

# Disorder at the LHC

Matthew Low

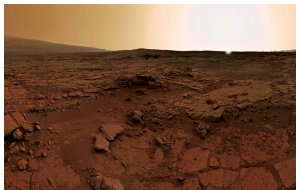
Institute for Advanced Study

*in progress with*

Raffaele Tito D'Agnolo

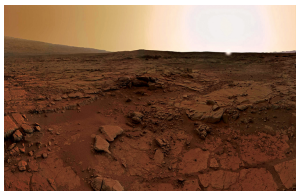
- ▶ Why haven't we seen anything at the LHC?

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  - ▶ *Possibility 1*: There is nothing to see.

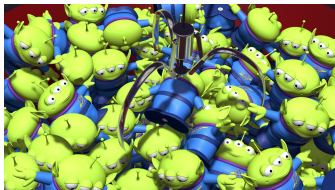


# Motivation

- ▶ Why haven't we seen anything at the LHC?
  - ▶ *Possibility 1:* There is nothing to see.



- ▶ *Possibility 2:* There are too many new particles.

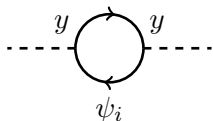


# Large $N$

- ▶ Consider  $N$  new states



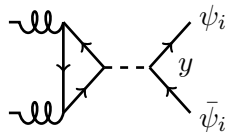
- ▶ With  $\mathcal{L}_{\text{int}} = hy\bar{\psi}_i\psi_i$



- ▶ For perturbativity

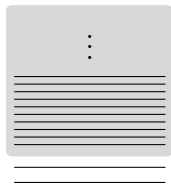
$$y \sim \frac{1}{\sqrt{N}}$$

- ▶ Production rate can still be sizable



$$\sigma_{\text{tot}} \sim \sum_{i=1}^N y^2 \sigma_h \sim (Ny^2) \sigma_h$$

out of reach



- ▶ Perhaps only the bottom of the spectrum is accessible

$$\sigma_{\text{acc}} \sim y^2 \sigma_h \sim \frac{1}{N} \sigma_h$$

- ▶ Contribution to invisible Higgs width reduced

$$\Gamma_{\text{inv}} \sim \frac{y^2}{y_b^2} \Gamma_{bb} \sim \frac{1}{N} \Gamma_{bb}$$

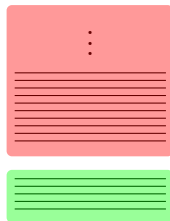
- ▶ **Heavy** states lead to cascade decays

(given couplings among new sector)

- ▶ Many low  $p_T$  final state particles
- ▶ Diffuse signal, hard for triggers

- ▶ **Light** states have short decay chains

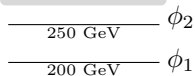
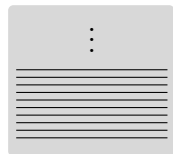
- ▶ Little visible and little missing energy
- ▶ Hard to extract from backgrounds



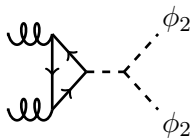
# Toy Model

- Consider a toy model with two accessible real scalars  $\phi_1$  and  $\phi_2$

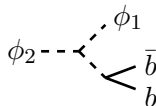
out of reach



$$\mathcal{L}_{\text{int}} = (y_1\phi_1^2 + y_{12}\phi_1\phi_2 + y_2\phi_2^2)|H|^2$$



- Decays to  $\phi_1$  and offshell Higgs



- Have  $pp \rightarrow \phi_2\phi_2 \rightarrow (\phi_1 b\bar{b})(\phi_1 b\bar{b})$   
 $\rightarrow 4b + \cancel{E}_T$

- ▶ Consider  $N$  real scalars  $\phi_i$ ,  $i = 1, \dots, N$

$$\mathcal{L} \supset \frac{1}{2} \partial_\mu \phi_i \partial^\mu \phi_i - \frac{m_{ij}^2}{2} \phi_i \phi_j - \lambda_{ijkl} \phi_i \phi_j \phi_k \phi_l$$

- ▶ Let the mass term be randomly-determined

$$m^2 = \begin{pmatrix} m_{11}^2 & m_{12}^2 & \cdots & m_{1N}^2 \\ m_{21}^2 & m_{22}^2 & \cdots & m_{2N}^2 \\ \vdots & \vdots & \ddots & \\ m_{N1}^2 & m_{N2}^2 & & m_{NN}^2 \end{pmatrix}$$

- ▶ *Example 1:*  $m_{ij}^2 = 0$  for  $i \neq j$ .

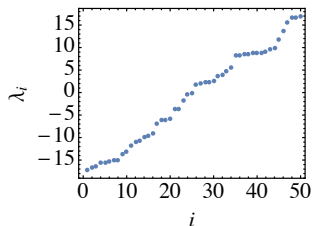
$$m^2 = \begin{pmatrix} m_1^2 & 0 & \cdots & 0 \\ 0 & m_2^2 & \cdots & 0 \\ \vdots & \vdots & \ddots & \\ 0 & 0 & & m_N^2 \end{pmatrix}$$



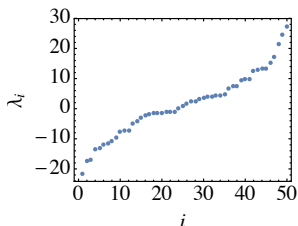
- ▶ Example 1:  $m_{ij}^2 = 0$  for  $i \neq j$ .

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- ▶ From  $Unif(\mu = 0, \sigma = 10)$



- ▶ From  $Gaus(\mu = 0, \sigma = 10)$

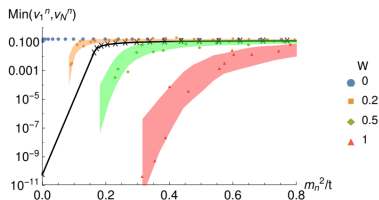


- ▶ Spectrum can be shifted by non-zero mean  $\mu$

- ▶ *Example 2: Anderson localization*
  - ▶ Nearest neighbor structure to mass matrix
  - ▶ Randomness in diagonal terms,  $\epsilon_i$

$$m^2 = \begin{pmatrix} \epsilon_0 + t & -t & 0 & \cdots & 0 \\ -t & \epsilon_1 + 2t & -t & \cdots & 0 \\ 0 & -t & \epsilon_2 + 2t & \cdots & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & \cdots & \epsilon_{N-1} + 2t & -t \\ & & & & -t & \epsilon_N + t \end{pmatrix}$$

- ▶ Mass spectrum

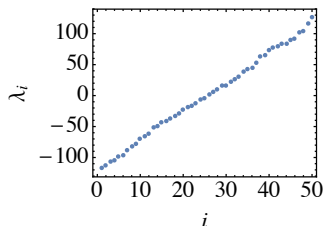


- ▶ *Example 3:* Fully-populated  $m_{ij}^2$  drawn from  $Gaus(\mu = 0, \sigma)$

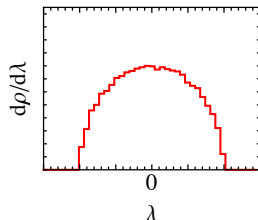
$$m^2 = \begin{pmatrix} 29.9 & 19.2 & 3.5 & -1.3 \\ 19.2 & 27.6 & -31.5 & -14.7 \\ 3.5 & -31.5 & 19.1 & 14.4 \\ -1.3 & -14.7 & 14.4 & -9.6 \end{pmatrix}$$

$$\lambda = 65.6, 31.2, -15.8, -13.9$$

- ▶ Eigenvalue distribution



- ▶ Probability density

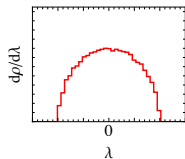


- ▶ In large  $N$  limit, distribution goes to Wigner semicircle

- ▶ Wigner semicircle distribution
  - ▶ Distributions has endpoints at  $\pm\sqrt{2N}\sigma$
  - ▶ Non-zero  $\mu$  does not shift spectrum

$$\lambda_N \sim N\mu$$

$$\lambda_i \in (-\sqrt{2N}\sigma, \sqrt{2N}\sigma) \quad i = 1, \dots, N - 1$$



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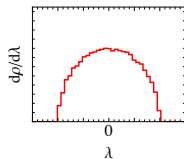
$$\lambda_i \in (-\sqrt{2N}\sigma, \sqrt{2N}\sigma) \quad i = 1, \dots, N-1$$

- ▶ Consider  $\mu \gg \sigma$ , matrix is approximately constant

$$\lambda_N \approx N\mu$$

$$\lambda_i \approx 0 \quad i = 1, \dots, N-1$$

$$\text{trace} = \sum \lambda_i = N\mu$$



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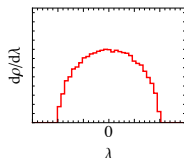
$$\lambda_N \approx N\mu$$

$$\text{trace} = \sum \lambda_i = N\mu$$

$$\lambda_i \approx 0 \quad i = 1, \dots, N-1$$

- ▶ Wigner distribution still holds for general  $\mu$  and  $\sigma$

- ▶ A model with  $\sigma$  and  $\mu$  independent could yield light states



- ▶ Possible portals (assuming no light mediators)

$$\phi|H|^2 \quad \phi^2|H|^2 \quad LH\psi$$

- ▶ We use the following model

$$\mathcal{L} \supset -\frac{1}{2}m_{ij}^2\phi_i\phi_j - \lambda_{ijkl}\phi_i\phi_j\phi_k\phi_l - \lambda_{Hij}|H|^2\phi_i\phi_j$$

- ▶ For simplicity, for all scalars, take

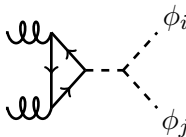
$$\lambda_{ijkl} = \frac{1}{N} \quad \lambda_{Hij} = \frac{1}{N}$$

- ▶ For simplicity, take masses uniformly distributed
- ▶ Lightest scalar,  $\phi_1$ , is stable

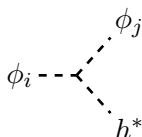
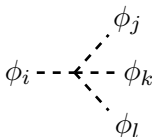
- We use the following model

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- Production from gluon fusion

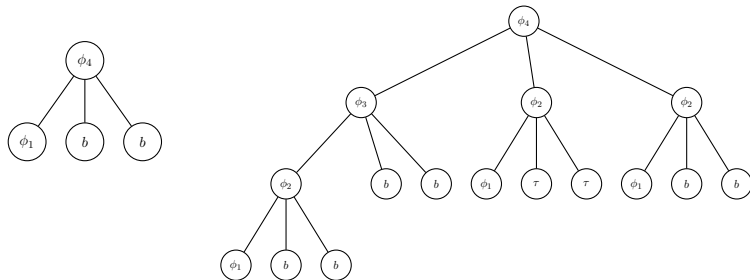


- Decays within scalar sector and via offshell Higgs

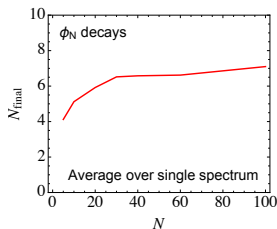




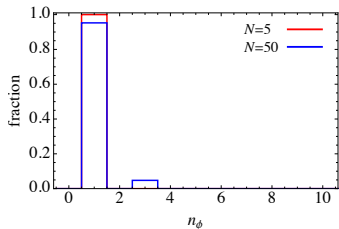
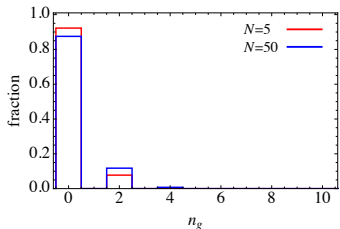
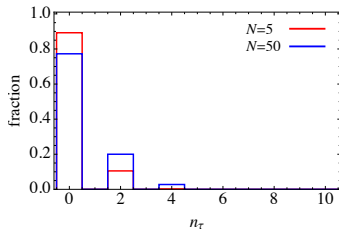
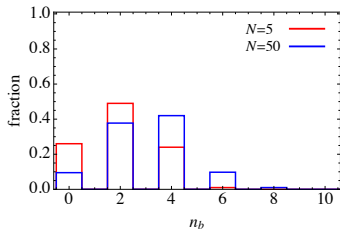
- ▶ When heavy scalars are produced, they cascade down through the scalar sector



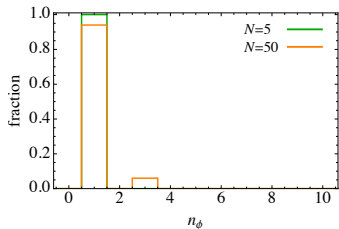
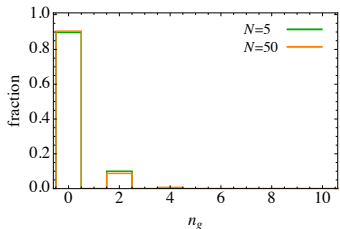
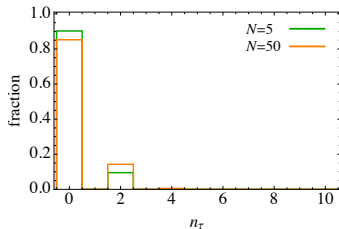
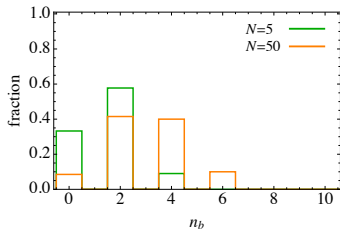
- ▶ Maximum height of tree  $\sim \log N$
- ▶ Maximum final state particles  $\sim N$
- ▶ In practice, spectrum-dependent



- ▶ Number of particles in  $\phi_N$  decay (spectrum A)



► Number of particles in  $\phi_N$  decay (spectrum B)



# Outlook

- ▶ New sectors are plausible new physics scenarios
- ▶ Have a large  $N$  number of new particles does not mean it is necessarily easy to observe
- ▶ Can disorder answer any other interesting questions in model building?

