

BENCHMARKS

for EW lineshape, loops + boxes corrections in different codes

EW settings

Table 45: Different options for EW parameters setting in $\mathcal{A}^{Born+EW}$ matrix element as defined in formula (64). Indicated is the value of the form-factors. Only those fulfilling Standard Model tree-level relation $\frac{G_\mu \cdot M_Z^2 \cdot \Delta^2}{\sqrt{2} \cdot 8\pi \cdot \alpha} = 1$ can be used with default version of DYTURBO. $\Delta = \sqrt{16 \cdot s_W^2 \cdot (1 - s_W^2)}$.

**EW LO
reference**



α_0 (LO) (EWopt=1)	LEP (EWopt=3)	LEP with improved norm. (EWopt=4)
$G_\mu = 1.166389 \cdot 10^{-5}$	$G_\mu = 1.166389 \cdot 10^{-5}$	$G_\mu = 1.166389 \cdot 10^{-5}$
$M_Z = 91.1876 \text{ GeV}$	$M_Z = 91.1876 \text{ GeV}$	$M_Z = 91.1876 \text{ GeV}$
$\alpha = 1/137.03604$	$\alpha = 1/128.8667$	$\alpha = 1/128.8667$
$s_W^2 = 0.21211538$	$s_W^2 = 0.23152$	$s_W^2 = 0.23152$
$\rho_{\ell f} = 1.0$	$\rho_{\ell f} = 1.0$	$\rho_{\ell f} = 1.005$
$K_i = 1.0$	$K_i = 1.0$	$K_i = 1.0$
$\frac{G_\mu \cdot M_Z^2 \cdot \Delta^2}{\sqrt{2} \cdot 8\pi \cdot \alpha} = 1$	$\frac{G_\mu \cdot M_Z^2 \cdot \Delta^2}{\sqrt{2} \cdot 8\pi \cdot \alpha} = 1$	$\frac{G_\mu \cdot M_Z^2 \cdot \Delta^2}{\sqrt{2} \cdot 8\pi \cdot \alpha} = 1$
α_0 (Dizet) (EWopt=2)	α_0 (Dizet) + EW (loops) (EWopt=5)	α_0 (Dizet) + EW (loops + boxes) (EWopt=6)
$G_\mu = 1.166389 \cdot 10^{-5}$	$G_\mu = 1.166389 \cdot 10^{-5}$	$G_\mu = 1.166389 \cdot 10^{-5}$
$M_Z = 91.1876 \text{ GeV}$	$M_Z = 91.1876 \text{ GeV}$	$M_Z = 91.1876 \text{ GeV}$
$\alpha = 1/137.03604$	$\alpha = 1/137.03604$	$\alpha = 1/137.03604$
$s_W^2 = 0.22351946$	$s_W^2 = 0.22351946$	$s_W^2 = 0.22351946$
$\rho_{\ell f} = 1.0$	$\rho_{\ell f}(s)$	$\rho_{\ell f}(s, t)$
$\rho_{\ell f} = 1.0$	$K_i(s)$	$K_i(s, t)$
$\frac{G_\mu \cdot M_Z^2 \cdot \Delta^2}{\sqrt{2} \cdot 8\pi \cdot \alpha} \neq 1$	$\frac{G_\mu \cdot M_Z^2 \cdot \Delta^2}{\sqrt{2} \cdot 8\pi \cdot \alpha} \neq 1$	$\frac{G_\mu \cdot M_Z^2 \cdot \Delta^2}{\sqrt{2} \cdot 8\pi \cdot \alpha} \neq 1$

**EW HO
„best”**

Benchmark table: EW 1loop

Table 47: Results of calculations: masses, couplings, etc. *Tuned* comparison: predictions with EW 1-loop corrections, configuration of Dizet 6.21 shown in Table 39.

Parameter	Powheg_ew (EW 1-loop)	Dizet 6.21 (EW 1-loop)	Description
m_W	?	80.4912759	mass of W boson
s_W^2	?	0.22084586	$s_W^2 = 1 - m_W^2/m_Z^2$
c_W^2	?	0.77915414	
$\alpha_{QED}(M_Z)$?	0.007534	
$1/\alpha_{QED}(M_Z)$?	132.7316	
$\alpha_s(M_Z)$?	0.1250	
$\alpha_s(m_t)$?	0.1134	
$ZPAR(1) = \Delta r$?	0.0286451	the loop corrections to G_μ
$ZPAR(2) = \Delta r_{rem}$?	0.0	the remainder contribution $O(\alpha)$
$ZPAR(17)$?	1.04204102	
$ZPAR(18)$		1.05001972	
$ZPAR(19)$		1.04144942	
$ZPAR(20)$		1.03347072	
$ZPAR(21)$		1.04205621	
$ZPAR(22)$		1.04986350	
$ZPAR(23)$		1.04144942	
$ZPAR(24)$		1.03347072	
$ZPAR(25)$		1.04204134	
$ZPAR(26)$		1.05002004	
$ZPAR(27)$		1.04201352	
$ZPAR(28)$		1.02762415	
$ZPAR(29)$		-0.00002603	
$ZPAR(30)$		0.21220659	

Benchmark table: EW 1loop

Table 48: Common for Dizet 6.21 and MCSANC convention, real part of form-factors for $\sqrt{s} = m_Z$ and $\cos\theta = 0$. Shown are *tuned* (EW 1-loop) and best default predictions of both programs.

Form-factors	Powheg_ew (EW 1-loop)	Dizet 6.21 (EW 1-loop)	Dizet 6.21 Default (KKMC)	Comments
ρ_{eu}		1.0064962	1.005408	
K_e	?	1.0415753	1.036649	
K_u		1.0411421	1.036172	
K_{eu}		1.0827174	1.074146	
ρ_{ed}		1.0069491	1.005894	
K_e	?	1.0415753	1.036649	
K_d		1.0406299	1.035603	
K_{ed}		1.0822052	1.073556	

Benchmark table: best predictions

Table 46: Results of calculations: masses, couplings, etc. Default (best) predictions of both programs.

Parameter	Powheg_ew (EW best)	Dizet 6.21 Default (KKMC)	Description
m_W		80.353 GeV	mass of W boson
s_W^2	?	0.22351946	$s_W^2 = 1 - m_W^2/m_Z^2$
c_W^2			
$\alpha_{QED}(M_Z)$	0.007751	0.007759	
$1/\alpha_{QED}(M_Z)$	129.02086	128.88259	
$ZPAR(1) = \Delta r$?	0.03694272	the loop corrections to G_μ
$ZPAR(2) = \Delta r_{rem}$		0.01169749	the remainder contribution $O(\alpha)$
$\sin^2\theta_W^{eff}$ neutrino		0.23137443	$ZPAR(5)$
$\sin^2\theta_W^{eff}$ electron		0.23175590	$ZPAR(6)$
$\sin^2\theta_W^{eff}$ muon		0.23175590	$ZPAR(7)$
$\sin^2\theta_W^{eff}$ tau		0.23175590	$ZPAR(8)$
$\sin^2\theta_W^{eff}$ u-quark	?	0.23164930	$ZPAR(9)$
$\sin^2\theta_W^{eff}$ d-quark		0.23152214	$ZPAR(10)$
$\sin^2\theta_W^{eff}$ c-quark		0.23164930	$ZPAR(11)$
$\sin^2\theta_W^{eff}$ s-quark		0.23152214	$ZPAR(12)$
$\sin^2\theta_W^{eff}$ b-quark		0.23307588	$ZPAR(14)$

Lineshape: ratio to EW LO

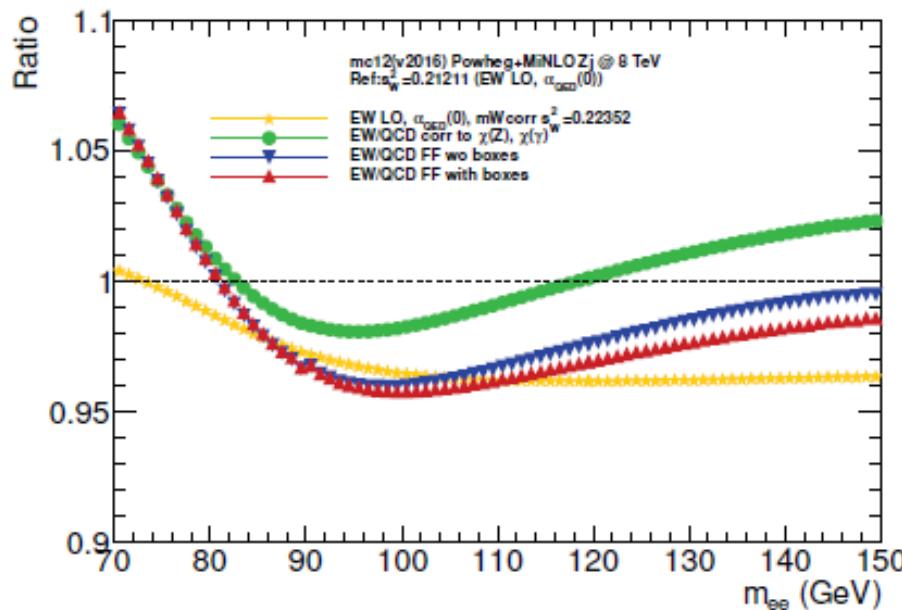


Table 49: Normalisation corrections to the total cross-section. Shown are individual normalisation factors vs predicted with EW LO $\alpha(0)$ scheme EWopt = 1.

Version	Corrections to cross-section	$m_{ee} = 89 - 93$ GeV	$m_{ee} = 80 - 100$ GeV	Comments	Comments
$EW_{\text{opt}} = 1$	EW $\alpha(0)$	1.0	1.0	reference	$s_W^2 = 0.21211$
$EW_{\text{opt}} = 2$	EW $\alpha(0) + m_W \text{corr}$	0.97102	0.98245		$s_W^2 = 0.22352$
	EW $\alpha(0) + m_W \text{corr} + \text{FF (prop.)}$	0.98245	0.98344		$s_W^2 = 0.22352$
$EW_{\text{opt}} = 5$	EW $\alpha(0) + m_W \text{corr} + \text{FF (loops)}$	0.96526	0.96645		$s_W^2 = 0.22352$
$EW_{\text{opt}} = 6$	EW $\alpha(0) + m_W \text{corr} + \text{FF (loops + boxes)}$	0.96526	0.96648		$s_W^2 = 0.22352$

SPARES

From Zfitter/Dizet documentation

D. Bardin et al., Comp. Phys. Comun. 133 (2001) 229-395

Zfitter is a **semi-analytical program** for calculating total cross-sections and pseudo-observables (eg. A_{fb} , $\sin^2\theta_W^{eff}$), used by LEP1, and to a lesser degree by LEP2.

DIZET is a library for calculating form-factors and some other corrections. Provides complete EW $O(\alpha)$ weak-loop corrections supplemented with selected higher order terms (eg. vacuum polarisation, $\alpha_{QED}(Q^2)$).

For analyses at LEP1, LEP2 used always in parallel with **MC generators (KoralZ, KoralW)** eg. to evaluate systematics of simplified cuts used in analysis integration.

$$\mathcal{A}_Z^{OLA}(s, t) = i\sqrt{2}G_\mu I_e^{(3)} I_f^{(3)} M_Z^2 \chi_Z(s) \rho_{ef}(s, t) \left\{ \gamma_\mu(1 + \gamma_5) \otimes \gamma_\mu(1 + \gamma_5) \right. \\ - 4|Q_e| s_W^2 \kappa_e(s, t) \gamma_\mu \otimes \gamma_\mu(1 + \gamma_5) - 4|Q_f| s_W^2 \kappa_f(s, t) \gamma_\mu(1 + \gamma_5) \otimes \gamma_\mu \\ \left. + 16|Q_e Q_f| s_W^4 \kappa_{e,f}(s, t) \gamma_\mu \otimes \gamma_\mu \right\}. \quad (A.4.75)$$

one loop amplitude

$$A_\gamma^{OLA} = i\chi_\gamma(s) \alpha(s) \gamma_\mu \otimes \gamma_\mu. \quad (2.2.36)$$

Dyson summation leads to the change of α into $\alpha(s)$:

$$\alpha(s) = \frac{\alpha(0)}{1 - \Delta\alpha^{fer}(s)} = \frac{\alpha(0)}{1 - \boxed{\Delta\alpha^{(5)}(s) - \Delta\alpha^t(s) - \Delta\alpha^{\alpha\alpha_s}(s)}}, \quad (2.2.37)$$

Vacuum polarisation corrections

From Zfitter/Dizet documentation

D. Bardin et al., Comp. Phys. Comun. 133 (2001) 229-395

After some trivial algebra one derives the final expressions:

$$\begin{aligned} \rho_{ef} &= 1 + \frac{g^2}{16\pi^2} \left\{ -\Delta\rho_z^F + D_z^F(s) + \frac{5}{3}B_0^F(-s; M_w, M_w) - \frac{9}{4}\frac{c_w^2}{s_w^2} \ln c_w^2 - 6 \right. \\ &\quad + \frac{5}{8}c_w^2(1+c_w^2) + \frac{1}{4c_w^2}(3v_e^2 + a_e^2 + 3v_f^2 + a_f^2)\mathcal{F}_z(s) + \hat{\mathcal{F}}_w^0(s) + \hat{\mathcal{F}}_w(s) \\ &\quad \left. - \frac{r_t}{4}[B_0^F(-s; M_w, M_w) + 1] - c_w^2(R_z - 1)s\hat{B}_{ww}^d(s, t) \right\}, \end{aligned} \quad (\text{A.4.80})$$

$$\begin{aligned} \kappa_e &= 1 + \frac{g^2}{16\pi^2} \left\{ -\frac{c_w^2}{s_w^2} \Delta\rho^F - \Pi_{Z\gamma}^F(s) - \frac{1}{6}B_0^F(-s; M_w, M_w) - \frac{1}{9} - \frac{v_e\sigma_e}{2c_w^2}\mathcal{F}_z(s) \right. \\ &\quad - \hat{\mathcal{F}}_w^0(s) + (R_z - 1) \left[\frac{|Q_f|}{2} (1 - 4|Q_f|s_w^2) \mathcal{F}_z(s) + c_w^2 [\hat{\mathcal{F}}_{wn}(s) \right. \\ &\quad \left. \left. - |Q_f| \mathcal{F}_{wa}(s) + s\hat{B}_{ww}^d(s, t)] \right] \right\}, \end{aligned} \quad (\text{A.4.81})$$

$$\begin{aligned} \kappa_f &= 1 + \frac{g^2}{16\pi^2} \left\{ -\frac{c_w^2}{s_w^2} \Delta\rho^F - \Pi_{Z\gamma}^F(s) - \frac{1}{6}B_0^F(-s; M_w, M_w) - \frac{1}{9} - \frac{v_f\sigma_f}{2c_w^2}\mathcal{F}_z(s) \right. \\ &\quad \left. - \hat{\mathcal{F}}_w(s) + (R_z - 1) \left[\frac{|Q_e|}{2} (1 - 4|Q_e|s_w^2) \mathcal{F}_z(s) + c_w^2 [\hat{\mathcal{F}}_{wn}^0(s) \right. \right. \\ &\quad \left. \left. - |Q_e| \mathcal{F}_{wa}(s) + s\hat{B}_{ww}^d(s, t)] \right] \right\}, \end{aligned}$$

$$\text{Interference} \quad - \frac{r_t}{4}[B_0^F(-s; M_w, M_w) + 1], \quad (\text{A.4.82})$$

$$\begin{aligned} \kappa_{ef} &= 1 + \frac{g^2}{16\pi^2} \left\{ -2\frac{c_w^2}{s_w^2} \Delta\rho^F - 2\Pi_{Z\gamma}^F(s) - \frac{1}{3}B_0^F(-s; M_w, M_w) - \frac{2}{9} \right. \\ &\quad - \frac{1}{4c_w^2} \left[\frac{\delta_e^2 + \delta_f^2}{s_w^2} (R_w - 1) + 3v_e^2 + a_e^2 + 3v_f^2 + a_f^2 \right] \mathcal{F}_z(s) \\ &\quad - \hat{\mathcal{F}}_w^0(s) - \hat{\mathcal{F}}_w(s) - \frac{r_t}{4}[B_0^F(-s; M_w, M_w) + 1] \\ &\quad + c_w^2(R_z - 1) \left[\frac{2}{3} - \hat{\Pi}_{\gamma\gamma}^{\text{bos}, F}(s) + s\hat{B}_{ww}^d(s, t) \right] \Big\}. \end{aligned} \quad (\text{A.4.83})$$

E. Richter-Was, IFJU
Fermionic loops in γ^* propagator

EWPrecisionWS, Paris, 24.05.2018

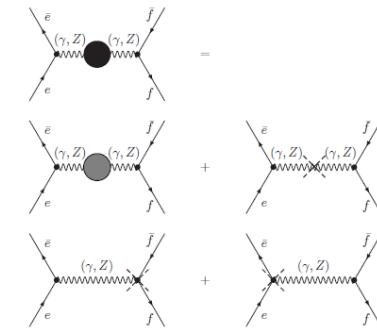


Figure A.11. Bosonic self-energies and bosonic counter-terms for $e\bar{e} \rightarrow (Z, \gamma) \rightarrow f\bar{f}$



Figure A.10. Electron (a) and final fermion (b) vertices in $e\bar{e} \rightarrow (Z) \rightarrow f\bar{f}$

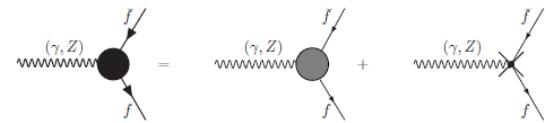


Figure A.6. Off-shell $Z f\bar{f}$ and $\gamma f\bar{f}$ vertices

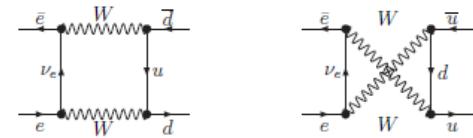


Figure A.7. The WW boxes

etc. etc.

EW $\alpha(0)$ scheme (Dizet)

1) Start with tree parameters at their best measured values :

$$G_\mu, \alpha(0), M_Z$$

2) Calculate m_W using iterative formula including higher-order corrections

Parameter	Value
M_Z	91.1876 GeV
m_h	125.0 GeV
m_t	173.0 GeV
m_b	4.7 GeV
$1/\alpha(0)$	137.0359895(61)
G_μ	$1.166389(22) \cdot 10^{-5} \text{ GeV}^{-2}$

$$M_W = \frac{M_Z}{\sqrt{2}} \sqrt{1 + \sqrt{1 - \frac{4A_0^2}{M_Z^2(1 - \Delta r)}}}$$

$$A_0 = \sqrt{\frac{\pi\alpha(0)}{\sqrt{2}G_\mu}}$$

$$s_W^2 = 1 - M_W^2/M_Z^2$$



On-mass-shell relation
valid to all orders

M_W	80.353 GeV
s_W^2	0.22351946

$\alpha(0)$ important to consistently separate QED corrections from genuine EW ones. No need to introduce two α scales.

Spin amplitude: EW Improved Born (IBA)

$$\mathcal{A}^{Born+EW} = \frac{\alpha}{s} \{ [\bar{u} \gamma^\mu v g_{\mu\nu} \bar{v} \gamma^\nu u] \cdot (q_\ell \cdot q_f) \cdot \boxed{\Gamma_{V\Pi}} \cdot \chi_\gamma(s) \\ + [\bar{u} \gamma^\mu v g_{\mu\nu} \bar{v} \gamma^\nu u \cdot (v_\ell \cdot v_f \cdot v v_{\ell f}) + \bar{u} \gamma^\mu v g_{\mu\nu} \bar{v} \gamma^\nu \gamma^5 u \cdot (v_\ell \cdot a_f) \\ + \bar{u} \gamma^\mu \gamma^5 v g_{\mu\nu} \bar{v} \gamma^\nu u \cdot (a_\ell \cdot v_f) + \bar{u} \gamma^\mu \gamma^5 v g_{\mu\nu} \bar{v} \gamma^\nu \gamma^5 u \cdot (a_\ell \cdot a_f)] \cdot \boxed{Z_{V\Pi}} \cdot \chi_Z(s) \}$$

$$\chi_\gamma(s) = 1$$

$$\chi_Z(s) = \frac{G_\mu \cdot M_z^2 \cdot \Delta^2}{\sqrt{2} \cdot 8\pi \cdot \alpha} \cdot \frac{s}{s - M_Z^2 + i \cdot \Gamma_Z \cdot M_Z}$$

$$\boxed{Z_{V\Pi}} = \rho_{e,f}(s,t)$$

$$\boxed{\Gamma_{V\Pi}} = \frac{1}{2 - (1 + \boxed{\Pi_{\gamma\gamma}(s)})}$$

$$v_\ell = (2 \cdot T_3^\ell - 4 \cdot q_\ell \cdot s_W^2 \cdot \boxed{K_\ell(s,t)}) / \Delta$$

$$v_f = (2 \cdot T_3^f - 4 \cdot q_f \cdot s_W^2 \cdot \boxed{K_f(s,t)}) / \Delta$$

$$a_\ell = (2 \cdot T_3^\ell) / \Delta$$

$$a_f = (2 \cdot T_3^f) / \Delta$$

$$\Delta = \sqrt{16 \cdot s_W^2 \cdot (1 - s_W^2)}$$

EW form-factors , functions of $(s,t)=(m_{||}, \cos\theta)$
Calculated with Dizet 6.21 library

$$v v_{\ell f} = \frac{1}{v_\ell \cdot v_f} [(2 \cdot T_3^\ell)(2 \cdot T_3^f) - 4 \cdot q_\ell \cdot s_W^2 \cdot \boxed{K_f(s,t)}[2 \cdot T_3^\ell] - 4 \cdot q_f \cdot s_W^2 \cdot \boxed{K_\ell(s,t)}[2 \cdot T_3^f)]$$

$$+ (4 \cdot q_\ell \cdot s_W^2)(4 \cdot q_f \cdot s_W^2) \boxed{K_{\ell f}(s,t)} \frac{1}{\Delta^2}$$

Vacuum polarisation corrections, used low-energy experiment input.

Warning: problem for analytic continuation

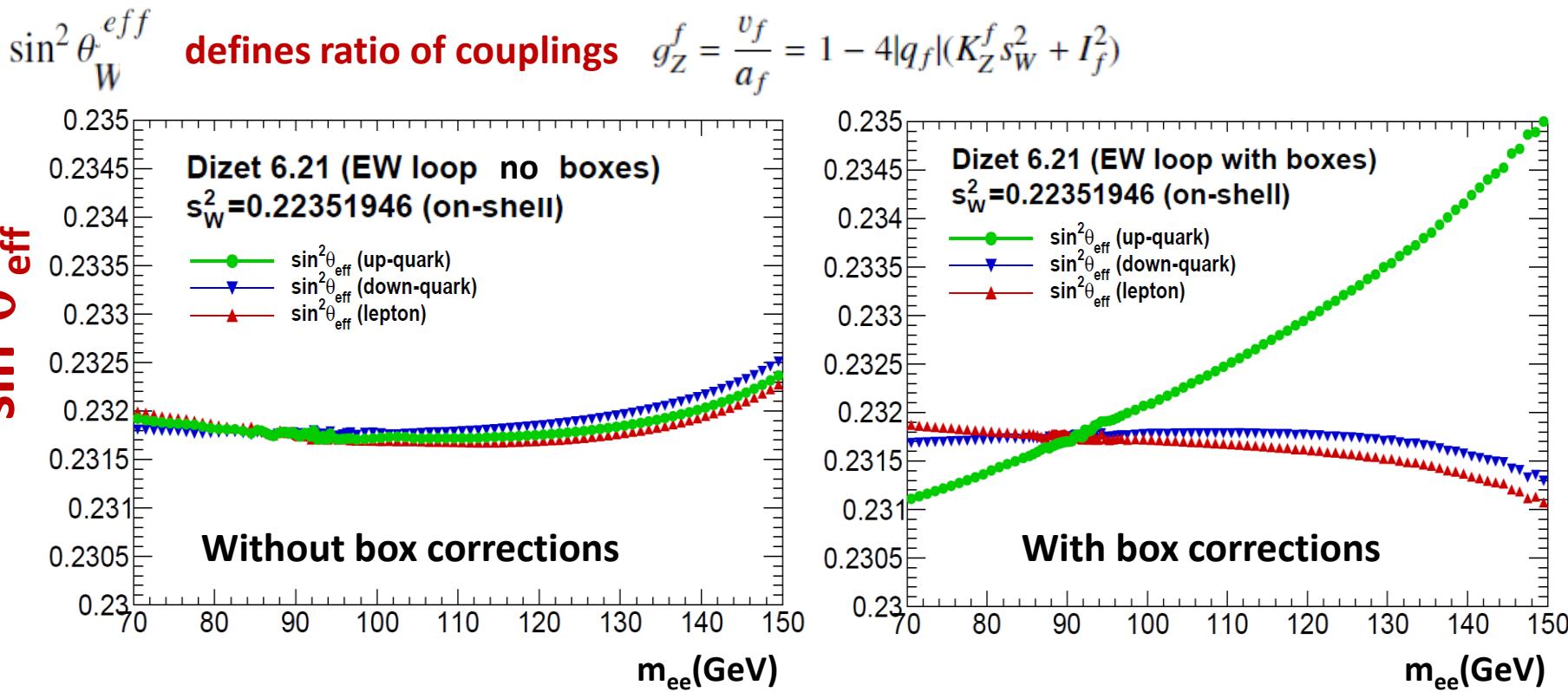
Effective weak mixing angle

Here convoluted with line-shape and $\cos\theta^*$ distribution of MC events.

$$\sin^2 \theta_W^{eff} = Re(K_Z^f) s_W^2 + I_f^2$$

$$I_f^2 = \alpha^2(s) \frac{35}{18} [1 - \frac{8}{3} Re(K_Z^f) s_W^2] = \sim 10^{-4}$$

$$s_W^2 = 1 - M_W^2/M_Z^2$$



De-mystifying EW schemes

EW $\alpha(0)$ scheme:

Input: $\alpha(0)$, M_Z , G_μ

$$M_Z = 91.1876 \text{ GeV}$$

$$\alpha = 1/137.03599$$

$$G_\mu = 1.1663787 \cdot 10^{-5} \text{ GeV}^{-2}$$

$$M_W = 80.353 \text{ GeV}$$

$$s_W^2 = 0.22351946$$

EW G_μ scheme:

Input: M_Z , M_W , G_μ

$$M_Z = 91.1876 \text{ GeV}$$

$$M_W = 80.385 \text{ GeV}$$

$$G_\mu = 1.1663787 \cdot 10^{-5} \text{ GeV}^{-2}$$

$$\alpha = 1/132.2332$$

$$s_W^2 = 0.222897$$

EW $\alpha(M_Z)$ scheme:

Input: $\alpha(M_Z)$, M_Z , G_μ

$$M_Z = 91.1876 \text{ GeV}$$

$$\alpha = 1/128.8793$$

$$G_\mu = 1.1663787 \cdot 10^{-5} \text{ GeV}^{-2}$$

$$M_W = 79.9968 \text{ GeV}$$

$$s_W^2 = 0.2311538$$

$$M_W = \frac{M_Z}{\sqrt{2}} \sqrt{1 + \sqrt{1 - \frac{4A_0^2}{M_Z^2(1 - \Delta r)}}}$$

$$A_0 = \sqrt{\frac{\pi\alpha(0)}{\sqrt{2}G_\mu}}$$

$$s_W^2 = 1 - M_W^2/M_Z^2$$

$$\alpha = G_\mu \sqrt{2} M_W^2 (1 - M_W^2/M_Z^2)/\pi$$

$$s_W^2 = 1 - M_W^2/M_Z^2$$

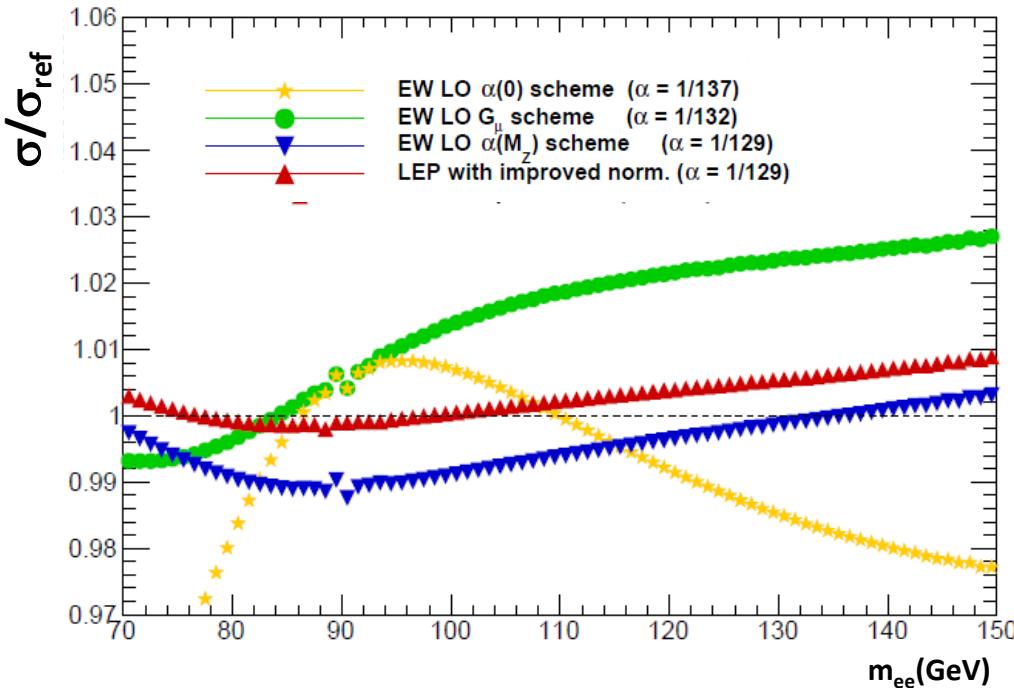
**Most commonly used scheme
in MC generators todays.**

$$s_W^2 = \frac{1 - \sqrt{1 - 4 \frac{\alpha\pi}{G_\mu \sqrt{2}M_Z^2}}}{2}$$

$$M_W = M_Z \cdot \sqrt{1 - s_W^2}$$

Line-shape in different EW schemes

Ref.: best predictions with EW $O(\alpha)$ loop+boxes corrections



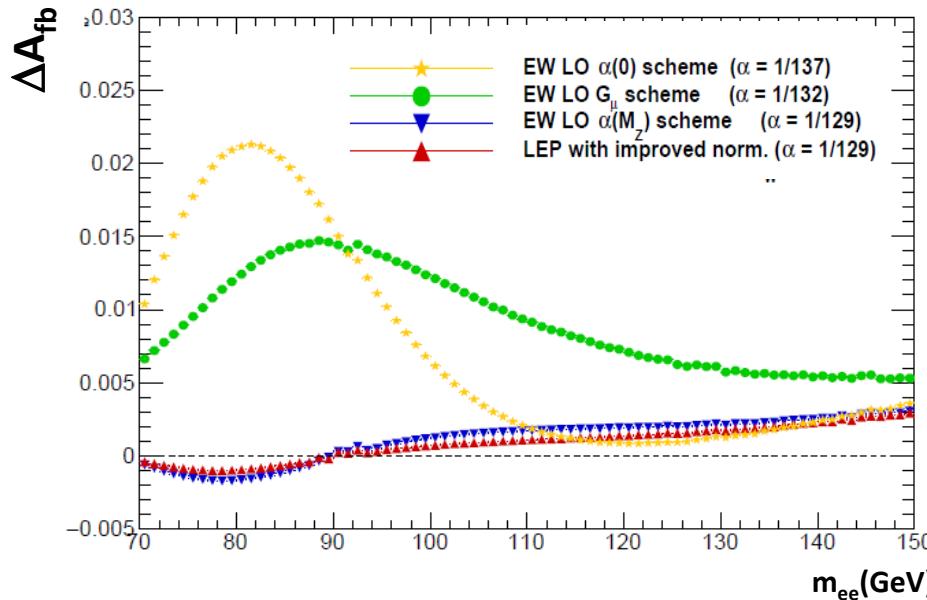
„LEP with improved norm.”
(red) predicts very well
 shape and normalisation

EW LO G_μ (green) requires
 0.6% corrections at Z-pole
 and below , 1.7% above

σ/σ_{ref}	$m_{ee} = 70\text{-}80 \text{ GeV}$	$m_{ee} = 80\text{-}100 \text{ GeV}$	$m_{ee} = 100\text{-}125 \text{ GeV}$
EW LO $\alpha(0)$ scheme	0.9641	1.0052	1.0019
EW LO G_μ scheme	0.9942	1.0061	1.0173
EW LO $\alpha(M_Z)$ scheme	0.9934	0.9891	0.9930
LEP with improv. norm.	1.0007	0.9990	1.0017

A_{fb} in different EW schemes

Reference: best predictions with EW $O(\alpha)$ loop+boxes corrections.



„LEP with improved norm.”
(red) predicts remarkably well
 shape and normalisation

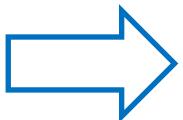
EW LO G_μ (green) requires
 large corrections, above 50%
 around Z-pole!

$A_{fb} - A_{fb,ref}$	$m_{ee} = 70\text{-}80 \text{ GeV}$ $A_{fb} = -0.1990$	$m_{ee} = 80\text{-}100 \text{ GeV}$ $A_{fb} = 0.0262$	$m_{ee} = 100\text{-}125 \text{ GeV}$ $A_{fb} = 0.1928$
EW LO $\alpha(0)$ scheme	0.0175	0.0146	0.0033
EW LO G_μ scheme	0.0096	0.0143	0.0102
EW LO $\alpha(M_Z)$ scheme	0.0015	0.0002	0.0016
LEP with improv. norm.	-0.0009	0.0001	0.0010

EW $\alpha(0)$ scheme

Tree level inputs and predictions

Often used for configuring
MC generators at LHC



EW scheme input	EW LO predictions
$M_Z = 91.1876 \text{ GeV}$ $\alpha = 1/137.0359895$ $G_\mu = 1.166389 \cdot 10^{-5} \text{ GeV}^{-2}$	$M_W = 80.93875 \text{ GeV}$ $s_W^2 = 0.2121538$
$M_Z = 91.1876 \text{ GeV}$ $\alpha = 1/128.8667$ $s_W^2 = 0.23152$	$M_W = 79.93777 \text{ GeV}$ $G_\mu = 1.165224 \cdot 10^{-5} \text{ GeV}^{-2}$
$M_Z = 91.1876 \text{ GeV}$ $\alpha = 1/127.934$ $s_W^2 = 0.23113$	$M_W = 79.958 \text{ GeV}$ $G_\mu = 1.175057 \cdot 10^{-5} \text{ GeV}^{-2}$

Input for DIZET (beyond tree level)

- 1) Starts with tree parameters at best measured value : $G_\mu, \alpha(0), M_Z$
- 2) Calculate m_W using iterative formula including higher-order corrections

Parameter	Value
M_Z	91.1876 GeV
m_h	125.0 GeV
m_t	173.0 GeV
m_b	4.7 GeV
$1/\alpha(0)$	137.0359895(61)
G_μ	$1.166389(22) \cdot 10^{-5} \text{ GeV}^{-2}$
M_W	80.353 GeV
s_W^2	0.22351946