SEXPARERGY Science Office of

Discovering or Falsifying Predictive Thermal Dark Matter (<GeV)

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Overview

1) What's **great** about thermal DM?

2) What's **different** about light thermal DM (< GeV)?

3) How can we test **all** predictive models?

[few known examples]

Requires nonstandard cosmology

Q: What's so great about equilibrium? A: Generic and easy to achieve

Compare interaction rate to Hubble expansion

$$
\mathcal{L}_{\text{eff}} = \frac{g^2}{\Lambda^2} (\bar{\chi} \gamma^\mu \chi) (\bar{f} \gamma_\mu f)
$$

$$
H \sim n \sigma v \quad \Longrightarrow \quad \frac{T^2}{m_{Pl}} \sim \frac{g^2 T^5}{\Lambda^4} \bigg|_{T=m_\chi}
$$

Equilibrium is reached in the early universe if

$$
g \gtrsim 10^{-8} \left(\frac{\Lambda}{10\,{\rm GeV}} \right)^2 \left(\frac{\rm GeV}{m_\chi} \right)^{3/2}
$$

Nearly all testable models feature equilibrium at early times

Griest et. al. 1992

Q: What's so great about equilibrium? **A: Insensitive to unknown high energy physics**

Initial condition known

Calculable and independent of inflation, reheating, baryogengesis etc.

Mass & couplings set abundance

A discovery would directly probe early universe cosmology

Only *other* **UV insensitive mechanism is "freeze-in"**

- Ad hoc initial condition $n_{\chi}(0) = 0$
- DM produced through tiny couplings, **very hard to test**

Thermal Equilibrium *Thermal Equilibrium* **A: Narrows Viable Mass Range (!)** Q: What's so great about equilibrium?

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Light DM vs. WIMPs

LDM must be SM neutral Otherwise would have been discovered at LEP

Overproduced without comparably light, neutral "mediators" LDM requires light new forces

Annihilation through renormalizable interactions Higher dimension operators have same problem as electroweak mediators

Light mediators are not optional; they're essential

Who's Heavier: DM or Mediator?

Abundance set by g_{χ} **No clear experimental target**

Mediator decays **visibly Motivates hidden force searches**

 $m_{\chi} < m_{\text{med}}$

 g_{χ} Abundance depends on g_{SM} **Predictive thermal targets**

Mediator decays **invisibly* Motivates missing energy probes**

What Kind of Mediator?

Neutrality and Renormalizability require "portal" interactions

 $\phi H^{\dagger}H \longrightarrow$ Scalar ϕ mixes with Higgs after EWSB Couples to SM masses $\epsilon \phi \frac{m_f}{m_f}$ $\epsilon \phi H^{\dagger}H$ $\frac{\partial f}{\partial y} \bar{f} f$

 F_{μ}^{\prime} $\mu\nu^{\prime}F^{\mu\nu}$ **Dark photon** *A'* **mixes with SM photon** Couples to **EM** current $\epsilon A'_\mu J_{\rm EM}^\mu$ ✏

 $V_\mu J_{\rm SM}^\mu$ **Vector V directly couples to DM & SM** Couples to different current $J^\mu_{\rm SM}$ ✏

Anomaly free options $B - L$, $L_i - L_j$, $B - 3L_i$

Vector models all similar, but also couple to neutrinos

Higgs Portal Direct-Annihilation Ruled Out!

Conclusion independent of DM candidate Similar situation for pseudo-scalar mediator etc.

GK arXiv:1512.04119

What Kind of Mediator?

Vector models all similar, but also couple to neutrinos

Classify DM by Annihilation During CMB Era

Planck Collaboration 1502.01589

 T D are exit of equilibrium kare out-of-equilibriatmanial \sim region ultimately accessible by a cosmic variance limited experiment with angular resolution comparable to that of *Planck*. $DM \sim SM$ \rightharpoonup \rightharpoon **annihilation channels. The data grey circles show the best-fitting** $\text{DM} < 10 \text{ GeV}$ *predictive Rare out-of-equilibrium annihilation ionizes H (z=1100) $\bigcap_{i=1}^n$ $\bigcap_{i=1}^n$ semi-annihilation, …) CMB photons pass through more ions (modifies peaks)

Classify DM by Annihilation During CMB Era

tiny annihilation rate at CMB

No observable indirect detection for < GeV thermal DM

Safe models require either:

Scalar or Majorana

P-wave annihilation Different DM population @ CMB Asymmetric Dirac or Pseudo-Dirac

Representative Scenario: Dark Photon Mediator *A'*(*m < m*)

$$
\mathcal{L} = -\frac{1}{4} F'_{\mu\nu} F'_{\mu\nu} + \frac{m_{A'}^2}{2} A'_{\mu} A'^{\mu} + A'_{\mu} J^{\mu}_{\chi} + \epsilon A'_{\mu} J^{\mu}_{\text{EM}}
$$

Not the only model, but qualitatively similar to viable variations Main difference for other scenarios: $J_{\text{EM}}^{\mu} \rightarrow J_{B-L}^{\mu}$, $J_{L_i-L_i}^{\mu}$ \cdots Overview

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Accelerator Searches

- Beam Dumps [DM production + detection]

-Missing Energy/Momentum [DM production only]

Fig. 2. In electron at electron and proton beam dump experiments via dark bremsstrahlung and meson decay. The resulting $\mathbf{r} = \mathbf{r} \cdot \mathbf{r}$

electrons can expand the sensitivity to sub-GeV DM that C
T $\overline{}$ \mathbf{F} Batell, Pospelov, Ritz 0903.0363 direct detection In this Letter, we report on the first dedicated search of deNiverville, Pospelov, Ritz 1107.4580 the two winds where $\frac{4041}{77040}$ and the two states of the second temperature of the second second second temperature and the second temperature of the second temperature of the second second temperature and second tem Frugiuele 1701.05464 Coloma, Dobrescu, Francescu, Francescu, 1512.03852 Batell, deNiverville, McKeen, Pospelov, Ritz 1405.7049 which capeler, the decaded to decade the detector of the detector to detector the detector to detector to detector to detector to detector to detector the detector of the detector of the detector of the detector of the det Target *e* $\chi_1\chi_2$ $e^- \longrightarrow$ $v = 0$ \sim 000 \sim 2.03852

relativistic radiativistic process (Relativistic direct detection

e

Active Target (ECAL/HCAL) *b*) ... but a we control flux!

e

Neutrino Experiments: MiniBooNE-DM Collaboration Search با
أ

Neutrino Experiments: Superior Probes of Coannihilation *Z* α ^{*a*}) *a*_{robes} d

Jordan, Kahn, GK, Moschella, Spitz 1806.05185 Izaguirre, Kahn, GK, Moschella 1703.06881

dark matter searches with beam dump experiments.

 $\mathcal{C} \supseteq \mathcal{C} = \Lambda' \bar{\mathcal{Z}} \circ \mathcal{C}^{\mu}$ $\mathcal{L} \supset \mathcal{Y} \mathcal{Y} \mathcal{Y} \mathcal{Y} \mathcal{Y} \mathcal{Y} \mathcal{Y} \mathcal{Y} \mathcal{Y} \mathcal{Y}$

Proton Fixed Target: SeaQuest @ Fermilab *Z* d Target SeaO Ł

 ν_{p} $\Omega_0 \cap \overline{\Lambda}$ in $\overline{\Lambda}$ and $\overline{\Lambda}$ and $\overline{\Lambda}$ in $\overline{\Lambda}$ Δ udev. Iu $\hphantom{00}$ – Iu $\hphantom{0}$ I $\textbf{18.3:}$ $\textbf{19.4:}$ and sensitivity interactions) to signals of $\textbf{19.4:}$ $E = 190 \text{ C} \text{V}$ $10²⁰$ to $10³⁰$ to $10²⁰$ to models in θ $L_p \sim$ 120 GeV, 10 $-$ 10 Γ Ol (0.5). Along the black contour, the abundance of ¹ matches the observed dark matter energy density. The shaded regions $U = 100 \Omega V$ 10^{18} 10^{20} $D\Omega$ $\mathcal{D}_{p}\sim$ 120 GeV , 10 $\,$ $-$ 10 $\,$ FO 1 $\,$ $E_p \sim 120 \,\, \mathrm{GeV} \,\, , \,\, 10^{18} - 10^{20} \,\, \mathrm{POT}$

Berlin, Gori, Schuster, Toro 1801.05805, 1804.00661 compared to proton beam dumps, Sea discovery potential for long-lived particles below the do, Tou4.0000 I $\mathbf{1}$ are excluded by LEP [61, 62], BaBar [63, 64], dark matter scattering at LSND [37, 65], E137 [66, 67], and MiniBooNE [25], α , Gori, Scriuster, Toro Tou Losous, Tou4.0000 for visible decays at α

BDX: Dark Photon & Leptophilic DM Mediators

Freeze out via coannihilation signatures (red dashed) for 10²² electrons on target mediated by *leptophilic* gauged *^L^e ^L^µ* (top row) and *^L^e ^L*⌧

Note Asymmetric DM models are viable anywhere above the targets Izaguirre, Kahn, GK, Moschella 1703.06881 \mathcal{L} such columns are shown: In all four plots, the blue curve represents, the blue curve represents, the blue curve represents of \mathcal{L} the parameter space is continued to parameter space $\frac{1}{2}$ and *full density* and *full densit* But still double taxation for beam dumps. How do we improve? BDX Collaboration 1712.01518

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- **-Missing Energy/Momentum [DM production only]**

B-Factory Searches: BaBar + Belle II

 $\sqrt{s} = 10.58 \text{ GeV} [\Upsilon(4\text{s})]$

Resonance search for *A'* $m_{A'}^2 = (p_\gamma - p_{e^+} - p_{e^-})^2$

BABAR: 53/fb, Belle II: 50/ab (!)

Izaguirre, GK, Schuster, Toro 1307.6554 Essig, Mardon, Papucci, Volansky Zhong 1309.5084 C. Hearty, Cosmic Visions Workshop Talk BABAR Collaboration 1702.03327

Electron Beam Missing Momentum Strategy: LDMX

Comprehensive Coverage: Dark Photon Mediator *A'*

Comprehensive Coverage: Dark Photon Mediator *A'*

Near resonance, the targets depend on A' decay width: hardest case to cover Feng & Smolinsky 1707.03835 **Berlin, Blinov GK, Schuster, Toro: 1807.01730** Berlin, Blinov GK, Schuster, Toro: 1807.01730

Comprehensive Coverage: Other Viable Mediators

FIG. 7: As in Fig. 2, thermal targets for the representative data matter can be interested at \sim III A but instead of \sim A but in stead of Sec. III A but in stead of Sec. III A but in stead of \sim A but in stead of $\$ Berlin, Blinov GK, Schuster, Toro arXiv: 1807.01730 Where are the blind spots?

So far we have covered nearly all **predictive** < GeV models *M I M <i>M I m I*

Dark photon

 $F = 1 \quad 1 \quad 1 \quad 1$ **Dark photon Anomaly free U(1)** R_I R_2 R_3 _{θ} R_4 hilation cross section is independent of the *A*⁰ coupling to visible matter. In the direct annihilation regime, *^e[±] ^Z*⁰ *B-L, B-3Le …* etc.

What about mediators w/ mainly 2nd & 3rd generation couplings? *<u>∤</u>* $\frac{1}{2}$ *, 1, 0,*

A

discovery or falsification (see Fig. 5).

Only one anomaly free U(1) group

$$
\begin{array}{ccc}\n\text{we} & \text{U(1) group} & x \\
\text{and} & \text{U(1)} & \text{U(1)} \\
x & \text{U(1)} & \text{U(1)} \\
x & \text{U(1)} & \text{U(1)}\n\end{array}
$$

A

^e[±] ^e[±]

- \sim 100-200 GeV muon beam \sim 10s meter baseline \sim \sim 10 $\sim 10^{11} - 10^{12}$ u $\frac{1}{\sqrt{2\pi}}$ $\frac{1}{\sqrt{2\pi}}$ $\mathbf{1000000}$ \sim 100.700 is carried and state the scattered muon, which restricts the total energy is transmitted by the \sim \mathcal{L} , the veto counters V1 and V2, and V2, and \mathcal{L} decay penetrates them without interactions resulting in a zero-energy signature in the dashed line represents r \sim institution muon is measured muon is measured muon is momentum of the second muon is measured muo ~ 10s meter baseline $\sim 10^{-7} - 10^{-7}$ μ $r\sim 10^{11}-10^{12}$ μ radia
- $\overline{\text{M}}$ 1) Measure E in/out $\sum_{i=1}^n \sum_{i=1}^n \sum_{i=1}^n \sum_{j=1}^n \sum_{j$ z) ingger on nussing energy $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ \sim previously measured parameter space, it is expected parameter space, it is expected parameter space, it is expected by 2) Trigger on missing energy and the state
- \sim interactions in the momentum of the mome the income is vectoral support in \mathcal{S}) veto additional SM activity α is in the range α of the range α 0.1 α for the massess α 0.1 α for the massess α $\sum I(t)$ and $\sum I(t)$ 10−6 $\sum I(t)$ the veto additional p \mathcal{O}_f roce characterization of 3) Veto additional SM activity 10

Gninenko, Krasnekov, Mateev, arXiv: 1412. function of the Z^µ mass. Using the relation n90% Shinenko, Krasnekov, Mateey $\mathcal{L}_{\mathcal{L}}$ of the first (last \sim 1412, 1400 $\,$ ₩ 100 × 100 mm2 (⇔ 400 × 40 Gninenko, Krasnekov, Mateev, arXiv: 1412.1400

M^3 Muon Missing Momentum

Kahn, GK, Tran, Whitbeck 1804.03144

Covers Predictive Muon-Philic Models

 $Gauped$ $L_u - L_{\tau}$ Interaction χ Figure $L_{\mu} = L_{\tau}$ Interaction χ_{Δ} Gauged $L_{\mu}-L_{\tau}$ Interaction

Also resolve muon g-2 with light physics $\left\{\begin{array}{ccc} \sim & \sim & \sim & \sim & \sim \sim \end{array}\right\}$ $\begin{array}{c|c}\n\hline\n\end{array}$. Also plotted are constraints from the neutrino trident process from the \sim $\text{Conjugative parameter space for freedom}$ (see Appendix A). Also plotted are constraints from the neutrino trident process from the CCFR experiment process from the Compatible parameter space for freeze-out X

 ρ_{max} or $-\mu$ - such that the smaller such that the that the theorem is the theorem over μ $\bigwedge_{\Lambda} \bigwedge^{\mathbb{Z}} \bigwedge^{\mathbb{Z}} \mathbb{L}, \mathcal{T}, \mathcal{V}_{\mathbb{U}}, \mathcal{L}$ $\rho_{\text{max}} - \mu$ or smaller such that the that the that the that the theorem is the theor Z^{\prime} ⌫¯*µ,*⌧ α ,

NB: annihilation to neutrinos also CMB safe *s*-wave, so this process is ruled out by CMB energy injection bounds for *m > m^µ* [52]. *s* (*b*, mathematically is reduced one can be seen NB: annihilation to neutrinos also CMB safe Summary

 \bf{A} **Modest Proposal** $\Gamma(DM \leftrightarrow SM) > H$ **Thermodynamics Set Initial Condition** Rate beats Hubble expansion at *some* point [easy to realize] **Predicts Min. Annihilation Rate** $\sigma v \gtrsim 10^{-26} \text{cm}^3 \text{s}^{-1}$ $n_{\rm DM} \sim T^3$ Insensitive to unknown high scales [inflation, baryogenesis…] Equilibrium overproduces DM, must deplete with non-gravitational force **Viable Window In Our Neighborhood** $MeV \sim m_e$ GeV $\sim m_p$ **"WIMPs"** $m_{Z,h}$ **LDM** BBN ΔN_{eff} **but if the LDM** $\alpha_{\chi} > \Omega_{\text{DM}}$ Coincidentally in broad vicinity of the electroweak scale ~ 10 s TeV

Thanks!