Putting Einstein to the test with Gravitational Wave Observations



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- Theoretical Landscape
- Observational landscape
- First look at dynamical strong field gravity
- Polarization
- Future prospects

Outline

Theoretical Landscape

Lovelock theorem: GR is the only Higher dimensions "nice" purely metric theory in 4D. In 4D, the only divergence free symmetric rank-2 tensor constructed only by the metric and its derivatives up to 2nd order and preserving diffeomorphism invariance is the Einstein tensor plus a constant Extra fields Dynamical fields Nondynamical fields (SEP violations) Palatini f(R) Eddington-Born-Infeld Scalars Vectors Scalar-tensor, Metric f(R) Einstein-Aether Horndeski, galileons Horava-Lifshitz Quadratic gravity, n-DBI



[Berti et al, arXiv:1501.07274]

Rouges Gallery of Alternative Theories

Theory	Field	Strong	Massless	Lorentz	Linear	Weak
2	content		graviton	symmetry	$\perp_{\mu\nu}$	EF
Extra scalar field						
Scalar-tensor	S	×	\checkmark	\checkmark	\checkmark	\checkmark
Multiscalar	S	×	\checkmark	\checkmark	\checkmark	\checkmark
Metric $f(R)$	\mathbf{S}	×	\checkmark	\checkmark	1	1
Quadratic gravity						
Gauss-Bonnet	S	×	\checkmark	\checkmark	1	1
Chern-Simons	Р	×	\checkmark	\checkmark	1	\checkmark
Generic	S/P	×	\checkmark	\checkmark	1	\checkmark
Horndeski	S	×	\checkmark	\checkmark	~	\checkmark
Lorentz-violating						
Æ-gravity	\mathbf{SV}	×	\checkmark	×	\checkmark	1
Khronometric/	See 15					
Hořava-Lifshitz	S	×	\checkmark	×	1	1
n-DBI	S	×	\checkmark	×	1	1
Massive gravity	1					
dRGT/Bimetric	SVT	X	×	\checkmark	\checkmark	1
Galileon	S	×	\checkmark	\checkmark	~	1
Nondynamical fields		1750			00	
Palatini $f(R)$	_		\checkmark	\checkmark	×	\checkmark
Eddington-Born-Infeld		1	\checkmark	\checkmark	×	1

[Berti et al, arXiv:1501.07274]

Popular Examples

Ultraviolet Modifications (e.g. EdGB, dCS)

$$S = \frac{1}{16\pi} \int \sqrt{-g} \, d^4x \Big[R - 2\nabla_\mu \phi \nabla^\mu \phi - V(\phi) + S_{\text{mat}} \left[\Psi, \, \gamma(\phi) g_{\mu\nu} \right] \\ + f_1(\phi) R^2 + f_2(\phi) R_{\mu\nu} R^{\mu\nu} + f_3(\phi) R_{\mu\nu\rho\sigma} R^{\mu\nu\rho\sigma} + f_4(\phi)^* RR \Big]$$

Infrared (screened) Modifications (e.g. Horndeski)

$$\begin{split} S &= \int d^4 x \sqrt{-g} \Big\{ K(\phi, X) - G_3(\phi, X) \Box \phi \\ &+ G_4(\phi, X) R + G_{4,X}(\phi, X) \left[(\Box \phi)^2 - (\nabla_\mu \nabla_\nu \phi) (\nabla^\mu \nabla^\nu \phi) \right] \\ &+ G_5(\phi, X) G_{\mu\nu} \nabla^\mu \nabla^\nu \phi - \frac{G_{5,X}(\phi, X)}{6} \left[(\Box \phi)^3 - 3 \Box \phi (\nabla_\mu \nabla_\nu \phi) (\nabla^\mu \nabla^\nu \phi) \right. \\ &+ 2 (\nabla_\mu \nabla_\nu \phi) (\nabla^\mu \nabla_\sigma \phi) (\nabla^\nu \nabla^\sigma \phi) \right] \Big\} \end{split}$$

Observational Landscape



[Baker, Psaltis, Skordis, arXiv:1412.3455]

- Lacking compelling alternatives, prefer generic null tests
- Do GR models leave behind a residual signal?
- Parameterize possible departures from GR waveforms
- Search for additional polarization states
- Ringdown GR predicts unique relationship between harmonics

Null Tests, Generic Tests

Morphology-independent signal reconstruction



[Cornish & Littenberg, arXiv:1410.3835]

Extracting signals without templates



The gravitational-wave event GW150914 observed by the LIGO Hanford (H1, left column panels) and Livingston (L1, olumn panels) detectors. Times are shown relative to September 14, 2015 at 09:50:45 UTC. For visualization, all time series

94% match with GR waveform

(discrepancy consistent with what we expect from noise)

[LVC, PRL**116**, 061102 (2016); PRL **116** 221101, (2016)]

Search for Residual Signals

mon



data

$\mathcal{F} = \sqrt{1 - \frac{\mathrm{SNR}_{\mathrm{res}}^2}{\mathrm{SNR}^2}} > 0.96$

Deviations from GR < 4%



signal





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Uber Gravitationswellen.

Die wichtige Frage, wie die Ausbreitung der Gravitationsfelder erfolgt, ist schon vor anderthalb Jahren in einer Akademiearbeit von mir behandelt worden¹. Da aber meine damalige Darstellung des Gegenstandes nicht genügend durchsichtig und außerdem durch einen bedauerlichen Rechenfehler verunstaltet ist, muß ich hier nochmals auf die Angelegenheit zurückkommen.

Gesamtsitzung vom 14. Februar 1918. — Mitteilung vom 31. Januar

Von A. EINSTEIN.

(Vorgelegt am 31. Januar 1918 [s. oben S. 79].)



Gravitational Wave Observations

- predicted by GR?
- Is the graviton massless?
- Are gravitational waves transverse?
- Additional polarization states?
- Did anyone hear an echo?

Do gravitational waves travel at the speed of light?

Is the emission of energy and angular momentum as

Gravitational Waves Travel at the Speed of Light





LIGO





Frequency (Hz)

[LIGO/Virgo/Fermi, ApJ Lett 848 (2017)]

Gamma rays, 50 to 300 keV

GRB 170817A

 $c_{\rm gw} = 1^{+7 \times 10^{-16}}_{-3 \times 10^{-15}}$

2

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Gravitational-wave strain

-4

GW170817

Time from merger (seconds)

0

-2



Bad news for many dark energy alternatives



[Creminelli Vernizzi, arXiv:1710.05877, Sakstein, Jain, arXiv:1710.05893, Baker et al, arXiv:1710.06394]

Many models that could potentially explain the accelerated expansion yet evade solar system constraints via screening have been ruled out

One way out might be to build models that change behavior with length scale rather than density since $\lambda_{
m gw} \ll R_H, ~D_L \ll R_H$

[Battye, Pace, Trinh, arXiv:1802.09447]



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Is the emission of energy and angular momentum as predicted by GR?

MM

Post-Newtonian Inspiral

$$\mathcal{A}_{\rm GR}(f) = \sqrt{\frac{5}{96}} \frac{\mathcal{M}^{5/6}}{D \pi^{2/3}} f^{-7/6}$$

 $h(f) = \mathcal{A}(f) e^{i\Psi(f)}$

$$\Psi_{\rm GR} = 2\pi f t_c + \Phi_c + \frac{\pi}{4} + \sum_{k=-5} \psi_k \, u^k$$

 $u = (\pi \mathcal{M}f)^{1/3} \sim v$



Modified Waveforms

$$\mathcal{A}(f) = \sqrt{\frac{5}{96}} \frac{\mathcal{M}^{5/6}}{D \pi^{2/3}} f^{-7/6} \left(1 - \frac{5}{512} \dot{G} \mathcal{M} u^{-8} + \left(\frac{743}{672} + \frac{11}{8} \eta \right) \eta^{-2/5} u^2 + \dots \right)$$

$$\text{Variable G} \qquad \text{Scalar Field}$$

$$\mathbb{P}(f) = 2\pi f t_c - \Phi_c - \frac{\pi}{4} + \frac{3}{128} u^{-5} \left[1 - \frac{25}{1536} \dot{G} \mathcal{M} u^{-8} - \frac{5}{84} \frac{S^2}{\omega_{\text{BD}}} \eta^{3/5} u^{-2} + \left(\frac{3715}{756} + \frac{55}{9} \eta \right) \eta^{-2/5} u^2 - 16\pi \eta^{-3/5} u^3 - \frac{128}{3} \frac{\pi^2 D \mathcal{M}}{\lambda_g^2 (1+z)} u^2 + \dots \right]$$

$$= \sqrt{\frac{5}{96}} \frac{\mathcal{M}^{5/6}}{D \pi^{2/3}} f^{-7/6} \left(1 - \frac{5}{512} \dot{G} \mathcal{M} u^{-8} + \left(\frac{743}{672} + \frac{11}{8} \eta \right) \eta^{-2/5} u^2 + \dots \right)$$
Variable G
$$= 2\pi f t_c - \Phi_c - \frac{\pi}{4} + \frac{3}{128} u^{-5} \left[1 - \frac{25}{1536} \dot{G} \mathcal{M} u^{-8} - \frac{5}{84} \frac{S^2}{\omega_{\rm BD}} \eta^{3/5} u^{-2} + \left(\frac{3715}{756} + \frac{55}{9} \eta \right) \eta^{-2/5} u^2 - 16\pi \eta^{-3/5} u^3 - \frac{128}{3} \frac{\pi^2 D \mathcal{M}}{\lambda_g^2 (1+z)} u^2 + \dots \right]$$

$$h(f) = (1 + \delta \mathcal{A}(f)) e^{i\delta\Psi(f)} h_{\rm GR}(f)$$

Parametrized Post-Einsteinian

Massive Graviton

[Yunes, Pretorius, arXiv:0909.3328] [Arun, Iyer, Qusailah, Sathyaprakash, arXiv:gr-qc/0604018]

No deviations seen from GR phasing



 $\frac{\delta\psi_k}{\psi_k}$

[LVC, arXiv:1602.03841, arXiv:1606.04856]

Merger-ringdown

Caveat emptor:

These bounds are term-by-term (not jointly allowed to vary)

Bounds would be weaker if all terms allowed to vary together

[Sampson, Cornish, Yunes, arXiv:1303.1185]



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Is the graviton massless?

"Chirp Squeezing" - higher frequency signal from near the merger travels faster and arrive earlier

$$\Psi(f) = 2\pi f t_c - \Phi_c - \frac{\pi}{4} + \frac{3}{128} u^{-5} \left[1 + \left(\frac{3715}{756} \right) \right]$$

$$m_g < 1.2 \times 10^{-22} \text{ eV}$$

[LVC, arXiv:1602.03841, arXiv:1606.04856]



Partially degenerate with changing the total mass Degeneracy mostly broken by higher PN terms and merger

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Alternative Theories Predict additional Polarization States

 \hat{i}



 \hat{k} /

 \mathcal{U}

 $e_{TT}^{+} = u \otimes u - v \otimes v$ $e_{TT}^{\times} = u \otimes v + v \otimes u$ $e_{ST}^{\odot} = u \otimes u + v \otimes v$ $e_{SL}^{\odot} = k \otimes k$

 $\mathbf{e}_{\mathrm{VL}}^{u} = \mathbf{u} \otimes \mathbf{k} + \mathbf{k} \otimes \mathbf{u}$ $\mathbf{e}_{\mathrm{VL}}^{v} = \mathbf{v} \otimes \mathbf{k} + \mathbf{k} \otimes \mathbf{v}$

Antenna Patterns Pulsar Timing

Interferometers





[Yunes & Siemens, Living Rev.Rel. 16, 9 (2013)]



Binary Black Hole Merger GW170814 - First Constraints on Polarization



[LIGO/Virgo PRL 119, 141101 (2017)]

week ending 6 OCTOBER 2017

Binary Black Hole Merger GW170814 - First Constraints on Polarization



[Isi & Weinstein, arXiv:1710.03794 (2017)]

[LIGO/Virgo PRL 119, 141101 (2017)]

Tensor favored over Vector by 200:1 Tensor favored over Scalar by 1000:1



Binary Black Hole Merger GW170814 - BayesWave Constraints on Polarization



Bayes Factors between models similar to those form template based analysis

Binary Black Hole Merger GW170814 - First Constraints on Polarization







LIGO/Virgo Antenna Patterns



First search for nontensorial gravitational waves from known pulsars



[LIGO/Virgo PRL 120, 031104 (2018)]



No detections so far, so bounds are weak

Note: The analysis assumed that $f_{gw} = 2f_{orb}$, which is not always the case for dipole emission





A Search for Tensor, Vector, and Scalar Polarizations in the Stochastic Gravitational-Wave Background

$$\langle \tilde{s}_1(f) \tilde{s}_2^*(f') \rangle = \delta(f - f') \sum_A \gamma_A(f) H^A(f).$$

$$\Omega(f) = \frac{20\pi^2}{3H_0^2} f^3 H(f)$$

$$\Omega(f) = \Omega_0^T \left(\frac{f}{f_0}\right)^{\alpha_T} + \Omega_0^V \left(\frac{f}{f_0}\right)^{\alpha_V} + \Omega_0^S \left(\frac{f}{f_0}\right)^{\alpha_S}$$

[LIGO/Virgo arXiv:1802.10194 (2018)]



A Search for Tensor, Vector, and Scalar Polarizations in the Stochastic Gravitational-Wave Background

$$\Omega(f) = \Omega_0^T \left(\frac{f}{f_0}\right)^{\alpha_T} + \Omega_0^V \left(\frac{f}{f_0}\right)^{\alpha_V} + \Omega_0^S \left(\frac{f}{f_0}\right)^{\alpha_S}$$

LIGO O1 Limits: (Marginalized over spectral slope) $f_0 = 25 \text{Hz}$

$$\Omega_0^T < 2.0 \times 10^{-7} \qquad \Omega_0^V$$

[LIGO/Virgo arXiv:1802.10194 (2018)]

 2.5×10^{-7} $\Omega_0^S < 8.4 \times 10^{-7}$



PulsarTiming



(c) M.KrJmer



PTA Two-Point Correlations



$$S_{ab}(f) = \frac{\Gamma_{ab}^{\mathrm{TT}} \mathcal{A}_{\mathrm{Q}}^2 (f/f_y)^{-4/3} + (\Gamma_{ab}^{\mathrm{ST}} \mathcal{A}_{\mathrm{ST}}^2 + \Gamma_{ab}^{\mathrm{VL}} \mathcal{A}_{\mathrm{ST}}^2}{(1 + \kappa^2 (f/f_y)^{-2/3})^{-2/3}}$$



[Cornish, O'Beirne, Taylor, Yunes, PRL (2018)]

Upper Limit using NANOGrav 9-year results

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 $m_g < 1.2 \times 10^{-22} \text{ eV}$

no evidence to contrary

Echoes from the abyss?

BH perturbation theory - structure of QN modes depends on inner boundary conditions. Any reflections from inside the light ring produce echoes



[Cardoso, Hopper, Macedo, Palenzuela, Pani, arXiv:1608.08637]



Echoes from the abyss?



I'm waiting for the analysis to be re-done using the specialized BayesWave search

 2.5σ hint of detection claimed using LIGO O1 detections [Abedi, Dykaar, Afshordi, arXiv:1612.00266]

More careful reanalysis by some LIGO researchers found nothing [Westerweck et al, arXiv:1712.09966]

 4.2σ claim of detection using BNS GW170817 data [Abedi, Afshordi, arXiv:1803.10454]

Detection of echoes also claimed for GW151226, GW170104, GW170606, GW170814, GW170817 [Conklin, Holdom, Ren, arXiv:1712.06517]

[Tsang et al, arXiv:1804.04877]

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Future Prospects

Next steps - a worldwide network













3rd and 4th generation ground-based instruments



A+: aLIGO upgrade, freq. dep. squeezing, heavier mirrors, more powerful lasers
Voyager: aLIGO upgrade, same facility, cryogenic, more powerful lasers
Einstein Telescope: Underground, 10 km, triangular, cryogenic
Cosmic Explorer: New facility, 40 km arms, squeezing etc

The International Pulsar Timing Array

















Next steps: Chime, FAST, MeerKAT, and the SKA



June 2017, LISA mission selected by ESA

 10^{-17} 10^{-18} 10^{-18} 10^{-19} 10^{-20} 10^{-20} 10^{-21} 10^{-21} 10^{-21} $0^{\text{Deservatory}}$ Characteristic Strain $- \cdot \text{Total}$ 10^{-4}

10⁻¹⁶

Lead Proposer Prof. Dr. Karsten Danzmann

