

Small Rock Seminar

Study of jet fragmentation and inclusive jet production in heavy-ion collisions with the ATLAS experiment



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9.4.2018

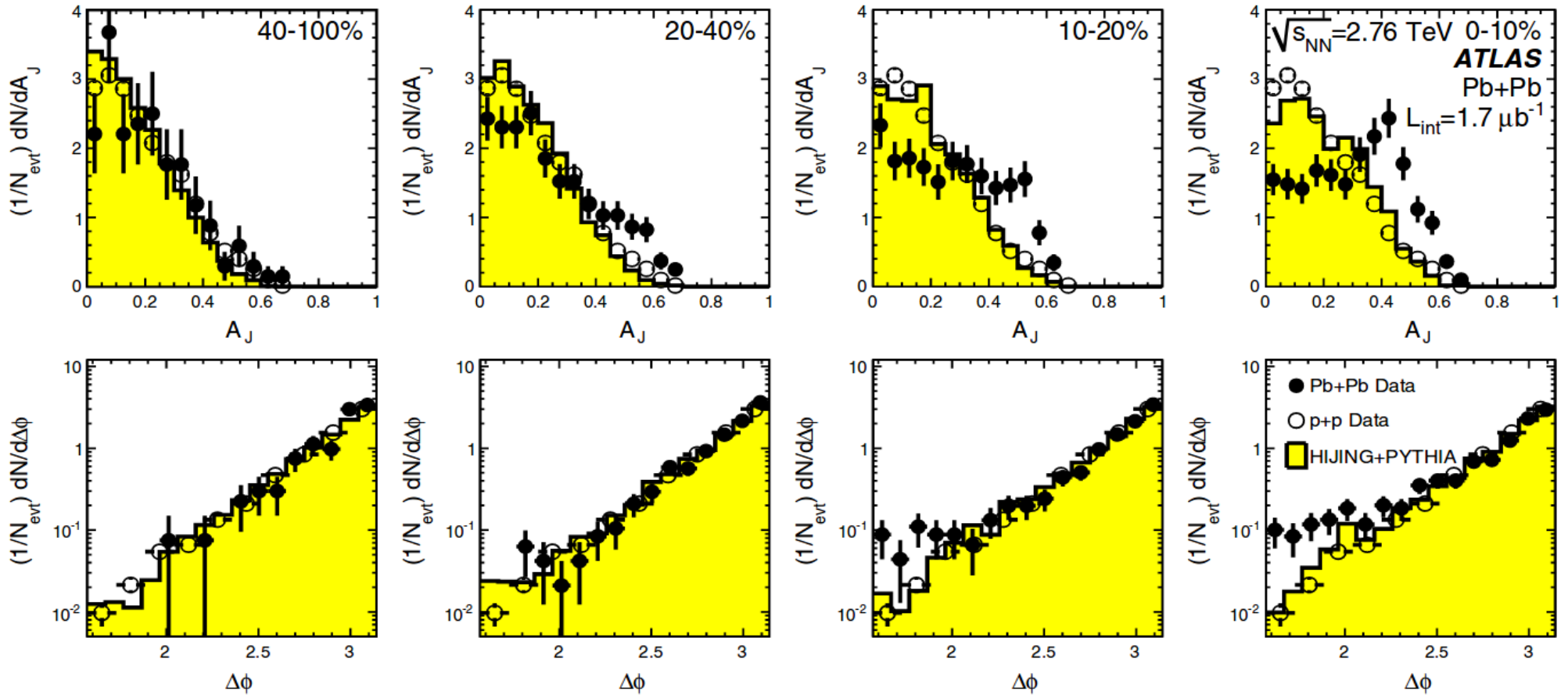


Jets in pp , $p+Pb$, $Pb+Pb$ collisions



- In heavy-ion collisions the strongly interacting deconfined matter (Quark-Gluon plasma - QGP) is created
- The main motivation is to understand better the modification of jet properties: a parton with high momentum is traversing the hot and dense medium and loses its energy by radiating gluons = **jet quenching**
- To fully understand jet quenching phenomenon, in addition to jet suppression it is necessary to measure jet fragmentation i.e. possible modifications of parton showers through interactions in the plasma
- We obtained the final result on the jet Fragmentation Functions (FF) for **Pb+Pb @ 2.76 TeV** and **p+Pb @ 5.02 TeV** with new pp @ 5.02 TeV reference and first preliminary result for **Pb+Pb @ 5.02 TeV** with pp reference at corresponding energy
- **Pb+Pb @ 5.02 TeV** was done with large p_T range
- We obtained preliminary R_{AA} @ 5.02 TeV for jets

Discovery of jet quenching on LHC



PRL 105 (2010) 252303 (>500 citations)

$$A_J = \frac{E_{T1} - E_{T2}}{E_{T1} + E_{T2}}$$

Used data and MC



- **Pb+Pb @ 2.76 TeV** (2011), $L_{\text{int}} = 0.15 \text{ nb}^{-1}$
 - **pp @ 2.76 TeV** (2013), $L_{\text{int}} = 4 \text{ pb}^{-1}$
 - Pb+Pb data are compared to MC, where MC PYTHIA 6 di-jet events were embedded into real MB Pb+Pb events
- **Pb+Pb @ 5.02 TeV** (2015), $L_{\text{int}} = 0.49 \text{ nb}^{-1}$
 - **pp @ 5.02 TeV** (2015), $L_{\text{int}} = 25 \text{ pb}^{-1}$
 - Pb+Pb utilize POWHEG+PYTHIA8 with overlaid MinBias Pb+Pb collisions
 - pp utilize PYTHIA8 with the A14 ATLAS tune and the NNPDF23LO pdfs
- **p+Pb @ 5.02 TeV** (2013), $L_{\text{int}} = 28 \text{ nb}^{-1}$
 - **pp @ 5.02 TeV** (2015), $L_{\text{int}} = 25 \text{ pb}^{-1}$
 - p+Pb and pp measurements utilize PYTHIA 6 (embedded into real MB p+Pb), PYTHIA 8 and HERWIG++ samples

Jet fragmentation in Pb+Pb

@ 2.76 TeV and @ 5.02 TeV, p+Pb @ 5.02 TeV

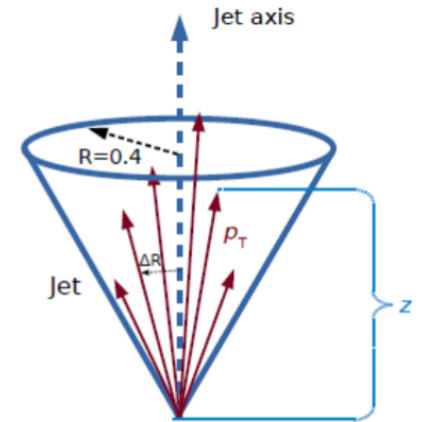
arXiv:1702.00674

ATLAS-CONF-2017-005

ATLAS-CONF-2017-004

- Jet fragmentation functions (FF) are defined as:

$$D(p_T) = \frac{1}{N_{\text{jet}}} \frac{dN_{\text{ch}}}{dp_T^{\text{ch}}} \quad D(z) = \frac{1}{N_{\text{jet}}} \frac{dN_{\text{ch}}}{dz} \quad z = \frac{p_T}{p_T^{\text{jet}}} \cos \Delta R$$

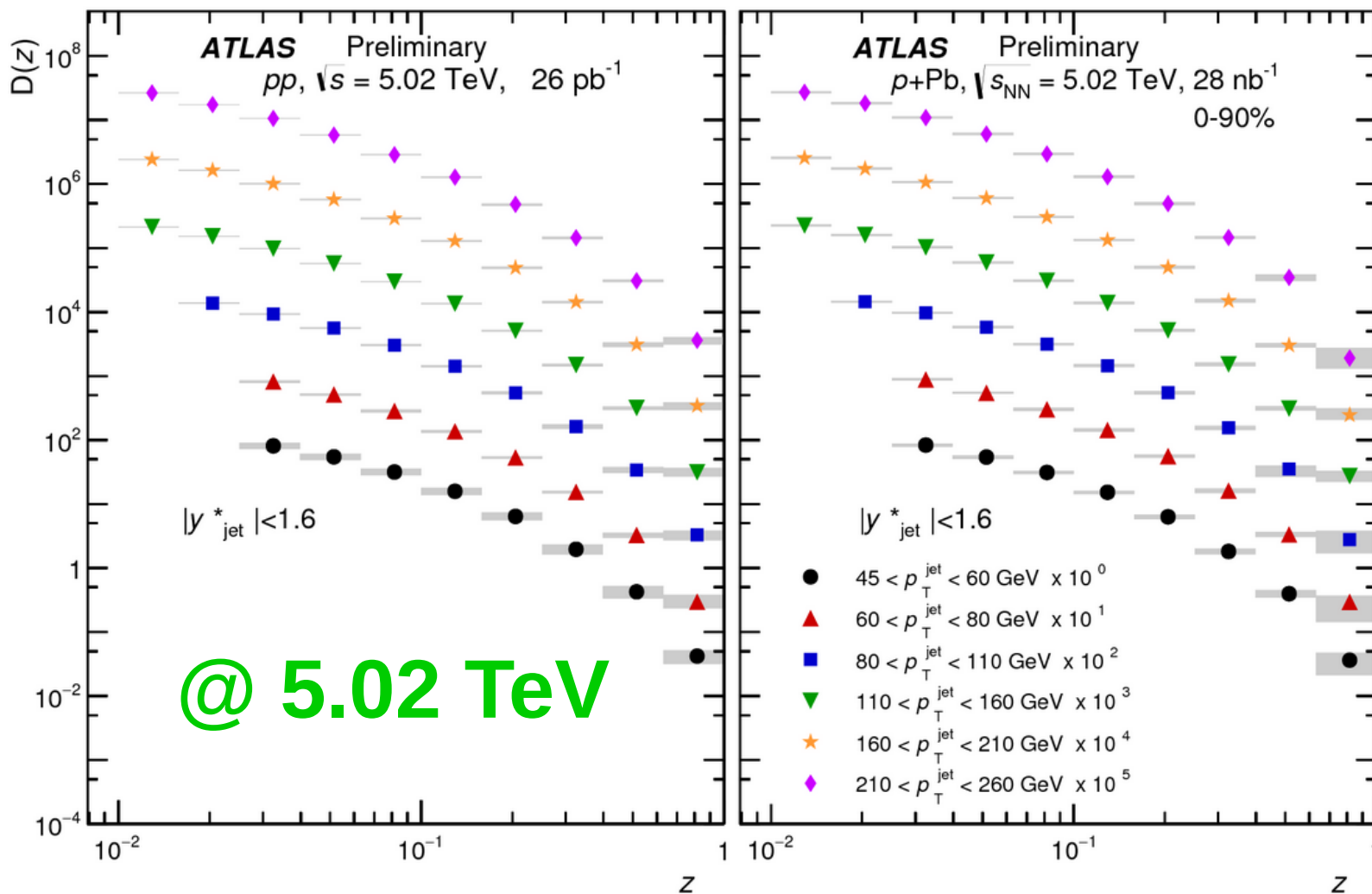


- N_{ch} is the number of charged particles associated to a jet
- Measurement was done for $R = 0.4$ jets differentially in y and p_T for Pb+Pb and differentially in p_T for p+Pb
- Jets measured using charged tracks starting at $p_T = 1$ GeV for p+Pb @ 5.02 TeV and Pb+Pb @ 2.76 TeV and $p_T = 4$ GeV for Pb+Pb @ 5.02 TeV
- FF are background subtracted, corrected for tracking efficiency and fully unfolded with 2D Bayesian unfolding
- The jet spectra for the normalisation were unfolded using 1D Bayesian unfolding for both Pb+Pb FF and using bin-by-bin correction for p+Pb

Unfolded fragmentation functions pp and $p+Pb$



ATLAS-CONF-2017-004



Spectra have similar shape, ratios are needed to see the modifications:

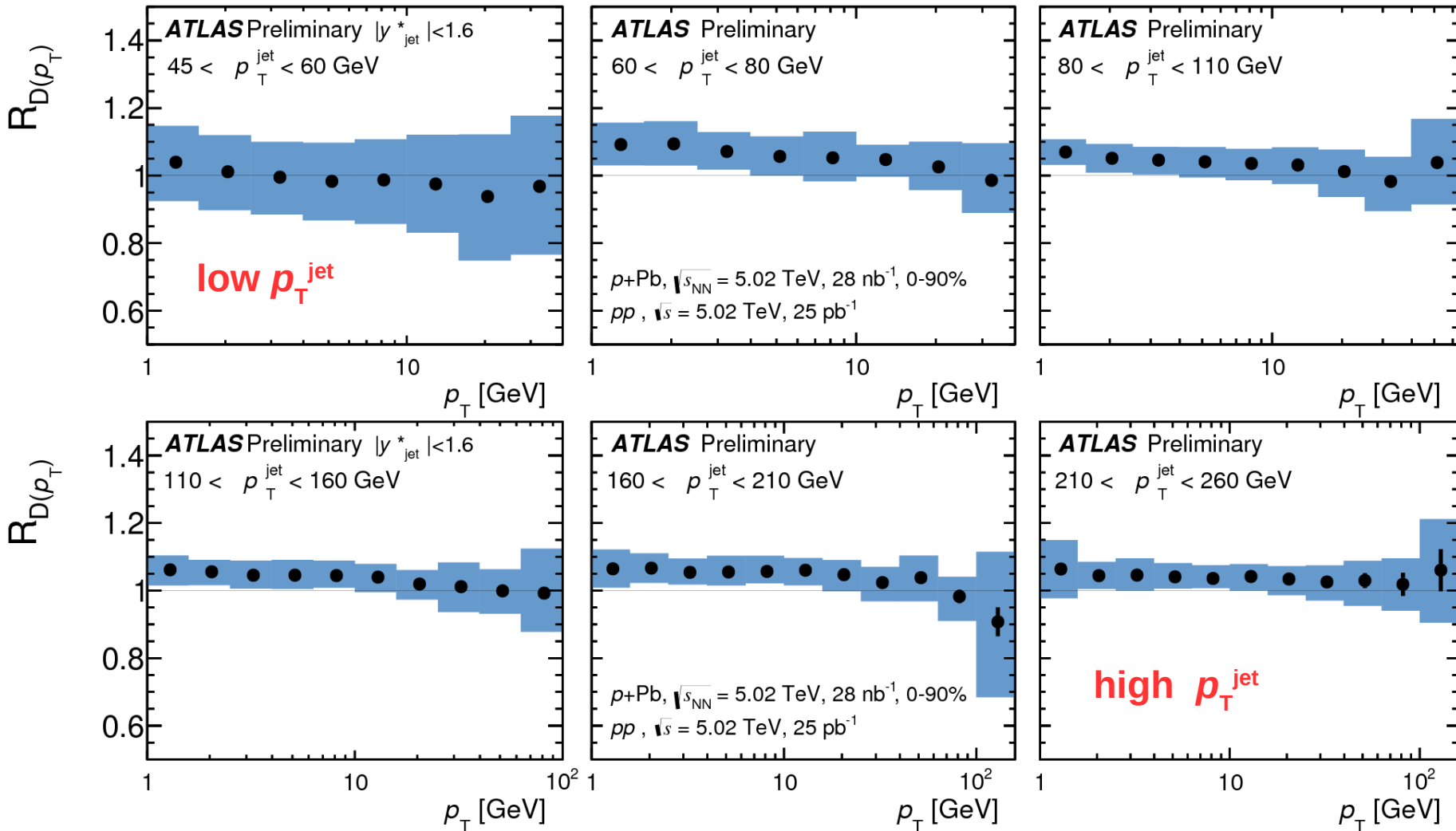
$$R_D(z) \equiv \frac{D(z)_{pPb}}{D(z)_{pp}}$$

$R_{D(p_T)}$ for different p_T bins

ATLAS-CONF-2017-004

$p+Pb$ @ 5.02 TeV with new **pp** reference @ 5.02 TeV

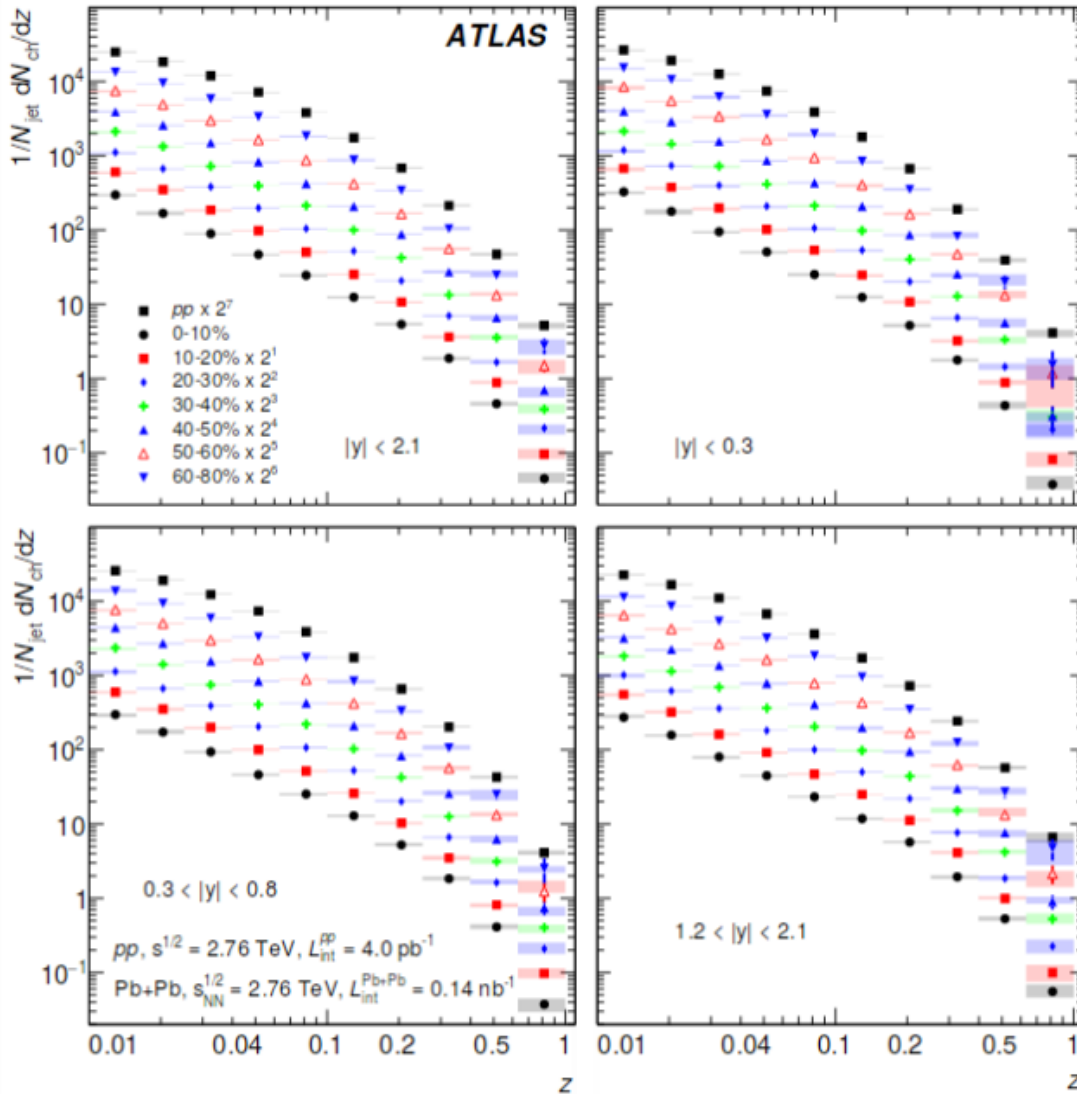
- No modification of the fragmentation functions is observed



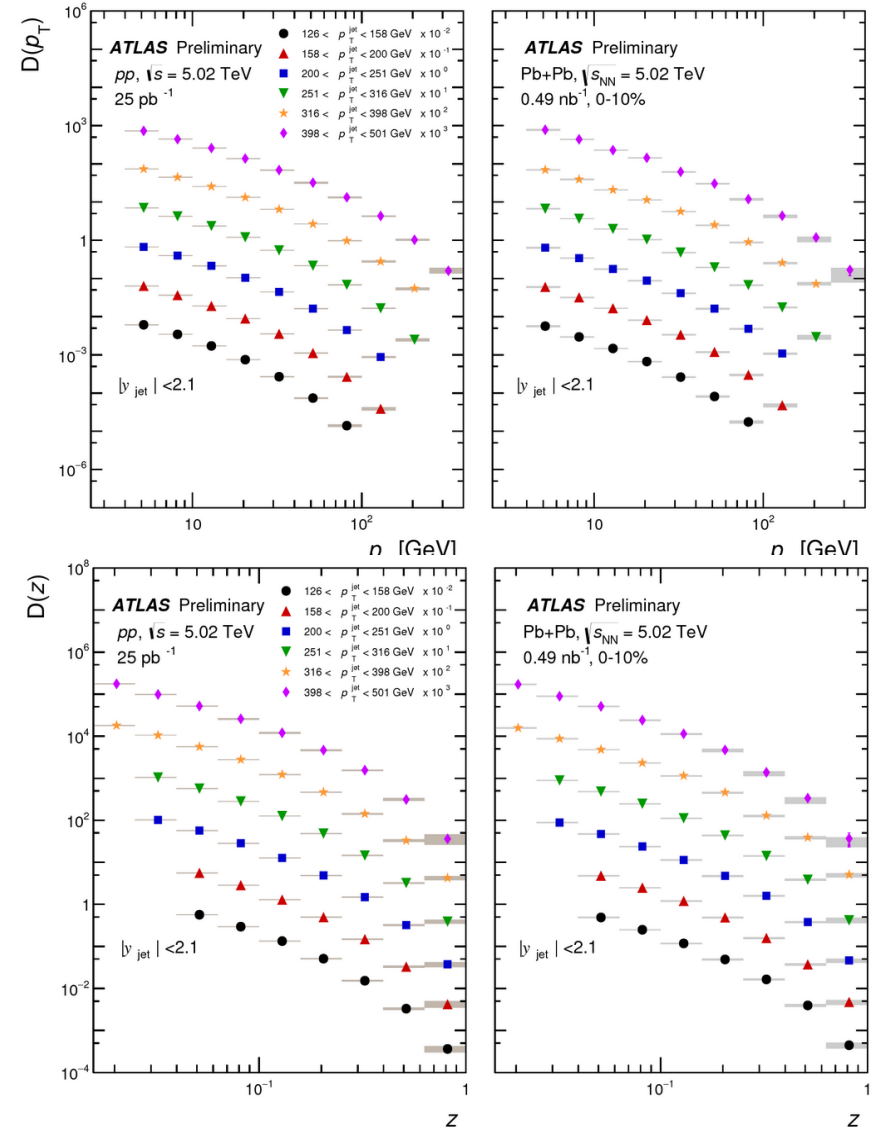
Comprehensive Pb+Pb fragmentation functions

Pb+Pb and pp @ 2.76 TeV

Pb+Pb and pp @ 5.02 TeV



arXiv:1702.00674



ATLAS-CONF-2017-005

Spectra have similar shape, ratios are needed to see the modifications:

$$R_{D(z)} \equiv \frac{D(z)_{\text{PbPb}}}{D(z)_{\text{pp}}}$$

Jet fragmentation ratios Pb+Pb

Ratios of $D(z)$ for 4 centralities

$$D(z) \equiv \frac{1}{N_{\text{jet}}} \frac{dN_{\text{ch}}}{dz}$$

@ 2.76 TeV

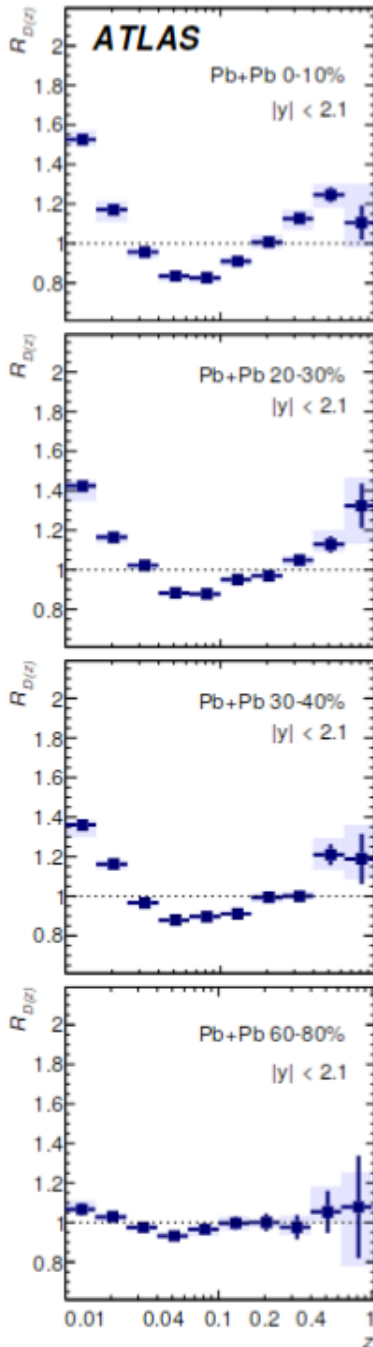
arXiv:1702.00674

$$z \equiv \frac{p_{\text{T}}}{p_{\text{T}}^{\text{jet}}} \cos \Delta R$$

Centrality dependence

- Enhancement at low z
- Suppression at intermediate z
- Enhancement at high z

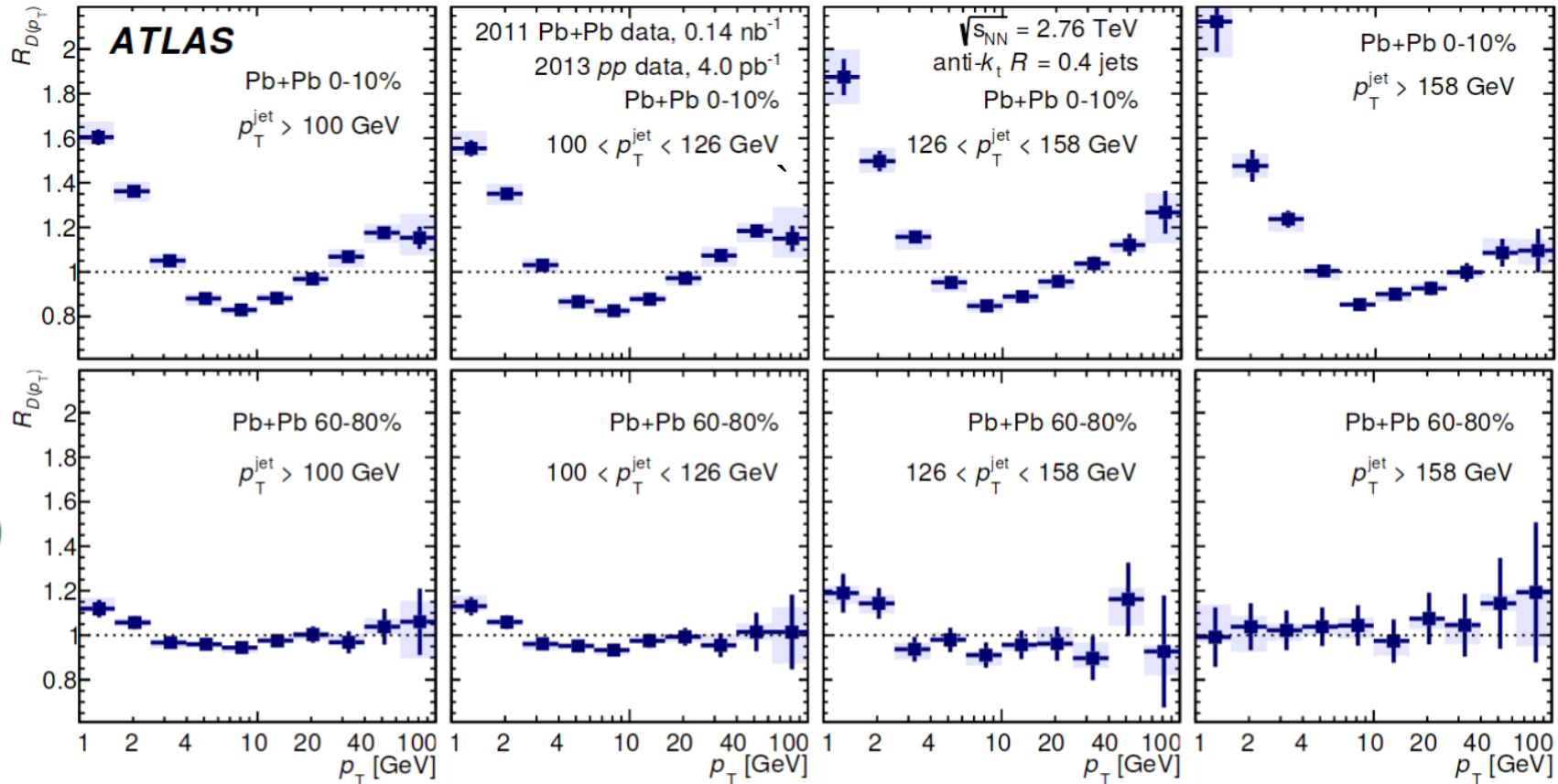
Similar observation for $D(p_{\text{T}})$



Jet fragmentation ratios Pb+Pb

Ratios of $D(p_T)$ for 2 centralities and 4 p_T bins

@ 2.76 TeV



Jet p_T dependence

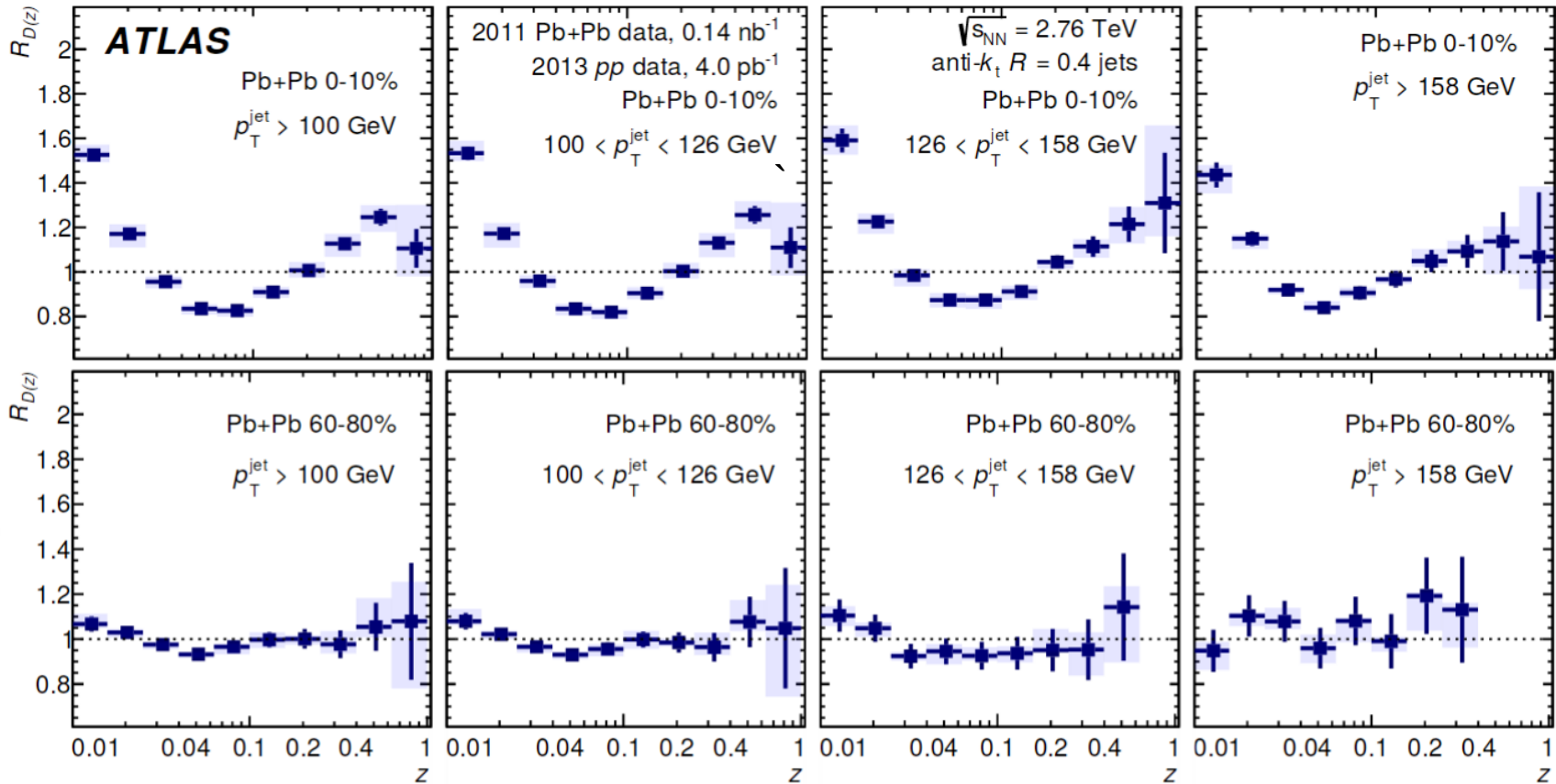
- No significant dependence on jet p_T
- Enhancement at low p_T is larger for bigger p_T^{jet} interval in the central events

$100 < p_T^{\text{jet}} < 398$ GeV
 $|y| < 2.1$

Jet fragmentation ratios Pb+Pb

Ratios of $D(z)$ for 2 centralities and 4 p_T bins

@ 2.76 TeV



Jet p_T dependence

- No significant dependence on jet p_T

$100 < p_T^{\text{jet}} < 398$ GeV
 $|y| < 2.1$

Jet fragmentation – flow of particles

Pb+Pb @ 2.76 TeV

arXiv:1702.00674

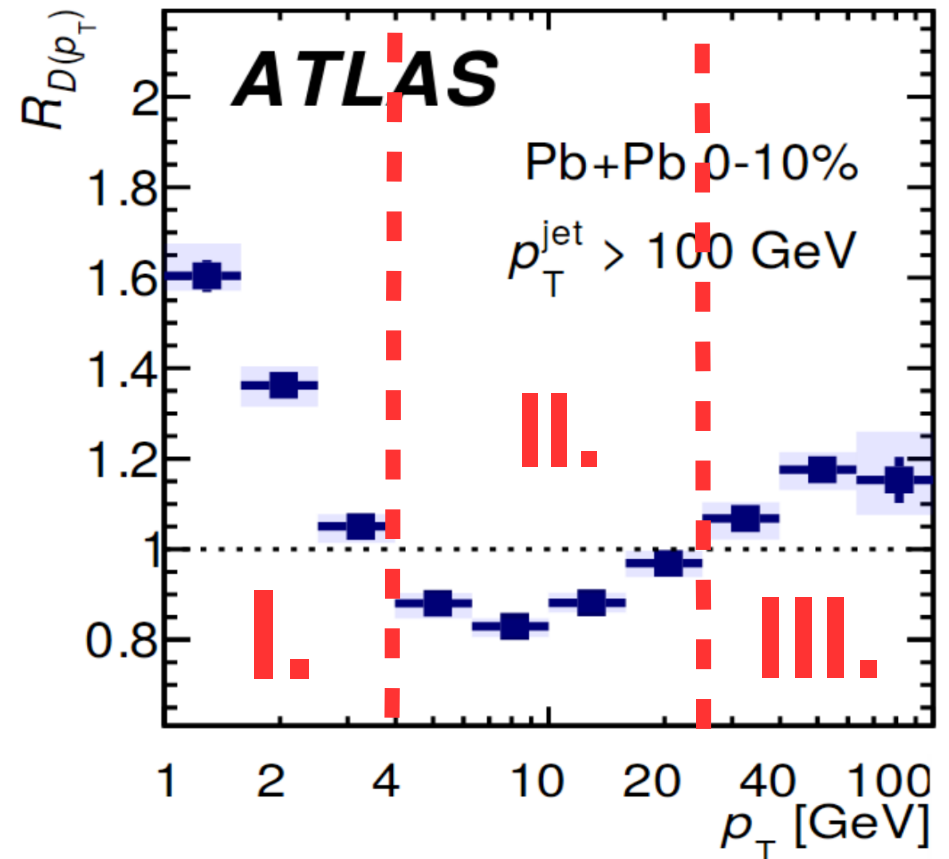
- To quantify the flow of particles as a function of N_{part} :

$$N^{\text{ch}} \equiv \int_{p_{T,\text{min}}}^{p_{T,\text{max}}} \left(D(p_T)|_{\text{cent}} - D(p_T)|_{\text{pp}} \right) dp_T$$

- Values $p_{T,\text{min}}$ and $p_{T,\text{max}}$ set accordingly where $R_{D(p_T)}$ is above or under unity

- Three p_T intervals:

- I. 1 – 4 GeV
- II. 4 – 25 GeV
- III. 25 – 100 GeV



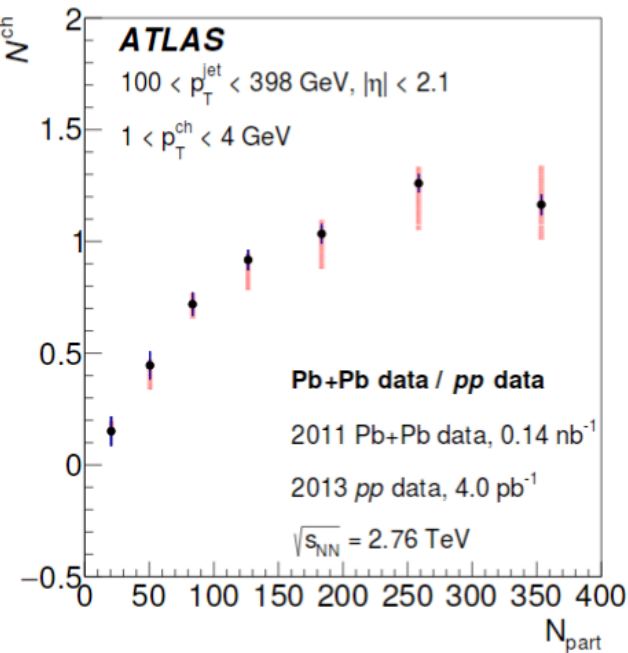
Jet fragmentation – flow of particles

Pb+Pb @ 2.76 TeV

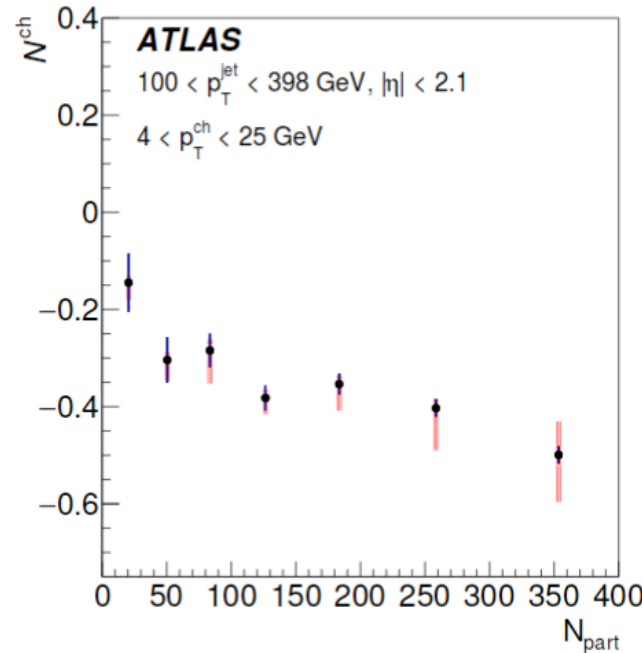
arXiv:1702.00674

- To quantify the flow of particles as a function of N_{part} :
$$N^{\text{ch}} \equiv \int_{p_{T,\text{min}}}^{p_{T,\text{max}}} \left(D(p_T)|_{\text{cent}} - D(p_T)|_{\text{pp}} \right) dp_T$$

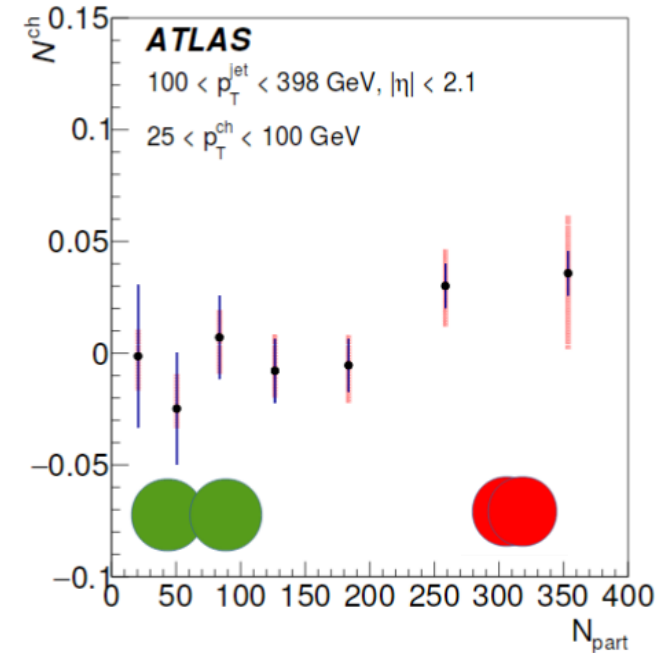
I.



II.



III.



- Tells us how many extra/missing particles are present in a given p_T range
- Observed a clear increase of yields of particles with low transverse momenta as the collision's centrality increases
- Particles with $p_T > 25 \text{ GeV}$ do not exhibit noticeable variations with centrality

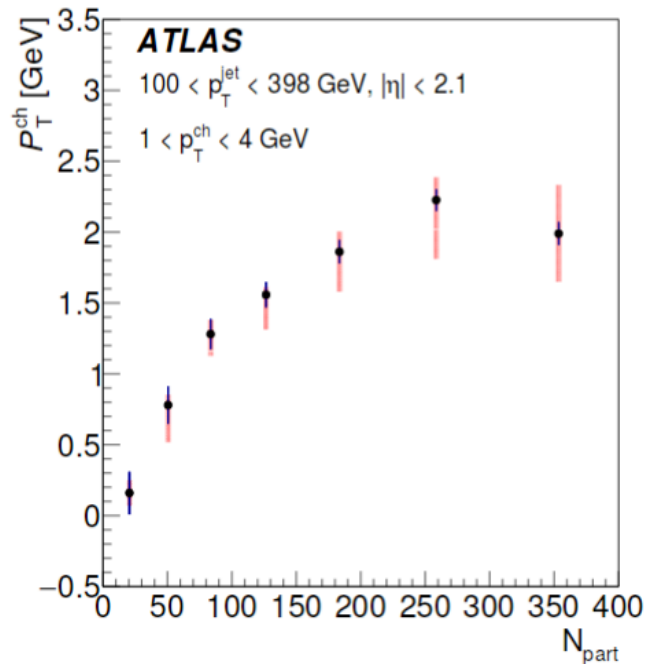
Jet fragmentation – flow of momentum

Pb+Pb @ 2.76 TeV

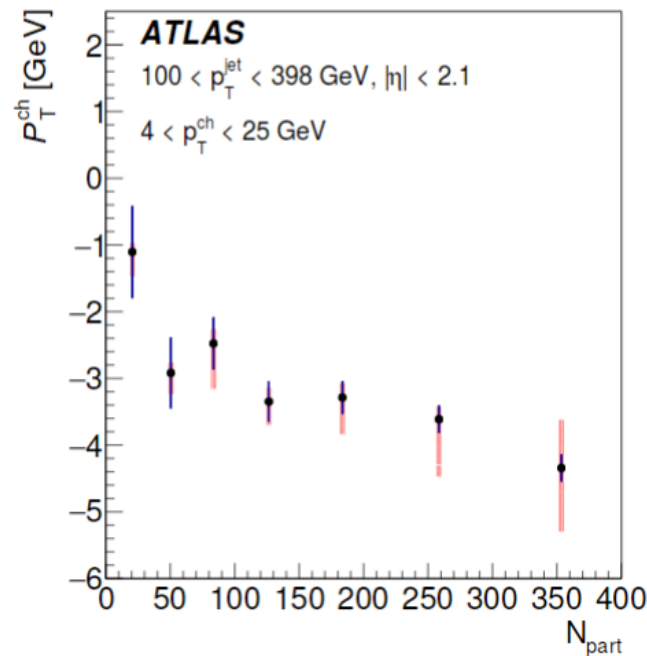
arXiv:1702.00674

- The flow of momentum as a function of N_{part} :
$$P_T^{\text{ch}} \equiv \int_{p_{T,\text{min}}}^{p_{T,\text{max}}} \left(D(p_T)|_{\text{cent}} - D(p_T)|_{\text{pp}} \right) p_T dp_T$$

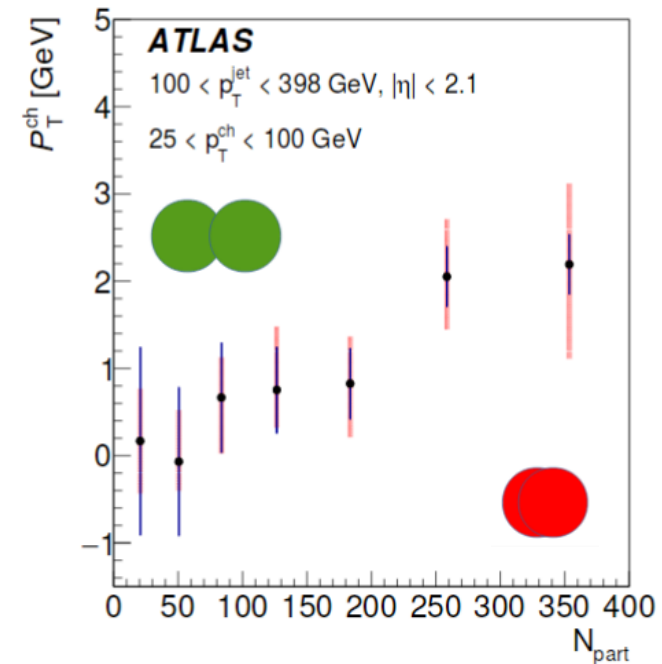
I.



II.



III.

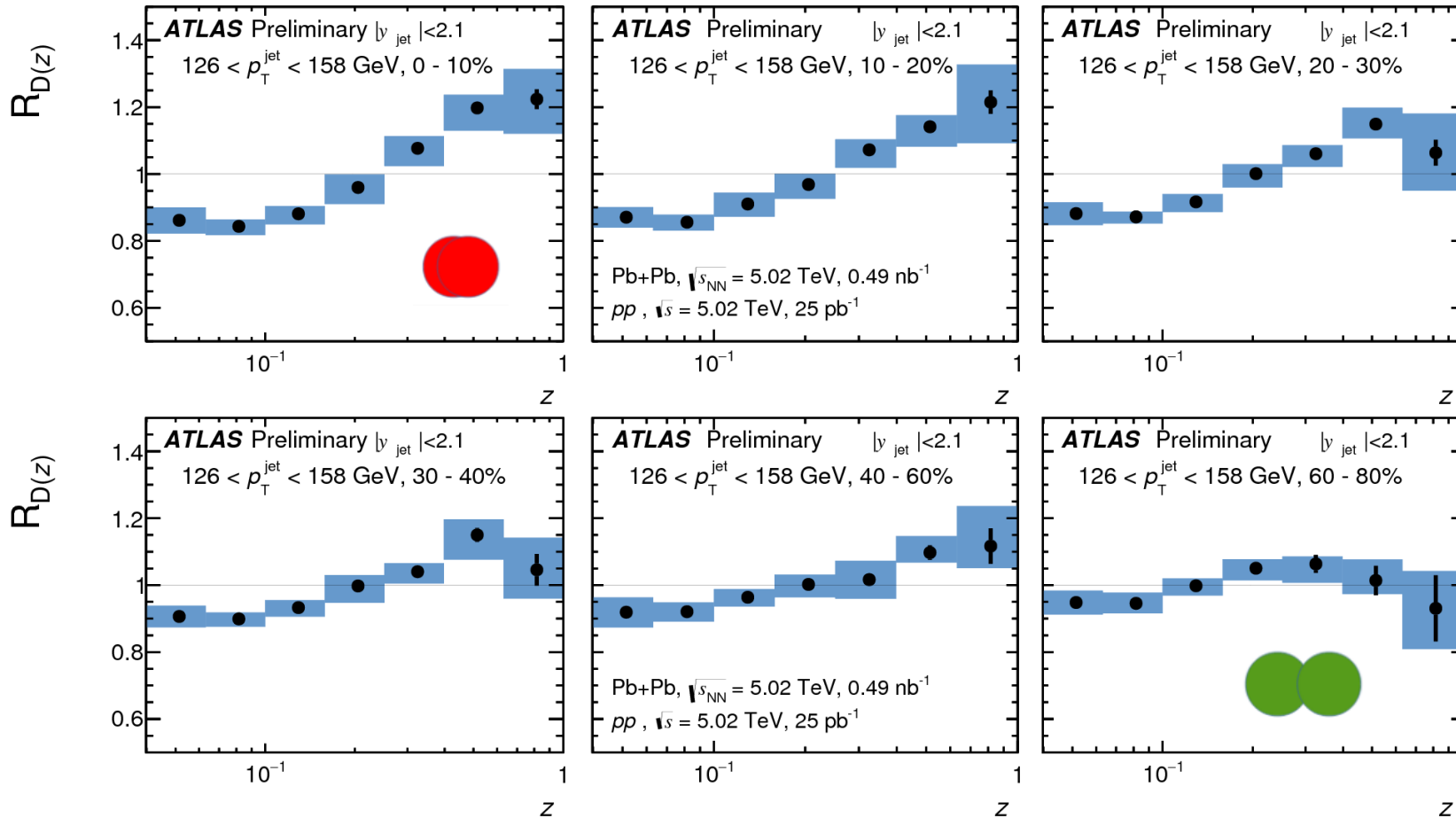


- Tells us how much p_T is carried by extra/missing particles in a given p_T range
- The changes in the total transverse momentum follow the trends seen in the yields.

Jet fragmentation ratios Pb+Pb

Ratios of $D(z)$ for **6 centralities** in one p_T bin

@ 5.02 TeV



Centrality dependence

- Enhancement at high z
- Suppression at intermediate z

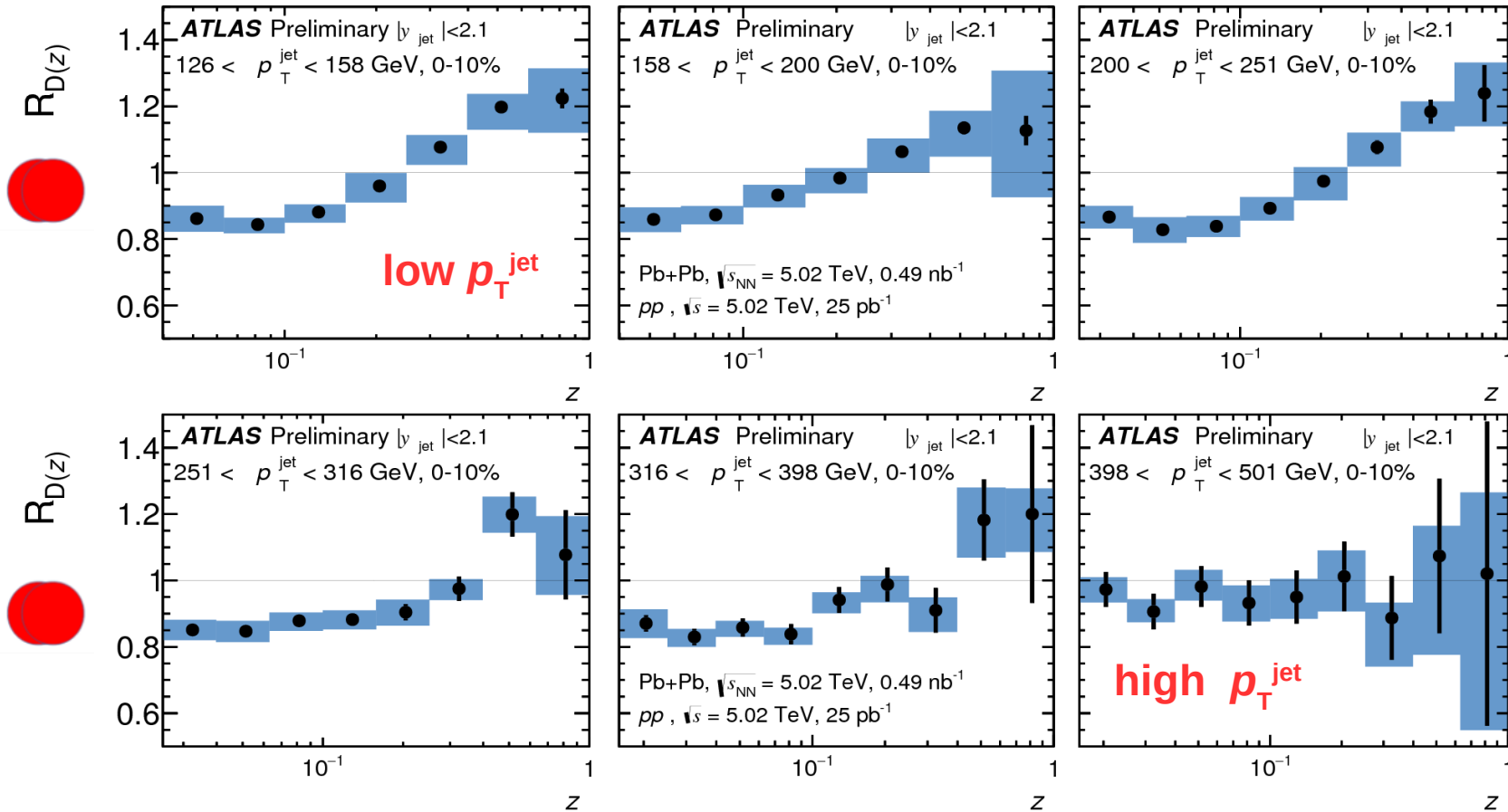
$126 < p_T < 158$ GeV

- Cut on $p_{T}^{\text{trk}} > 4$ GeV
- Missing the low p_T enhancement

Jet fragmentation ratios Pb+Pb

Ratios of $D(z)$ for central events and **6 p_T bins**

@ 5.02 TeV



Centrality dependence

- Enhancement at high z
- Suppression at intermediate z

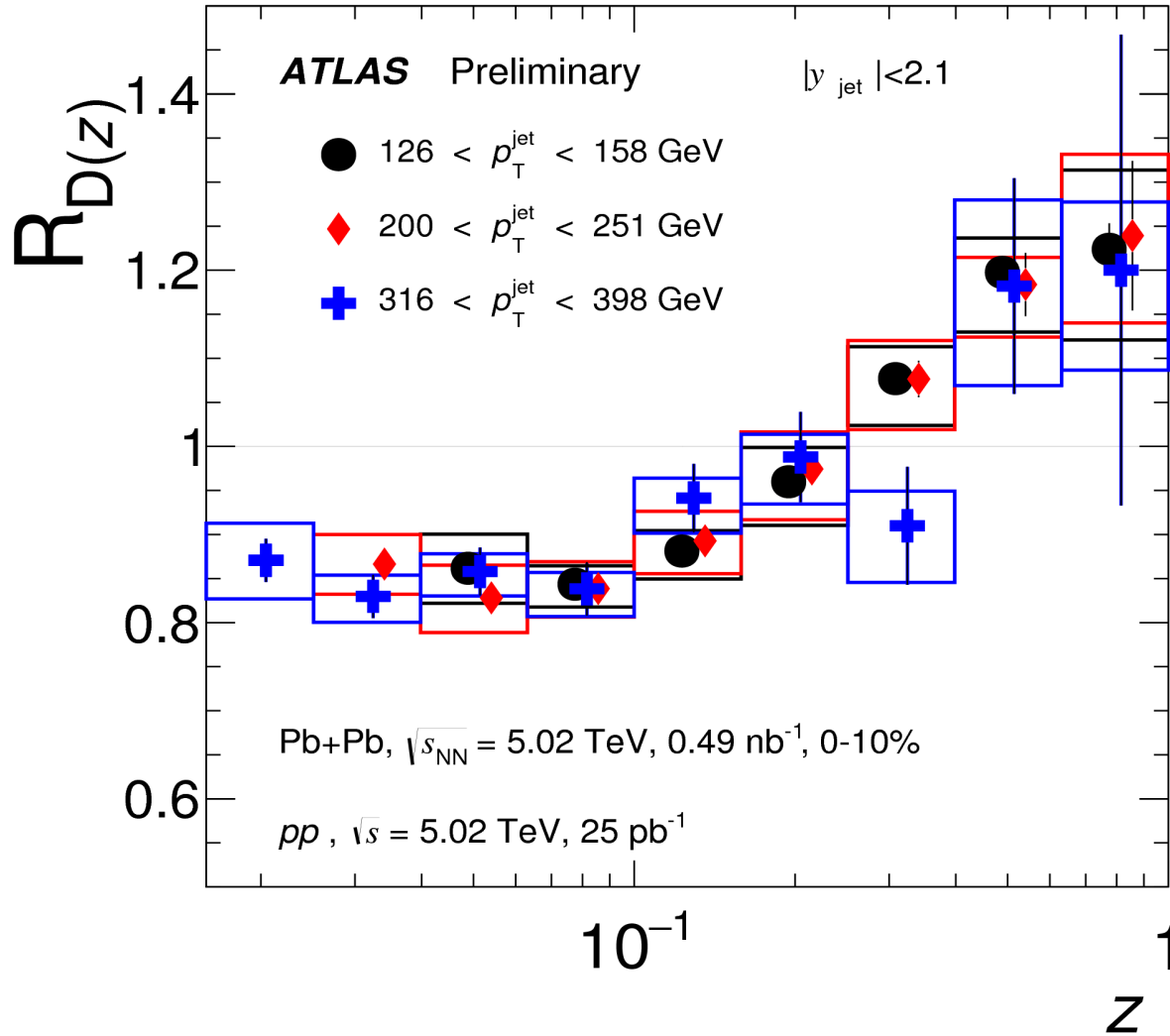
Jet p_T dependence

- No significant dependence on jet p_T

$126 < p_T < 501 \text{ GeV}$

- Cut on $p_{T, \text{trk}} > 4 \text{ GeV}$
- Missing the low p_T enhancement

Jet fragmentation ratios Pb+Pb



@ 5.02 TeV

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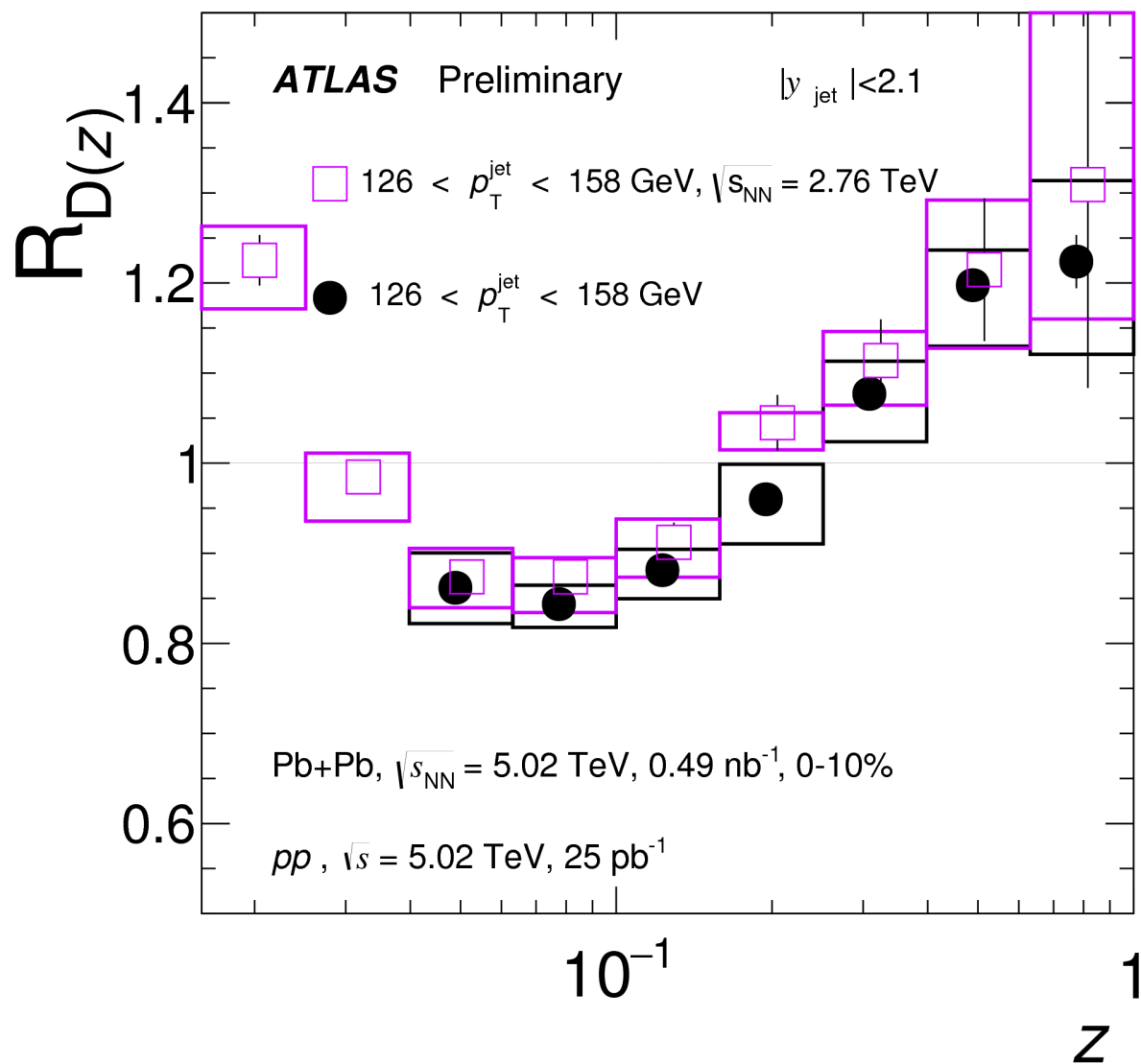
Cut on $p_{\text{T}}^{\text{trk}} > 4 \text{ GeV}$

$126 < p_{\text{T}}^{\text{jet}} < 398 \text{ GeV}$

No significant dependence on jet p_{T}

Comparison of $R_{D(z)}$ ratios of three $p_{\text{T}}^{\text{jet}}$ bin selections as a function of z .

Comparison of $R_{D(z)}$ in Pb+Pb



@ 2.76 TeV

@ 5.02 TeV

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Cut on $p_{\text{T}}^{\text{trk}} > 1 \text{ GeV}$
for **2.76 TeV**

Cut on $p_{\text{T}}^{\text{trk}} > 4 \text{ GeV}$
for **5.02 TeV**

$126 < p_{\text{T}}^{\text{jet}} < 158 \text{ GeV}$

No observed dependence on \sqrt{s}

Inclusive jets suppression

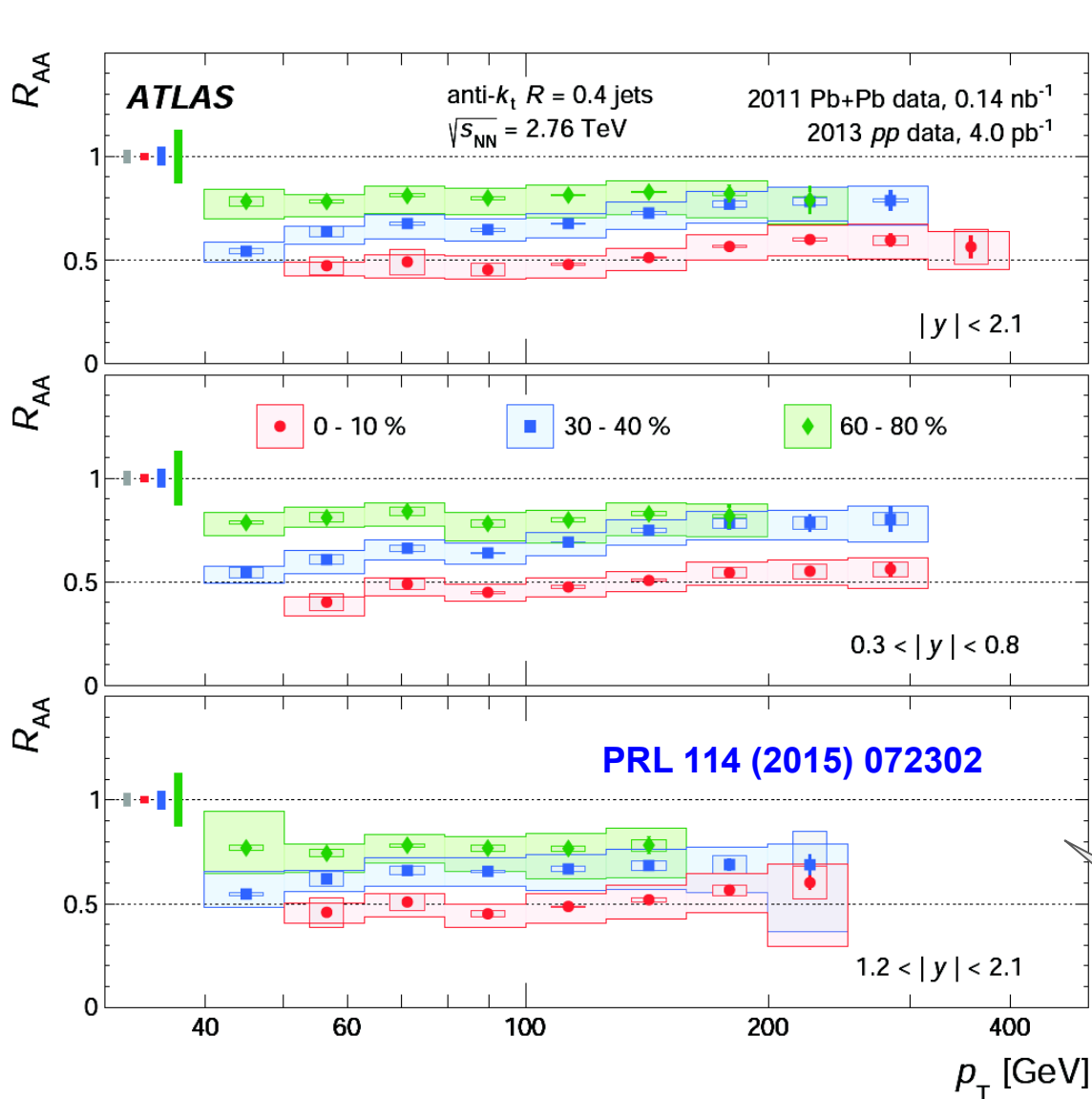
$$R_{AA} = \frac{\frac{1}{N_{evnt}} \frac{d^2 N_{jet}^{PbPb}}{dp_T dy} \Big|_{cent}}{\langle T_{AA} \rangle_{cent} \times \frac{d^2 \sigma_{jet}^{pp}}{dp_T dy}}$$

Jet yield in heavy ion collisions
Nuclear thickness function
Jet cross-section in pp collisions

Number of expected jets per event of a given centrality

Nuclear modification factor quantifies the magnitude of the jet suppression which is dominantly due to final state interactions with constituents of the medium

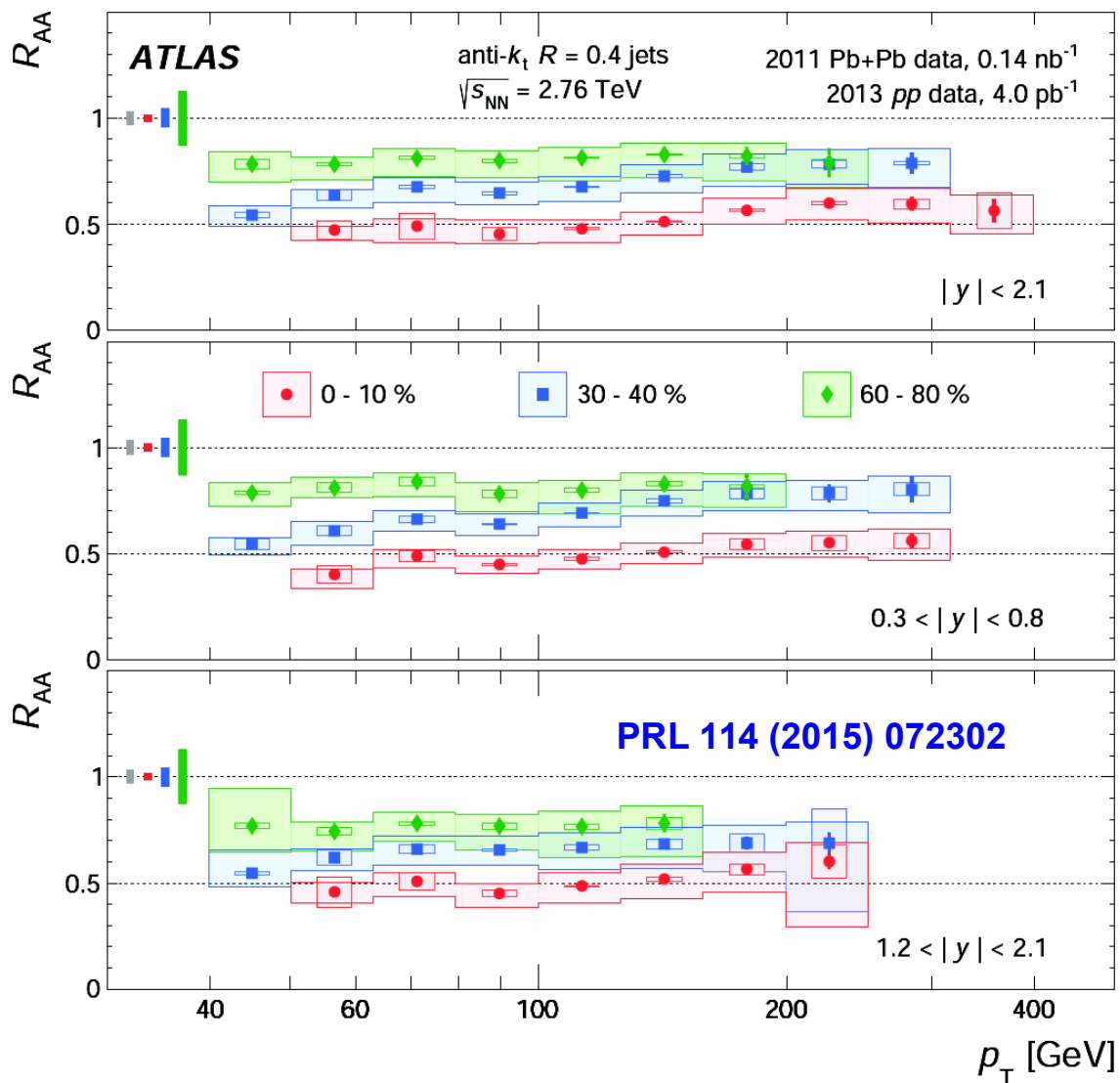
Jet R_{AA} @ 2.76 TeV



$$R_{AA} = \frac{1}{N_{\text{evnt}}} \frac{d^2 N_{\text{jet}}^{\text{PbPb}}}{dp_T dy} \Big|_{\text{cent}}}{\langle T_{AA} \rangle_{\text{cent}} \times \frac{d^2 \sigma_{\text{jet}}^{\text{pp}}}{dp_T dy}}$$

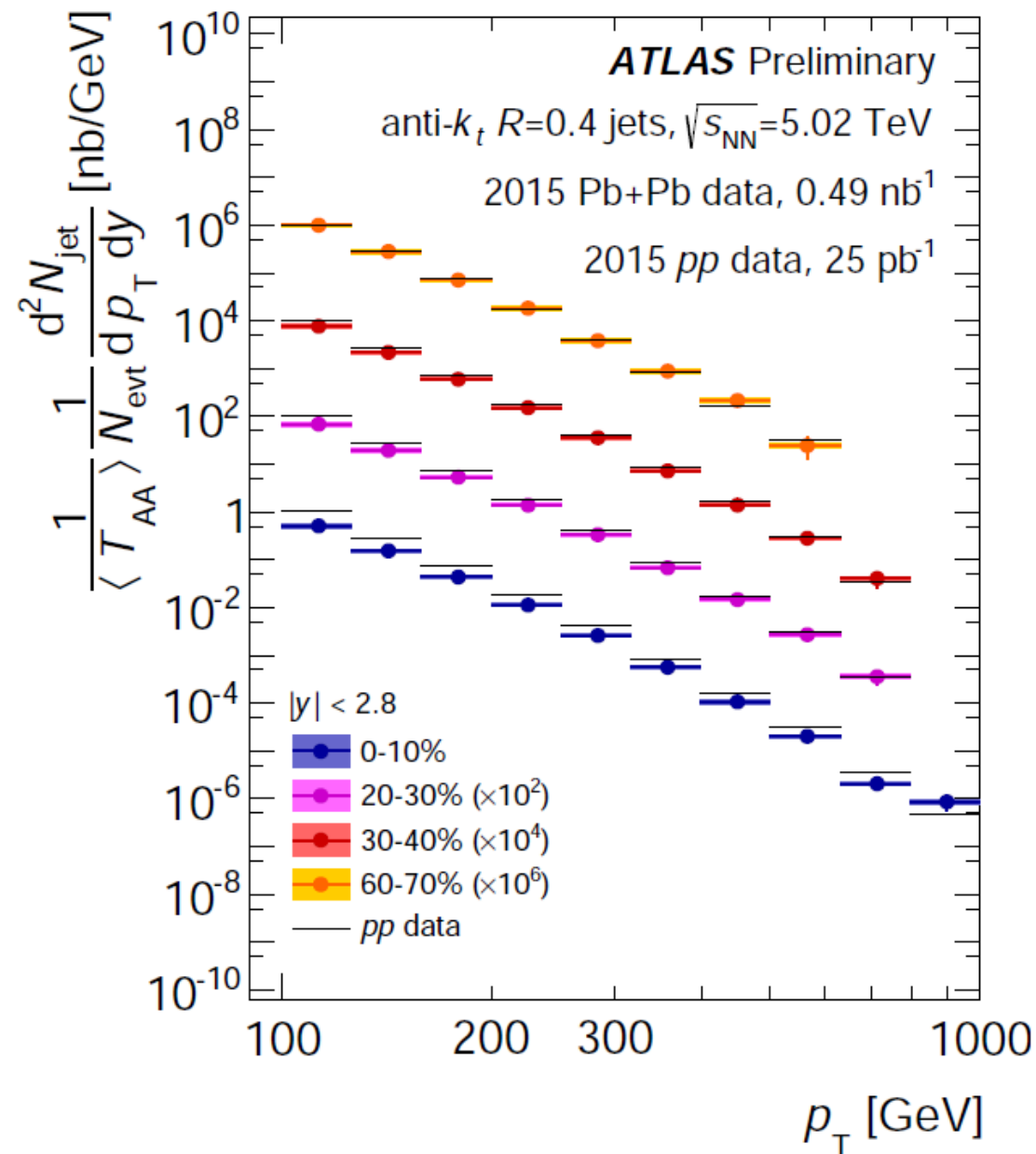
- A modest grow of jet R_{AA} with increasing jet p_T
- Still significant suppression even for 60 – 80% centrality bin
- No rapidity dependence observed

Angerami, Cole,
Kosek, Spousta



- Goal of the new analysis:
 - Repeat measurement for new energy @ 5.02 TeV
 - Extend the measurement to higher jet momenta
 - Extend the measurement to higher jet rapidities

Slovák, Spousta,
Havener

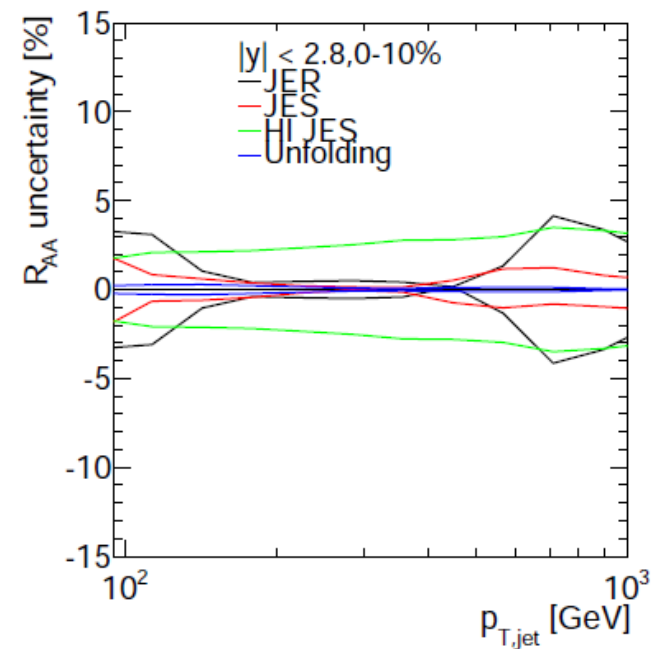
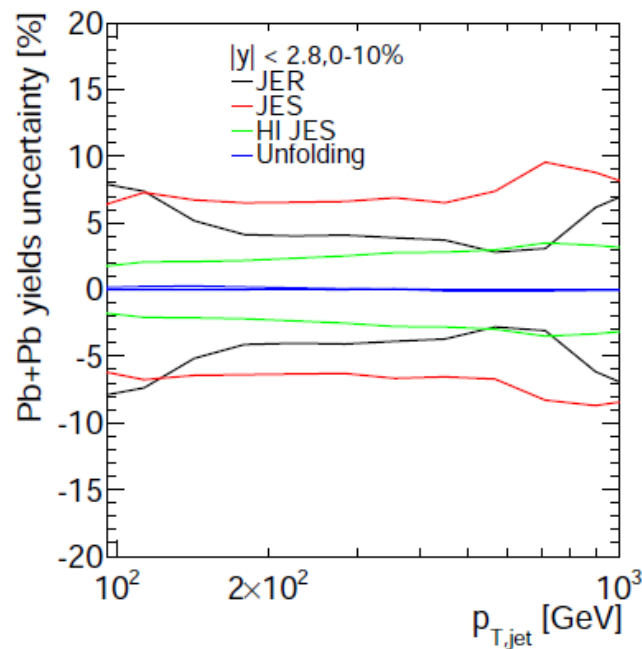
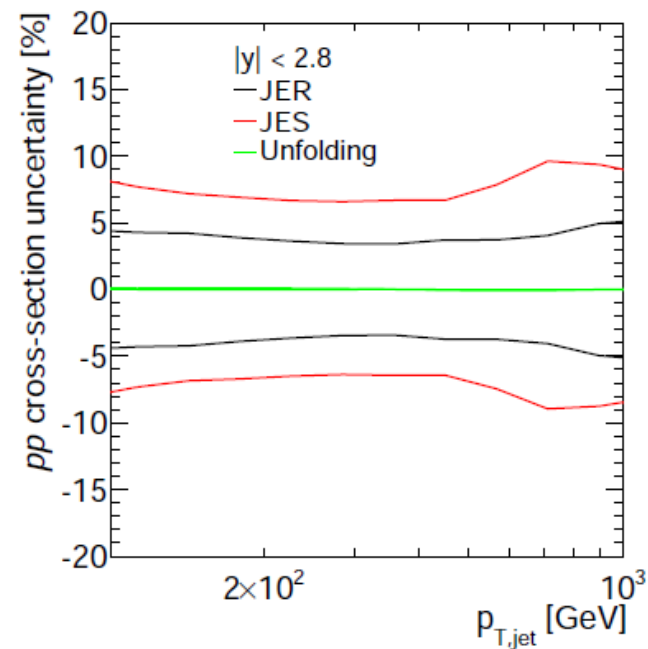


Kinematics region:

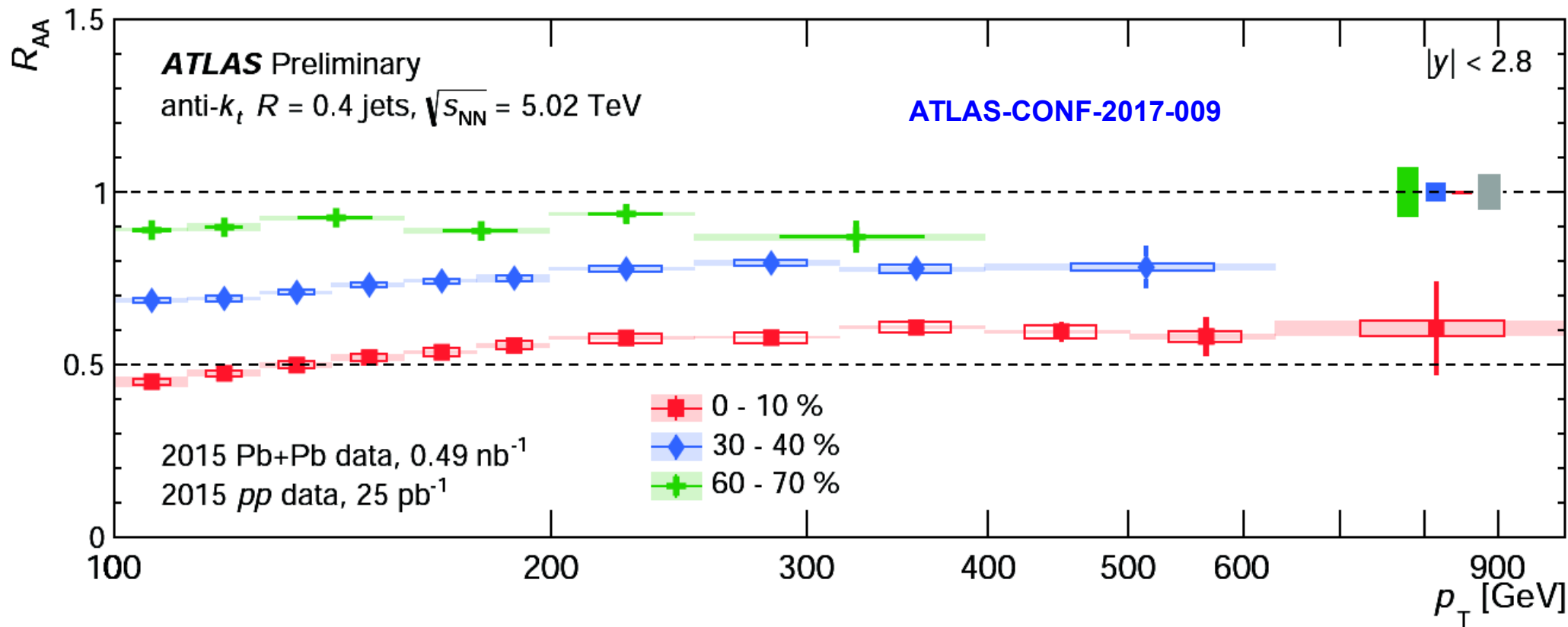
- $R = 0.4$ jets with $p_T > 100$ GeV
- Six bins in rapidity: $|y| < 0.3$, $0.3 < |y| < 0.8$, $0.8 < |y| < 1.2$, $1.2 < |y| < 1.6$, $1.6 < |y| < 2.1$, $2.1 < |y| < 2.8$, $|y| < 2.1$, $|y| < 2.8$
- Eight bins in centrality: 0-10% to 70-80%
- 1D Bayesian unfolding

Systematical uncertainties

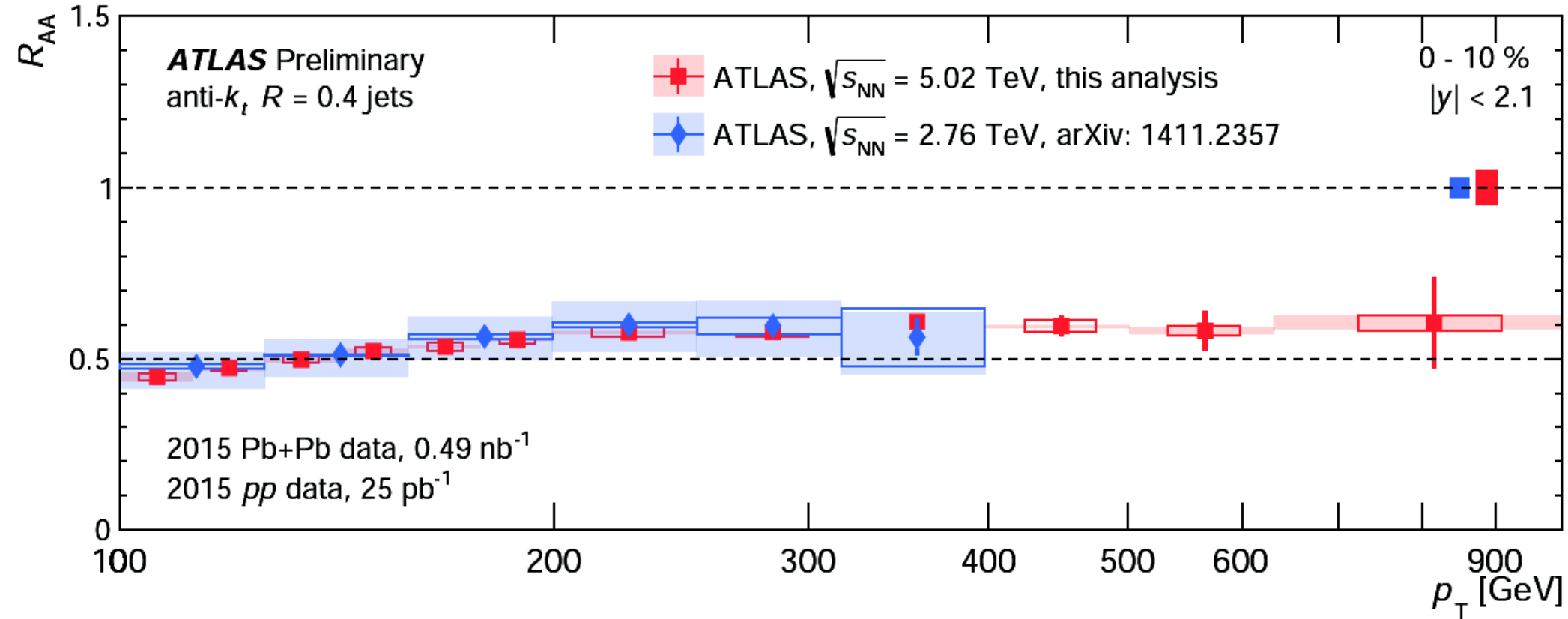
- Jet energy scale:
 - Standard pp component + 5 TeV flavor a HI cross-calibration
 - HI uncertainties due to jet quenching
- Jet energy resolution - standard pp a HI part
- Luminosity
- Uncertainty on T_{AA}
- Unfolding – results unfolded with not weighted response matrix



Jet R_{AA} @ 5.02 TeV p_T -dependence

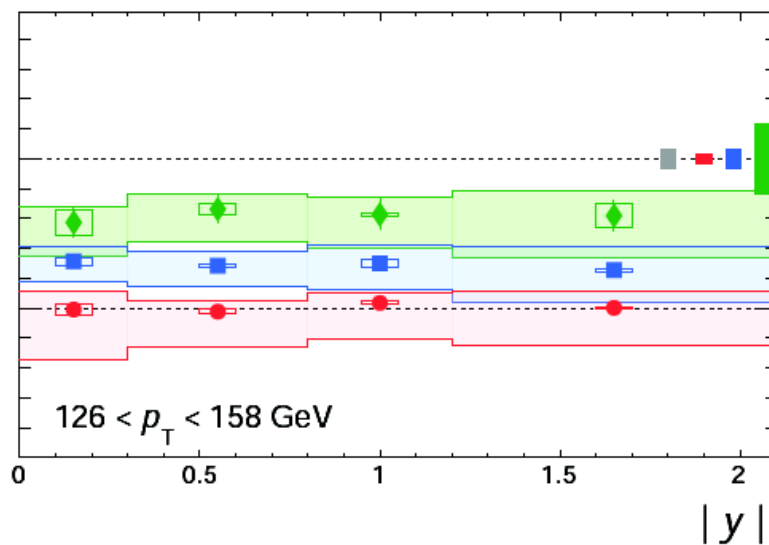
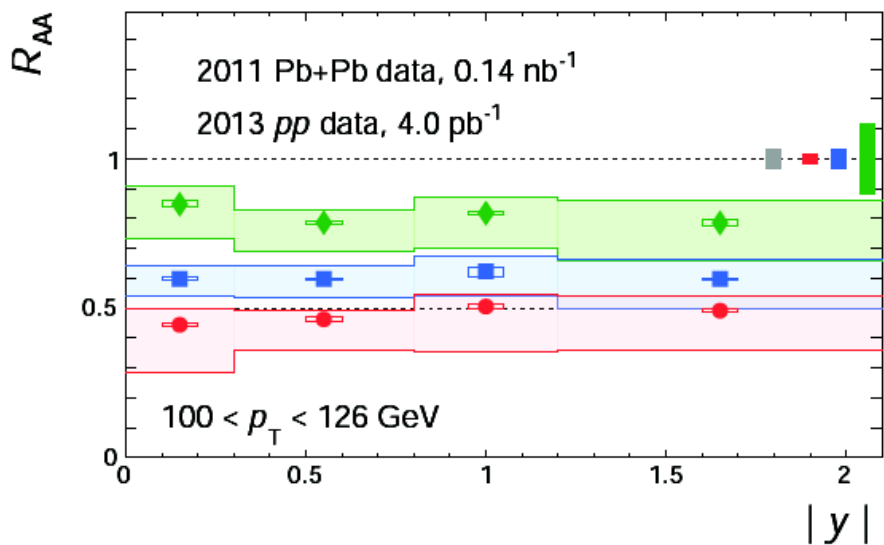
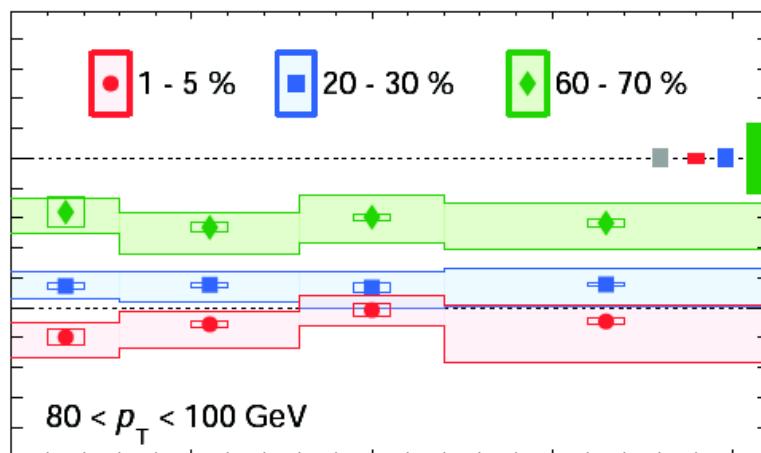
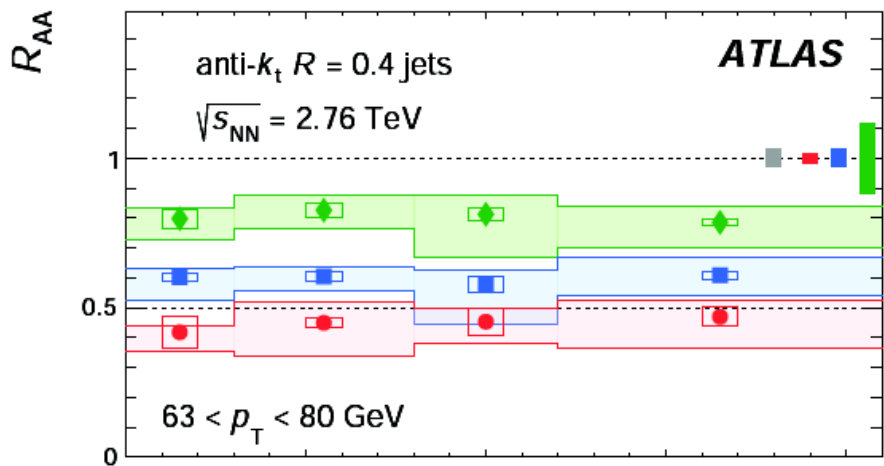


- R_{AA} measured for jets with p_T **100 GeV - 1 TeV** in rapidity $|y| < 2.8$.
- The central events suppressed by factor ~ 2 . For jets within 100 ~ 200 GeV we observe a modest grow, then R_{AA} is flat.



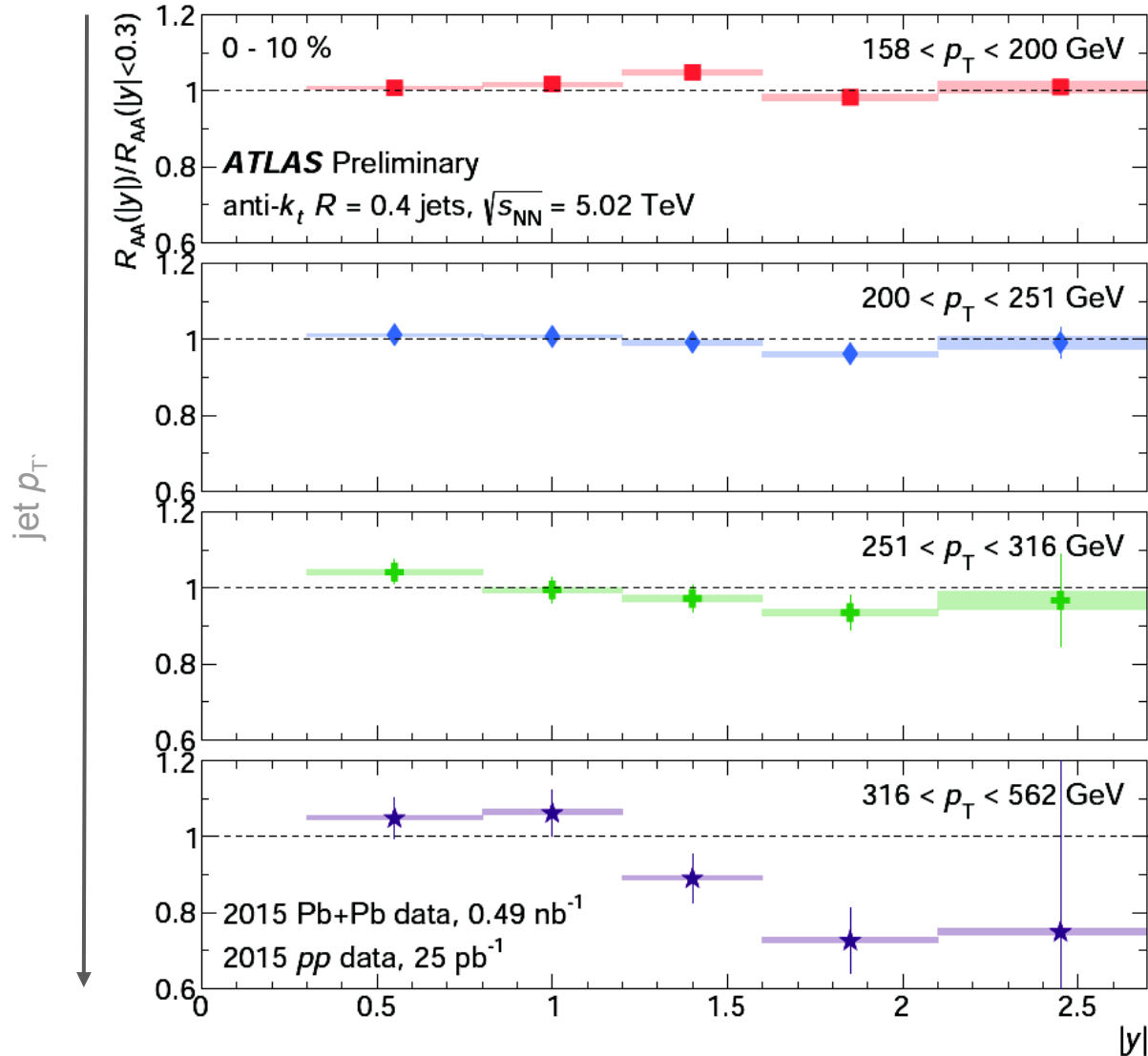
- The magnitude of jet suppression is the same within the systematic uncertainties for both measurement
- A significant reduction of systematic uncertainties for the **new measurement** mainly due to the fact than pp and Pb+Pb data were taken in the same period.

Jet R_{AA} @ 2.76 TeV – dependence on y



No rapidity dependence observed

Jet R_{AA} @ 5.02 TeV – dependence on y

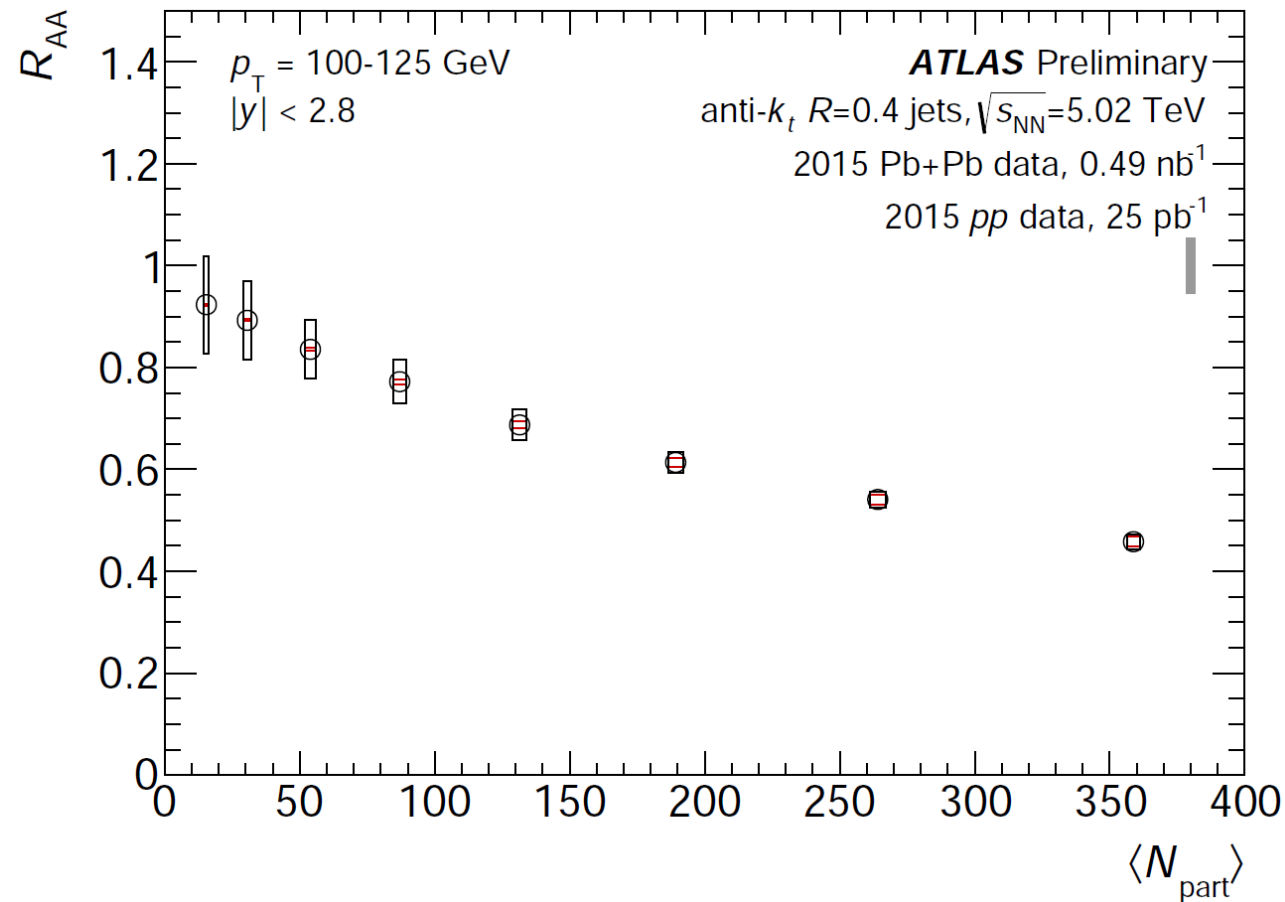


The vertical axis describes the ratio of R_{AA} for a given rapidity and R_{AA} for jets with $|y| < 0.3$.

$$\frac{R_{AA}(y)}{R_{AA}(|y| < 0.3)}$$

With increasing momentum R_{AA} is getting smaller in the forward rapidity as compared to the mid-rapidity region.

Jet R_{AA} @ 5.02 TeV – dependence on N_{part}



$100 < p_{T}^{jet} < 125$ GeV
 $|y| < 2.8$

$\langle N_{part} \rangle$ = the average number of participating nucleons for a given centrality interval.

Smooth decrease of R_{AA} with increasing centrality.

Summary

- We have presented a measurement of fragmentation variables in Pb+Pb collisions @ 2.76 TeV, @ 5.02 TeV, p+Pb collision @ 5.02 TeV and pp collisions for the reference measurement at corresponding energy
- Pb+Pb @ 2.76 TeV and @ 5.02 TeV: - clear centrality dependence
- no observed dependence on \sqrt{s}
- Pb+Pb @ 2.76 TeV – hint of rapidity dependence
- Observed an increase of yields of particles and transverse momentum as centrality increases for particles with low transverse momenta
- Pb+Pb @ 5.02 TeV - we observe the lack of p_T dependence
- We have presented a measurement of nuclear modification factor in Pb+Pb collisions @ 5.02 TeV and pp collisions for the reference measurement at corresponding energy where we observed:
 - a significant suppression up to ~ 1 TeV
 - rapidity dependence of R_{AA} for jets with high momentum
 - a quantitative comparison with the measurement at 2.76 TeV

Back up

ATLAS

Inner Detector
 $|η| < 2.5$

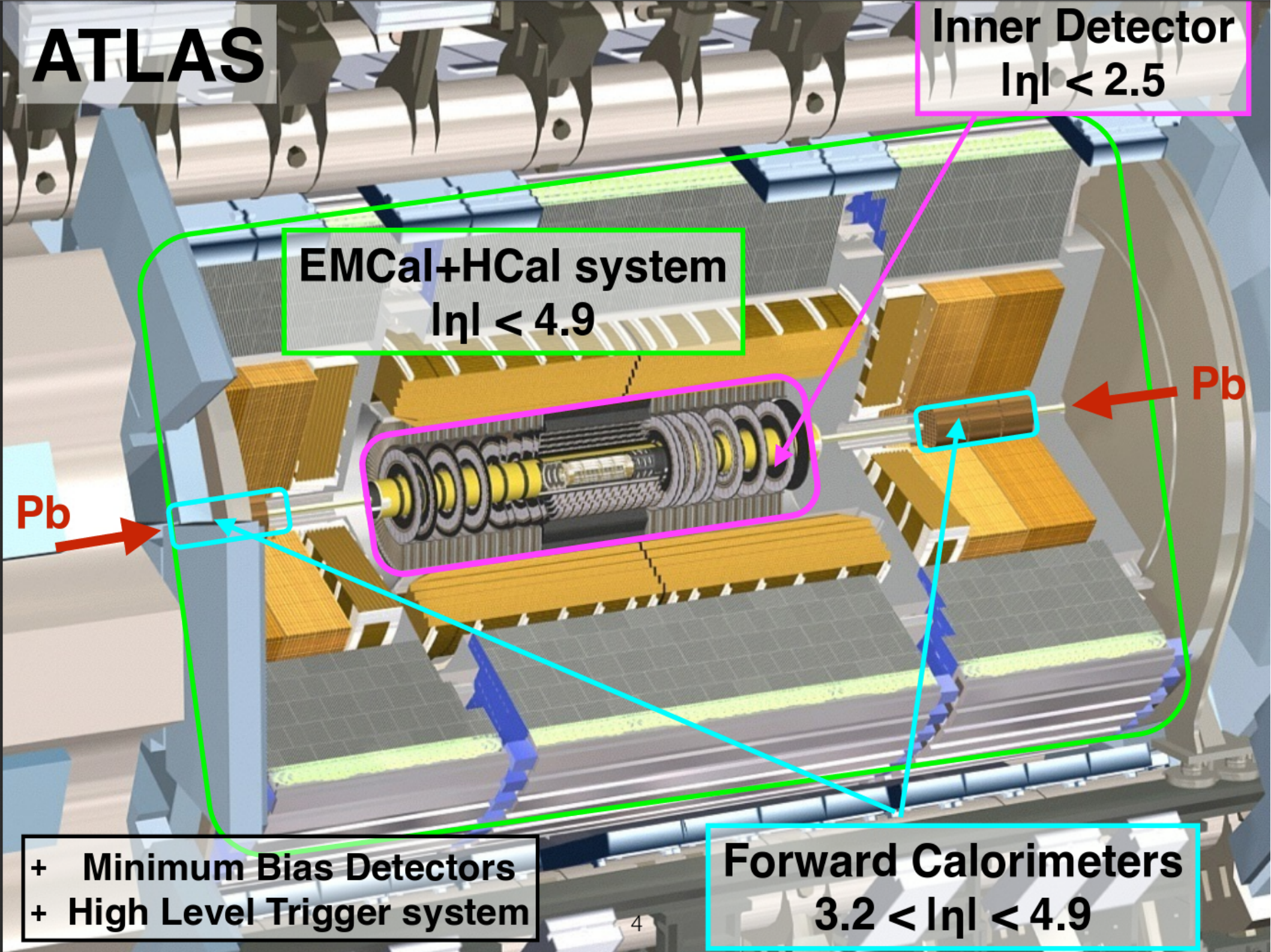
EMCal+HCal system
 $|η| < 4.9$

Pb

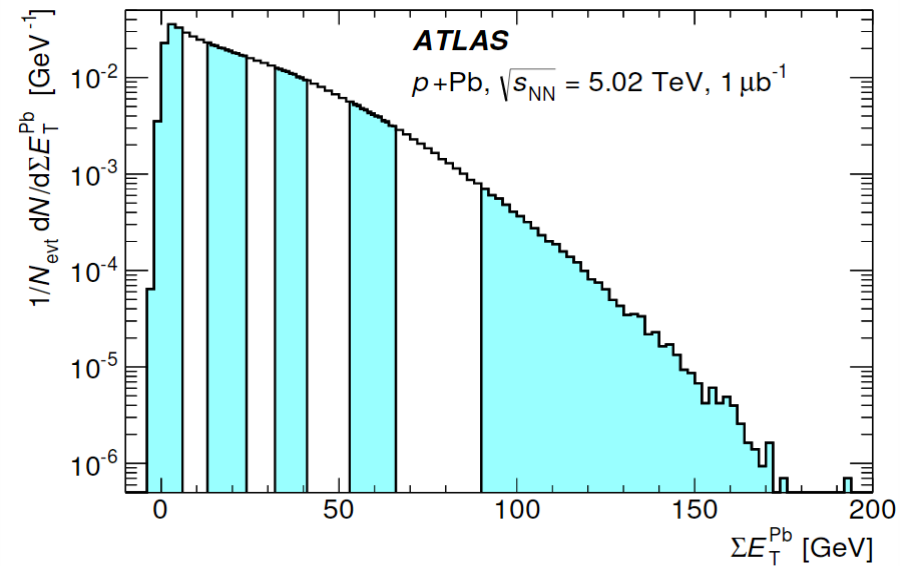
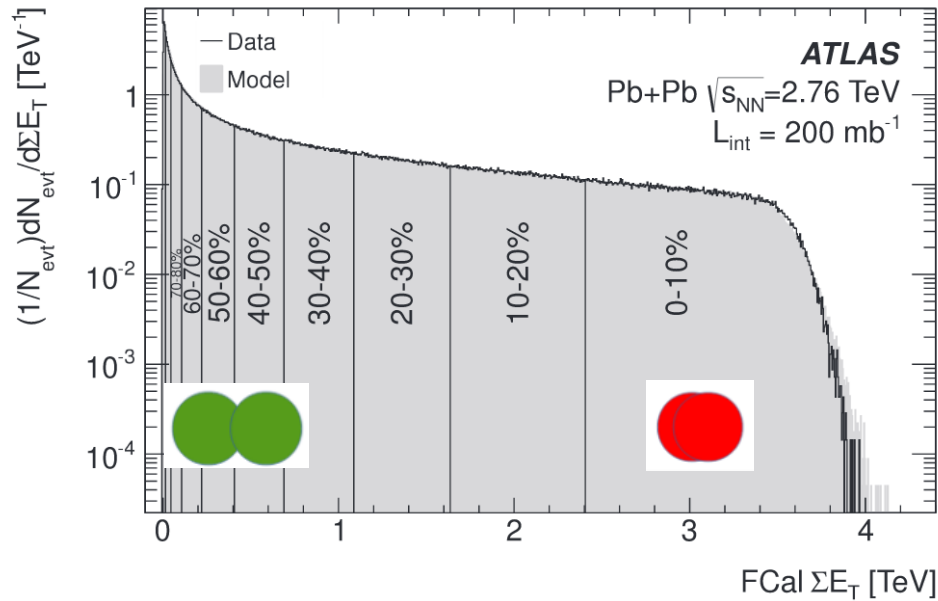
Pb

+ Minimum Bias Detectors
+ High Level Trigger system

Forward Calorimeters
 $3.2 < |η| < 4.9$



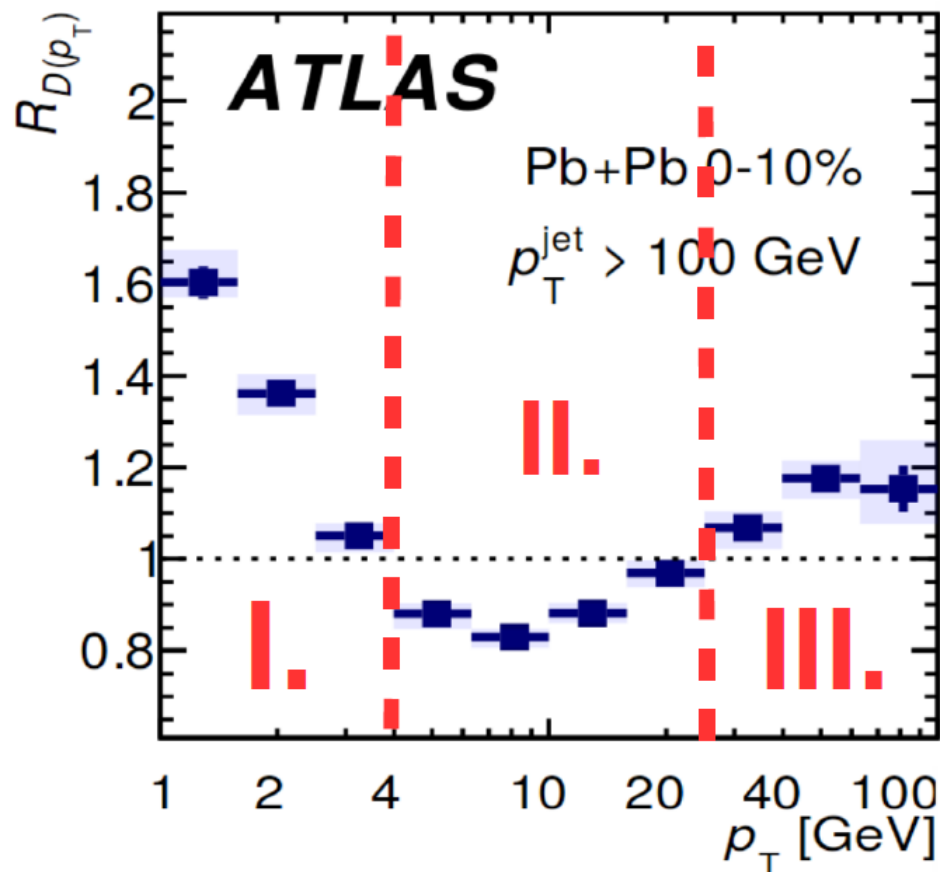
Centrality



- Centrality express measure of overlap of two colliding nuclei
- Determined by the sum of the transverse energy deposited in the Forward calorimeters
- It is closely related to the average number of participant nucleons and number of binary inelastic collisions
- Events divided into successive percentiles of the $\sum E_T^{FCal}$
- In Pb+Pb collisions use sum of the transverse energy in both sides
- In p+Pb collisions use sum of the transverse energy on Pb-going side only

Jet Fragmentation

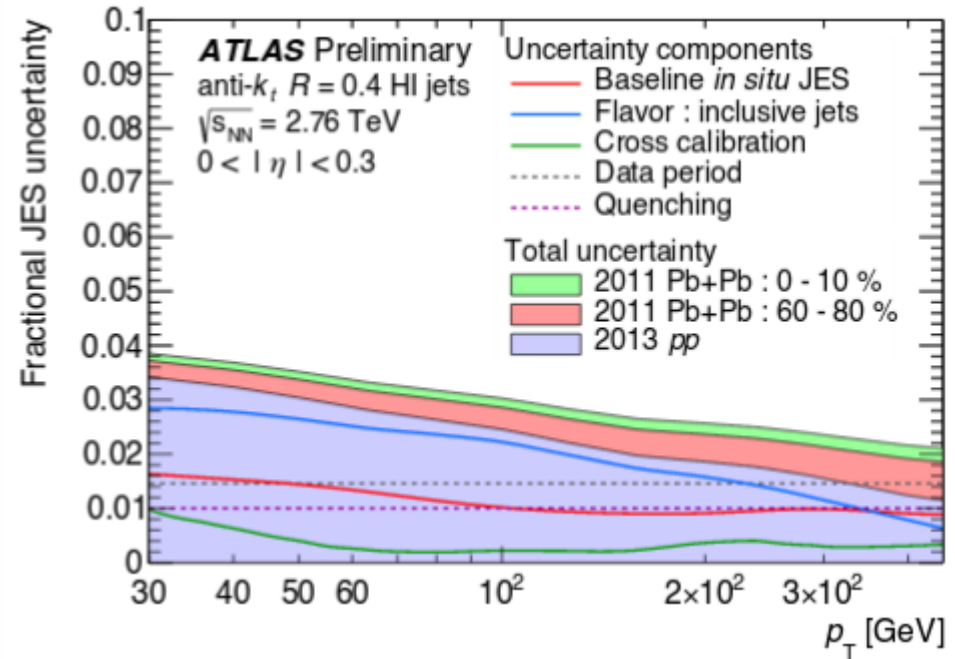
- Energy lost by parton is transferred mainly to the soft particles (I.)
- The modification on the middle and high p_T is caused by the different quenching of quark and gluon jets (II. + III.)



JES/JER uncertainty

JES:

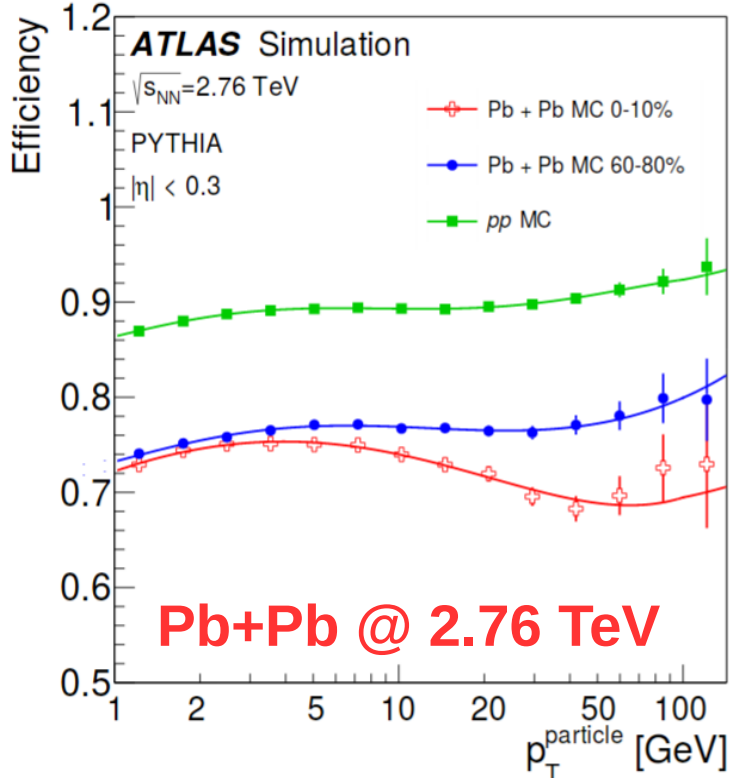
- Use the baseline 8 nuisance parameters from *in situ* calibration
- Additional parameters due to flavor response and composition and cross calibration
- Two additional parameters for Pb+Pb due to the difference in the data taking period and detector response to quenched jets



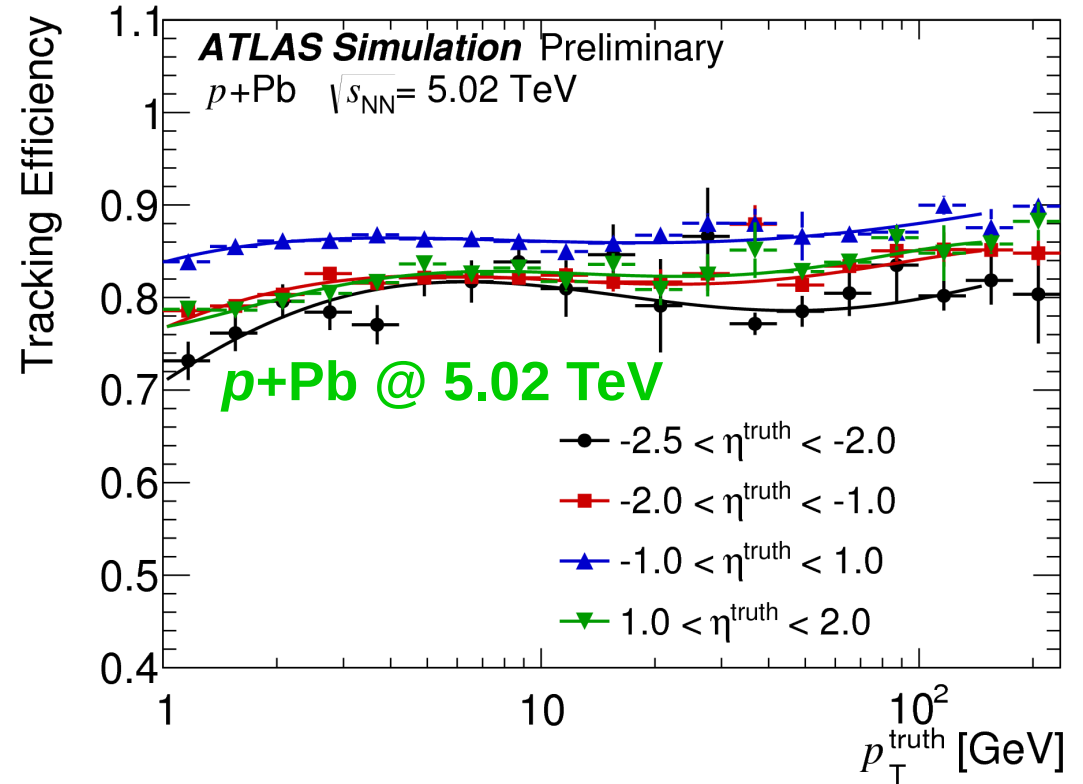
JER:

- Standard centrality-independent JER uncertainties
- Additional centrality dependent uncertainty for possible disagreement between fluctuations term in JES in the MC independent analysis of fluctuations in data
- This is very small because MC sample is data overlay

arXiv:1702.00674



ATLAS-CONF-2017-004



- Tracks with different p_T are used :

$p_T > 1$ GeV for **Pb+Pb @ 2.76 TeV** and **p+Pb @ 5.02 TeV**

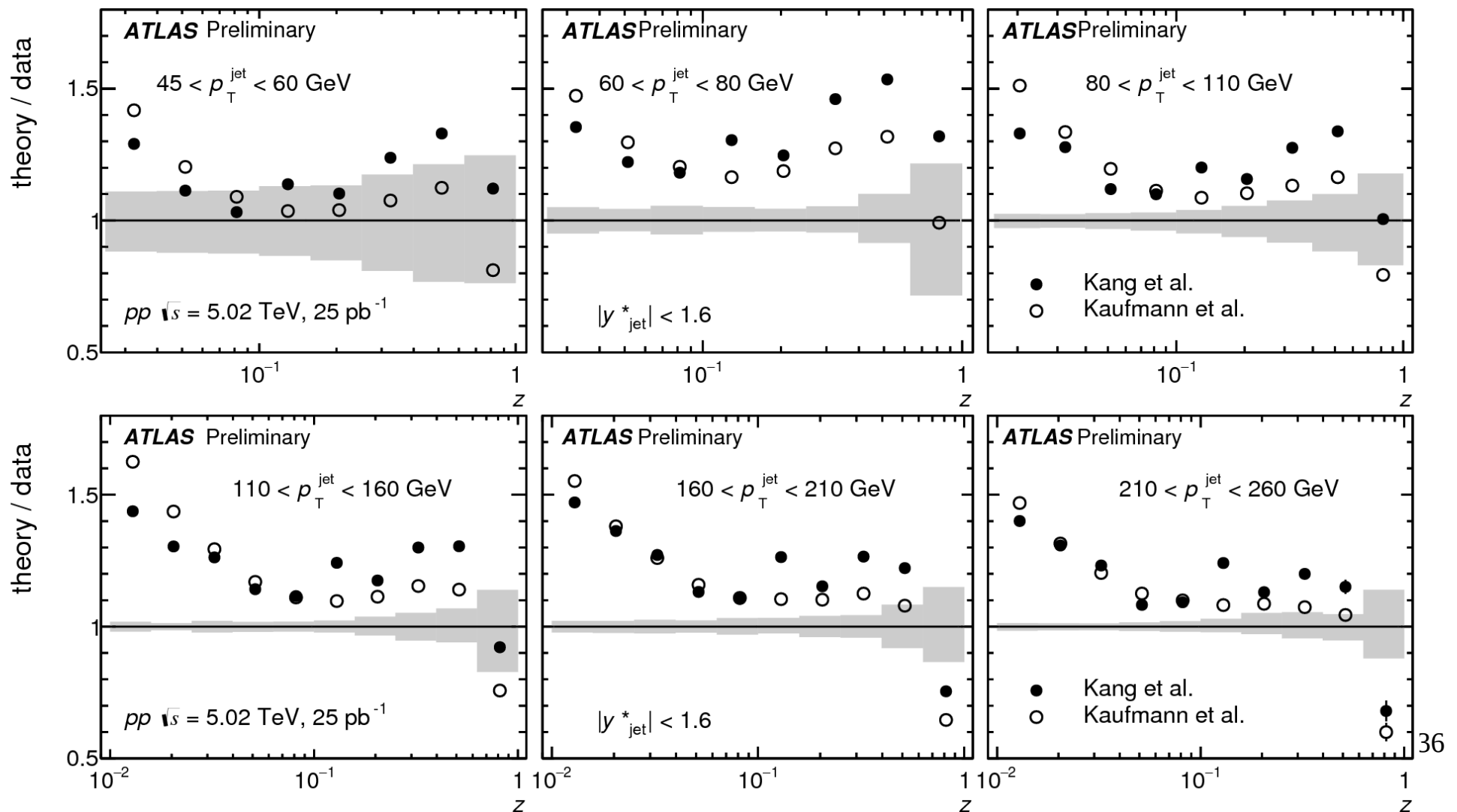
$p_T > 4$ GeV for **Pb+Pb @ 5.02 TeV**

- Tracking efficiency is parametrized as a function of track p_T and track η

Comparison of pp results to different theoretical calculation

pp @ 5.02 TeV

Significant differences between theory and data [ATLAS-CONF-2017-004](#)

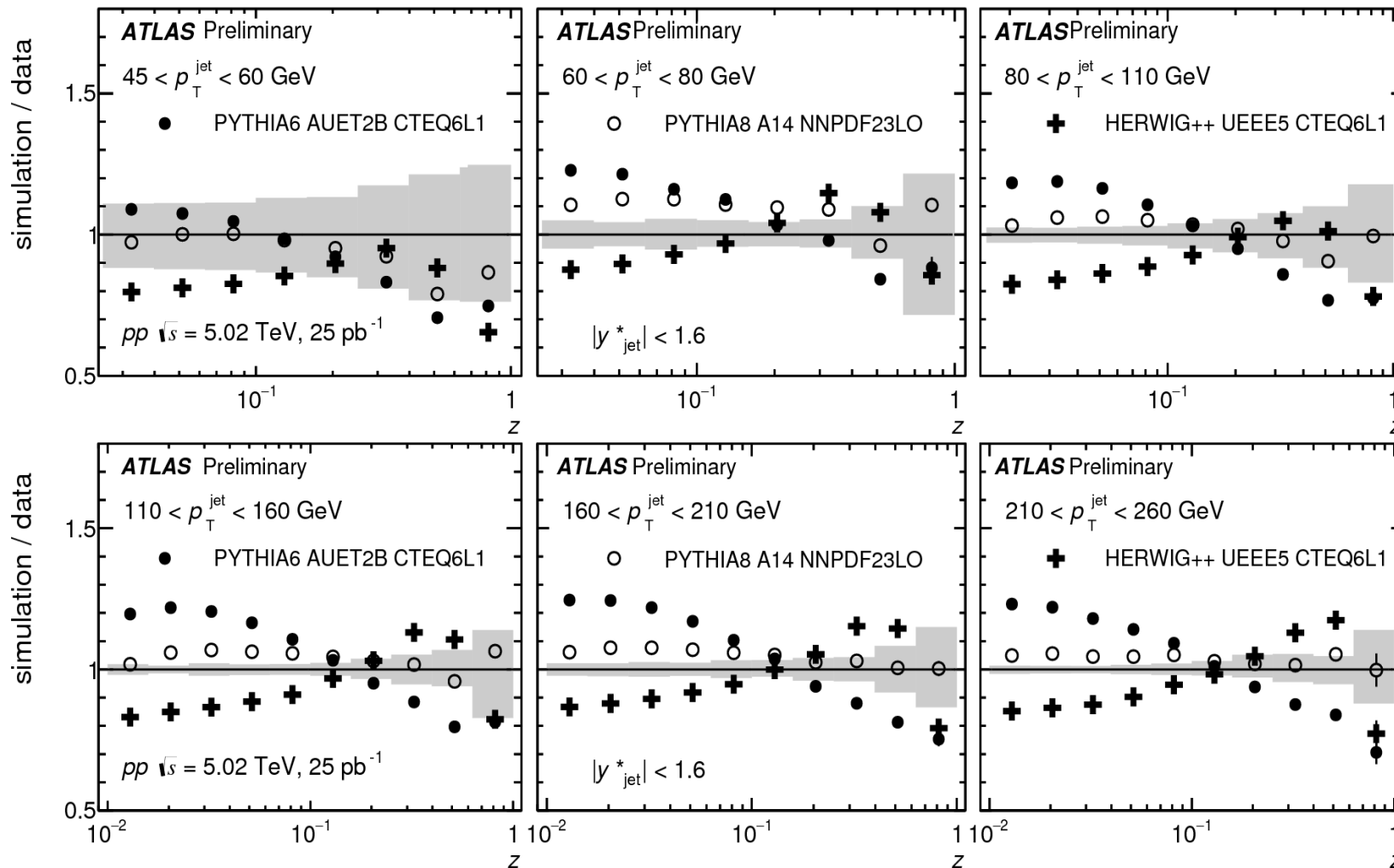


Comparison of pp results to different theoretical calculation

pp @ 5.02 TeV

- Significant differences between MC models
- PYTHIA8 describes the data in the best way

ATLAS-CONF-2017-004



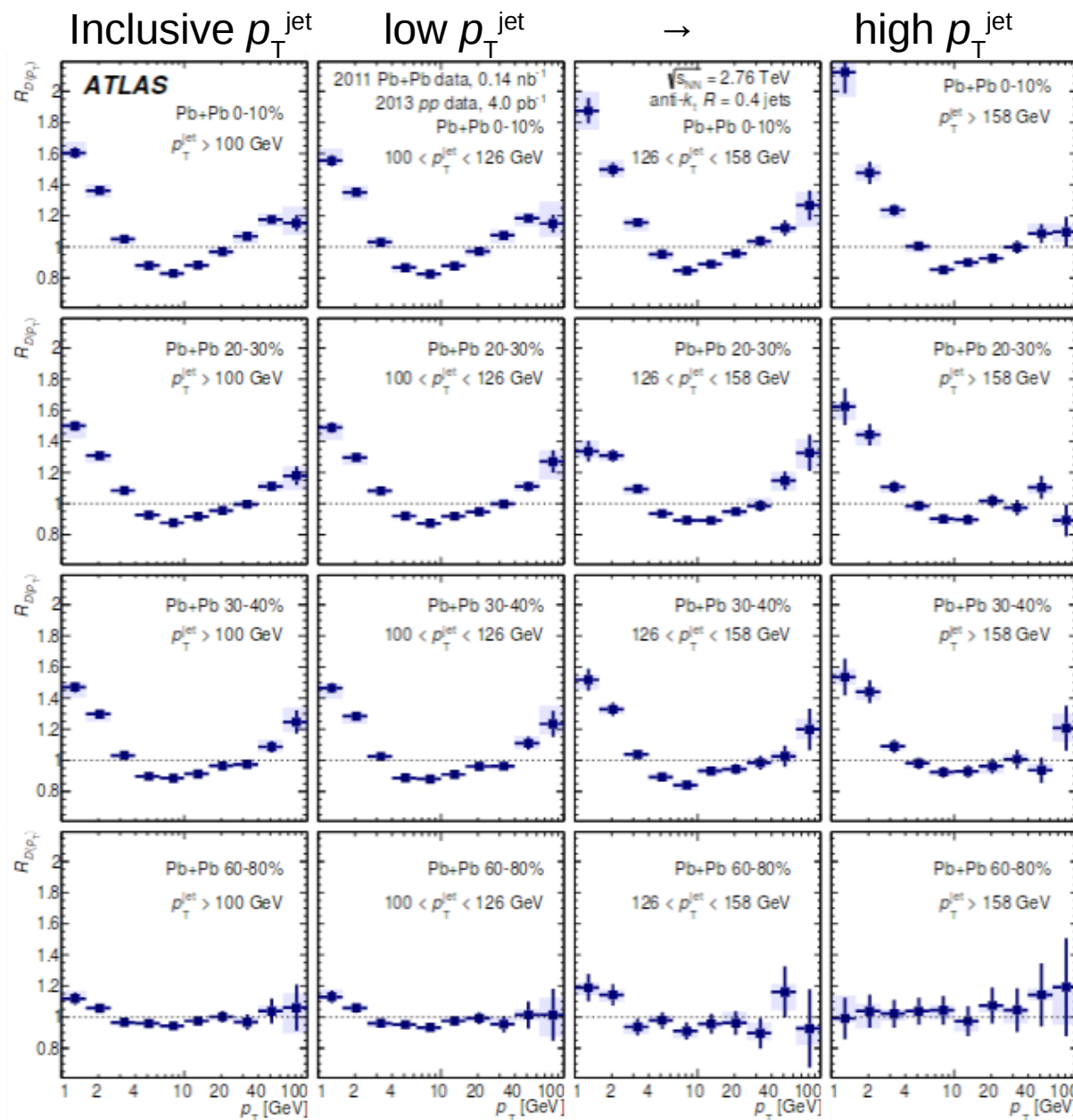
Jet fragmentation ratios Pb+Pb

Ratios of $D(p_T)$ for 4 centralities in 4 p_T bins

@ 2.76 TeV

[arXiv:1702.00674](https://arxiv.org/abs/1702.00674)

$100 < p_T < 300$ GeV



Centrality dependence

- Enhancement at low and high p_T
- Suppression at intermediate p_T

Jet p_T dependence

- No significant dependence on jet p_T

Similar observation for $D(z)$

Jet fragmentation ratios Pb+Pb

Ratios of $D(p_T)$ for 4 centralities in 4 **pseudorapidity** bins

$|\eta| < 2.1$ less forward \rightarrow more forward

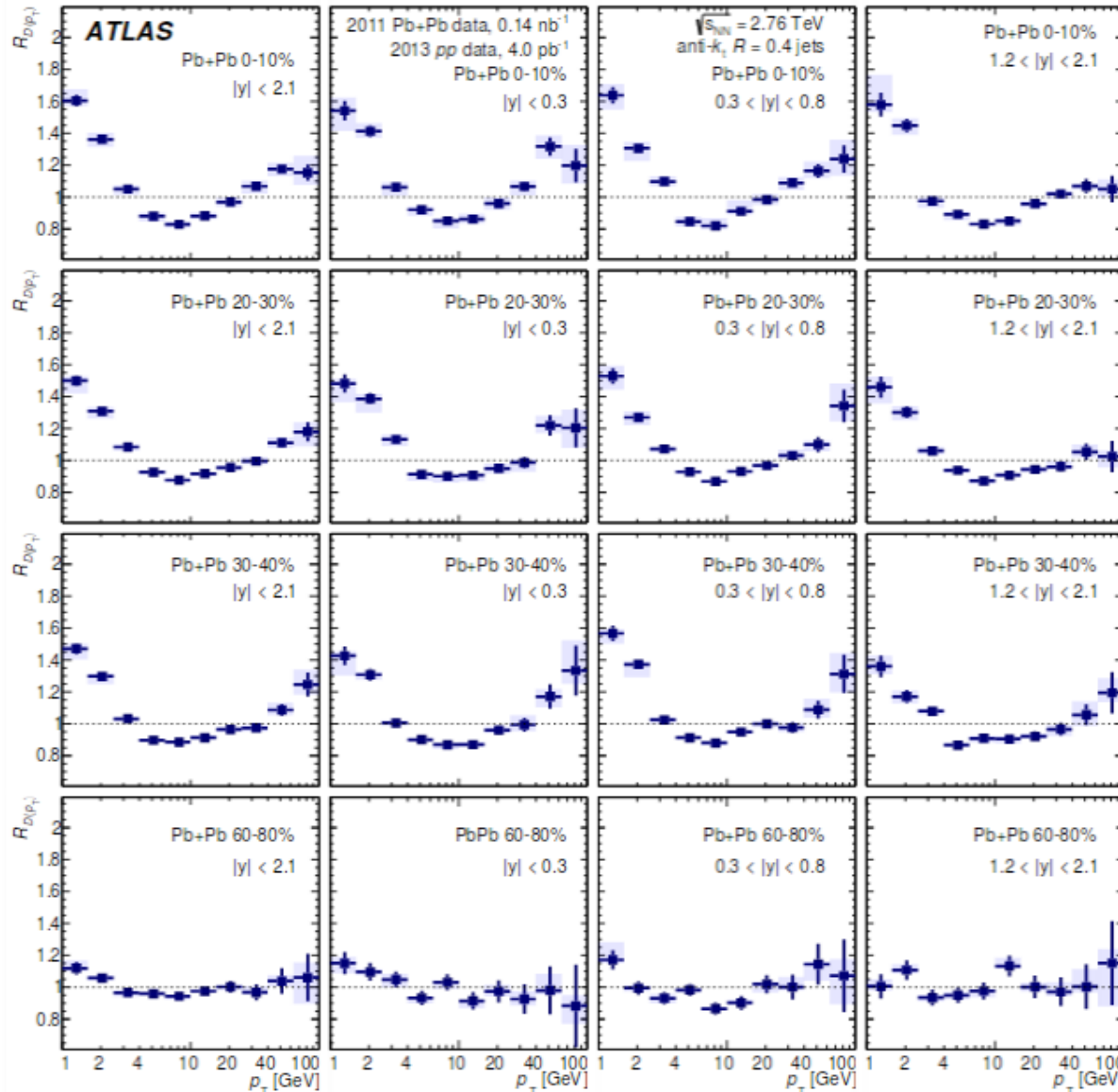
@ 2.76 TeV

[arXiv:1702.00674](https://arxiv.org/abs/1702.00674)

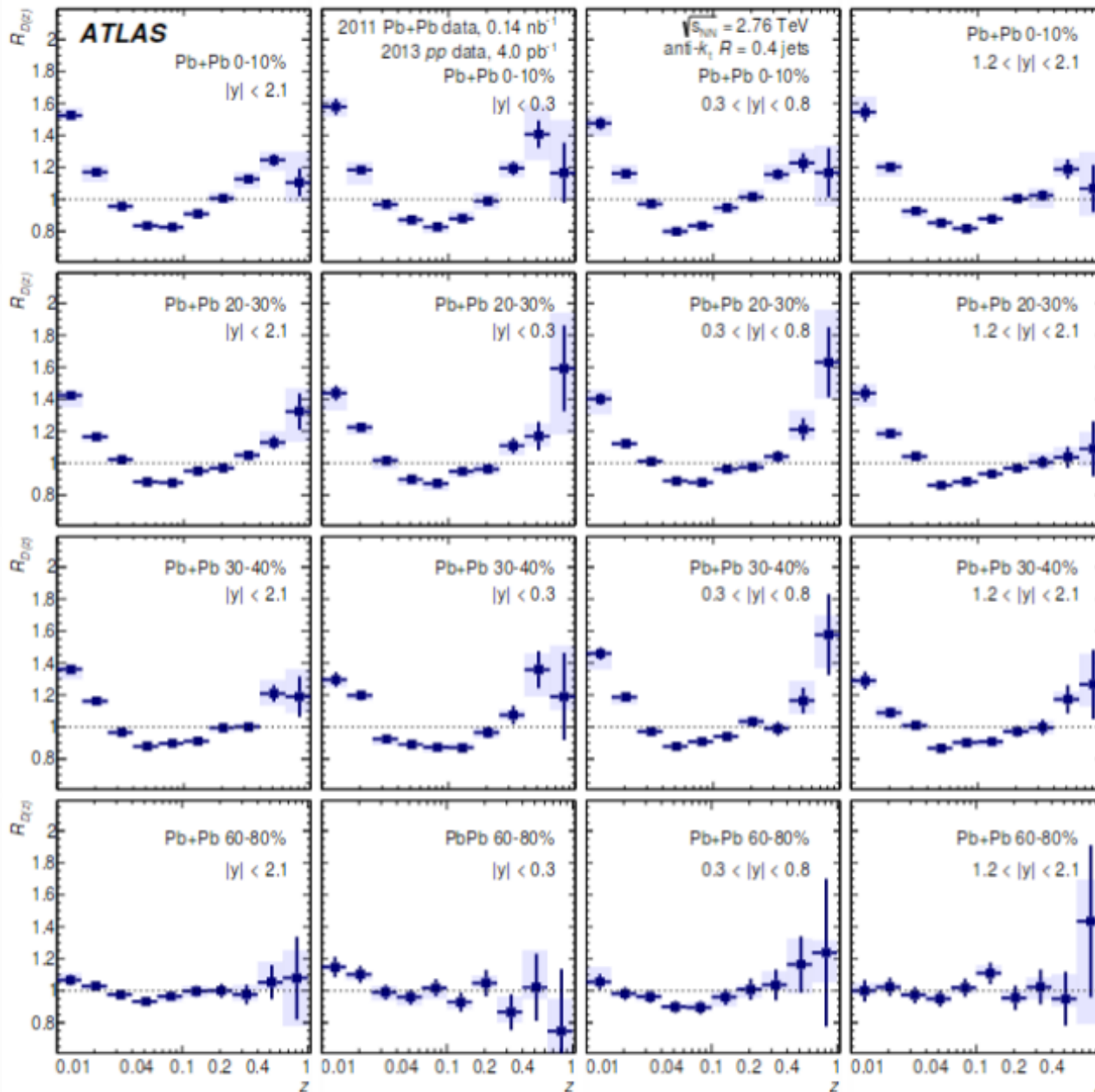
Rapidity dependence

- Hint of rapidity dependence
- Consistent with change of flavour composition as suggested in [ArXiv:1504.05169](https://arxiv.org/abs/1504.05169)

Same observation for $D(z)$

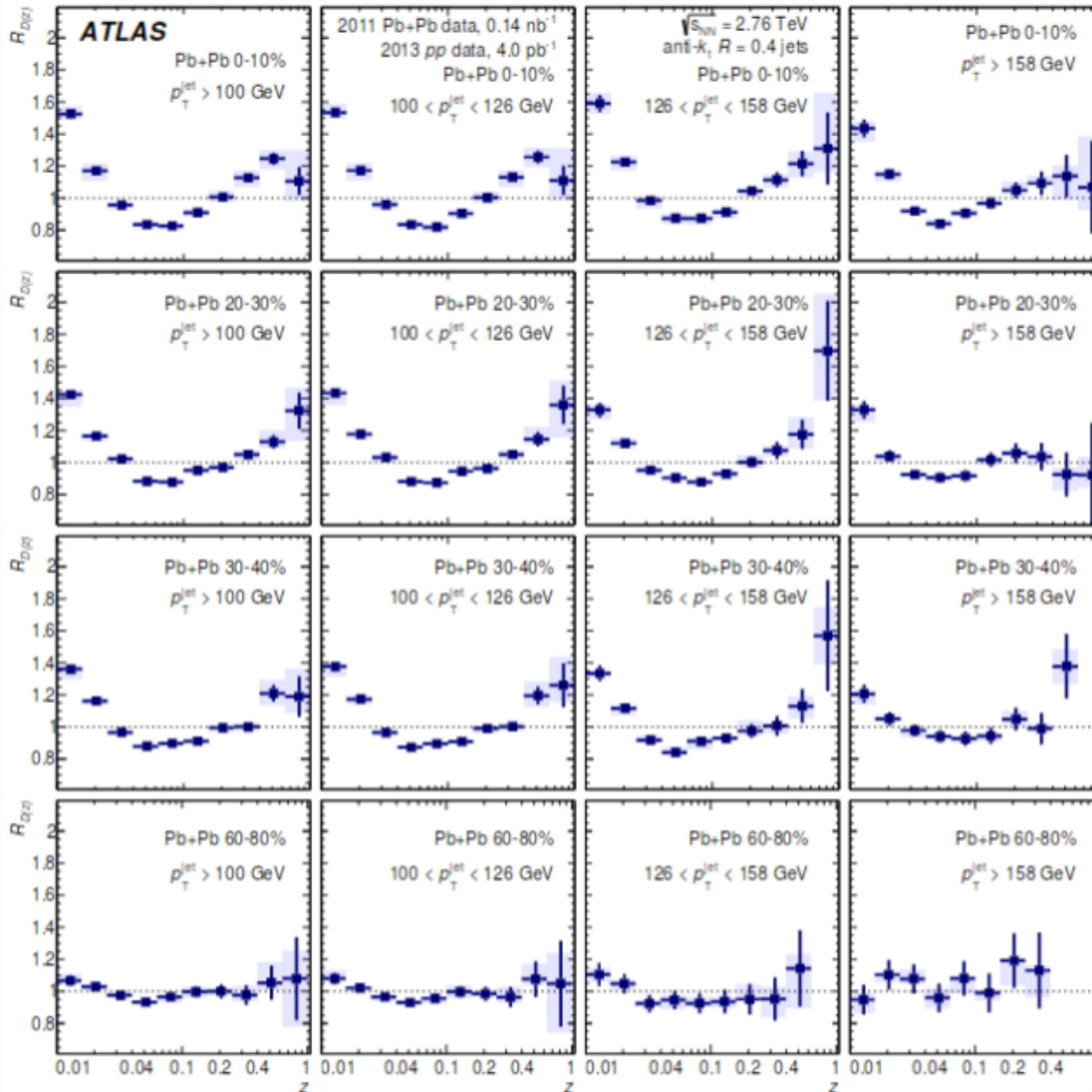


Jet fragmentation Pb+Pb @ 2.76 TeV



- Ratios of FF $D(z)$ for different centrality and rapidity bins

Jet fragmentation Pb+Pb @ 2.76 TeV

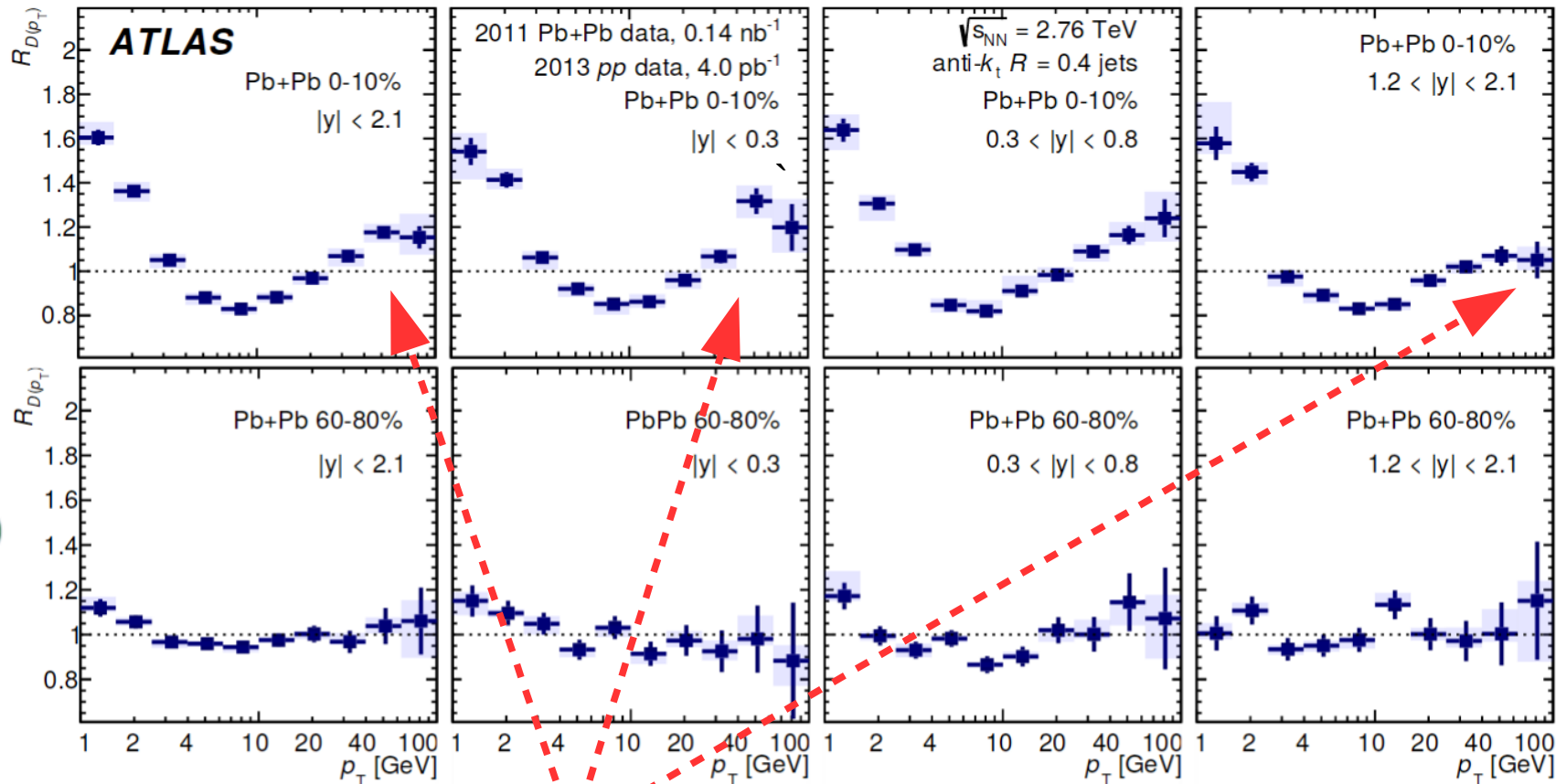


- Ratios of FF $D(z)$ for different centrality and p_T bins

Jet fragmentation ratios Pb+Pb

Ratios of $D(p_T)$ for 2 centralities and 4 *pseudorapidity* bins

@ 2.76 TeV



Rapidity dependence

- Hint of rapidity dependence
- Consistent with change of flavour composition as suggested in ArXiv:1504.05169

$100 < p_T < 300$ GeV

Same observation for $D(z)$

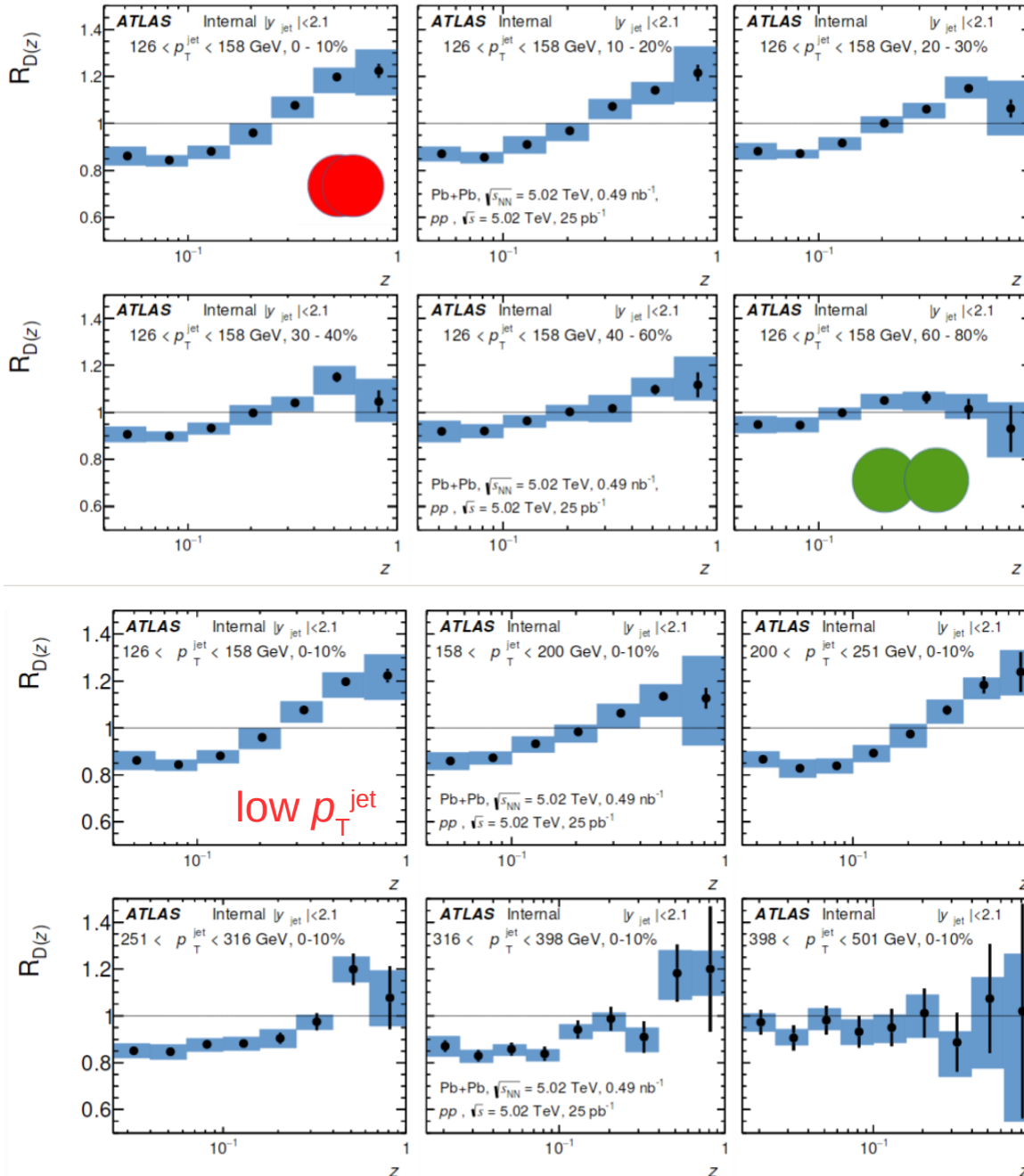
Jet fragmentation ratios Pb+Pb



Ratios of $D(z)$ for 6 centralities and 6 p_T bins

@ 5.02 TeV

ATLAS-CONF-2017-005



$126 < p_T < 501 \text{ GeV}$

- Cut on $p_T^{\text{trk}} > 4 \text{ GeV}$
- Missing the low p_T enhancement

Centrality dependence

- Enhancement at high p_T
- Suppression at intermediate p_T

Jet p_T dependence

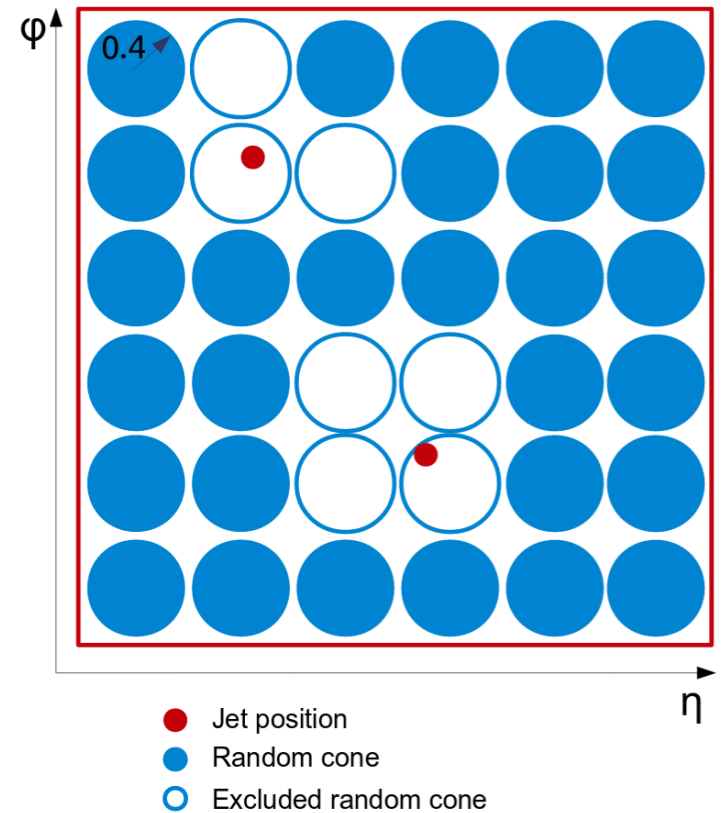
- No significant dependence on jet p_T

high p_T^{jet}

UE subtraction

$$\frac{dn_{\text{ch}}^{\text{UE}}}{dp_{\text{T}}^{\text{ch}}} = \frac{1}{N_{\text{cone}}} \frac{1}{\varepsilon} \frac{\Delta N_{\text{ch}}^{\text{cone}}(p_{\text{T}}^{\text{ch}}, p_{\text{T}}^{\text{jet}}, \eta^{\text{jet}})}{\Delta p_{\text{T}}^{\text{ch}}}$$

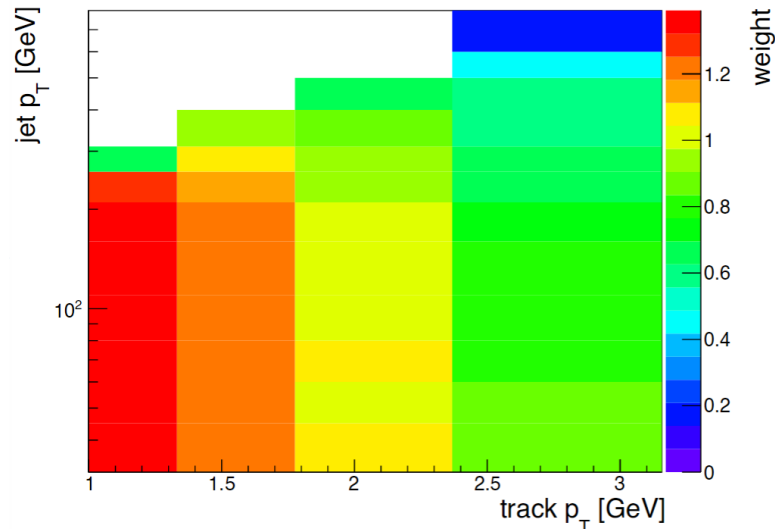
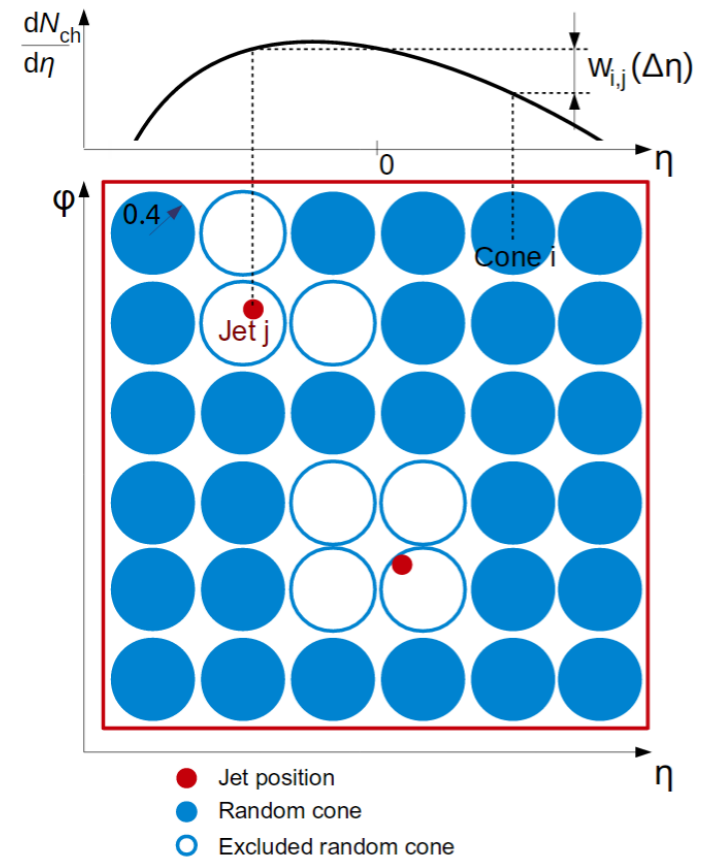
- Use grid of cones to estimate UE below a certain threshold: 3.5 GeV for p+Pb, 8 GeV for Pb+Pb
- Above those thresholds there is negligible UE, verified with overlay MC samples
- Cones with a track p_{T} bigger than certain threshold are omitted to exclude real jets



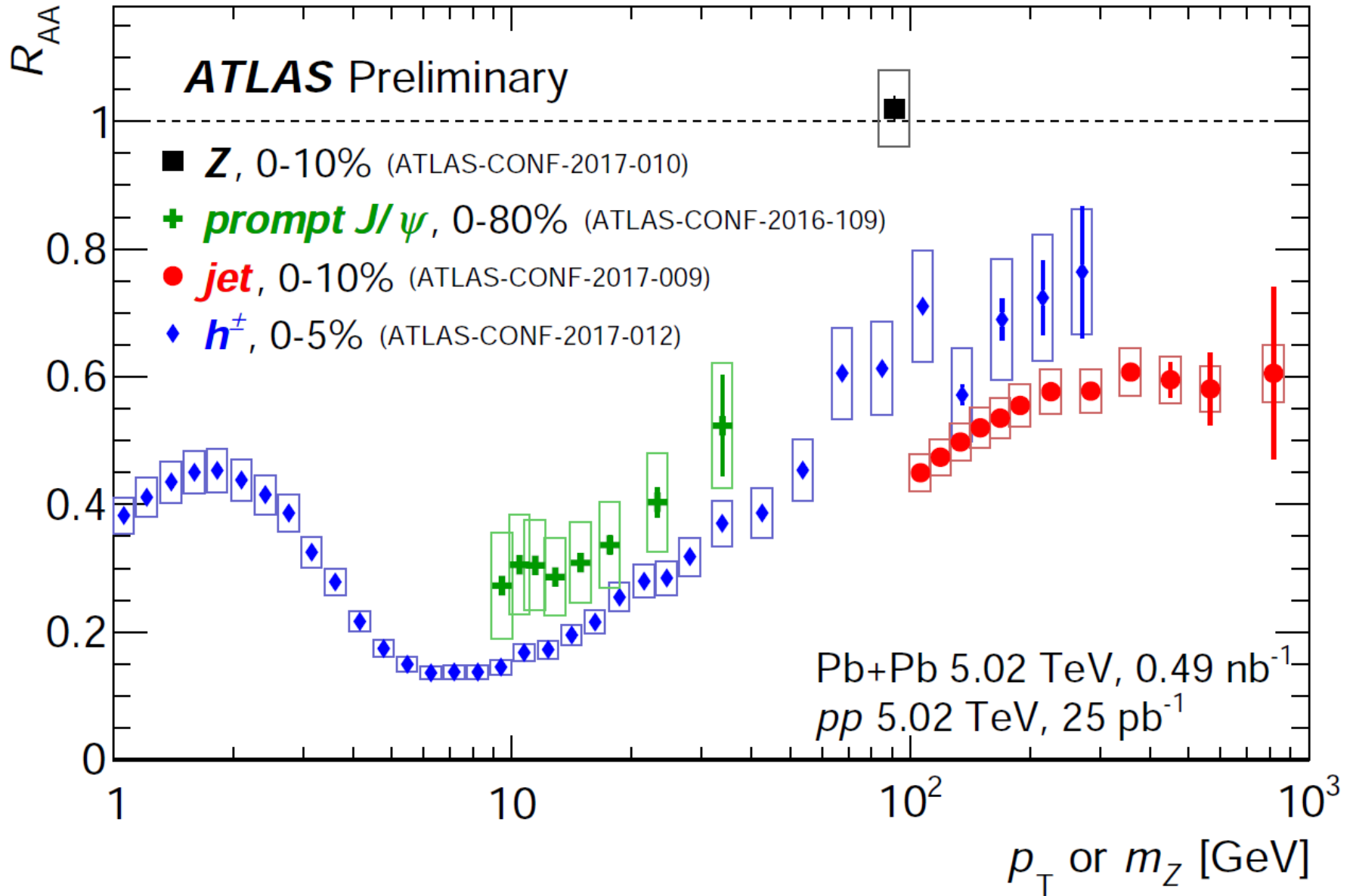
UE subtraction

- Correct UE distribution for difference between jet rapidity and cone rapidities
- Correlation between UE and JER corrected for by a multiplicative factor on the UE, $w_{UE}(z)$
 - Events with larger UE have worse JER
 - Also corrects for fakes below a certain threshold

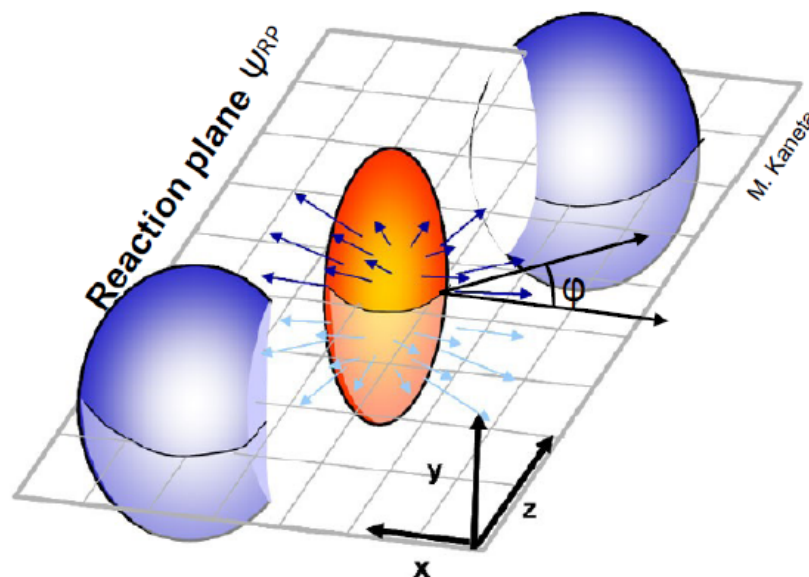
$$\frac{dn_{ch}^{UE}}{dp_T^{ch}} = \frac{1}{N_{cone}} \frac{1}{\varepsilon} \frac{\Delta N_{ch}^{cone}(p_T^{ch}, p_T^{jet}, \eta^{jet})}{\Delta p_T^{ch}}$$



$$w_{UE}(p_T) = \frac{dn_{ch}^{UE}/dp_T^{ch}}{D(p_T)_f} \Big|_{MC}$$



Eliptický tok

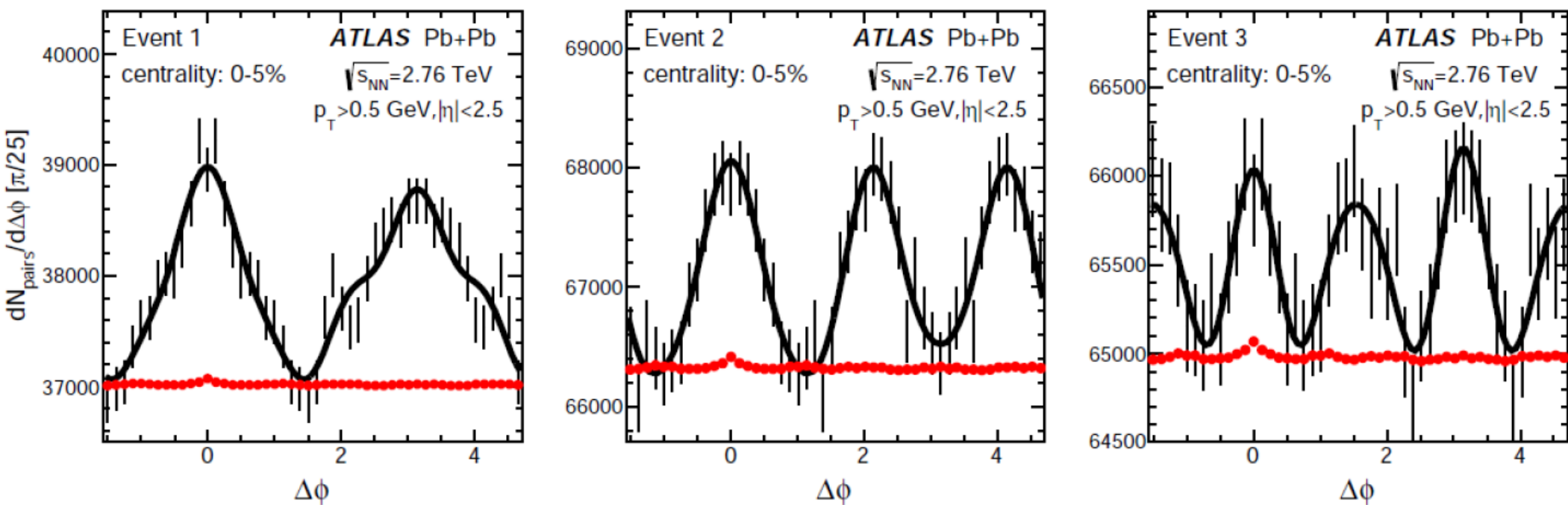


- V dekonfinovaná hmotě panují gradienty tlaku díky počáteční prostorové nesymetrii.
- Dekonfinovaná hmota se rozpíná **v různých směrech různě**.
- Tato různost lze charakterizovat například pomocí Fourierova rozvoje v úhlové vzdálenosti vzhledem k rovině interakce. Omezíme se na první netriviální člen tohoto rozvoje:

$$\frac{dN}{d\phi} = N_0 \left(1 + 2v_2 \cos 2(\phi - \Phi^{RP}) \right) \quad v_2 = \left\langle \cos 2(\phi - \Phi^{RP}) \right\rangle$$

... jeho velikost se nazývá eliptický tok (v_2)

Eliptický tok



- Tato modulace je typicky vyhodnocována statisticky (tedy průměr z vícero událostí).
- Je však často viditelná i “pouhým okem” v každé události zvlášť ... tj. velmi silný rys!
- Kdyby těžko-intové srážky byly prostou superpozicí mnoha pp srážek, tak bychom nic takového nepozorovali.