



Comparison of QED-FSR tools Horace, Photos, Pythia-QED

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Few comparisons of QED-FSR tools available in the literature

- HORACE-old vs WINHAC [hep-ph/0402235](#)
- PHOTOS vs HORACE-old [arXiv:1510.02458](#)
- PYTHIA-QED vs PHOTOS (vs HORACE-old) [arXiv:1612.02841](#)

Disclaimer: I refer to the papers for the details about the individual plots

codes evolved from 2004 to 2016 → non-uniform comparison

Comparison of tools simulating QED-FSR to all orders: generalities

All QED tool describing FSR to all orders share some general features.

- The inclusion to all orders of radiation in the soft limit allows to recover the “classical” limit, i.e. the vanishing probability of observing a bare, naked, charged particle.
This behaviour is due to the exponential suppression of the Sudakov form factor.
- The Kinoshita-Lee-Nauenberg theorem states the cancellation of mass (collinear) singularities for observables inclusive over radiation.
The sensitivity to these large logarithms remains for exclusive observables (e.g. bare leptons)
- Comparison of tools according to their:
 - fixed-order accuracy
 - logarithmic accuracy
- Differences arise in higher-orders because of:
 - different resummation formalism → different inclusion of subleading terms
 - algorithmic implementations
- Matching fixed NLO-EW and all-orders QED-PS results pushes the differences to NNLO-EW reducing some discrepancies present before matching

The comparison of tools provides hints about the size of higher-order terms → of missing h.o.

Classification of EW radiative corrections

Fermion masses regulate collinear emission \rightarrow mass logarithm $L \equiv \log\left(\frac{s}{m_l^2}\right)$
 Final state lepton masses are physical parameters

Perturbative expansion in α

at each order, classification w.r.t. powers of L

$$\sigma = \underbrace{a_0 + a_1 \alpha L + b_1 \alpha}_{\substack{\text{Photos} \\ \text{HORACE-matched} \\ \text{DKM} \\ \text{Winhac}}} + \underbrace{a_2 \alpha^2 L^2 + b_2 \alpha^2 L + c_2 \alpha^2}_{\substack{\text{WGRAD} \\ \text{DK} \\ \text{HORACE-matched} \\ \text{SANC}}} + \underbrace{a_3 \alpha^3 L^3 + b_3 \alpha^3 L^2 + c_3 \alpha^3 L + d_3 \alpha^3}_{\text{various codes include different subsets}} + \dots$$

numerical simulation of final state QED multiple photon emission via Parton Shower (Photos, HORACE)
 matching of NLO-EW results with complete QED Parton Shower (HORACE)

different FSR models handle in a different way the columns of this expansion

HORACE vs WINHAC comparison: hep-ph/0402235

C.M. Carloni Calame, S. Jadach, G. Montagna, O. Nicrosini, W. Płaczek

- HORACE and WINHAC are two Monte Carlo event generators
- HORACE is run in its “old” mode, i.e. including only QED-FSR via Parton Shower

LL accuracy,
 energy extracted according to $P_{q\gamma}(z)$,
 angles extracted according to $1/(1-\beta\cos\theta)$
 in 2004 muon mass neglected, electron mass kept

effects studied at

- $O(\alpha)$ (truncated PS)
- to all orders (best)

- WINHAC includes FSR via $O(\alpha)$ YFS exponentiation

exact to all orders in the IR limit

up to $O(\alpha)$ for the non IR contributions

non-IR photonic terms beyond $O(\alpha)$ are approximated

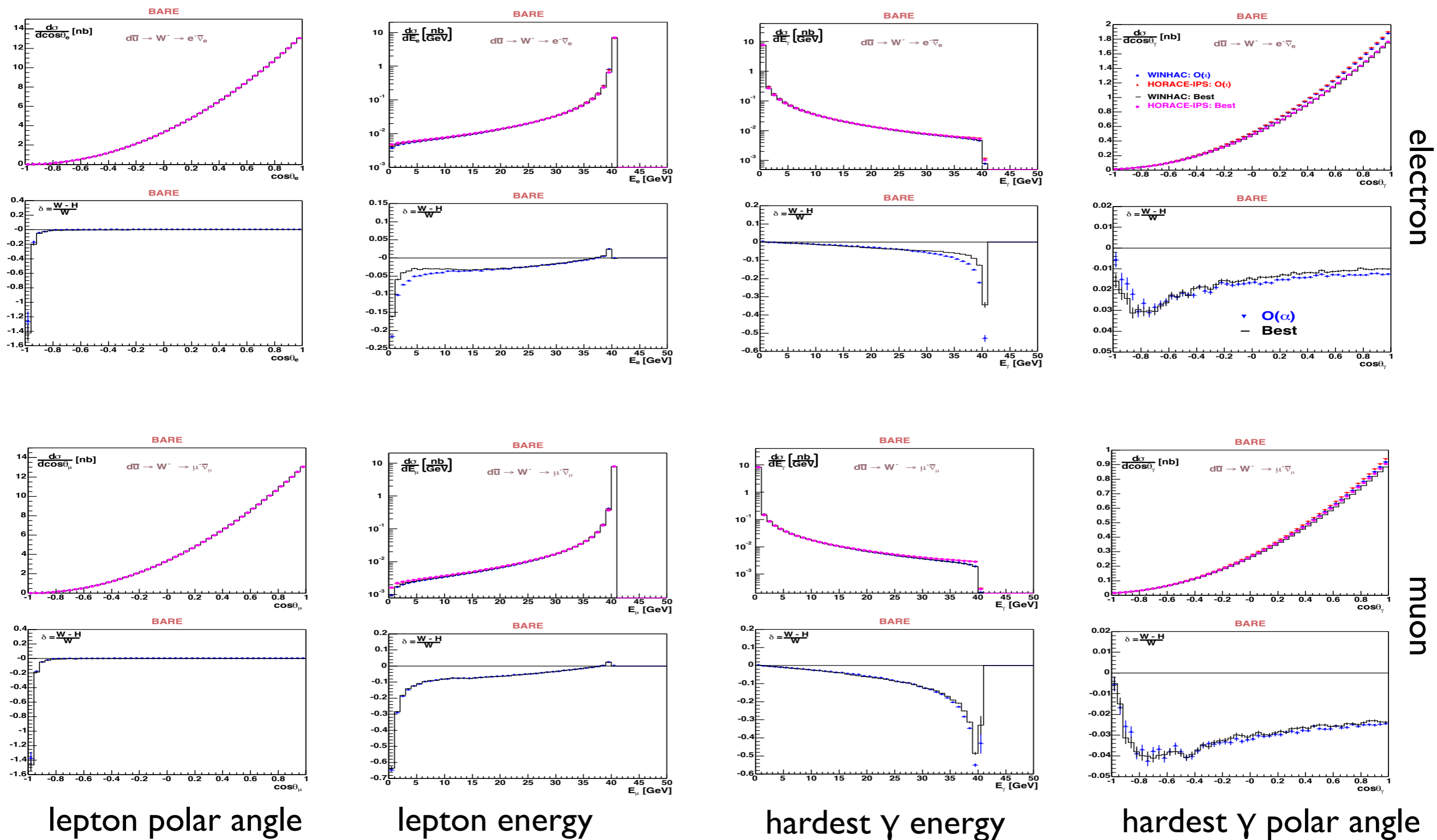
subleading logarithmic terms included in the YFS scheme

Program	σ^{tot} [nb]: WITH CUTS		
	Born	$O(\alpha)$	Best
$W^- \rightarrow e^- \bar{\nu}_e$			
HORACE	3.23633 (12)	3.18707 (13)	3.18696 (13)
WINHAC	3.23629 (09)	3.18779 (07)	3.18765 (06)
$\delta = (W - H)/W$	$-1.2 (4.6) \times 10^{-5}$	$2.3 (0.5) \times 10^{-4}$	$2.2 (0.5) \times 10^{-4}$
$W^- \rightarrow \mu^- \bar{\nu}_\mu$			
HORACE	3.23632 (12)	3.15990 (12)	3.16013 (13)
WINHAC	3.23630 (07)	3.16418 (06)	3.16409 (05)
$\delta = (W - H)/W$	$-0.6 (4.3) \times 10^{-5}$	$1.35 (0.05) \times 10^{-3}$	$1.25 (0.05) \times 10^{-3}$
$W^+ \rightarrow e^+ \nu_e$			
HORACE	4.39341 (16)	4.32186 (17)	4.32187 (18)
WINHAC	4.39328 (13)	4.32286 (10)	4.32273 (08)
$\delta = (W - H)/W$	$-3.0 (4.7) \times 10^{-5}$	$2.3 (0.5) \times 10^{-4}$	$2.0 (0.5) \times 10^{-4}$
$W^+ \rightarrow \mu^+ \nu_\mu$			
HORACE	4.39340 (16)	4.28255 (16)	4.28326 (16)
WINHAC	4.39336 (10)	4.28837 (08)	4.28848 (08)
$\delta = (W - H)/W$	$-0.9 (4.3) \times 10^{-5}$	$1.36 (0.05) \times 10^{-3}$	$1.22 (0.05) \times 10^{-3}$

HORACE vs WINHAC comparison: hep-ph/0402235

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parton-level comparison



electron

muon

lepton polar angle

lepton energy

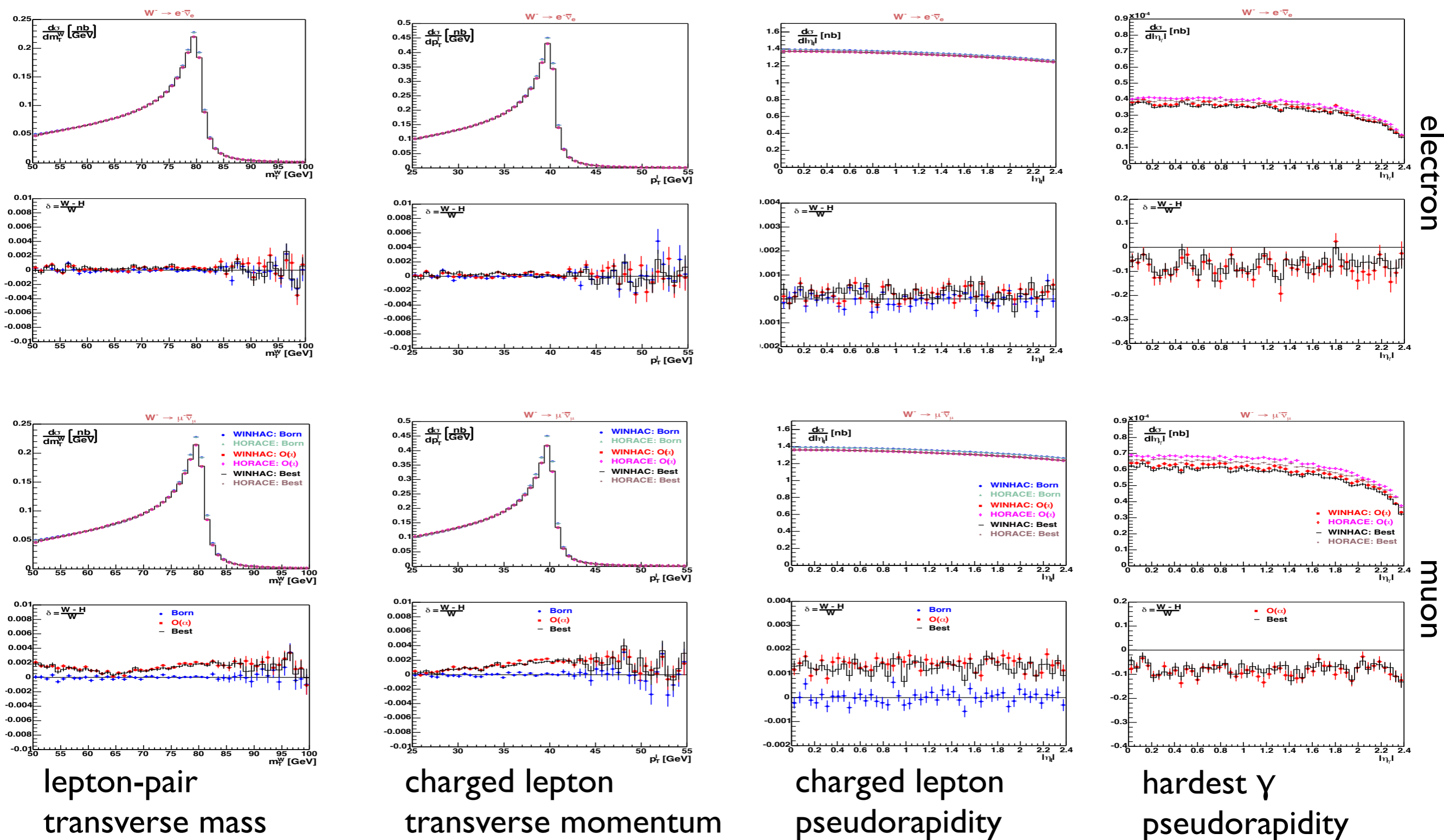
hardest γ energy

hardest γ polar angle

HORACE vs WINHAC comparison: hep-ph/0402235

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hadron-level comparison



HORACE vs WINHAC comparison: hep-ph/0402235

C.M. Carloni Calame, S. Jadach, G. Montagna, O. Nicrosini, W. Płaczek

- Very good agreement for the fiducial xsecs
- Differential distributions agree in general at the 0.1% level
exceptions (discrepancies up to 10%) occur in very suppressed phase-space corners in the parton level comparison
- when a $O(0.2\%)$ systematic difference is observed at hadron level, it is due to the different treatment of subleading $O(\alpha)$ terms while higher-order corrections are in very good agreement
- neglecting the muon mass, in the 2004 version of HORACE, was the main cause of the $O(0.2\%)$ systematic difference for leptonic observables
- photonic observables could differ at the $O(\%)$ level, again because of different treatment of subleading $O(\alpha)$ terms

HORACE vs PHOTOS comparison: arXiv:1510.02458

A.V. Kotwal, B. Jayatilaka

- HORACE is run in its “old” mode, i.e. including only QED-FSR via Parton Shower LL accuracy,
energy extracted according to $P_{q\gamma}(z)$,
angles extracted according to $1/(1-\beta\cos\theta)$
lepton masses (both muon and electron) are kept (at variance with hep-ph/0402235)
effects studied to all orders (best)
- PHOTOS uses the exact first-order matrix element of W and Z boson decay for the photon emission kernel.

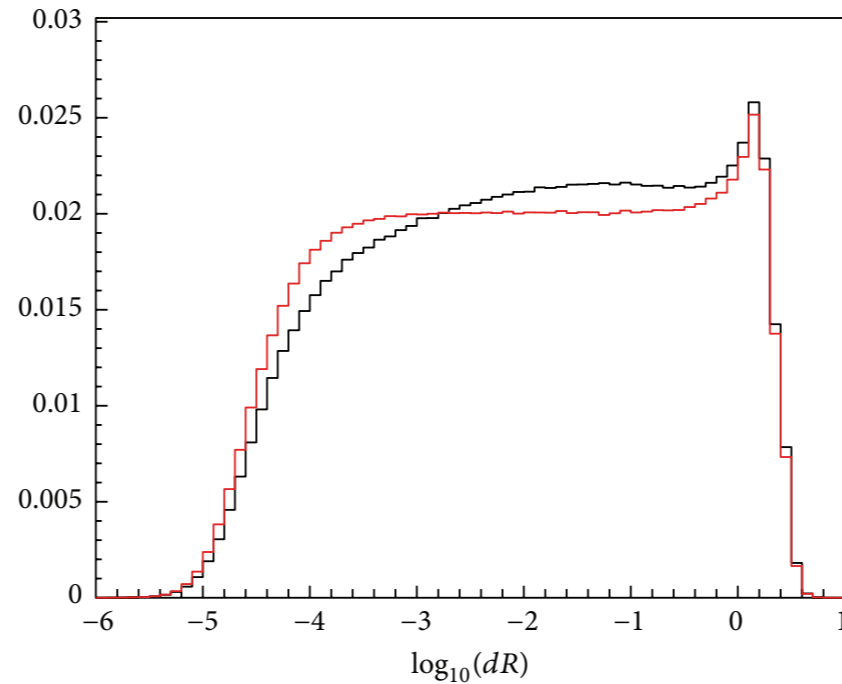
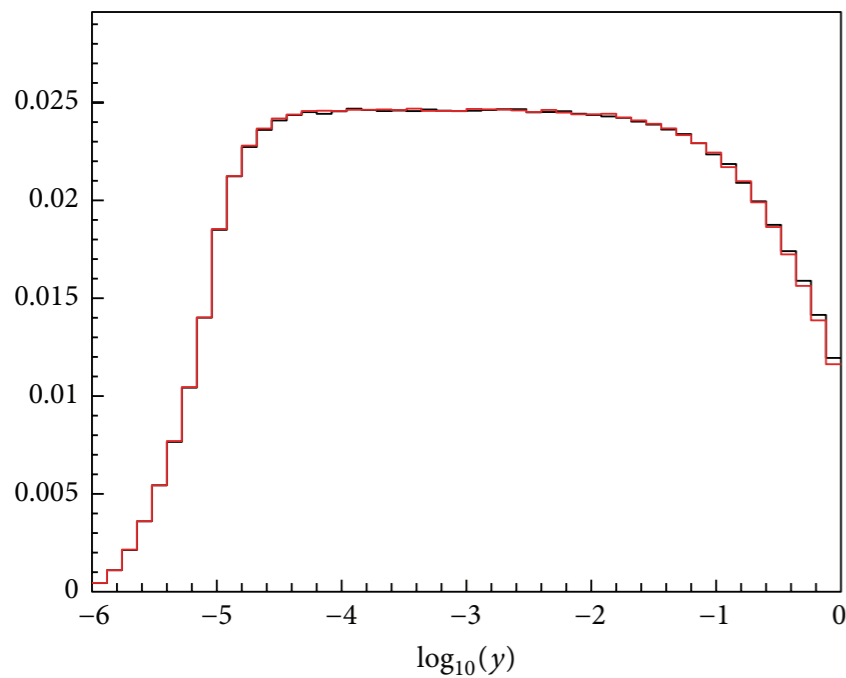
For multiphoton radiation, PHOTOS uses an iterative solution for this kernel, developed on the basis of an exact and complete phase-space parametrization.

This ensures not only resummation of leading-logarithm contributions of higher orders, but also the infrared region of the phase space being accurately simulated.

HORACE vs PHOTOS comparison: arXiv:1510.02458

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$pp \rightarrow e^+e^- + n\gamma$

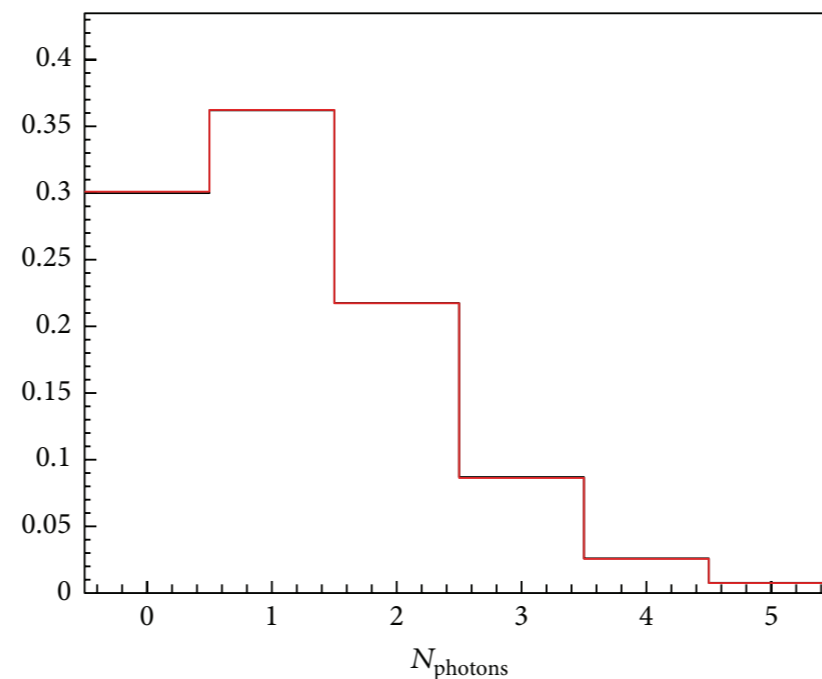
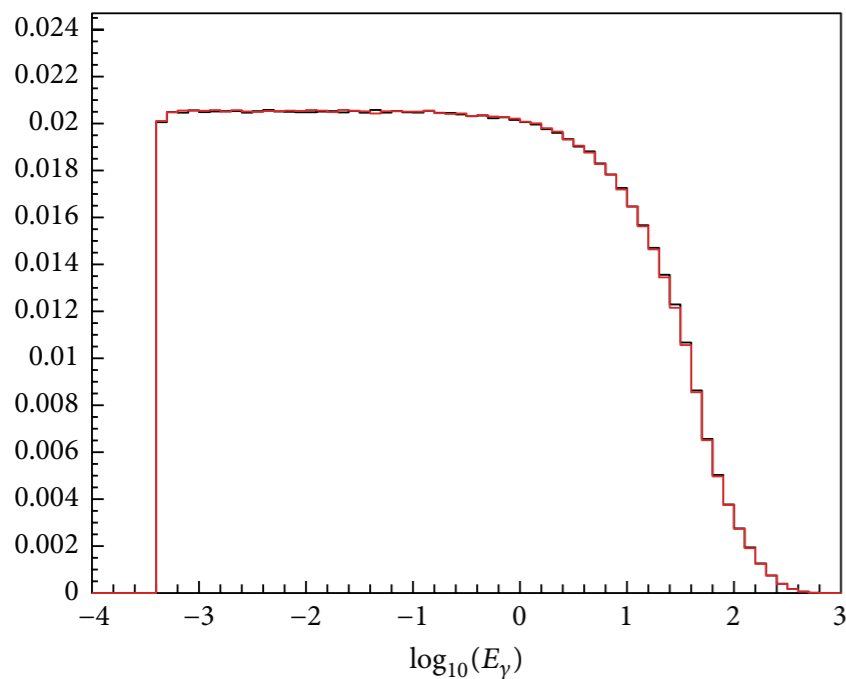


$$y_\gamma = \frac{E_\gamma}{E_\gamma + E_l}$$

$$dR = \sqrt{(\Delta\eta_{l\gamma})^2 + (\Delta\phi_{l\gamma})^2}$$

— Born+PHOTOS
— OLD HORACE

— Born+PHOTOS
— OLD HORACE



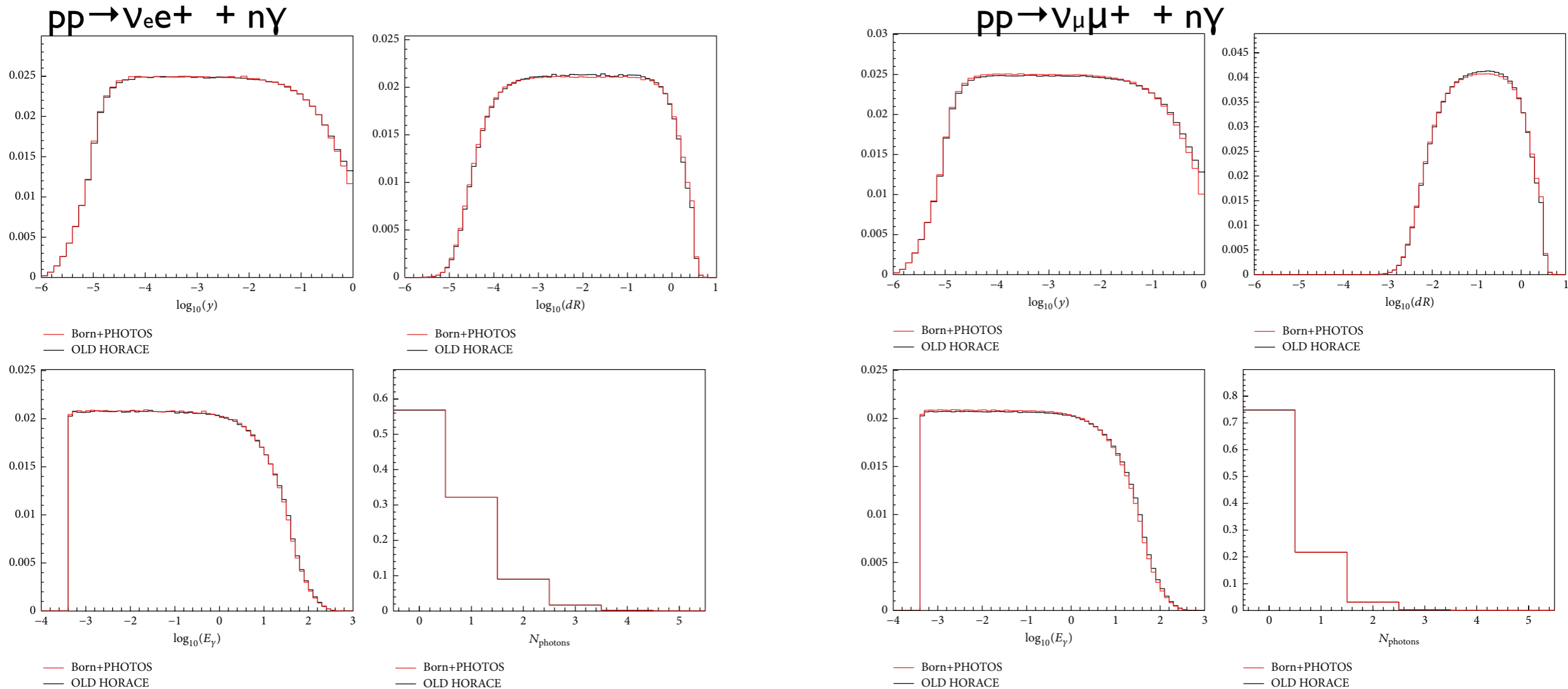
— Born+PHOTOS
— OLD HORACE

— Born+PHOTOS
— OLD HORACE

the discrepancy for $\log(dR)$
disappears in the muon case

HORACE vs PHOTOS comparison: arXiv:1510.02458

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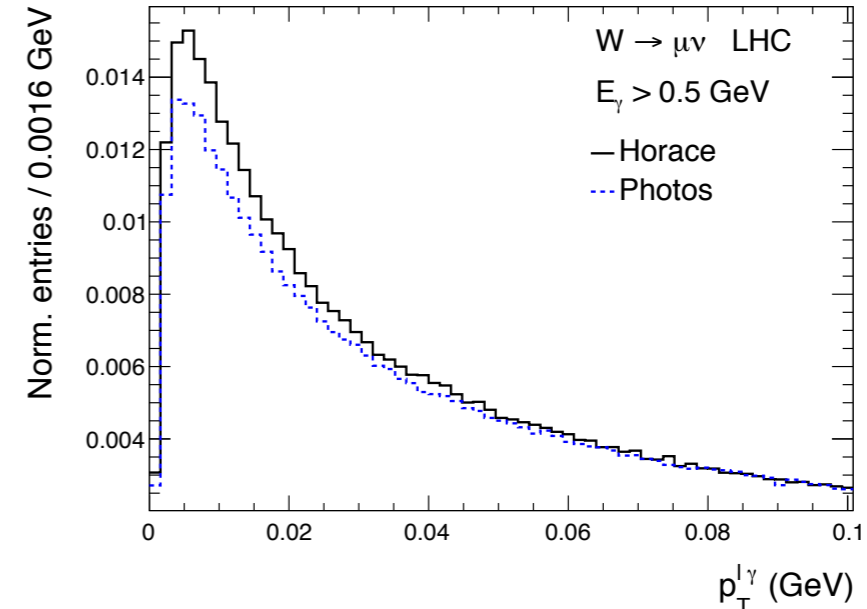
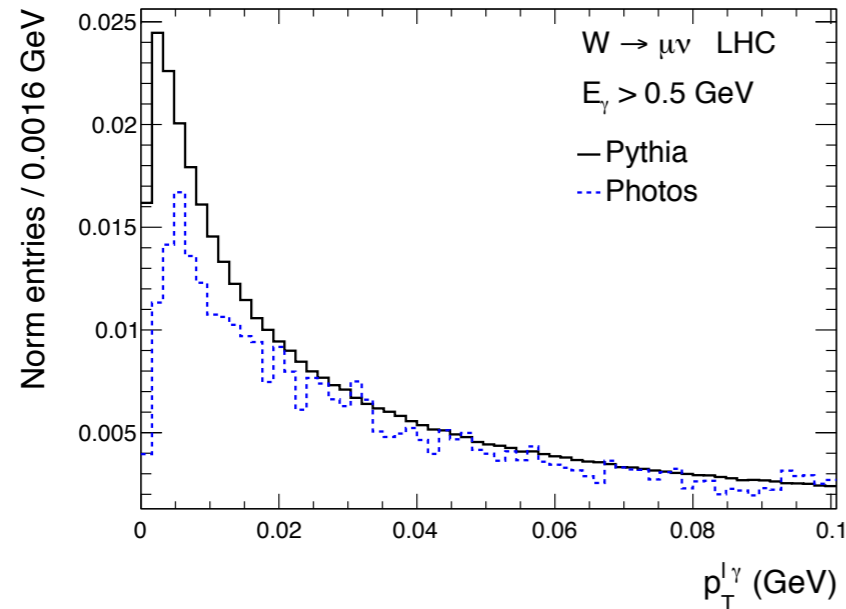
the agreement in the W case, for both electron and muon channels, is remarkable

the $Z \rightarrow l^+l^- n\gamma$ has a stronger (w.r.t. W) radiation pattern, that enhances possible differences visible in the muon case, almost negligible in the electron case

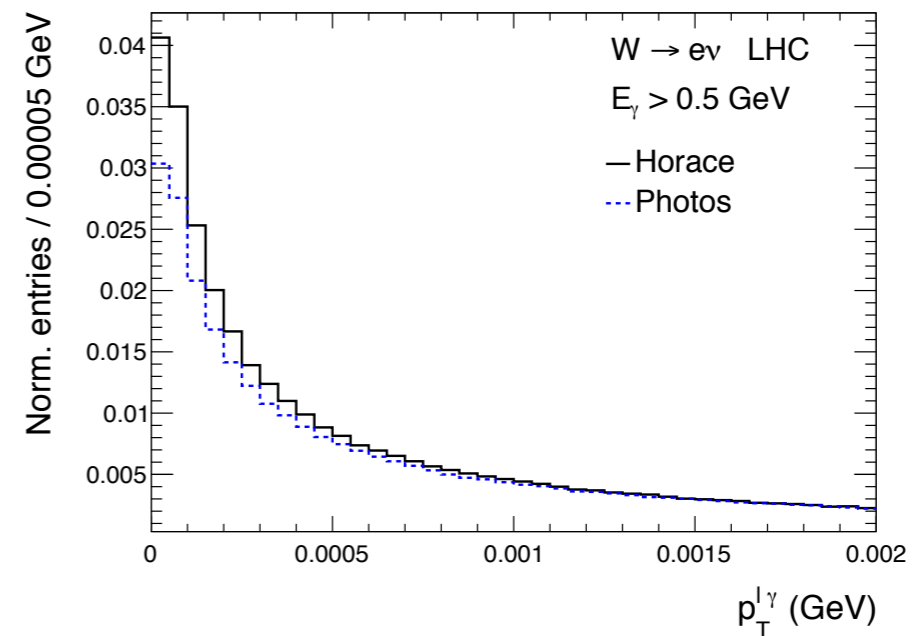
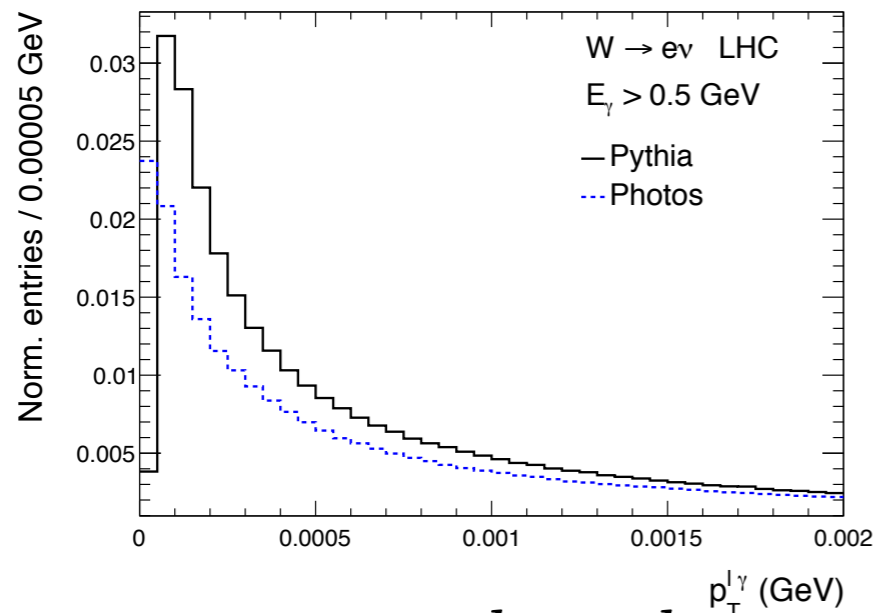
PHOTOS vs PYTHIA-QED vs HORACE comparison: arXiv:1612.02841

C.M. Carloni Calame, M. Chiesa, H. Martinez, G. Montagna, O. Nicrosini, F. Piccinini, A. Vicini

relative lepton-photon p_T , in presence of a cut $E_\gamma > 0.5$ GeV



muon



electron

PYTHIA-QED

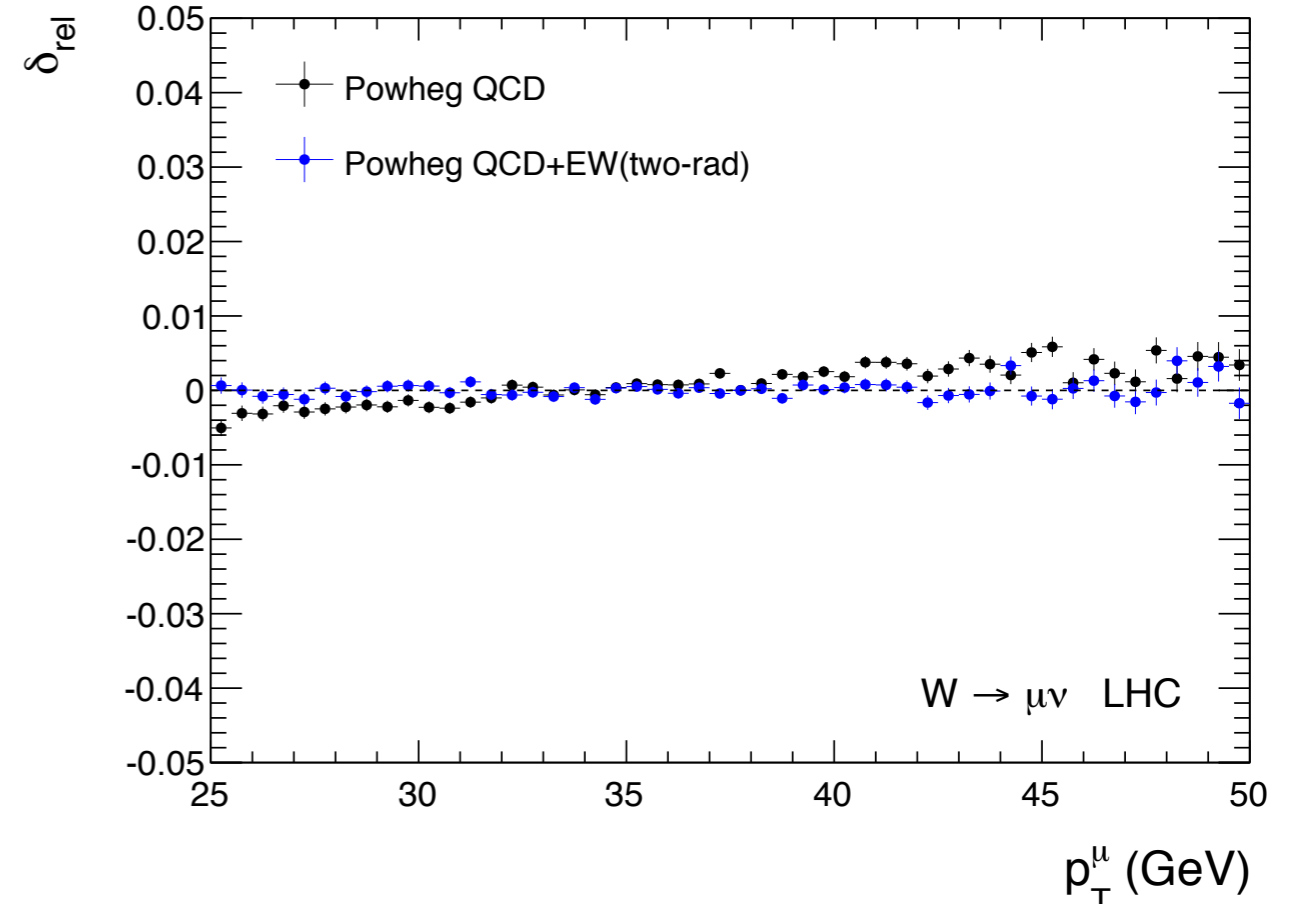
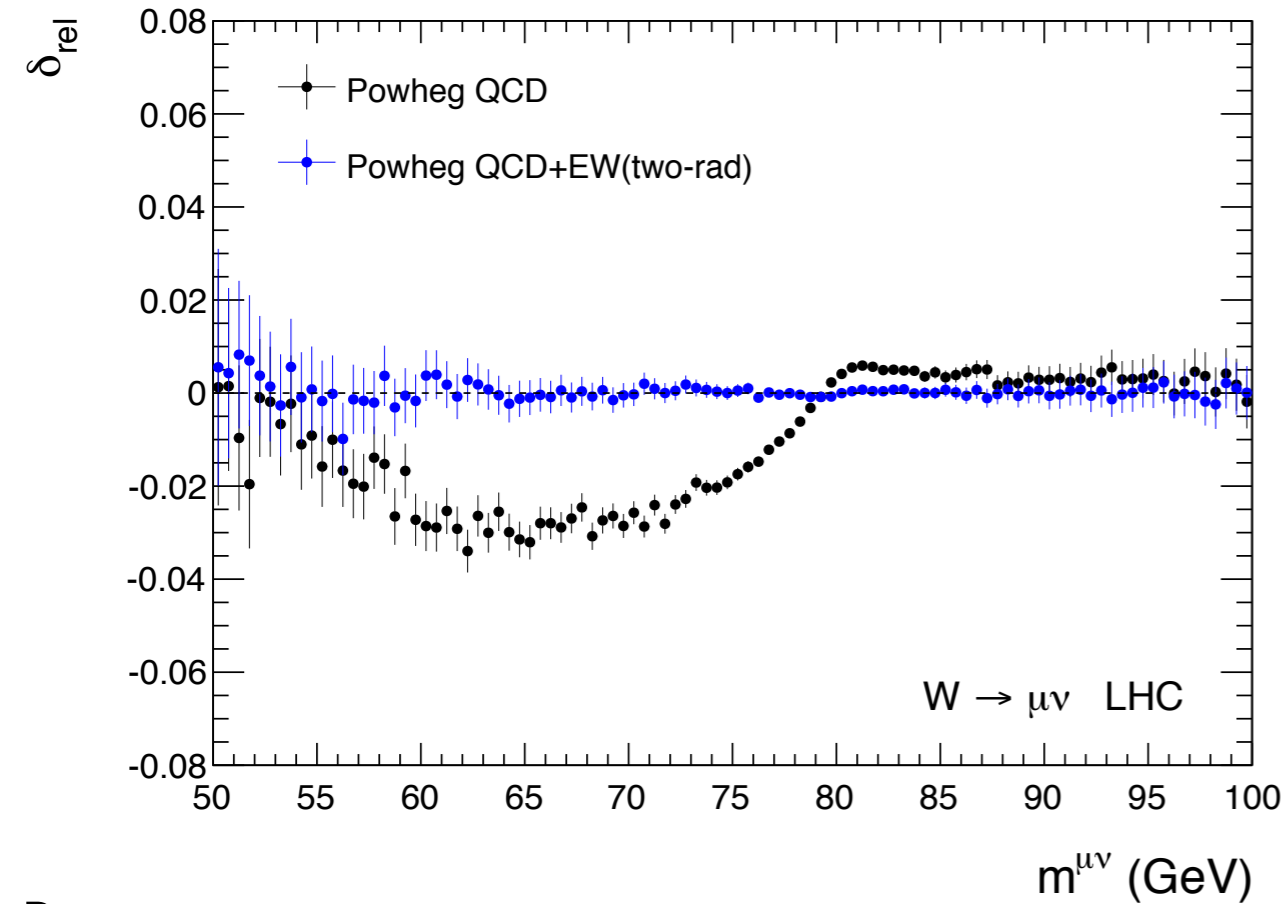
$$dp_T^{l\gamma} / p_T^{l\gamma}$$

HORACE
 PHOTOS

$$d\sigma / d \cos \theta_{l\gamma} \propto 1 / (1 - \beta \cos \theta_{l\gamma})$$

PHOTOS vs PYTHIA-QED comparison: arXiv:1612.02841

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Ratio:

$$\frac{(\text{POWHEG} \otimes \text{PYTHIA-QED})}{(\text{POWHEG} \otimes \text{PHOTOS})} \delta_{rel}$$

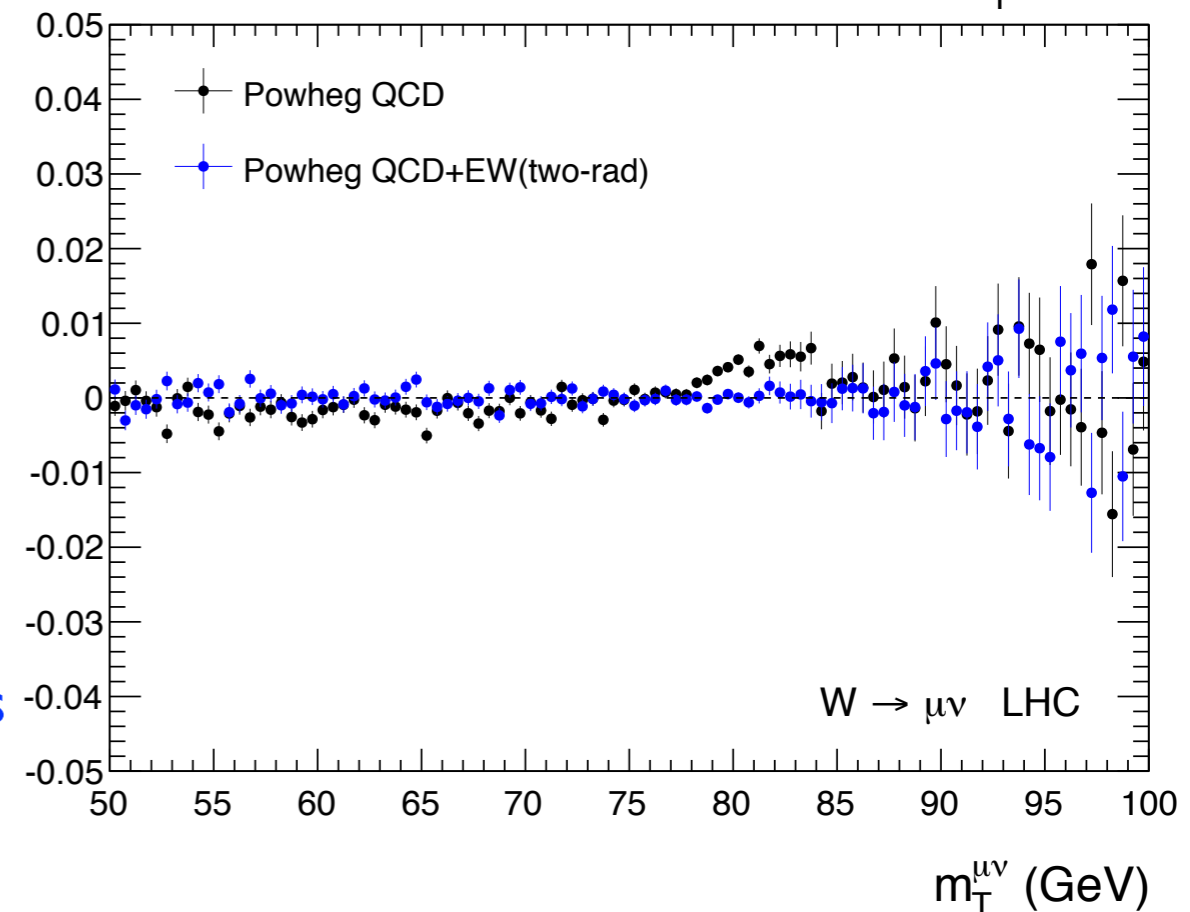
POWHEG versions (identical QCD shower):

black: POWHEG-QCD

blue: POWHEG-(QCD+EW)

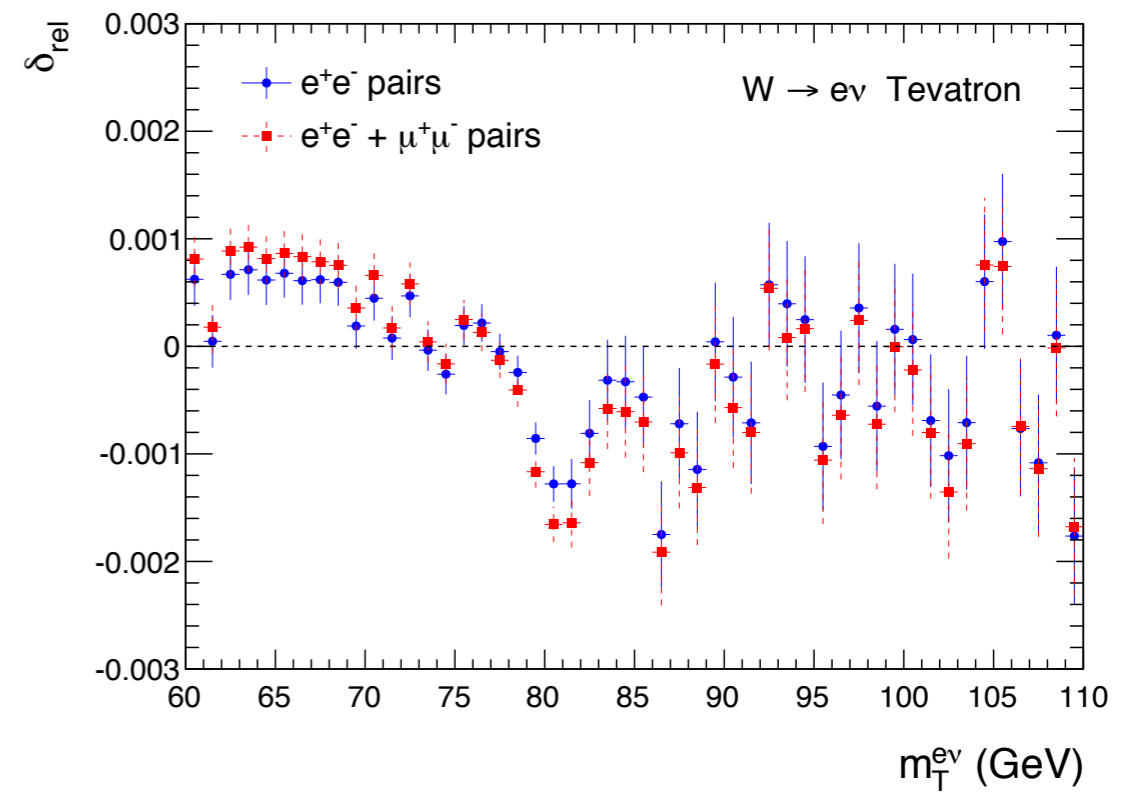
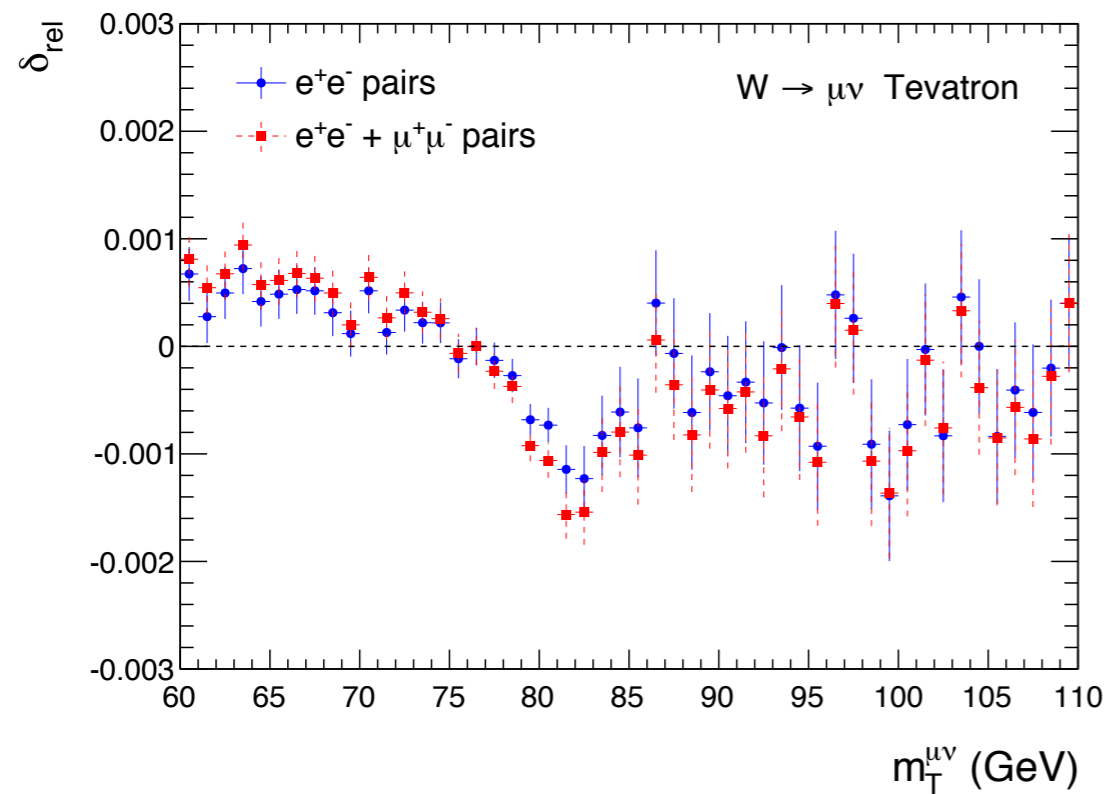
black dots show the impact of $O(\alpha)$ subleading terms

blue dots (after matching) show $O(\alpha^2)$ subleading effects



Leading $O(\alpha^2)$ term: light-pair production in the HORACE shower

C.M. Carloni Calame, M. Chiesa, H. Martinez, G. Montagna, O. Nicrosini, F. Piccinini, A. Vicini



$$\delta = [(\text{QED-FSR with additional pairs}) / (\text{QED-FSR without additional pairs}) - 1]$$

additional class of higher-order terms of $O(\alpha^2 L^2)$ not due to the treatment of photon kinematics

Conclusions

Higher-order corrections available in various codes should all be included in a “best ideal” tool

The spread of different predictions provides hints about the size of the uncertainty

Uncertainty defined as “size of missing higher orders” can only be guessed

QED-FSR models differ by the inclusion of subleading terms

Matching with a given fixed-order matrix element (LO-EW vs NLO-EW)

sets the overall scale of ambiguities:

→ at LO ambiguities appear at $O(\alpha)$ and are relevant in view of precision measurements

→ after matching with NLO-EW ambiguities of $O(\alpha^2)$ appear to be phenomenologically small

Future comparisons

different implementations of

matching of full NLO-EW matrix elements with QED-PS from all external charged legs

differ

not only because of the QED-PS model

but also because of the matching recipe (e.g. HORACE vs POWHEG)

