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### Higgs mass, strong CP problem, GUT

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#### $V = \lambda_{\rm SM} |H|^4 - m_H^2 |H|$ 2

#### Today's topic  $V = \left( \lambda_{\rm SM} \right) H \vert^4 - m_H^2 \vert H \vert^4$ 2

Assume that the SM is valid up to high energy scale



#### Small boundary condition



# Some new physics to explain λ = 0?



 $H \leftrightarrow H'$ Introduce Z2 symmetry  $SU(2) \leftrightarrow SU(2)$ <sup>'</sup>

#### $V(H, H') = \lambda(|H|^2 + |H'|^2)$  $(2)^2 + \lambda' |H|^2 |H'|^2 - m^2(|H|^2 + |H'|^2)$ 2)

Let us assume  $m >> v_{EW}$ 

 $V(H, H') = \lambda(|H|^2 + |H'|^2)$  $(2)^2 + \lambda' |H|^2 |H'|^2 - m^2(|H|^2 + |H'|^2)$ 2)

$$
\langle H' \rangle^2 = \frac{m^2}{2\lambda} \qquad m_H^2 \simeq 0 \to \lambda' \simeq 0
$$

 $V(H, H') \simeq \lambda(|H|^2 + |H'|^2)$  $(2)^2 - m^2(|H|^2 + |H'|^2)$ 2)

Accidentally SU(4) symmetric

 $SU(4)$  ->  $SU(3)$  by  $\langle H' \rangle$ 

SM Higgs is a Nambu-Goldstone boson

 $\lambda_{\rm SM}=0$ 





### Fine-tuning

 $V(H, H') = \lambda(|H|^2 + |H'|^2)$  $(2)^2 + \lambda' |H|^2 |H'|^2 - m^2(|H|^2 + |H'|^2)$ 2)



Same as that of SM

#### Fermions, gauge groups

 $q \leftrightarrow q' = (\bar{u}, \bar{d}), \ell \leftrightarrow \ell' \supset \bar{e}$ 

 $SU(3)_c \times SU(2)_L \times SU(2)_R \times U(1)_{B-L} \supset SO(10)$ 

 $q,\bar{u},\bar{d},q',\bar{u}',\bar{d}',\cdots$ 

 $SU(3)_c \times SU(3)_c'$  $SU(3)_c$  *U*(1)  $U(1) \times U(1)$  $\propto SU(2)_L \times SU(2)' \times \text{or}$ 



### $Z_2$  from  $SO(10)$

#### Remnant of SO(10) *SO*(10)  $H, H' \subset 16$  $q, \ell, q', \ell' = 16$  $q'=(\bar{u},\bar{d})$  $d$ ,  $\ell' \supset \overline{e}$

 $SU(3)_c \times SU(2)_L \times SU(2)_R \times U(1)_{B-L}$  $SU(3)_c \times SU(2)_L \times U(1)_Y$  $\langle H' \rangle \neq 0$ Left-right symmetry  $H \leftrightarrow H'$ 





### Top-down perspective

SUSY GUT

3 parameters 4 parameters  $g_{GUT}$ ,  $M_{GUT}$ ,  $m_{SUSY}$  *g*<sub>1</sub>, *g*<sub>2</sub>, *g*<sub>3</sub>, *v*<sub>EW</sub>



GUT here

4 parameters

 $g_{\text{GUT}},~M_{\text{GUT}},~v',~y_t$ 



5 parameters

 $g_1, g_2, g_3, y_t, \lambda_{\text{higgs}}$ 

### Top-down perspective

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GUT here Altanative to SUSY GUT ?

4 parameters

 $g_{\text{GUT}},~M_{\text{GUT}},~v',~y_t$ 



5 parameters

 $g_1, g_2, g_3, y_t, \lambda_{\text{higgs}}$ 

#### Intermediate Pati-Salam *SO*(10)  $H, H' \subset 16$  $q, \ell, q', \ell' = 16$  $q'=(\bar{u},\bar{d})$  $d$ ,  $\ell' \supset \bar{e}$

 $SU(4) \times SU(2)_L \times SU(2)_R$ 

 $SU(3)_c \times SU(2)_L \times U(1)_Y$  $\langle H'\rangle =$  $\sqrt{2}$  $0 \quad 0 \quad 0 \quad v'$ 000 0  $H(1,2,1,-\frac{1}{2}) \subset (4,2,1)$ 1 2  $) \subset (4, 2, 1)$  $H'({\bf 1}, 1, 2,$ 1 2  $) \subset (\bar{4}, 1, 2)$ 

### Coupling Unification







v' / GeV

~2σ smaller top mass





### Parity and the strong CP problem

(SO(10) is not required)

Parity

 $H(t, x) \leftrightarrow H'(t, -x)$ 



 $q(t, x) \leftrightarrow i\sigma_2q$  $\overline{\phantom{a}}$  $*(t, -x)$ 

Assume  $SU(3) \leftrightarrow SU(3)$ 

 $G\tilde{G} \rightarrow -G\tilde{G}$ 

 $\theta_{\rm QCD}=0$ 

#### Yukawa coupling?

Ex. Left-Right symmetry

 $SU(3)_c \times SU(2)_L \times SU(2)_R \times U(1)_{B-L}$ 

$$
\frac{c_{ij}}{M}HH'q_iq'_j \qquad q(t,x) \leftrightarrow i\sigma_2q'^{*}(t,-x)
$$

$$
c=c^\dagger, \arg(\det[c])=0
$$

Also, H and H' has no physical phase dof.

### Parity solutions

1978, Beg and Tsao, Mohapatra and Senjanovic

Parity can solve the strong CP problem,  $H(2,2)$ . Dangerous contribution from complex phase in the Higgs vev (1991, Barr, Chang and Senjanovic)

1989, Babu and Mohapatra

with soft Z<sub>2</sub> breaking  $H(2,1) + H'(1,2)$ 

#### Embedding into SO(10)

 $q(t,x) \leftrightarrow q'(t,x)$ Part of  $SO(10)$ 

 $q(t, x) \leftrightarrow i\sigma_2 q^*(t, -x)$ 

 $q(t, x) \leftrightarrow i\sigma_2q$  $\overline{\phantom{a}}$  $*(t, -x)$ 

 $SO(10) \times CP \stackrel{\phi_{45}^-}{\longrightarrow} SU(3) \times SU(2)_L \times SU(2)_R \times U(1)_{B-L} \times P_{LR}$ 

#### Loop correction to  $\theta$

#### Suppressed by loop factors, flavor mixing

 $\delta\theta \sim 10^{-11}$ 



#### Fermions

Doublets have Z2 partners

 $q, \ell \leftrightarrow q', \ell'$ TABLE II. Doublet fields: the four possible *SU*(3)*<sup>c</sup>* ⇥ *SU*(2)*<sup>L</sup>* ⇥ *SU*(2)<sup>0</sup> ⇥ *U*(1) assignments for the

$$
q({\bf 3},2,1,\frac{1}{6}), \ell({\bf 1},2,1,-\frac{1}{2}), H({\bf 1},2,1,-\frac{1}{2})
$$

*Z*<sup>2</sup> partners of *q*(3*,* 2*,* 1*,*

$$
\begin{array}{|c|c|c|c|c|} \hline & {\rm A}(-,-) & {\rm B}(+,-) & {\rm C}(-,+) & {\rm D}(+,+) \\ \hline q' & (\mathbf{\bar{3}},1,2,-\tfrac{1}{6}) & (\mathbf{3},1,2,-\tfrac{1}{6}) & (\mathbf{\bar{3}},1,2,\tfrac{1}{6}) & (\mathbf{3},1,2,\tfrac{1}{6}) \\ \ell',H' & (\mathbf{1},1,2,\tfrac{1}{2}) & (\mathbf{1},1,2,\tfrac{1}{2}) & (\mathbf{1},1,2,-\tfrac{1}{2}) & (\mathbf{1},1,2,-\tfrac{1}{2}) \\ \hline \end{array}
$$

*Z*<sup>2</sup> partners of *q*(3*,* 2*,* 1*,* <sup>6</sup> ), `(1*,* <sup>2</sup>*,* <sup>1</sup>*,* <sup>1</sup> <sup>2</sup> ) and *<sup>H</sup>*(1*,* <sup>2</sup>*,* <sup>1</sup>*,* <sup>1</sup> 2 ). A(*,* ) B(+*,* ) C(*,* +) D(+*,* +) *<sup>q</sup>*<sup>0</sup> (¯3*,* <sup>1</sup>*,* <sup>2</sup>*,* <sup>1</sup> <sup>6</sup> ) (3*,* <sup>1</sup>*,* <sup>2</sup>*,* <sup>1</sup> <sup>6</sup> ) (¯3*,* <sup>1</sup>*,* <sup>2</sup>*,* 1 <sup>6</sup> ) (3*,* 1*,* 2*,* 1 6 ) `0 *, H*0 (1*,* 1*,* 2*,* 1 <sup>2</sup> ) (1*,* 1*,* 2*,* 1 <sup>2</sup> ) (1*,* <sup>1</sup>*,* <sup>2</sup>*,* <sup>1</sup> <sup>2</sup> ) (1*,* <sup>1</sup>*,* <sup>2</sup>*,* <sup>1</sup> 2 ) *q*(3*,* 2*,* 1*,* 1 6 )*,* `(1*,* <sup>2</sup>*,* <sup>1</sup>*,* <sup>1</sup> 2 )*, H*(1*,* <sup>2</sup>*,* <sup>1</sup>*,* <sup>1</sup> 2 )

TABLE II. Doublet fields: the four possible *SU*(3)*<sup>c</sup>* ⇥ *SU*(2)*<sup>L</sup>* ⇥ *SU*(2)<sup>0</sup> ⇥ *U*(1) assignments for the

q', l' are identified with  $SU(2)_L$  singlet SM fermions A: almost vector-like and anomaly free.

of *SU*(2)*<sup>L</sup>* as *SU*(2)*R*. In the second sub-section we study models B, C and D.  $SU(2)' = SU(2)_R, U(1) \sim U(1)_{B-L}$ 

$$
\mathcal{L} = \frac{1}{M} (q \tilde{y_u} q') H^{\dagger} H^{'\dagger} + \frac{1}{M} (q \tilde{y_d} q') H H'
$$

1

 $\mathcal{Y}$  yukawa couplings, there are interactions between fermions and scalars at dimension  $\mathcal{Y}$ 

1

*H*0

1

*Z*<sup>2</sup> partners of *q*(3*,* 2*,* 1*,* <sup>6</sup> ), `(1*,* <sup>2</sup>*,* <sup>1</sup>*,* <sup>1</sup> <sup>2</sup> ) and *<sup>H</sup>*(1*,* <sup>2</sup>*,* <sup>1</sup>*,* <sup>1</sup> 2 ). A(*,* ) B(+*,* ) C(*,* +) D(+*,* +) *<sup>q</sup>*<sup>0</sup> (¯3*,* <sup>1</sup>*,* <sup>2</sup>*,* <sup>1</sup> <sup>6</sup> ) (3*,* <sup>1</sup>*,* <sup>2</sup>*,* <sup>1</sup> <sup>6</sup> ) (¯3*,* <sup>1</sup>*,* <sup>2</sup>*,* 1 <sup>6</sup> ) (3*,* 1*,* 2*,* 1 6 ) `0 *, H*0 (1*,* 1*,* 2*,* 1 <sup>2</sup> ) (1*,* 1*,* 2*,* 1 <sup>2</sup> ) (1*,* <sup>1</sup>*,* <sup>2</sup>*,* <sup>1</sup> <sup>2</sup> ) (1*,* <sup>1</sup>*,* <sup>2</sup>*,* <sup>1</sup> 2 ) *q*(3*,* 2*,* 1*,* 1 6 )*,* `(1*,* <sup>2</sup>*,* <sup>1</sup>*,* <sup>1</sup> 2 )*, H*(1*,* <sup>2</sup>*,* <sup>1</sup>*,* <sup>1</sup> 2 )

TABLE II. Doublet fields: the four possible *SU*(3)*<sup>c</sup>* ⇥ *SU*(2)*<sup>L</sup>* ⇥ *SU*(2)<sup>0</sup> ⇥ *U*(1) assignments for the

c C C and D and D and D and D and the first sub-section we study model A  $N$  form i.o. notes that *Z*  $N$  **Example 1.** B, C, D: needs extra fermion which are identified with  $SU(2)$ <sub>L</sub> singlet SM fermions

of *SU*(2)*<sup>L</sup>* as *SU*(2)*R*. In the second sub-section we study models B, C and D.

$$
\mathcal{L} = Hq\bar{u} + H'q'\bar{u}' + \cdots
$$

$$
\bar{u} \leftrightarrow \bar{u}'
$$

1

 $\mathcal{Y}$  yukawa couplings, there are interactions between fermions and scalars at dimension  $\mathcal{Y}$ 

1

*H*0

1

#### Parity and the strong CP problem  $q(t, x) \leftrightarrow i\sigma_2q$  $\overline{\phantom{a}}$  $^*(t, -x)$

#### Model A

$$
\mathcal{L} = \frac{1}{M} (q\tilde{y_u}q')H^{\dagger}H^{'\dagger} + \frac{1}{M} (q\tilde{y_d}q')HH' + \text{h.c.}
$$
  

$$
\tilde{y}^{\dagger} = \tilde{y}, \text{ real } \det \tilde{y}
$$
  

$$
\theta G\tilde{G} : \theta = 0
$$

Strong CP problem is solved!

#### Parity and the strong CP problem  $q(t, x) \leftrightarrow i\sigma_2q$  $\overline{\phantom{a}}$  $^*(t, -x)$

Model B,C

 $\mathcal{L} = yHQ\bar{u} + y^*H'Q'\bar{u}' + y^*H^\dagger Q^\dagger\bar{u}^\dagger + yH^\dagger Q^{'\dagger}\bar{u}^\dagger$  $det y \times det y^*$  is real

#### Cancellation between the SM and partners

$$
Parity and the strong CP problem
$$
\n
$$
q(t, x) \leftrightarrow i\sigma_2 q' * (t, -x)
$$

Model D

$$
\frac{q({\bf 3},2,1,\frac{1}{6}),\ell({\bf 1},2,1,-\frac{1}{2}),H({\bf 1},2,1,-\frac{1}{2})}{A(-,-)}\frac{B(+,-)}{B(+,-)}\frac{C(-,+)}{C(-,+)}\frac{D(+,+)}{D(+,+)}\\ q'\left| (\mathbf{\bar{3}},1,2,-\frac{1}{6})\right|({\bf \bar{3}},1,2,-\frac{1}{6})}(\mathbf{\bar{3}},1,2,\frac{1}{6})\frac{({\bf \bar{3}},1,2,\frac{1}{6})}{({\bf 1},1,2,-\frac{1}{2})}\frac{({\bf 3},1,2,\frac{1}{6})}{({\bf 1},1,2,-\frac{1}{2})}
$$

in models B, C and D. In the first sub-section we study Model A and identify the *Z*<sup>2</sup> partner  $\mathcal{L} = yHq\bar{u} + y^*H'q'\bar{u}' + \lambda Hq\bar{u}' + \lambda^*H'q'\bar{u} + \text{h.c.}$ 

$$
\det\begin{pmatrix}y & \lambda \\ \lambda^* & y^*\end{pmatrix} = \det\begin{pmatrix}\lambda & y \\ y^* & \lambda^*\end{pmatrix} = \det\begin{pmatrix}y^* & \lambda^* \\ \lambda & y\end{pmatrix} = \det\begin{pmatrix}y & \lambda \\ \lambda^* & y^*\end{pmatrix}^*
$$

## Yukawa couplings

X

TABLE III. Possible *X* particles for generating Yukawa couplings in Model A.



### Yukawa couplings

X

Small enough not to blow up the gauge coupling



### CKM phase

 $SO(10) \times CP \stackrel{\phi_{45}^{-}}{\longrightarrow} SU(3) \times SU(2)_{L} \times SU(2)_{R} \times U(1)_{B-L} \times P_{LR}$ 

Real yukawa without CP symmetry breaking

A simple example

 $\mathcal{L} = \left( M^{ij} + i \lambda^{ij} \phi_{45} \right) X_{10,i} X_{10,j}$ 

 $m^2 \ll \Lambda_c^2$ cut

#### More Fine-tuned than SM?

No.

$$
\frac{v_{\text{EW}}^2}{m^2} \times \frac{m^2}{\Lambda_{\text{cut}}^2} \sim \frac{v_{\text{EW}}^2}{\Lambda_{\text{cut}}^2}
$$

$$
y \simeq 2\lambda \quad m^2 \ll \Lambda_{\text{cut}}^2
$$

#### How non-trivial?

#### Ex. SUSY GUT

3 parameters

*g*GUT*, M<sub>GUT</sub>*, *m<sub>SUSY</sub>* 



4 parameters

*g*1*, g*2*, g*3*, v*EW

#### How non-trivial?

4 parameters

 $g_{\text{GUT}},~M_{\text{GUT}},~v',~y_t$ 



5 parameters

 $g_1, g_2, g_3, y_t, \lambda_{\text{higgs}}$ 

Altanative to SUSY GUT ?

### Coupling Unification





Correction to the gauge coupling unification by high dimensional operator

 $SO(10)$   $\overset{\phi_{210}}{\longrightarrow}$  $\longrightarrow$  *SU*(3)  $\times$  *SU*(2)<sub>*L*</sub>  $\times$  *SU*(2)<sub>*R*</sub>  $\times$  *U*(1)<sub>*B*-*L*</sub>  $\times$  *C*<sub>*LR*</sub>

$$
\frac{210^{abcd}}{M_*}F_{10}^{ab}F_{10}^{cd} \qquad \Delta\left(\frac{2\pi}{\alpha}\right) \lesssim 10
$$

 $SO(10) \times CP \stackrel{\phi_{45}}{\longrightarrow} SU(3) \times SU(2)_L \times SU(2)_R \times U(1)_{B-L} \times P_{LR}$ 

$$
\frac{45^{ac}}{M_{*}}\frac{45^{bd}}{M_{*}}F_{10}^{ab}F_{10}^{cd} \qquad \Delta\left(\frac{2\pi}{\alpha}\right) \lesssim 1
$$

Correction to the gauge coupling unification by high dimensional operator

 $SO(10)$   $\frac{\phi_{54}}{2}$  $\longrightarrow$  *SU*(4)  $\times$  *SU*(2)<sub>*L*</sub>  $\times$  *SU*(2)<sub>*R*</sub>  $\times$  *C<sub>LR</sub>* 

$$
\frac{54^{ab}}{M_{*}}F_{10}^{ac}F_{10}^{bc} \qquad \Delta\left(\frac{2\pi}{\alpha}\right) \lesssim 1
$$

 $SO(10) \times CP \stackrel{\phi_{210}}{\longrightarrow} SU(4) \times SU(2)_L \times SU(2)_R \times P_{LR}$ 

210  $M_*$ 210  $M_*$  $F_{10}F_{10}$   $\Delta$ 

$$
\Delta\left(\frac{2\pi}{\alpha}\right)\ll 1
$$