

# See-saw mediators in t-channel

Margherita Ghezzi

PAUL SCHERRER INSTITUT



BSM direct searches at CLIC, 21 February 2018

# Outline

- 1 Introduction: the doubly charged scalar
- 2 Low energy: EFT and current limits
- 3 High energy: LHC searches
- 4 High energy: future colliders
- 5 Summary

# Outline

- 1 Introduction: the doubly charged scalar
- 2 Low energy: EFT and current limits
- 3 High energy: LHC searches
- 4 High energy: future colliders
- 5 Summary

# The doubly charged scalar from the $SU(2)_L$ -triplet scalar

## Type-II see-saw model

$$S = \begin{pmatrix} S^+ & \sqrt{2}S^{++} \\ \sqrt{2}S^0 & -S^+ \end{pmatrix}$$

$$\langle S \rangle_0 = \begin{pmatrix} 0 & 0 \\ w & 0 \end{pmatrix}$$

Yukawa term with the triplet:

$$\Delta\mathcal{L}_Y = f_{ij} L_i^T C^{-1} i\tau_2 S L_j + \text{h.c.}$$

Majorana mass term for neutrinos:

$$m_{ij} \bar{\nu}_{iL}^c \nu_{jL} \quad m_{ij} = w f_{ij} = m_{ji}$$

T. P. Cheng and L. F. Li, Phys. Rev. D 22 (1980) 2860

W. Grimus, R. Pfeiffer and T. Schwetz, Eur. Phys. J. C 13 (2000) 125

E. Ma, M. Raidal and U. Sarkar, Nucl. Phys. B 615 (2001) 313

A. G. Akeroyd and M. Aoki, Phys. Rev. D 72 (2005) 035011

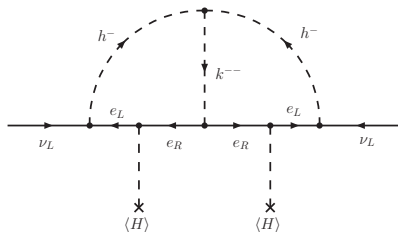
# The doubly Charged $SU(2)_L$ -singlet scalar

## Zee-Babu model

SM + 2  $SU(2)_L$ -singlet scalars:

- a singly charged scalar which couples to left-handed leptons:  $h^\pm$
- a doubly charged scalar which couples to right-handed leptons:  $k^{\pm\pm}$

It generates mass terms for the neutrinos at two loops:



A. Zee, Nucl. Phys. B **264** (1986) 99

K. S. Babu, Phys. Lett. B **203**, 132 (1988)

M. Nebot, J. F. Oliver, D. Palao and A. Santamaria, Phys. Rev. D **77** (2008) 093013

# The doubly Charged $SU(2)_L$ -singlet scalar

## Minimal model for neutrino masses

SM + 1  $SU(2)_L$ -singlet doubly charged scalar:  $S_R^{\pm\pm}$

It couples only with right-handed charged leptons:

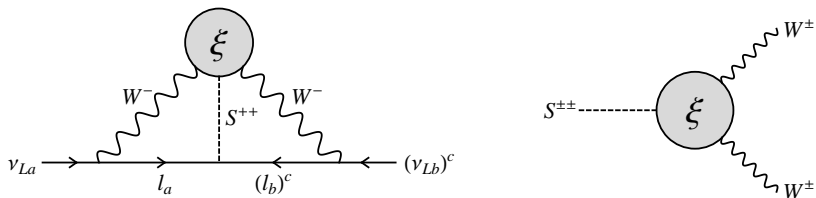
$$\begin{aligned}\mathcal{L}_{UV} = & \mathcal{L}_{\text{SM}} + (D_\mu S^{++})^\dagger (D^\mu S^{++}) \\ & + \left( \lambda_{ab} \overline{(\ell_R)_a^c} \ell_{Rb} S^{++} + \text{h.c.} \right)\end{aligned}$$

$\lambda_{ab}$  consist of 6 independent parameters and allow for **LFV processes**

S. F. King, A. Merle and L. Panizzi, JHEP 1411 (2014) 124

# The doubly charged $SU(2)_L$ -singlet scalar

Neutrino mass terms are generated at three loop:



S. F. King, A. Merle and L. Panizzi, JHEP 1411 (2014) 124

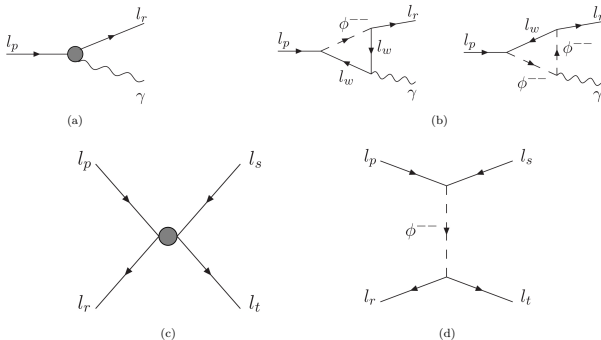
# Outline

- 1 Introduction: the doubly charged scalar
- 2 Low energy: EFT and current limits**
- 3 High energy: LHC searches
- 4 High energy: future colliders
- 5 Summary



# Low-energy effective Lagrangian and the matching

Feynman diagrams representing the UV-complete contributions that match to the dipole and four-fermion operators.



- Diagrams in Fig. (b) match into the diagram in Fig. (a) (dipole interaction)
- Diagram in Fig. (d) matches into the diagram in Fig. (c) (contact interaction)

# Low-energy effective Lagrangian and the matching

Dipole			
$Q_{e\gamma}$	$em_r(\bar{l}_p\sigma^{\mu\nu}P_L l_r)F_{\mu\nu} + \text{H.c.}$		
Scalar/Tensorial		Vectorial	
$Q_S$	$(\bar{l}_p P_L l_r)(\bar{l}_s P_L l_t) + \text{H.c.}$	$Q_{VLL}$	$(\bar{l}_p \gamma^\mu P_L l_r)(\bar{l}_s \gamma_\mu P_L l_t)$
		$Q_{VLR}$	$(\bar{l}_p \gamma^\mu P_L l_r)(\bar{l}_s \gamma_\mu P_R l_t)$
		$Q_{VRR}$	$(\bar{l}_p \gamma^\mu P_R l_r)(\bar{l}_s \gamma_\mu P_R l_t)$
$Q_{Slq(1)}$	$(\bar{l}_p P_L l_r)(\bar{q}_s P_L q_t) + \text{H.c.}$	$Q_{VlqLL}$	$(\bar{l}_p \gamma^\mu P_L l_r)(\bar{q}_s \gamma_\mu P_L q_t)$
$Q_{Slq(2)}$	$(\bar{l}_p P_L l_r)(\bar{q}_s P_R q_t) + \text{H.c.}$	$Q_{VlqLR}$	$(\bar{l}_p \gamma^\mu P_L l_r)(\bar{q}_s \gamma_\mu P_R q_t)$
$Q_{Tlq}$	$(\bar{l}_p \sigma^{\mu\nu} P_L l_r)(\bar{q}_s \sigma_{\mu\nu} P_L q_t) + \text{H.c.}$	$Q_{VlqRL}$	$(\bar{l}_p \gamma^\mu P_R l_r)(\bar{q}_s \gamma_\mu P_L q_t)$
		$Q_{VlqRR}$	$(\bar{l}_p \gamma^\mu P_R l_r)(\bar{q}_s \gamma_\mu P_R q_t)$

Dimension-six operators that allow for effective leptonic transitions below the EW scale

# Current low-energy experimental limits

$$\text{Br} [\tau^\mp \rightarrow e^\mp e^\pm e^\mp] \leq 1.4 \times 10^{-8}$$

$$\text{Br} [\tau^\mp \rightarrow \mu^\mp \mu^\pm \mu^\mp] \leq 1.2 \times 10^{-8}$$

$$\text{Br} [\tau^\mp \rightarrow e^\mp \mu^\pm \mu^\mp] \leq 1.6 \times 10^{-8}$$

$$\text{Br} [\tau^\mp \rightarrow \mu^\mp e^\pm \mu^\mp] \leq 9.8 \times 10^{-9}$$

$$\text{Br} [\tau^\mp \rightarrow \mu^\mp e^\pm e^\mp] \leq 1.1 \times 10^{-8}$$

$$\text{Br} [\tau^\mp \rightarrow e^\mp \mu^\pm e^\mp] \leq 8.4 \times 10^{-8}$$

$$\text{Br} [\mu^\mp \rightarrow e^\mp e^\pm e^\mp] \leq 1.0 \times 10^{-12}$$

$$\text{Br} [\tau \rightarrow e \gamma] \leq 3.3 \times 10^{-8}$$

$$\text{Br} [\tau \rightarrow \mu \gamma] \leq 4.4 \times 10^{-8}$$

$$\text{Br} [\mu \rightarrow e \gamma] \leq 4.2 \times 10^{-13}$$

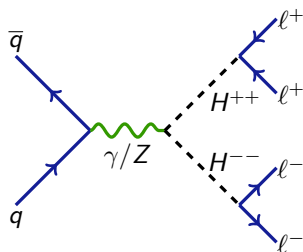
$$\text{BR}(l_p^\pm \rightarrow l_r^\pm \gamma) \simeq \frac{\alpha m_p^5}{(24\pi^2)^2 m_\phi^4 \Gamma_p} \left| \sum_{w=1}^3 \lambda_{pw} \lambda_{rw}^* \right|^2$$

$$\text{BR}(l_p^\pm \rightarrow l_r^\pm l_s^\mp l_t^\pm) \simeq \frac{m_p^5 |\lambda_{ps}|^2 |\lambda_{rt}|^2}{s_{rt} 6(4\pi)^3 m_\phi^4 \Gamma_p}$$

# Outline

- 1 Introduction: the doubly charged scalar
- 2 Low energy: EFT and current limits
- 3 High energy: LHC searches**
- 4 High energy: future colliders
- 5 Summary

# Direct searches at LHC



- Signature: same-sign lepton pairs
- Assumptions on the branching ratios
- Narrow width approximation

## ATLAS 7 TeV:

- Eur.Phys.J. C72 (2012) 2244

## CMS 7 TeV:

- Eur.Phys.J. C72 (2012) 2189

## ATLAS 13 TeV:

- CERN-EP-2017-198

## CMS 13 TeV:

- CMS-PAS-HIG-16-036

# Current limits from LHC

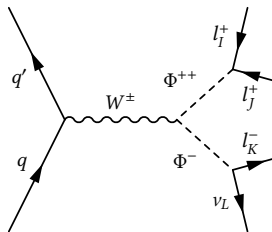
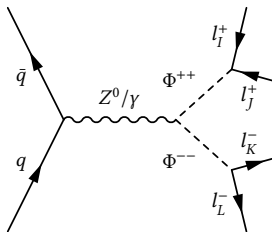
## CMS searches

Search for a scalar triplet  $S = \begin{pmatrix} S^+ & \sqrt{2}S^{++} \\ \sqrt{2}S^0 & -S^+ \end{pmatrix}$  with degenerate masses.

12.9 fb<sup>-1</sup> of integrated luminosity at 13 TeV

Channels:

- Pair production with decays  $S^{++}S^{--} \rightarrow \ell^+\ell^+\ell^-\ell^-$
- Associated production with decays  $S^{\pm\pm}S^\mp \rightarrow \ell^\pm\ell^\pm\ell^\mp\nu$



# Current limits from LHC

## CMS searches

- $S_L^{\pm\pm}$  decaying at 100% to  $ee, \mu\mu, \tau\tau, e\mu, e\tau, \mu\tau$ ;
- Benchmark points:

Benchmark Point	$ee$	$e\mu$	$e\tau$	$\mu\mu$	$\mu\tau$	$\tau\tau$
BP1	0	0.01	0.01	0.30	0.38	0.30
BP2	1/2	0	0	1/8	1/4	1/8
BP3	1/3	0	0	1/3	0	1/3
BP4	1/6	1/6	1/6	1/6	1/6	1/6

Lower bounds on the mass of the  $S_L^{\pm\pm}$  - observed (expected) 95% CL:

Benchmark	AP [GeV]	PP [GeV]	Combined [GeV]
100% $\Phi^{\pm\pm} \rightarrow ee$	734 (720)	652 (639)	800 (785)
100% $\Phi^{\pm\pm} \rightarrow e\mu$	750 (729)	665 (660)	820 (810)
100% $\Phi^{\pm\pm} \rightarrow \mu\mu$	746 (774)	712 (712)	816 (843)
100% $\Phi^{\pm\pm} \rightarrow e\tau$	568 (582)	481 (543)	714 (658)
100% $\Phi^{\pm\pm} \rightarrow \mu\tau$	518 (613)	537 (591)	643 (708)
100% $\Phi^{\pm\pm} \rightarrow \tau\tau$	479 (483)	396 (419)	535 (544)
Benchmark 1	613 (649)	519 (548)	723 (715)
Benchmark 2	670 (671)	465 (554)	716 (723)
Benchmark 3	706 (682)	531 (562)	761 (732)
Benchmark 4	639 (639)	496 (539)	722 (704)

$S_R^{\pm\pm}$  may have similar kinematic properties, but potentially very different production cross sections. No associate production.

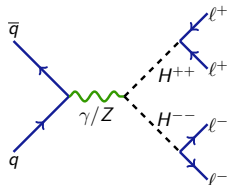
# Current limits from LHC

## ATLAS searches

36.1 fb<sup>-1</sup> of integrated luminosity at 13 TeV.

Scenarios:

- $\sum_{i,j=e,\mu} \mathcal{B}(S^{\pm\pm} \rightarrow \ell_i \ell_j) = 100\%$ 
  - $m(S_L^{\pm\pm})$  between 770 GeV and 870 GeV @ 95% C.L.
  - $m(S_R^{\pm\pm})$  between 660 GeV and 760 GeV @ 95% C.L.
- $\mathcal{B}(S^{\pm\pm} \rightarrow \ell_i \ell_j) > 10\%$  (decays to  $\tau$  and  $W$  are possible)
  - $m(S_L^{\pm\pm})$  larger than 450 GeV @ 95% C.L.
  - $m(S_R^{\pm\pm})$  larger than 320 GeV @ 95% C.L.





# Expected discovery power of HiLumi-LHC



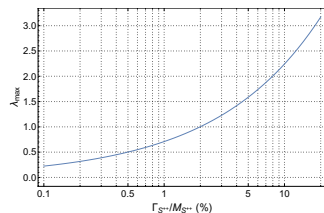
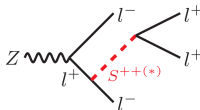
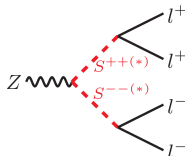
Expected lower limits on the mass - projections at 3000 fb<sup>-1</sup>

CMS 12.9 fb <sup>-1</sup>	$S_L^{\pm\pm}$	[600, 800] GeV	→	[2400, 3200] GeV:
	$S_L^{\pm\pm}, \text{PP only}$	[400, 700] GeV	→	[1600, 2800] GeV
ATLAS 36.1 fb <sup>-1</sup>	$S_L^{\pm\pm}, \mathcal{B}_{\ell\ell} = 100\%$	~ 800 GeV	→	~ 2400 GeV
	$S_L^{\pm\pm}, \mathcal{B}_{\ell\ell} > 10\%$	~ 450 GeV	→	~ 1350 GeV
	$S_R^{\pm\pm}, \mathcal{B}_{\ell\ell} = 100\%$	~ 700 GeV	→	~ 2100 GeV
	$S_R^{\pm\pm}, \mathcal{B}_{\ell\ell} > 10\%$	~ 300 GeV	→	~ 900 GeV

# Width effects

- No *production*  $\times$  *decay* approximation;
- some topologies that are negligible in the NWA can become relevant;
- assumption: gauge sector not modified, i.e.  $S_R^{\pm\pm}$  coupling to  $Z$  is not a free coupling;
- $\Gamma_S$  is considered as a free parameter and  $\sum_{ab,cd} \Gamma_S^{\text{part}} \leq \Gamma_S$

$$\sigma_{PP \rightarrow l_a^+ l_b^+ l_c^- l_d^-}(M_S, \Gamma_S, \lambda_{ab}, \lambda_{cd}) = \lambda_{ab}^2 \lambda_{cd}^2 \hat{\sigma}(M_S, \Gamma_S)$$

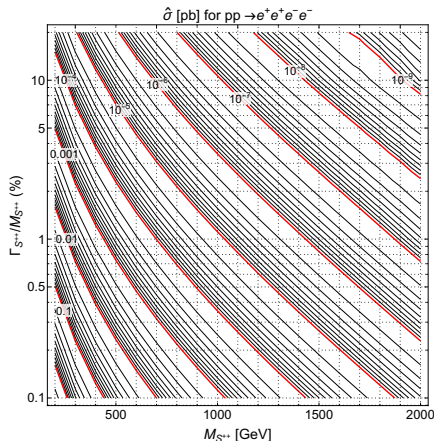


Crivellin, MG, Panizzi, Pruna, Signer, work in progress

# Width effects: results

Very good approximation for light leptons:

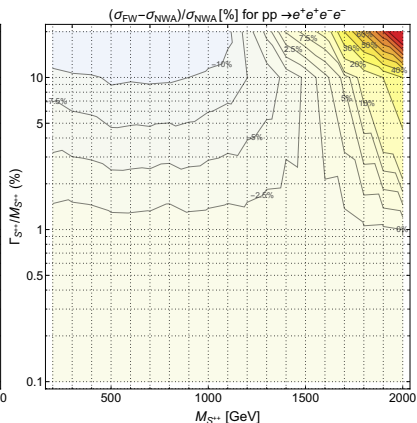
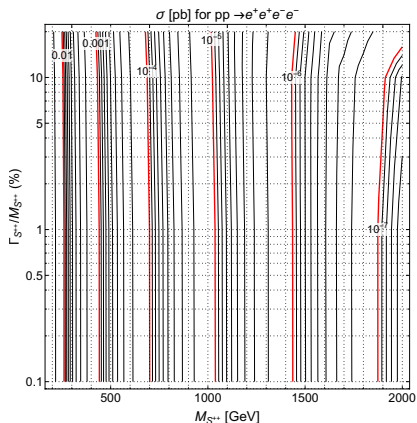
$$\sigma_{PP \rightarrow l_a^+ l_b^+ l_c^- l_d^-}(M_S, \Gamma_S, \lambda_{ij}) = \kappa_{ab,cd} \lambda_{ab}^2 \lambda_{cd}^2 \hat{\sigma}_{PP \rightarrow 2e^+ 2e^-}(M_S, \Gamma_S),$$



Crivellin, MG, Panizzi, Pruna, Signer, work in progress

# Width effects: results

- Cross-section corresponding to the maximum coupling values;
- relative ratio between cross-sections in the FW regime and NWA.



Crivellin, MG, Panizzi, Pruna, Signer, work in progress

# Outline

- 1 Introduction: the doubly charged scalar
- 2 Low energy: EFT and current limits
- 3 High energy: LHC searches
- 4 High energy: future colliders**
- 5 Summary

# Perspective of searches at future colliders

Crivellin, MG, Panizzi, Pruna, Signer, work in progress

(Preliminary plot)

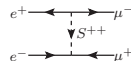
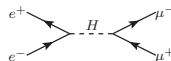
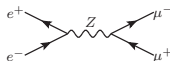
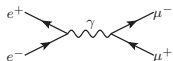
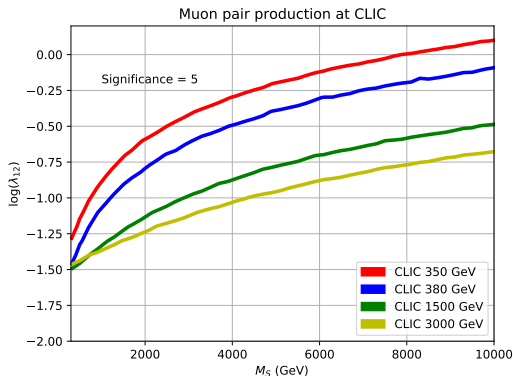
- $S = \frac{N_s}{\sqrt{N_s + N_b}}$

- Beamstrahlung

- Standard acceptance cuts:

$$E(\mu^\pm) > 10 \text{ GeV}$$

$$|\cos(\theta)| < 0.95$$



# Perspective of searches at future colliders

Crivellin, MG, Panizzi, Pruna, Signer, work in progress

(Preliminary plot)

- Electron beam polarization:

$$P_{e^-} = 0.4$$

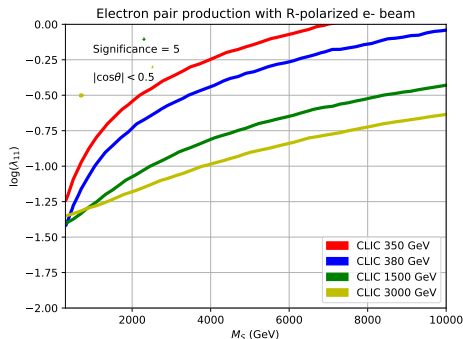
$$P_{e^+} = 0$$

- Angular cut:

$$|\cos(\theta)| < 0.5$$

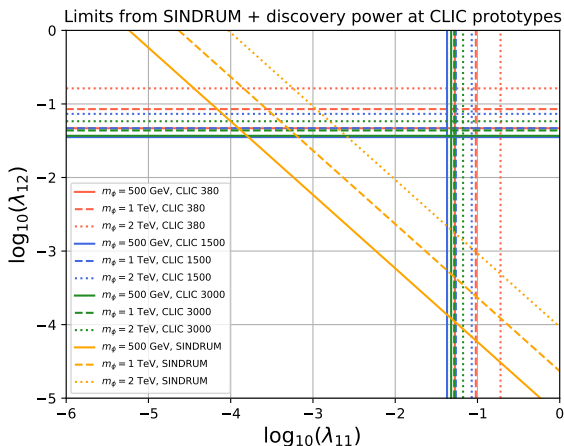
- Integrated luminosity:

350 GeV	380 GeV	1.5 TeV	3 TeV
100 fb <sup>-1</sup>	500 fb <sup>-1</sup>	1500 fb <sup>-1</sup>	3000 fb <sup>-1</sup>



# Limits from low energy and discovery power of CLIC

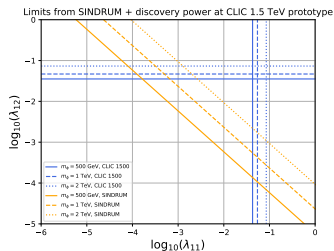
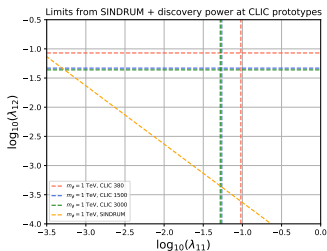
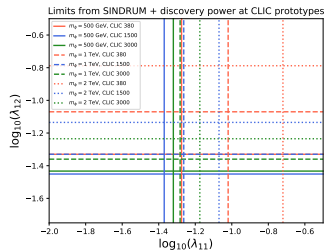
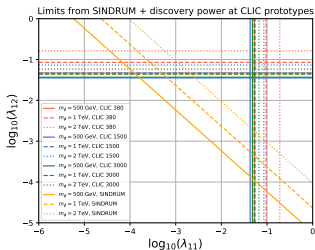
(Preliminary plot)



Crivellin, MG, Panizzi, Pruna, Signer, work in progress



# Limits from low energy and discovery power of CLIC

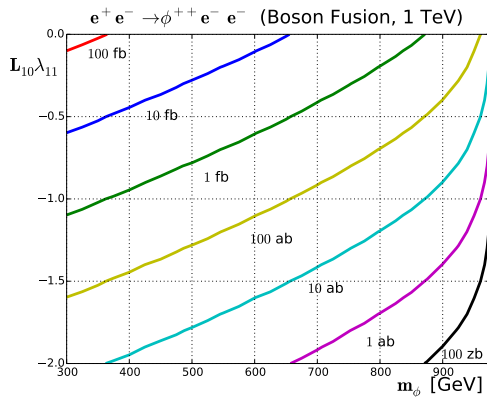


Crivellin, MG, Panizzi, Pruna, Signer, work in progress

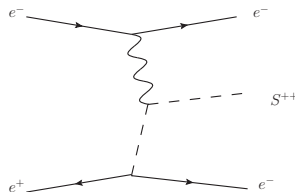
# Direct production

## Single production at ILC

(Preliminary plot)



$$e^+e^- \rightarrow \phi^{++}e^-e^-$$



Crivellin, MG, Panizzi, Pruna, Signer, work in progress

# Outline

- 1 Introduction: the doubly charged scalar
- 2 Low energy: EFT and current limits
- 3 High energy: LHC searches
- 4 High energy: future colliders
- 5 Summary**

# Summary

- Doubly charged scalars arise in **many BSM models**, in triplets or singlets under  $SU(2)_L$ , often in connection with the **neutrino masses**;
- **LFV low energy** processes set strong limits on combination of the DCS couplings to leptons;
- future  $e^+e^-$  **colliders** can provide **complementary bounds**;
- due to the production of the DCS in the **t-channel**, future  $e^+e^-$  colliders can be sensitive to mass scales of several TeV;
- direct searches have been performed at **LHC** by both ATLAS and CMS, setting limits on the **DCS mass** in the range (320, 870) GeV depending on the assumptions;
- a moderately **large width** ( $\Gamma_S/m_S \sim \text{few}\%$ ) can have 10-20% effect on the cross section compared to the NWA;
- further investigations of the DCS phenomenology are ongoing and the results will be published soon.