



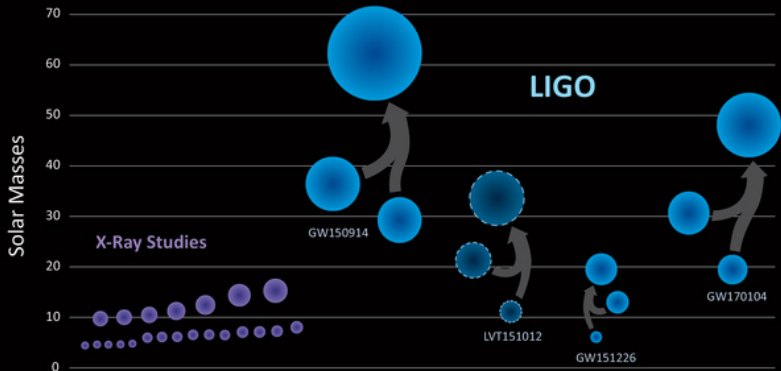
# Primordial Black Holes in the wake of LIGO

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# Black Holes of Known Mass



BH Name

GW170814

GW170608

GW170104

GW151226

LVT151012

GW150914

GS 2023+338

GS 2000+251

Cyg X-1

GRS 1915+105

V4641 Sgr

H1705-250

GROJ 1655-40

XTEJ 1550-564

4U 1543-47

GS 1354-64

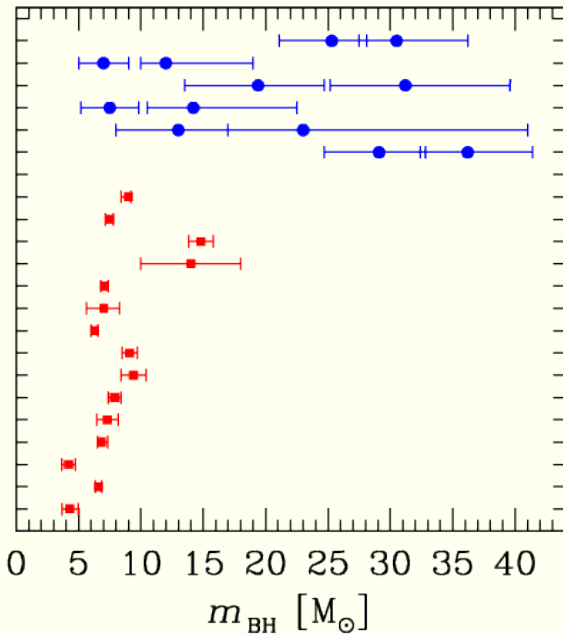
GS 1124-683

XTE J 1118+480

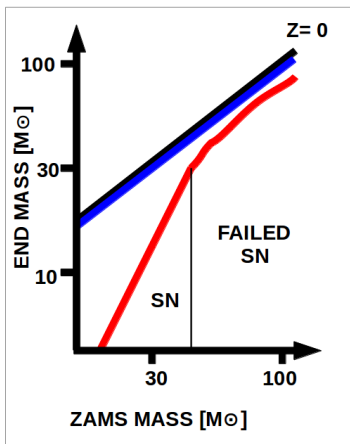
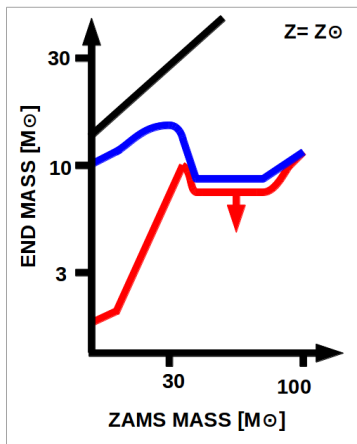
GRS 1009-45

A 0620-003

GRO J 0422+32

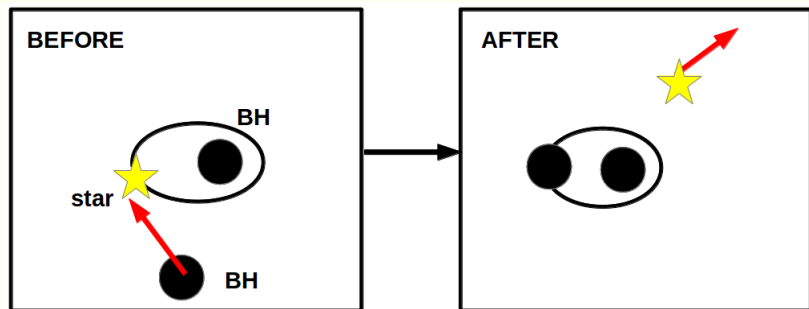


## Unexpected/surprising?



Most astrophysical models did not predict BHs with  $M \gtrsim 20M_{\odot}$ .  
But, large BHs masses can be generated from  $\geq 40M_{\odot}$   
metal-free stars undergoing direct collapse.

## How are binaries formed?



Could work in young star clusters or in nuclear star clusters surrounding SMBHs. Unlike isolated binaries, spins are misaligned/isotropic. But, three body encounters (necessary to harden the binary) can eject the system.

The astrophysical picture is largely incomplete:

- ▶ The formation channels of merging BH binaries are still uncertain. Major simplifications are adopted in dynamical simulations, and the statistics about BHs in young star clusters is small.
- ▶ A global picture of the BH merger history as a function of redshift is missing.

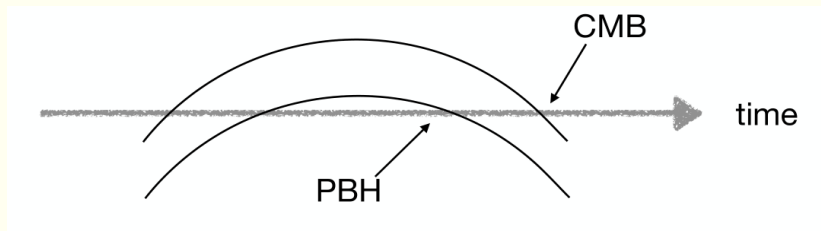
The LIGO/Virgo horizon is  $z \sim 0.1 - 0.2$ , but third-generation ground-based GW detectors (e.g. Einstein Telescope) will be able to observe binary mergers up to  $z \sim 10$ .

## Another (more massive) puzzle



SMBHs reaching  $\gtrsim 10^{10} M_{\odot}$  are present in the centers of most massive galaxies, even at large redshifts.

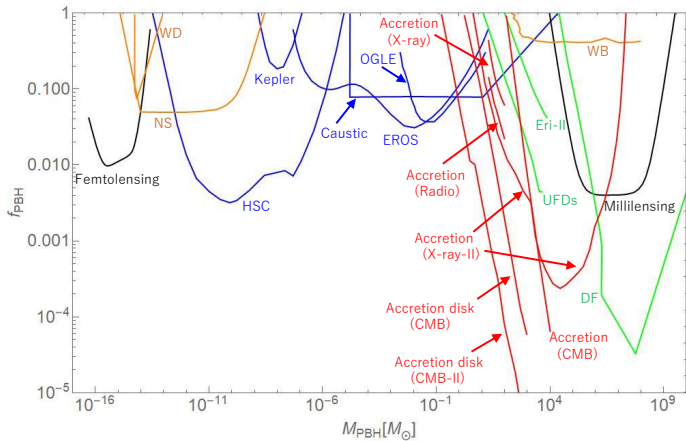
## Could they be primordial?



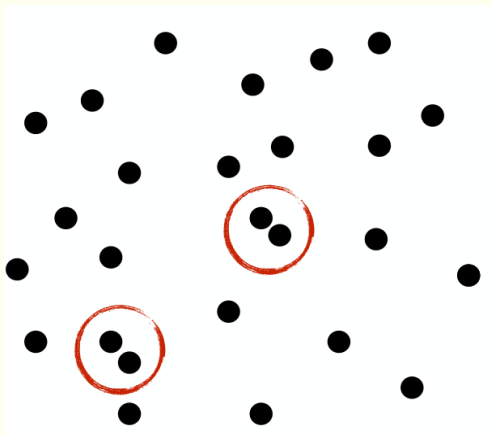
Rare Hubble scale perturbations can collapse into BHs:

$$\beta \approx \text{erfc} \left( \frac{\delta_c}{\sqrt{2}\sigma} \right)$$





## Binary formation



PBHs are randomly distributed, but some pairs are close enough to decouple from Hubble flow.

Most of the BH pairs that merge today form in the early universe, deep in the radiation era. Pairs form due to the chance proximity of PHB pairs and merge on a time-scale:

$$t_{merge} = \frac{3c^5}{170G_N^3} \frac{a^4(1 - e^2)^{7/2}}{M_{pbh}^3}$$

Several processes (torques due to other BHs, encounters with other BHs, DM spikes around PBHs, ...) influence the merger rate that is measured by LIGO.

Clustering might substantially change the picture.

Ali-Haïmoud, Kovetz & Kamionkowski, 1709.06576

Kavanagh, Gaggero & Bertone, 1805.09034

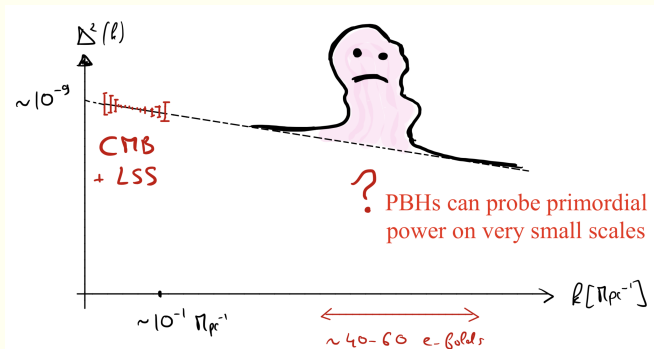
# Pair formation in present day halos

Binary BHs can also form in present day halos from GW emission. These binaries are very tight and highly eccentric so that they coalesce within a very short timescale. In principle this population gives a subdominant contribution to the LIGO observed events, but:

- ▶ PBHs could be clumped around SMBH spikes
- ▶ Merger rates could be boosted
- ▶ The cross-section is strongly velocity dependent,

$$\sigma \propto v_{\text{rel}}^{-18/7}$$

# Relics from inflation?



Picture by Ali-Haimoud

The power spectrum should be enhanced by a factor of  $10^3$  on scales  $\ll$  CMB!

# Alternative mechanisms?

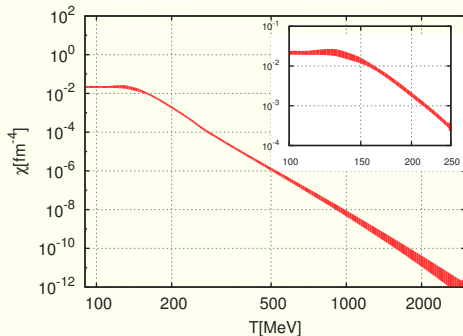
Phase transitions in the early universe provide a potential avenue:

- ▶ Several violent phenomena naturally occur that can assist in generating large overdensities that gravitationally collapse into BHs: bubble collisions, topological defects, ...
- ▶ We will consider axionic string-wall networks.

At finite  $T$  the axion field potential can be modelled by

$$V(a, T) = \chi(T) \left[ 1 - \cos \left( N_{\text{DW}} \frac{a}{v_{\text{PQ}}} \right) \right],$$

where  $v_{\text{PQ}} = f_a N_{\text{DW}}$  is the PQ breaking scale, and  $\chi(T)$  determines the axion mass  $m_a(T)^2 = \chi(T)/f_a^2 \propto T^{-8.16}$ .



# Cosmological evolution

Important distinction whether PQ symmetry is broken before or after inflation:

- ▶ Pre-inflationary PQ breaking  $\rightarrow$  the axion has a single uniform initial value  $a_i$  within the observable universe.
- ▶ In the post-inflationary case the axion takes different values in different regions.

In the latter case when the axion gets its mass, around the QCD phase transition, a hybrid string-domain wall network is formed. Eventually, the network has to decay. Otherwise, the energy density would be quickly dominated by domain walls.

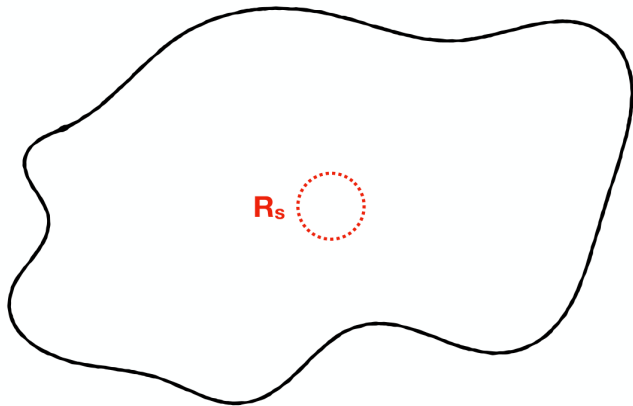


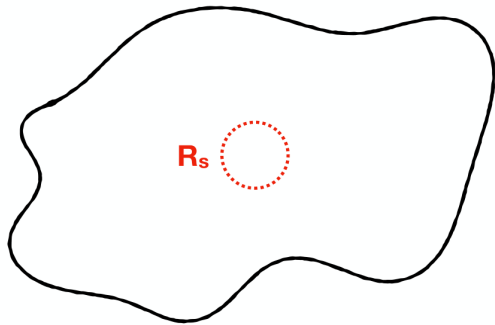
The collapse of closed domain walls, which belong to the hybrid string-wall network can lead to the formation of PBHs.

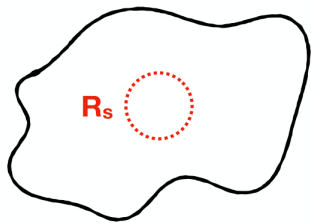
T. Vachaspati, 1706.03868

- ▶ This mechanism does not rely on (nor complicate) the physics of inflation.
- ▶ GW astronomy can potentially probe the physics of axions.

It is crucial that the annihilation of the network proceeds slowly.









# PBHs from string-wall defects

A closed DW of size  $R_*$  will rapidly shrink because of its own tension, once  $R_* \sim H^{-1} \approx g_{\text{eff}}(T_*)^{-1/2} M_p / T_*^2$ .

Its mass has contributions from the wall tension and from any difference in energy density between the two regions separated by the DW:

$$M_* = 4\pi\sigma R_*^2 + \frac{4}{3}\pi\Delta\rho R_*^3 \approx 4\pi\sigma H_*^{-2} + \frac{4}{3}\pi\Delta\rho H_*^{-3}$$

$\Rightarrow$  Heavier black holes form from DW which collapse later in cosmological history.

# PBHs from string-wall defects

The Schwarzschild radius of the collapsing defect is  $R_{S,*} = 2G_N M_*$ , and the *figure of merit* for PBH formation is:

$$p \equiv R_{S,*}/R_* \sim \frac{\sigma H_*^{-1}}{M_p^2} + \frac{\Delta\rho H_*^{-2}}{3M_p^2}$$

⇒ As the temperature decreases it becomes more likely to form a black hole.

Two regimes:

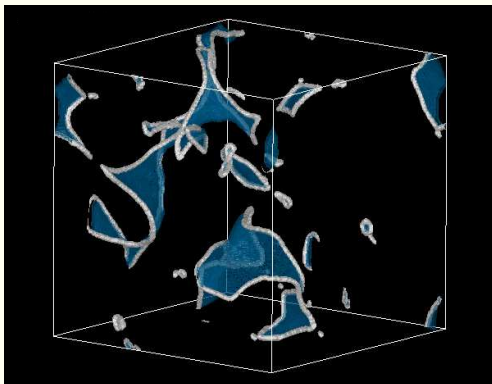
- ▶ When the tension dominates,  $M_* \sim T_*^{-4}$  and  $\rho \sim T^{-2}$ .
- ▶ When the energy density dominates,  $M_* \sim T_*^{-6}$  and  $\rho \sim T^{-4}$ .

(Deviations from spherical symmetry, radiation friction during collapse can partly modify this picture.)



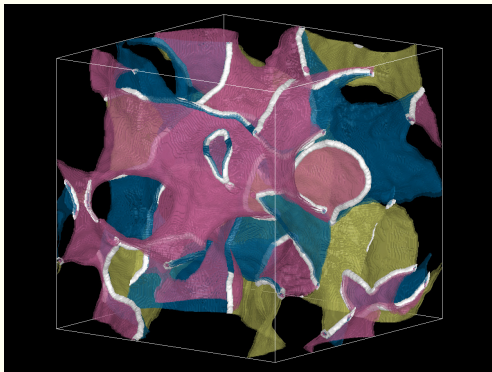
$$N_{\text{DW}} = 1$$

Only one domain wall is attached to each string. Such topological configurations quickly annihilate leaving behind a population of barely relativistic axions.

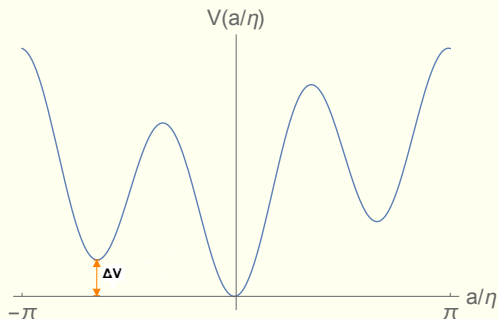


$$N_{\text{DW}} > 1$$

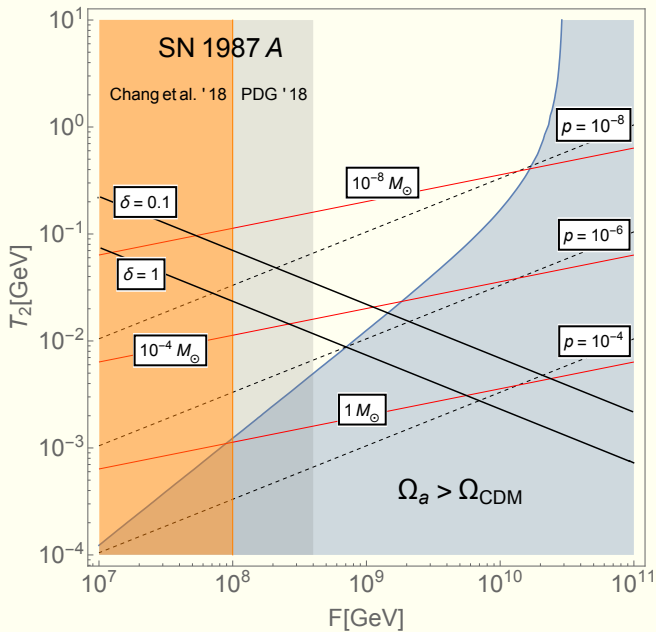
There are  $N_{\text{DW}}$  domain walls attached to every string, each one pulling in a different direction. The network can actually be stable, and dominate the universe.



Lift the degeneracy of axionic vacua by introducing a bias term (dark QCD?). The energy difference between the different minima acts as a pressure force on the corresponding domain walls.



The annihilation occurs at  $T_2 > T_*$  set by  $\frac{\Delta V}{\sigma}$ . There can be a significant separation between formation  $T_1$  and  $T_2$ .



# Axion-QCD vs ALPs

- ▶ For the QCD axion we find an interesting region around

$$f_a \sim 10^9 \text{ GeV.}$$

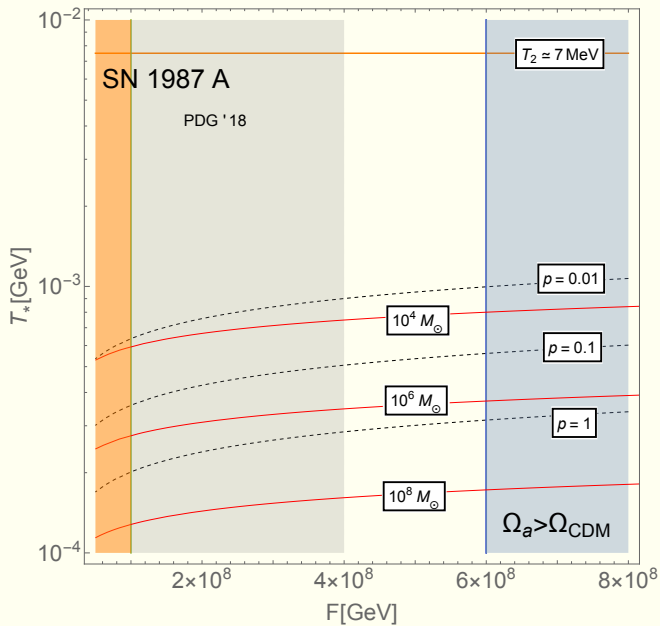
PBHs of mass  $10^{-4} M_\odot$  can form with  $p \sim 10^{-6}$ .

- ▶ For generic ALPs we can reach larger probabilities  $p \sim 10^{-3}$  in scenarios where

$$T_2 \sim \text{keV.}$$

Interestingly much larger BHs,  $\lesssim 10^8 M_\odot$  could be formed.

# Late collapses



# Late collapses

Most of the axionic string-wall network disappears at  $T_2$ , which is when the vacuum contribution starts dominating, and both  $\rho$  and  $M_*$  increase steeply.

But, 1 – 10% of the walls survive until  $\sim 0.1 T_2$ , when:

- ▶  $\rho \sim 1$

- ▶  $M_* \sim 10^6 M_\odot$

⇒ A fraction  $f \sim 10^{-6}$  of the DM end up forming SMBHs!

We have not said much about the bias term . . .

Planck suppressed operators are unlikely.

A dark gauge sector with  $\Lambda_B \sim MeV$  is an interesting possibility.



# Conclusions

- ▶ LIGO has confirmed the existence of BH binaries that are able to merge within a Hubble time.
- ▶ The observed BHs mass  $\gtrsim 20M_{\odot}$  is surprising from the astrophysics point of view.
- ▶ Alternatively, BHs could be primordial. A fraction, but not all, of the DM could be made of black holes.
- ▶ Axionic topological defects with  $N_{\text{DW}} > 1$  lead to a new Network Annihilation epoch that can potentially generate PBHs of up to  $10^6 M_{\odot}$ .