

# Theory status of offshell Higgs production

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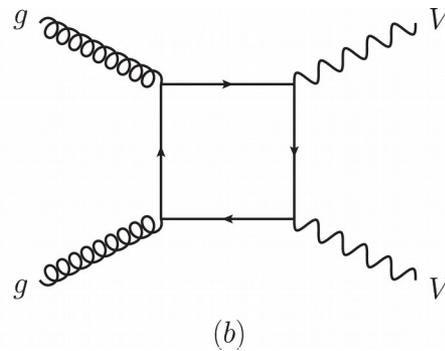
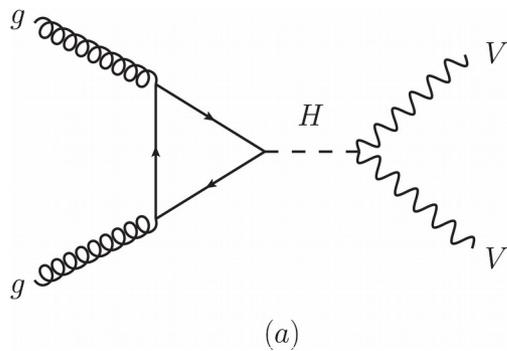
KARLSRUHE INSTITUTE OF TECHNOLOGY

with F. Caola, M. Dowling, K. Melnikov, L. Tancredi, hep-ph/1605.04610

LHC Higgs Cross Section Working Group, Offshell subgroup meeting  
CERN, 22 February 2018

# Offshell Higgs production at LO

Consistent treatment in **offshell region** requires both **signal**  $gg \rightarrow H \rightarrow ZZ$  and **background**  $gg \rightarrow ZZ$  amplitudes



LO: amplitudes known with full mass dependence.

“signal”  $A_H$

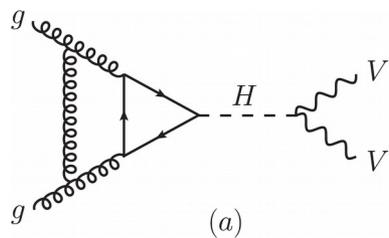
“background”  $A_b$

$$|A_{ZZ}|^2 = |A_H|^2 + |A_b|^2 + 2\text{Re}[A_H A_b^*]$$

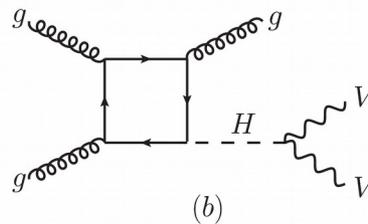
$$\rightarrow \sigma_{\text{full}} = \sigma_{\text{sigl}} + \sigma_{\text{bkgd}} + \sigma_{\text{intf}}$$

# Offshell Higgs production at NLO

Known

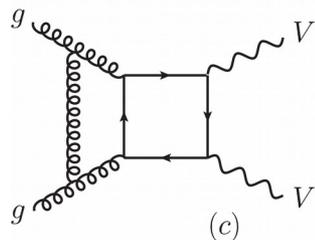


Known



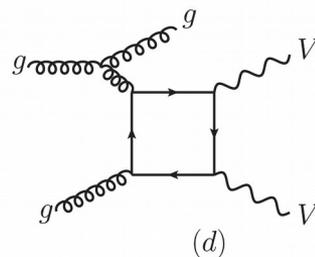
Massless: known

Massive: extremely difficult\*

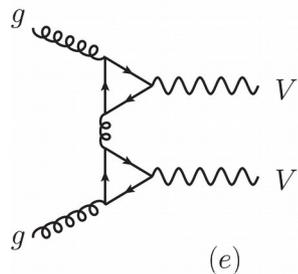


Massless: easy

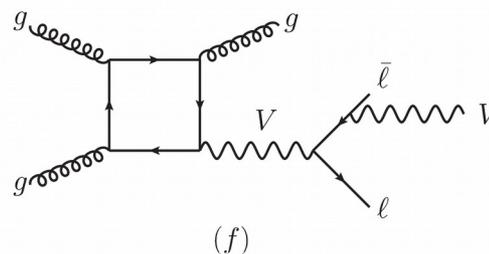
Massive: moderate



Known



Known



\* but possibly within reach now

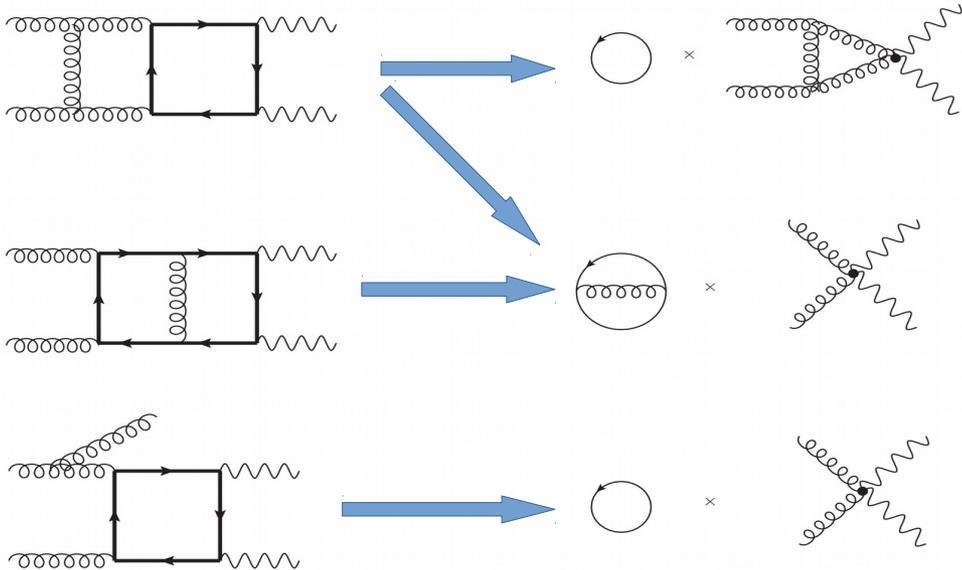
Spira, Djouadi, Graudenz, Zerwas '95; Harlander, Kant '05; Aglietti, Bonciani, Degrossi, Vicini '07; Ellis, Hinchliffe, Soldate, v.d. Bij '88; Caola et al '15, v. Manteuffel, Tancredi '15; Hagiwara, Kuruma, Yamada '91; Campbell, Ellis, Zanderighi '07; v.d. Bij, Glover '89;

# NLO calculations

<p>Campbell, Czakon, Ellis, Kirchner,        hep-ph/1605.01380</p>	<p>Caola, Dowling, Melnikov, RR, Tancredi        hep-ph/1605.04610</p>
<p>Only <b>interference</b>.</p>	<p><b>Signal</b>, <b>background</b>, and <b>interference</b>.</p>
<p>Onshell Z bosons, so <math>m_{ZZ} &gt; 2m_Z</math>.</p>	<p>Offshell Z bosons.</p>
<p>Massive two-loop amplitudes computed in mass expansion to <math>(s/m_t^2)^6</math></p>	<p>Massive two-loop amplitudes computed in mass expansion to <math>(s/m_t^2)^4</math></p>
<p>Massive real emission amplitudes computed exactly – no need for jet veto.</p>	<p>Massive real emission amplitudes computed in mass expansion – <b>need to veto energetic jets</b>.</p>
<p>Results extended beyond <math>2m_t</math> threshold using Padé approximant.</p>	

# $gg \rightarrow (H) \rightarrow ZZ$ : Top Mass Expansion

Expand in  $s/m_t^2$

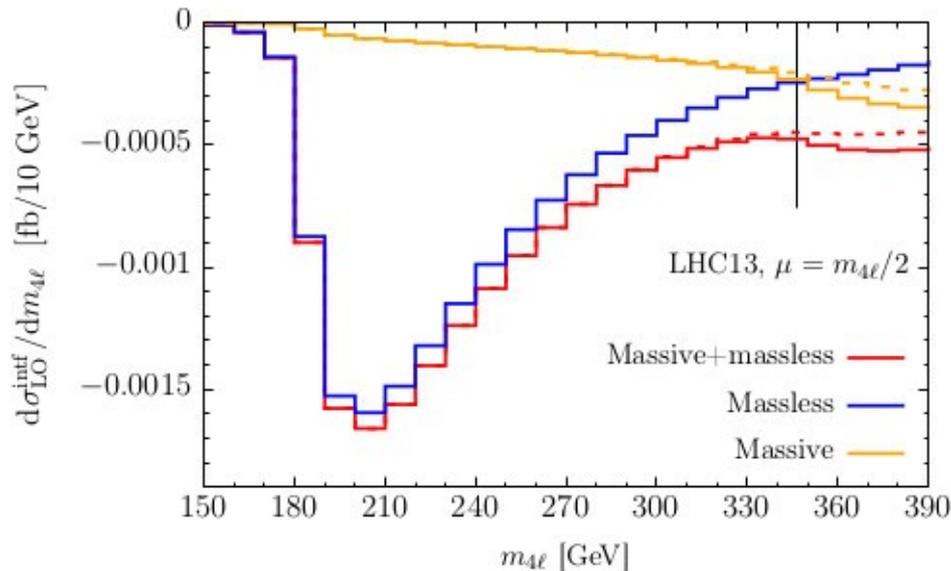


- Keep terms to  $(s/m_t^2)^4$
- Expect to be valid for partonic energies  $s \lesssim 4m_t^2$

Dowling, Melnikov '15

# Validity of top mass expansion

Can check validity at LO – Interference for exact  $m_t$  dependence vs mass expansion



Good approximation below  $2m_t$  threshold



Restricted to  $m_{4\ell} \leq 2m_t$

- Cannot give reliable description of interference effects in high-mass region.

- Large window  $150 \text{ GeV} \lesssim m_{4\ell} \leq 2m_t$  where Higgs is offshell and we can study **interference effects at NLO.**

# Necessity of jet veto

- Real radiation corrections computed in **large mass expansion**.
- These can give rise to **energetic jet**:
  - can lead to  $\sqrt{s} > 2m_t$  even though  $m_{4\ell} < 2m_t$
  - Expansion not valid.
- Comparing to interference effects for  $H+j$  with exact top mass: Campbell, Ellis, Furlan, RR '14
  - ➔ jet requirement  $p_{T,j} < 150 \text{ GeV}$
- **Not** required in calculation of Campbell *et al.*

# Quark-gluon channels

- **Quark-gluon fusion**  $qg \rightarrow ZZq$  appears in NLO corrections.
- **Separately gauge invariant** so *in principle* they could be included.
- *In practice*, they cannot be distinguished from quark-gluon contributions to  $pp \rightarrow ZZ$  at NNNLO.
- **Not included in our calculation** – suppressed by ratio gluon/quark luminosities.
- Contribution to interference  $\sim 4$  smaller than  $gg$  channel  
[Campbell, Ellis, Furlan, RR '14]
- Warrants further study – **any possible sensitivity to this channel should be approached with caution.**

# Parameters

- $gg \rightarrow ZZ \rightarrow e^+e^-\mu^+\mu^-$  at 13 TeV LHC
- Dynamical scale  $\mu_F = \mu_R = \{m_{4\ell}/4, m_{4\ell}/2, m_{4\ell}\}$
- Minimal cuts:
  - $150 \text{ GeV} \leq m_{4\ell} \leq 340 \text{ GeV}$
  - $p_{T,j} < 150 \text{ GeV}$
  - $60 \text{ GeV} \leq m_{\ell\ell} \leq 120 \text{ GeV}$

# $gg \rightarrow (H) \rightarrow ZZ$ Results: Cross Sections

$$\sigma_{\text{LO}}^{\text{signal}} = 0.043_{-0.009}^{+0.012} \text{ fb}, \quad \sigma_{\text{NLO}}^{\text{signal}} = 0.074_{-0.008}^{+0.008} \text{ fb}$$

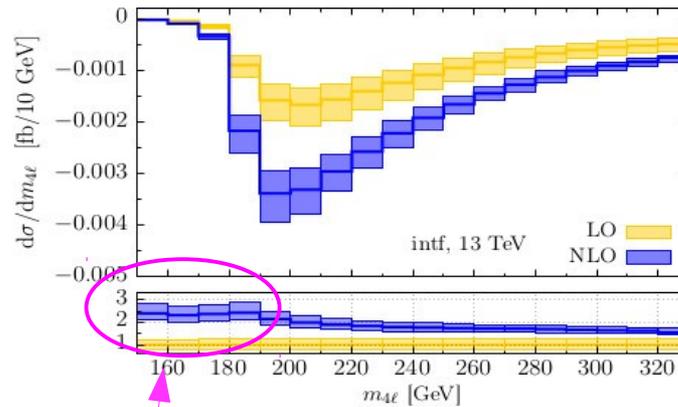
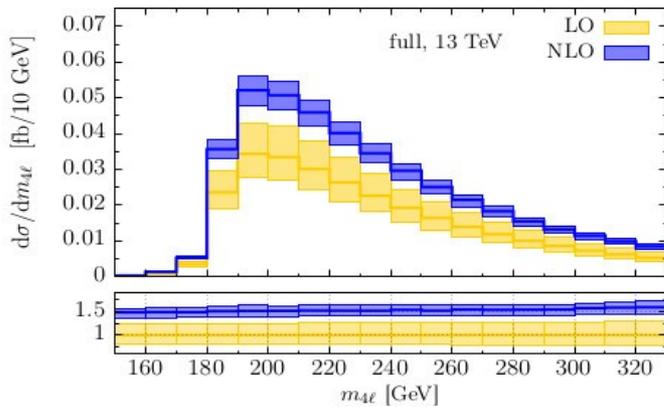
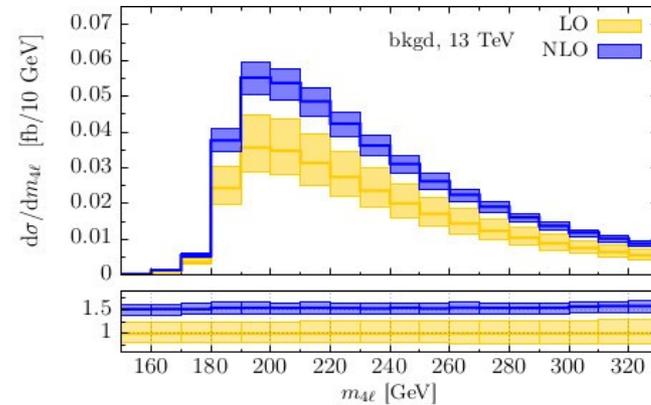
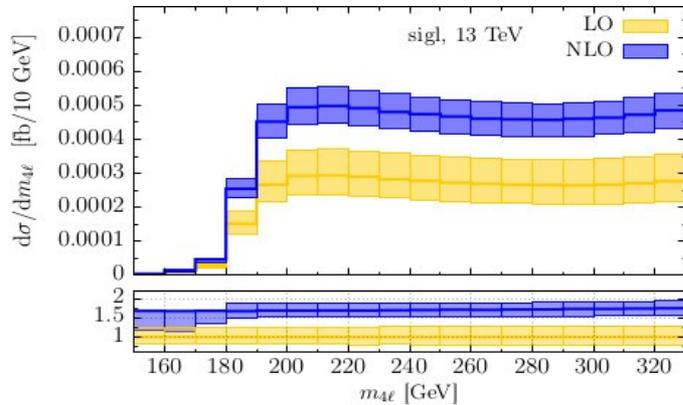
$$\sigma_{\text{LO}}^{\text{bkgd}} = 2.90_{-0.58}^{+0.77} \text{ fb}, \quad \sigma_{\text{NLO}}^{\text{bkgd}} = 4.49_{-0.38}^{+0.34} \text{ fb}$$

$$\sigma_{\text{LO}}^{\text{intf}} = -0.154_{-0.04}^{+0.031} \text{ fb}, \quad \sigma_{\text{NLO}}^{\text{intf}} = -0.287_{-0.037}^{+0.031} \text{ fb}$$

$$\sigma_{\text{LO}}^{\text{full}} = 2.79_{-0.56}^{+0.74} \text{ fb}, \quad \sigma_{\text{NLO}}^{\text{full}} = 4.27_{-0.35}^{+0.32} \text{ fb},$$

- **~ 13k events** at HL-LHC ( $3 \text{ ab}^{-1}$ )
- **Destructive interference ~ 5%**
  - ~ 4 x larger than signal, order of magnitude smaller than background
  - Can use specialized cuts needed to enhance relative to signal and background
- Scale uncertainty: **20%-30% at LO**, **10% at NLO**
- $K_{\text{sigl}} = 1.72$     $K_{\text{bkgd}} = 1.55$     $K_{\text{intf}} = 1.86$

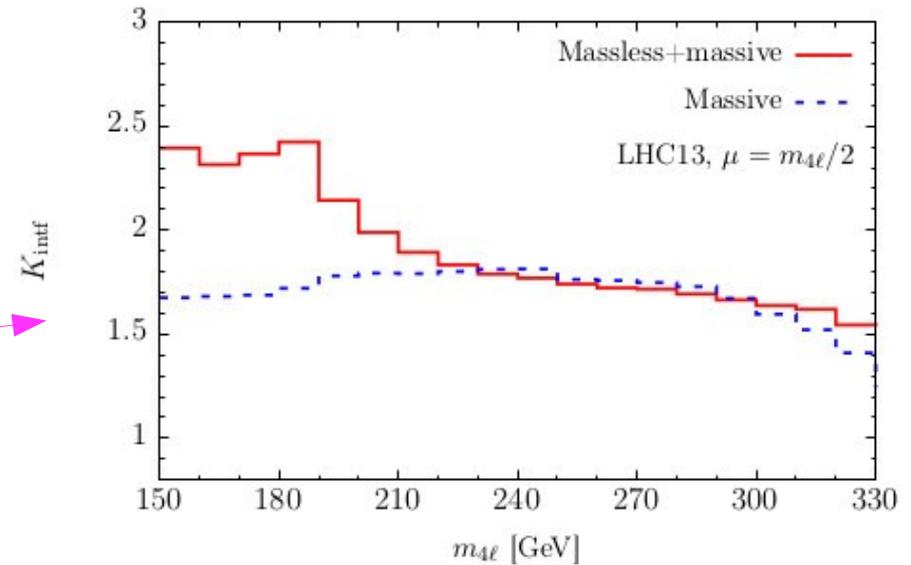
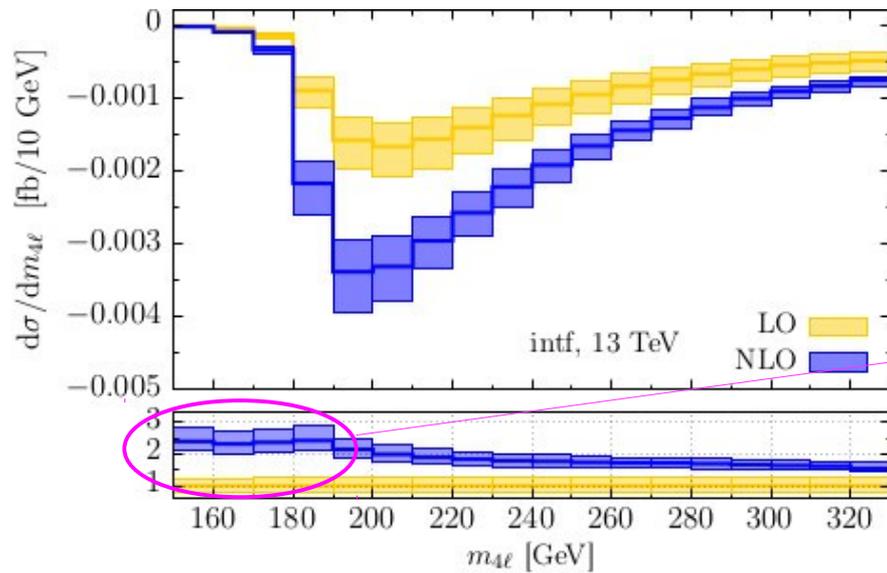
# $gg \rightarrow (H) \rightarrow ZZ$ Results: Mass distributions



K-factors  $\sim 1.5$   
near  $2m_t$   
threshold

- Differential k-factors **relatively flat**...
- Except for interference near  $2m_z$  threshold

# $gg \rightarrow (H) \rightarrow ZZ$ Results: Differential k-factor



- **Massless loop** dominates near  $2m_Z$  threshold, **drives k-factor behavior**

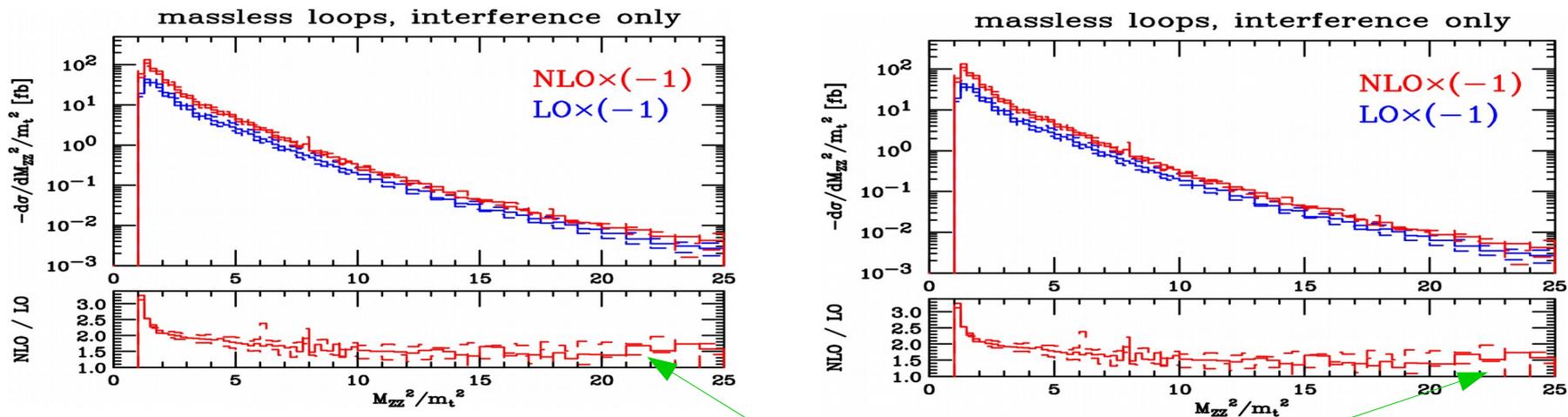
# Going above the top threshold (I)

- Can consider interference with **massless loops** only:
  - Remove restrictions on  $m_{4\ell}$  and jet veto.
  - But **massive loops dominate** interference above  $2m_t$  threshold.
  - Unclear how to estimate uncertainty from neglecting massive loops.
- Numerical calculation of two-loop amplitudes **with full mass dependence** appears feasible:
  - $gg \rightarrow HH$  computed at NLO with full mass dependence.  
[Borowka *et al*, '16]
  - $H+j$  computed at NLO with full top mass dependence.  
[Jones, Kerner, Luisoni '18]

# Going above the top threshold (II)

- Use Padé approximants to extend mass expansion above threshold

Campbell, Czakon, Ellis, Kirchner, '16



K-factor flat in **high-energy tail**

- Comparison of Padé approximants with exact results for  $gg \rightarrow HH$  [Gröber, Maier, Rauh '17]:
  - Padé approximants based on mass expansion do not reproduce exact results very well.
  - Include threshold expansion → dramatic improvement in agreement with exact results.

# Beyond NLO...

NLO is known to be insufficient for onshell Higgs – **is it okay for offshell?**

- NLO results lie **outside** LO scale uncertainty bands
- Do we trust the **NLO scale uncertainty?**
- **Offshell NNLO** impossible at present
- Use NNLO results for infinite  $m_t$  to get **approximate NNLO** k-factor offshell?

# Conclusions

- Difficulty of computing massive two-loop amplitudes  **top mass expansion.**
- Cannot use **massive** interference results in **high energy regimes** (high invariant mass or energetic jets).
- Below top threshold:
  - **Interference** k-factors  $\sim 1.86$ , slightly larger than signal or background.
  - **Interference** k-factor  $\sim 2.5$  around  $m_{ZZ}$  peak; decreases to  $\sim 1.5$  at top threshold.
- Above top threshold: **massless** interference only, **Padé approximants**? *Full top mass dependence within reach???*