\left(P_{1}^{2} - Q_{1}^{2}\right)Q_{g3} + \frac{5Q_{1}P_{g4}}{4P_{g4}} + \frac{5P_{1}Q_{g4}}{4P_{g4}} + \frac{6P_{1}P_{g3}}{4P_{g3}} + \frac{6P_{1}P_{g3}}{4P_{g3}} + 10\left(P_{1}^{2} - Q_{1}^{2}\right)Q_{g3} + \frac{5Q_{1}P_{g4}}{4P_{g3}} + \frac{6P_{1}P_{g4}}{4P_{g3}} +$$

The terms colored with are regarded as ZERO, becouse of unmeasurable relateve moments.

$$\begin{split} P_{2} &= C_{2} \frac{R_{C2P2u}^{2}}{2} - C_{2} \frac{R_{C2P2u}^{2}}{R_{C2P2d}^{2}} \Big\{ \Big(P_{1}^{2} - Q_{1}^{2}\Big) + P_{g2} \Big\} \\ &+ C_{2} \frac{R_{C2P2u}^{2}}{R_{C2P4d}^{4}} \Big\{ \Big(P_{1}^{4} - 6P_{1}^{2}Q_{1}^{2} + Q_{1}^{4}\Big) + 6\Big(P_{1}^{2} - Q_{1}^{2}\Big)P_{g2} - 12P_{1}Q_{1}Q_{g2} + 4P_{1}P_{g3} - 4Q_{1}Q_{g3} + P_{g4} \Big\} \\ &+ \frac{R_{C2P2u}^{2}}{R_{C2P4u}^{4}} \Big\{ \Big(P_{1}^{4} - 6P_{1}^{2}Q_{1}^{2} + Q_{1}^{4}\Big) + 6\Big(P_{1}^{2} - Q_{1}^{2}\Big)P_{g2} - 12P_{1}Q_{1}Q_{g2} + 4P_{1}P_{g3} - 4Q_{1}Q_{g3} + P_{g4} \Big\} \\ Q_{2} &= S_{2} \frac{R_{S2Q2u}^{2}}{2} + S_{2} \frac{R_{S2Q2u}^{2}}{R_{S2P2d}^{2}} \Big\{ \Big(P_{1}^{2} - Q_{1}^{2}\Big) + P_{g2} \Big\} \\ &- S_{2} \frac{R_{S2Q2u}^{2}}{R_{S2P4d}^{4}} \Big\{ \Big(P_{1}^{4} - 6P_{1}^{2}Q_{1}^{2} + Q_{1}^{4}\Big) + 6\Big(P_{1}^{2} - Q_{1}^{2}\Big)P_{g2} - 12P_{1}Q_{1}Q_{g2} + 4P_{1}P_{g3} - 4Q_{1}Q_{g3} + P_{g4} \Big\} \\ &- \frac{R_{S2Q2u}^{2}}{R_{S2P4d}^{4}} \Big\{ \Big(P_{1}^{4} - 6P_{1}^{2}Q_{1}^{2} + Q_{1}^{4}\Big) + 6\Big(P_{1}^{2} - Q_{1}^{2}\Big)P_{g2} - 12P_{1}Q_{1}Q_{g2} + 4P_{1}P_{g3} - 4Q_{1}Q_{g3} + P_{g4} \Big\} \\ &- \frac{R_{S2Q2u}^{2}}{R_{S2Q4u}^{4}} \Big\{ \Big(P_{1}^{4} - 6P_{1}^{2}Q_{1}^{2} + Q_{1}^{4}\Big) + 6\Big(P_{1}^{2} - Q_{1}^{2}\Big)P_{g2} - 12P_{1}Q_{1}Q_{g2} + 4P_{1}P_{g3} - 4Q_{1}Q_{g3} + P_{g4} \Big\} \\ &- \frac{R_{S2Q2u}^{2}}{R_{S2Q4u}^{4}} \Big\{ \Big(P_{1}^{4} - 6P_{1}^{2}Q_{1}^{2} + Q_{1}^{4}\Big) + 6\Big(P_{1}^{2} - Q_{1}^{2}\Big)P_{g2} - 12P_{1}Q_{1}Q_{g2} + 4P_{1}P_{g3} - 4Q_{1}Q_{g3} + P_{g4} \Big\} \\ &- \frac{R_{S2Q2u}^{2}}{R_{S2Q4u}^{4}} \Big\{ \Big(P_{1}^{4} - 6P_{1}^{2}Q_{1}^{2} + Q_{1}^{4}\Big) + 12P_{1}Q_{1}P_{g2} + 6\Big(P_{1}^{2} - Q_{1}^{2}\Big)Q_{g2} + 4Q_{1}P_{g3} + 4P_{1}Q_{g3} + Q_{g4} \Big\} \\ &Q_{3} = S_{3} \frac{R_{3}^{3}}{2} \Big\} \end{split}$$

The terms colored with are regarded as ZERO.

How do we calculate the absolute (or relative) moments?

For single variable, e.g. x, suppose an equation can be written as x = F(x). If F(x) is continuous, smooth, monotone and IdF(x)/dxI < 1 in the typical interest region, a successive iteration method can be employed. That is for proper initial x_0 and iteration number i, solution of x can be obtained as the limit at infinity $x = \lim_{i \to \infty} x_i$.

$$x_{1} = F(x_{0}), \quad x_{2} = F(x_{1}), \quad x_{3} = F(x_{2}), \quad x_{4} = F(x_{3}), \quad x_{5} = F(x_{4}), \quad x_{6} = F(x_{5}), \dots$$

$$y = x_{0}$$

$$y = f(x)$$

$$y = f(x)$$

$$x_{0} = x_{2}x_{4}x_{3}x_{1}$$

$$x_{1} = x_{2}$$

$$x_{1} = x_{2}$$

$$x_{2} = x_{1} + x_{2}$$

$$x_{2} = x_{2} + x_{3} + x_{3}$$

$$x_{1} = x_{2}$$

In the case of the multiple variable we solve following equations.

 $\begin{cases} P_{1} = F_{P1}(P_{1}, Q_{1}, P_{g2}, Q_{g2}, Q_{g3}) \\ Q_{1} = F_{Q1}(P_{1}, Q_{1}, P_{g2}, Q_{g2}, Q_{g3}) \\ P_{2} = F_{P2}(P_{1}, Q_{1}, P_{g2}, Q_{g2}, Q_{g3}) \\ Q_{2} = F_{Q2}(P_{1}, Q_{1}, P_{g2}, Q_{g2}, Q_{g3}) \\ Q_{3} = F_{Q3}(P_{1}, Q_{1}, P_{g2}, Q_{g2}, Q_{g3}) \end{cases}$

If all function $F_{P1}(P_1, \dots)$, $F_{Q1}(P_1, \dots)$, \dots are continuous, smooth, monotone and, all absolute values of partial differential $I\partial F_{P1}/\partial P_1 I$, $I\partial F_{P1}/\partial Q_1 I$, \dots are smaller than 1, the successive iteration method may be applicable.



Iteration scheme

• First we calculate P₁, Q₁, P₂, Q₂, Q₃ using no correction term as the initial values of P₁₀, Q₁₀, P₂₀ (P_{g20}), Q₂₀ (Q_{g20}), Q₃₀ (Q_{g30}).

$$P_{10} = C_{1} \frac{P_{C1}P_{10}}{2}$$

$$Q_{10} = S_{1} \frac{R_{S101u}}{2} \qquad P_{10}$$

$$Q_{10} = S_{1} \frac{R_{S101u}}{2} \qquad P_{10}$$

$$Q_{11} = F_{P1}(P_{10},Q_{10},P_{g20},Q_{g20},Q_{g30}) \qquad P_{11}$$

$$Q_{11} = F_{P1}(P_{10},Q_{10},P_{g20},Q_{g20},Q_{g30}) \qquad Q_{11}$$

$$P_{12} = F_{P1}(P_{11},Q_{11},P_{g21},Q_{g21},Q_{g31}) \qquad P_{12}$$

$$Q_{12} = F_{P1}(P_{12},Q_{12},P_{g22},Q_{g31}) \qquad Q_{13}$$

$$P_{20} = C_{2} \frac{R_{22P2u}^{2}}{2} \qquad P_{g20}$$

$$Q_{20} = S_{2} \frac{R_{22O2u}^{2}}{2} \qquad Q_{g30}$$

$$Q_{21} = F_{P2}(P_{10},Q_{10},P_{g20},Q_{g20},Q_{g30}) \qquad Q_{g21}$$

$$Q_{21} = F_{O2}(P_{10},Q_{10},P_{g20},Q_{g20},Q_{g30}) \qquad Q_{g31}$$

$$Q_{31} = F_{O3}(P_{10},Q_{10},P_{g20},Q_{g20},Q_{g30}) \qquad Q_{g31}$$

$$I = 1$$

$$I = 1$$

$$I = 2$$

$$I = 2$$

$$I = 2$$

$$I = 3$$

Simulation of successive iteration (convergence)



Comparison of calculation error between the cases of no correction, up to 3rd-order and up to 5th-order



We supposed zero beam size, i.e.,

all relative momennts are zero,

$$\mathsf{P_{g2}} = \mathsf{Q}_{g2} = \mathsf{Q}_{g3} = \cdots = \mathsf{0}.$$

If attenuation values of electrode (or cable, •••) are not corret, the error distribution is

asymmetrical.

-> Entire calibration

Electron beam is swept horizontally and vertically, then we determine relative attenuation factors of Electrode 2 electrode channels so that relative moments are constant. P_{a2} = constant, Q_{a2} = constant, Q_{a3} = 0 Electrode 1 10mm-Electrode/3 ST LS 2 V P \leq 10mm -10^{mm} Steering Magnet Electrode 4 -10mm+ Electrode/6 Electrode 5 -2 Current of Measured -3 Beam Position at BPM QM LSBT 10 BPM LSBT 9 2014/5/25 15:28:00 2014/5/25 15:32:30 2014/5/25 15:37:00 QM LSBT 1 BPM LSBT_2 Date · Time QM LSBT 11 **BPM_LSBT_8** SB PM LSBT 1 ST LSBT 1 BM LSBT QM LSBT 2 ST_LSBT_2 LSB1 Matching section (end of SPring-8 1GeV linac)

Estimated

noise level : $150 \times 10^{-9} [V / \sqrt{Hz}] \cdot 1 \times 10^{3} [\sqrt{Hz}] / 0.02 [V / dB] = 0.008 dB -> 0.1\%$ Error (1 Electrode)

Suppose $V \approx V_1 \approx V_3 \approx V_4 \approx V_6$, $\Delta V \approx \Delta V_1 \approx \Delta V_3 \approx \Delta V_4 \approx \Delta V_6$. Take $(\Delta P_1)^2$ or $(\Delta Q_2)^2$, then calculate deviation $\sigma_{P_1}^2$ or $\sigma_{Q_2}^2$. Because $\frac{\partial C_1}{\partial V_i} \left(\text{or } \frac{\partial S_2}{\partial V_i} \right)$ and ΔV_j $(i \neq j)$ are uncorrelated statistically,

the cross term $\frac{\partial C_1}{\partial V_i} \Delta V_j$ or $\frac{\partial S_2}{\partial V_i} \Delta V_j$ ($i \neq j$) is vanish.

Finally we obtain following relations;

$$\sigma_{P1}^{2} \approx \left(\frac{R_{C1P1u}}{4}\right)^{2} \frac{\sigma_{V}^{2}}{V^{2}} \Rightarrow \sigma_{P1} \approx \frac{R_{C1P1u}}{4} \frac{\sigma_{V}}{V}, \ \sigma_{Q2}^{2} \approx \left(\frac{R_{S2Q2u}^{2}}{4}\right)^{2} \frac{\sigma_{V}^{2}}{V^{2}} \Rightarrow \sigma_{Q2} \approx \frac{R_{S2Q2u}^{2}}{4} \frac{\sigma_{V}}{V},$$

Because of $\frac{\sigma_V}{V} \approx 0.001$, $R_{C1P1u} = 18.69$ [mm], $R_{S2Q2u} = 17.59$ [mm],

estimated resolutions $\sigma_{P1} \approx 0.005 \text{ [mm]}$ and $\sigma_{Q2} \approx 0.077 \text{ [mm^2]}$ are obtained.

Resolution measurement by means of three BPM method in the drift space (on beam)

Resolution measurement by means of signal generator N5181A (off beam)

Summary

- SPring-8 BPM detects wall current longitudinally and transversely
- Transverse wall current distribution is determined by beam moments and geometrical factor.
- Effective aperture radius is calculated from geometrical factor and relates to an amplitude of BPM output.
- Absolute moment is <u>measurable physical quantity</u>. Relative moment which depends on the beam shape is <u>exracted from the absolute moment</u>.
- As the moment-correction scheme a recursive method was employed.
- To calibrate the attenuation factors of channels the entire calibration method was developed.
- Measurement resolutions of horizontal beam positions are <u>5µm for CW RF</u> measurement, but <u>15µm for pulse RF</u> measurement.